

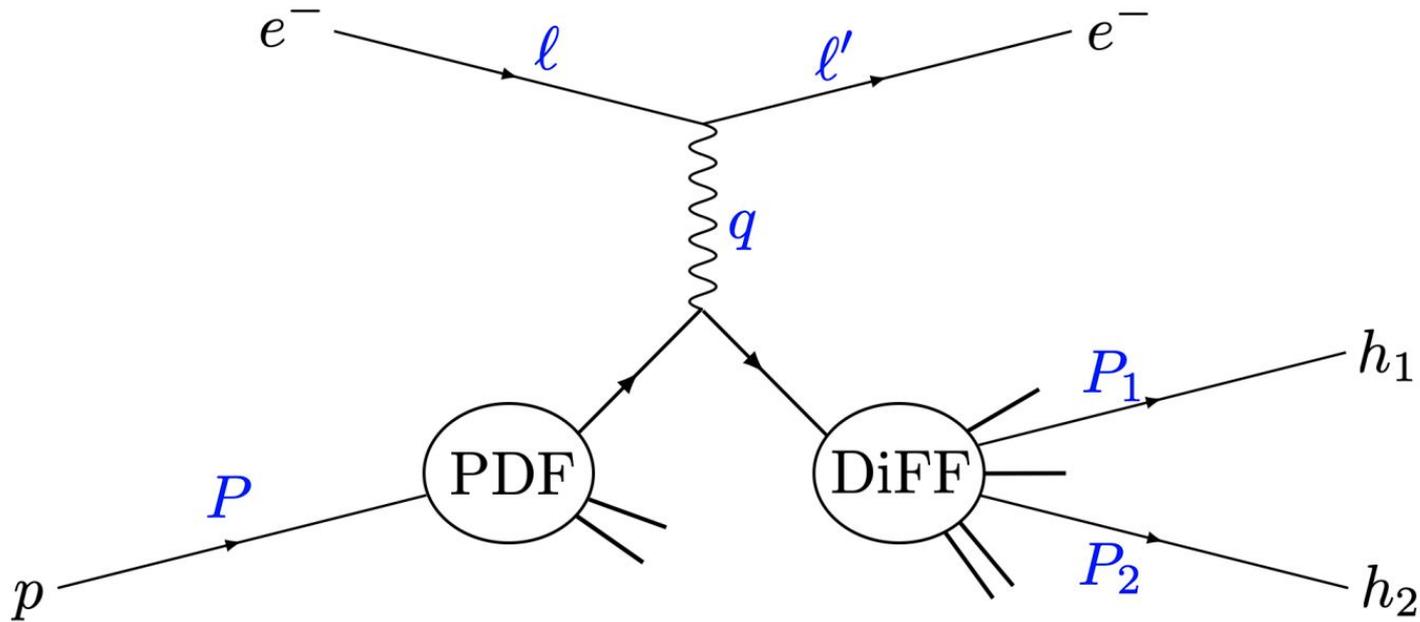
# Partial Wave Decomposition of $\pi\pi$ Beam Spin Asymmetries

Gregory Matousek



# Semi-Inclusive Dihadron Production

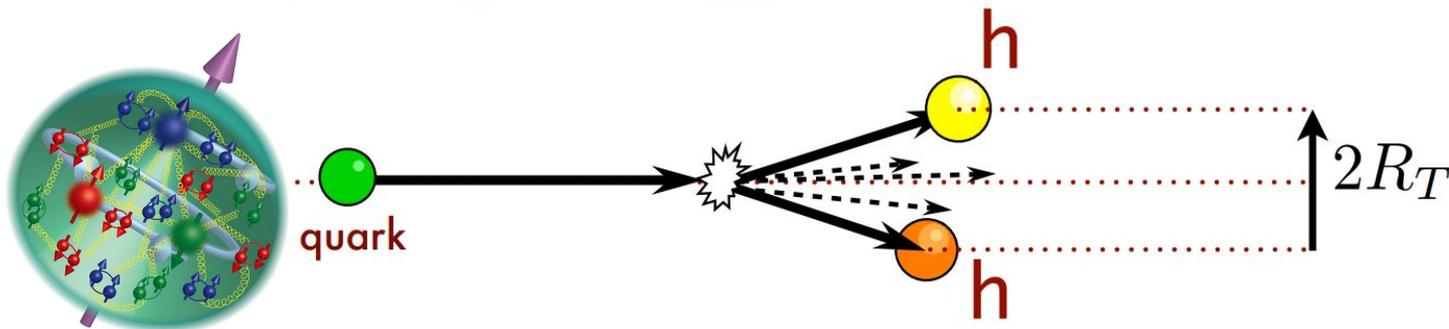
$$e(\ell) + p(P) \rightarrow e(\ell') + h_1(P_1) + h_2(P_2) + X$$



# Semi-Inclusive Dihadron Production

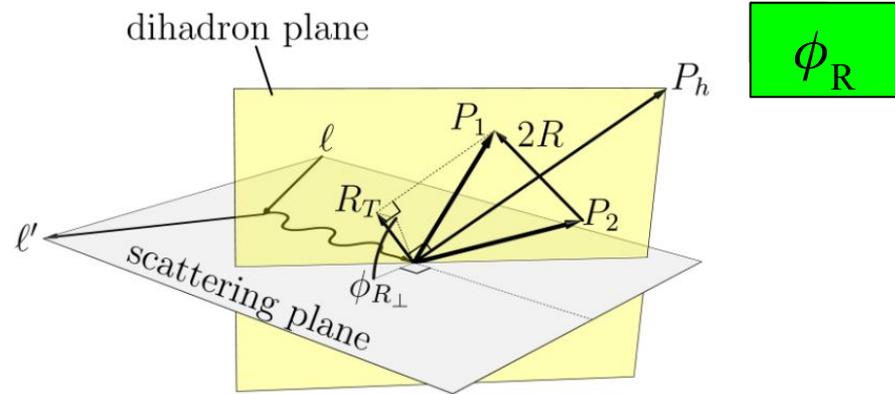
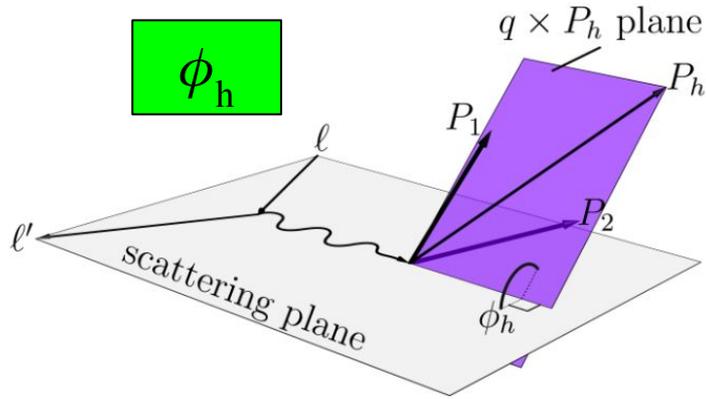
$$e(\ell) + p(P) \rightarrow e(\ell') + h_1(P_1) + h_2(P_2) + X$$

- Dihadron SIDIS production cross section:  $\sigma = \hat{\sigma} \otimes \text{PDF} \otimes \text{DiFF}$
- **Parton Distribution Function:** Probability to find a quark  $q$  with momentum fraction  $x$  in proton
  - **Dihadron Fragmentation Function:** Probability for a quark  $q$  to fragment into hadrons  $h_1$  and  $h_2$  with energy fraction  $z$  and invariant mass  $M_h$

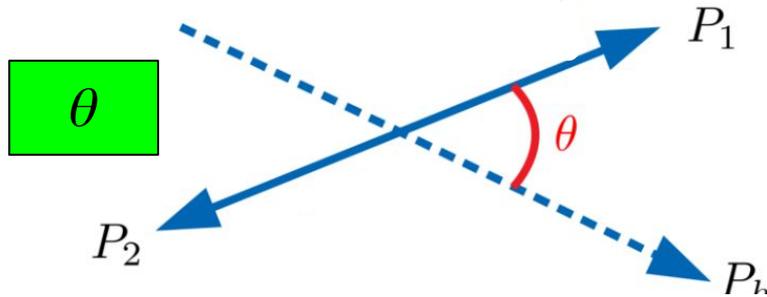


[1]

# Dihadron Angular Variables



## Dihadron CoM frame



★ Our beam spin asymmetries are modulated by these 3 angles ★

# Dihadron Beam Spin Asymmetries

- Beam-spin dependent cross section contains twist-2 and twist-3 terms...

$$d\sigma_{LU} \propto \lambda_e \sum_{\ell=0}^{\ell_{\max}} \left\{ C(\epsilon, y) \sum_{m=1}^{\ell} 2 P_{\ell, m}(\cos \theta) \sin[m(\phi_h - \phi_R)] F_{LU, \text{tw.2}}^{|\ell, m\rangle} \right. \\ \left. + W(\epsilon, y) \sum_{m=-\ell}^{\ell} P_{\ell, m}(\cos \theta) \sin[(1-m)\phi_h + m\phi_R] F_{LU, \text{tw.3}}^{|\ell, m\rangle} \right\}.$$

- Total of 12 independent modulations (7 when  $\theta$ -integrated)

$$F_{LU, \text{tw.2}}^{|\ell, m\rangle} \rightarrow f_1 \otimes G_1^\perp | \ell, m \rangle \quad \left| \quad F_{LU, \text{tw.3}}^{|\ell, m\rangle} \rightarrow e \otimes H_1^\perp | \ell, m \rangle$$

# The Twist-3 PDF $e(x)$

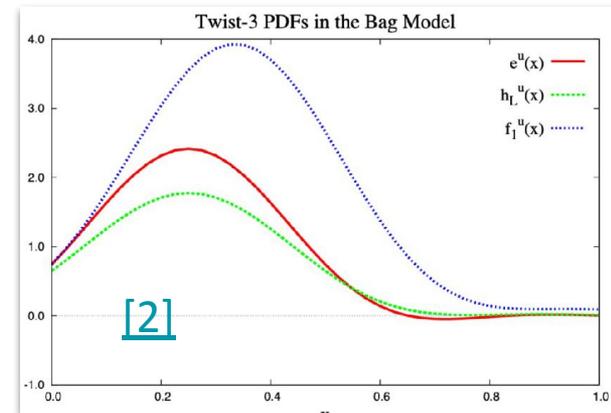
“Quark-gluon interactions help generate fundamental proton properties”

- Twist-2 PDFs  $f_1(x)$ ,  $g_1(x)$ , and  $h_1(x)$  describe **structure** → probabilistic interpretation
- Twist-3 PDFs such as  $e(x)$  capture **dynamics** → no number-density interpretation
  - Involves incoherent scattering off a **quark-gluon** pair in amplitude
  - Suppressed by  $1/Q$  in the cross section → JLab optimal!

“Why do we care about  $e(x)$ ?”

- $x^1$  moment related to pion-nucleon sigma term →  $\chi$ -PT
- $x^2$  moment related to transverse force experienced by T-polarized quark in a T-polarized  $N$  [1]

will not appear in deep-inelastic scattering if quark masses are ignored. In that we know of no practical way to measure  $e(x)$  but we include it here for completeness and because its properties are interesting. Our discussion about



Jaffe, Ji 92

# The Twist-3 PDF $e(x)$

★ **QUESTION:** “Why study 2-hadron SIDIS?” ★

★ **ANSWER:** In the dihadron BSA,  $e(x)$  appears without transverse momentum-dependence! ★

$$d^7\sigma_{LO} = \frac{\alpha^2}{2\pi Q^2 y} \lambda \sum_a e_a^2 W(y) \sin\phi_R \frac{|\vec{R}_T|}{Q} \left[ \frac{M}{M_h} x e(x) H_1^{\not{A}}(z, \zeta, M_h^2) + \frac{1}{z} f_1(x) \tilde{G}^{\not{A}}(z, \zeta, M_h^2) \right]$$

★ This is not the case in 1-hadron SIDIS —  $e(x)$  extraction requires TMD modeling *and* appears in a structure function with 4 PDF⊗FF pairs ★

$$F_{LU}^{\sin\phi_h} = \frac{2M}{Q} \mathcal{C} \left[ -\frac{\hat{\mathbf{h}} \cdot \mathbf{k}_T}{M_h} \left( x e H_1^\perp + \frac{M_h}{M} f_1 \frac{\tilde{G}^\perp}{z} \right) + \frac{\hat{\mathbf{h}} \cdot \mathbf{p}_T}{M} \left( x g^\perp D_1 + \frac{M_h}{M} h_1^\perp \frac{\tilde{E}}{z} \right) \right],$$

# Dihadron Fragmentation Functions (DiFFs)

- The leading twist fragmentation quark-quark correlator (below) contains 4 **DiFFs** [1]

$$D_1 + i H_1^{\triangleleft} \frac{\not{R}_T}{M_h} + i H_1^{\perp} \frac{\not{P}_T}{M_h} + G_1^{\perp} \frac{\epsilon_T^{\mu\nu} R_{T\mu} p_{T\nu}}{M_h^2} \gamma_5$$

Projection along “good” light-cone component

- $G_1^{\perp}$  = measures how much the hadron pair “swirls” clockwise or anti-clockwise about the jet axis (helicity-dependent) – *has no single hadron SIDIS counterpart!*
- $H_1^{\perp}$  and  $H_1^{\triangleleft}$  = measures how the hadron pair’s “R” and “h” planes tilt left-or-right with respect to the quark’s transverse spin — *chiral odd*

★ DiFFs are expanded into a basis of spherical harmonics with respect to  $\theta$ .  
Each partial wave is assigned an  $|l,m\rangle$  angular momentum eigenstate ★

# Event Cuts (RG-A dataset)

## SIDIS Channel

- $Q^2 > 1 \text{ GeV}^2$
- $y < 0.8$
- $W > 2 \text{ GeV}$
- $5 < \theta < 35 \text{ [deg]}$
- Pass QA

## Electron Cuts

- $-8 < v_z < 3 \text{ [cm]}$
- $E_{\text{PCAL}} > 0.07 \text{ [GeV]}$
- $E_{\text{ECIN}} > 0, E_{\text{COUT}} > 0$
- Pcal coords ( $9 < lu, lw < 400$ )
- SF cut
- DC fiducial cut
- Scattered  $e^-$  max energy with  $-3000 < \text{status} < -2000$

## Photon Cuts

- $E > 0.2 \text{ [GeV]}$
- $0.9 < \beta < 1.1$
- $\alpha(e, \gamma) > 8 \text{ [deg]}$
- $E_{\text{PCAL}} > 0$
- Pcal coords ( $14 < lv, lw < 400$ )
- **classifier output  $\rightarrow p > 0.78$  for M.L** 

## Pion Cuts

- $|v_z^e - v_z^\pi| < 20 \text{ [cm]}$
- DC fiducial cut
- $\square_{\text{pid}}^2$  cut
- **$P(\pi^\pm) > 1.25 \text{ [GeV]}$  only for pass-1 analysis**
- Cannot be associated with CD !( $4000 \leq \text{status} < 5000$ )

## Dihadron Cuts

- $M_X > 1.5 \text{ GeV}$
- $x_F(\pi) > 0$
- $z < 0.95$

## Inbending Datasets

$\pi^+\pi^+, \pi^+\pi^-, \pi^+\pi^0$

## Outbending Datasets

$\pi^-\pi^-, \pi^-\pi^+$

# RG-A Datasets

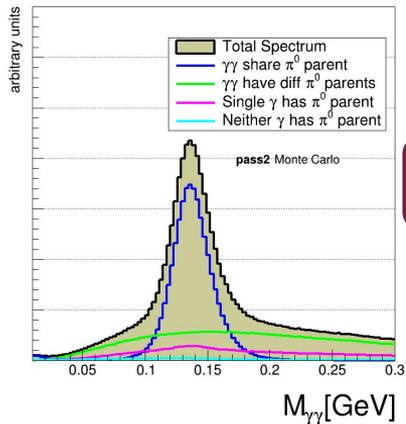
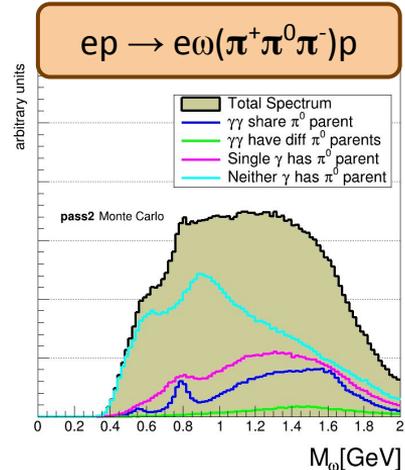
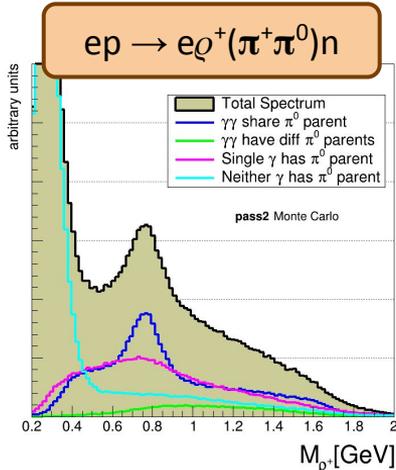
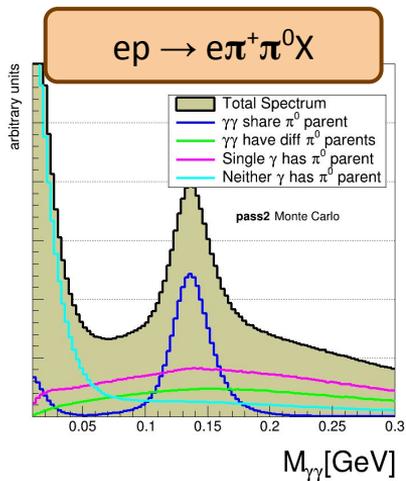
Dataset	Channel	Total	Legacy	ML
MC Inbending	$\pi^+\pi^+$	4.8 M	—	—
	$\pi^+\pi^-$	6.3 M	—	—
	$\pi^-\pi^-$	236 K	—	—
	$\pi^+\pi^0$	8.2 M	340 K	3.0 M
	$\pi^-\pi^0$	2.3 M	87 K	602 K
MC Outbending	$\pi^+\pi^+$	80 K	—	—
	$\pi^+\pi^-$	613 K	—	—
	$\pi^-\pi^-$	158 K	—	—
	$\pi^+\pi^0$	366 K	12 K	108 K
	$\pi^-\pi^0$	497 K	17 K	200 K
RG-A Inbending	$\pi^+\pi^+$	5.6 M	—	—
	$\pi^+\pi^-$	8.0 M	—	—
	$\pi^-\pi^-$	306 K	—	—
	$\pi^+\pi^0$	9.2 M	414 K	4.0 M
	$\pi^-\pi^0$	2.5 M	95 K	783 K
RG-A Outbending	$\pi^+\pi^+$	512 K	—	—
	$\pi^+\pi^-$	4.4 M	—	—
	$\pi^-\pi^-$	1.2 M	—	—
	$\pi^+\pi^0$	2.2 M	83 K	732 K
	$\pi^-\pi^0$	3.0 M	115 K	1.4 M

**pass-1**

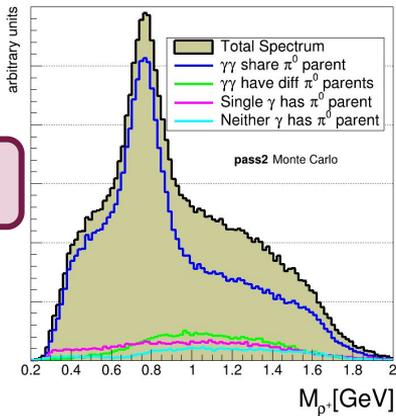
Dataset	Channel	Total	Legacy	ML
MC Inbending	$\pi^+\pi^+$	3.4 M	—	—
	$\pi^+\pi^-$	4.6 M	—	—
	$\pi^-\pi^-$	177 K	—	—
	$\pi^+\pi^0$	5.2 M	248 K	1.1 M
	$\pi^-\pi^0$	1.5 M	63 K	252 K
MC Outbending	$\pi^+\pi^+$	998 K	—	—
	$\pi^+\pi^-$	7.4 M	—	—
	$\pi^-\pi^-$	2.0 M	—	—
	$\pi^+\pi^0$	3.9 M	149 K	801 K
	$\pi^-\pi^0$	5.4 M	218 K	1.1 M
RG-A Inbending	$\pi^+\pi^+$	41.6 M	—	—
	$\pi^+\pi^-$	59.2 M	—	—
	$\pi^-\pi^-$	2.3 M	—	—
	$\pi^+\pi^0$	61.0 M	3.0 M	15.8 M
	$\pi^-\pi^0$	16.3 M	695 K	3.6 M
RG-A Outbending	$\pi^+\pi^+$	6.2 M	—	—
	$\pi^+\pi^-$	52.9 M	—	—
	$\pi^-\pi^-$	14.4 M	—	—
	$\pi^+\pi^0$	24.3 M	1.0 M	5.6 M
	$\pi^-\pi^0$	34.0 M	1.5 M	8.1 M

**pass-2**

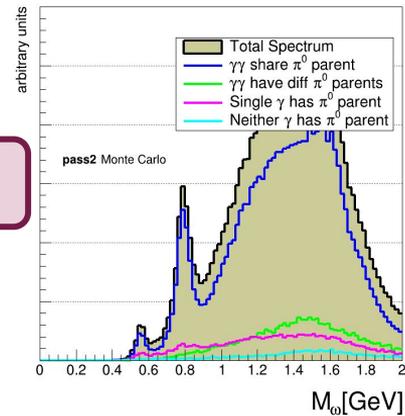
# CLAS12 Photon AI ... SIDIS $\pi^+\pi^0$ ... Exclusive $\rho^+$ , $\omega$



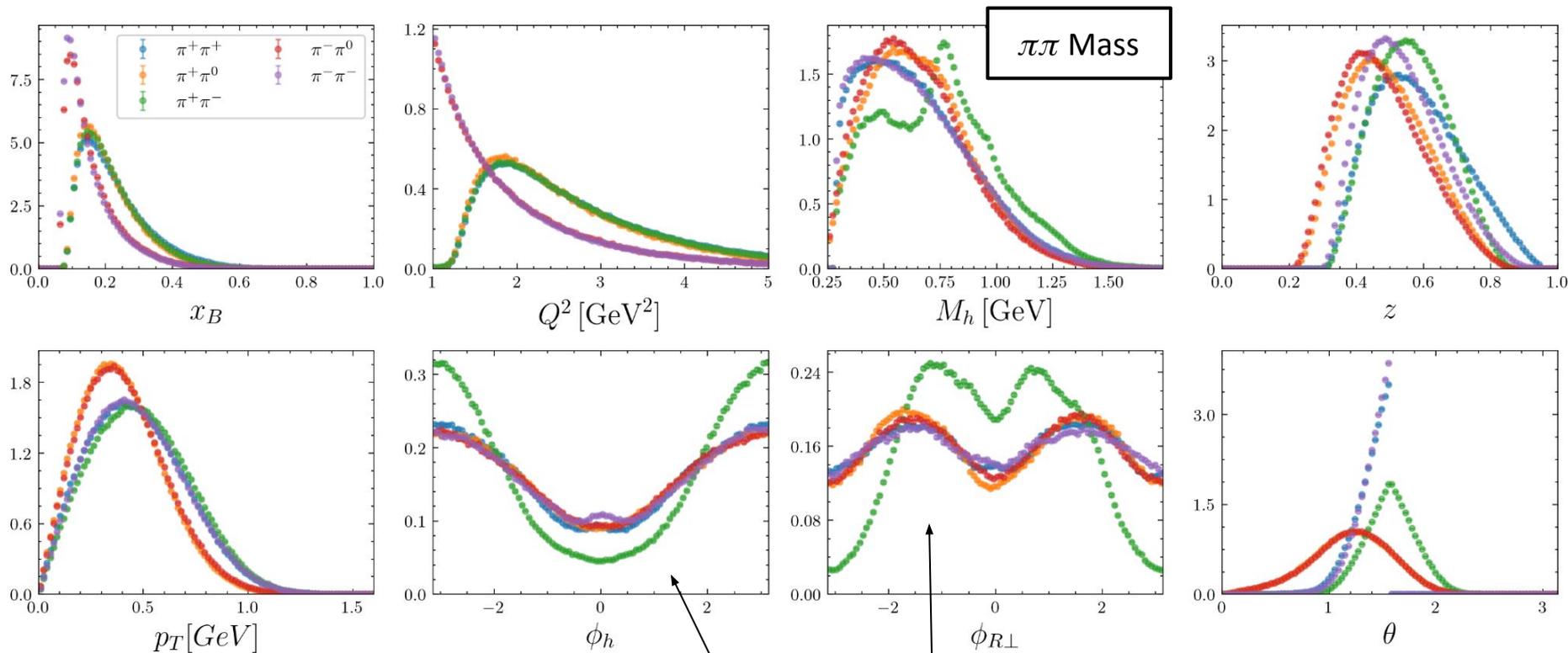
★ ML ★



★ ML ★



# Kinematic Distributions



$\pi^+\pi^-$  differences tied to sharp  $\phi$  peak

# Asymmetry Extraction

- An unbinned maximum log-likelihood fit motivated by the dihadron cross section is used to extract the  $A_{LU}$ 's.

- 12 w/o  $\theta$ -integration
- 7 with  $\theta$ -integration

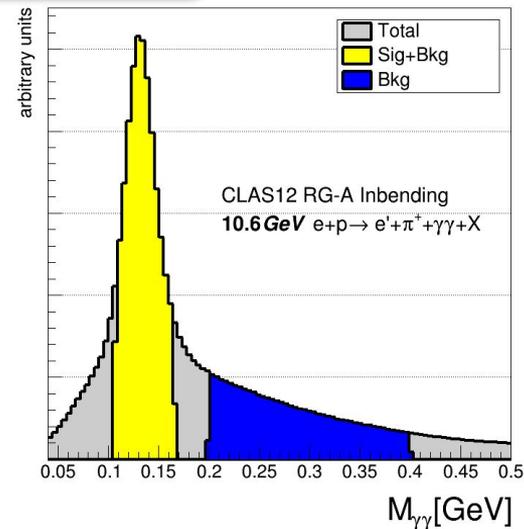
$$P \propto 1 + \lambda_e P_b \left[ A_{LU}^{\Psi_1}(\phi_h, \phi_R, \theta) \Psi_1(\phi_h, \phi_R, \theta) + \dots \right]$$

- For  $\pi^0$ -having dihadrons, **signal+background** & **background** regions are defined where separate asymmetries are extracted.

- ★ Correlations of the  $\pi^0$  purity “ $u$ ” with  $\phi_h, \phi_R$  require us to bin the purity...

- Without this, asymmetries from one modulation can “pull at” the asymmetry in other modulations
- Verified with MC injection studies (**extra slides**)

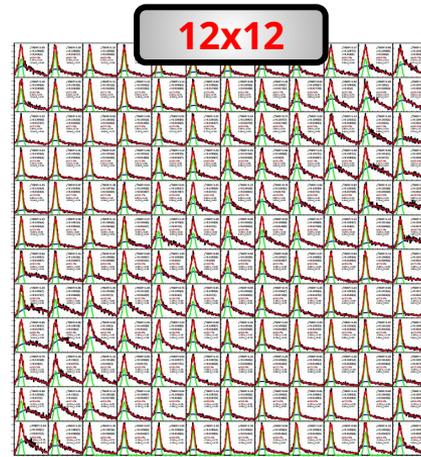
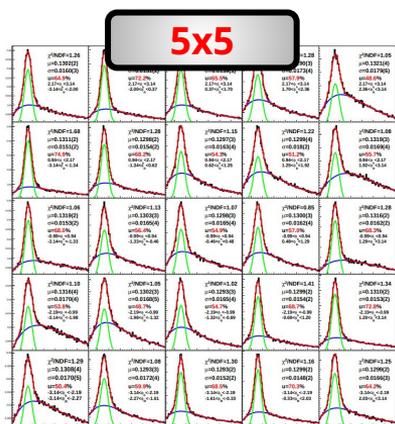
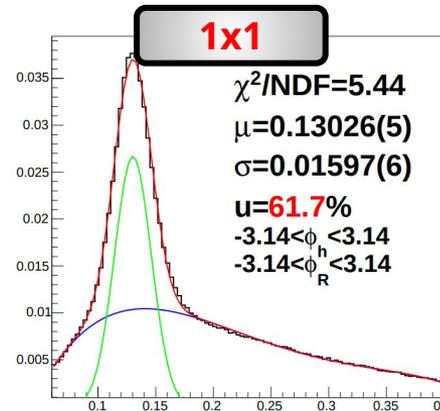
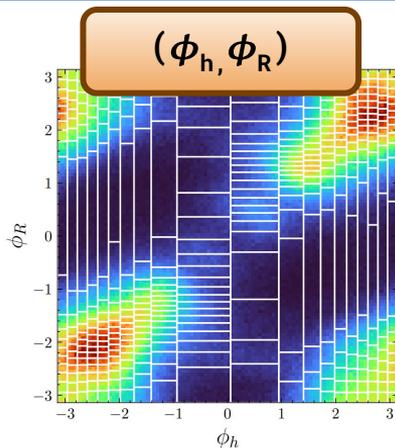
$$\mathcal{P}_{\pm}(\phi_h, \phi_R; A_{\ell}) = \frac{L_{\pm}}{L} \left( 1 \pm P_b \sum_{\ell} \psi_{\ell}(\phi_h, \phi_R) \left[ u(\phi_h, \phi_R) A_{LU,\ell}^{\text{sig}} + (1 - u(\phi_h, \phi_R)) A_{LU,\ell}^{\text{bkg}} \right] \right).$$



# Calculating Event-By-Event Purity

$u(\phi_h, \phi_R)$  Procedure:

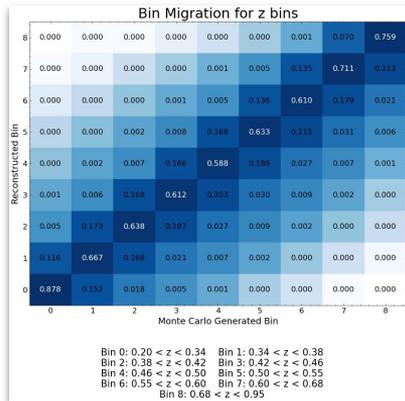
1. For each kinematic bin ( $x, z$ , etc.) subdivide the  $(\phi_h, \phi_R)$  space into  $N \times M$  asymmetric bins such that each bin contains *roughly* the same # events
2.  $f(M_{\gamma\gamma}) = \text{gaus} + \text{pol4}$  in each sub-bin to calculate its purity (signal integral between  $0.106 < M_{\gamma\gamma} < 0.166$ )
3. For the **MLM** in the sig+bkg region, estimate the event-by-event purity as the purity bin for which the event's  $(\phi_h, \phi_R)$  is inside



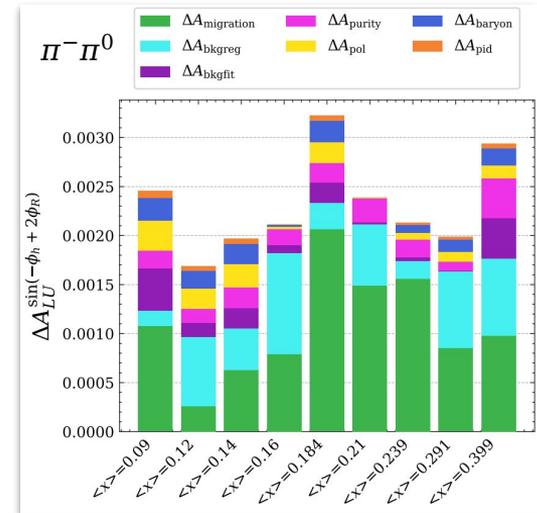
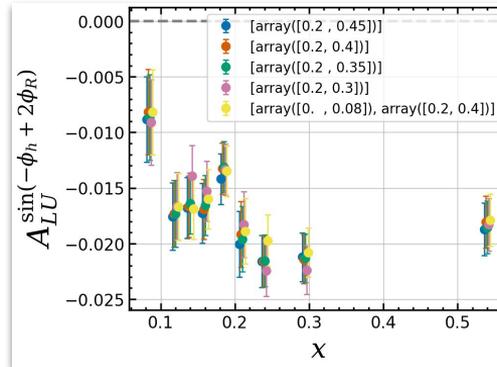
# Systematic Uncertainties

- One source of normalization error (beam polarization)
- Six sources of point-to-point errors explored
  - Bin Migration
  - Density of the purity-binning scheme
  - Baryonic decay contamination (ex:  $\Lambda \rightarrow p\pi^-$ )
  - Particle misidentification
  - Background polynomial fit degree ( $\pi^0$  purity)
  - Sideband region definition ( $\pi^0$  background asymmetry)

★ Ex:  $\pi^- \pi^0$  systematic error breakdown for  $A_{LU}$  twist-3  $|2,2\rangle$ . ★



← Bin migration systematic



←  $\pi^0$  sideband-region systematic

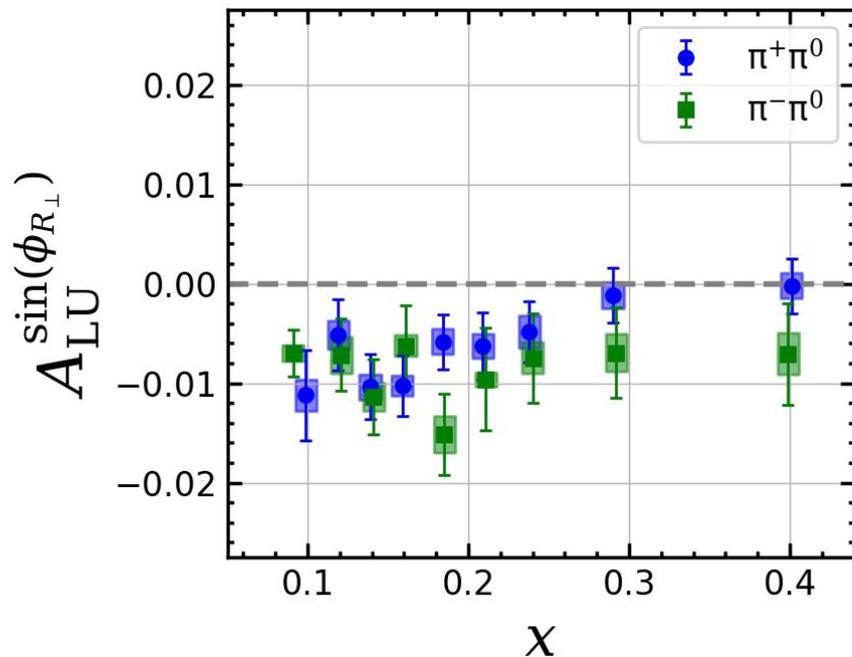
# PRL Plots: $\pi^\pm\pi^0$ Twist-3

Non-zero result measured for both channels

Second, independent avenue for extracting  $e(x)$  !

$A_{LU} < 0$  as opposed to large positive signal found for the  $\pi^+\pi^-$  at CLAS12 [1] (due to sign of  $H_1^\zeta$ ?)

Twist-3 DiFF  $\tilde{G}^\zeta$  believed to be negligible due to Wandzura Wilcek approximation (to be measured with complementary RG-C dihadron studies)



$$d^7\sigma_{LO} = \frac{\alpha^2}{2\pi Q^2 y} \lambda \sum_a e_a^2 W(y) \sin\phi_R \frac{|\vec{R}_T|}{Q} \left[ \frac{M}{M_h} x e(x) H_1^\zeta(z, \zeta, M_h^2) + \frac{1}{z} f_1(x) \tilde{G}^\zeta(z, \zeta, M_h^2) \right]$$

$$\zeta \equiv \zeta(\theta)$$

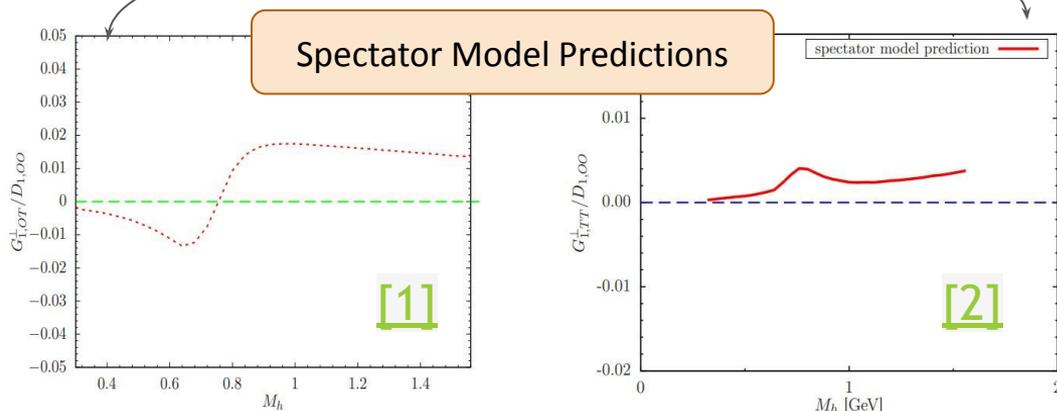
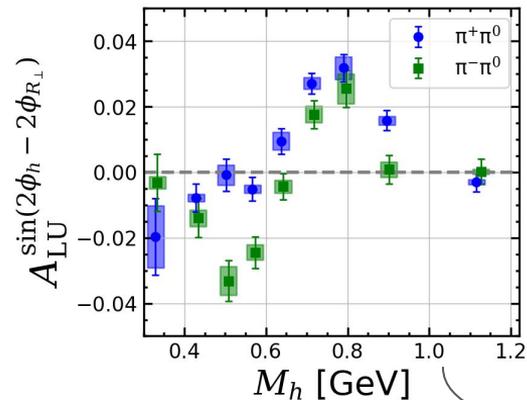
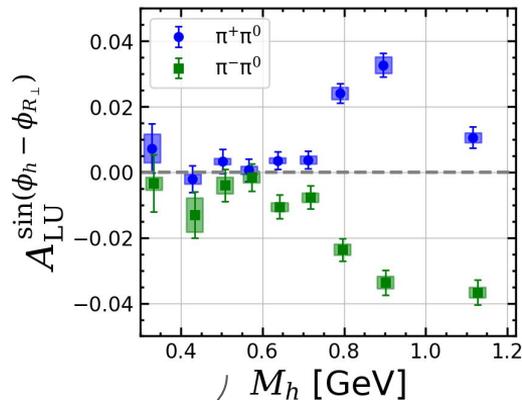
# PRL Plots: $\pi^\pm\pi^0$ Twist-2

Clear resonant behavior near  $\rho$ -mass, indicative of s- and p- wave dihadron interference

First evidence of isospin-dependence in  $G_1^\perp$ , depends on the partial wave!

Very limited model predictions for non- $\pi^+\pi^-$ , future phenomenology?

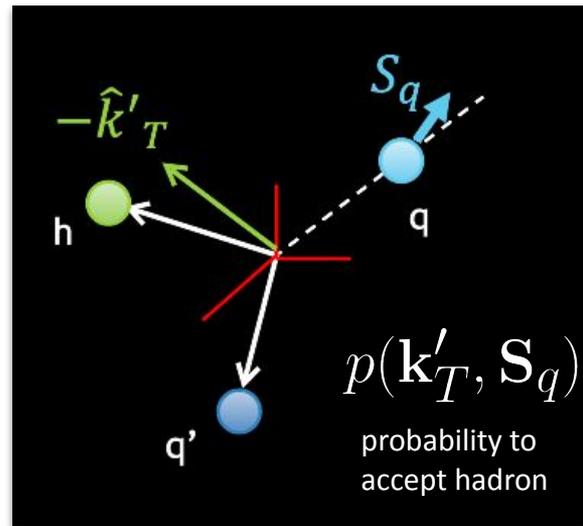
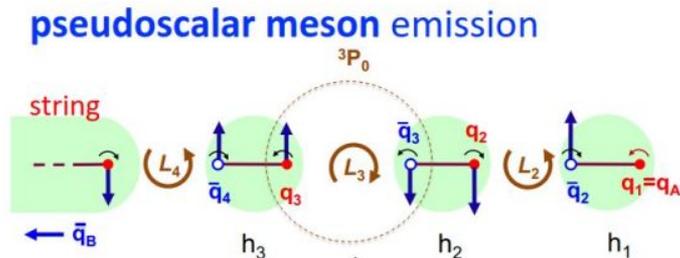
$$F_{LU,tw.2}^{|\ell,m\rangle} \rightarrow f_1 \otimes G_1^\perp^{|\ell,m\rangle}$$



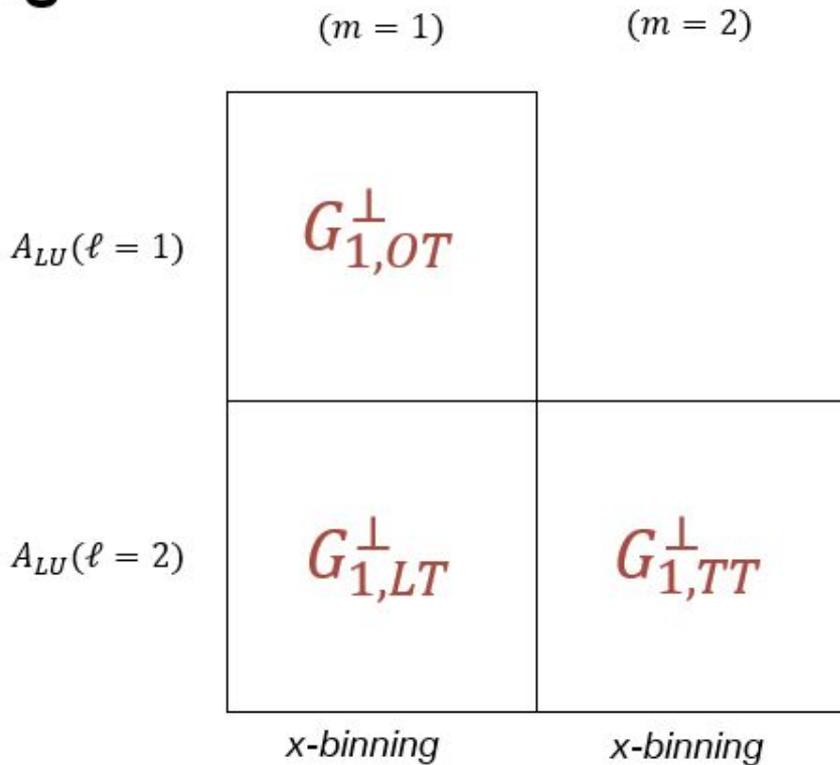
# Spin-dependent fragmentation in PYTHIA8

- **StringSpinner** is a plugin for the introduction of the spin effects in the hadronization part of PYTHIA8 [1], [2]
  - Replicates COMPASS  $\pi$ , K Collins asymmetries – Replicates BELLE b2b hadron asymmetries [3]
- Quark-spin propagated in hadronization using the String+3P0 fragmentation model [4]
  - Initialize struck quark with longitudinal polarization
  - Generate q-qbar with  $L=S=1, J=0$  (vacuum quantum #'s)
  - Bias hadrons with  $\mathbf{k}_T$  relative to fragmenting quark  $\mathbf{S}_q$
  - Propagate spin to next quark  $\mathbf{q}'$  and (if produced) vector mesons

★ Can **StringSpinner** MC reproduce **CLAS12** partial waves? ★

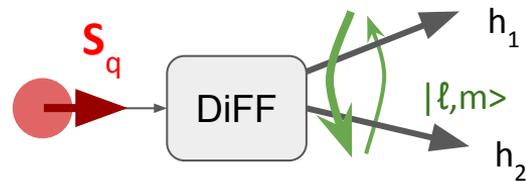


# Plot Formatting (Twist-2)

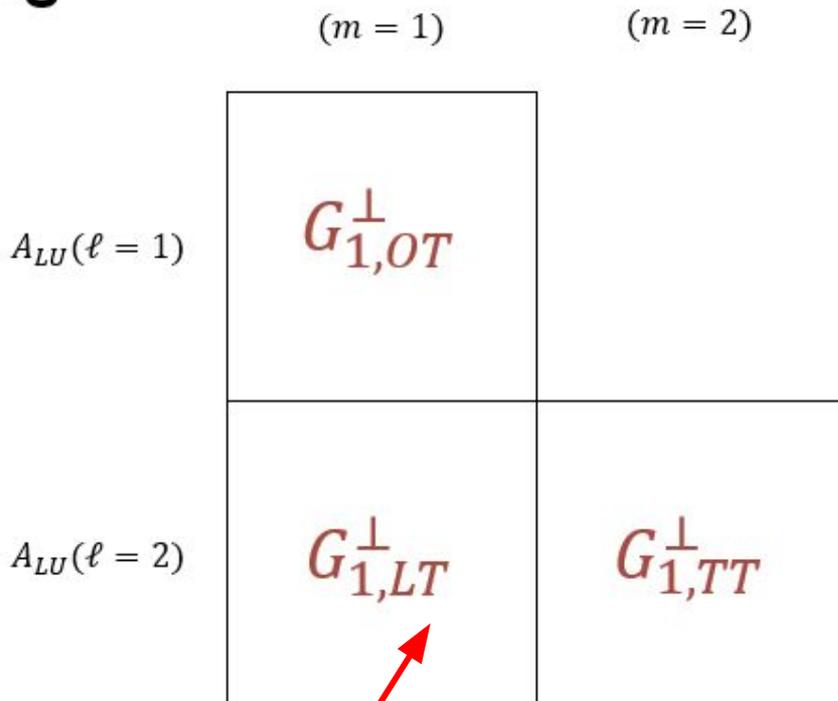


$ \ell, m\rangle$	Twist-2 Amplitude	Twist-3 Amplitude
$ 0, 0\rangle$	—	$P_{0,0} \sin(\phi_h)$
$ 1, 1\rangle$	$P_{1,1} \sin(\phi_h - \phi_R)$	$P_{1,1} \sin(\phi_R)$
$ 1, 0\rangle$	—	$P_{1,0} \sin(\phi_h)$
$ 1, -1\rangle$	—	$P_{1,-1} \sin(2\phi_h - \phi_R)$
$ 2, 2\rangle$	$P_{2,2} \sin(2(\phi_h - \phi_R))$	$P_{2,2} \sin(-\phi_h + 2\phi_R)$
$ 2, 1\rangle$	$P_{2,1} \sin(\phi_h - \phi_R)$	$P_{2,1} \sin(\phi_R)$
$ 2, 0\rangle$	—	$P_{2,0} \sin(\phi_h)$
$ 2, -1\rangle$	—	$P_{2,-1} \sin(2\phi_h - \phi_R)$
$ 2, -2\rangle$	—	$P_{2,-2} \sin(3\phi_h - 2\phi_R)$

$$A_{LU}^{tw2} \propto f_1 \otimes G_1^\perp$$

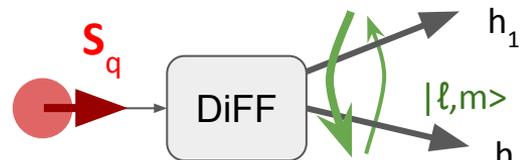


# Plot Formatting (Twist-2)



$ \ell, m\rangle$	Twist-2 Amplitude	Twist-3 Amplitude
$ 0, 0\rangle$	—	$P_{0,0} \sin(\phi_h)$
$ 1, 1\rangle$	$P_{1,1} \sin(\phi_h - \phi_R)$	$P_{1,1} \sin(\phi_R)$
$ 1, 0\rangle$	—	$P_{1,0} \sin(\phi_h)$
$ 1, -1\rangle$	—	$P_{1,-1} \sin(2\phi_h - \phi_R)$
$ 2, 2\rangle$	$P_{2,2} \sin(2(\phi_h - \phi_R))$	$P_{2,2} \sin(-\phi_h + 2\phi_R)$
$ 2, 1\rangle$	$P_{2,1} \sin(\phi_h - \phi_R)$	$P_{2,1} \sin(\phi_R)$
$ 2, 0\rangle$	—	$P_{2,0} \sin(\phi_h)$
$ 2, -1\rangle$	—	$P_{2,-1} \sin(2\phi_h - \phi_R)$
$ 2, -2\rangle$	—	$P_{2,-2} \sin(3\phi_h - 2\phi_R)$

$$A_{LU}^{tw2} \propto f_1 \otimes G_1^\perp$$

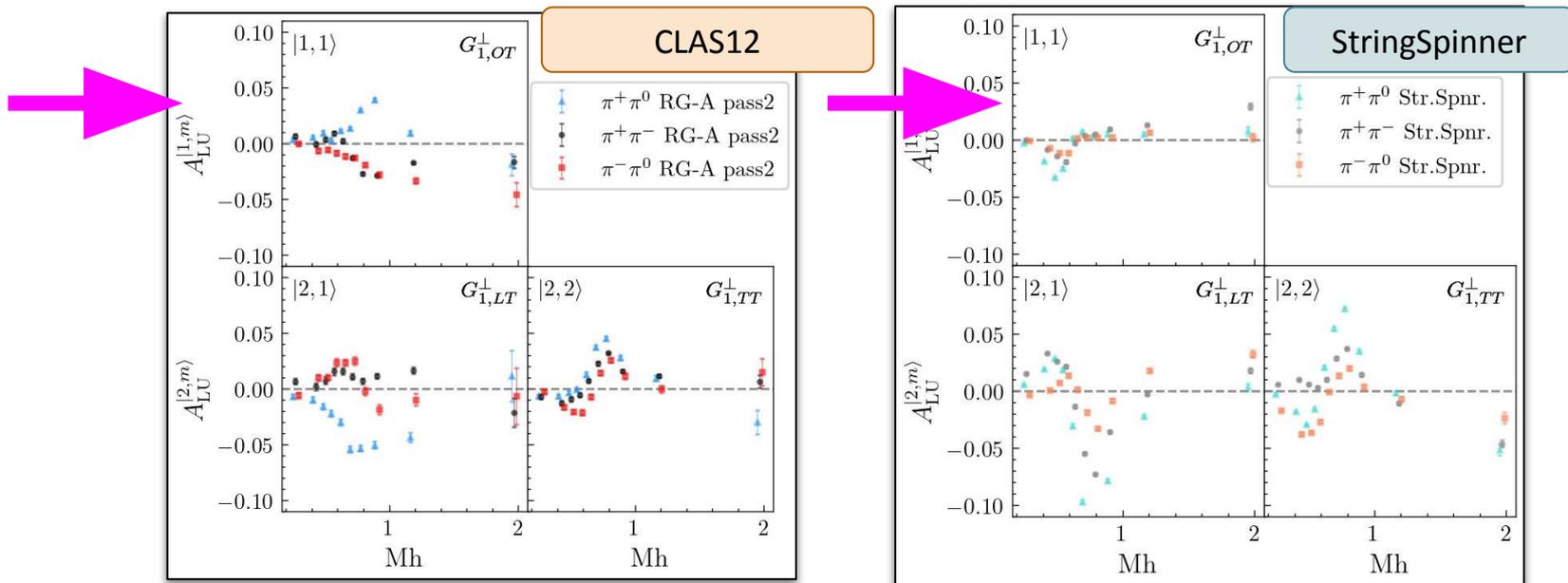


$$d\sigma_{LU} = \lambda_e(\dots + A_{LU,tw.2}^{[2,1]} \sin 2\theta \sin[\phi_h - \phi_R] + \dots)$$

# Twist-2 Partial Waves ( $\pi^+\pi^-$ , $\pi^+\pi^0$ , $\pi^-\pi^0$ )

 $G_{1,OT}^\perp$ 

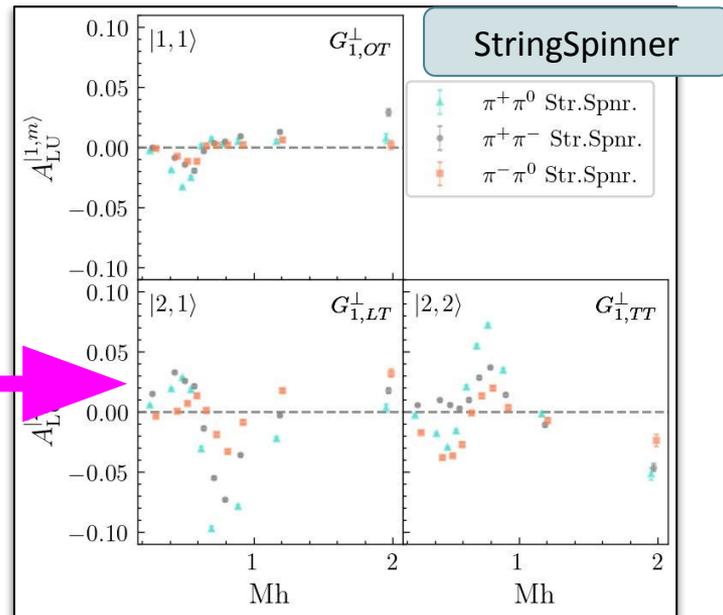
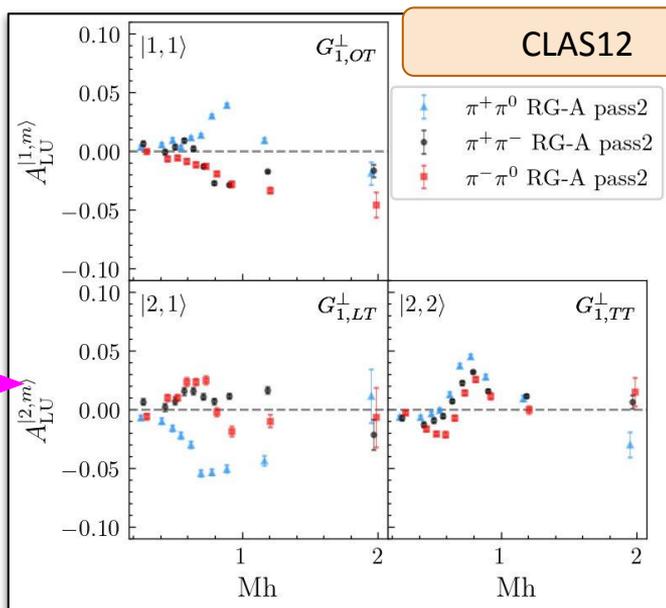
- Sign change only in **CLAS12**  $\pi^+\pi^-$  (predicted in Spectator Model [\[1\]](#))
- $\omega \rightarrow \pi\pi X$  decay creates negative **StringSpinner** signal. Why CLAS12  $\pi^+\pi^0$  so positive?



# Twist-2 Partial Waves ( $\pi^+\pi^-$ , $\pi^+\pi^0$ , $\pi^-\pi^0$ )

 $G_{1,LT}^\perp$ 

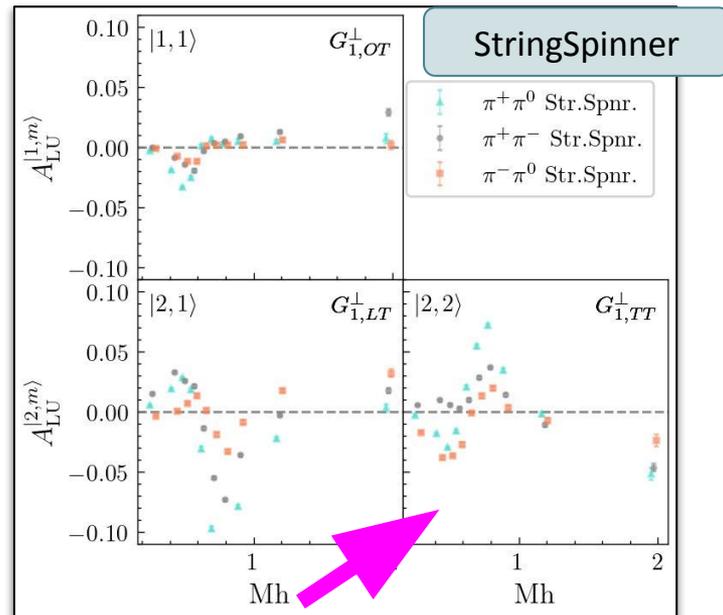
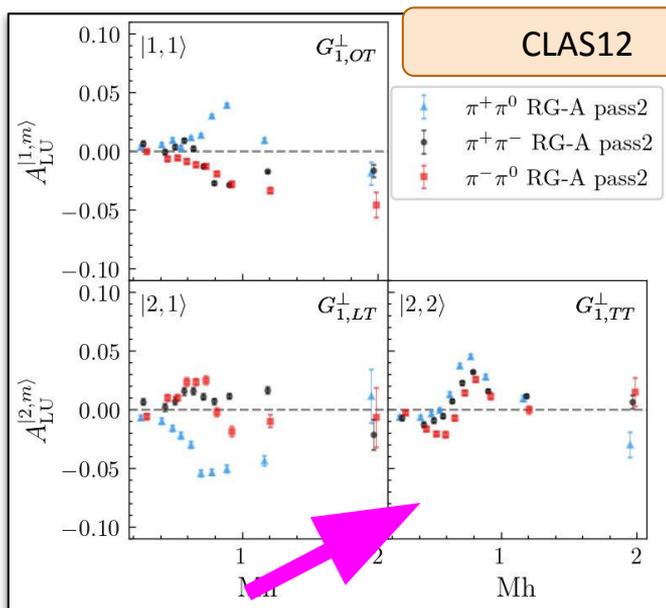
- Unexplored in literature → large **CLAS12** asymmetries observed!
- Interference of **L** and **T** polarized VM's create effect in **StringSpinner**
  - Mechanism for  $\pi^+\pi^-$  staying positive still unknown...



# Twist-2 Partial Waves ( $\pi^+\pi^-$ , $\pi^+\pi^0$ , $\pi^-\pi^0$ )

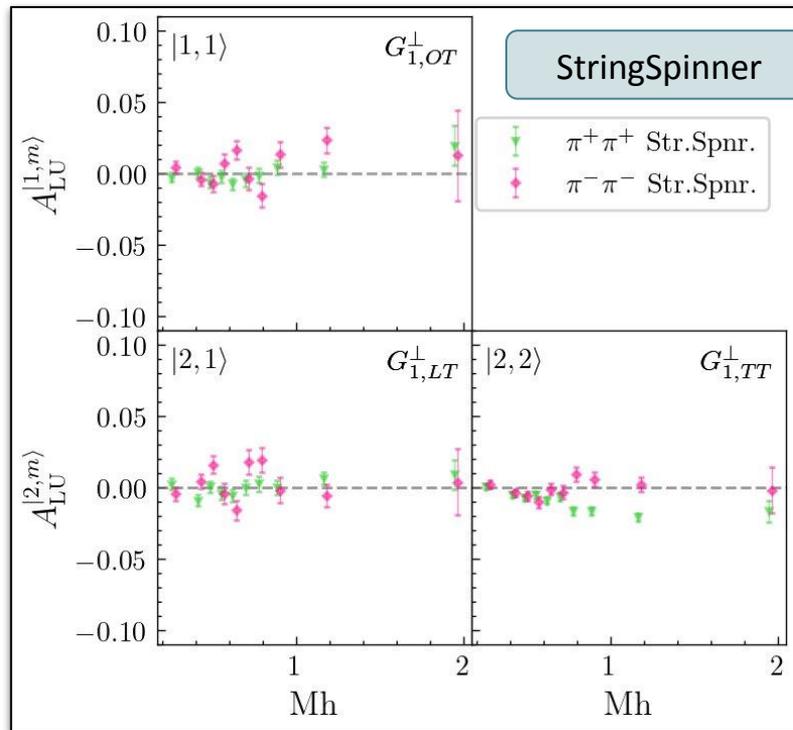
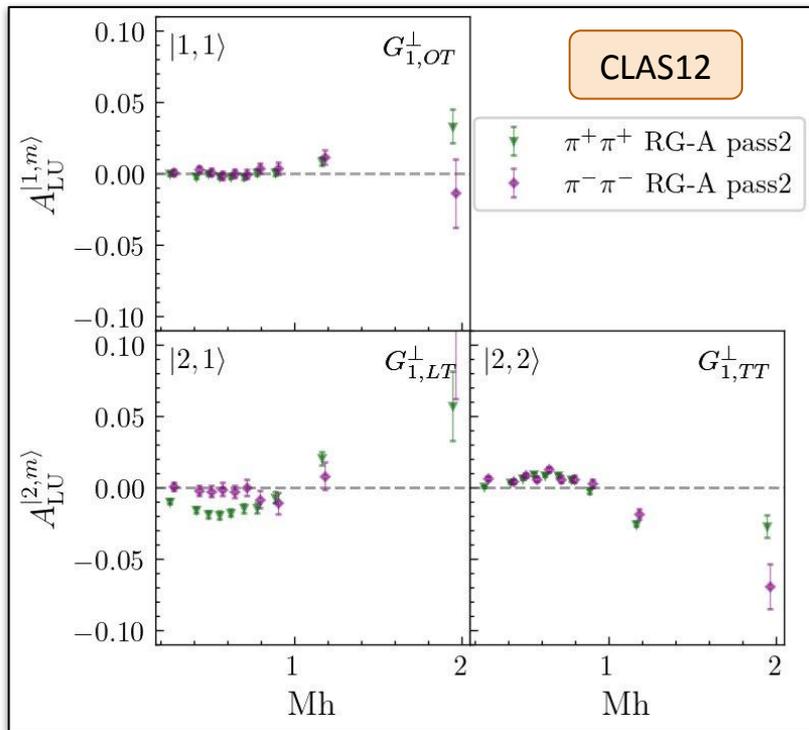
 $G_{1,TT}^\perp$ 

- Excellent agreement between **CLAS12** and **StringSpinner**
- Asymmetry generated by interference of polarized  $\rho$ 's  $\rightarrow$  captures relative size!



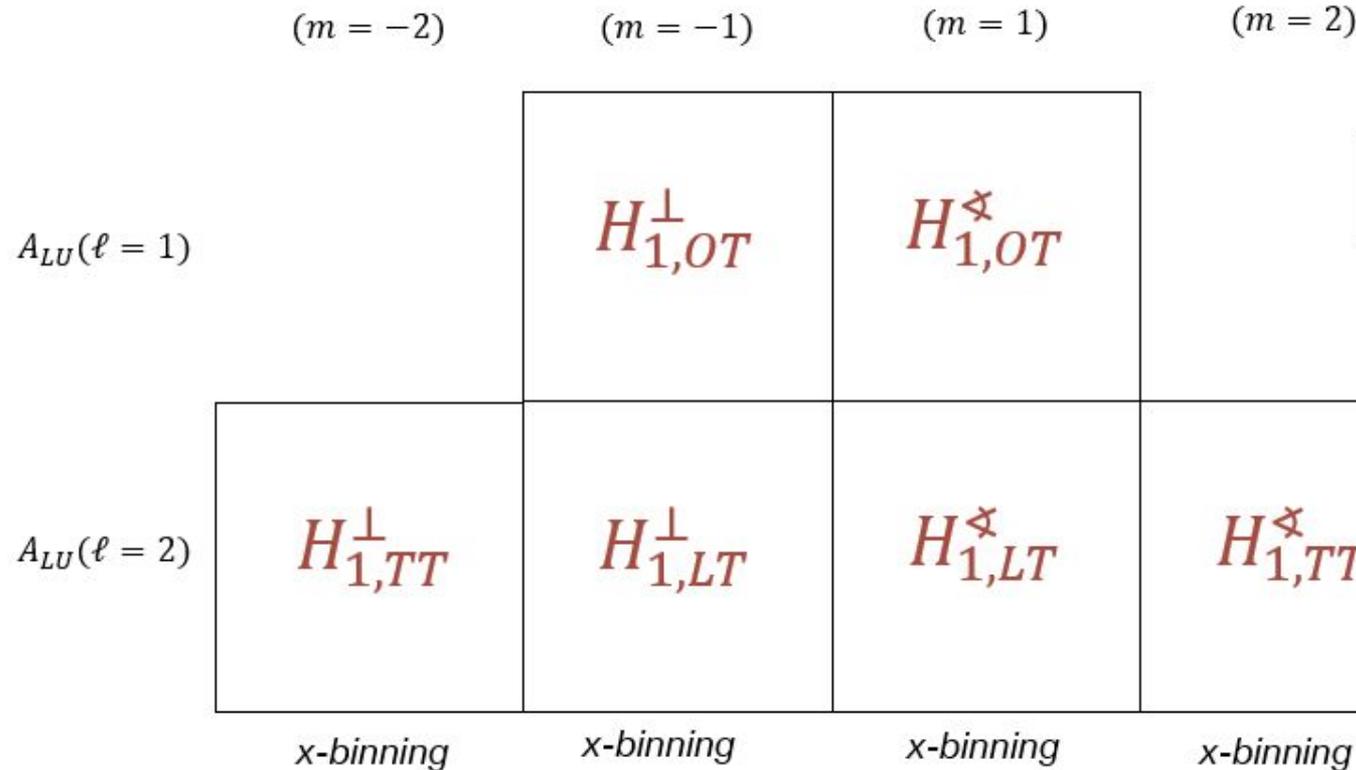
# Twist-2 Partial Waves ( $\pi^+\pi^+$ , $\pi^-\pi^-$ )

- High  $M_h$  asymmetry in like-pion pairs *mysterious* - high angular momentum states?

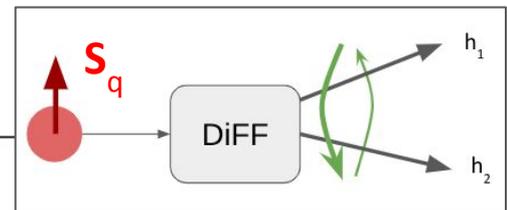


# Plot Formatting (Twist-3)

$ \ell, m\rangle$	Twist-2 Amplitude	Twist-3 Amplitude
$ 0, 0\rangle$	—	$P_{0,0} \sin(\phi_h)$
$ 1, 1\rangle$	$P_{1,1} \sin(\phi_h - \phi_R)$	$P_{1,1} \sin(\phi_R)$
$ 1, 0\rangle$	—	$P_{1,0} \sin(\phi_h)$
$ 1, -1\rangle$	—	$P_{1,-1} \sin(2\phi_h - \phi_R)$
$ 2, 2\rangle$	$P_{2,2} \sin(2(\phi_h - \phi_R))$	$P_{2,2} \sin(-\phi_h + 2\phi_R)$
$ 2, 1\rangle$	$P_{2,1} \sin(\phi_h - \phi_R)$	$P_{2,1} \sin(\phi_R)$
$ 2, 0\rangle$	—	$P_{2,0} \sin(\phi_h)$
$ 2, -1\rangle$	—	$P_{2,-1} \sin(2\phi_h - \phi_R)$
$ 2, -2\rangle$	—	$P_{2,-2} \sin(3\phi_h - 2\phi_R)$



$$A_{LU}^{tw3} \propto e \otimes H_1^{\perp, \chi}$$



Omitting  $m=0$  to save some space. Also, twist-3  $|2,0\rangle$  and  $|0,0\rangle$  correlate with one another

$$P_{0,0} = 1 \leftrightarrow P_{2,0} = \frac{1}{2} (3 \cos^2 \theta - 1)$$

# Plot Formatting (Twist-3)

$ \ell, m\rangle$	Twist-2 Amplitude	Twist-3 Amplitude
$ 0, 0\rangle$	—	$P_{0,0} \sin(\phi_h)$
$ 1, 1\rangle$	$P_{1,1} \sin(\phi_h - \phi_R)$	$P_{1,1} \sin(\phi_R)$
$ 1, 0\rangle$	—	$P_{1,0} \sin(\phi_h)$
$ 1, -1\rangle$	—	$P_{1,-1} \sin(2\phi_h - \phi_R)$
$ 2, 2\rangle$	$P_{2,2} \sin(2(\phi_h - \phi_R))$	$P_{2,2} \sin(-\phi_h + 2\phi_R)$
$ 2, 1\rangle$	$P_{2,1} \sin(\phi_h - \phi_R)$	$P_{2,1} \sin(\phi_R)$
$ 2, 0\rangle$	—	$P_{2,0} \sin(\phi_h)$
$ 2, -1\rangle$	—	$P_{2,-1} \sin(2\phi_h - \phi_R)$
$ 2, -2\rangle$	—	$P_{2,-2} \sin(3\phi_h - 2\phi_R)$

$(m = -2)$

$(m = -1)$

$(m = 0)$  \$money plot\$  $(m = 2)$

$A_{LU}(\ell = 1)$

$H_{1,OT}^\perp$

\$\$\$\$\$\$\$\$

$H_{1,OT}^\chi$

\$\$\$\$\$\$\$\$

$$A_{LU}^{tw3} \propto e \otimes H_1^{\perp, \chi}$$

Direct access to collinear  $e(x)$  without TMD-modeling

$A_{LU}(\ell = 2)$

$H_{1,TT}^\perp$

$H_{1,LT}^\perp$

$H_{1,LT}^\chi$

$H_{1,TT}^\chi$

$x$ -binning

$x$ -binning

$x$ -binning

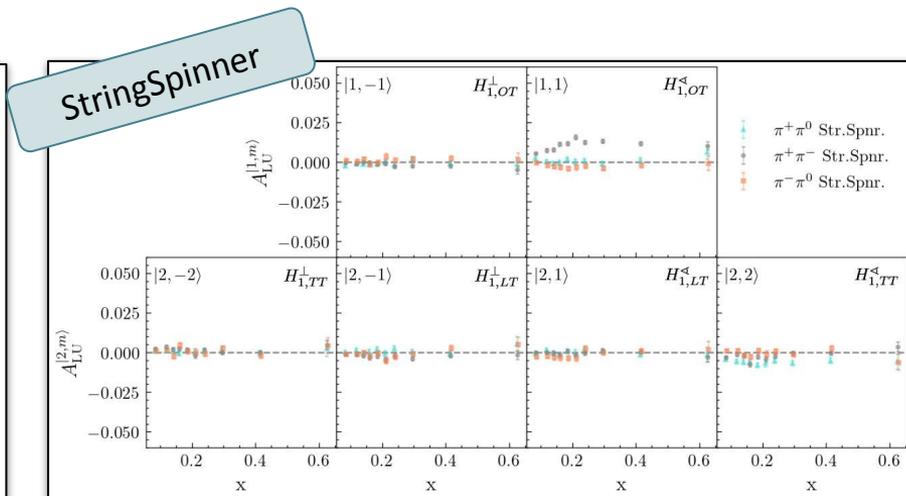
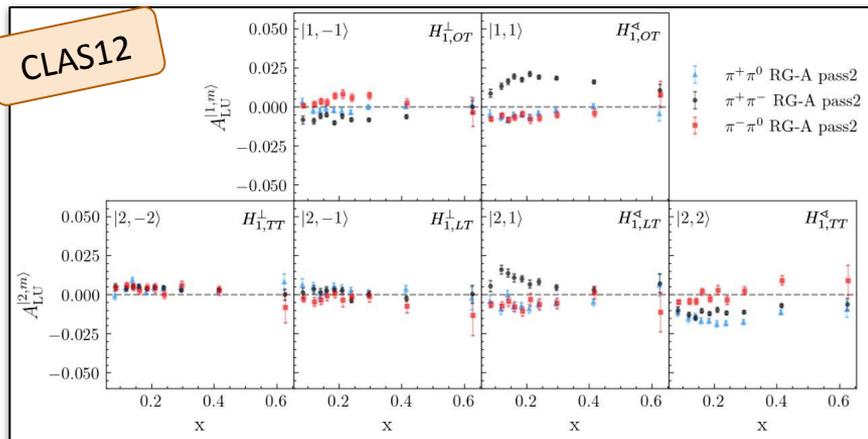
$x$ -binning

Omitting  $m=0$  to save some space. Also, twist-3  $|2,0\rangle$  and  $|0,0\rangle$  correlate with one another

$$P_{0,0} = 1 \leftrightarrow P_{2,0} = \frac{1}{2}(3 \cos^2 \theta - 1)$$

# Twist-3 Partial Waves ( $\pi^+\pi^-$ , $\pi^+\pi^0$ , $\pi^-\pi^0$ )

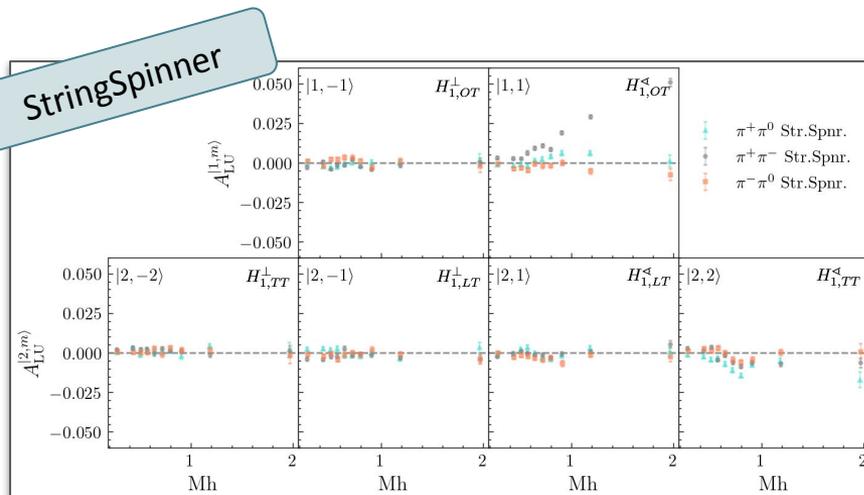
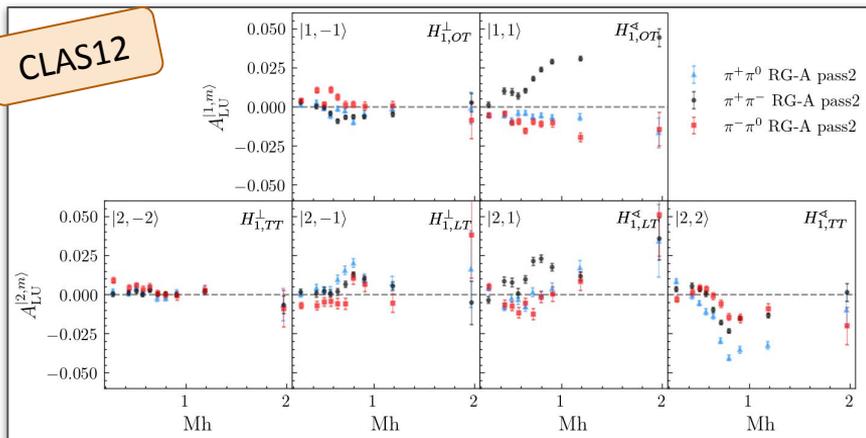
- The **CLAS12**  $e(x)$  \$money plot\$  $|1,1\rangle$  asymmetry replicated by **StringSpinner**
  - Negative asymmetry seen in **CLAS12**  $\pi^\pm\pi^0$  – pion-pair dependence of  $H_1$ ? [1]
  - Two new avenues of accessing  $e(x)$  – limited knowledge of  $H_1$  for  $\pi^\pm\pi^0$  however
  - **We still don't understand how StringSpinner recreates this asymmetry, has no  $e(x)$  parameterization!**
- Large effects in other tw.3 partial waves seen at **CLAS12**
  - Relatively muted in **StringSpinner**



# Twist-3 Partial Waves ( $\pi^+\pi^-$ , $\pi^+\pi^0$ , $\pi^-\pi^0$ )

- Never-before seen  $\varrho$ -induced resonances in  $H_1$  visible
  - Both **CLAS12** and **StringSpinner** match for  $|2,2\rangle$  partial wave (similar to tw.2 result)
- For **CLAS12** the  $|1,1\rangle$   $M_h$  dependence in  $\pi^+\pi^-$  is far stronger than  $\pi^\pm\pi^0$

★ *It is abundantly clear that VM's play a significant role towards interpreting our asymmetries* ★



# Outlook and Conclusion

- Analysis Note for  $\pi^+ \pi^0$  and  $\pi^- \pi^0$  BSA's **APPROVED** (May 2025)
  - Pass-1 RG-A data
  - $\theta$ -integrated results = 7 modulations
  - PRL paper first draft written (waiting for CLAS committee to form)
  
- Partial Wave Decomposition of five  $\pi\pi$  channels (excluding  $\pi^0\pi^0$ ) *in progress*
  - Pass-2 RG-A data
  - Full  $\theta$ -dependence captured = 12 modulations/partial waves
  - Collaborating with Chris Dilks → in development of  $\pi^+ \pi^-$  PW analysis note
  - Meeting with theorists to interpret results → comparing with theoretical models

★ *These analyses will provide the most comprehensive catalog of dihadron fragmentation to date.*

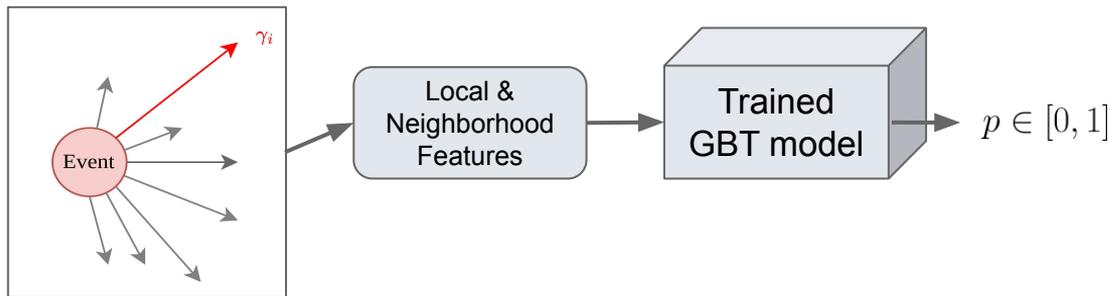


# Extra Slides

# Photon GBT Classifier

Train on intrinsic ( $E_{\text{dep}}$ ,  $\theta$ , calo-shape, etc.) and nearest neighbor (angular separation with N-nearest charged hadron, neutral particle, etc.) features (total = 16 feat.)

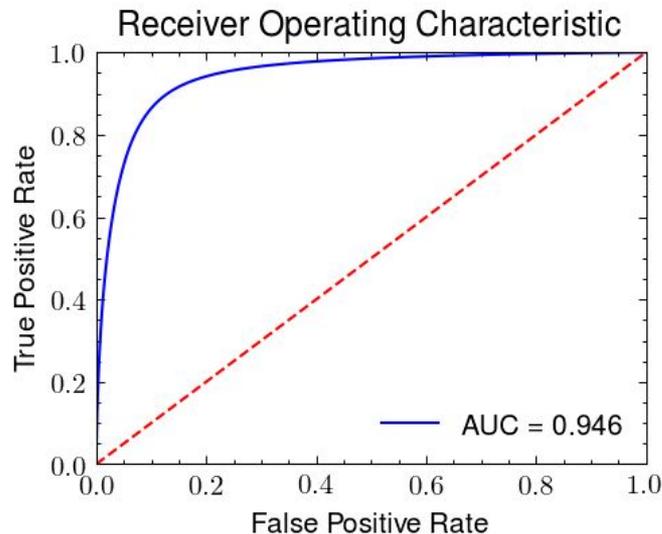
photon of interest (POI)



$p \approx 0 \rightarrow$  Photon does not have a MC match (background)

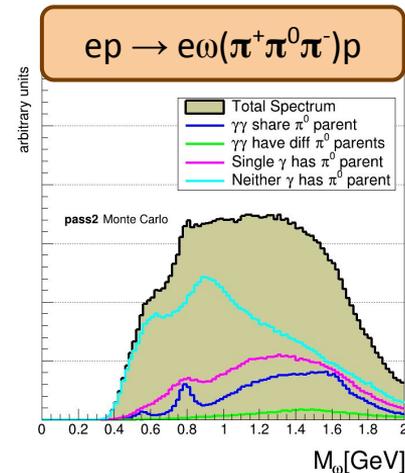
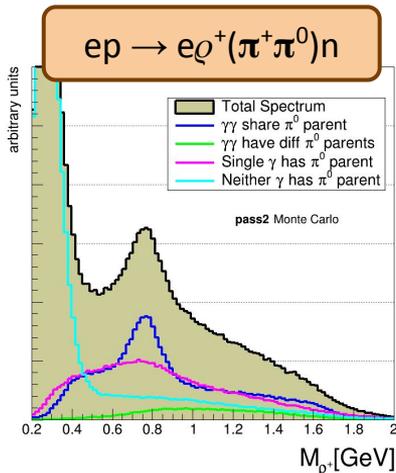
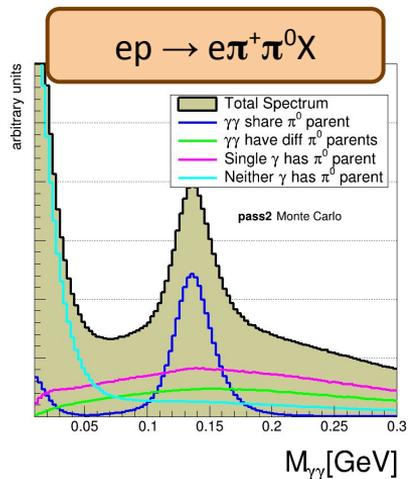
$p \approx 1 \rightarrow$  Photon likely has MC match (signal)

- LogLoss Evaluation Metric
- 75/25 split with 0.1 LR
- Symmetric Growth policy (10 depth)
- 50 generations early stopping



**Avoids learning resonant structure**

# Impact on ... SIDIS $\pi^+\pi^0$ ... Exclusive $\rho^+$ , $\omega$

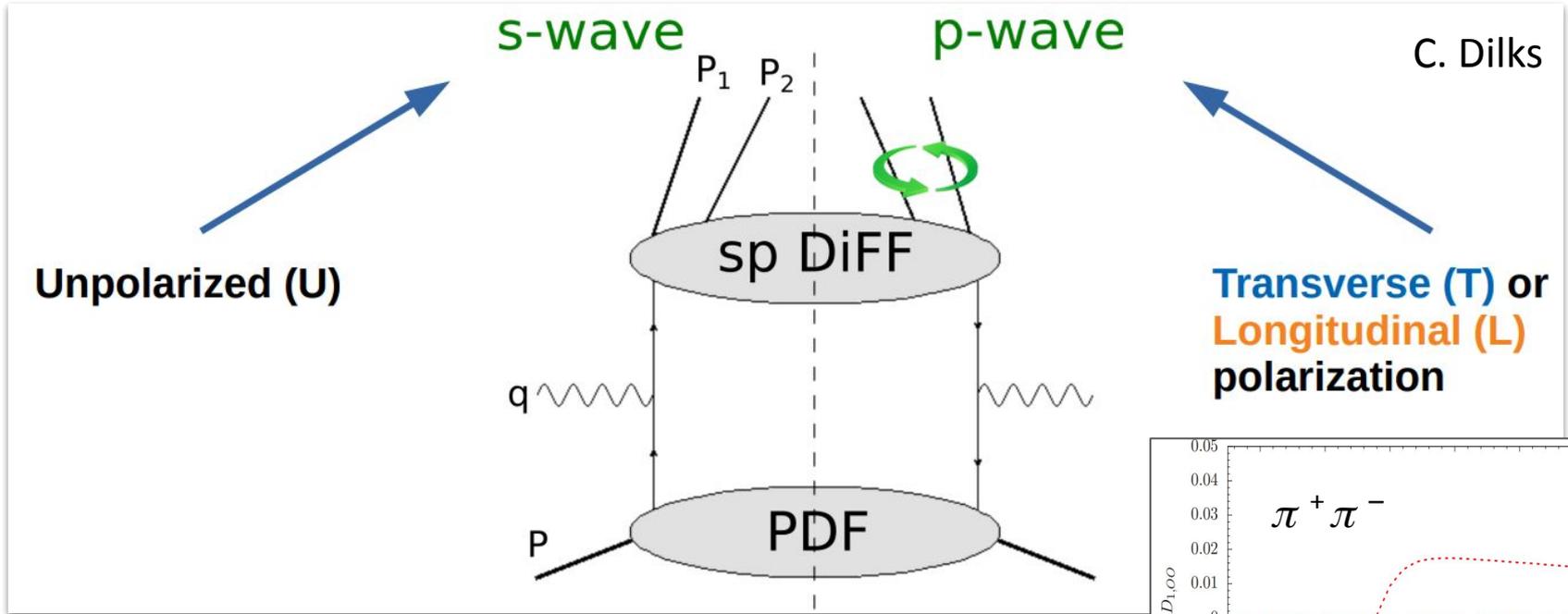


$\pi^0 \rightarrow \gamma\gamma$  background is a mix of true combinatoric (**LIME GREEN**) and false combinatoric (**MAGENTA and TEAL**)

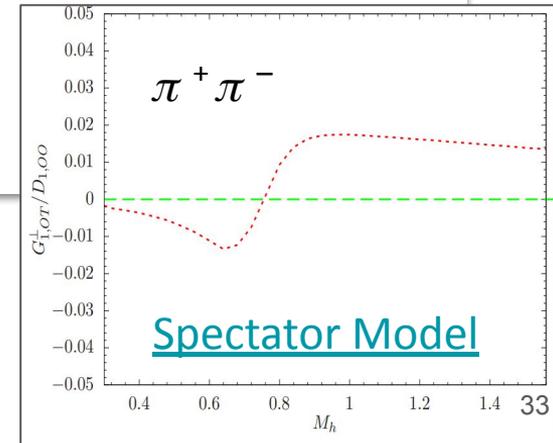
Exclusive  $\rho^+$  ( $M_{\text{miss}} < 1.2$  GeV) region is *dominated* by false combinatoric backgrounds (**MAGENTA and TEAL**)

Exclusive  $\omega$  ( $M_{\text{miss}} < 1.2$  GeV) region is *dominated* by false combinatoric backgrounds (**MAGENTA and TEAL**)

# Partial Wave Decomposition



★ Ex:  $G_{1,0T}^\perp$  is an  $|1,1\rangle$  DiFF partial wave — interference between an unpolarized pair and a transversely polarized pair ★



# Dihadron Fragmentation Functions (DiFFs)

$$\begin{aligned}
 & \left[ \begin{array}{c} |0,0\rangle \\ \nearrow \nearrow \\ \text{q} \end{array} + \begin{array}{c} |1,1\rangle \\ \nearrow \nearrow \\ \text{q} \end{array} + \begin{array}{c} |1,0\rangle \\ \nearrow \nearrow \\ \text{q} \end{array} + \dots \right]^2 \\
 = & \begin{array}{c} |0,0\rangle \quad |0,0\rangle \\ \nearrow \nearrow \nearrow \nearrow \\ \text{q} \quad \text{q}' \end{array} + \begin{array}{c} |0,0\rangle \quad |1,1\rangle \\ \nearrow \nearrow \nearrow \nearrow \\ \text{q} \quad \text{q}' \end{array} + \dots \\
 & \boxed{|L=0, M=0\rangle} \qquad \boxed{|L=1, M=1\rangle}
 \end{aligned}$$

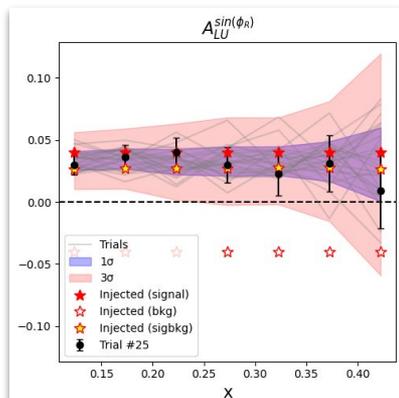
# False Asymmetry Injection Studies

- ★ By default, Monte Carlo has **helicity** =  $\mathbf{0}$
- ★ By defining false  $A^{\text{sig}}$  and  $A^{\text{bkg}}$  we can probabilistically assign **helicity** =  $\pm 1$  event-by-event

$$f(\phi_h, \phi_R; A_{\text{inj}}^{\text{sig}}, A_{\text{inj}}^{\text{bkg}}) = \begin{cases} \frac{1}{2} \left( 1 + \sum_{\ell} A_{\text{inj},\ell}^{\text{sig}} \times \psi_{\ell}(\phi_h^{\text{true}}, \phi_R^{\text{true}}) \right) & \text{if event e's } \gamma\gamma \text{ came from a } \pi^0 \\ \frac{1}{2} \left( 1 + \sum_{\ell} A_{\text{inj},\ell}^{\text{bkg}} \times \psi_{\ell}(\phi_h^{\text{true}}, \phi_R^{\text{true}}) \right) & \text{otherwise.} \end{cases}$$

Allows us to test the accuracy of our asymmetry extraction

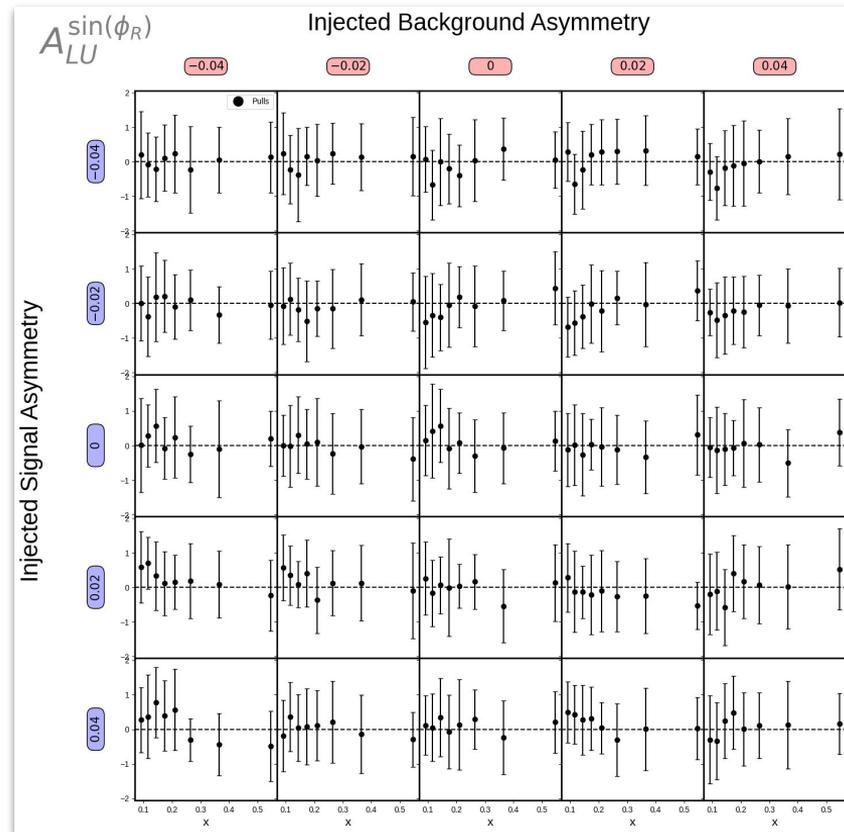
- ★ We study the pulls of with only one injection (**Case A**)



$$A_{\text{sin}(\phi_R)}^{\text{sig,bkg}} \in [-0.04, -0.02, 0, 0.02, 0.04]$$

$$A_{\text{other}}^{\text{sig,bkg}} = 0$$

Only one injected mod  
 → Only need one purity  
 (no correlations)



**1x1 purity binning**

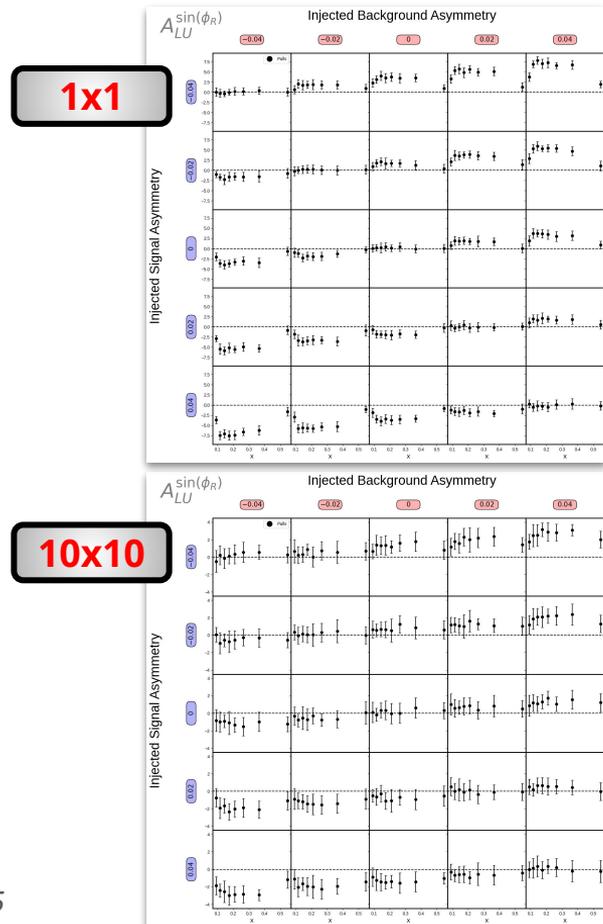
# False Asymmetry Injection Studies

- ★ In the extreme case, simultaneously inject all mods. (**Case B**)

$$A_{\text{all}}^{\text{sig,bkg}} \in [-0.04, -0.02, 0, 0.02, 0.04]$$

- ★ In the **1x1** purity binning case, the average pull across many trials climbs to  $\sim 5-8$
- ★ In the **10x10** purity binning case, the average pull across many trials climbs to  $\sim 2-3$
- ★ We find that purity binning is *more important* when the **signal** and **background** asymmetries are farther apart
  - In the analysis note, **signal** and **background** asymmetries are compared. They at most have  $\sim 0.02$  separation which purity binning cleans up

We assign a systematic error based on the choice of the purity binning scheme later



## Monte Carlo Injection Analysis Slides (**Analysis Note Review**)

# Injection Analysis — Addressing GBT vs. Legacy

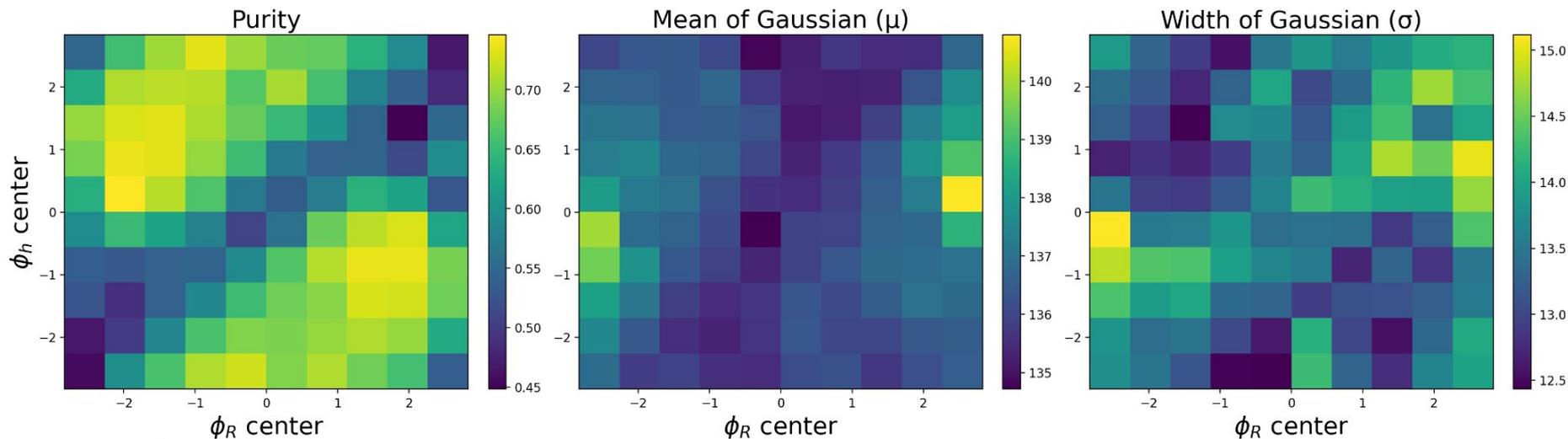
- **Observation:** Monte Carlo studies with *Legacy cuts* show that simultaneous asymmetry injections do not cause significant pulls in the extracted asym.
    - Why does this happen with the GBT analysis?
  - **Solution:** As argued in the analysis note, if the “purity X modulation<sub>i</sub>”  $u(\phi_H, \phi_R)\Psi_i$  correlates with “purity X modulation<sub>j</sub>”  $u(\phi_H, \phi_R)\Psi_j$  then they affect one another when simultaneously injected
- Interestingly, the legacy purity  $u(\phi_H, \phi_R)$  is far less symmetric than the GBT purity  $u(\phi_H, \phi_R)$ . There are *still* false photons in the legacy analyses, so the  $\phi_H, \phi_R$  distribution is more “smeared” and “uncorrelated”. Without a clear correlation, the pulls shouldn’t affect one another, as observed.

# Procedure

1. Using Monte Carlo simulated data calculate purity  $u(\phi_H, \phi_R)$  in NxN bins
2. Create toy  $\pi^+\pi^0$  data
  - a. Randomly generated phi variables
  - b. Inject a +0.04 signal asymmetry and -0.04 background asymmetry for all 7 modulations
  - c. Using the Monte Carlo purity table, generate a  $\phi_H, \phi_R$  dependent  $M_{\gamma\gamma}$  spectrum
3. Use the *same* maximum-likelihood analysis method to extract the injected asymmetries (no variable smearing). Only use **1** purity bin when extracting

**Hypothesis:** Since the toy model extraction results for **legacy purity** and **GBT purity** will only be dependent on their  $u(\phi_H, \phi_R)$  binning, a systematic difference in the extraction performance **must** be connected to it.

# purity binning gbt cuts



The **CLAS12** acceptance reveals an interesting symmetry along  $\phi_H = \phi_R$  for the  $M_{\gamma\gamma}$  distribution.

# Results of using the gbt purity binning

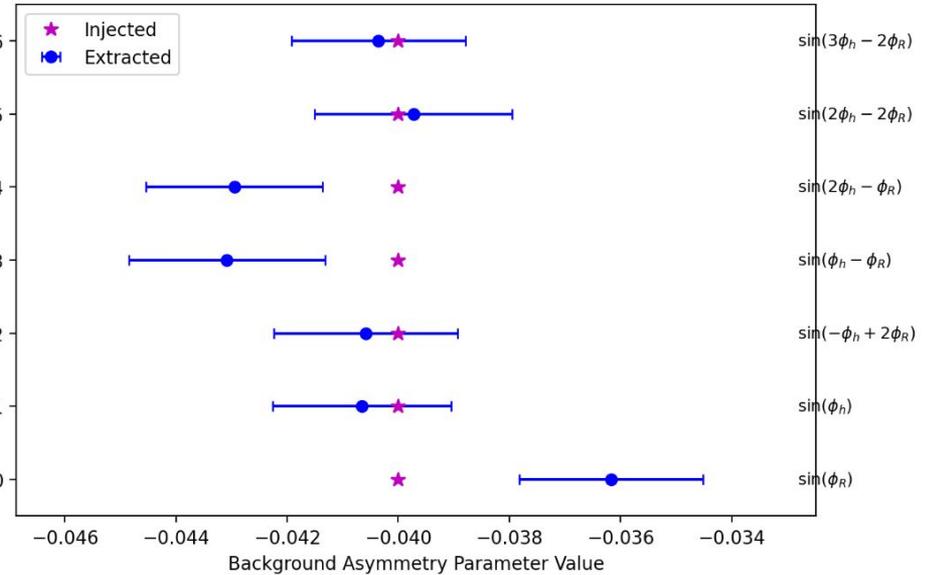
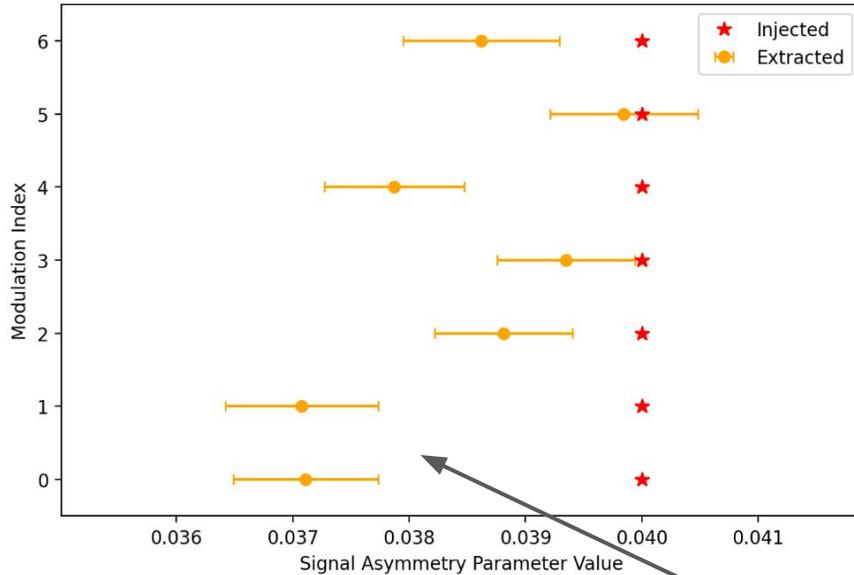
**SIGNAL INJ.**

Asymmetry Parameter Extraction Summary (100 Trials)

**BKG INJ.**

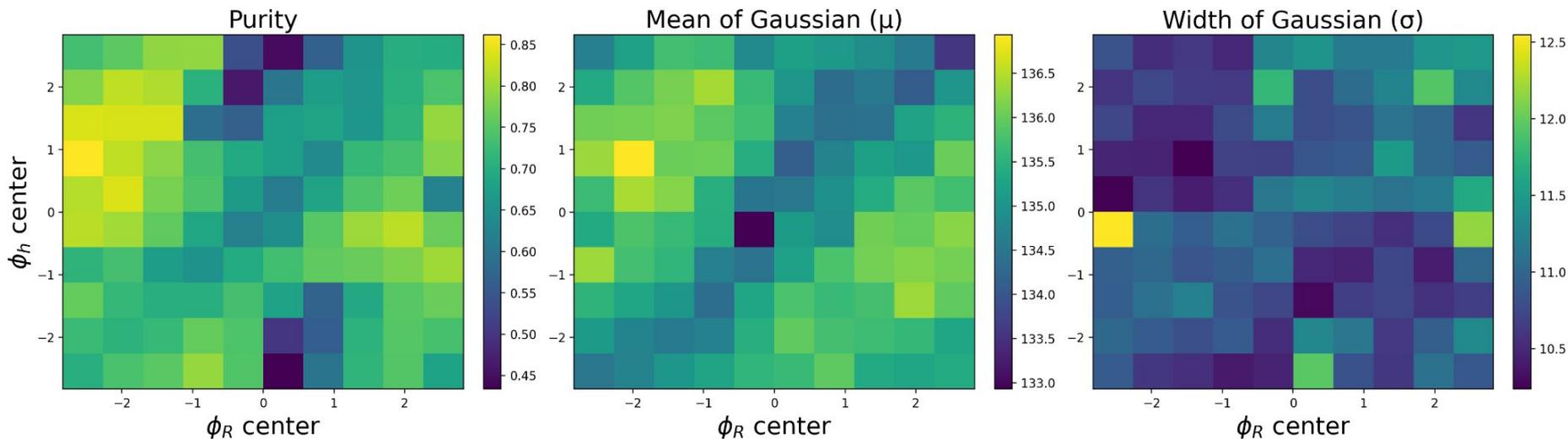
Num Modulations = 7  
Successful Trials = 100/100  
Events/Trial = 1.0e+05

Error Bars =  $\sigma_{100 \text{ trials}}/\sqrt{100}$   
Purity Mode: Single Value (Reco)



Huge pulls seen in, ex:  $\sin(\phi_h)$  asymmetries

# purity binning legacy cuts



Legacy cuts still have false photons, purity binning is more smeared (less physical), but *some* correlation still.

# Results of using the legacy purity binning

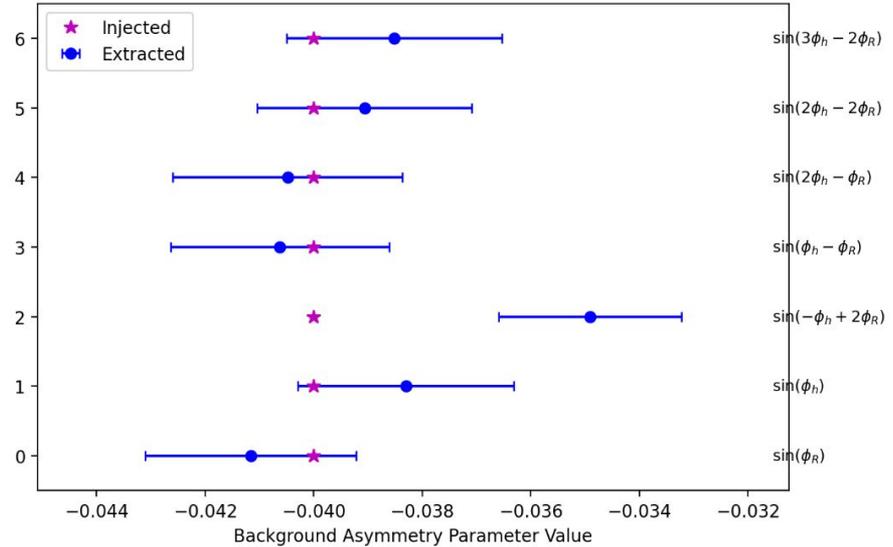
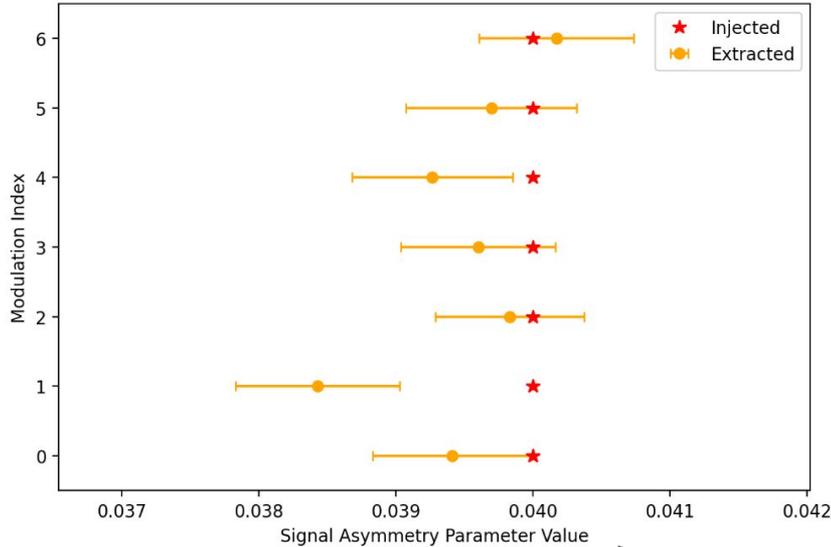
**SIGNAL INJ.**

Asymmetry Parameter Extraction Summary (100 Trials)

**BKG INJ.**

Num Modulations = 7  
Successful Trials = 100/100  
Events/Trial = 1.0e+05

Error Bars =  $\sigma_{100 \text{ trials}}/\sqrt{100}$   
Purity Mode: Single Value (Reco)



← Pulls are much more in line with +/- 1 StD Dev

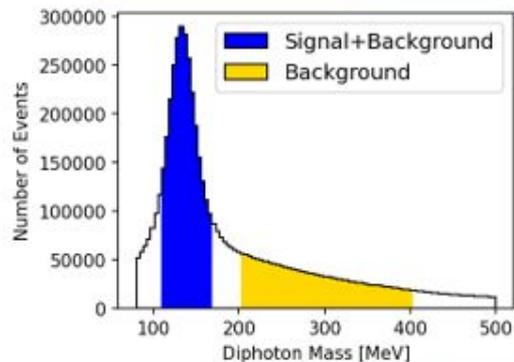
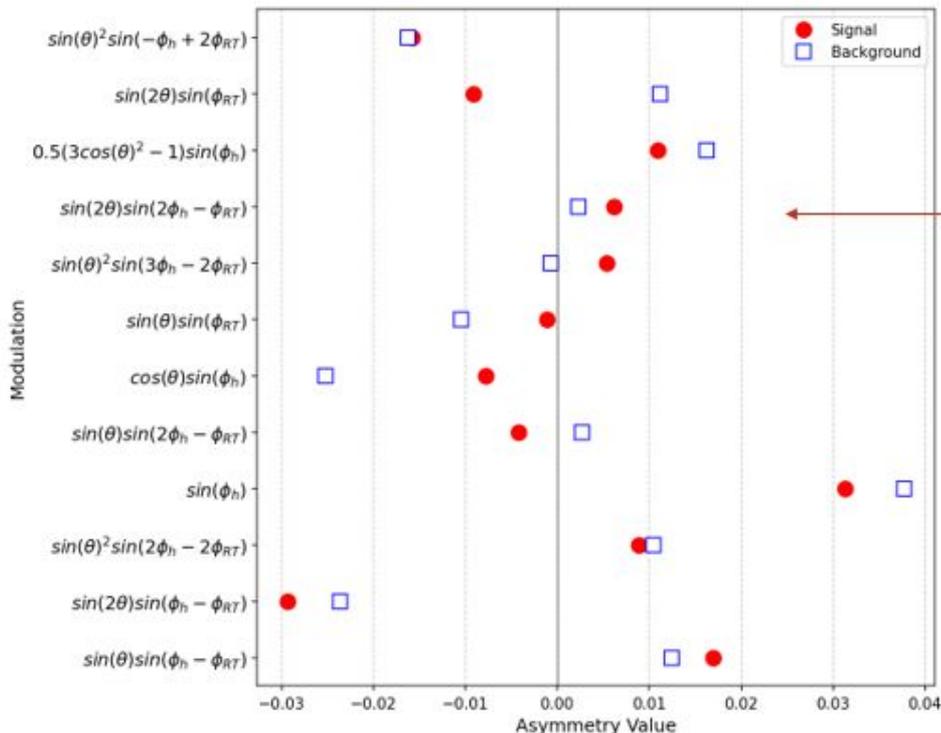
# Conclusion

- With CLAS12 acceptance, the true phi-dependent purity  $u(\phi_H, \phi_R)$  contains sharp symmetries which affect the pulls of simultaneously injected asymmetries
- The legacy purity  $u(\phi_H, \phi_R)$  does not have nearly as sharp a symmetry as the GBT purity, which we expect to better represent the actual SIDIS events

An asymmetry extraction toy model was put together with the purity binning as the **only independent** variable, validating our hypothesis.

# $\pi^0$ sideband subtraction

- Combinatorial background asymmetry (sideband) is close to signal asymmetry for most modulations (shown below for  $\pi^+\pi^0$ )



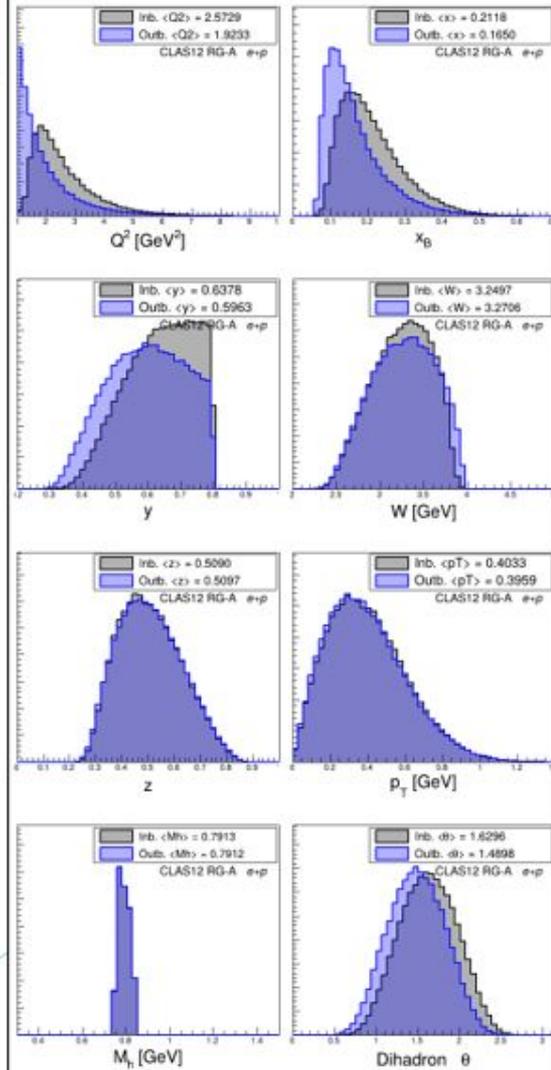
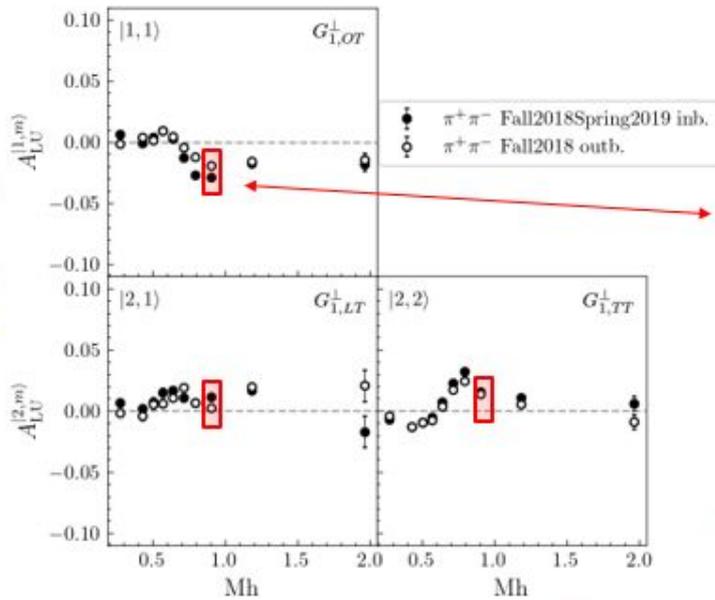
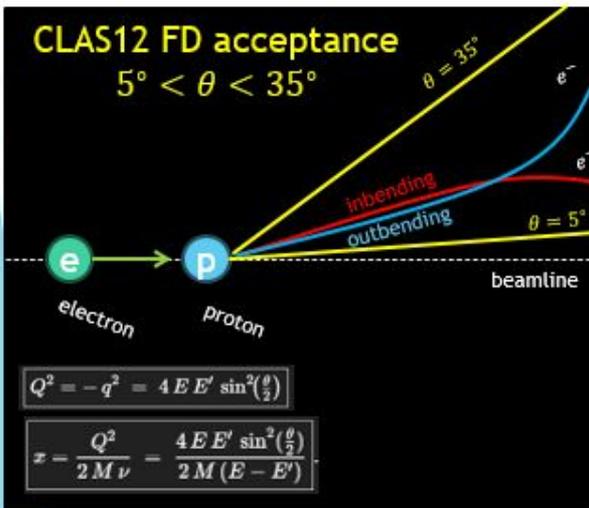
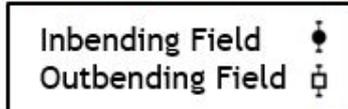
Fully-integrated results  
(no x, Mh, z binning)

*How can we explain why some modulations have a vastly different background asymmetry, while some are nearly identical?*

# $\pi^+\pi^-$ Twist-2 (proton)

## CLAS12 Inbending vs. Outbending

- ▶ CLAS12's toroidal magnet polarization effects charged particle acceptance
  - ▶ *Inbending*  $\rightarrow$  Negatively charged particles bent towards beamline
- ▶ Mainly effects electron acceptance ( $x, Q^2, y$ )  $\rightarrow$  different  $PDF(x), D(y)$ 
  - ▶ Also seems to effect  $\theta$  for some bins, could explain differences in some PW's?



## CLAS12 Run Group A vs. B

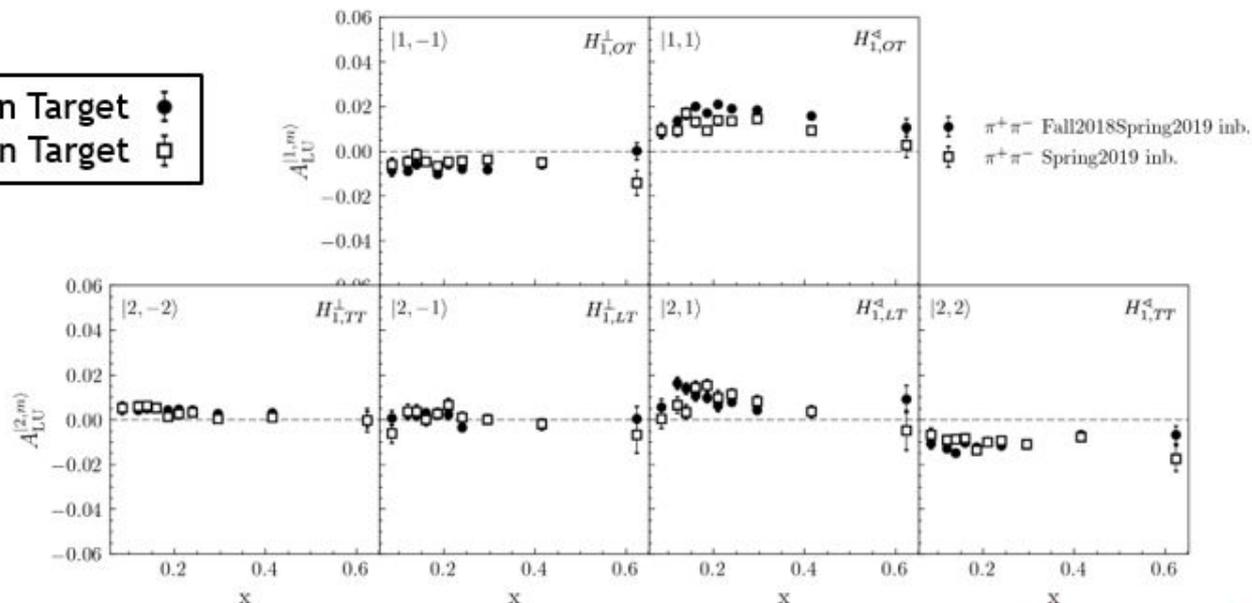
- Existing data at CLAS12 may be enough to disentangle flavor composition of  $e(x)$ 
  - We observe a slightly smaller signal in twist-3 |1,1> BSA

$$A_{LU,p}^{\text{twist } 3} \propto 4xe^{uv}(x) - xe^{dv}(x)$$

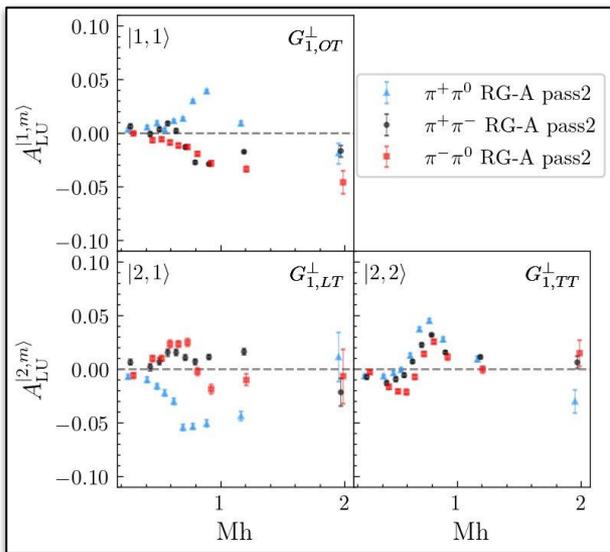
$$A_{LU,d}^{\text{twist } 3} \propto xe^{uv}(x) + xe^{dv}(x)$$

$$e \otimes H_1^\perp$$

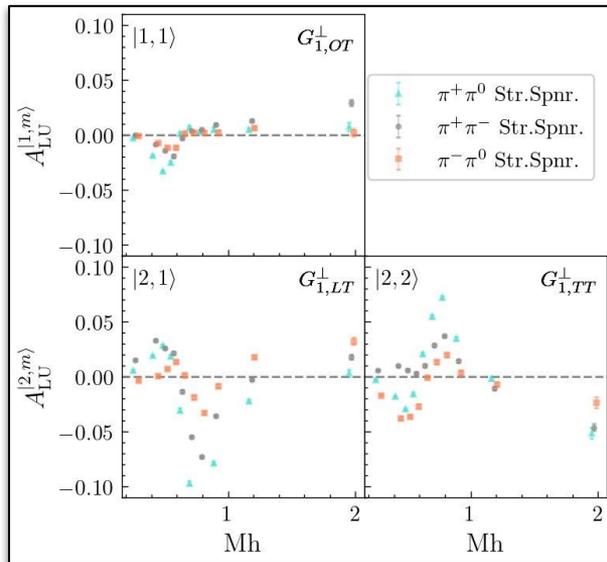
Proton Target  $\bullet$   
Deuteron Target  $\square$



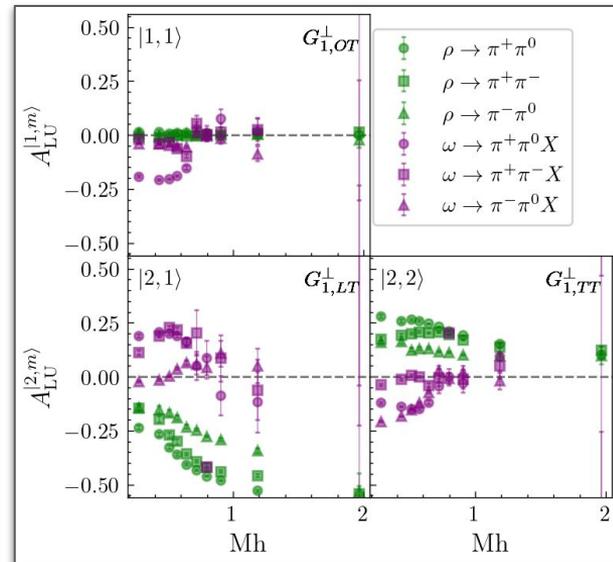
# Contributions of VM's to Partial Waves



CLAS12



StringSpinner

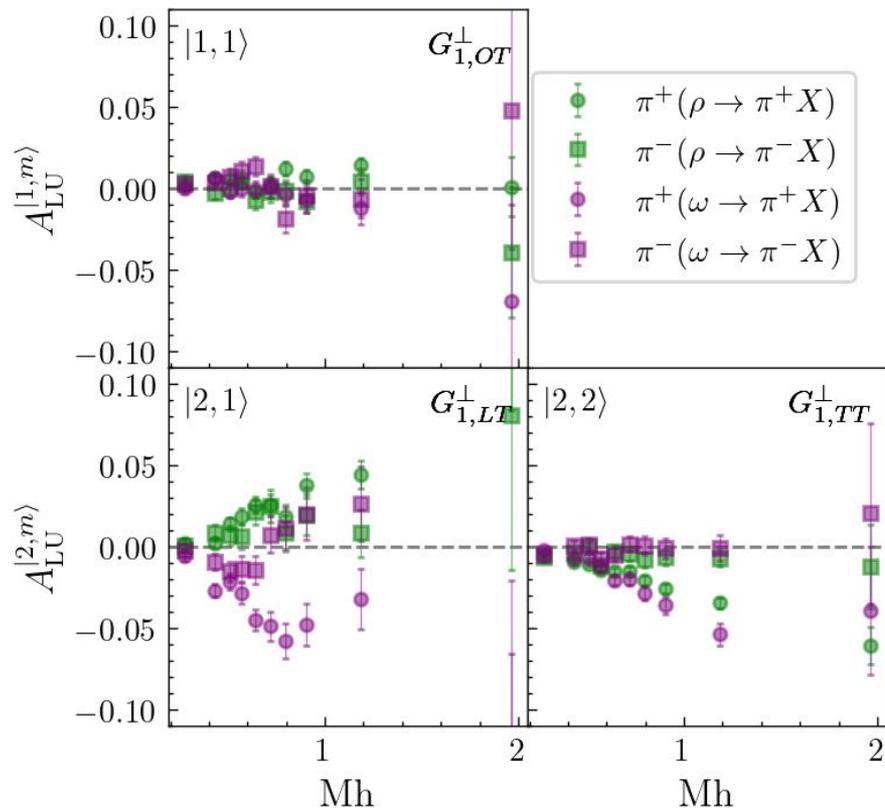


StringSpinner VM's

# VM Origin's of Like-Pion Partial Waves

**Premise:** VM decays only one of the pions (other *can* come from a VM, but not the same)

Looks as if coincidenting on the  $\omega$  can generate asymmetries in like-pion pairs



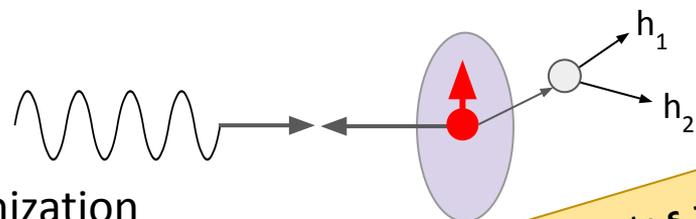
# Dihadron Production

$$e(\ell) + p(P) \rightarrow e(\ell') + h_1(P_1) + h_2(P_2) + X$$

➤ Dihadron SIDIS production cross section:  $\sigma = \hat{\sigma} \otimes \text{PDF} \otimes \text{DiFF}$

➤ **(PDFs)** Probe quark-gluon dynamics in proton

- Twist-3 PDF  $e(x)$  directly accessible without  $p_T$  convolutions
- $e(x)$  related to forces between quarks and gluons and  $\sigma_{\text{TTN}}$



➤ **(DiFFs)** Probe more complex correlations in hadronization

- Interference Fragmentation Functions ( $h_1, h_2$  produced in **s** and **p** wave)
- Impact of quark polarization on hadronization
- Comparing different pion channels (ex:  $\pi^+\pi^-$  and  $\pi^+\pi^0$ )  $\rightarrow$  compare with theory
- Separate Resonant production ( $\rho^0, \rho^{+/-}$ ) vs. Non-Resonant fragmentation

# Systematic Error (1/6) Beam Polarization

- ★ Uncertainty does not depend on the kinematic bin or the modulation studied

$$\frac{\sigma_{A_{\text{true}}}}{A_{\text{true}}} = \frac{\sigma_P}{P}$$

- ★ 2% systematic error in the reported polarization assumed
- ★ For  $\pi^+\pi^0$  we use the **Total Inbending**, for  $\pi^+\pi^0$  we use the **Fall2018 Outbending**

Run Period	Beam Polarization	Number of Triggers	$\Delta_{\text{ALU}}/\text{ALU}$
Fall 2018 Inbending	$85.92\% \pm 1.29\% \pm 2\%$	$1.04\text{e}^{10}$	2.8%
Fall 2018 Outbending	$89.22\% \pm 2.509\% \pm 2\%$	$1.346\text{e}^{10}$	3.6%
Spring 2019 Inbending	$84.53\% \pm 1.474\% \pm 2\%$	$1.262\text{e}^{10}$	2.9%
Total RG-A Inbending	$85.16\% \pm 1.7\% \pm 2\%$	$2.302\text{e}^{10}$	3.1%

TABLE 18: Collection of beam polarizations for the RG-A experiment and their impact on the extracted BSA uncertainties.

# Systematic Error (2/6) Bin Migration

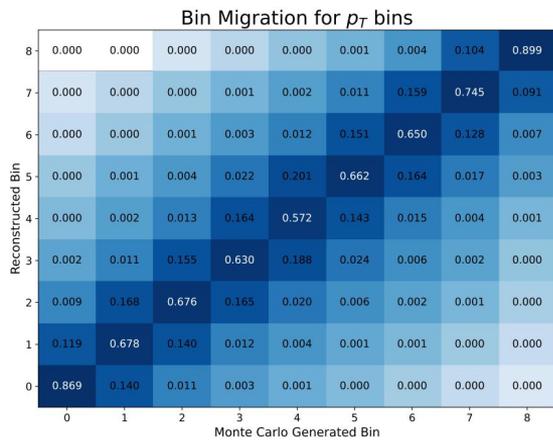
- ★ Relative uncertainty depends on the kinematic bin studied and modulation #

$$\Delta A_i = \sum_{k=1}^K (f_{i+k,i} A_{i+k} + f_{i-k,i} A_{i-k} - f_{i+k,i} A_i - f_{i-k,i} A_i)$$

Error in bin “i”

Asymmetries leaking into bin “i”

Asymmetry leaking out of bin “i”



Bin 0:  $0.00 < p_T < 0.18$  Bin 1:  $0.18 < p_T < 0.25$   
 Bin 2:  $0.25 < p_T < 0.32$  Bin 3:  $0.32 < p_T < 0.38$   
 Bin 4:  $0.38 < p_T < 0.43$  Bin 5:  $0.43 < p_T < 0.50$   
 Bin 6:  $0.50 < p_T < 0.57$  Bin 7:  $0.57 < p_T < 0.68$   
 Bin 8:  $0.68 < p_T < 1.82$

- ★  $f_{i,j} \rightarrow$  fraction of reconstructed events in bin “i” that were generated in bin “j”

Modulation	$PT_0$	$PT_1$	$PT_2$	$PT_3$	$PT_4$	$PT_5$	$PT_6$	$PT_7$	$PT_8$
$\sin(\phi_R)$	0.0032	0.041	0.021	0.029	0.018	0.06	0.15	0.0062	0.015
$\sin(\phi_h)$	0.089	0.0088	0.027	0.012	0.1	0.064	0.022	0.043	0.034
$\sin(-\phi_h + 2\phi_R)$	0.11	0.015	0.013	0.017	0.07	0.07	0.0071	0.063	0.057
$\sin(\phi_h - \phi_R)$	0.46	0.057	0.12	0.049	0.1	0.036	0.075	0.033	0.029
$\sin(2\phi_h - \phi_R)$	0.19	0.098	0.53	0.097	0.47	0.14	0.14	0.055	0.034
$\sin(2\phi_h - 2\phi_R)$	0.0094	0.26	0.05	0.038	0.049	0.078	0.048	0.02	0.055
$\sin(3\phi_h - 2\phi_R)$	0.016	0.036	0.012	0.017	0.51	0.25	0.041	0.0066	0.026

$\Delta A/A$  for  $\pi^+\pi^0$   
 binned in  $p_T$

# Systematic Error (3/6) Particle Misidentification

- ★ Relative uncertainty depends on the kinematic bin studied
- ★ Assumes that the fractional asymmetry from misidentified events is much less than that of correctly identified events

$$\frac{\Delta A}{A} = \frac{f_m}{1 - f_m} \quad \star \quad f_m \rightarrow \text{fraction of events with a misidentified } e, \gamma, \text{ or } \pi$$

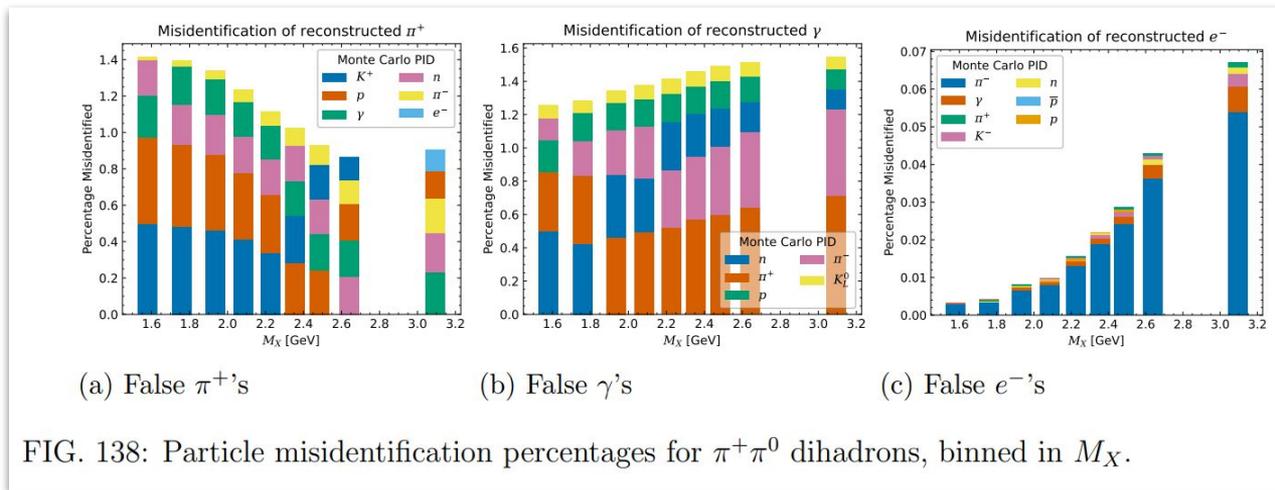


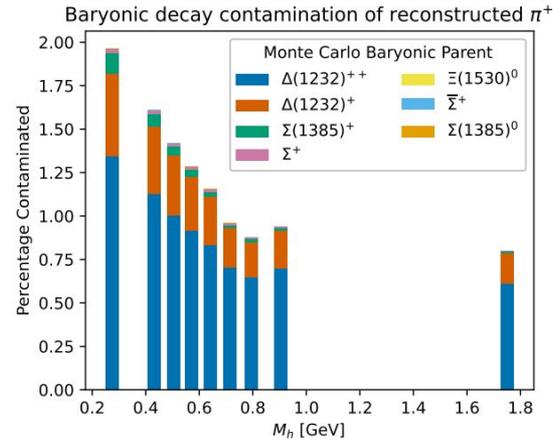
FIG. 138: Particle misidentification percentages for  $\pi^+\pi^0$  dihadrons, binned in  $M_X$ .

# Systematic Error (4/6) Baryon Decays

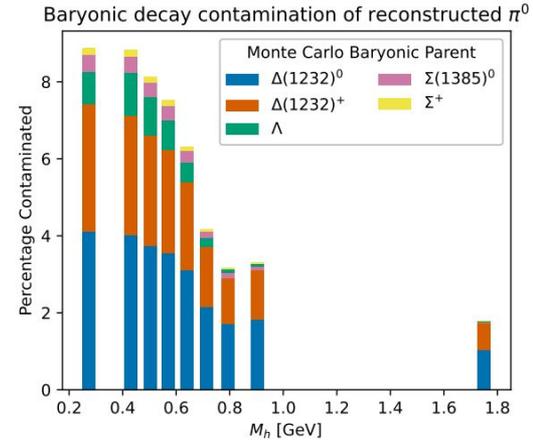
- ★ Relative uncertainty depends on the kinematic bin studied
- ★  $e p \rightarrow e \pi^{\pm} \pi^0 X$  where either  $\pi$  originates from a baryon's decay is background
- ★ Assumes that the fractional asymmetry from baryonic decays is much less than that from traditional fragmentation

$$\frac{\Delta A}{A} = \frac{f_b}{1 - f_b}$$

- ★  $f_b \rightarrow$  fraction of events where one or more of the pions originated from a baryon's decay



(a) Parents of  $\pi^+$ 's

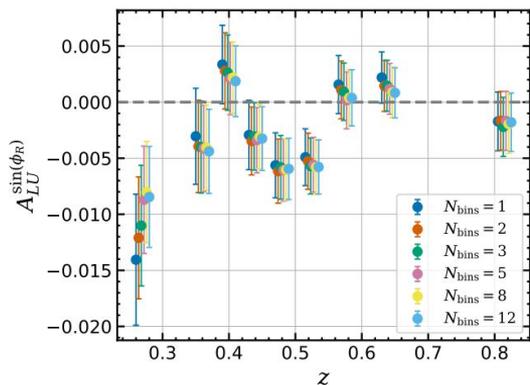


(b) Parents of  $\pi^0$ 's

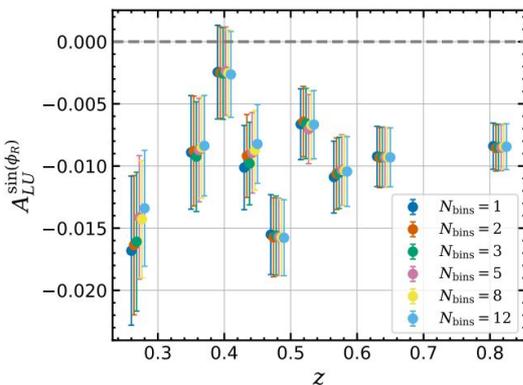
FIG. 144: Percentages of baryonic parents for  $\pi^{\pm}\pi^0$  dihadrons, binned in  $M_h$ .  
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# Systematic Error (5/6) Purity Binning

- ★ Relative uncertainty depends on the kinematic bin studied and modulation
- ★ The choice of how many bins to subdivide  $(\phi_h, \phi_R)$  can have a slight impact on the asymmetry extracted
- ★ Estimate the systematic by taking half the range across multiple options



(a)  $\pi^+\pi^0$  dihadron pairs



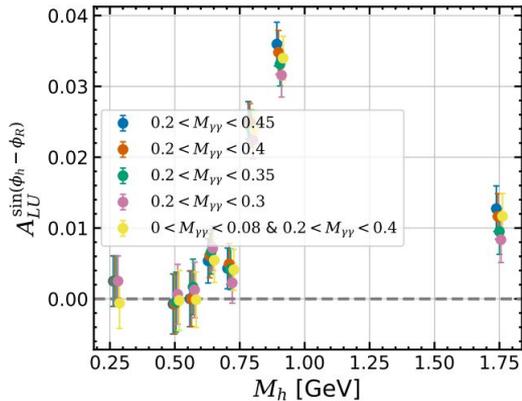
(b)  $\pi^-\pi^0$  dihadron pairs

FIG. 208:  $A_{LU}^{\sin(\phi_R)}$  results for changing the number of rows and columns in the  $u(\phi_h, \phi_R)$  purity binning scheme, binned in  $z$ .

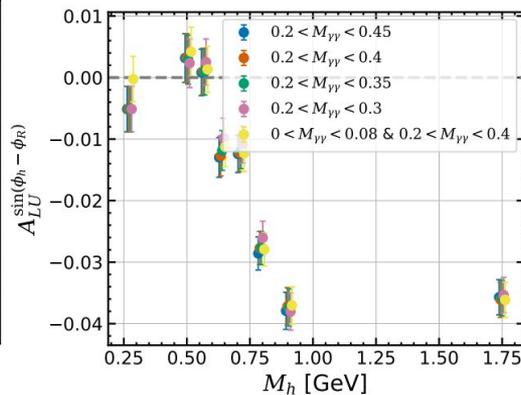
In this hand-picked example, the purity binning has a notable impact at low  $z$ , with nearly an error bar of separation between a **1x1** and **8x8** grid

# Systematic Error (6/6) Sideband Region

- ★ Relative uncertainty depends on the kinematic bin studied and modulation
- ★ The choice of where to define the “**background**” (a.k.a sideband) region may have a non-negligible impact on the  $A^{\text{bkg}}$  and therefore  $A^{\text{sig}}$  extracted
- ★ Estimate the systematic by taking half the range across multiple options

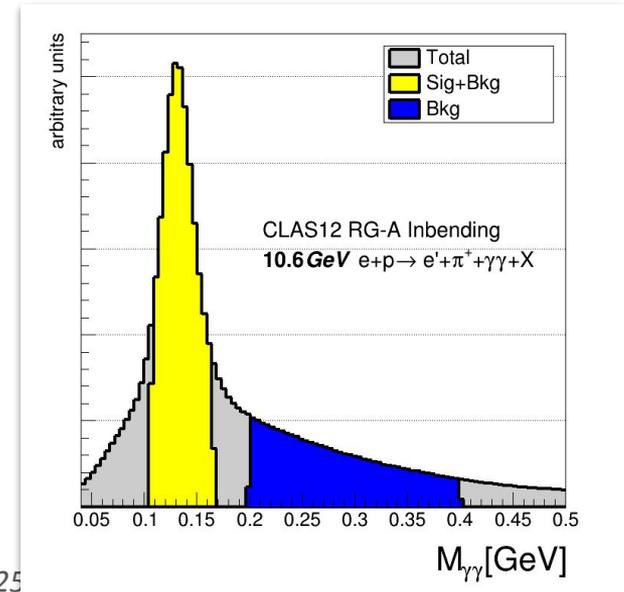


(a)  $\pi^+\pi^0$  dihadron pairs



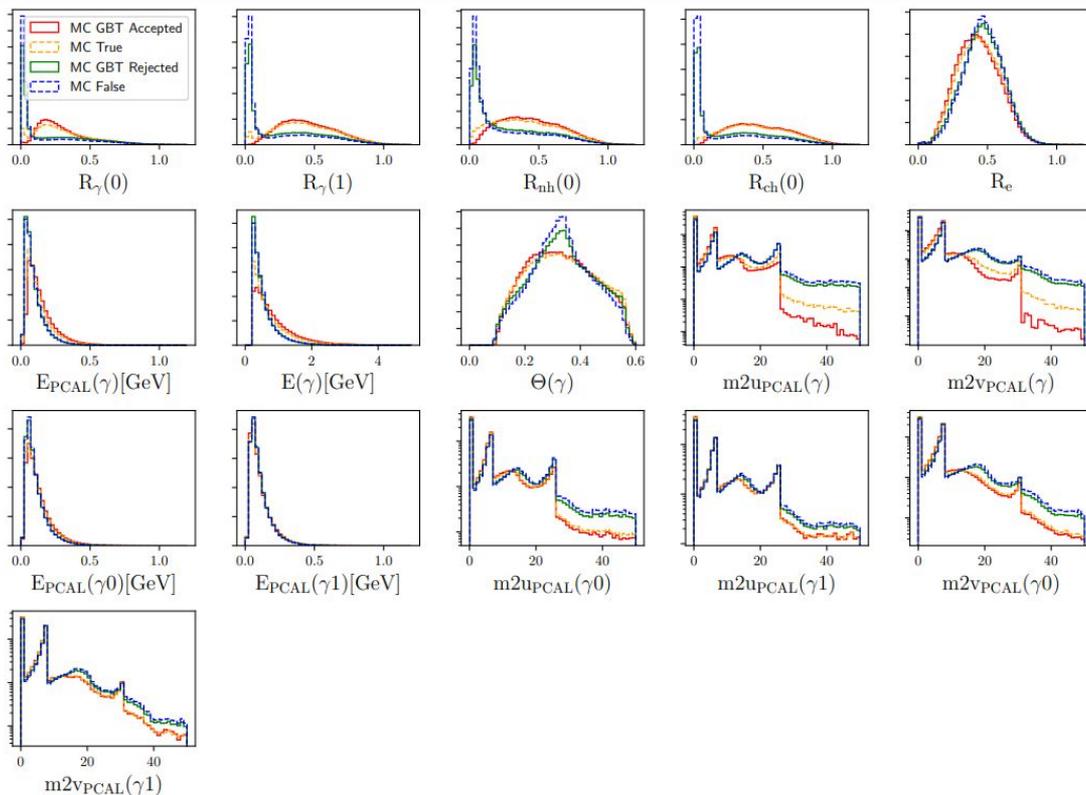
(b)  $\pi^-\pi^0$  dihadron pairs

FIG. 260:  $A_{LU}^{\sin(\phi_h - \phi_R)}$  results for changing the definition of the  $M_{\gamma\gamma}$  sideband region(s) for the background asymmetries, binned in  $M_h$ .



# Input Features (Monte Carlo Inbending)

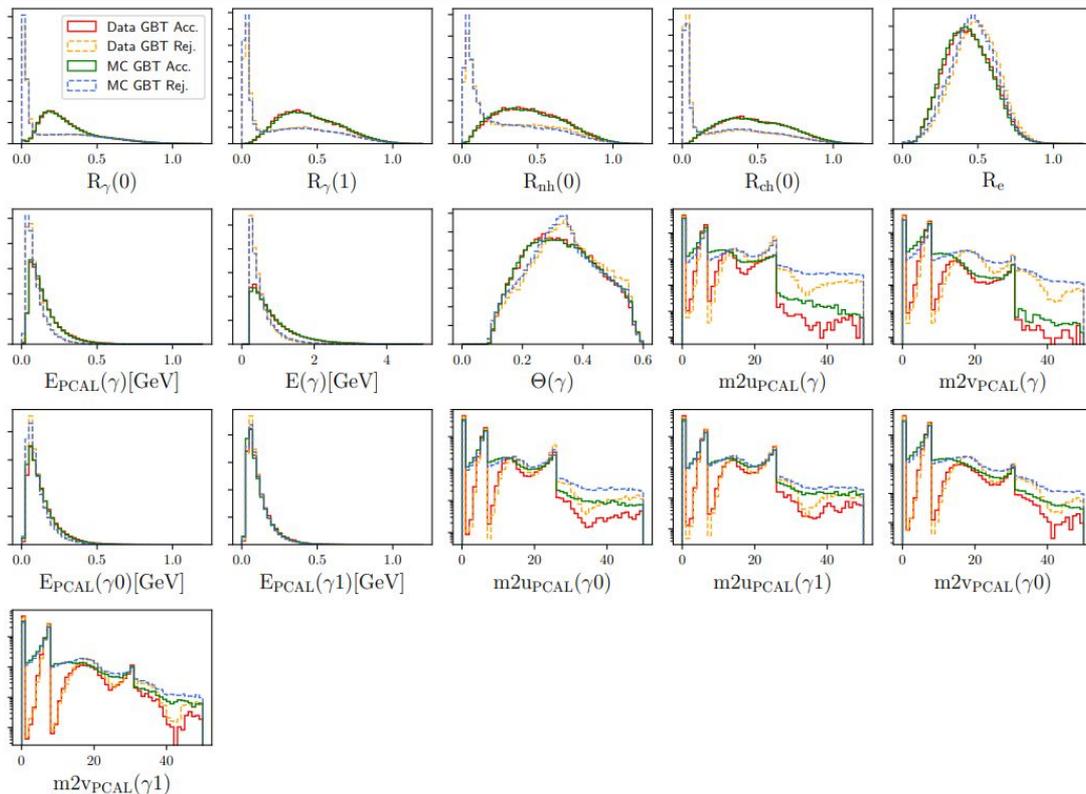
$R_n(\#) \rightarrow$  Angle between nearest “#” neighbor and photon of interest



...

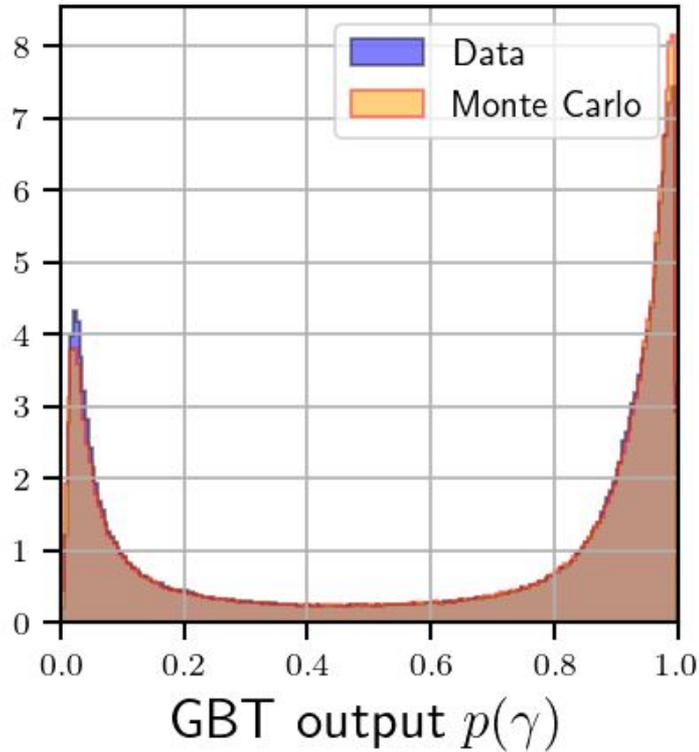
# Input Features (Monte Carlo vs Data)

$R_n(\#) \rightarrow$  Angle between nearest “#” neighbor and photon of interest



...

# GBT Output



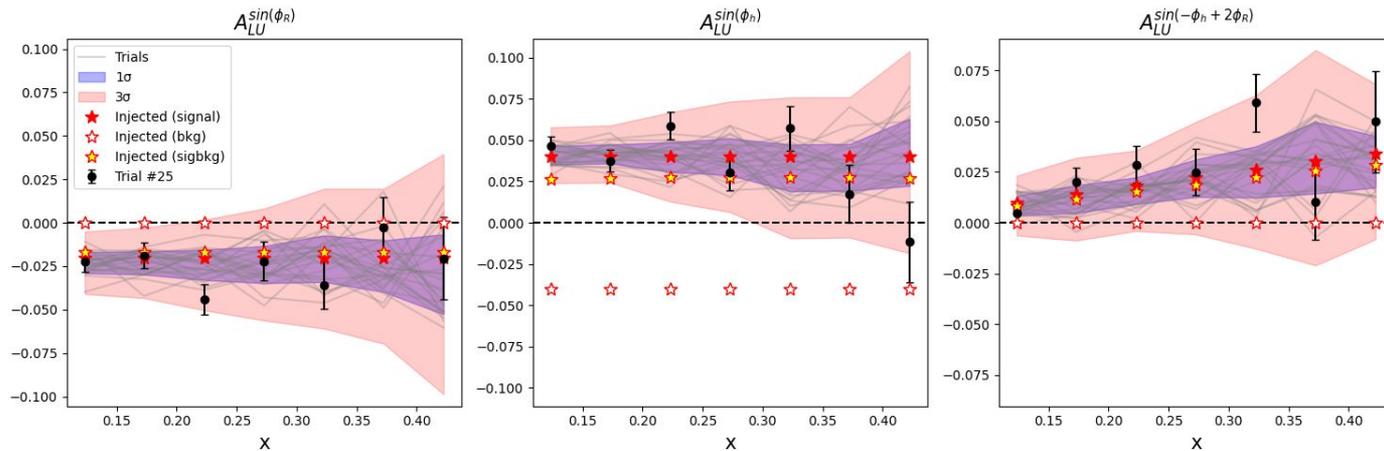
- ★ For each photon in **data** and **Monte Carlo** we histogram the GBT output value
  - We see that the aggregate outputs are *very similar* → indicates that the feature spaces are very similar
- ★ Results speaks to our confidence that the predictions made on data can be trusted

# Example Injection (3 modulations)

$A_{\text{signal}}$        $A_{\text{bkg}}$       Modulation

```
injector.load_weight_funcs( [  
  ["-0.02" , "-0.00", "sin(:phi_R0:)"] ,  
  ["0.04"  , "-0.04", "sin(:phi_h:)"] ,  
  ["0.08*:x:", "0.00", "sin(-:phi_h:+2*:phi_R0:)"] ,  
  ] )
```

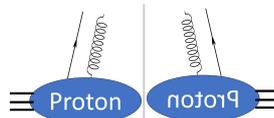
```
injector.addSideband("signal+bkg",0.115,0.15)  
injector.addSideband("bkg",0.24,0.4)
```



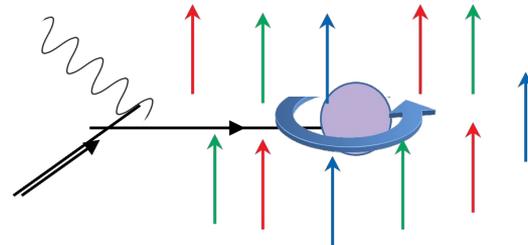
# The Twist-3 PDF $e(x)$

➤ Quark-gluon interactions play a significant role in describing the proton

- Encoded in higher twist PDFs such as  $e(x)$
- Interference of coherent scattering of  $qg$  pair
- We rely on quark-gluon interactions to manifest proton spin and mass



[Phys.Rev.D 88 \(2013\) 114502](#)



➤  $e(x)$  is better understood through its moments

- $x^2$  moment sensitive to transverse force experienced by transversely polarized quark in a transversely polarized  $N$  [1]

★ **QUESTION:** “Why study 2-hadron SIDIS?” ★

★ **ANSWER:** In the dihadron BSA,  $e(x)$  appears without transverse momentum-dependence! ★

$$d^7\sigma_{LO} = \frac{\alpha^2}{2\pi Q^2 y} \lambda \sum_a e_a^2 W(y) \sin \phi_R \frac{|\vec{R}_T|}{Q} \left[ \frac{M}{M_h} x e(x) H_1^{\triangleleft}(z, \zeta, M_h^2) + \frac{1}{z} f_1(x) \tilde{G}^{\triangleleft}(z, \zeta, M_h^2) \right]$$

★ This is not the case in 1-hadron SIDIS —  $e(x)$  extraction requires TMD modeling *and* appears in a structure function with 4 PDF⊗FF pairs ★