



Measurement of DVCS observables with ey detection

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INTRODUCTION Generalized Parton Distributions

Motivation

1.INTRODUCTION

At low energy QCD, we describe the nucleon structure in terms of structure functions including:

- **1**. Form Factors describe transverse position of partons
- 2. Parton Distribution Functions describe longitudinal momentum distributions
- **3**. Generalized Parton Distributions (GPDs) correlates transverse position and longitudinal momentum distributions



GPDs encode information of the nucleon structure such as:



Contributions to the nucleon total spin.

- X. Ji, Phy.Rev.Lett.78,610(1997) $=\frac{1}{2}\Delta\Sigma+\Delta L+\Delta G$
- Quark contribution is not the main contribution \rightarrow Spin Crisis
- Quark's orbital angular momentum is accessed through GPDs
- Gluon contribution



Access to Gravitational Form Factors.



- V. D. Burkert et al. Nature 557.7705 (2018): 396
 - Mass/Energy distribution inside the nucleon
 - Forces distribution.
 - Nucleon radius
 - Shear forces and pressure distribution

1. INTRODUCTION



The proton chiral-even GPDs can be accessed via electro-production of a real photon.

GPDs enters the DVCS-BH interference term through **Compton Form Factors** $\mathcal{F}(\boldsymbol{\xi},t) = \sum_{q} e_q^2 \left\{ \mathcal{P} \int_{-1}^{1} dx \ \boldsymbol{F}^q(x,\boldsymbol{\xi},t) \left[\frac{1}{x-\boldsymbol{\xi}} + \frac{1}{x+\boldsymbol{\xi}} \right] - i\pi \left[\boldsymbol{F}^q(\boldsymbol{\xi},\boldsymbol{\xi},t) - \boldsymbol{F}^q(-\boldsymbol{\xi},\boldsymbol{\xi},t) \right] \right\}$





For a spin $\frac{1}{2}$ particle, there are four chiraleven GPDs

 $F=H,E,\widetilde{H},\widetilde{E}$

 $= \sum_{q} e_{q} \left\{ r \int_{-1}^{ax} r^{-(x,\xi,t)} \left[\frac{1}{x-\xi} + \frac{1}{x+\xi} \right] = i\pi \left[r^{-(\xi,\xi,t)} - r^{-(-\xi,\xi,t)} \right] \right\}$ which are measured in the Beam Spin Asymmetry

$$\mathbf{BSA} = rac{1}{P_{ ext{beam}}} rac{d^4 \sigma^+ - d^4 \sigma^-}{d^4 \sigma^+ + d^4 \sigma^-} \ \propto \Im \mathfrak{m} \left\{ F_1 \mathcal{H} + \xi' (F_1 + F_2) \widetilde{\mathcal{H}} - rac{t}{4M_N^2} F_2 \mathcal{E}
ight\}$$
5

1.INTRODUCTION

1. Detecting the proton is ideal, but by not detecting it:

- There is access to the small $-t=(p-p')^2$ region
- Larger statistics can be achieved, as the proton is not constrained by detector acceptance
- It is complementary to the existing measurements and studies:
 - pass1 BSA RG-A results [Christiaens, G., et al. Phys. Rev. Lett. 130.21 (2023): 211902.]
 - pass1 cross-section measurements (Sangbaek)
 - pass2 BSA with proton detection (Timothy)

2. However,

- There are background contributions from the whole SIDIS spectra.
- There are reduced options for cuts. Missing proton mass is the only exclusivity variable available

3. Therefore,

- - Then applied to the no-proton case



• We move to a ML approach for channel selection. It is first validated by including the proton information.



DATA SELECTION



2. DATA SELECTION

Data from the fall-2018 RG-A dataset.

- Inbending and outbending torus configurations.
- Pass-2 reconstruction

KINEMATICS

- W>2 GeV
- Q²>1 GeV²
- q' >2 GeV
- e'>1 GeV
- (p' > 0.3 GeV)

EXCLUSIVITY CUTS

(if proton detected)

- $\Delta \Phi = |\Phi(p') \Phi(\gamma)| < 2^{\circ}$
- $\Delta t = |t(p')-t(\gamma)| < 2 \text{ GeV}^2$
- Pmiss<1 GeV

 $\phi(p')$ uses γ^* and p' $\phi(\gamma)$ uses γ^* and γ $t(p') = (p-p')^2$

$$\mathsf{t}(\mathbf{\gamma}) = \frac{Q^2 M_p + 2\nu M_p \left(\nu - \sqrt{\nu^2 + Q^2} \cos \theta_{\gamma^* \gamma}\right)}{\sqrt{\nu^2 + Q^2} \cos \theta_{\gamma^* \gamma} - \nu - M_p}$$



The event is built with the (e,γ,p) set with minimum missing ep → eγp missing mass

2. DATA SELECTION





2. CHANNEL SELECTION





-0.2

-0.3

-0.1

0

0.1

0.2

0.3

BDT response

After **BDT** classification, residual background is removed with well known methods

2. CHANNEL SELECTION: BACKGROUND SUBTRACTION

BDT classification leave a leftover contamination



The entirety of the background cannot be removed with a BDT

To estimate and remove the residual background, we estimate the 1y/2y - phase space ratio using two methods

Method 1: Let us define:

> $\square n_{MC/Data}^{1\gamma} = \text{Number of}$ simulated π^0 events that pass the DVCS analysis.

 $\square n_{MC/Data}^{2\gamma} = \text{Number of}$ simulated π^0 events that are reconstructed.

The contamination is then:

 $n_{Data}^{1\gamma} = \left(\frac{n_{MC}^{1\gamma}}{n_{MC}^{2\gamma}}\right) n_{Data}^{2\gamma}$

 $rac{(N_{DVMP}^{MC})_{ep
ightarrow ep \gamma \gamma}}{(N_{DVMP}^{MC})_{ep
ightarrow ep \gamma \gamma}} = rac{(N_{DVMP}^{Data})_{ep
ightarrow ep \gamma \gamma}}{(N_{DVMP}^{Data})_{ep
ightarrow ep \gamma \gamma}}$

Method 2:

- **1.** Reconstruct π^0 events.

- **2.** For each π^0 , generate 1500 decays.
- 3. If the event pass the DVCS analysis with any photon, fill histograms.
- 4. If the event pass the DVMP analysis, increment $n_{MC}^{2\gamma}$ by the reconstruction efficiency.
- 5. At the end of the decays, DVCS events are normalized by $1/n_{MC}^{2\gamma}$.

2. CHANNEL SELECTION



simulation and data.







DVCS MEASUREMENTS

Beam Spin Asymmetry Cross-section



3. DVCS BSA



Results compatible with the published pass1 RG-A results



Pass-2 ey, eyp and RG-A measurements are compatible







The eyp selection favors large -t measurements



 $1 < \! Q^2 ({
m GeV}^2) < 1.6$ $0.17 < \! x_B < 0.27$ $0.25 < - t ({
m GeV}^2) < 0.4$



The ey gives unique access to the small -t region





Systematic uncertainties come from
Selection cuts
BDT classification
Background subtraction
Bin migration

Statistical and systematic errors can be at the same level for for approaches

BSA amplitude as a function of -t





• The eγ selection gives access to the small -t region Results compatible with the KM15 model prediction

4. DVCS MEASUREMENTS: BEYOND BSA

We foresee cross section measurements for a separate work.





Final normalization should account for electron and photon detection efficiency. Yet, preliminary results are in a good track





4. SUMMARY

- GPDs are crucial for understanding the nucleon structure as they encode a variety of properties.
- DVCS is a golden channel for GPD studies due to its direct access to GPDs.
- While detecting the DVCS recoil proton is ideal, by not detecting it
 - There is a boost in statistics
 - There is access to a larger phase space
- Most recent BSA measurements were presented.
 - the ey and eyp selection are complementary to each other
 - \circ BSA measurements are compatible with the KM15 model predictions at low (ey) and high $(e_{\gamma}p)$ -t values.
 - Systematic errors were estimated
 - Analysis note is currently is ready for a 3rd third round of comments.
- Cross-section measurements with e_{γ} selection are foreseen for a separate work.



THANKS

Systematic error

- PID selection:
 - $-13 < v_z^{electron} < 12$,
 - FD photons and electrons detected in different sectors.
 - $|v_z^{electron} v_z^{proton}| < 20$ cm for photons as recommended on [10].
 - For electrons with momentum greater than 4.5 GeV:

$$\frac{E_{dep}^{ECIN}}{p_e} > 0.2 - \frac{E_{dep}^{PCAL}}{p_e}$$

• Exclusivity cuts:

$$- |\Delta t| < 2 \text{ GeV}^2 \rightarrow |\Delta t| < 1.5 \text{ GeV}^2$$
$$- |\Delta \phi| < 2^\circ \rightarrow |\Delta \phi| < 1.5^\circ$$
$$- P_{miss} < 1 \text{ GeV} \rightarrow P_{miss} < 0.8 \text{ GeV}.$$

BDT-Score cut:

– BDT-Score> $0.0 \rightarrow \text{BDT-Score} > 0.04$.



$$\sigma^{bkg} = \frac{A^{raw} - A^{\pi^0}}{(1-f)^2} \delta f$$

$$\sigma^{cut} = \frac{\frac{A_{cut}}{\sigma(A_{cut})} - \frac{A}{\sigma(A)}}{\frac{1}{\sigma(A_{cut})} + \frac{1}{\sigma(A)}}.$$

$$\sigma^{cut} = \sqrt{\frac{(A_+ - A_0)^2 + (A_- - A_0)^2}{2}}$$

$$A_{\pm} = \frac{\frac{A_{inb} \pm \sigma_{inb}^{cut}}{\sigma(A_{inb})^2} + \frac{A_{outb} \pm \sigma_{outb}^{cut}}{\sigma(A_{outb})^2}}{\frac{1}{\sigma(A_{inb})^2} + \frac{1}{\sigma(A_{outb})^2}}$$





4. DVCS MEASUREMENTS: CROSS-SECTION

If the eg selection allows BSA measurements

$$\frac{d^{4}\sigma}{dx_{B}dQ^{2}dtd\phi} = \frac{1}{F_{eff}} \times \left\{ \frac{1}{F_{Bin}} \times \frac{1}{F_{RC}} \times \frac{1}{\mathcal{L}\Delta\Omega} \right\}$$



Bin centering and RC corrections are significant.



