

DVCS on a longitudinally polarised proton target

Introduction

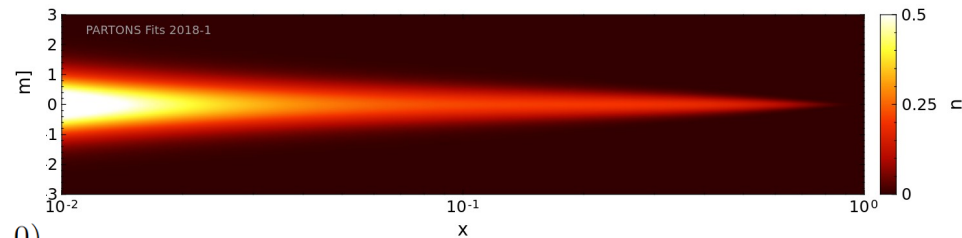
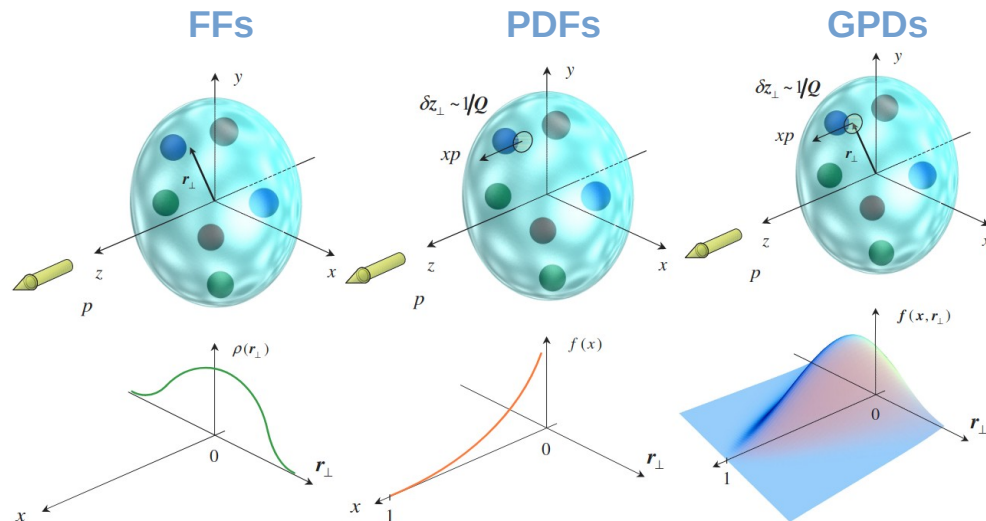
- **Outstanding question : How does the nucleon's mass and spin arise from partons ?**
- First step is to map out the nucleon
- GPDs (Generalized Parton Distributions) describe correlations between the longitudinal momentum and transverse position of partons in the nucleon

3D structure

- Access to its mechanical properties
- Access to the spin decomposition of the nucleon

$$\frac{1}{2} = \sum_q J^q + J^g$$

$$= \sum_q \frac{1}{2} \int_{-1}^1 dx x (H^q(x, \xi, 0) + E^q(x, \xi, 0)) + \frac{1}{2} \int_{-1}^1 dx H^g(x, \xi, 0) + E^g(x, \xi, 0)$$



[Phys.Rept.418:1-387] [Eur. Phys. J. C78 no. 11, 890]

Deeply Virtual Compton Scattering

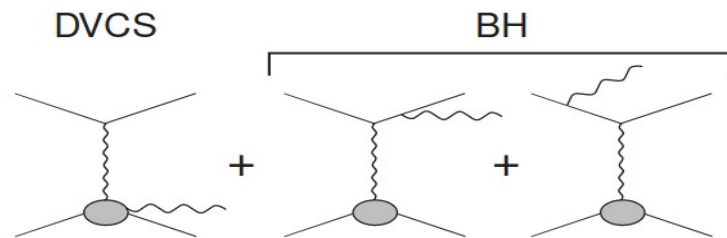
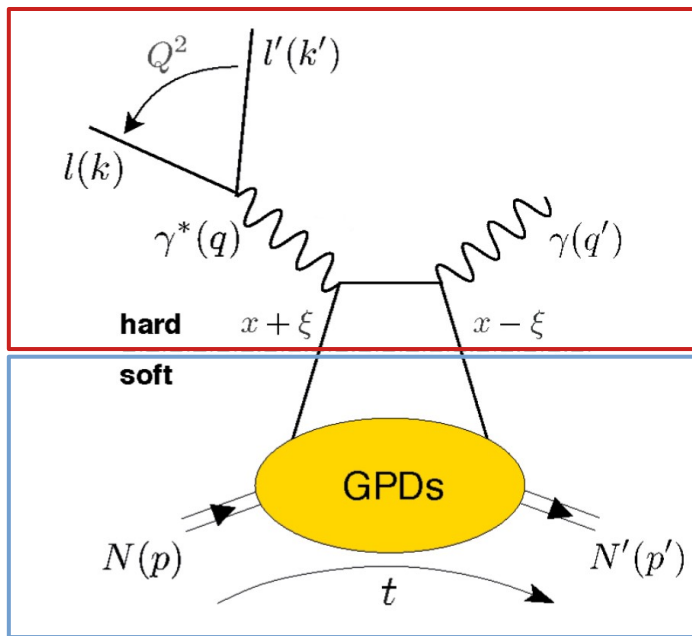
- DVCS offers the most straightforward access to GPDs
- DVCS can be factorised into :
 - **Hard part** $\gamma^* q$ scattering computed in perturbative QCD
 - **Soft part** described by 4 GPDs $H, \tilde{H}, E, \tilde{E}$ at leading order & twist
- Two indistinguishable processes, DVCS and Bethe-Heitler

$$|T|^2 = |T_{\text{DVCS}}|^2 + |T_{\text{BH}}|^2 + \underbrace{T_{\text{DVCS}} T_{\text{BH}}^* + T_{\text{DVCS}}^* T_{\text{BH}}}_I$$

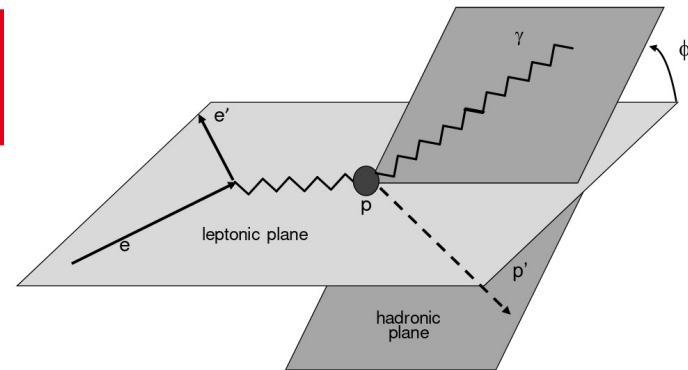
The amplitude is expressed as a function of FFs and CFFs which are functions of GPDs

$$\mathcal{F} = \int_{-1}^1 dx F(\mp x, \xi, t) \left[\frac{1}{x - \xi + i\epsilon} \pm \frac{1}{x + \xi - i\epsilon} \right]$$

[Eur. Phys. J. A 52, 151]



Observables



- Asymmetries in the DVCS cross section are sensitive to CFFs
- Beam spin asymmetry (BSA), polarised electron and unpolarised proton

$$A_{LU}(\phi) \sim \frac{s_{1,unp}^{\mathcal{I}} \sin \phi}{c_{0,unp}^{BH} + (c_{1,unp}^{BH} + c_{1,unp}^{\mathcal{I}} + \dots) \cos \phi \dots} \quad s_{1,unp}^{\mathcal{I}} \propto \Im[F_1 \mathcal{H} + \xi(F_1 + F_2) \tilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E}].$$

- Target spin asymmetry (TSA), unpolarised electron and longitudinally polarised proton

$$A_{UL}(\phi) \sim \frac{s_{1,LP}^{\mathcal{I}} \sin \phi}{c_{0,unp}^{BH} + (c_{1,unp}^{BH} + c_{1,unp}^{\mathcal{I}} + \dots) \cos \phi + \dots} \quad s_{1,LP}^{\mathcal{I}} \propto \Im[F_1 \tilde{\mathcal{H}} + \xi(F_1 + F_2)(\mathcal{H} + \frac{x_b}{2} \mathcal{E}) - \xi(\frac{x_b}{2} F_1 + \frac{t}{4M^2} F_2) \tilde{\mathcal{E}}]$$

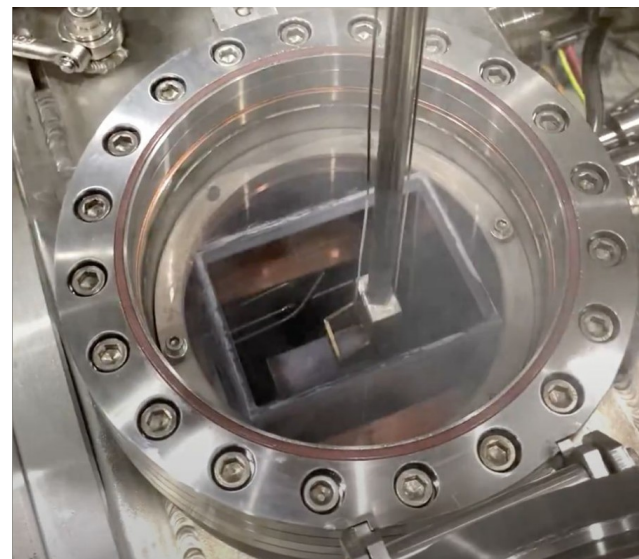
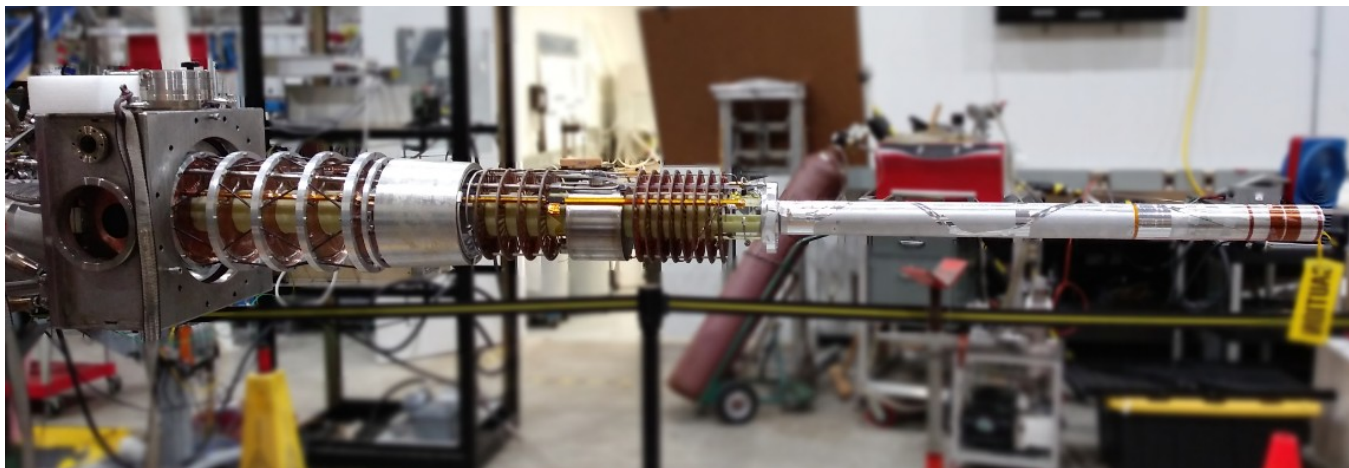
- Double spin asymmetry (DSA), aligned and anti-aligned electron and proton spin

$$A_{LL}(\phi) \sim \frac{c_{0,LP}^{BH} + c_{0,LP}^{\mathcal{I}} + (c_{1,LP}^{BH} + c_{1,LP}^{\mathcal{I}}) \cos \phi}{c_{0,unp}^{BH} + (c_{1,unp}^{BH} + c_{1,unp}^{\mathcal{I}} + \dots) \cos \phi \dots} \quad c_{0,LP}^{\mathcal{I}}, c_{1,LP}^{\mathcal{I}} \propto \Re[F_1 \tilde{\mathcal{H}} + \xi(F_1 + F_2)(\mathcal{H} + \frac{x_b}{2} \mathcal{E}) - \xi(\frac{x_b}{2} F_1 + \frac{t}{4M^2} F_2) \tilde{\mathcal{E}}]$$

The RGC experiment

- New polarised target APOLLO, cryogenic solid target
- Dynamically polarised hydrogen or deuterium in NH₃ or ND₃ cells
- Data taking from June 2022 to March 2023 in three run periods, we will mainly focus on Summer 22 data
- **Goal of this analysis: Measure BSA, TSA and DSA on polarised proton**

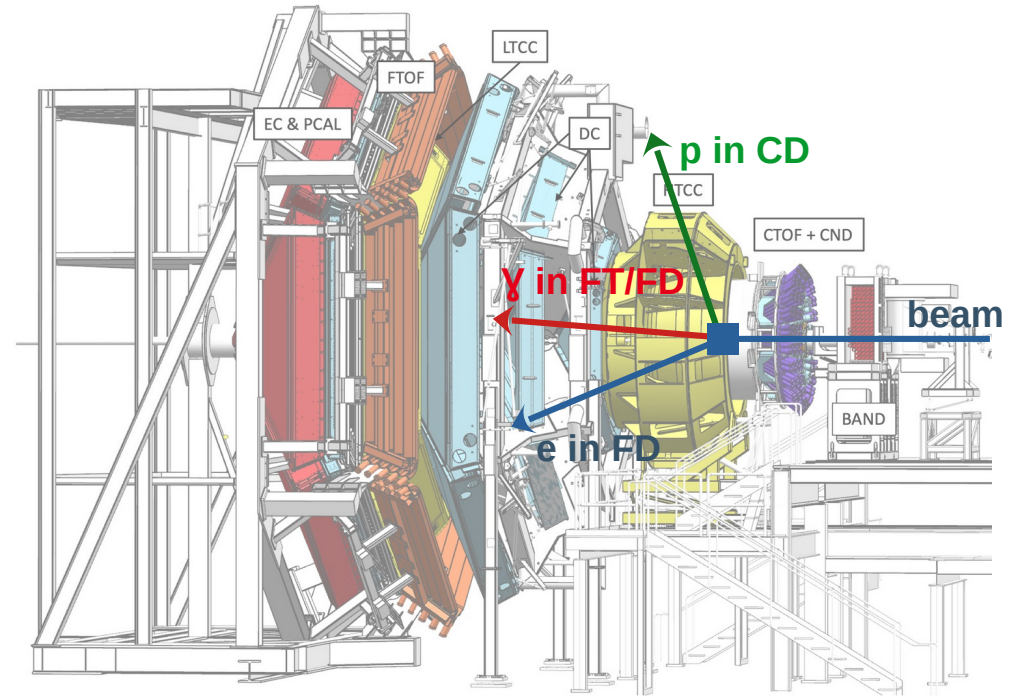
	Summer 2022	Fall 2022	Spring 2023
FT	ON	OFF	ON
Magnets (solenoid, torus)	-1, -1	-1, -1 / +1, -1	-1, -1 / -1, +1
Run range	16089 - 16788	16843 - 17183 / 17185 - 17408	17477 - 17768 / 17769 - 17811
PAC days	28	30	22
Total charge on NH ₃ [mC]	4.5	3.5 / 2.2	1.7 / 0.3



[DOI: 10.25777/36yz-ft35]

DVCS candidate selection

- Request at least one electron, one proton, and one photon in the event
- Kinematic cuts:
 - $Q^2 > 0.95 \text{ GeV}^2/c^4$
 - $W > 1.95 \text{ GeV}/c^2$
 - Electron $P > 1 \text{ GeV}/c$
 - Photon $E > 2 \text{ GeV}$
- Apply DC, ECAL, FTCAL fiducial cuts and PID cuts
- QADB selection of events fit for asymmetry measurements
- Application of GEMC based momentum corrections
- All three outgoing particles are detected so we can build exclusivity variables



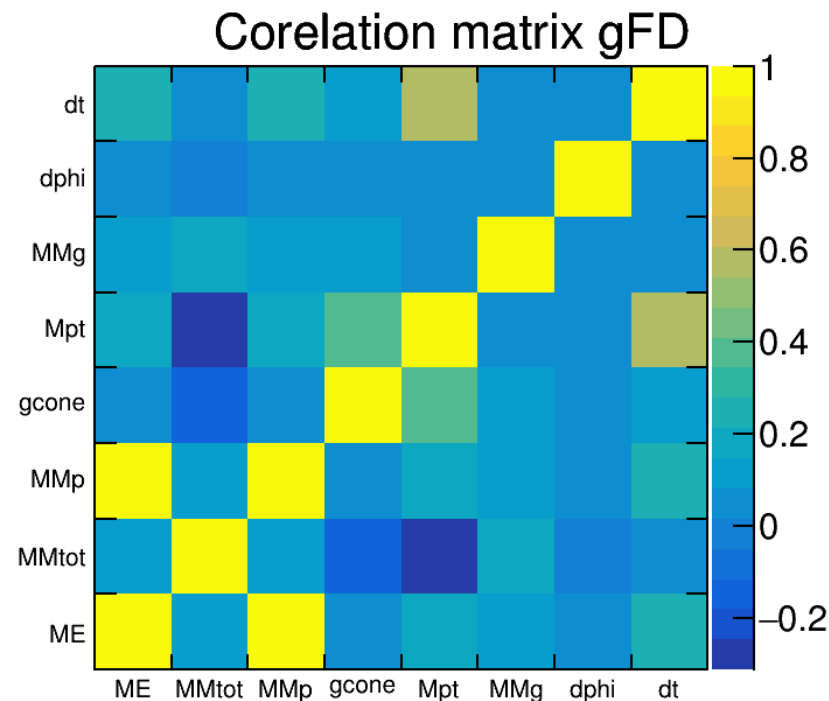
DVCS event selection

- Application of broad exclusivity cuts on 8 variables (level 2 cuts)
- Instead of applying fine cuts on all variable we compute a distance from each event to an ideal exclusive event ($ME=0$, $\Delta\phi=0$, ...)
- The relative weights of each variable needs to be scaled to account for different width of the distributions, different units ...
 - Mahalanobis distance: distance weighted by the covariance matrix between the different variables

$$d(\vec{x}, \vec{y}) = \sqrt{(\vec{x} - \vec{y})^T \Sigma^{-1} (\vec{x} - \vec{y})}$$

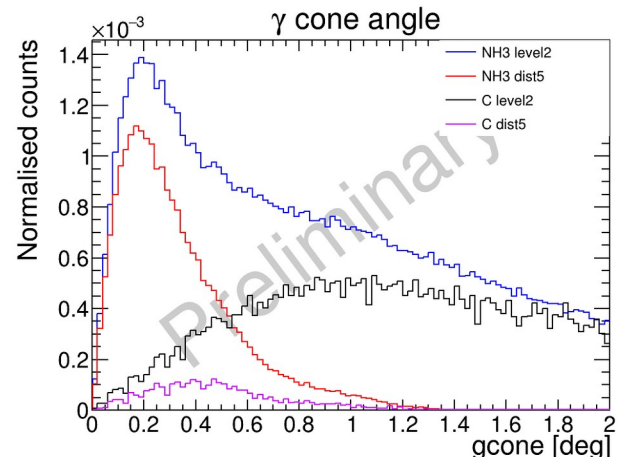
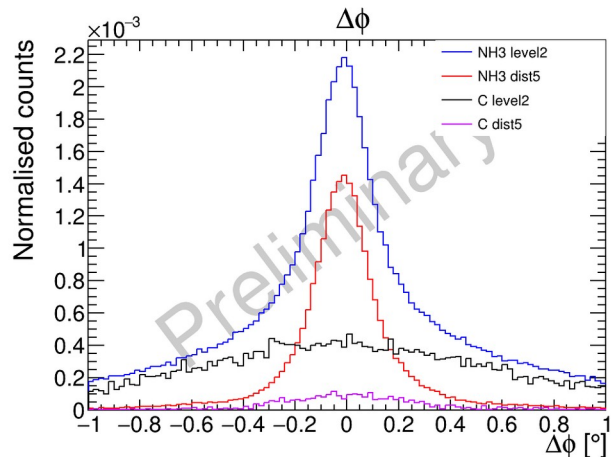
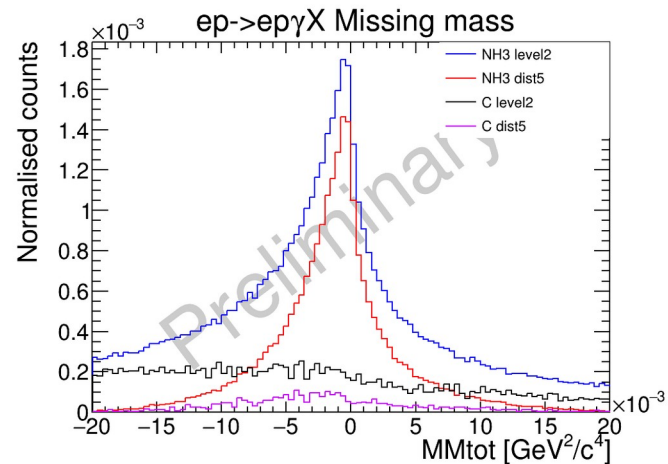
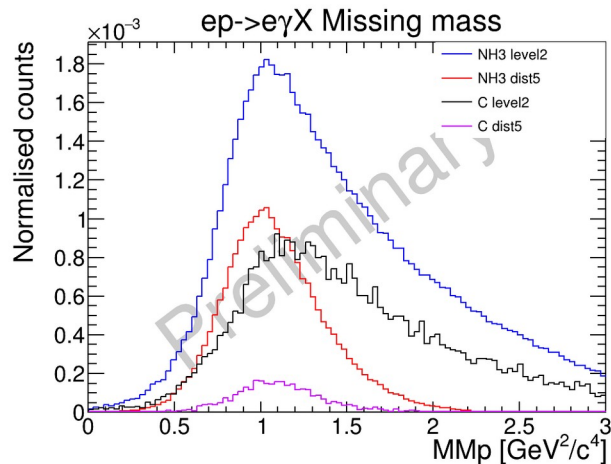
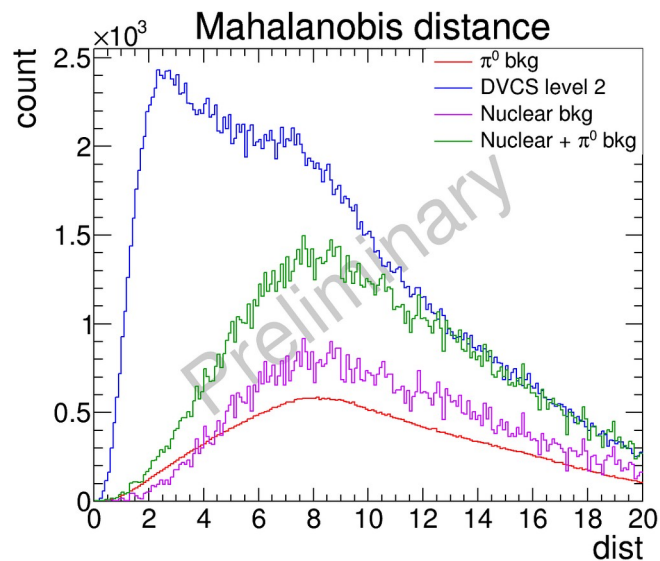
- Covariance matrices computed from a GEMC simulation of DVCS events on a proton target (separated matrices for photons in the FD and FT)
- Apply an exclusivity cut on one variable instead of 8 → simpler to adjust and improved efficiency

Variable	level 2 cuts
MMtot [GeV^2/c^4]	$[-0.1, 0.1]$
MMp [GeV^2/c^4]	$[0, 3]$
ME [GeV]	$[-1.5, 1.5]$
Mpt [GeV/c]	$[0, 1]$
MMg [GeV^2/c^4]	$[-0.8, 0.8]$
g_cone [deg]	$[0, 2]$
$\Delta\phi$ [deg]	$[-3, 3]$
Δt [GeV^2/c^4]	$[0, 2]$



DVCS event selection

- Nuclear background
 - Data taken on a carbon target of similar area density to NH3 → background estimate
- Cut at distance < 5 , able to cut almost all nuclear background



Nuclear background

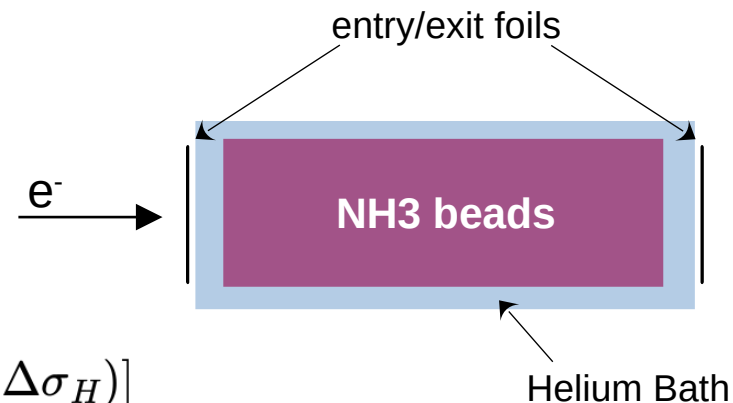
$$A_{UL} = \frac{1}{D_f} \frac{N^{++} + N^{-+} - N^{+-} - N^{--}}{P_t^-(N^{++} + N^{-+}) + P_t^+(N^{+-} + N^{--})}$$

- To take into account the remaining nuclear background, the TSA and DSA are scaled by the dilution factor

$$n_{NH3} = f_c[(L - l_{NH3})\rho_{He}\Delta\sigma_{He} + l_{NH3}\rho_{NH3}(\frac{7}{6}\Delta\sigma_C + 3\Delta\sigma_H)]$$

- The dilution factor is the fraction of the total number of events that comes from the polarised free protons

$$Df = \frac{3l_{NH3}\Delta\sigma_H}{(L - l)\rho_{He}\Delta\sigma_{He} + l_{NH3}(\frac{7}{6}\Delta\sigma_C + 3\Delta\sigma_H)}$$

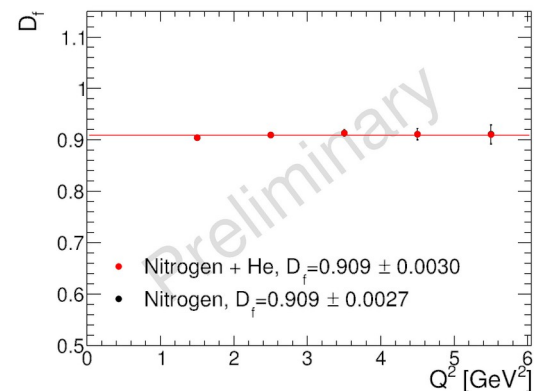
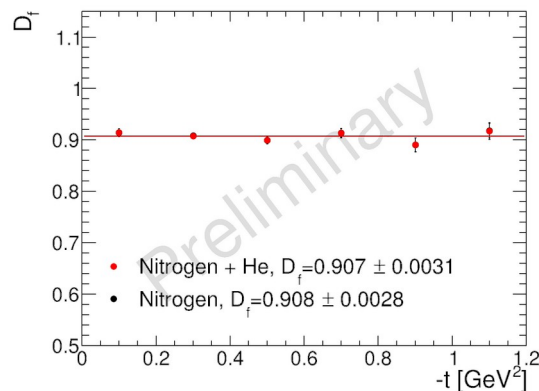
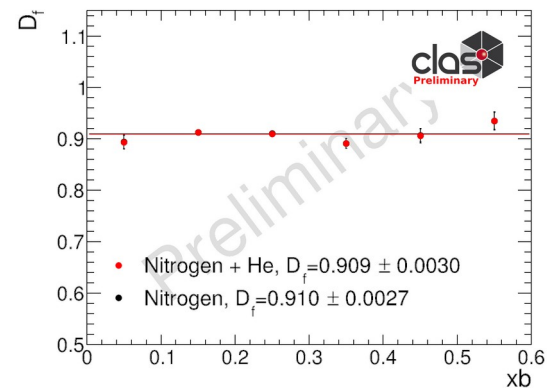
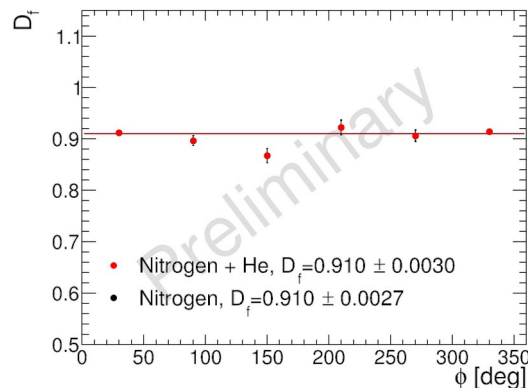


RGC took data with multiple targets:

- Empty target → Foil contribution
- He bath → He contribution
- Carbon → Nitrogen contribution
- CH2 → Hydrogen contribution

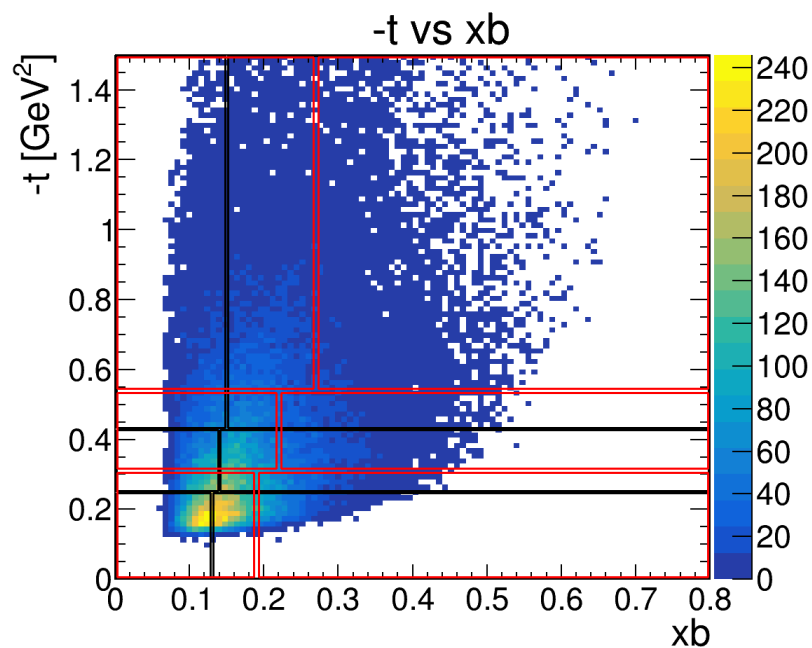
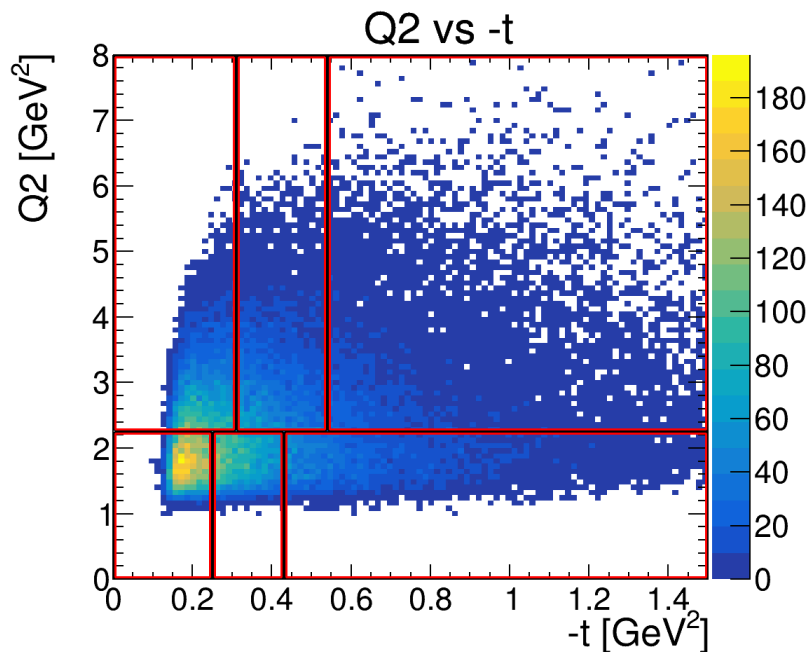
Dilution factor

- Foil contribution is neglected because less than 10 DVCS events collected in the empty target runs
- Helium contribution is very small
- Dilution factor is stable across all kinematic variables at $91 \pm 1\%$
- Packing fraction, $58 \pm 1\%$



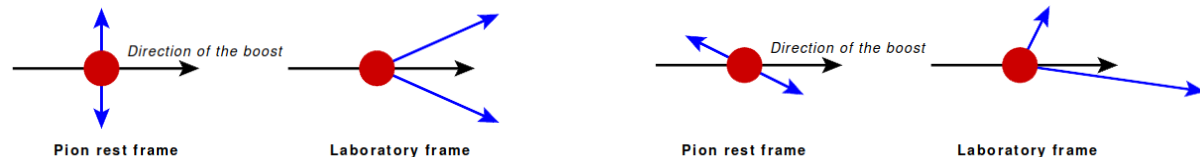
Binning

- Two bins in Q^2 and x_b and three bins in t , same number of DVCS events in each bin
- The binning in ϕ follows the distribution in each bin, with the following requirements
 - A minimum of 300 events per ϕ bin
 - At least 15° wide



π^0 subtraction method

- A $\pi^0 \rightarrow \gamma\gamma$ decay can pass the DVCS selection if one of the photons carries most of the momentum



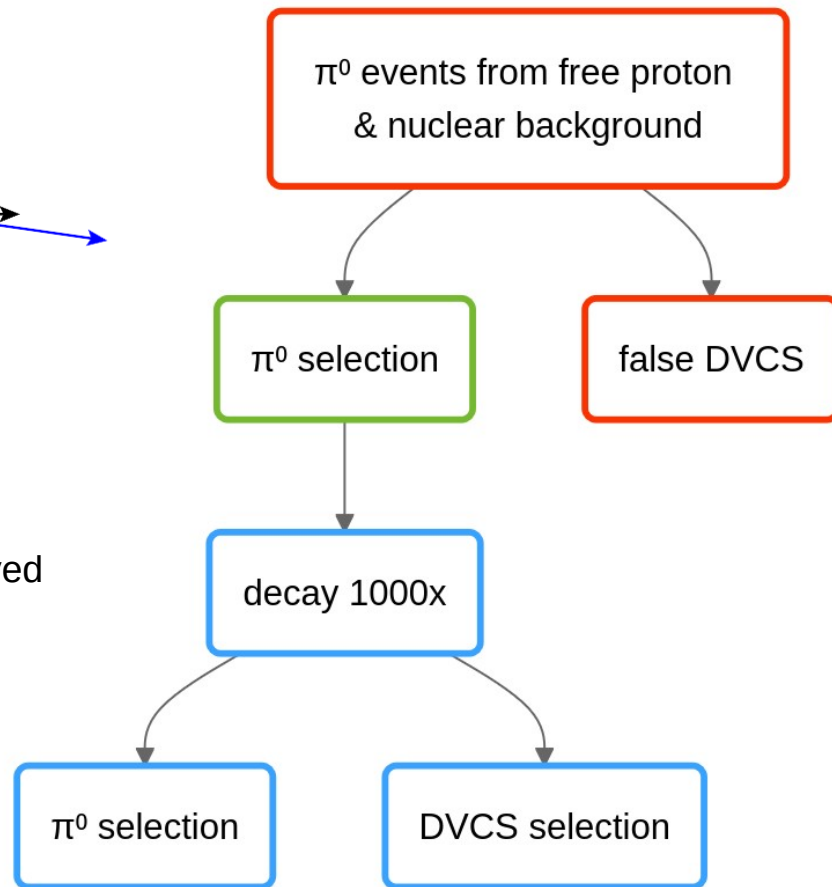
- Significant background that needs to be subtracted

$$N^{bt} = \frac{Y^{bt}}{\text{Charge}^{bt}} (1 - R^{bt})$$

- R is estimated from a data π^0 sample, each π^0 is decayed 1000x and passed through the CLAS12 simulation.

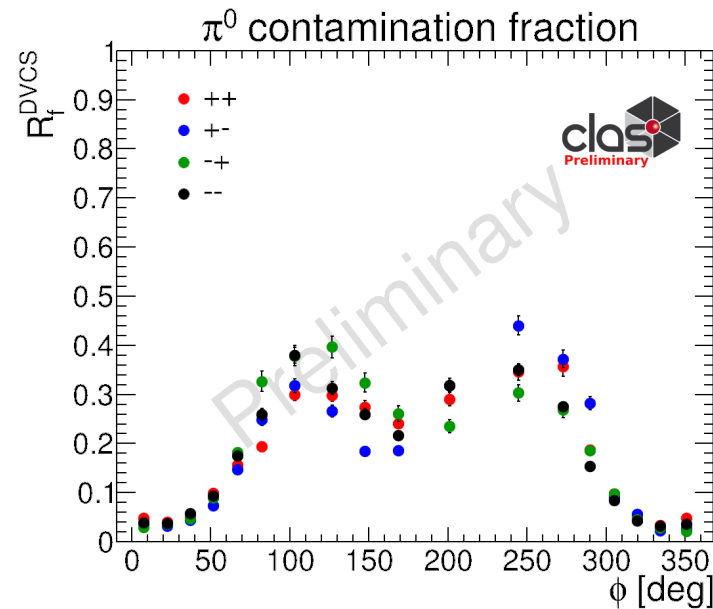
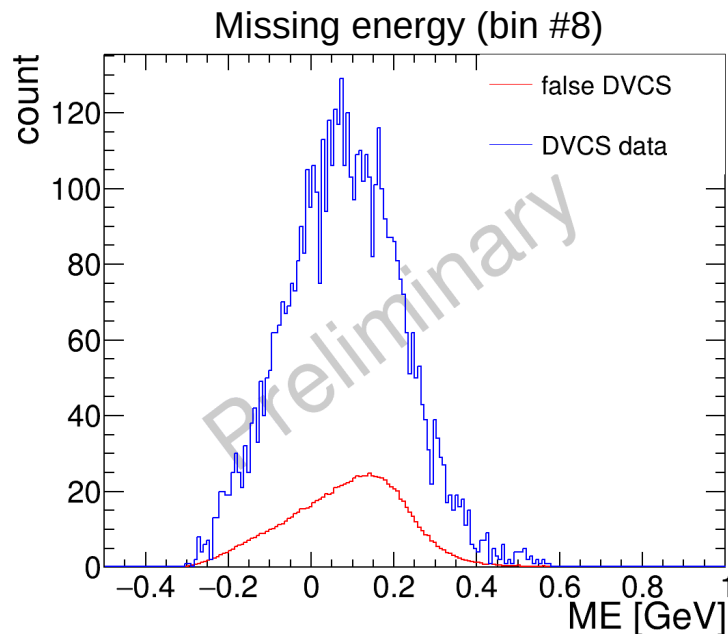
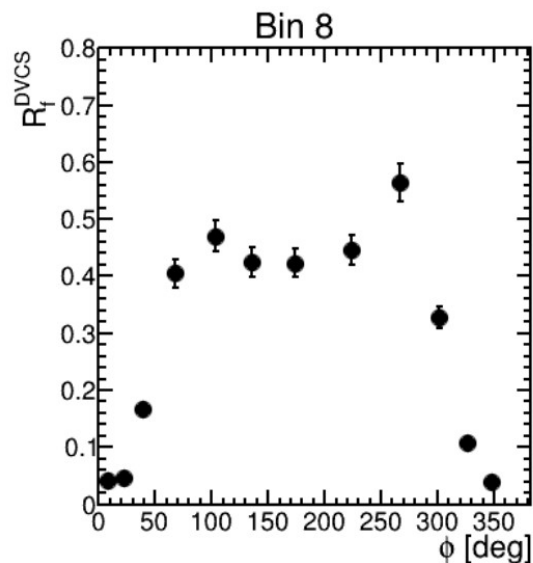
$$R = \frac{S^{\text{fDVCS}}}{S^{\pi^0}} \times \frac{Y^{\pi^0}}{Y^{\text{DVCS}}}$$

Y^i are the raw yields, S^i are the weighted number of decays that pass the DVCS or π^0 selection



π^0 contamination in data

- Contamination of up to 60% in some kinematic bins (largest contamination in bin #8)
- The contamination is larger in central ϕ bins where the Bethe-Heitler cross-section is small



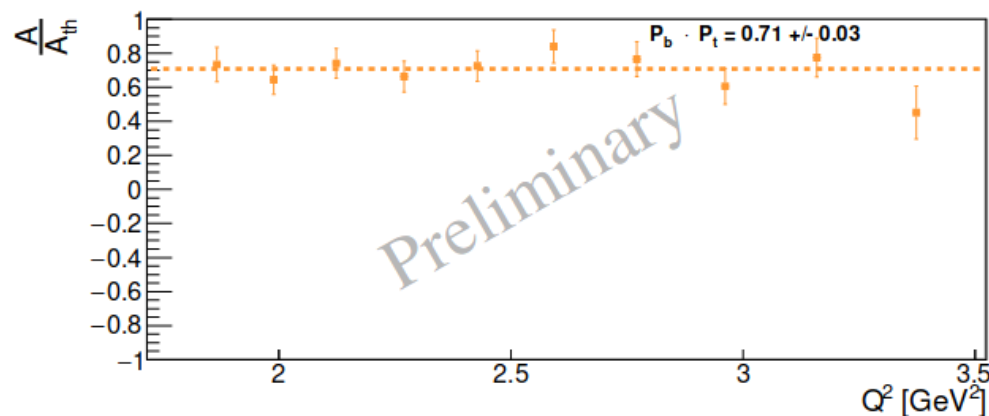
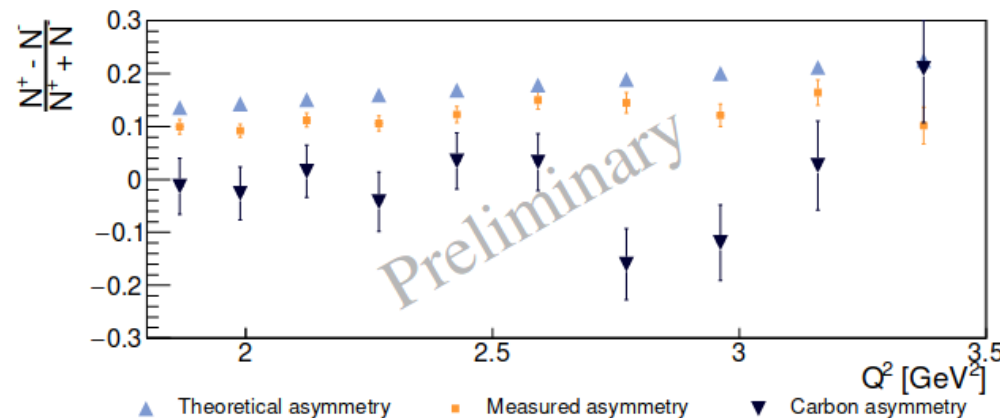
Target polarisation

N. Pilleux

- The $ep \rightarrow e'p'$ **elastic** double spin asymmetry is well known

$$A_{th} = \frac{A_{exp}}{P_b P_t} \quad A_{exp} = \frac{N^+ - N^-}{D_f(N^+ + N^-)}$$

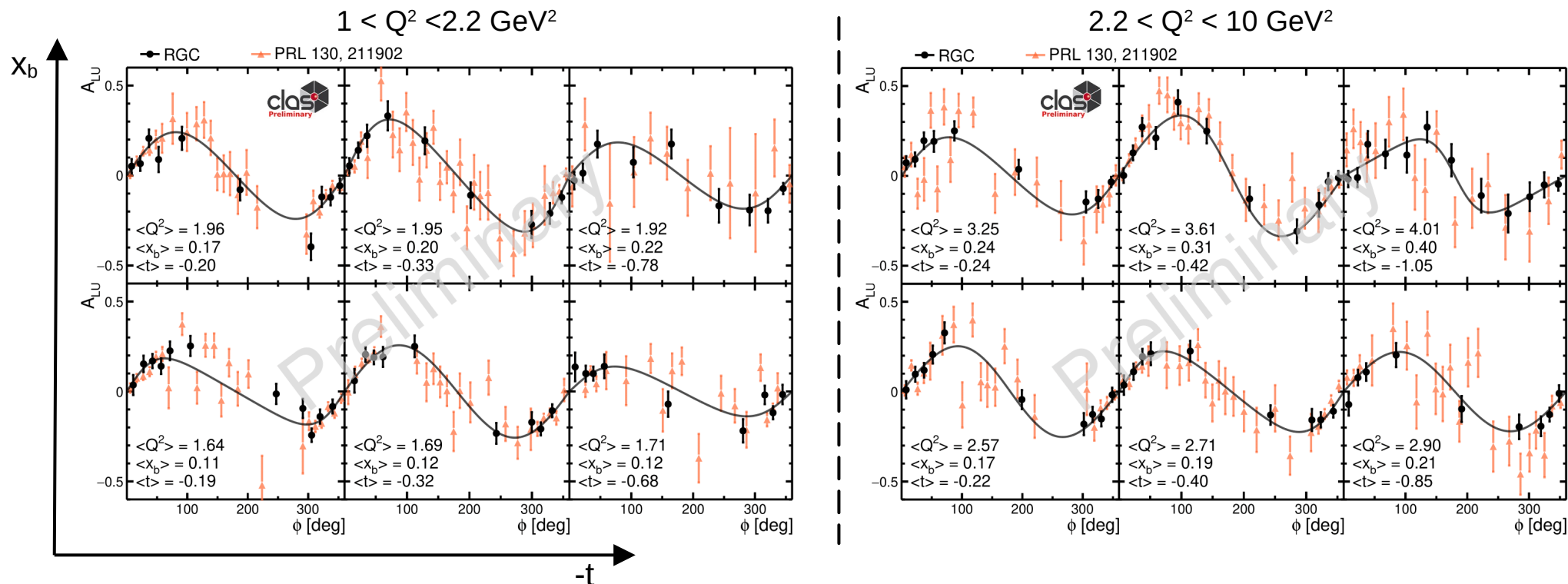
- The target polarisation is extracted from the ratio of theoretical to measured elastic asymmetry
 - $P_t^+ = 86 \pm 3 \%$
 - $P_t^- = 80 \pm 3 \%$
- Efforts are ongoing to cross-check elastic results with target polarisation measured from DIS asymmetries, and from MNR readings



Comparison with RGA BSA on unpolarised proton

- RGA BSA measurement on a liquid hydrogen target, [Phys. Rev. Lett. 130, 211902]
 - Same kinematic reach and no nuclear background → Good cross check

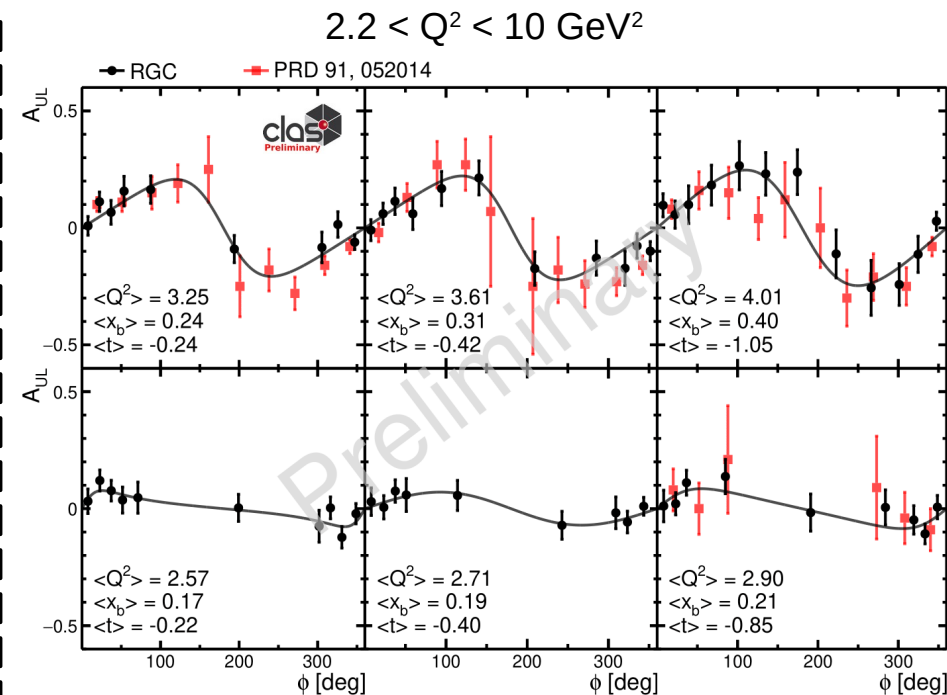
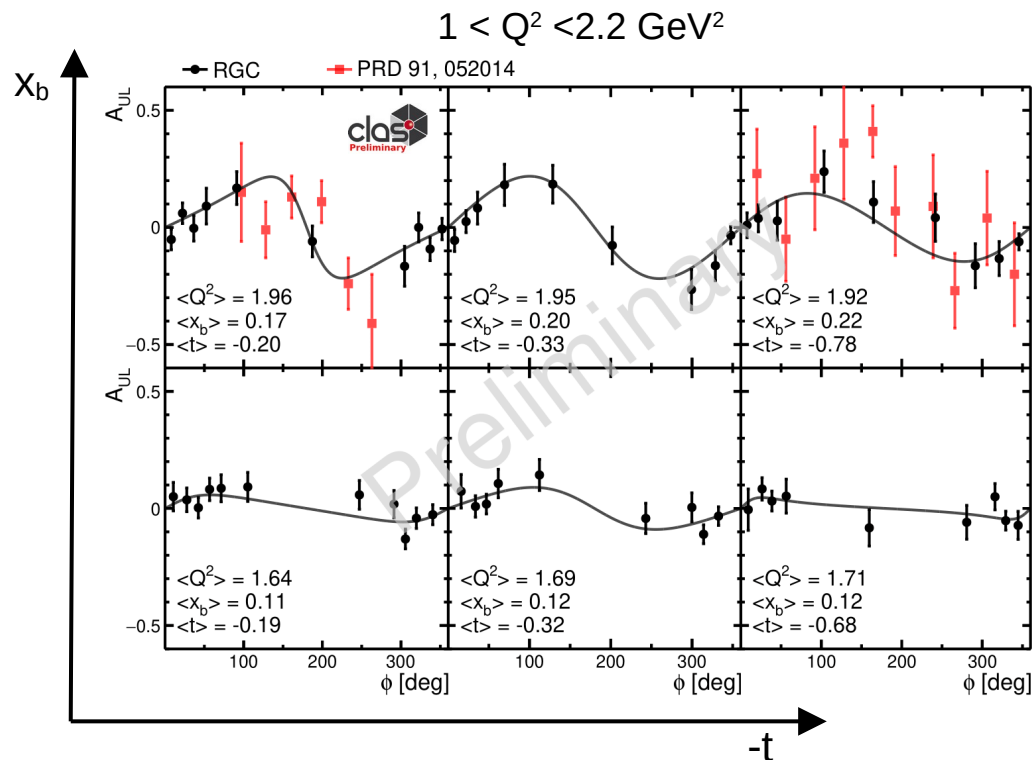
⚠ This work only includes part of the systematics uncertainties



Comparison with CLAS6 TSA

- CLAS6 eg1-dvcs TSA measurement [Phys. Rev. D 91, 052014] → Different beam energy and kinematic reach

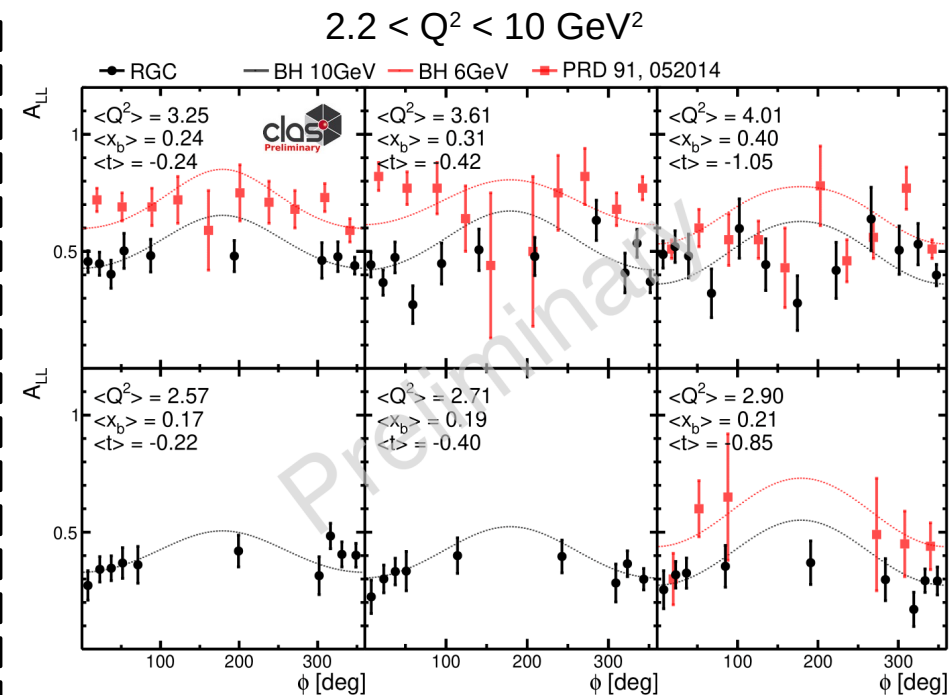
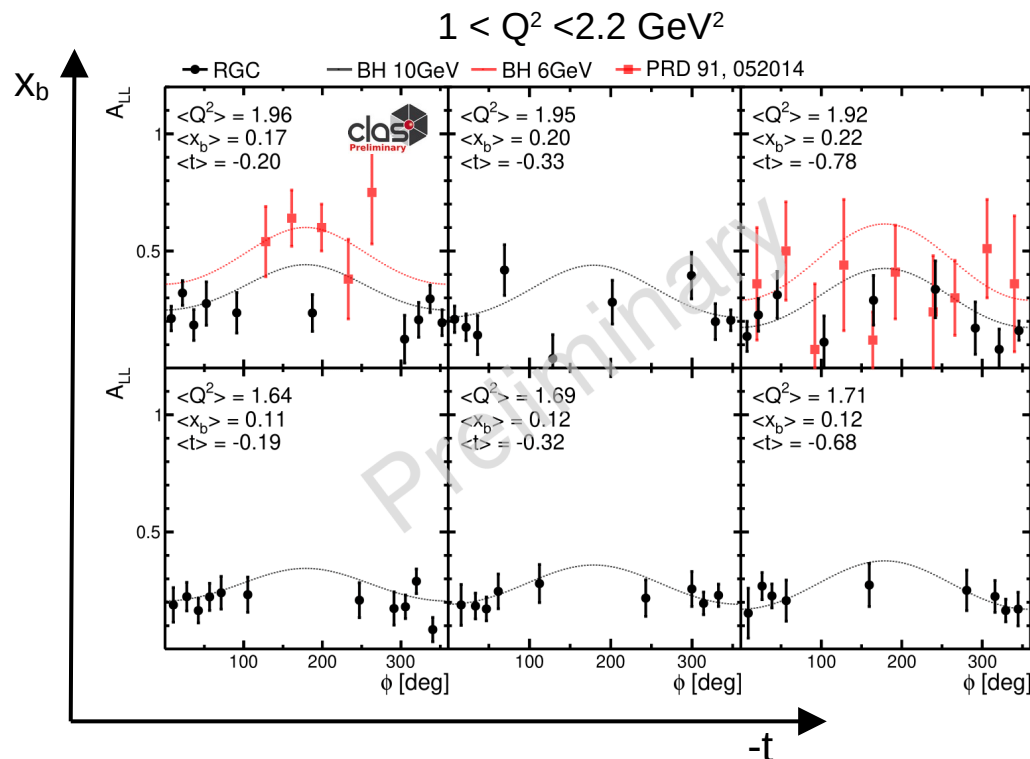
⚠ This work only includes part of the systematics uncertainties



Comparison with CLAS6 DSA

- The main contribution to the DSA is from the Bethe-Heitler asymmetry which differs with the beam energy

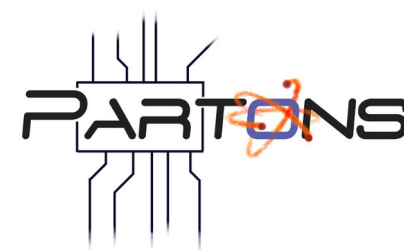
⚠ This work only includes part of the systematics uncertainties



Comparison with models

- In the next slides the asymmetries are compared to models from the [Partons](#) and [Gepard](#) libraries
- Three models/fits are compared:
 - KM15 implemented in Gepard
 - GK19 implemented in Partons
 - CFFNN, 2019 CFF fit with a neural network parametrisation implemented in Partons. The comparison is restricted to bins with $-t/Q^2 < 0.2$, where the fit is constrained
- The model points are computed at the average kinematics of each ϕ bin, and over 100 ϕ points at the average kinematics of the Q^2 , x_b , t bin to have a continuous line.

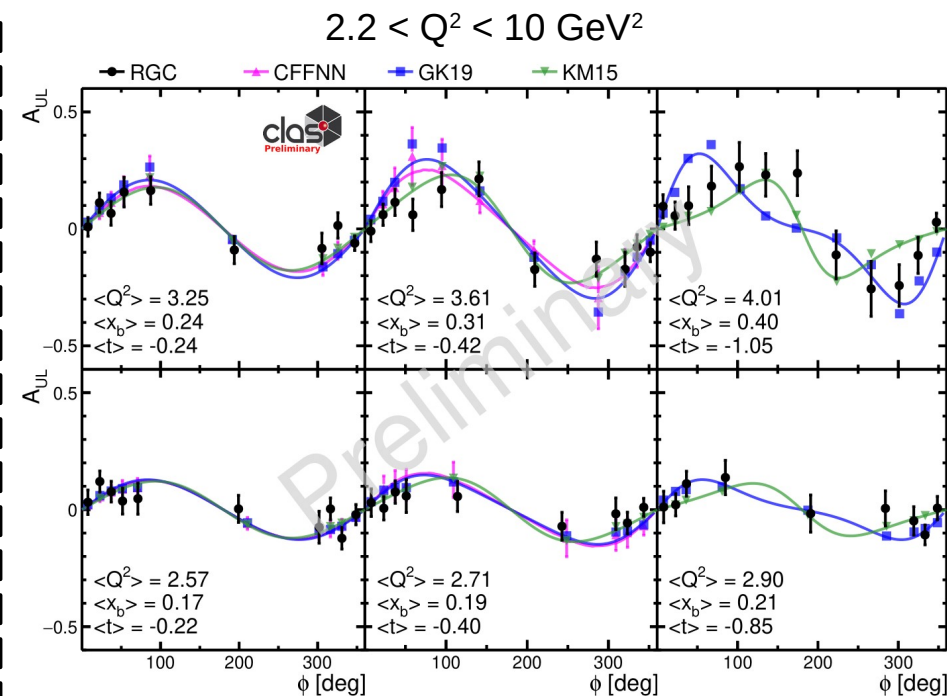
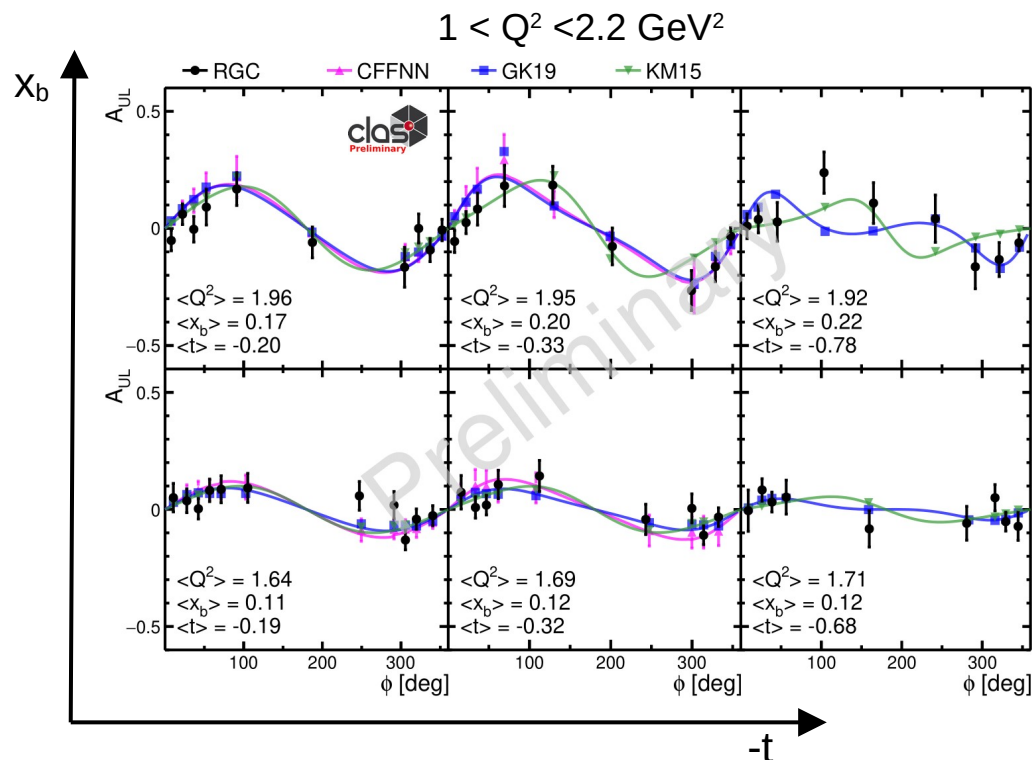
Gepard



Thanks to K. Kumerički and P. Sznajder for their help !

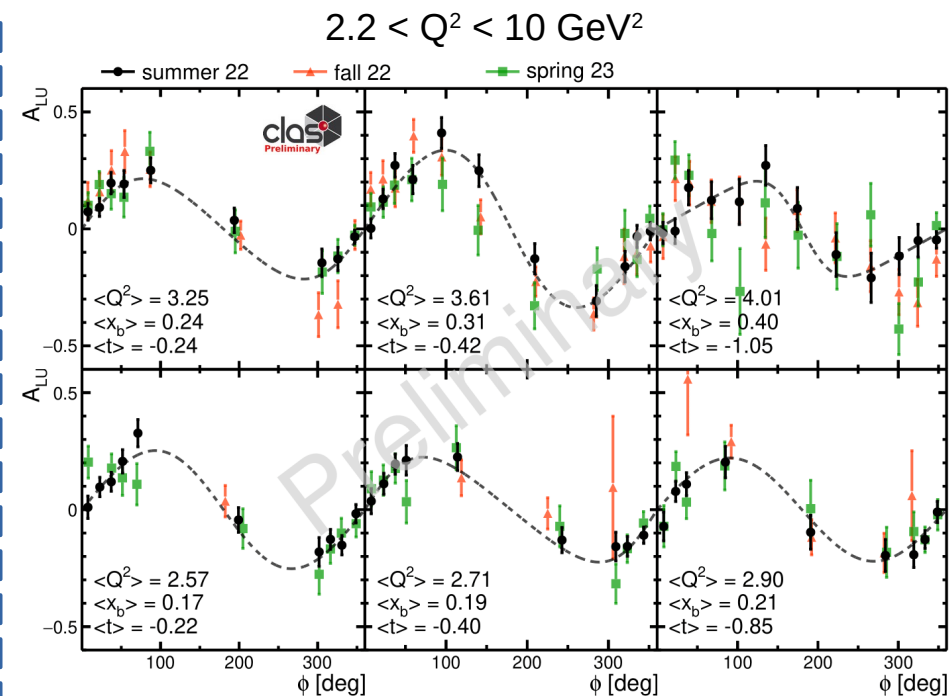
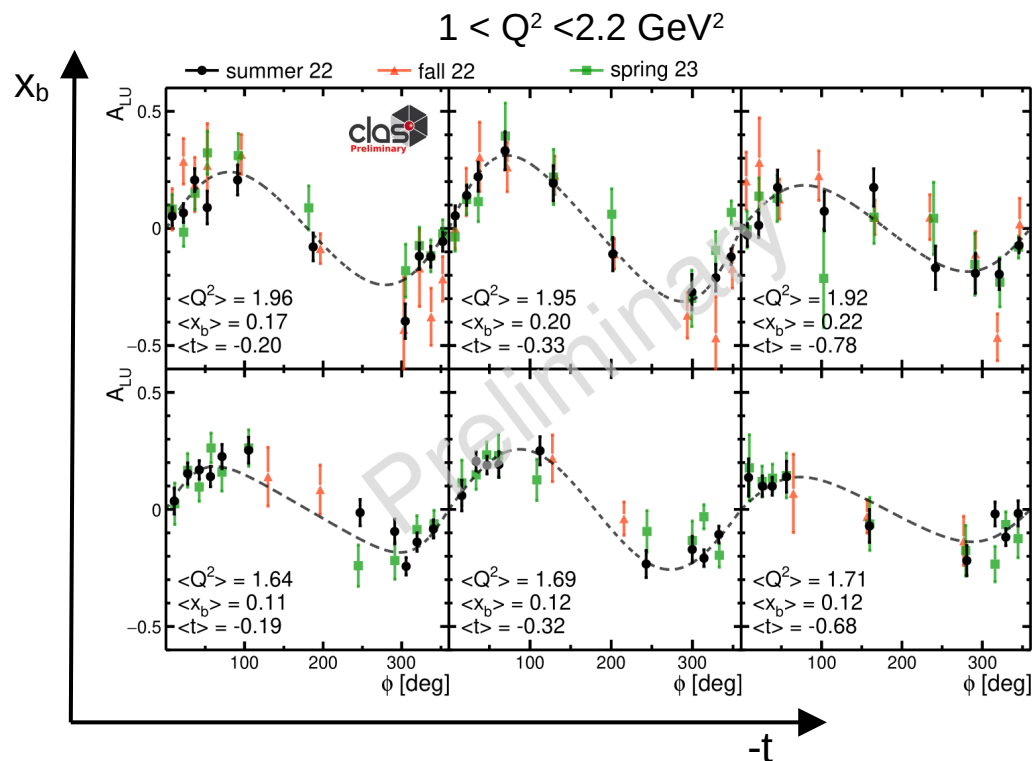
Comparison with models, TSA

- The model points are computed at the average kinematics of each specific phi bin, and the model line is computed at the kinematics of the Q^2 , x_b , t bin

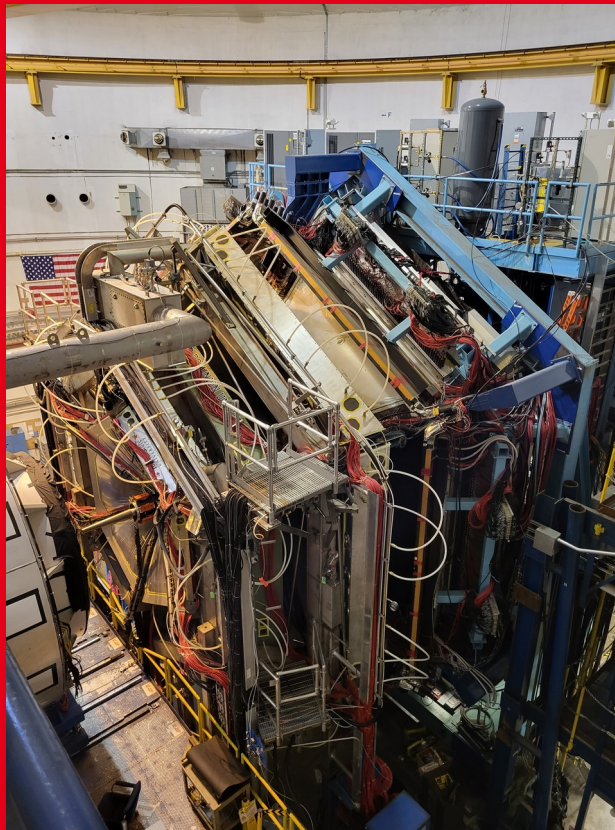


First look at the Fall 22 and Spring 23 BSA

- The QADB timelines are not available for Fall 22 and Spring 23 and no careful run selection has been done yet
- Fall 22 in the FTOFF configuration → different kinematic reach



Summary & outlook



- The analysis of DVCS spin asymmetries on the polarised proton with RGC data is well advanced
- Tools are in place to take into account nuclear and π^0 backgrounds
- Preliminary results are consistent with previous measurements by the CLAS collaboration
- The main remaining steps are:
 - Improvement of the momentum correction
 - Careful systematic uncertainty evaluation
 - Study how the different run period could be merged

Analysis overview

To measure the BSA, TSA, and DSA we need the following elements:

- **Corrected DVCS yields N^{bt}** in each **b**eam and **t**arget spin configurations
 - Selection of DVCS candidates
 - Subtract π^0 contamination
- **Dilution factor D_f** ,
 - Accounts for the nuclear background in the NH3 target
- **Target polarisation P_t**
 - Measured from elastic scattering
- **Beam polarisation P_b**
 - Measured with a Moeller polarimeter in the experimental hall, $P_b = 82.5 \pm 0.2\%$

$$A_{LU} = \frac{1}{P_b} \frac{P_t^- (N^{++} - N^{-+}) + P_t^+ (N^{+-} - N^{--})}{P_t^- (N^{++} + N^{-+}) + P_t^+ (N^{+-} + N^{--})}$$

$$A_{UL} = \frac{1}{D_f} \frac{N^{++} + N^{-+} - N^{+-} - N^{--}}{P_t^- (N^{++} + N^{-+}) + P_t^+ (N^{+-} + N^{--})}$$

$$A_{LL} = \frac{1}{D_f P_b} \frac{N^{++} + N^{--} - N^{+-} - N^{-+}}{P_t^- (N^{++} + N^{-+}) + P_t^+ (N^{+-} + N^{--})}$$

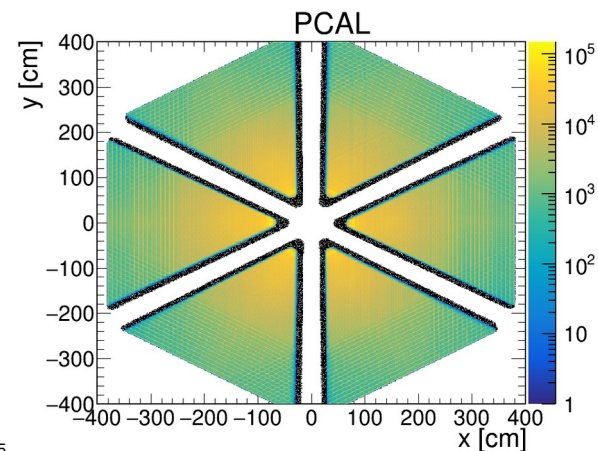
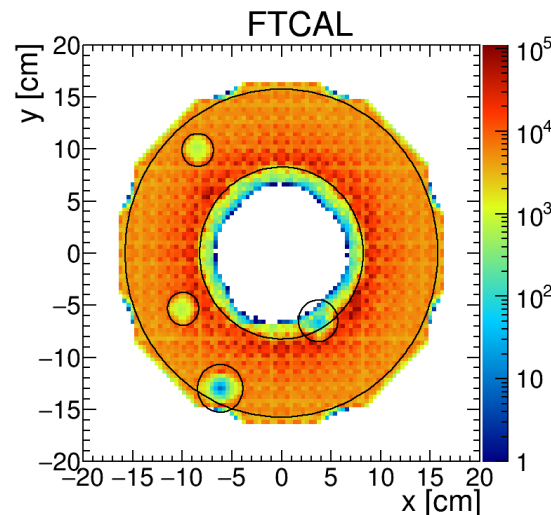
Fiducial cuts

- DC fiducial cuts on the edge variable

Torus polarity	Region 1 [cm]	Region 2 [cm]	Region 3 [cm]
Inbending	4	3	7
Outbending	3	3	10

- PCAL, medium fiducial cuts, $l_u > 29$, $l_v, l_v > 14$ cm
- FTCAL, cut using the the x,y coordinate of the track at the surface of FTCAL, $8.25 < \sqrt{x^2+y^2} < 15.75$ cm

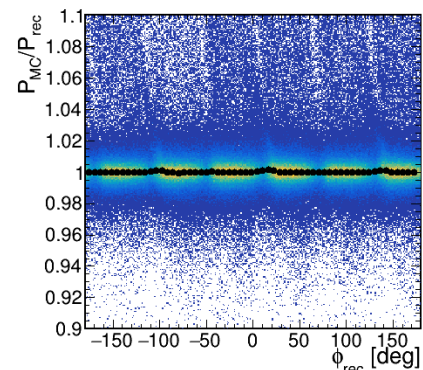
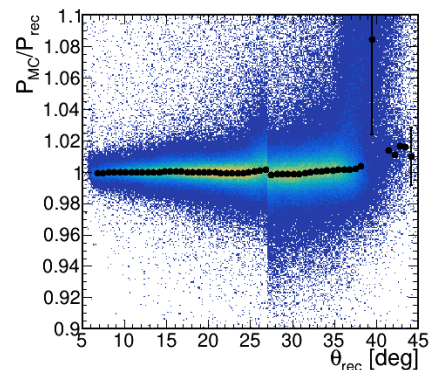
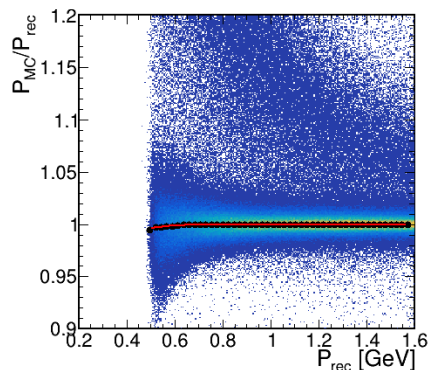
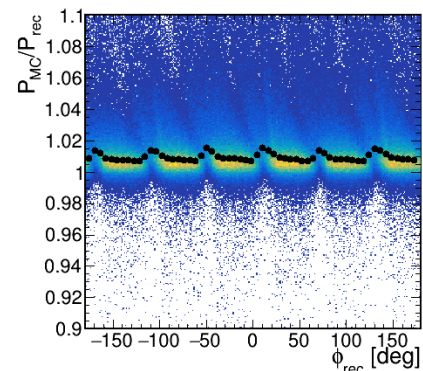
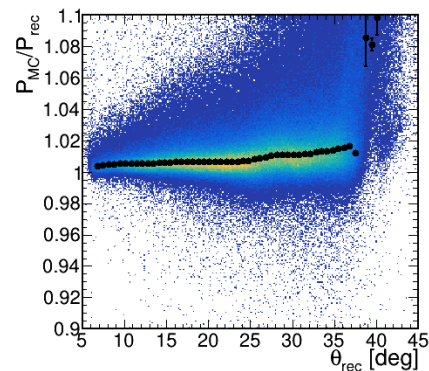
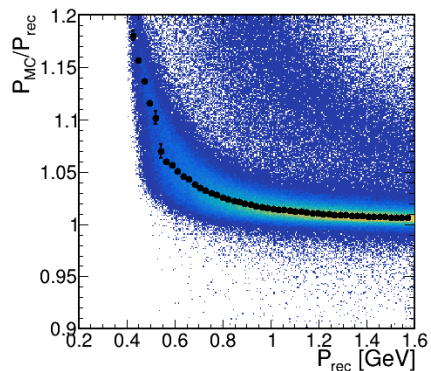
Centre x [cm]	Centre y [cm]	Radius [cm]
-8.42	9.89	1.6
-9.89	-5.33	1.6
-6.15	-13	2.3
3.7	-6.5	2



Energy loss corrections

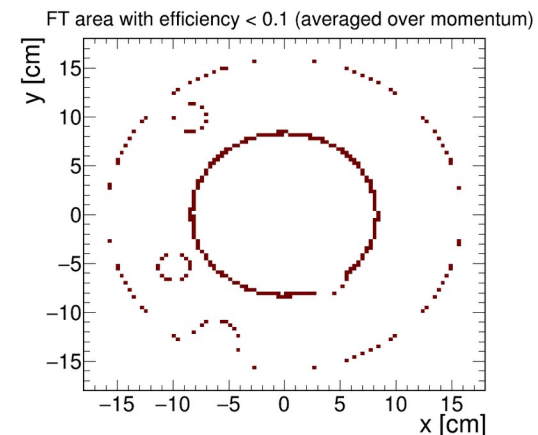
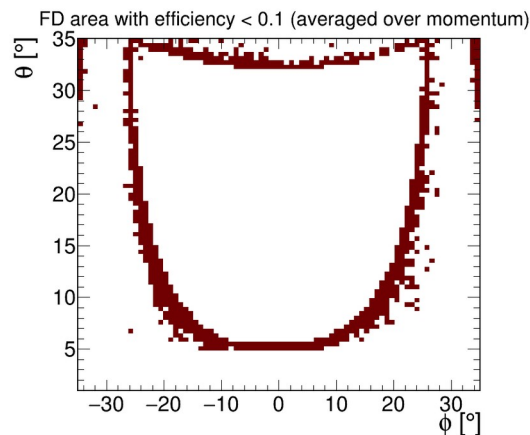
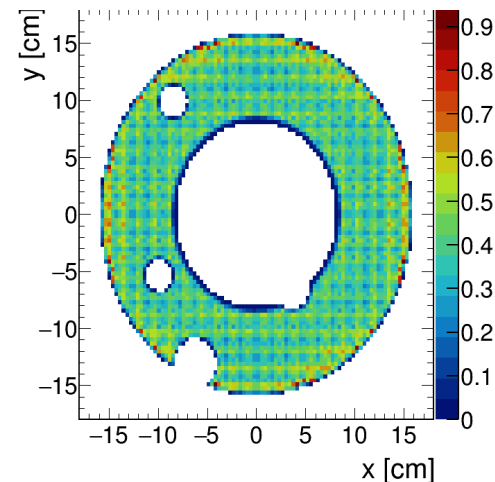
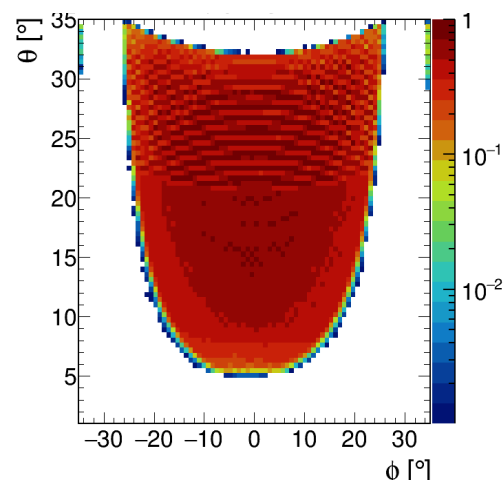
- Monte-Carlo based momentum corrections on the electron and proton
- Small correction for the electron and proton in the CD (<1%)
- Proton in FD, correction split in two θ region at 27° to account for the difference in material crossed

$$f(p) = a + b/p + c/p^2 + d/p^3$$



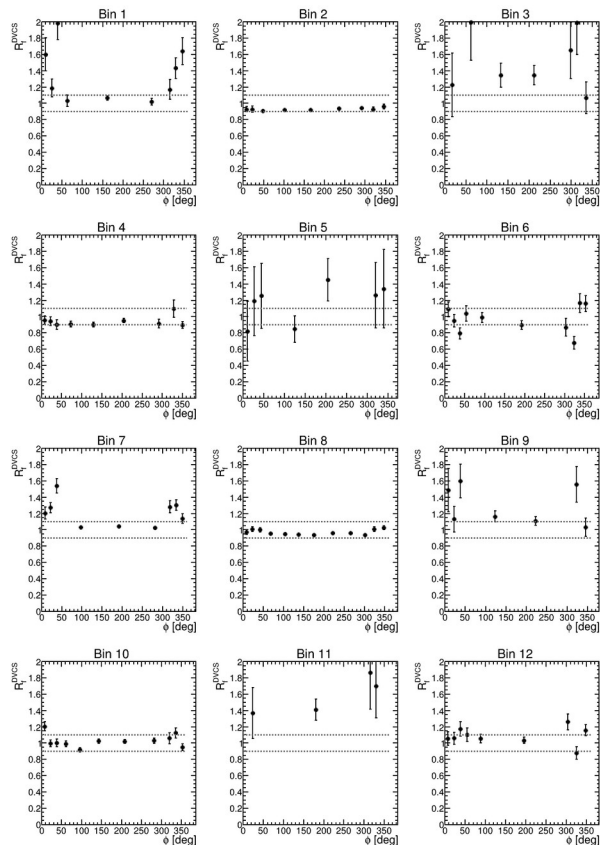
π^0 acceptance maps

- The DVCS analysis needs to be restricted to areas of the detector where π^0 can be measured so that the contamination can be estimated
- Flat generation of events of π^0 decays between 0 and 10 GeV over FD sector 1 and FT
- Count the number of well reconstructed π^0 and generated π^0 in bins of the leading photon position and momentum
- If a DVCS photon is in a bin where the leading photon of a π^0 decay has a less than 10% chance being detected, the event is cut
- Survival rate of DVCS events: 97%

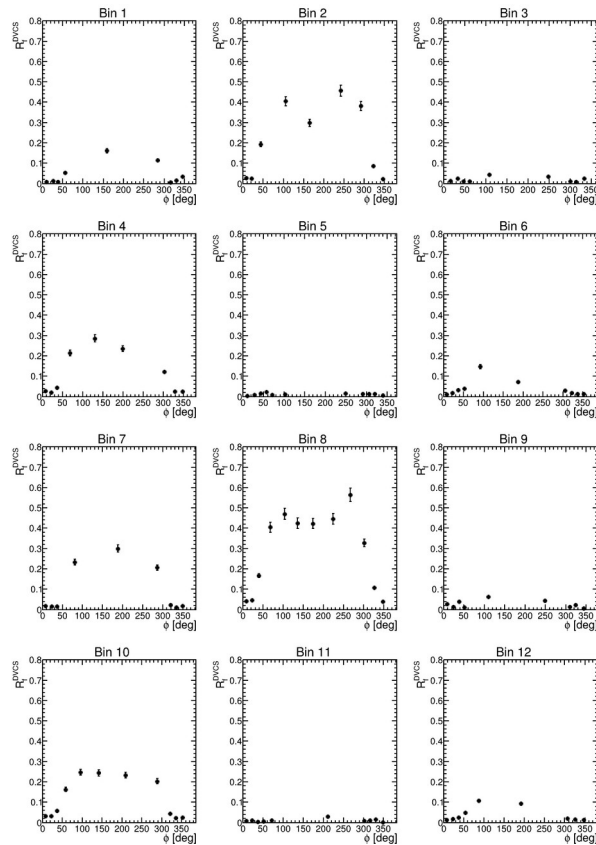


Contamination fraction in data

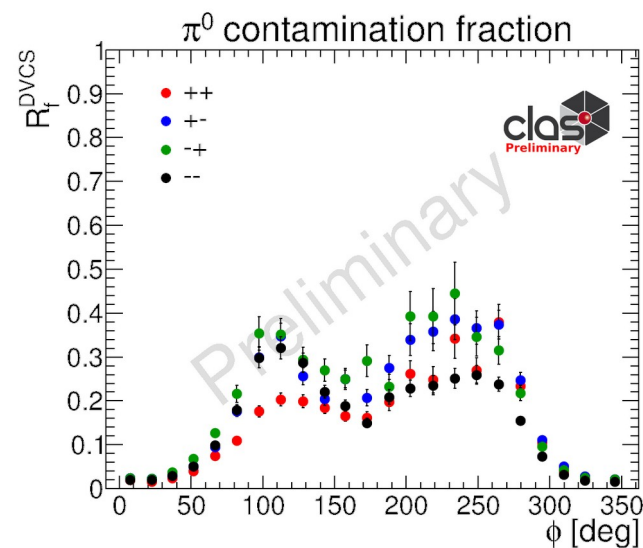
Validation



Data



Contamination fraction in data,
integrated over Q^2 , x_b , t



Asymmetries and uncertainties

- Two sources of uncertainty are included:
 - Statistical uncertainty on the DVCS yield
 - Statistical uncertainty in the π^0 yield used in the π^0 subtraction
- No systematic uncertainty on the dilution factor, on the target/beam polarisation or on the π^0 “acceptance ratio” is included

$$A_{LU} = \frac{1}{P_b} \frac{P_t^- (N^{++} - N^{-+}) + P_t^+ (N^{+-} - N^{--})}{P_t^- (N^{++} + N^{-+}) + P_t^+ (N^{+-} + N^{--})}$$

$$A_{UL} = \frac{1}{D_f} \frac{N^{++} + N^{-+} - N^{+-} - N^{--}}{P_t^- (N^{++} + N^{-+}) + P_t^+ (N^{+-} + N^{--})}$$

$$A_{LL} = \frac{1}{D_f P_b} \frac{N^{++} + N^{--} - N^{+-} - N^{-+}}{P_t^- (N^{++} + N^{-+}) + P_t^+ (N^{+-} + N^{--})}$$

$$N = \frac{1}{FC} \left(Y^D - \frac{S^{\text{fD}}}{S^{\pi^0}} Y^{\pi^0} \right)$$

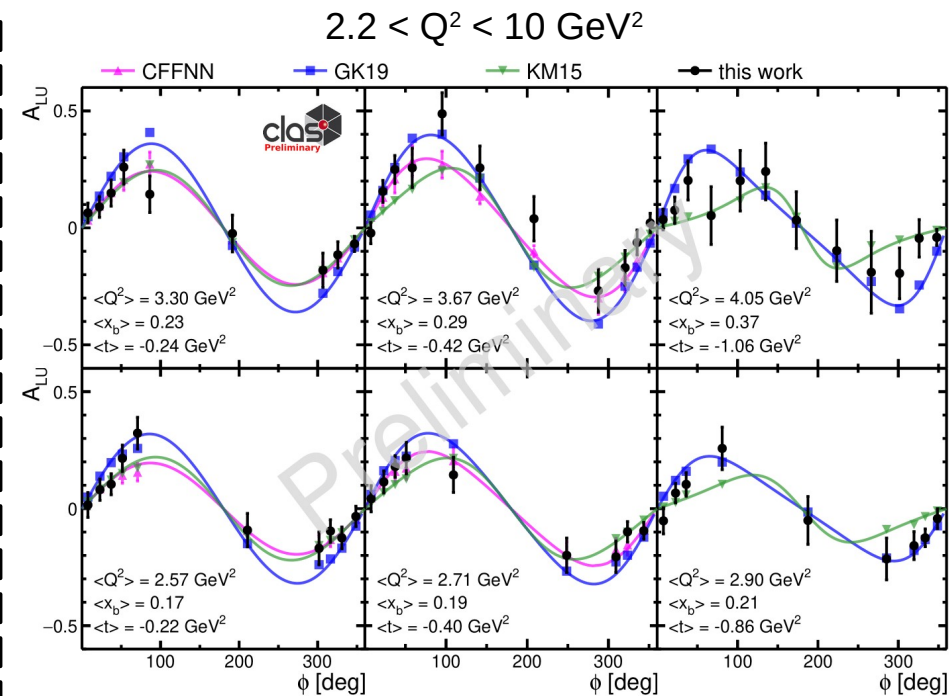
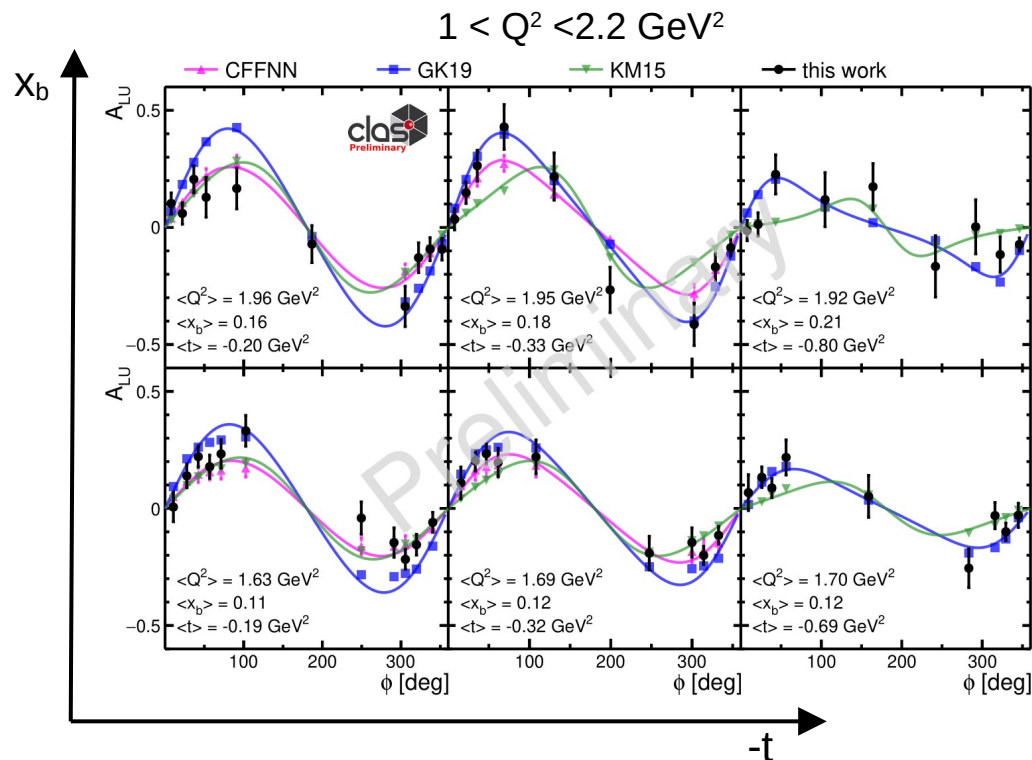
$$\sigma_N = \frac{1}{FC} \sqrt{\left(\frac{\partial N}{\partial Y^D} \right)^2 \sigma_{Y^D}^2 + \left(\frac{\partial N}{\partial Y^{\pi^0}} \right)^2 \sigma_{Y^{\pi^0}}^2}$$

$$\sigma_N = \frac{1}{FC} \sqrt{\sigma_{Y^D}^2 + \left(\frac{S^{\text{fD}}}{S^{\pi^0}} \right)^2 \sigma_{Y^{\pi^0}}^2}$$

$$\sigma A = \sqrt{\left(\frac{\partial A}{\partial N^{++}} \right)^2 \sigma_{N^{++}}^2 + \left(\frac{\partial A}{\partial N^{+-}} \right)^2 \sigma_{N^{+-}}^2 + \left(\frac{\partial A}{\partial N^{-+}} \right)^2 \sigma_{N^{-+}}^2 + \left(\frac{\partial A}{\partial N^{--}} \right)^2 \sigma_{N^{--}}^2}$$

Comparison to models, BSA

- The model points are computed at the average kinematics of each specific phi bin, and the model line is computed at the kinematics of the Q^2 , x_b , t bin

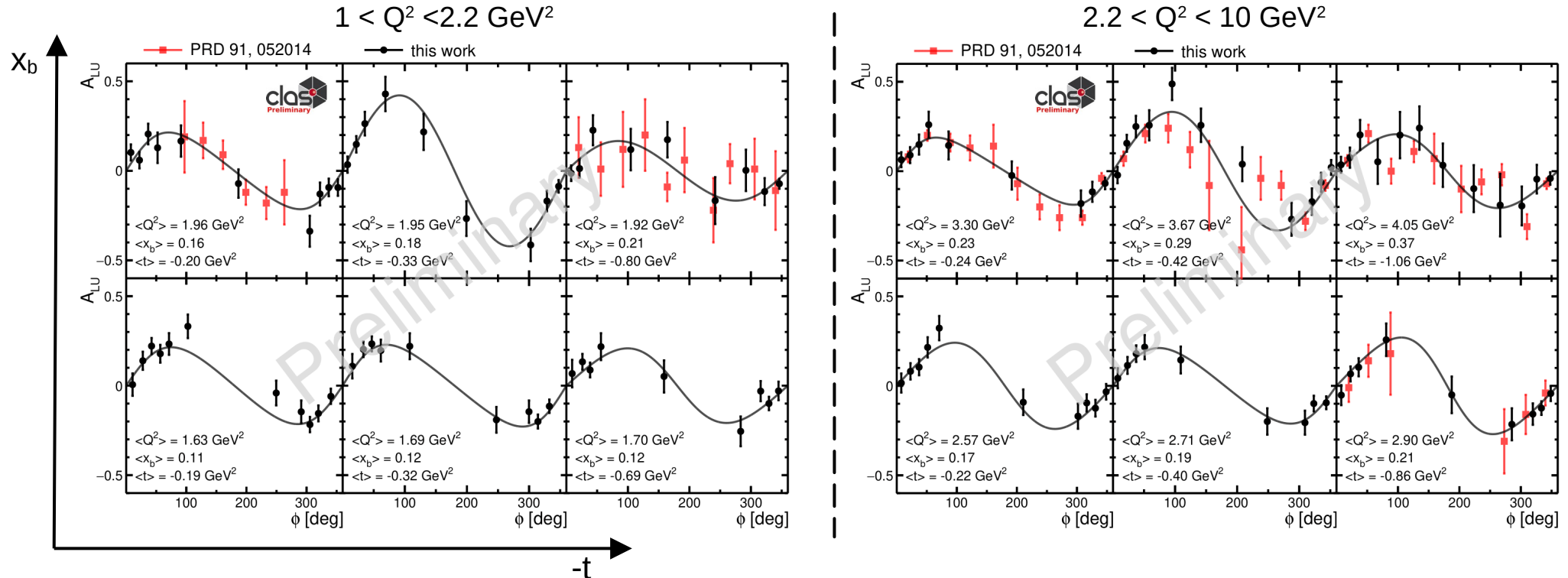


Comparison to CLAS6 BSA

- CLAS6 BSA measurement [Phys. Rev. D 91, 052014] at JLab 6GeV → Different beam energy and kinematic reach



The published results include systematic uncertainties and this work only includes statistical uncertainties

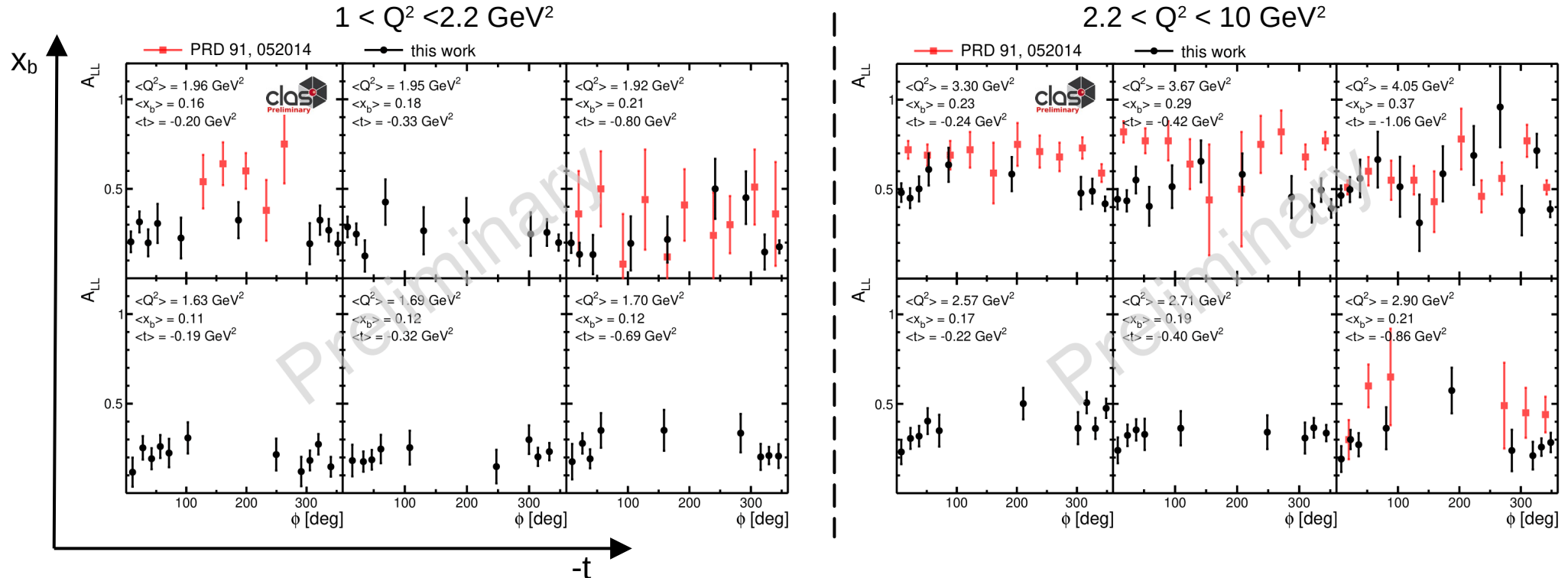


Comparison to CLAS6 DSA

- CLAS6 DSA measurement [Phys. Rev. D 91, 052014] at JLab 6GeV → Different beam energy and kinematic reach

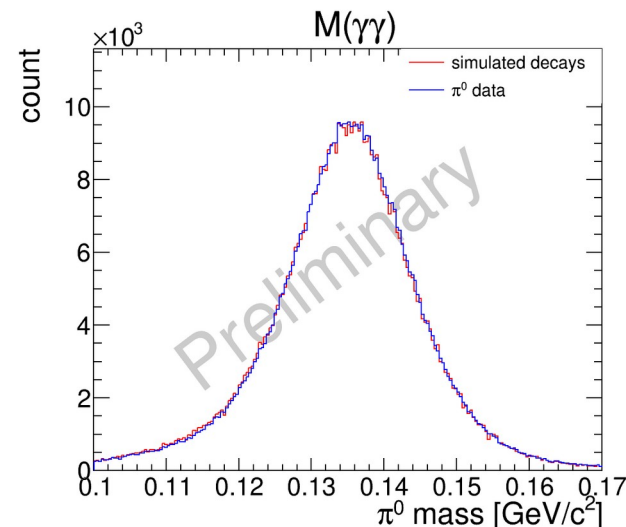
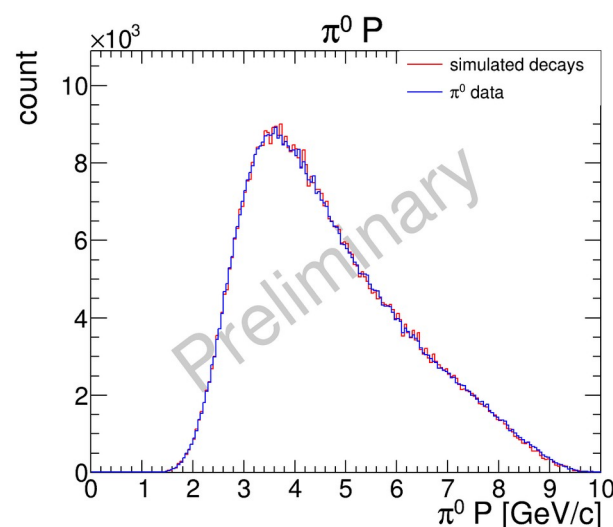
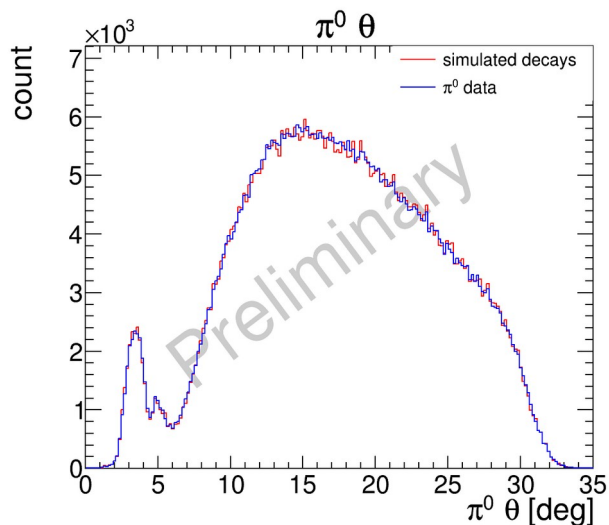
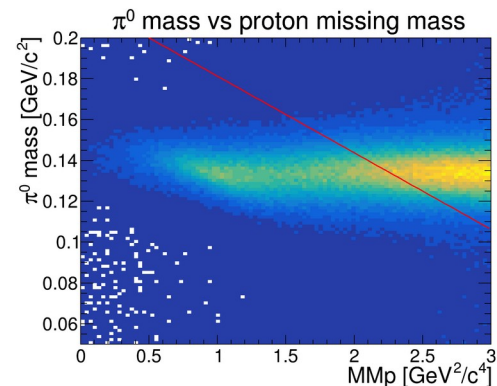


The published results include systematic uncertainties and this work only includes statistical uncertainties



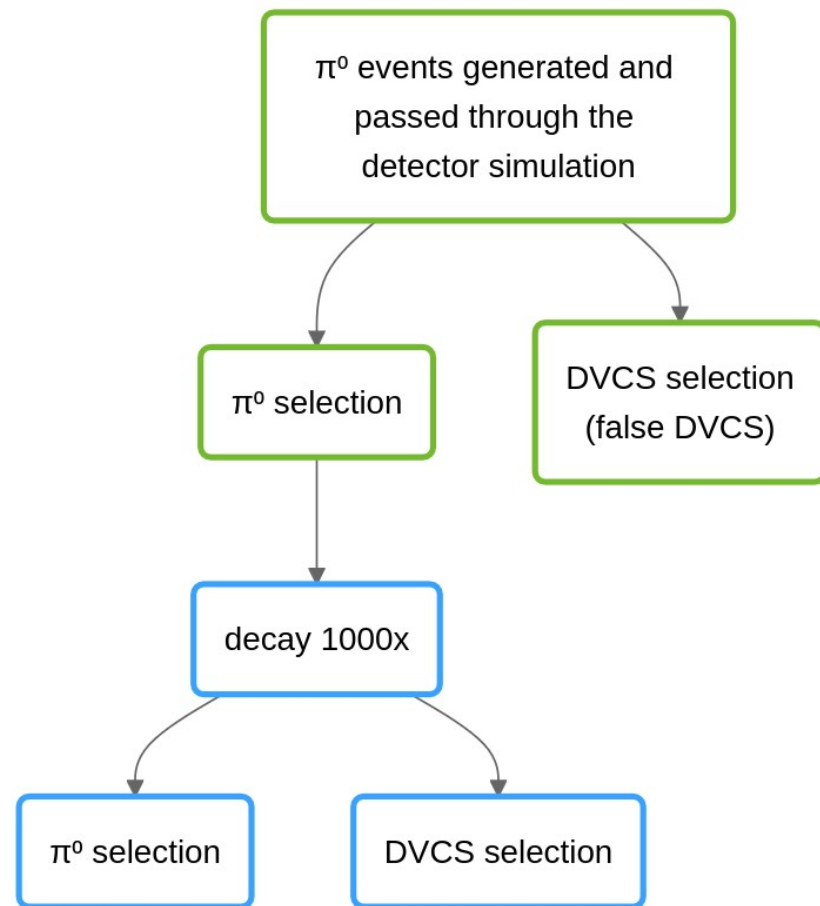
π^0 selection & decays

- Sample is selected with broad exclusivity cuts, to also include events from nuclear background and semi-inclusive channels
- Each π^0 is then decayed 1000x, decays are passed through the CLAS12 simulation software and weighted
- The π^0 selection is applied to the decays and weighted. The simulated decays and data distributions match.



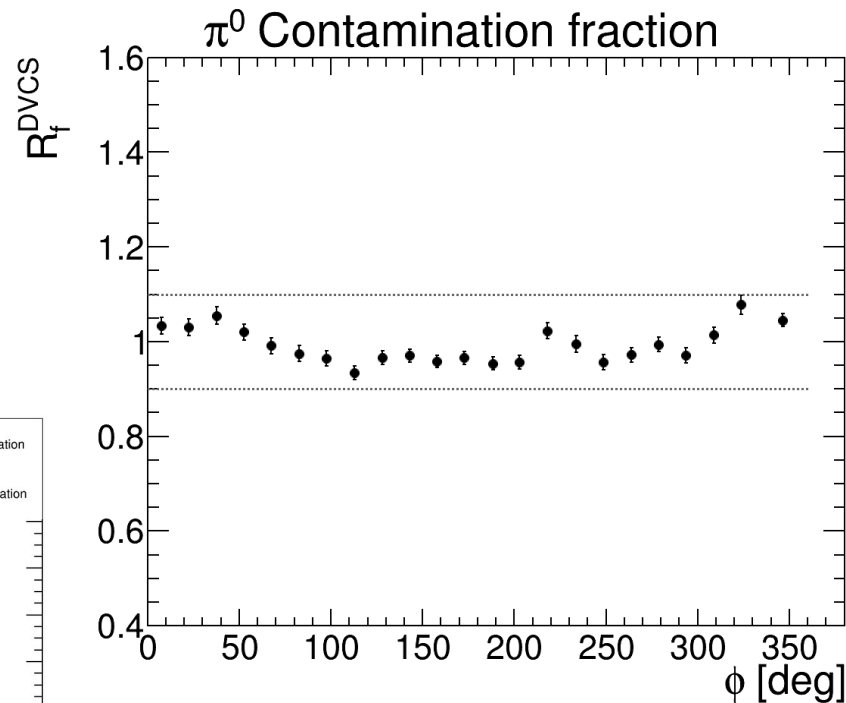
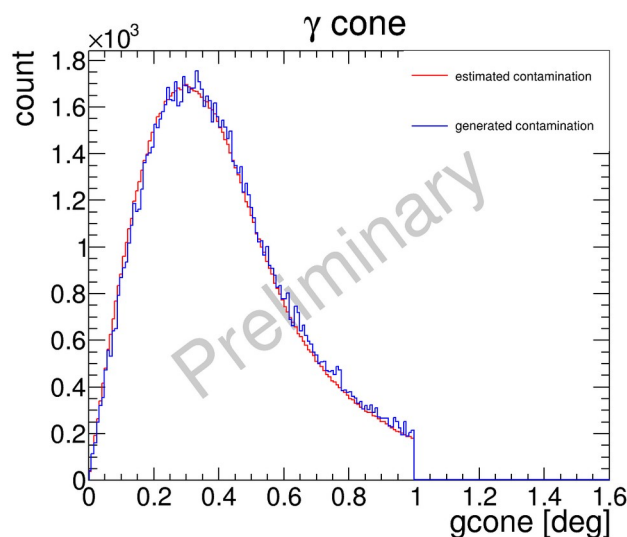
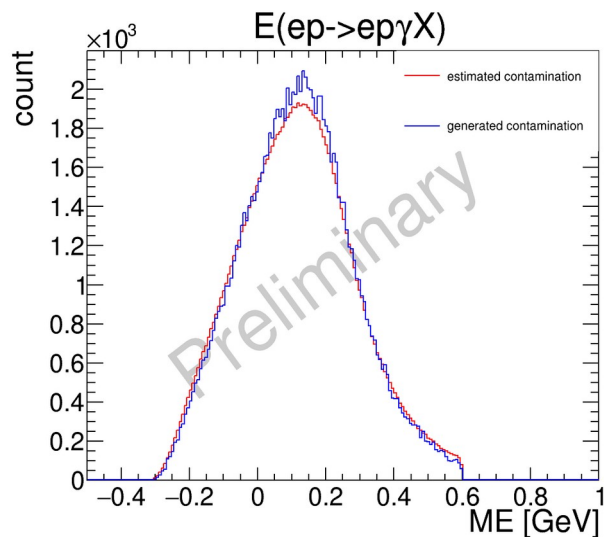
π^0 subtraction validation

- Apply the method to a π^0 simulation
- All events selected as DVCS are π^0 contamination
 - Contamination fraction of a 100%
- Check if the subtraction method is able to reproduce the generated false DVCS background



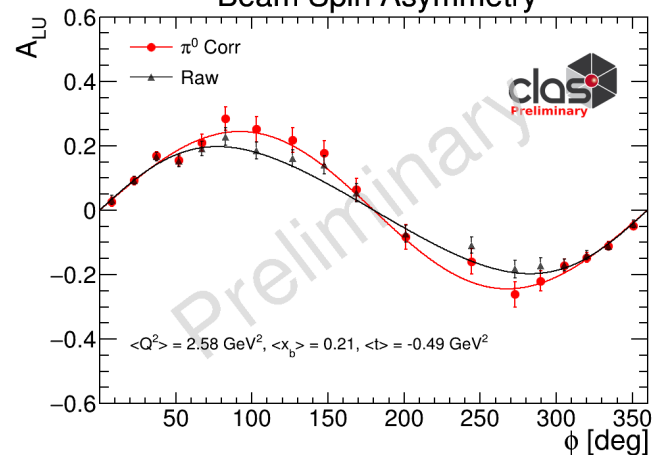
π^0 subtraction validation

- The generated π^0 contamination is recovered
- The false DVCS distribution estimated with the method matches the generated contamination well

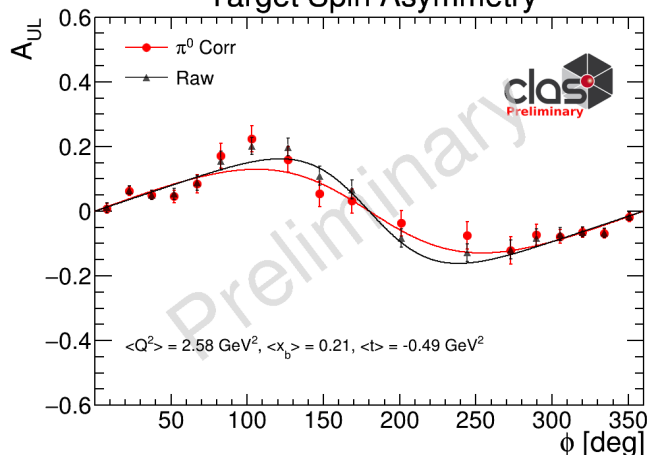


Integrated asymmetries

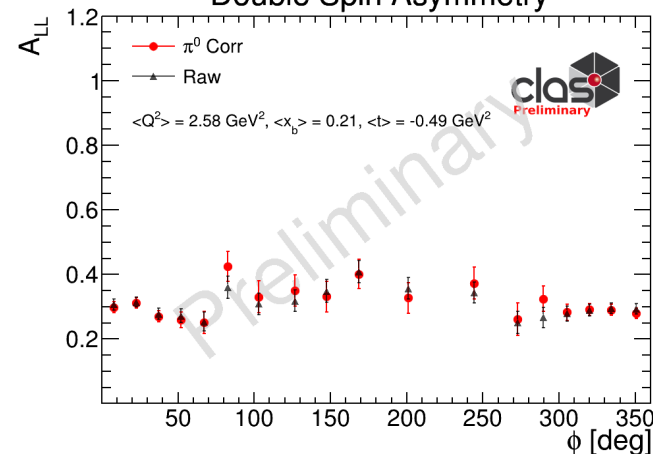
Beam Spin Asymmetry



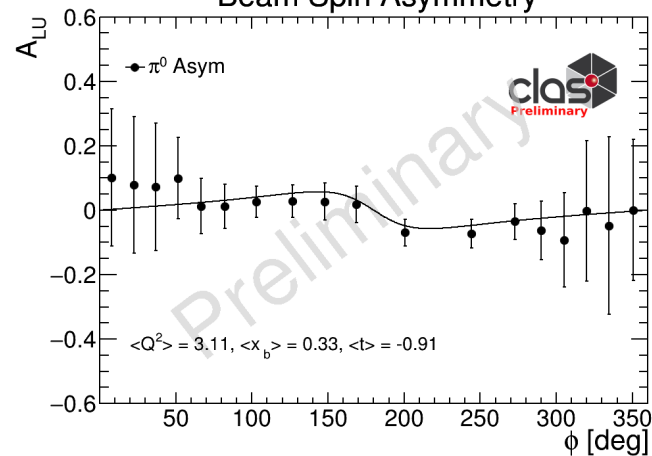
Target Spin Asymmetry



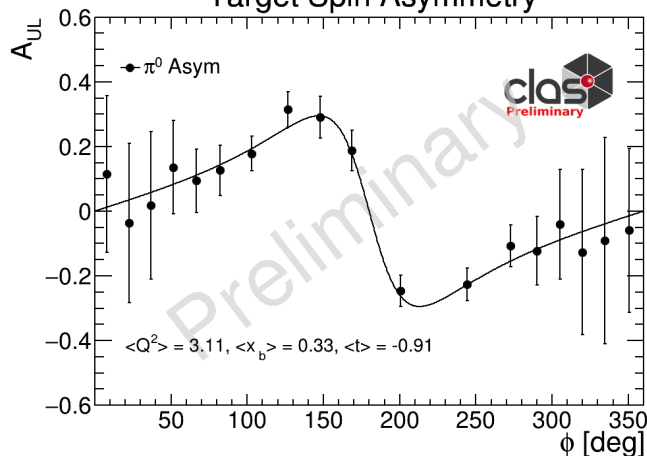
Double Spin Asymmetry



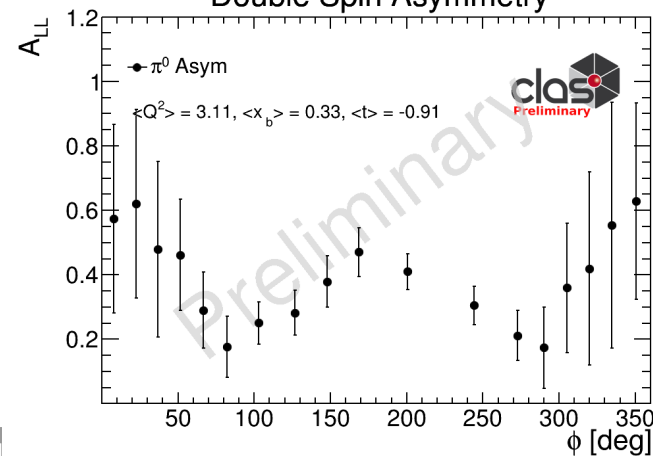
Beam Spin Asymmetry



Target Spin Asymmetry

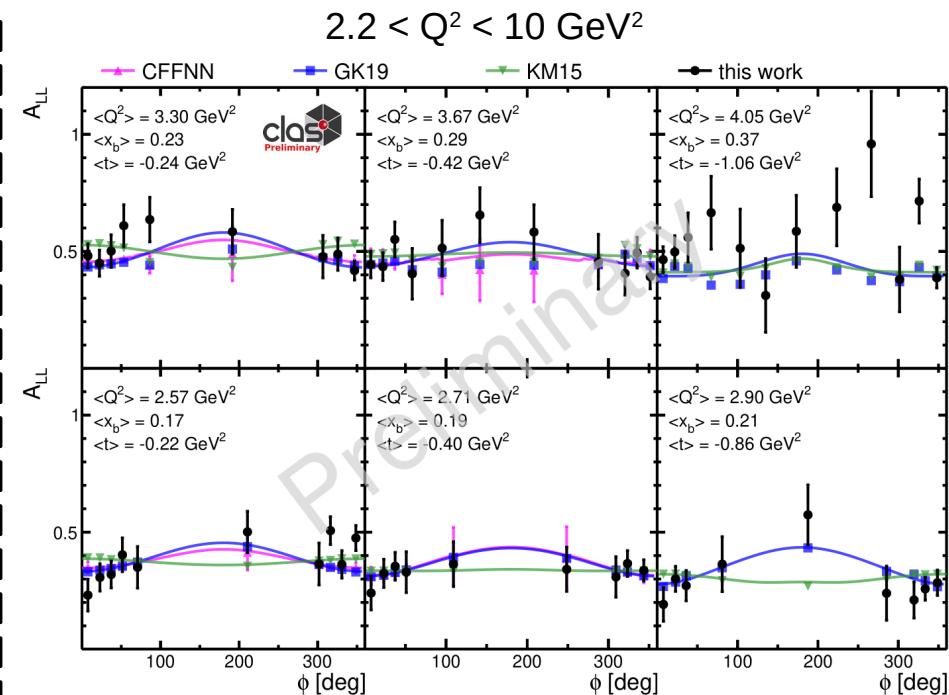
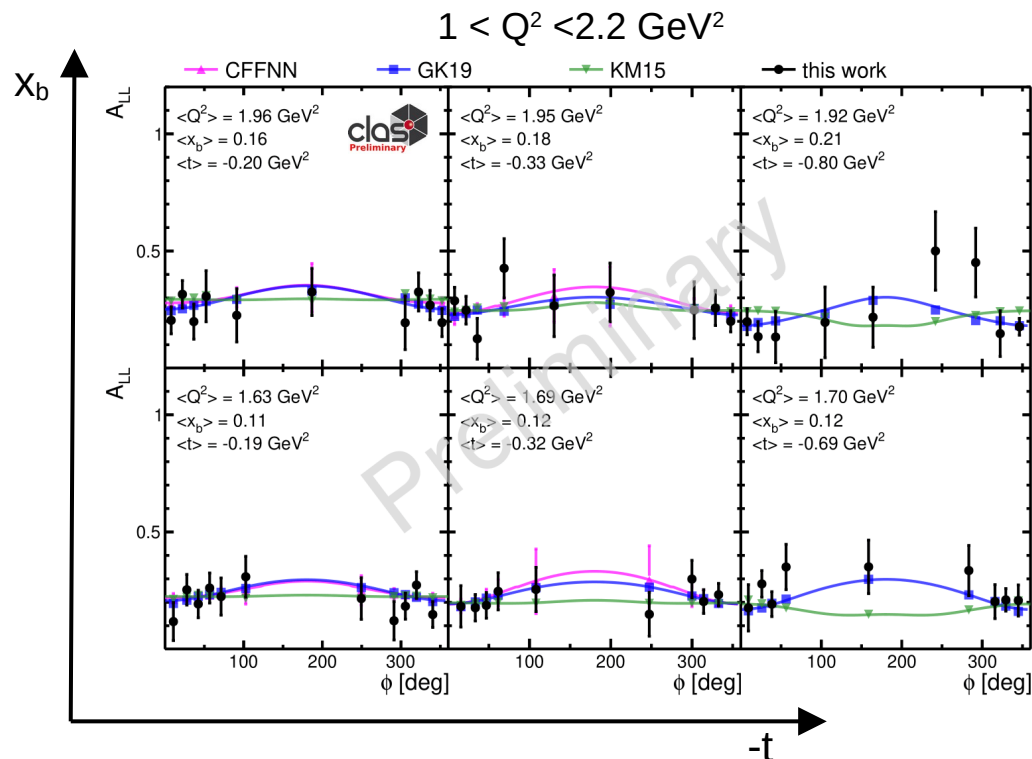


Double Spin Asymmetry

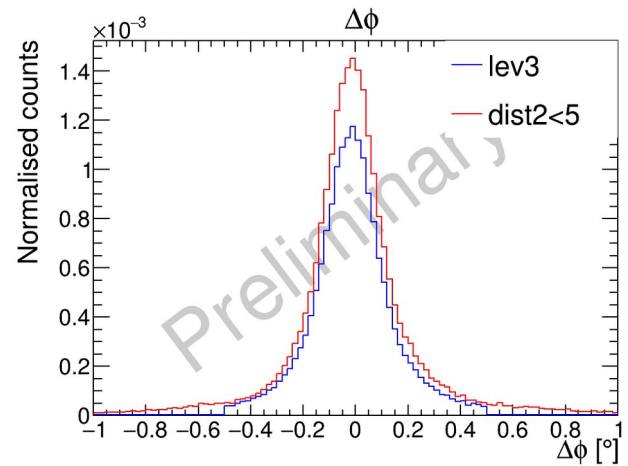
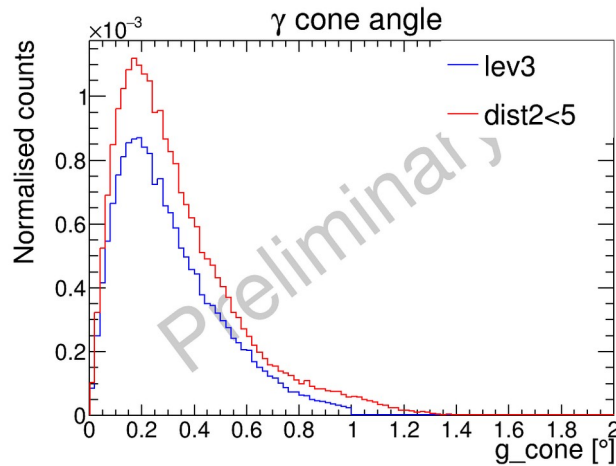
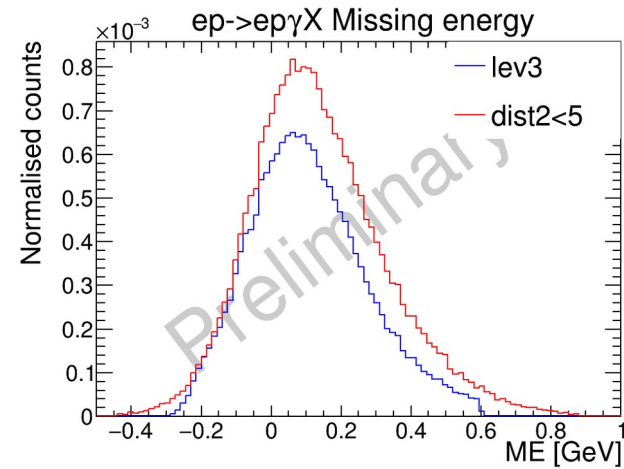
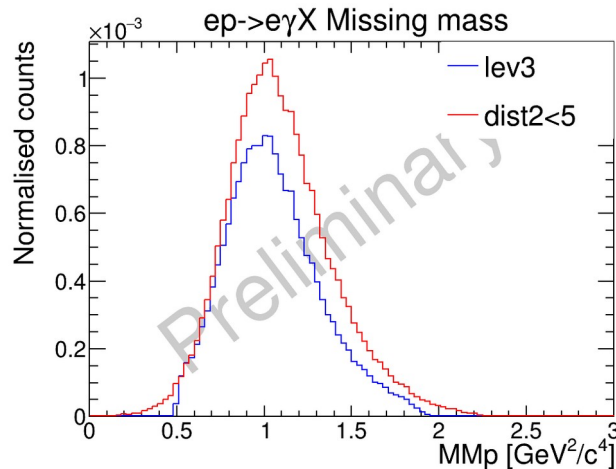


Comparison with models, DSA

- The model points are computed at the average kinematics of each specific phi bin, and the model line is computed at the kinematics of the Q^2 , x_b , t bin

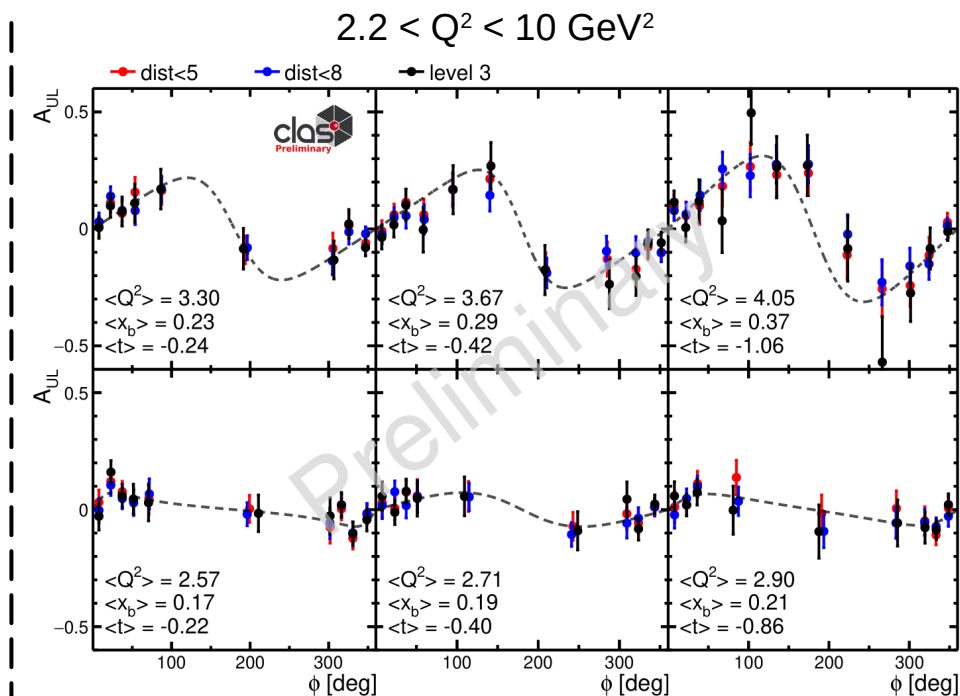
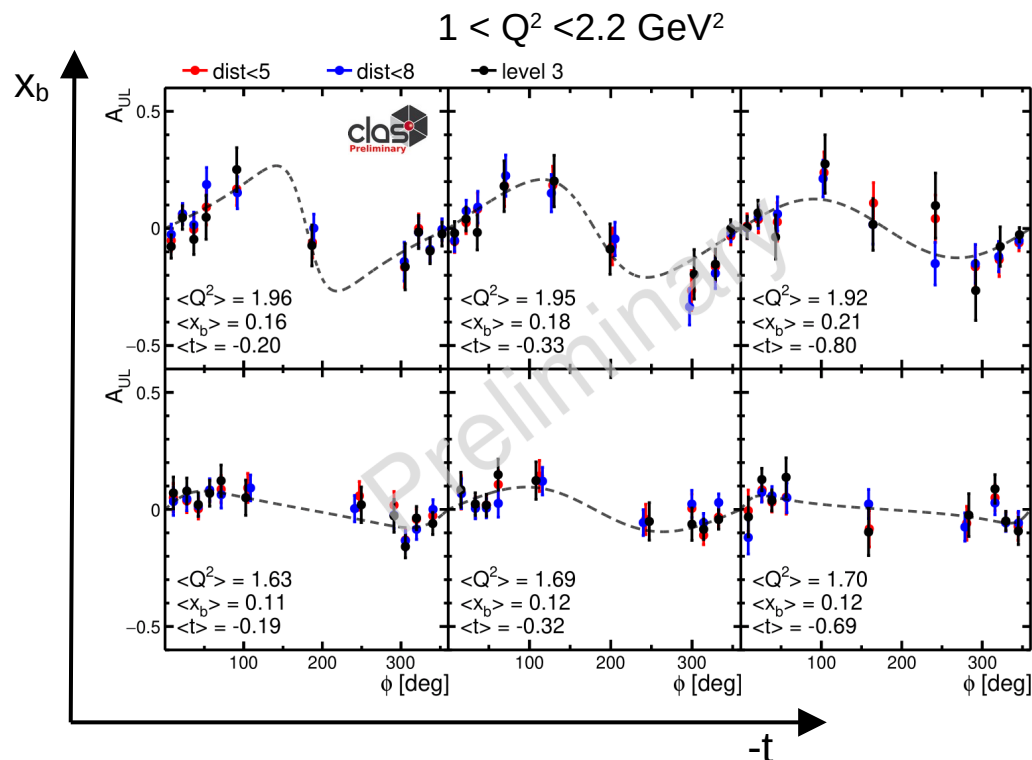


Cut comparison



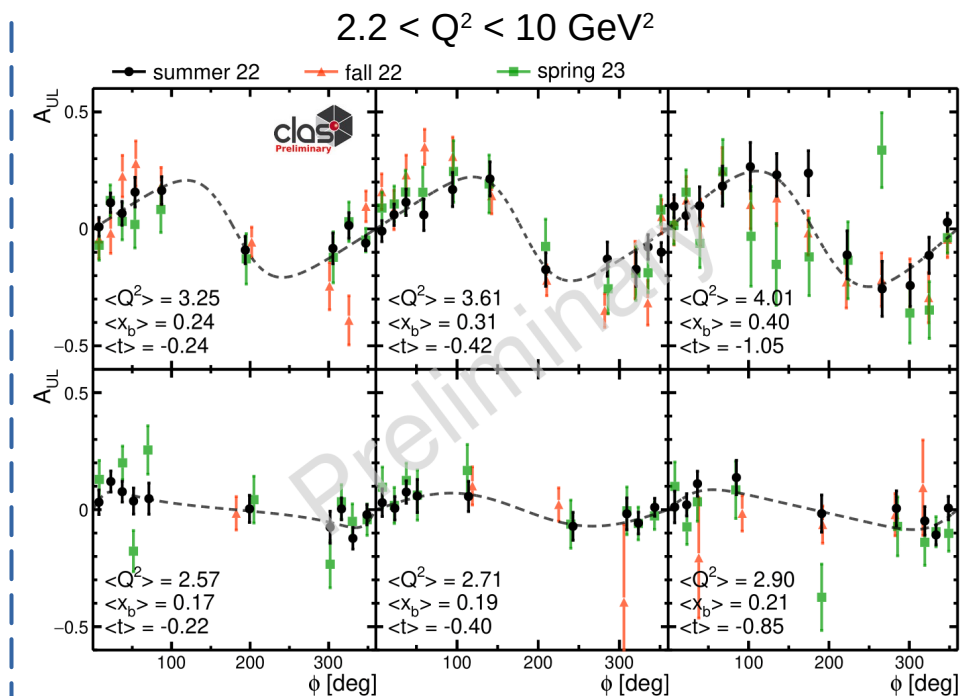
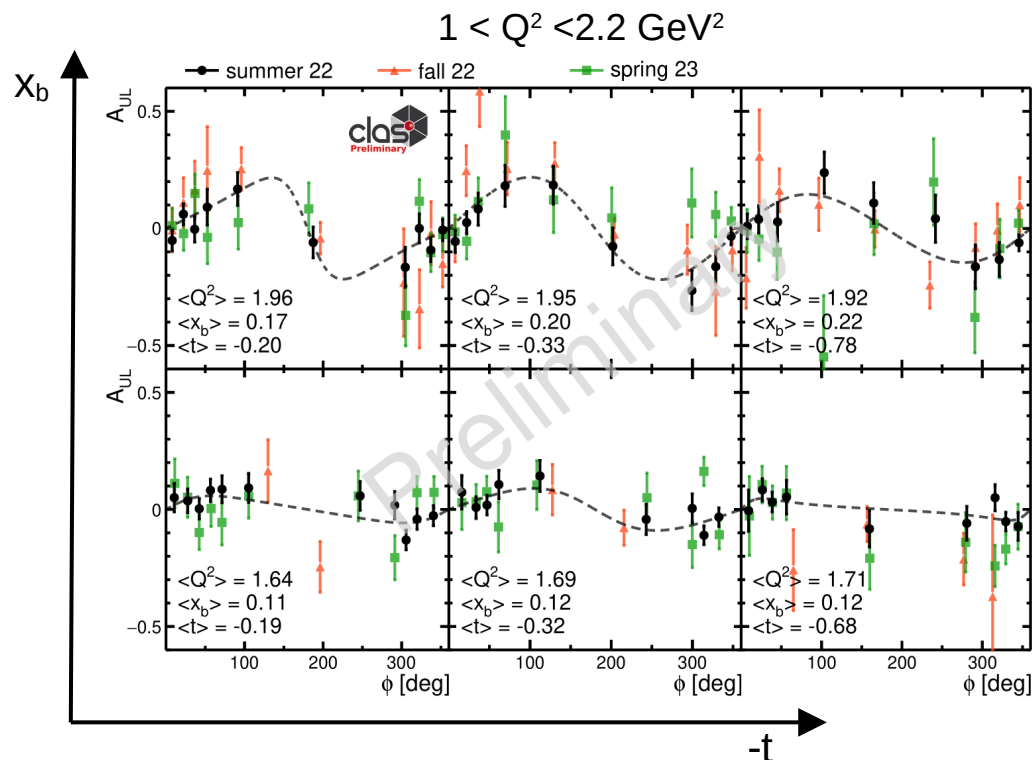
Variable	level 2 cuts	level 3 cuts
MMtot [GeV^2/c^4]	$[-0.1, 0.1]$	$[-0.008, 0.01]$
MMp [GeV^2/c^4]	$[0, 3]$	$[0.5, 2.2]$
ME [GeV]	$[-1.5, 1.5]$	$[-0.5, 0.6]$
Mpt [GeV/c]	$[0, 1]$	$[0, 0.1]$
MMg [GeV^2/c^4]	$[-0.8, 0.8]$	$[-0.25, 0.3]$
g_cone [deg]	$[0, 2]$	$[0, 1]$
$\Delta\phi$ [deg]	$[-3, 3]$	$[-0.5, 0.5]$
Δt [GeV^2/c^4]	$[0, 2]$	$[0, 0.2]$

Cut comparison, TSA



First look at the Fall 22 and Spring 23 TSA

- The QADB timelines are not available for Fall 22 and Spring 23 and no careful run selection has been done yet
- Fall 22 in the FTOFF configuration → different kinematic reach



First look at the Fall 22 and Spring 23 DSA

- The QADB timelines are not available for Fall 22 and Spring 23 and no careful run selection has been done yet
- Fall 22 in the FTOFF configuration → different kinematic reach

