

Vector & Beam-Target Asymmetries from b1/Azz Data

E. Long, EPJ A (in review)

Elena Long

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**University of
New Hampshire**



Polarization Cross Section

- General electron-deuteron scattering cross section:

$$\sigma(h_e, P, Q) = \sigma_u [1 + h_e (A_e + PA_{ed}^V + QA_{ed}^T) + PA_d^V + QA_d^T]$$


$$A_d^T = \frac{1}{2} A_{zz}$$

- For b1/Azz, we ideally want $h_e = 0$ & $P = 0$
 - But we can't get $P = 0$ w/ large Q using a DNP target

$$\sigma(0, P, Q) = \sigma_u [1 + PA_d^V + QA_d^T]$$

Tensor Asymmetry: 2 States w/ $h_e = 0$

- We extract A_d^T (& thus A_{zz} & b_1) w/ tensor asymmetries
- For 2 arbitrary polarization states, we have

$$\sigma_1 = \sigma(0, P_1, Q_1) = \sigma_u (1 + P_1 A_d^V + Q_1 A_d^T)$$

$$\sigma_2 = \sigma(0, P_2, Q_2) = \sigma_u (1 + P_2 A_d^V + Q_2 A_d^T)$$

- Making a typical asymmetry gives

$$\frac{\sigma_1 - \sigma_2}{\sigma_1 + \sigma_2} = \frac{(P_1 - P_2)A_d^V + (Q_1 - Q_2)A_d^T}{2 + P_1 A_d^V + P_2 A_d^V + Q_1 A_d^T + Q_2 A_d^T}$$

- Solving for A_d^T gives

$$A_d^T = \left(\frac{\sigma_1 - \sigma_2}{Q_1 \sigma_2 - Q_2 \sigma_1} \right) \left[1 - A_d^V \left(\frac{P_1 \sigma_2 - P_2 \sigma_1}{\sigma_1 - \sigma_2} \right) \right]$$

Gives original proposal Eq.
when 1 state unpolarized
($P_2 = Q_2 = 0$)

$$A_d^T = \left(\frac{\sigma_T - \sigma_u}{Q \sigma_u} \right) \left[1 - P A_d^V \left(\frac{\sigma_u}{\sigma_T - \sigma_u} \right) \right]$$

Formalism

$$A_d^T = \left(\frac{\sigma_1 - \sigma_2}{Q_1\sigma_2 - Q_2\sigma_1} \right) \left[1 - A_d^V \left(\frac{P_1\sigma_2 - P_2\sigma_1}{\sigma_1 - \sigma_2} \right) \right]$$

- We can generalize this equation
 - Tensor asymmetry defined by polarization configurations

Asym. to measure $\rightarrow A_d^T$

$$\begin{pmatrix} h_i \\ P_i \\ Q_i \end{pmatrix}_n$$

$\leftarrow$ Beam Helicity States
 $\leftarrow$ Target Vector States
 $\leftarrow$ Target Tensor States
 $\nwarrow$ # of configs

$$A_d^T \begin{pmatrix} 0, 0 \\ P_1, P_2 \\ Q_1, Q_2 \end{pmatrix}_2 = R_2^Q \left(1 - \frac{A_d^V}{R_2^P} \right)$$

Primary Measurement $\rightarrow R_2^Q$
 Contamination Term $\rightarrow R_2^P$

with

$$R_2^Q = \frac{\sigma_1 - \sigma_2}{Q_1\sigma_2 - Q_2\sigma_1}$$

$$R_2^P = \frac{\sigma_1 - \sigma_2}{P_1\sigma_2 - P_2\sigma_1}$$

Formalism

- Same formalism holds for many different configurations!

$$A_d^T = R_n^Q \left(1 - \frac{A_d^V}{R_n^P} - \frac{A_e}{R_n^h} - \frac{A_{ed}^V}{R_n^{hP}} - \frac{A_{ed}^T}{R_n^{hQ}} \right)$$

- But will get different R terms for each
- Can rearrange to get vector & beam-target asymmetries

4 State Tensor Asymmetry w/ $h_e = 0$

Target advances allow us to maintain large P while enhancing or suppressing Q

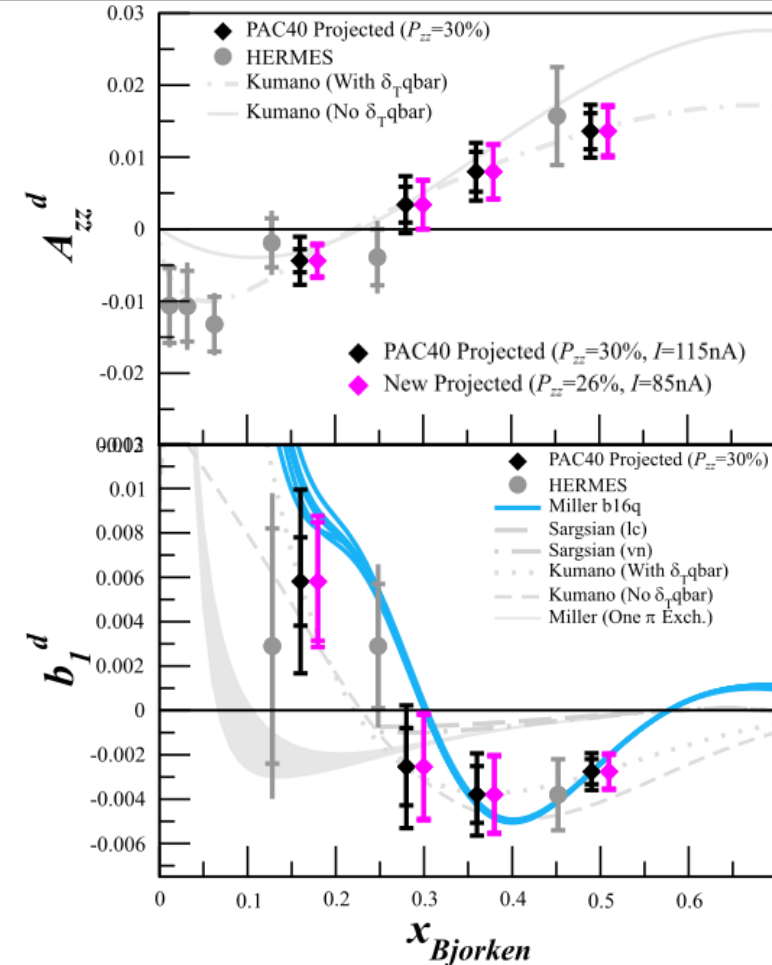
- Rapid tensor enhancement/suppression decreases systematic uncertainty

$$A_d^T \begin{pmatrix} 0, 0, 0, 0 \\ P, -P, P, -P \\ Q, Q, 0, 0 \end{pmatrix}_4 = R_4^Q \left(1 - \frac{A_d^V}{R_4^P} \right)$$

with

$$R_4^Q(Q, Q, 0, 0) = \left(\frac{1}{f} \right) \left(\frac{(Y_1 + Y_2) - (Y_3 + Y_4)}{Q(Y_3 + Y_4)} \right),$$

$$\frac{1}{R_4^P(P, -P, P, -P)} = 0.$$



Parasitic Vector Asymmetry (2 State)

With b1/Azz data, we can extract the vector asymmetry in two different 2-state configurations

$$A_d^V \begin{pmatrix} 0,0 \\ P,-P \\ 0,0 \end{pmatrix}_2 = R_2^P \left(1 - \frac{A_d^T}{R_2^Q} \right)$$

$$A_d^V \begin{pmatrix} 0,0 \\ P,-P \\ Q,Q \end{pmatrix}_2 = R_2^P \left(1 - \frac{A_d^T}{R_2^Q} \right)$$

or

$$R_2^P(P, -P) = \frac{Y_1 - Y_2}{fP(Y_1 + Y_2)},$$

$$\frac{1}{R_2^Q(0,0)} = 0.$$

Statistics ~ the ½
that of Azz/b1

$$R_2^P(P, -P) = \frac{Y_1 - Y_2}{fP(Y_1 + Y_2)},$$

$$\frac{1}{R_2^Q(Q,Q)} = -fQ$$

Statistics ~ the ½
that of Azz/b1

Parasitic Vector Asymmetry (4 State)

Can combine into a single 4-state measurement, increasing statistics but including tensor term

$$A_d^V \begin{pmatrix} 0, 0, 0, 0 \\ P, P, -P, -P \\ Q, 0, Q, 0 \end{pmatrix}_4 = R_4^P \left(1 - \frac{A_d^T}{R_4^Q} \right)$$

$$R_4^P(P, P, -P, -P) = \left(\frac{1}{fP} \right) \left(\frac{(Y_1 + Y_2) - (Y_3 + Y_4)}{(Y_1 + Y_2) + (Y_3 + Y_4)} \right),$$
$$\frac{1}{R_4^Q(Q, 0, Q, 0)} = -\frac{1}{2}fQ,$$

Parasitic Beam-Target Double Spin Asymmetries (8 states)

Since the JLab beam is polarized, we can split the beam helicities from our 4 state measurement to get 8 state measurements – Can extract vector & tensor DSA!

Tensor DSA

$$A_{ed}^T \left(\begin{array}{c} \pm h, \pm h, \pm h, \pm h \\ P, -P, P, -P \\ Q, Q, 0, 0 \end{array} \right)_{8'} = R_{8'}^{hQ} \left(1 - \frac{A_d^T}{R_{8'}^Q} - \frac{A_d^V}{R_{8'}^P} - \frac{A_e}{R_{8'}^h} - \frac{A_{ed}^V}{R_{8'}^{hP}} \right)$$

$$R_{8'}^{hQ} \left(\begin{array}{c} \pm h, \pm h, \pm h, \pm h \\ Q, Q, 0, 0 \end{array} \right) = \left(\frac{2}{fQh} \right) \left(\frac{Y_{-34}^{+12} - Y_{+34}^{-12}}{Y_{-34}^{+12} + Y_{+34}^{-12}} \right),$$

$$Y_{-34}^{+12} = Y_{+1} + Y_{+2} + Y_{-3} + Y_{-4},$$

$$Y_{+34}^{-12} = Y_{-1} + Y_{-2} + Y_{+3} + Y_{+4},$$

Statistics roughly the same as Azz/b1, but scaled by beam helicity

$$\begin{aligned} \frac{1}{R_{8'}^{hP} \left(\begin{array}{c} \pm h, \pm h, \pm h, \pm h \\ P, -P, P, -P \end{array} \right)} &= 0, \\ \frac{1}{R_{8'}^h(\pm h, \pm h, \pm h, \pm h)} &= 0, \\ \frac{1}{R_{8'}^Q(Q, Q, 0, 0)} &= -\frac{1}{2}fQ, \\ \frac{1}{R_{8'}^P(P, -P, P, -P)} &= 0, \end{aligned}$$

Parasitic Beam-Target Double Spin Asymmetries (8 states)

Since the JLab beam is polarized, we can split the beam helicities from our 4 state measurement to get 8 state measurements – Can extract vector & tensor DSA!

Vector DSA

$$A_{ed}^V \begin{pmatrix} \pm h, \pm h, \pm h, \pm h \\ P, P, -P, -P \\ Q, 0, Q, 0 \end{pmatrix}_{8'} = R_{8'}^{hP} \left(1 - \frac{A_d^T}{R_{8'}^Q} - \frac{A_d^V}{R_{8'}^P} - \frac{A_e}{R_{8'}^h} - \frac{A_{ed}^T}{R_{8'}^{hQ}} \right)$$

$$R_{8'}^{hP} \begin{pmatrix} \pm h, \pm h, \pm h, \pm h \\ P, P, -P, -P \end{pmatrix} = \left(\frac{2}{fPh} \right) \begin{pmatrix} Y_{-34}^{+12} - Y_{+34}^{-12} \\ Y_{-34}^{+12} + Y_{+34}^{-12} \end{pmatrix},$$

$$Y_{-34}^{+12} = Y_{+1} + Y_{+2} + Y_{-3} + Y_{-4},$$

$$Y_{+34}^{-12} = Y_{-1} + Y_{-2} + Y_{+3} + Y_{+4},$$

Statistics roughly the same as Azz/b1, but scaled by beam helicity

$$\frac{1}{R_{8'}^{hQ} \begin{pmatrix} \pm h, \pm h, \pm h, \pm h \\ Q, 0, Q, 0 \end{pmatrix}} = 0,$$

$$\frac{1}{R_{8'}^h(\pm h, \pm h, \pm h, \pm h)} = 0,$$

$$\frac{1}{R_{8'}^Q(Q, 0, Q, 0)} = -\frac{1}{2}fQ,$$

$$\frac{1}{R_{8'}^P(P, P, -P, -P)} = 0.$$

Complications

Things get messier when polarizations aren't exactly equal, but the algebra has all been worked out in upcoming EPJA paper

Also includes:

- Configurations better suited for atomic beam source targets/collider experiments
 - 3 state configurations are very useful w/ pure polarization states
- Full uncertainty equations

Uncertainties

$$\delta A_d^T = \left[\left(\left(1 - \frac{A_d^V}{R_n^P} - \frac{A_e}{R_n^h} - \frac{A_{ed}^V}{R_n^{hP}} - \frac{A_{ed}^T}{R_n^{hQ}} \right) \delta R_n^Q \right)^2 + \left(A_d^V \frac{R_n^Q}{(R_n^P)^2} \delta R_n^P \right)^2 \right. \\ + \left(\frac{R_n^Q}{R_n^P} \delta A_d^V \right)^2 + \left(A_e \frac{R_n^Q}{(R_n^h)^2} \delta R_n^h \right)^2 + \left(\frac{R_n^Q}{R_n^h} \delta A_e \right)^2 \\ + \left(A_{ed}^V \frac{R_n^Q}{(R_n^{hP})^2} \delta R_n^{hP} \right)^2 + \left(\frac{R_n^Q}{R_n^{hP}} \delta A_{ed}^V \right)^2 + \left(A_{ed}^T \frac{R_n^Q}{(R_n^{hQ})^2} \delta R_n^{hQ} \right)^2 \\ \left. + \left(\frac{R_n^Q}{R_n^{hQ}} \delta A_{ed}^T \right)^2 \right]^{(1/2)}$$

$$\frac{R_2^Q}{R_2^P} = \frac{P_1 Y_1 - P_2 Y_2}{Q_1 Y_1 - Q_2 Y_2},$$

$$\frac{R_3^Q}{R_3^P} = \frac{(P_1 + P_2) Y_0 - P_0 (Y_1 + Y_2)}{(Q_1 + Q_2) Y_0 - Q_0 (Y_1 + Y_2)},$$

$$\frac{R_{4'}^Q}{R_{4'}^P} = \frac{(P_1 + P_2 + P_3) Y_0 - P_0 (Y_1 + Y_2 + Y_3)}{(Q_1 + Q_2 + Q_3) Y_0 - Q_0 (Y_1 + Y_2 + Y_3)},$$

$$\frac{R_4^Q}{R_4^P} = \frac{(P_1 + P_2) (Y_3 + Y_4) - (P_3 + P_4) (Y_1 + Y_2)}{(Q_1 + Q_2) (Y_3 + Y_4) - (Q_3 + Q_4) (Y_1 + Y_2)}.$$

Polarization Choice Can Reduce
Uncertainty Terms

Existing & Upcoming Tensor Measurements

Observable	Method	Target Configuration	Target Type	Location and Reaction
$T_{20} \approx \sqrt{2} \frac{A_d^T}{P_2(\cos \theta)}$	R_2^Q	$\begin{pmatrix} h_1, h_2 \\ P_1, P_2 \\ Q_1, Q_2 \end{pmatrix} = \begin{pmatrix} 0, 0 \\ 0, 0 \\ +Q, -2Q \end{pmatrix}$	Atomic Beam Source	NIKHEF $d(e, e' d')$
A_d^T	R_2^Q	$\begin{pmatrix} h_1, h_2 \\ P_1, P_2 \\ Q_1, Q_2 \end{pmatrix} = \begin{pmatrix} 0, 0 \\ -P, 0 \\ +Q_{eq}, 0 \end{pmatrix}$	Solid ND ₃	Bonn $d(e, e' d')$
$T_{20} \approx \sqrt{2} \frac{A_d^T}{P_2(\cos \theta)}$	R_2^Q	$\begin{pmatrix} h_1, h_2 \\ P_1, P_2 \\ Q_1, Q_2 \end{pmatrix} = \begin{pmatrix} 0, 0 \\ P, 0 \\ +Q, 0 \end{pmatrix}$	Solid ND ₃	TRIUMF $d(\pi, 2p)$
$T_{20} \approx \sqrt{2} \frac{A_d^T}{P_2(\cos \theta)}$	R_2^Q	$\begin{pmatrix} h_1, h_2 \\ P_1, P_2 \\ Q_1, Q_2 \end{pmatrix} = \begin{pmatrix} 0, 0 \\ 0, 0 \\ +Q, -Q \end{pmatrix}$	Atomic Beam & Storage Cell	VEPP-2, VEPP-3 $d(e, e' d')$
A_d^T	R_3^Q	$\begin{pmatrix} h_1, h_2, h_0 \\ P_1, P_2, P_0 \\ Q_1, Q_2, Q_0 \end{pmatrix} = \begin{pmatrix} 0, 0, 0 \\ P, -P, 0 \\ +Q_{eq}, +Q_{eq}, -2Q \end{pmatrix}$	Atomic Beam Source	MIT-Bates $d(e, e' d')$, $d(e, e' p)$
$2A_d^T = A_{zz}$	$R_{4'}^Q$	$\begin{pmatrix} h_1, h_2, h_3, h_0 \\ P_1, P_2, P_3, P_0 \\ Q_1, Q_2, Q_3, Q_0 \end{pmatrix} = \begin{pmatrix} 0, 0, 0, 0 \\ P, -P, 0, 0 \\ Q, Q, Q, -2Q \end{pmatrix}$	Atomic Beam Source	DESY HERA $d(e^+, e^{+'})$
$2A_d^T = A_{zz}$	R_4^Q	$\begin{pmatrix} h_1, h_2, h_3, h_4 \\ P_1, P_2, P_3, P_4 \\ Q_1, Q_2, Q_3, Q_4 \end{pmatrix} = \begin{pmatrix} 0, 0, 0, 0 \\ P, -P, P, -P \\ Q, Q, 0, 0 \end{pmatrix}$	Solid ND ₃	JLab $d(e, e')$

Optimal Experimental Configurations

Target Type	Observable	Method	Target Configuration
ABS, DNP with H,D Molecule	A_d^T	Eqs. R_2^Q (58-60, 72-73)	$\begin{pmatrix} h_1, h_2 \\ P_1, P_2 \\ Q_1, Q_2 \end{pmatrix} = \begin{pmatrix} 0, 0 \\ P, 0 \\ +Q, 0 \end{pmatrix}$
ABS, DNP + ssRF	A_d^V	Eqs. R_2^P (83-85)	$\begin{pmatrix} h_1, h_2 \\ P_1, P_2 \\ Q_1, Q_2 \end{pmatrix} = \begin{pmatrix} 0, 0 \\ P, -P \\ 0, 0 \end{pmatrix}$
Dual-Cell DNP	A_d^T	Eqs. R_3^Q (76-78)	$\begin{pmatrix} h_1, h_2, h_0 \\ P_1, P_2, P_0 \\ Q_1, Q_2, Q_0 \end{pmatrix} = \begin{pmatrix} 0, 0, 0 \\ P, -P, 0 \\ Q, Q, 0 \end{pmatrix}$
ABS	A_d^T	Eqs. R_3^Q (64-66)	$\begin{pmatrix} h_1, h_2, h_0 \\ P_1, P_2, P_0 \\ Q_1, Q_2, Q_0 \end{pmatrix} = \begin{pmatrix} 0, 0, 0 \\ P, -P, 0 \\ Q, Q, -2Q \end{pmatrix}$
ABS	A_d^T	Eqs. $R_{4'}^Q$ (67-71)	$\begin{pmatrix} h_1, h_2, h_3, h_0 \\ P_1, P_2, P_3, P_0 \\ Q_1, Q_2, Q_3, Q_0 \end{pmatrix} = \begin{pmatrix} 0, 0, 0, 0 \\ P, -P, 0, 0 \\ Q, Q, Q, -2Q \end{pmatrix}$
DNP + ssRF	A_d^T	Eqs. R_4^Q (80-82)	$\begin{pmatrix} h_1, h_2, h_3, h_4 \\ P_1, P_2, P_3, P_4 \\ Q_1, Q_2, Q_3, Q_4 \end{pmatrix} = \begin{pmatrix} 0, 0, 0, 0 \\ P, -P, P, -P \\ Q, Q, 0, 0 \end{pmatrix}$
DNP + ssRF	A_d^V	Eqs. R_4^P (86-87)	$\begin{pmatrix} h_1, h_2, h_3, h_4 \\ P_1, P_2, P_3, P_4 \\ Q_1, Q_2, Q_3, Q_4 \end{pmatrix} = \begin{pmatrix} 0, 0, 0, 0 \\ P, P, -P, -P \\ Q, 0, Q, 0 \end{pmatrix}$
DNP + ssRF	A_{ed}^T	Eqs. $R_{8'}^{hQ}$ (88-91)	$\begin{pmatrix} h_1^\pm, h_2^\pm, h_3^\pm, h_4^\pm \\ P_1, P_2, P_3, P_4 \\ Q_1, Q_2, Q_3, Q_4 \end{pmatrix} = \begin{pmatrix} \pm h, \pm h, \pm h, \pm h \\ P, -P, P, -P \\ Q, Q, 0, 0 \end{pmatrix}$
DNP + ssRF	A_{ed}^V	Eqs. $R_{8'}^{hP}$ (92-95)	$\begin{pmatrix} h_1^\pm, h_2^\pm, h_3^\pm, h_4^\pm \\ P_1, P_2, P_3, P_4 \\ Q_1, Q_2, Q_3, Q_4 \end{pmatrix} = \begin{pmatrix} \pm h, \pm h, \pm h, \pm h \\ P, -P, P, -P \\ Q, Q, 0, 0 \end{pmatrix}$

Thank you!
