

Amplitude-Based Analyses at GlueX

Boris Grube*

on behalf of the  Collaboration

*Thomas Jefferson National Accelerator Facility

Workshop on AI for Hadron Spectroscopy at JLab

June 4, 2025

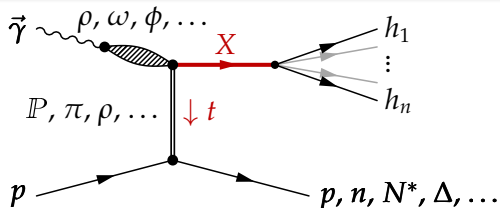
Jefferson Lab, Newport News, VA



U.S. DEPARTMENT OF
ENERGY

Office of
Science



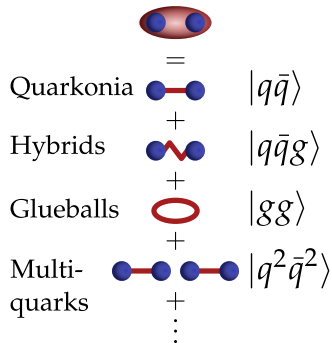
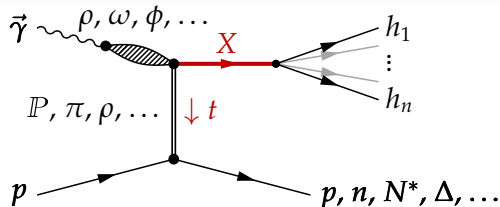


Rich process

- Light mesons with **all J^{PC} quantum numbers accessible** in various decay modes
- Produced by **various exchange processes**
- Linear beam polarization helps to **disentangle production mechanisms**
- Provides mostly **complementary information** to existing data

Goal: precision measurement of light-meson spectrum

- Confirm higher **excited conventional states**
- Complete **$SU(3)_{\text{flavor}}$ multiplets**
- Search for **exotic states beyond $q\bar{q}$**



Rich process

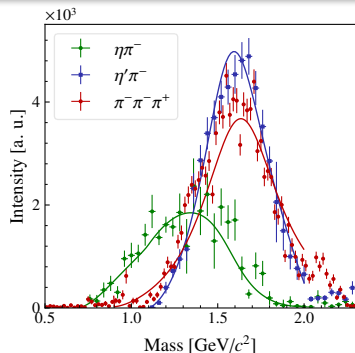
- Light mesons with all J^{PC} quantum numbers accessible in various decay modes
- Produced by various exchange processes
- Linear beam polarization helps disentangle production mechanisms
- Provides mostly complementary information to existing data

Goal: precision measurement of light-meson spectrum

- Confirm higher excited conventional states
- Complete $SU(3)_{\text{flavor}}$ multiplets
- Search for exotic states beyond $q\bar{q}$

Best exotic light-meson candidate: $\pi_1(1600)$

Spin-exotic $J^{PC} = 1^{-+}$ quantum numbers



- Best evidence in COMPASS $\eta\pi$, $\eta'\pi$, and $\rho(770)\pi$ data from pion diffraction

COMPASS, PLB **740** (2015) 303; PRD **105** (2022) 1012005; PRD **98** (2018) 092003; PRL **104** (2010) 241803

- JPAAC coupled-channel analysis: $\eta\pi$ and $\eta'\pi$ data can be described by single resonance pole consistent with $\pi_1(1600)$

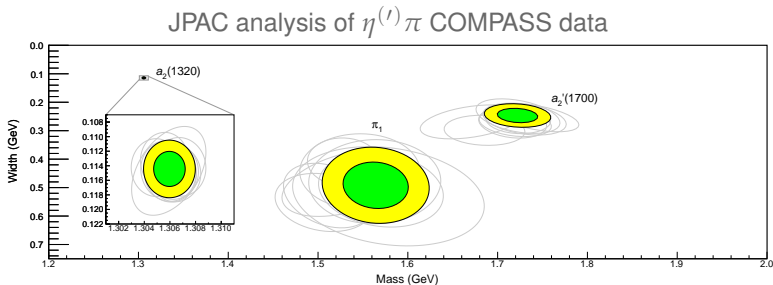
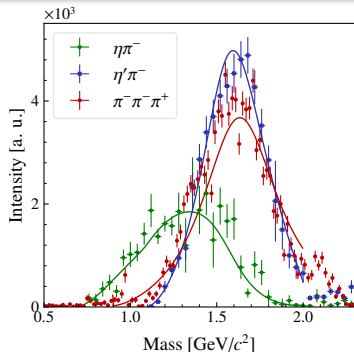
Rodas *et al.* [JPAC], PRL **122** (2019) 042002

- Recent find by BESIII: isoscalar partner $\eta_1(1855)$

BESIII, PRL **129** (2022) 192002

Best exotic light-meson candidate: $\pi_1(1600)$

Spin-exotic $J^{PC} = 1^{-+}$ quantum numbers



- Best evidence in COMPASS $\eta\pi$, $\eta'\pi$, and $\rho(770)\pi$ data from pion diffraction

COMPASS, PLB **740** (2015) 303; PRD **105** (2022) 1012005; PRD **98** (2018) 092003; PRL **104** (2010) 241803

- JPAC coupled-channel analysis: $\eta\pi$ and $\eta'\pi$ data can be described by single resonance pole consistent with $\pi_1(1600)$

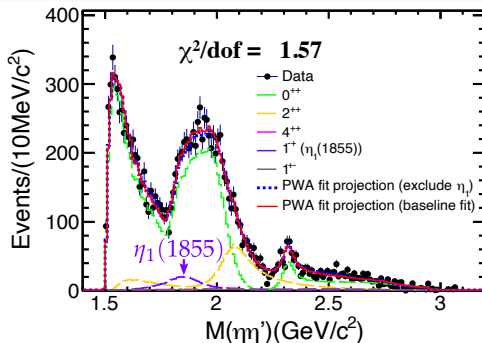
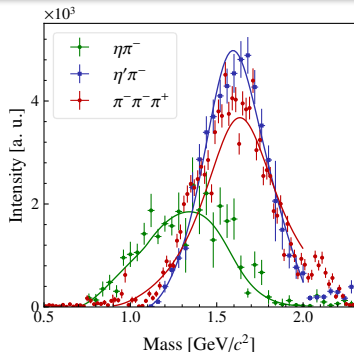
Rodas *et al.* [JPAC], PRL **122** (2019) 042002

- Recent find by BESIII: isoscalar partner $\eta_1(1855)$

BESIII, PRL **129** (2022) 192002

Best exotic light-meson candidate: $\pi_1(1600)$

Spin-exotic $J^{PC} = 1^{-+}$ quantum numbers



- Best evidence in COMPASS $\eta\pi$, $\eta'\pi$, and $\rho(770)\pi$ data from pion diffraction

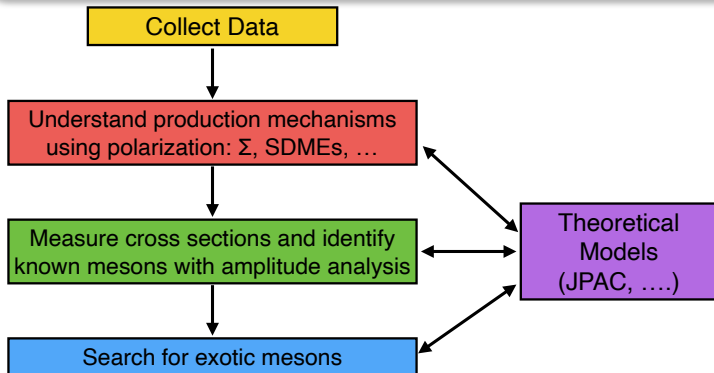
COMPASS, PLB **740** (2015) 303; PRD **105** (2022) 1012005; PRD **98** (2018) 092003; PRL **104** (2010) 241803

- JPAC coupled-channel analysis: $\eta\pi$ and $\eta'\pi$ data can be described by single resonance pole consistent with $\pi_1(1600)$

Rodas *et al.* [JPAC], PRL **122** (2019) 042002

- Recent find by BESIII: isoscalar partner $\eta_1(1855)$

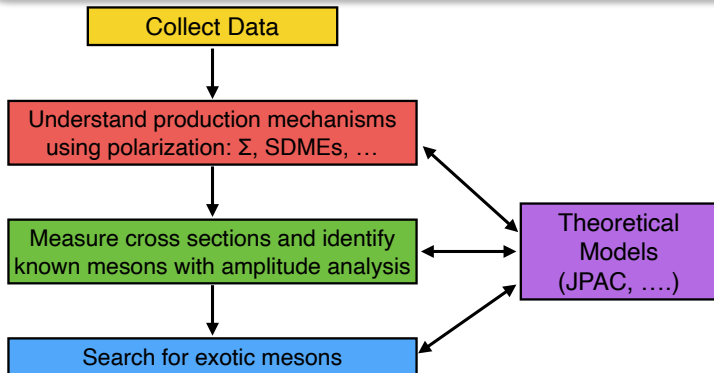
BESIII, PRL **129** (2022) 192002



- Amplitude analysis essential to extract meson spectrum
- *Strategy*: understand photoproduction of well-known states first and then use them as reference when searching for exotic states
- “Golden channels” for π_1 search: $\eta\pi$ and $\eta'\pi$

This talk: focus on common challenges encountered in amplitude analyses

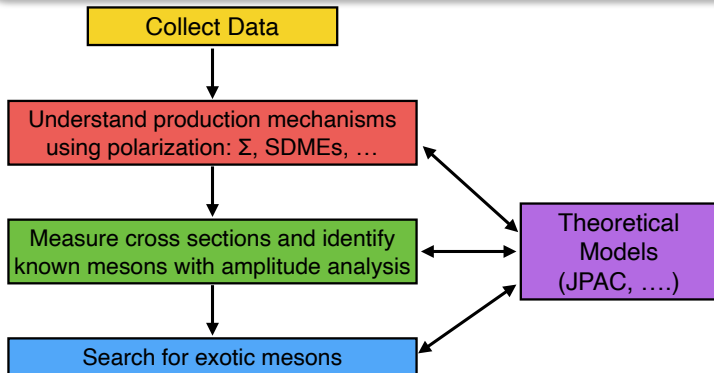
- Improving stability of fit results by imposing continuity constraints on amplitudes
- How to take into account non-resonant contributions?
- Moment analysis as an additional avenue



- **Amplitude analysis** essential to extract meson spectrum
- **Strategy: understand** photoproduction of **well-known states first** and then use them as **reference when searching for exotic states**
- “Golden channels” for π_1 search: **$\eta\pi$ and $\eta'\pi$**

This talk: focus on common challenges encountered in amplitude analyses

- Improving **stability of fit results** by imposing continuity constraints on amplitudes
- How to take into account **non-resonant contributions**?
- **Moment analysis** as an additional avenue



- Amplitude analysis essential to extract meson spectrum
- Strategy: understand photoproduction of well-known states first and then use them as reference when searching for exotic states
- “Golden channels” for π_1 search: $\eta\pi$ and $\eta'\pi$

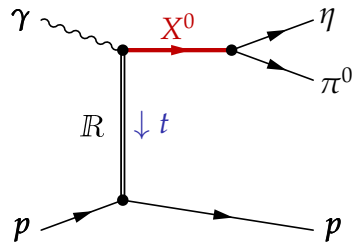
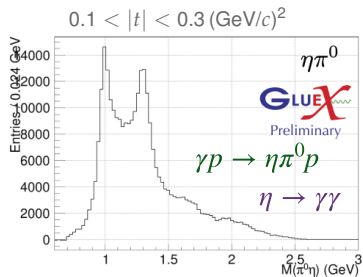
This talk: focus on common challenges encountered in amplitude analyses

- Improving stability of fit results by imposing continuity constraints on amplitudes
- How to take into account non-resonant contributions?
- Moment analysis as an additional avenue

Amplitude analysis of $\eta^{(\prime)}\pi$ at GlueX

Why $\eta^{(\prime)}\pi$?

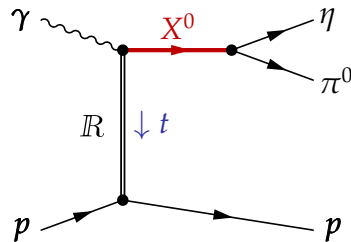
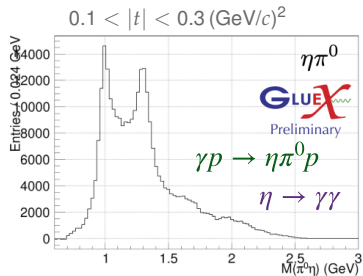
- Experimentally clean channels
- Two-body final states are easiest to model
- All waves with odd orbital angular momentum are spin-exotic, e.g. P -wave has $J^{PC} = 1^{-+}$
- Use well known $a_2(1320)$ as reference
- Multiple $\eta^{(\prime)}$ decay modes accessible, e.g. $\eta \rightarrow \gamma\gamma$ and $\pi^+\pi^-\pi^0$
 - Assess systematics from acceptance and backgrounds
- Test our understanding of production mechanisms:
 - Linear polarization of beam photons \Rightarrow separation of natural- and unnatural-parity exchange
 - Multiple production channels accessible, e.g.
 - $\gamma p \rightarrow \eta\pi^0 p$: mostly ρ and ω exchange
 - $\gamma p \rightarrow \eta\pi^- \Delta^{++}$: mostly π exchange



Amplitude analysis of $\eta^{(\prime)}\pi$ at GlueX

Why $\eta^{(\prime)}\pi$?

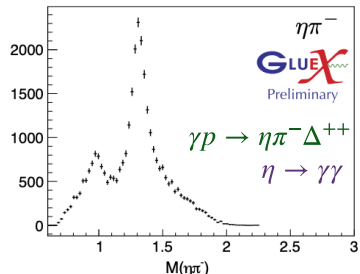
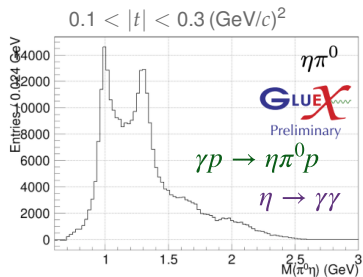
- Experimentally clean channels
- Two-body final states are easiest to model
- All waves with odd orbital angular momentum are spin-exotic, e.g. P -wave has $J^{PC} = 1^{-+}$
- Use well known $a_2(1320)$ as reference
- Multiple $\eta^{(\prime)}$ decay modes accessible, e.g. $\eta \rightarrow \gamma\gamma$ and $\pi^+\pi^-\pi^0$
 - Assess systematics from acceptance and backgrounds
- Test our understanding of production mechanisms:
 - Linear polarization of beam photons \Rightarrow separation of natural- and unnatural-parity exchange
 - Multiple production channels accessible, e.g.
 - $\gamma p \rightarrow \eta\pi^0 p$: mostly ρ and ω exchange
 - $\gamma p \rightarrow \eta\pi^- \Delta^{++}$: mostly π exchange



Amplitude analysis of $\eta^{(\prime)}\pi$ at GlueX

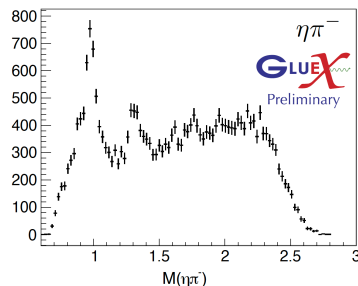
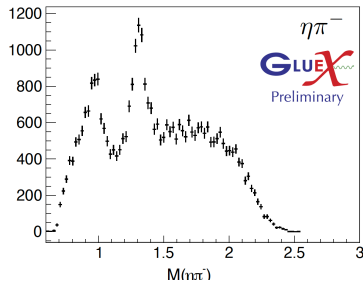
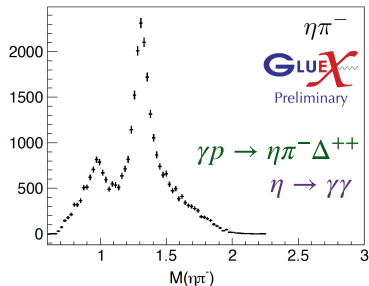
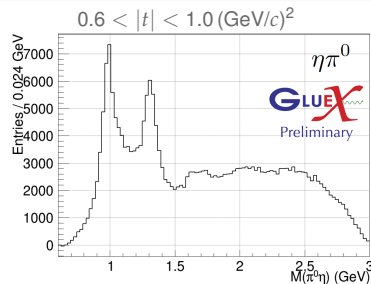
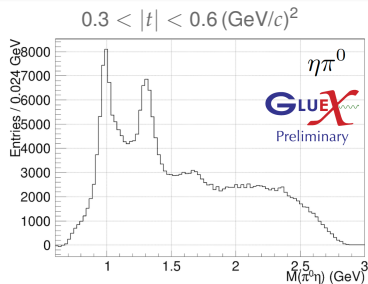
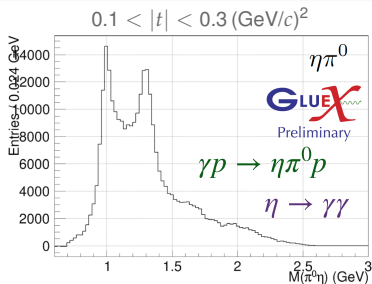
Why $\eta^{(\prime)}\pi$?

- Experimentally clean channels
- Two-body final states are easiest to model
- All waves with odd orbital angular momentum are spin-exotic, e.g. P -wave has $J^{PC} = 1^{-+}$
- Use well known $a_2(1320)$ as reference
- Multiple $\eta^{(\prime)}$ decay modes accessible, e.g. $\eta \rightarrow \gamma\gamma$ and $\pi^+\pi^-\pi^0$
 - Assess systematics from acceptance and backgrounds
- Test our understanding of production mechanisms:
 - Linear polarization of beam photons \Rightarrow separation of natural- and unnatural-parity exchange
 - Multiple production channels accessible, e.g.
 - $\gamma p \rightarrow \eta\pi^0 p$: mostly ρ and ω exchange
 - $\gamma p \rightarrow \eta\pi^- \Delta^{++}$: mostly π exchange



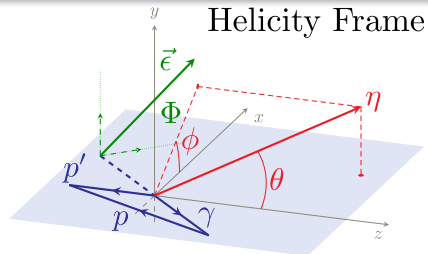
$\gamma p \rightarrow \eta \pi^0 p$ and $\gamma p \rightarrow \eta \pi^- \Delta^{++}; \eta \rightarrow \gamma \gamma$

Non-trivial dependence on 4-momentum transfer



3 Angles:

- $\Omega = (\theta, \phi)$ angles of η in $\eta\pi$ rest frame
- Φ is angle between γ polarization vector and production plane
- Linear beam polarization P_γ distinguishes between naturality (\pm) of exchange (= reflectivity)



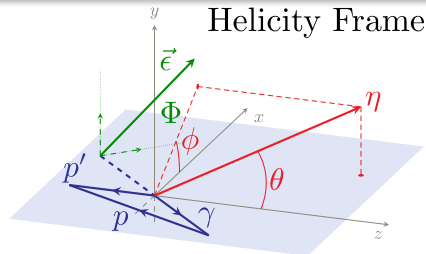
- Intensity model with angular amplitude $Z_\ell^m(\Omega, \Phi) \equiv Y_\ell^m(\Omega) e^{-i\Phi}$

$$\begin{aligned} \mathcal{I}(\Omega, \Phi) \propto & (1 - P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(-)} \operatorname{Re}[Z_\ell^m(\Omega, \Phi)] \right|^2 + (1 - P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(+)} \operatorname{Im}[Z_\ell^m(\Omega, \Phi)] \right|^2 \\ & + (1 + P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(+)} \operatorname{Re}[Z_\ell^m(\Omega, \Phi)] \right|^2 + (1 + P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(-)} \operatorname{Im}[Z_\ell^m(\Omega, \Phi)] \right|^2 \end{aligned}$$

- Wave set grows quickly with ℓ : $S_0^\pm, P_{-1,0,+1}^\pm, D_{-2,-1,0,+1,+2}^\pm, \dots$

• 3 Angles:

- $\Omega = (\theta, \phi)$ angles of η in $\eta\pi$ rest frame
- Φ is angle between γ polarization vector and production plane
- Linear beam polarization P_γ distinguishes between naturality (\pm) of exchange (= reflectivity)



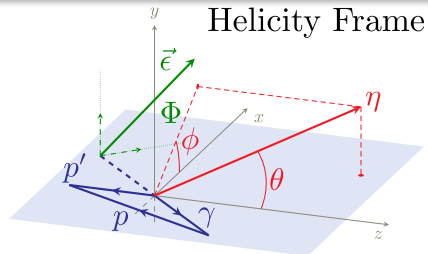
- Intensity model with angular amplitude $Z_\ell^m(\Omega, \Phi) \equiv Y_\ell^m(\Omega) e^{-i\Phi}$

$$\begin{aligned} \mathcal{I}(\Omega, \Phi) \propto & (1 - P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(-)} \operatorname{Re}[Z_\ell^m(\Omega, \Phi)] \right|^2 + (1 - P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(+)} \operatorname{Im}[Z_\ell^m(\Omega, \Phi)] \right|^2 \\ & + (1 + P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(+)} \operatorname{Re}[Z_\ell^m(\Omega, \Phi)] \right|^2 + (1 + P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(-)} \operatorname{Im}[Z_\ell^m(\Omega, \Phi)] \right|^2 \end{aligned}$$

- Wave set grows quickly with ℓ : $S_0^\pm, P_{-1,0,+1}^\pm, D_{-2,-1,0,+1,+2}^\pm, \dots$

3 Angles:

- $\Omega = (\theta, \phi)$ angles of η in $\eta\pi$ rest frame
- Φ is angle between γ polarization vector and production plane
- Linear beam polarization P_γ distinguishes between naturality (\pm) of exchange (= reflectivity)



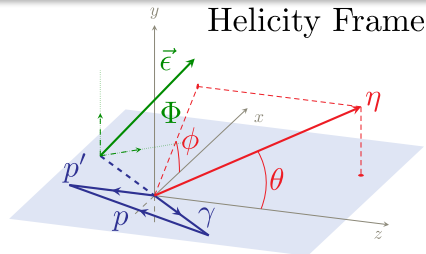
- Intensity model with angular amplitude $Z_\ell^m(\Omega, \Phi) \equiv Y_\ell^m(\Omega) e^{-i\Phi}$

$$\begin{aligned} \mathcal{I}(\Omega, \Phi) \propto & (1 - P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(-)} \operatorname{Re}[Z_\ell^m(\Omega, \Phi)] \right|^2 + (1 - P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(+)} \operatorname{Im}[Z_\ell^m(\Omega, \Phi)] \right|^2 \\ & + (1 + P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(+)} \operatorname{Re}[Z_\ell^m(\Omega, \Phi)] \right|^2 + (1 + P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(-)} \operatorname{Im}[Z_\ell^m(\Omega, \Phi)] \right|^2 \end{aligned}$$

- Wave set grows quickly with ℓ : $S_0^\pm, P_{-1,0,+1}^\pm, D_{-2,-1,0,+1,+2}^\pm, \dots$

3 Angles:

- $\Omega = (\theta, \phi)$ angles of η in $\eta\pi$ rest frame
- Φ is angle between γ polarization vector and production plane
- Linear beam polarization P_γ distinguishes between naturality (\pm) of exchange (= reflectivity)



- Intensity model with angular amplitude $Z_\ell^m(\Omega, \Phi) \equiv Y_\ell^m(\Omega) e^{-i\Phi}$

$$\begin{aligned} \mathcal{I}(\Omega, \Phi) \propto & (1 - P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(-)} \operatorname{Re}[Z_\ell^m(\Omega, \Phi)] \right|^2 + (1 - P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(+)} \operatorname{Im}[Z_\ell^m(\Omega, \Phi)] \right|^2 \\ & + (1 + P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(+)} \operatorname{Re}[Z_\ell^m(\Omega, \Phi)] \right|^2 + (1 + P_\gamma) \left| \sum_{\ell, m} [\ell]_m^{(-)} \operatorname{Im}[Z_\ell^m(\Omega, \Phi)] \right|^2 \end{aligned}$$

- Wave set grows quickly with ℓ : $S_0^\pm, P_{-1,0,+1}^\pm, D_{-2,-1,0,+1,+2}^\pm, \dots$

Mass-independent PWA of $\gamma p \rightarrow \eta \pi^0 p$; $\eta \rightarrow \gamma \gamma$

Low $|t|$: $0.1 < |t| < 0.3 \text{ (GeV/c)}^2$

- Waveset based on **tensor-meson dominance (TMD)**

model: $\{S_0^\pm, D_{-1}^-, D_0^\pm, D_{+1}^\pm, D_{+2}^+\}$

Mathieu *et al.* [JPAC], PRD **102** (2020) 014003

- Mostly **natural exchange**
- Sizable S_0^+ -**wave** contribution
 - $a_0(980)$ and $a_0(1450)$ signals?
 - Leakage from $a_2(1320)$?
 - **Non-resonant** contributions?
- Clear $a_2(1320)$ signal with $m = +2$ in **natural exchange** (= positive reflectivity)
- *Challenges:*
 - Model selection: what is the **optimal wave set**?
 - Large **bin-to-bin fluctuations** of partial-wave amplitudes of some waves

Mass-independent PWA of $\gamma p \rightarrow \eta \pi^0 p$; $\eta \rightarrow \gamma \gamma$

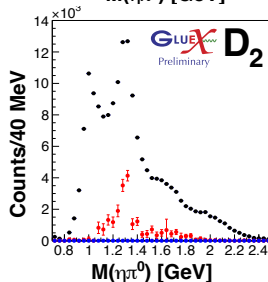
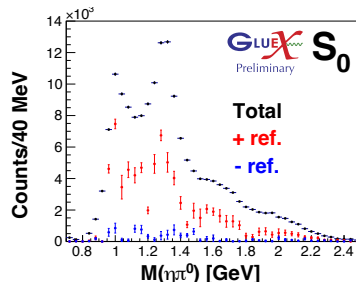
Low $|t|$: $0.1 < |t| < 0.3 \text{ (GeV/c)}^2$

- Waveset based on **tensor-meson dominance (TMD)**

model: $\{S_0^\pm, D_{-1}^-, D_0^\pm, D_{+1}^\pm, D_{+2}^\pm\}$

Mathieu *et al.* [JPAC], PRD **102** (2020) 014003

- Mostly **natural exchange**
- Sizable **S_0^+ -wave** contribution
 - $a_0(980)$ and $a_0(1450)$ signals?
 - Leakage from $a_2(1320)$?
 - Non-resonant** contributions?
- Clear $a_2(1320)$ signal with $m = +2$ in **natural exchange** (= positive reflectivity)
- Challenges:*
 - Model selection: what is the **optimal wave set**?
 - Large **bin-to-bin fluctuations** of partial-wave amplitudes of some waves



Mass-independent PWA of $\gamma p \rightarrow \eta \pi^0 p$; $\eta \rightarrow \gamma \gamma$

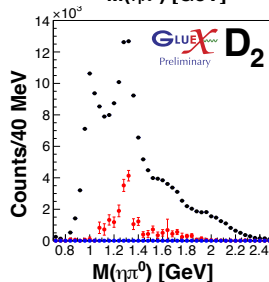
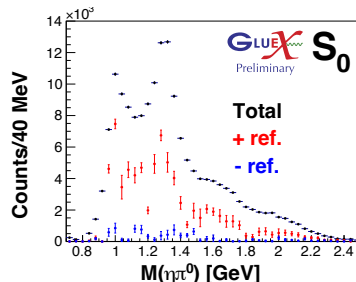
Low $|t|$: $0.1 < |t| < 0.3 \text{ (GeV/c)}^2$

- Waveset based on **tensor-meson dominance (TMD)**

model: $\{S_0^\pm, D_{-1}^-, D_0^\pm, D_{+1}^\pm, D_{+2}^\pm\}$

Mathieu *et al.* [JPAC], PRD **102** (2020) 014003

- Mostly **natural exchange**
- Sizable **S_0^+ -wave** contribution
 - $a_0(980)$ and $a_0(1450)$ signals?
 - Leakage from $a_2(1320)$?
 - Non-resonant** contributions?
- Clear $a_2(1320)$ signal with $m = +2$ in **natural exchange** (= positive reflectivity)
- Challenges:**
 - Model selection: what is the **optimal wave set**?
 - Large **bin-to-bin fluctuations** of partial-wave amplitudes of some waves



Ambiguities

- Spinless beam particle: mathematical ambiguities in form of Barrelet zeros
Chung, PRD **56** (1997) 7299
- Photoproduction:
 - No Barrelet zeros
Smith *et al.* [JPAC], PRD **108** (2023) 076001
 - Work in progress: find continuous mathematical ambiguities for special wave sets
 - Probably not relevant for most analyses

Local minima of negative log-likelihood function

- Input-output studies
 - Generate events from intensity distribution with known amplitudes and assuming 35 % beam polarization
 - Assume ideal case: fit generated events, i.e. no detector effects, with true wave set
 - Perform 100 PWA fit attempts with random initial values of amplitudes using MINUIT MIGRAD

Ambiguities

- Spinless beam particle: mathematical ambiguities in form of Barrelet zeros
Chung, PRD **56** (1997) 7299
- Photoproduction:
 - No Barrelet zeros
Smith *et al.* [JPAC], PRD **108** (2023) 076001
 - Work in progress: find continuous mathematical ambiguities for special wave sets
 - Probably not relevant for most analyses

Local minima of negative log-likelihood function

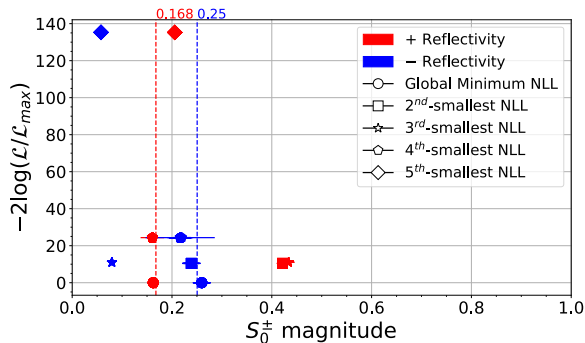
- Input-output studies
 - Generate events from intensity distribution with known amplitudes and assuming 35 % beam polarization
 - Assume ideal case: fit generated events, i.e. no detector effects, with true wave set
 - Perform 100 PWA fit attempts with random initial values of amplitudes using MINUIT MIGRAD

Interlude: two-pseudoscalar system

Local minima of negative log-likelihood function (NLL)

$\{S_0^\pm, P_{+1}^\pm, P_0^\pm, P_{-1}^\pm\}$ wave set, 10^6 events

Same wave set, 10^5 events



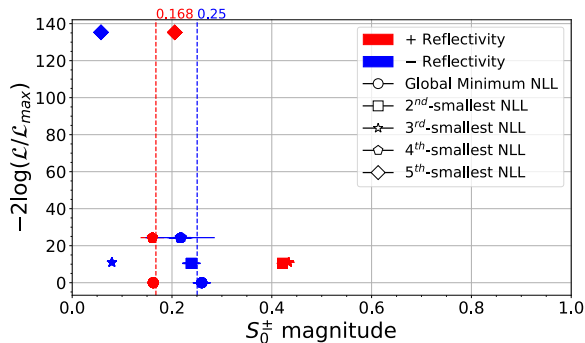
- **Global minimum:** estimated amplitudes consistent with input values
- **Additional local minima** close in likelihood
 - Amplitudes deviate significantly from input values

- Global minimum not consistent with input values
- **First local minimum:**
 - Only 0.5 units of log-likelihood worse
 - Consistent with input values

Interlude: two-pseudoscalar system

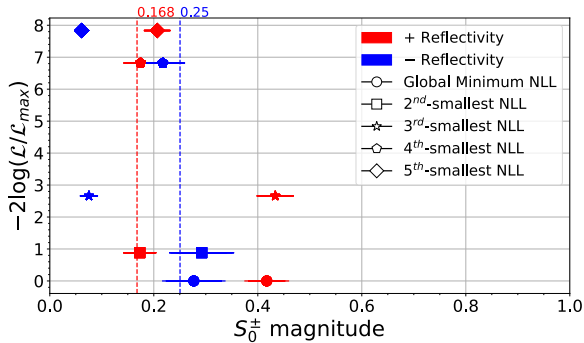
Local minima of negative log-likelihood function (NLL)

$\{S_0^\pm, P_{+1}^\pm, P_0^\pm, P_{-1}^\pm\}$ wave set, 10^6 events



- Global minimum: estimated amplitudes consistent with input values
- Additional local minima close in likelihood
 - Amplitudes deviate significantly from input values

Same wave set, 10^5 events



- Global minimum not consistent with input values
- First local minimum:
 - Only 0.5 units of log-likelihood worse
 - Consistent with input values

Interlude: two-pseudoscalar system

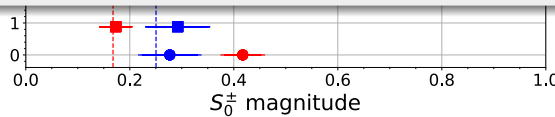
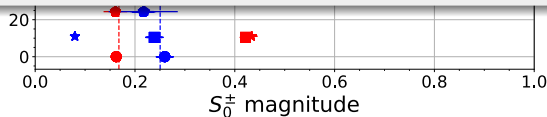
Local minima of negative log-likelihood function (NLL)

$\{S_0^\pm, P_{+1}^\pm, P_0^\pm, P_{-1}^\pm\}$ wave set, 10^6 events

Same wave set, 10^5 events



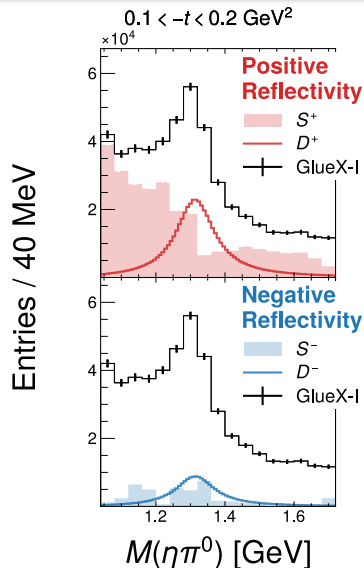
- Fits find **amplitudes far from true values** that describe data only slightly worse of even better than true values
- **More studies needed** to understand whether and how this affects real-data analyses
 - Could it (at least partly) explain observed large bin-to-bin fluctuations?



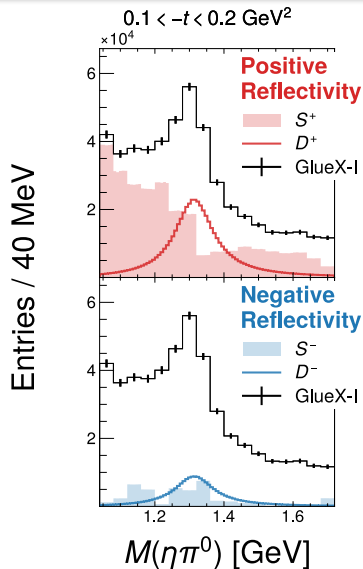
- **Global minimum:** estimated amplitudes **consistent** with input values
- **Additional local minima** close in likelihood
 - Amplitudes deviate significantly from input values

- **Global minimum not consistent** with input values
- **First local minimum:**
 - Only 0.5 units of log-likelihood worse
 - Consistent with input values

- **Constrain mass dependence** of selected amplitudes:
 - Model D -waves using sum of $a_2(1320)$ and $a_2(1700)$ **Breit-Wigner** amplitudes
- Keep mass-independent parametrization for S -waves
- **TMD wave set** $\{S_0^\pm, D_{-1}^\pm, D_0^\pm, D_{+1}^\pm, D_{+2}^\pm\}$: notable deviations from data at large $-t$
- Include all allowed S - and D -waves in fit, i.e. $\{S_0^\pm, D_{+2,+1,0,-1,-2}^\pm\}$
- For given naturality: **constrain production phases** of all m states of a given a_2 resonance to be identical



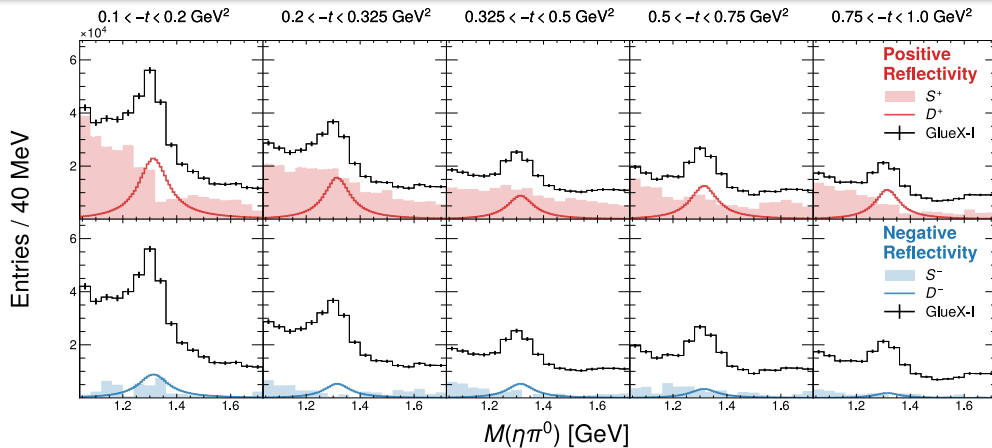
- **Constrain mass dependence** of selected amplitudes:
 - Model D -waves using sum of $a_2(1320)$ and $a_2(1700)$ **Breit-Wigner** amplitudes
- Keep mass-independent parametrization for S -waves
- **TMD wave set** $\{S_0^\pm, D_{-1}^-, D_0^\pm, D_{+1}^\pm, D_{+2}^\pm\}$: **notable deviations from data** at large $-t$
- Include **all allowed S - and D -waves** in fit, i.e. $\{S_0^\pm, D_{+2,+1,0,-1,-2}^\pm\}$
- For given naturality: **constrain production phases** of all m states of a given a_2 resonance to be identical



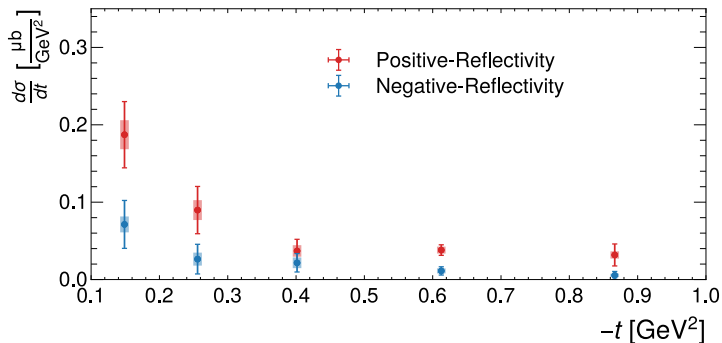
PWA of $\gamma p \rightarrow \eta \pi^0 p$; $\eta \rightarrow \gamma \gamma$

Hybrid PWA approach

arXiv:2501.03091

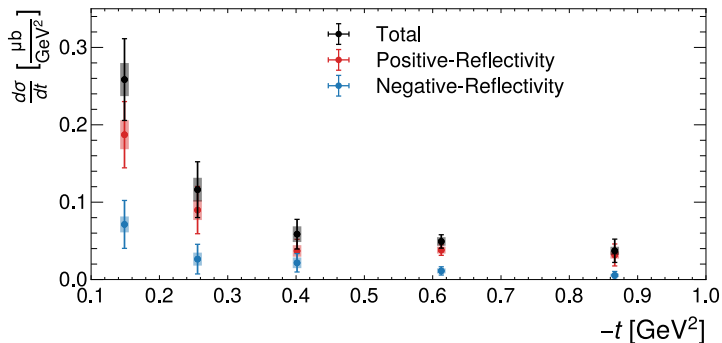


- Eliminates leakage between S - and D -waves
- Dominant contributions consistent with mass-independent PWA



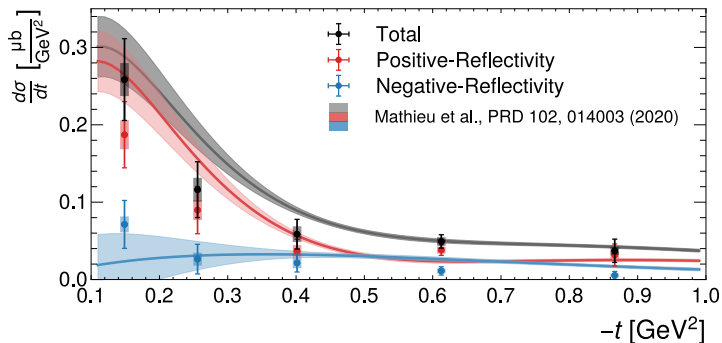
- First separation of contributions from **natural and unnatural exchanges**
 - Requires beam **polarization**; unique to GlueX
- Total cross section predicted by **TMD model** agrees well with data

Mathieu *et al.* [JPAC], PRD 102 (2020) 014003



- First separation of contributions from **natural and unnatural exchanges**
 - Requires beam **polarization**; unique to GlueX
- Total cross section predicted by **TMD model** agrees well with data

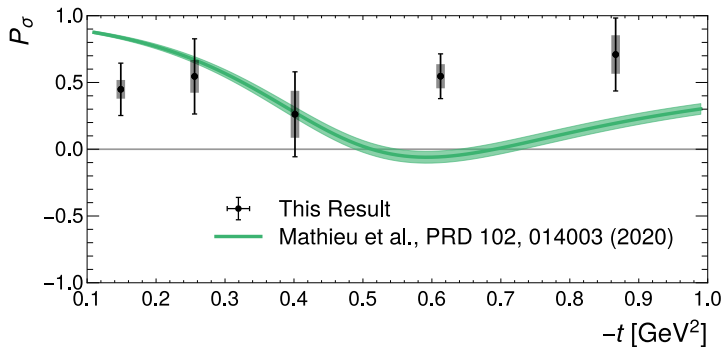
Mathieu *et al.* [JPAC], PRD 102 (2020) 014003



- First separation of contributions from **natural and unnatural exchanges**
 - Requires beam **polarization**; unique to GlueX
- Total cross section predicted by **TMD model** agrees well with data

Mathieu *et al.* [JPAC], PRD **102** (2020) 014003

$$P_\sigma = \frac{d\sigma^+/dt - d\sigma^-/dt}{d\sigma^+/dt + d\sigma^-/dt}$$



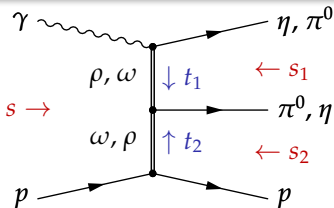
- We measure $P_\sigma \approx +0.5$, independent of t
- Significant deviation of **TMD model** from data for $-t \gtrsim 0.5 \text{ (GeV/c)}^2$

Mathieu *et al.* [JPAC], PRD **102** (2020) 014003

Challenge: double-Regge contributions

Example: $\gamma p \rightarrow \eta \pi^0 p$

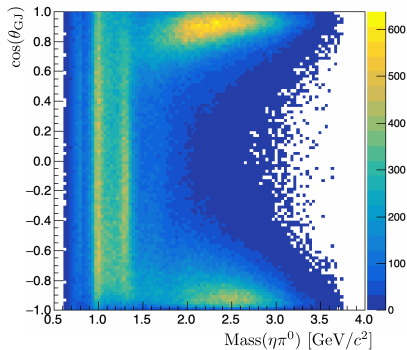
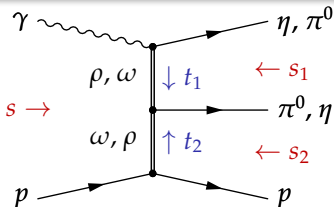
- Dominant at **large $\eta \pi^0$ mass and low $|t_2|$**
- Extend down to low-mass region and create **background for resonances**
- Important to **understand and model**
 - Can enhance spin-exotic odd- ℓ waves
 - Will **interfere with resonances**
 - Broad resonances such as $\pi_1(1600)$ may be masked, if not taken into account
- **Theory support indispensable**



Challenge: double-Regge contributions

Example: $\gamma p \rightarrow \eta \pi^0 p$

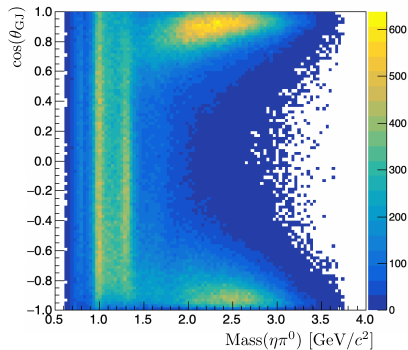
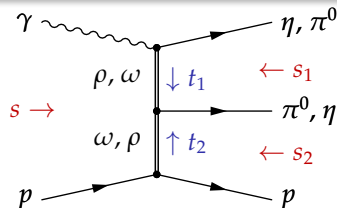
- Dominant at large $\eta\pi^0$ mass and low $|t_2|$
- Extend down to low-mass region and create background for resonances
- Important to understand and model
 - Can enhance spin-exotic odd- ℓ waves
 - Will interfere with resonances
 - Broad resonances such as $\pi_1(1600)$ may be masked, if not taken into account
- Theory support indispensable



Challenge: double-Regge contributions

Example: $\gamma p \rightarrow \eta \pi^0 p$

- Dominant at large $\eta \pi^0$ mass and low $|t_2|$
- Extend down to low-mass region and create background for resonances
- Important to understand and model
 - Can enhance spin-exotic odd- ℓ waves
 - Will interfere with resonances
 - Broad resonances such as $\pi_1(1600)$ may be masked, if not taken into account
- Theory support indispensable

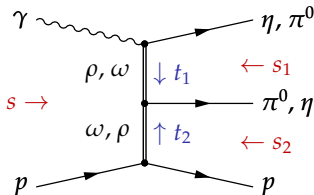


Modeling double-Regge processes

Example: $\gamma p \rightarrow \eta \pi^0 p$

- Close collaboration with JPAC

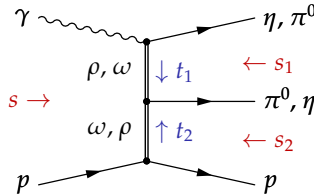
- Model describes **distribution in all phase-space variables**
- Based on Bibrzycki *et al.* [JPAC], EPJC 81 (2021) 647
- Improved description of vertex factors
- Fit regions with fast η and fast π^0 separately
- Reasonable agreement with data for $m_{\eta\pi} > 2 \text{ GeV}$
- Can we **extrapolate** model down to **resonance region**?



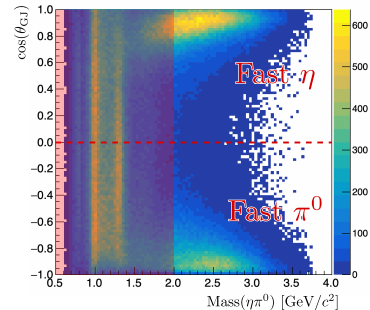
Modeling double-Regge processes

Example: $\gamma p \rightarrow \eta \pi^0 p$

- Close collaboration with JPAC
 - Model describes **distribution in all phase-space variables**
 - Based on Bibrzycki *et al.* [JPAC], EPJC 81 (2021) 647
 - Improved description of vertex factors
 - Fit regions with fast η and fast π^0 separately
 - Reasonable agreement with data for $m_{\eta\pi} > 2 \text{ GeV}$
- Can we **extrapolate** model down to **resonance region**?



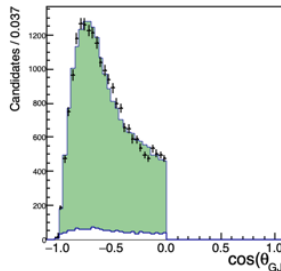
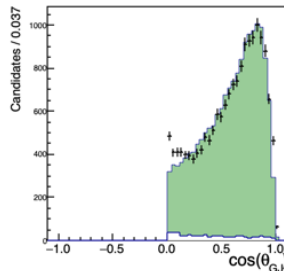
GLUEX
Internal



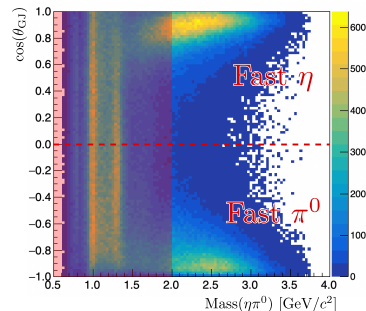
Modeling double-Regge processes

Example: $\gamma p \rightarrow \eta \pi^0 p$

- Close collaboration with JPAC
 - Model describes distribution in all phase-space variables
 - Based on Bibrzycki *et al.* [JPAC], EPJC 81 (2021) 647
 - Improved description of vertex factors
 - Fit regions with fast η and fast π^0 separately
 - Reasonable agreement with data for $m_{\eta\pi} > 2 \text{ GeV}$
- Can we extrapolate model down to resonance region?



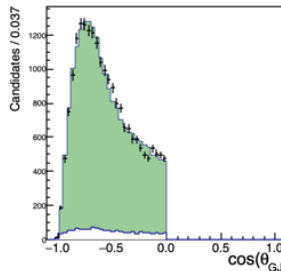
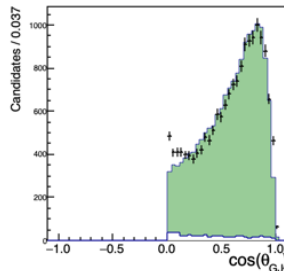
GLUEX
Internal



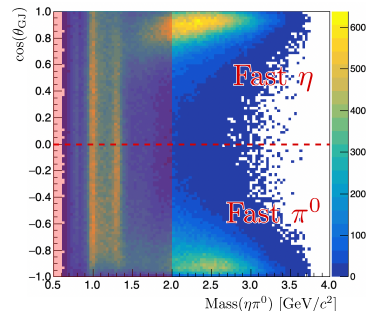
Modeling double-Regge processes

Example: $\gamma p \rightarrow \eta \pi^0 p$

- Close collaboration with JPAC
 - Model describes distribution in all phase-space variables
 - Based on Bibrzycki *et al.* [JPAC], EPJC 81 (2021) 647
 - Improved description of vertex factors
 - Fit regions with fast η and fast π^0 separately
 - Reasonable agreement with data for $m_{\eta\pi} > 2 \text{ GeV}$
- Can we extrapolate model down to resonance region?



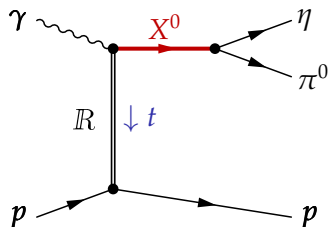
GLUEX
Internal



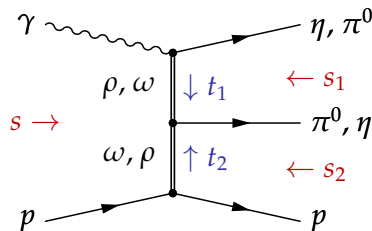
Additional challenge: baryon production

Example: $\gamma p \rightarrow \eta \pi^0 p$

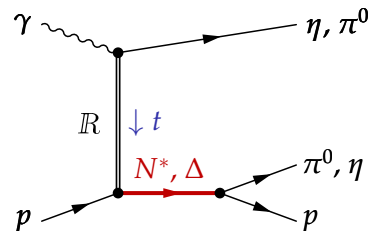
Meson production



Double-Regge



Baryon production



- Baryon production is **dominant at high $\eta\pi$ mass**
- Can be suppressed by **removing low $\pi^0 p$ masses for fast η and low ηp masses for fast π^0**
 - Cut strongly **distorts angular acceptance**
 - **Irreducible backgrounds** from higher baryon excitations remain

Moment decomposition of angular distribution in $\Omega = (\theta, \phi)$

$$\mathcal{I}(\Omega) = \sum_{LM} \sqrt{\frac{2L+1}{4\pi}} H(L, M) Y_L^M(\Omega)$$

with

$$H(L, M) = \sqrt{\frac{4\pi}{2L+1}} \int_{4\pi} d\Omega \mathcal{I}(\Omega) Y_L^{M*}(\Omega)$$

Advantages

- Moment decomposition is **unique**
- **Does not assume a model** \implies good way to pass experimental data to theorists

Disadvantages

- No direct access to **partial-wave amplitudes**
 - Would need to solve (overdetermined) second-order polynomial equation system

Moment decomposition of angular distribution in $\Omega = (\theta, \phi)$

$$\mathcal{I}(\Omega) = \sum_{LM} \sqrt{\frac{2L+1}{4\pi}} H(L, M) Y_L^M(\Omega)$$

with

$$H(L, M) = \sqrt{\frac{4\pi}{2L+1}} \int_{4\pi} d\Omega \mathcal{I}(\Omega) Y_L^{M*}(\Omega)$$

Advantages

- Moment decomposition is **unique**
- **Does not assume a model** \implies good way to pass experimental data to theorists

Disadvantages

- No direct access to **partial-wave amplitudes**
 - Would need to solve (overdetermined) second-order polynomial equation system

Moment decomposition of angular distribution in $\Omega = (\theta, \phi)$

$$\mathcal{I}(\Omega) = \sum_{LM}^{\infty} \sqrt{\frac{2L+1}{4\pi}} H(L, M) Y_L^M(\Omega)$$

with

$$H(L, M) = \sqrt{\frac{4\pi}{2L+1}} \int_{4\pi} d\Omega \mathcal{I}(\Omega) Y_L^{M*}(\Omega)$$

Advantages

- Moment decomposition is **unique**
- **Does not assume a model** \implies good way to pass experimental data to theorists

Disadvantages

- No direct access to **partial-wave amplitudes**
 - Would need to solve (overdetermined) second-order polynomial equation system

Moment analysis of two-body final states

Correcting for acceptance effects

- Measured intensity given by detector acceptance $\eta(\Omega)$: $\mathcal{I}_{\text{meas}}(\Omega) = \eta(\Omega) \mathcal{I}(\Omega)$

- Moment decomposition of $\mathcal{I}_{\text{meas}} \Rightarrow$ measured moments

$$H_{\text{meas}}(L, M) = \sqrt{\frac{4\pi}{2L+1}} \int_{4\pi} d\Omega \eta(\Omega) \mathcal{I}(\Omega) Y_L^{M*}(\Omega)$$

- Relate to physical moments $H(L', M')$ by inserting moment decomposition of physical intensity:

$$\mathcal{I}(\Omega) = \sum_{L' M'}^{\infty} \sqrt{\frac{2L'+1}{4\pi}} H(L', M') Y_{L'}^{M'}(\Omega)$$

- Acceptance mixes physical moments: $H_{\text{meas}}(L, M)$ are linear combinations of $H(L', M')$
- Matrix formulation: $\mathbf{H}_{\text{meas}} = \mathbf{I}_{\text{acc}} \mathbf{H} \Rightarrow$ physical moments: $\mathbf{H} = (\mathbf{I}_{\text{acc}})^{-1} \mathbf{H}_{\text{meas}}$
- Acceptance integral matrix \mathbf{I}_{acc} estimated using accepted phase-space events
- Estimate \mathbf{H}_{meas} from data by replacing integral by sum over N_{meas} measured events with weights $\{w_i\}$

$$H_{\text{meas}}(L, M) = \sqrt{\frac{4\pi}{2L+1}} \sum_{k=1}^{N_{\text{meas}}} w_k Y_L^{M*}(\Omega_k)$$

with $\sum_{k=1}^{N_{\text{meas}}} w_k =$ number of events after background subtraction

Moment analysis of two-body final states

Correcting for acceptance effects

- Measured intensity given by detector acceptance $\eta(\Omega)$: $\mathcal{I}_{\text{meas}}(\Omega) = \eta(\Omega) \mathcal{I}(\Omega)$
- Moment decomposition of $\mathcal{I}_{\text{meas}} \Rightarrow$ measured moments

$$H_{\text{meas}}(L, M) = \sqrt{\frac{4\pi}{2L+1}} \int_{4\pi} d\Omega \eta(\Omega) \mathcal{I}(\Omega) Y_L^{M*}(\Omega)$$

- Relate to physical moments $H(L', M')$ by inserting moment decomposition of physical intensity:

$$\mathcal{I}(\Omega) = \sum_{L' M'}^{\infty} \sqrt{\frac{2L'+1}{4\pi}} H(L', M') Y_{L'}^{M'}(\Omega)$$

- Acceptance mixes physical moments: $H_{\text{meas}}(L, M)$ are linear combinations of $H(L', M')$
- Matrix formulation: $\mathbf{H}_{\text{meas}} = \mathbf{I}_{\text{acc}} \mathbf{H} \Rightarrow$ physical moments: $\mathbf{H} = (\mathbf{I}_{\text{acc}})^{-1} \mathbf{H}_{\text{meas}}$
- Acceptance integral matrix \mathbf{I}_{acc} estimated using accepted phase-space events
- Estimate \mathbf{H}_{meas} from data by replacing integral by sum over N_{meas} measured events with weights $\{w_i\}$

$$H_{\text{meas}}(L, M) = \sqrt{\frac{4\pi}{2L+1}} \sum_{k=1}^{N_{\text{meas}}} w_k Y_L^{M*}(\Omega_k)$$

with $\sum_{k=1}^{N_{\text{meas}}} w_k =$ number of events after background subtraction

Moment analysis of two-body final states

Correcting for acceptance effects

- Measured intensity given by detector acceptance $\eta(\Omega)$: $\mathcal{I}_{\text{meas}}(\Omega) = \eta(\Omega) \mathcal{I}(\Omega)$
- Moment decomposition of $\mathcal{I}_{\text{meas}} \Rightarrow$ measured moments

$$H_{\text{meas}}(L, M) = \sqrt{\frac{4\pi}{2L+1}} \int_{4\pi} d\Omega \eta(\Omega) \mathcal{I}(\Omega) Y_L^{M*}(\Omega)$$

- Relate to physical moments $H(L', M')$ by inserting moment decomposition of physical intensity:

$$\mathcal{I}(\Omega) = \sum_{L' M'}^{\infty} \sqrt{\frac{2L'+1}{4\pi}} H(L', M') Y_{L'}^{M'}(\Omega)$$

- Acceptance mixes physical moments: $H_{\text{meas}}(L, M)$ are linear combinations of $H(L', M')$
- Matrix formulation: $\mathbf{H}_{\text{meas}} = \mathbf{I}_{\text{acc}} \mathbf{H} \Rightarrow$ physical moments: $\mathbf{H} = (\mathbf{I}_{\text{acc}})^{-1} \mathbf{H}_{\text{meas}}$
- Acceptance integral matrix \mathbf{I}_{acc} estimated using accepted phase-space events
- Estimate \mathbf{H}_{meas} from data by replacing integral by sum over N_{meas} measured events with weights $\{w_i\}$

$$H_{\text{meas}}(L, M) = \sqrt{\frac{4\pi}{2L+1}} \sum_{k=1}^{N_{\text{meas}}} w_k Y_L^{M*}(\Omega_k)$$

with $\sum_{k=1}^{N_{\text{meas}}} w_k =$ number of events after background subtraction

Moment analysis of two-body final states

Correcting for acceptance effects

- Measured intensity given by detector acceptance $\eta(\Omega)$: $\mathcal{I}_{\text{meas}}(\Omega) = \eta(\Omega) \mathcal{I}(\Omega)$
- Moment decomposition of $\mathcal{I}_{\text{meas}} \Rightarrow$ measured moments

$$H_{\text{meas}}(L, M) = \sqrt{\frac{4\pi}{2L+1}} \int_{4\pi} d\Omega \eta(\Omega) \mathcal{I}(\Omega) Y_L^{M*}(\Omega)$$

- Relate to physical moments $H(L', M')$ by inserting moment decomposition of physical intensity:

$$\mathcal{I}(\Omega) = \sum_{L' M'}^{\infty} \sqrt{\frac{2L'+1}{4\pi}} H(L', M') Y_{L'}^{M'}(\Omega)$$

- Acceptance mixes physical moments: $H_{\text{meas}}(L, M)$ are linear combinations of $H(L', M')$
- Matrix formulation: $\mathbf{H}_{\text{meas}} = \mathbf{I}_{\text{acc}} \mathbf{H} \Rightarrow$ physical moments: $\mathbf{H} = (\mathbf{I}_{\text{acc}})^{-1} \mathbf{H}_{\text{meas}}$
- Acceptance integral matrix \mathbf{I}_{acc} estimated using accepted phase-space events
- Estimate \mathbf{H}_{meas} from data by replacing integral by sum over N_{meas} measured events with weights $\{w_i\}$

$$H_{\text{meas}}(L, M) = \sqrt{\frac{4\pi}{2L+1}} \sum_{k=1}^{N_{\text{meas}}} w_k Y_L^{M*}(\Omega_k)$$

with $\sum_{k=1}^{N_{\text{meas}}} w_k =$ number of events after background subtraction

Moment analysis of two-body final states

Correcting for acceptance effects

- Measured intensity given by detector acceptance $\eta(\Omega)$: $\mathcal{I}_{\text{meas}}(\Omega) = \eta(\Omega) \mathcal{I}(\Omega)$
- Moment decomposition of $\mathcal{I}_{\text{meas}} \Rightarrow$ measured moments

$$H_{\text{meas}}(L, M) = \sqrt{\frac{4\pi}{2L+1}} \int_{4\pi} d\Omega \eta(\Omega) \mathcal{I}(\Omega) Y_L^{M*}(\Omega)$$

- Relate to physical moments $H(L', M')$ by inserting moment decomposition of physical intensity:

$$\mathcal{I}(\Omega) = \sum_{L' M'}^{\infty} \sqrt{\frac{2L'+1}{4\pi}} H(L', M') Y_{L'}^{M'}(\Omega)$$

- Acceptance mixes physical moments: $H_{\text{meas}}(L, M)$ are linear combinations of $H(L', M')$
- Matrix formulation: $\mathbf{H}_{\text{meas}} = \mathbf{I}_{\text{acc}} \mathbf{H} \Rightarrow$ physical moments: $\mathbf{H} = (\mathbf{I}_{\text{acc}})^{-1} \mathbf{H}_{\text{meas}}$
- Acceptance integral matrix \mathbf{I}_{acc} estimated using accepted phase-space events
- Estimate \mathbf{H}_{meas} from data by replacing integral by sum over N_{meas} measured events with weights $\{w_i\}$

$$H_{\text{meas}}(L, M) = \sqrt{\frac{4\pi}{2L+1}} \sum_{k=1}^{N_{\text{meas}}} w_k Y_L^{M*}(\Omega_k)$$

with $\sum_{k=1}^{N_{\text{meas}}} w_k =$ number of events after background subtraction

Moment analysis of two-body final states

Correcting for acceptance effects

- Measured intensity given by detector acceptance $\eta(\Omega)$: $\mathcal{I}_{\text{meas}}(\Omega) = \eta(\Omega) \mathcal{I}(\Omega)$
- Moment decomposition of $\mathcal{I}_{\text{meas}} \implies$ measured moments

$$H_{\text{meas}}(L, M) = \sqrt{\frac{4\pi}{2L+1}} \int_{4\pi} d\Omega \eta(\Omega) \mathcal{I}(\Omega) Y_L^{M*}(\Omega)$$

- Relate to physical moments $H(L', M')$ by inserting moment decomposition of physical intensity:

$$\mathcal{I}(\Omega) = \sum_{L' M'}^{\infty} \sqrt{\frac{2L'+1}{4\pi}} H(L', M') Y_{L'}^{M'}(\Omega)$$

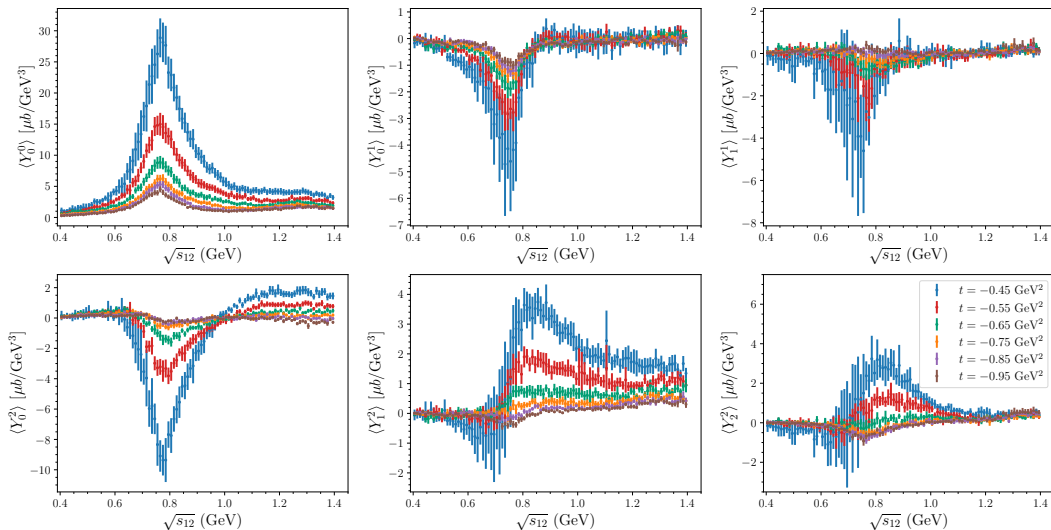
- Acceptance mixes physical moments: $H_{\text{meas}}(L, M)$ are linear combinations of $H(L', M')$
- Matrix formulation: $\mathbf{H}_{\text{meas}} = \mathbf{I}_{\text{acc}} \mathbf{H} \implies$ physical moments: $\mathbf{H} = (\mathbf{I}_{\text{acc}})^{-1} \mathbf{H}_{\text{meas}}$
- Acceptance integral matrix \mathbf{I}_{acc} estimated using accepted phase-space events
- Estimate \mathbf{H}_{meas} from data by replacing integral by sum over N_{meas} measured events with weights $\{w_i\}$

$$H_{\text{meas}}(L, M) = \sqrt{\frac{4\pi}{2L+1}} \sum_{k=1}^{N_{\text{meas}}} w_k Y_L^{M*}(\Omega_k)$$

with $\sum_{k=1}^{N_{\text{meas}}} w_k = \text{number of events after background subtraction}$

Moment analysis of unpolarized $\gamma p \rightarrow \pi^+ \pi^- p$

CLAS Results for $3.6 < E_\gamma < 3.8$ GeV



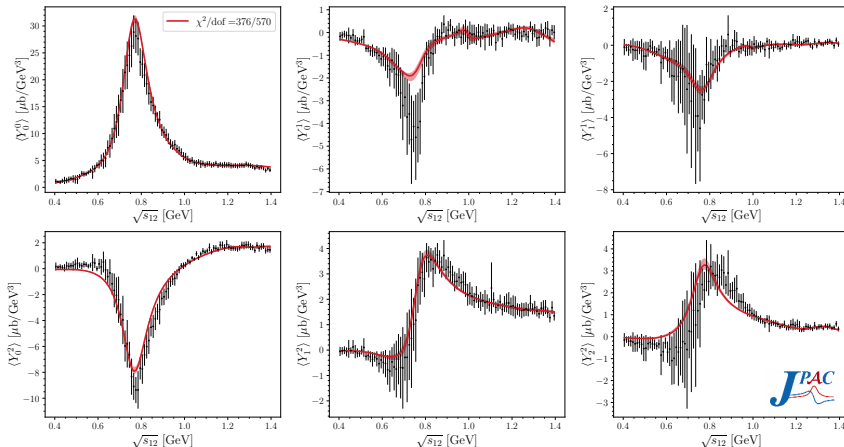
$$\sqrt{s_{12}} = m_{\pi\pi}$$

Battaglieri *et al.* [CLAS], PRD **80** (2009) 072005; Fig. from Bibrzycki *et al.* [JPAC], PRD **111** (2025) 014002

JPAC analysis of $\pi^+\pi^-$ moments from CLAS

Bibrzycki *et al.* [JPAC], PRD **111** (2025) 014002 CLAS, PRD **80** (2009) 072005

- JPAC: fit of **lowest 6 moments** in **highest energy bin** $3.6 < E_\gamma < 3.8$ GeV from CLAS analysis
- Model: **resonances** $f_0(500)$, $\rho(770)$, $f_0(980)$, $f_2(1270)$, $f_0(1370)$ + **Deck process** + **non-resonant polynomial** in $s_{12} = m_{\pi\pi}^2$
- Resonance production modeled using **Regge trajectories**: ρ and ω for $J^P = 0^+$ and 2^+ ; P , a_2 , and f_2 for $J^P = 1^-$



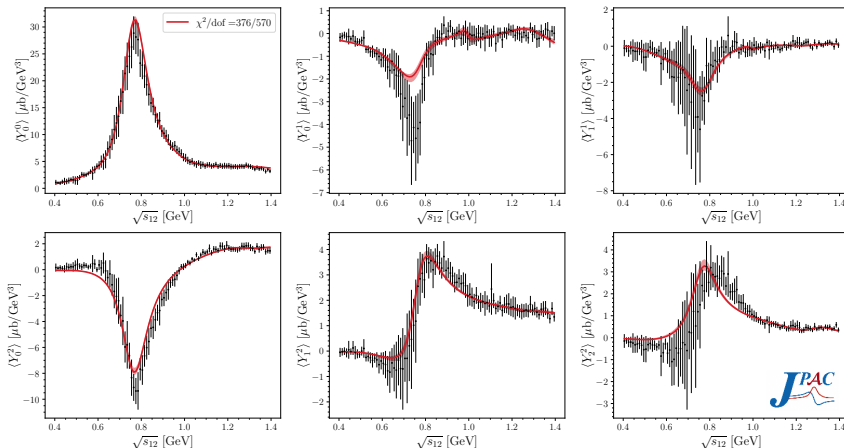
- Resonances described using **Breit-Wigner** amplitudes
- Resonance parameters fixed
- **30 free parameters**: relative strengths and phases of production mechanisms for each model component
- Each t bin fitted independently

$$0.4 < |t| < 0.5 \text{ GeV}^2$$

JPAC analysis of $\pi^+\pi^-$ moments from CLAS

Bibrzycki *et al.* [JPAC], PRD **111** (2025) 014002 CLAS, PRD **80** (2009) 072005

- JPAC: fit of **lowest 6 moments** in **highest energy bin** $3.6 < E_\gamma < 3.8$ GeV from CLAS analysis
- Model: **resonances** $f_0(500)$, $\rho(770)$, $f_0(980)$, $f_2(1270)$, $f_0(1370)$ + **Deck process** + **non-resonant polynomial** in $s_{12} = m_{\pi\pi}^2$
- Resonance production modeled using **Regge trajectories**: ρ and ω for $J^P = 0^+$ and 2^+ ; \mathbb{P} , a_2 , and f_2 for $J^P = 1^-$



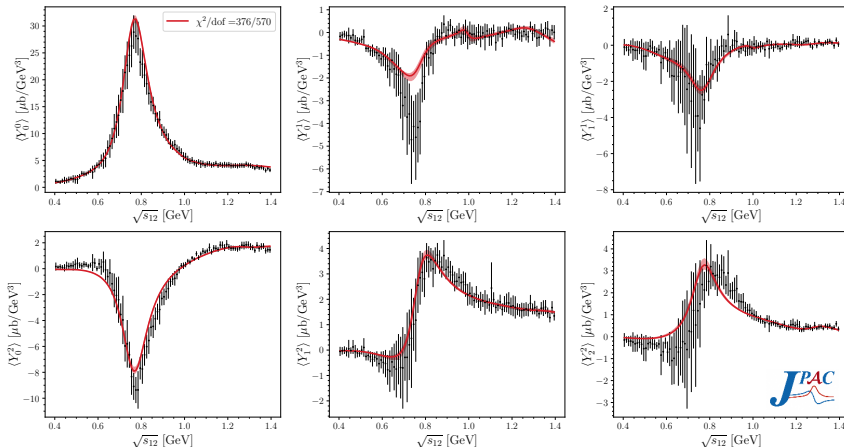
- Resonances described using **Breit-Wigner** amplitudes
- Resonance parameters fixed
- **30 free parameters**: relative strengths and phases of production mechanisms for each model component
- Each t bin fitted independently

$$0.4 < |t| < 0.5 \text{ GeV}^2$$

JPAC analysis of $\pi^+\pi^-$ moments from CLAS

Bibrzycki *et al.* [JPAC], PRD **111** (2025) 014002 CLAS, PRD **80** (2009) 072005

- JPAC: fit of **lowest 6 moments** in **highest energy bin** $3.6 < E_\gamma < 3.8$ GeV from CLAS analysis
- Model: **resonances** $f_0(500)$, $\rho(770)$, $f_0(980)$, $f_2(1270)$, $f_0(1370)$ + **Deck process** + **non-resonant polynomial** in $s_{12} = m_{\pi\pi}^2$
- Resonance production modeled using **Regge trajectories**: ρ and ω for $J^P = 0^+$ and 2^+ ; \mathbb{P} , a_2 , and f_2 for $J^P = 1^-$



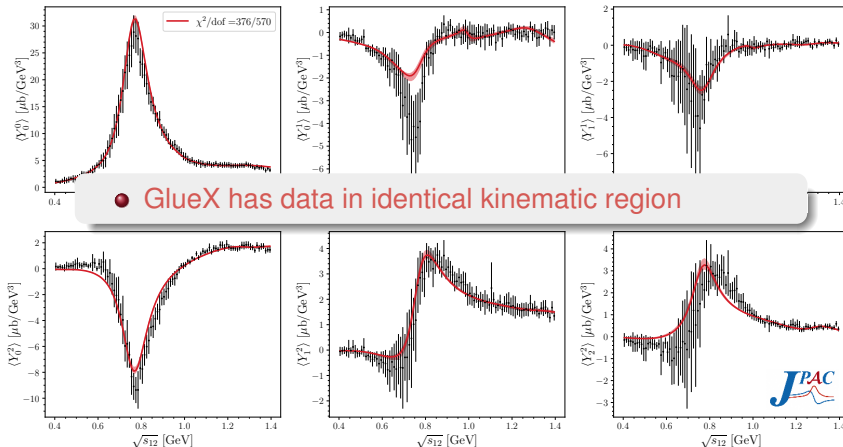
- Resonances described using **Breit-Wigner amplitudes**
- Resonance parameters fixed
- **30 free parameters**: relative strengths and phases of production mechanisms for each model component
- Each t bin fitted independently

$$0.4 < |t| < 0.5 \text{ GeV}^2$$

JPAC analysis of $\pi^+\pi^-$ moments from CLAS

Bibrzycki *et al.* [JPAC], PRD **111** (2025) 014002 CLAS, PRD **80** (2009) 072005

- JPAC: fit of **lowest 6 moments** in **highest energy bin** $3.6 < E_\gamma < 3.8$ GeV from CLAS analysis
- Model: **resonances** $f_0(500)$, $\rho(770)$, $f_0(980)$, $f_2(1270)$, $f_0(1370)$ + **Deck process** + **non-resonant polynomial** in $s_{12} = m_{\pi\pi}^2$
- Resonance production modeled using **Regge trajectories**: ρ and ω for $J^P = 0^+$ and 2^+ ; \mathbb{P} , a_2 , and f_2 for $J^P = 1^-$



- **GlueX has data in identical kinematic region**

- Resonances described using **Breit-Wigner amplitudes**
- Resonance parameters fixed
- **30 free parameters**: relative strengths and phases of production mechanisms for each model component
- Each t bin fitted independently

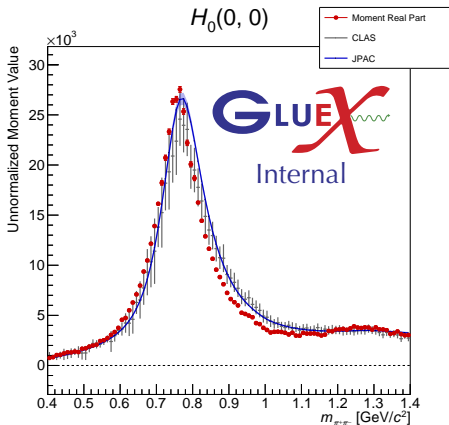
$$0.4 < |t| < 0.5 \text{ GeV}^2$$

- $$H(0,0) = \sqrt{4\pi} \int_{4\pi} d\Omega \mathcal{I}(\Omega) \underbrace{Y_0^{0*}(\Omega)}_{= 1/\sqrt{4\pi}} = \int_{4\pi} d\Omega \mathcal{I}(\Omega) = \text{number of acceptance-corrected events}$$

- All CLAS and JPAC moments scaled by same factor such that

$$\int dm_{\pi\pi} H_{\text{CLAS}}(0,0) = \int dm_{\pi\pi} H_{\text{GlueX}}(0,0)$$
- Error bars of CLAS points dominated by systematic uncertainties
- Only statistical uncertainties for GlueX
- Good qualitative agreement of GlueX results with CLAS/JPAC

$$\bullet \quad H(0,0) = \sqrt{4\pi} \int_{4\pi} d\Omega \mathcal{I}(\Omega) \underbrace{Y_0^{0*}(\Omega)}_{= 1/\sqrt{4\pi}} = \int_{4\pi} d\Omega \mathcal{I}(\Omega) = \text{number of acceptance-corrected events}$$



- All **CLAS** and **JPAC** moments scaled by same factor such that

$$\int dm_{\pi\pi} H_{\text{CLAS}}(0,0) = \int dm_{\pi\pi} H_{\text{GlueX}}(0,0)$$
- Error bars of **CLAS** points dominated by systematic uncertainties
- Only statistical uncertainties for **GlueX**
- Good qualitative agreement of **GlueX** results with **CLAS/JPAC**

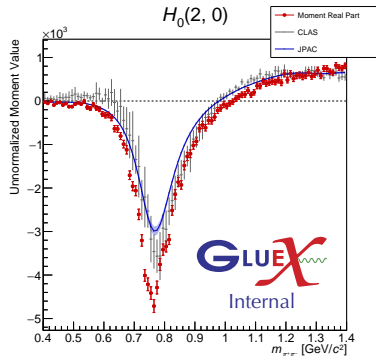
Moment analysis of unpolarized $\gamma p \rightarrow \pi^+ \pi^- p$

Example Results

$$H(2,0) = \dots -|P_{-1}|^2 + 2|P_0|^2 -|P_{+1}|^2 \dots$$

$$H(2,1) = \text{Re}[\dots -P_{-1}^* P_0 - P_{+1}^* P_0 \dots]$$

$$H(2,2) = \text{Re}[\dots -P_{-1}^* P_{+1} \dots]$$



- P_0 wave smaller than $P_{\pm 1}$ waves; P_0 and $P_{\pm 1}$ amplitudes not phase locked
- GlueX results agree qualitatively with CLAS/JPAC; deviations in details

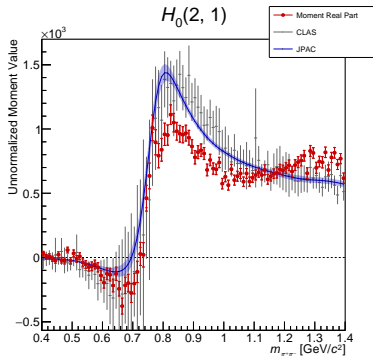
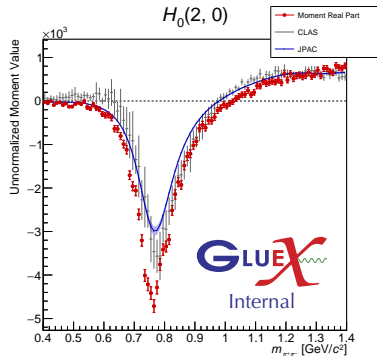
Moment analysis of unpolarized $\gamma p \rightarrow \pi^+ \pi^- p$

Example Results

$$H(2,0) = \dots -|P_{-1}|^2 + 2|P_0|^2 - |P_{+1}|^2 \dots$$

$$H(2,1) = \text{Re}[\dots - P_{-1}^* P_0 - P_{+1}^* P_0 \dots]$$

$$H(2,2) = \text{Re}[\dots - P_{-1}^* P_{+1} \dots]$$



- P_0 wave smaller than $P_{\pm 1}$ waves; P_0 and $P_{\pm 1}$ amplitudes not phase locked
- GlueX results agree qualitatively with CLAS/JPAC; deviations in details

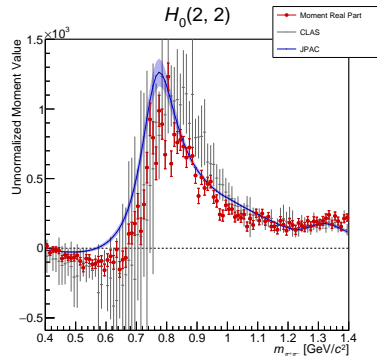
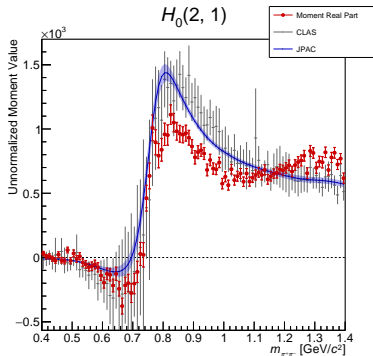
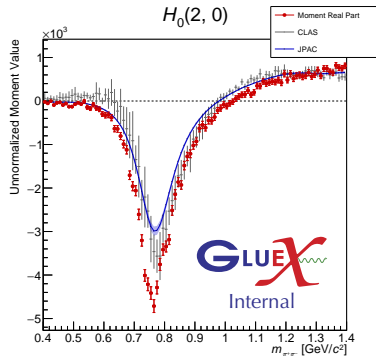
Moment analysis of unpolarized $\gamma p \rightarrow \pi^+ \pi^- p$

Example Results

$$H(2,0) = \dots -|P_{-1}|^2 + 2|P_0|^2 -|P_{+1}|^2 \dots$$

$$H(2,1) = \text{Re}[\dots - P_{-1}^* P_0 - P_{+1}^* P_0 \dots]$$

$$H(2,2) = \text{Re}[\dots - P_{-1}^* P_{+1} \dots]$$



- P_0 wave smaller than $P_{\pm 1}$ waves; P_0 and $P_{\pm 1}$ amplitudes not phase locked
- GlueX results agree qualitatively with CLAS/JPAC; deviations in details

Work in progress

- Use GlueX data to extend $\pi^+ \pi^-$ moment analysis to
 - Linearly polarized photon beams
 - Higher beam energies
 - Lower momentum transfers
 - Higher $\pi^+ \pi^-$ masses
 - Higher precision

Opportunity: moment analyses are good testing grounds for

- Surrogate models of detector acceptance
- Parameter estimation methods
 - Moment decomposition of intensity is unique
 - But moment values may lead to unphysical, i.e. negative intensities

Work in progress

- Use GlueX data to extend $\pi^+ \pi^-$ moment analysis to
 - Linearly polarized photon beams
 - Higher beam energies
 - Lower momentum transfers
 - Higher $\pi^+ \pi^-$ masses
 - Higher precision

Opportunity: moment analyses are good testing grounds for

- Surrogate models of detector acceptance
- Parameter estimation methods
 - Moment decomposition of intensity is unique
 - But moment values may lead to unphysical, i.e. negative intensities

Challenges in amplitude analyses provide opportunities for AI applications

- **Accelerating computations** by replacing expensive functions with **surrogate models**, e.g. for detector acceptance, would
 - allow **exploration of larger model spaces**
 - make advanced methods such as **Markov chain Monte Carlo** more feasible
 - allow more **detailed input-output studies**
 - ...
- **Model selection**
 - Determination of optimal wave sets
 - Determination of minimal resonance content
- Imposing **continuity and smoothness constraints** on amplitudes
- Alternative approaches for **parameter estimation**
- ...

Challenges in amplitude analyses provide opportunities for AI applications

- **Accelerating computations** by replacing expensive functions with **surrogate models**, e.g. for detector acceptance, would
 - allow **exploration of larger model spaces**
 - make advanced methods such as **Markov chain Monte Carlo** more feasible
 - allow more **detailed input-output studies**
 - ...
- **Model selection**
 - Determination of optimal wave sets
 - Determination of minimal resonance content
- Imposing **continuity and smoothness constraints** on amplitudes
- Alternative approaches for **parameter estimation**
- ...

Challenges in amplitude analyses provide opportunities for AI applications

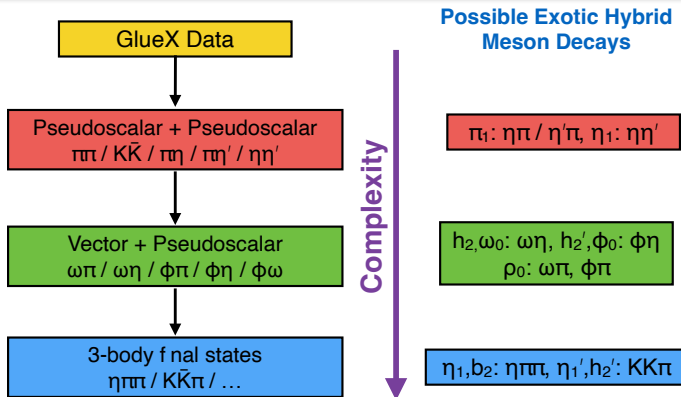
- **Accelerating computations** by replacing expensive functions with **surrogate models**, e.g. for detector acceptance, would
 - allow **exploration of larger model spaces**
 - make advanced methods such as **Markov chain Monte Carlo** more feasible
 - allow more **detailed input-output studies**
 - ...
- **Model selection**
 - Determination of optimal wave sets
 - Determination of minimal resonance content
- Imposing **continuity and smoothness constraints** on amplitudes
- Alternative approaches for **parameter estimation**
- ...

Part II

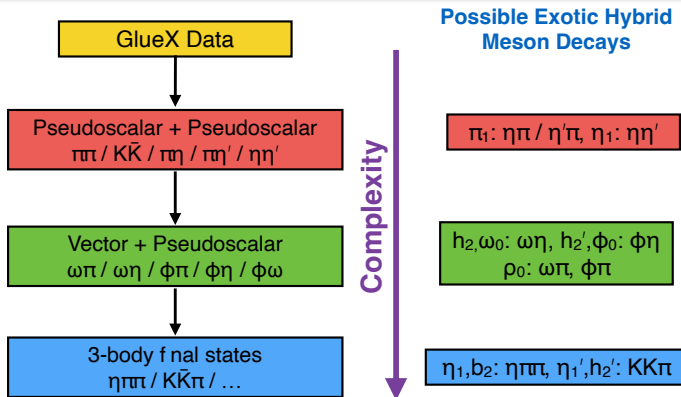
Backup Slides

1 Introduction

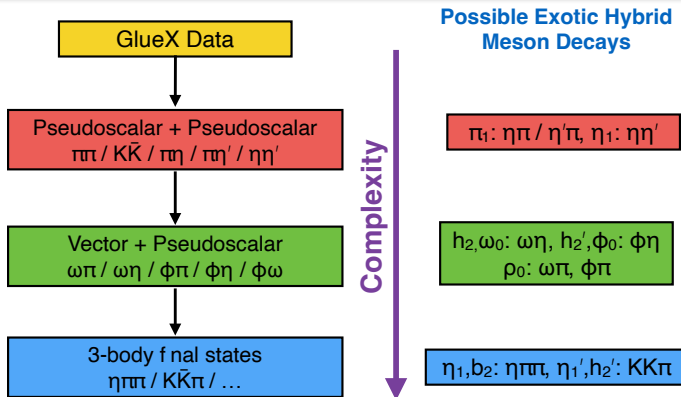
2 The GlueX experiment



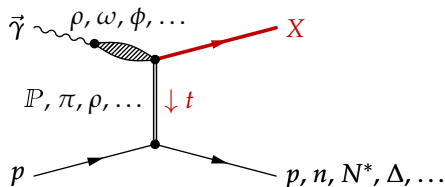
- *Strategy:* understand photoproduction of well-known states first and then use them as reference when searching for exotic states
- "Golden channels" for π_1 search: $\eta\pi$ and $\eta'\pi \Rightarrow$ focus of this talk



- *Strategy:* understand photoproduction of well-known states first and then use them as reference when searching for exotic states
- “Golden channels” for π_1 search: $\eta\pi$ and $\eta'\pi \implies$ focus of this talk



- *Strategy:* understand photoproduction of well-known states first and then use them as reference when searching for exotic states
- “Golden channels” for π_1 search: $\eta\pi$ and $\eta'\pi \implies$ focus of this talk

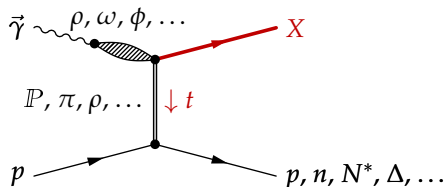


Exchange	Spin-exotic states		
P	0^{++}	b, h, h'	$0^{+-}, 2^{+-}$
π^0	0^{-+}	b_2, h_2, h'_2	2^{+-}
π^\pm	0^{-+}	π_1	1^{-+}
ω	1^{--}	π_1, η_1, η'_1	1^{-+}

- Wide variety of intermediate states X accessible
- Photon polarization provides constraints on production process
- Prerequisite: understanding of production mechanism
- Existing photoproduction data very limited

Goal of GlueX

- Confirm π_1 and η_1
- Establish full light-quark hybrid spectrum

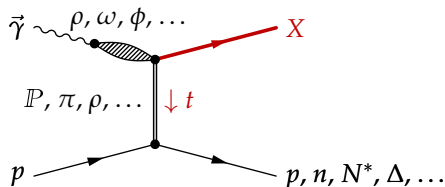


Exchange	Spin-exotic states		
P	0^{++}	b, h, h'	$0^{+-}, 2^{+-}$
π^0	0^{-+}	b_2, h_2, h'_2	2^{+-}
π^\pm	0^{-+}	π_1	1^{-+}
ω	1^{--}	π_1, η_1, η'_1	1^{-+}

- Wide variety of intermediate states X accessible
- Photon polarization provides constraints on production process
- *Prerequisite:* understanding of production mechanism
- Existing photoproduction data very limited

Goal of GlueX

- Confirm π_1 and η_1
- Establish full light-quark hybrid spectrum



Exchange	Spin-exotic states		
P	0^{++}	b, h, h'	$0^{+-}, 2^{+-}$
π^0	0^{-+}	b_2, h_2, h'_2	2^{+-}
π^\pm	0^{-+}	π_1	1^{-+}
ω	1^{--}	π_1, η_1, η'_1	1^{-+}

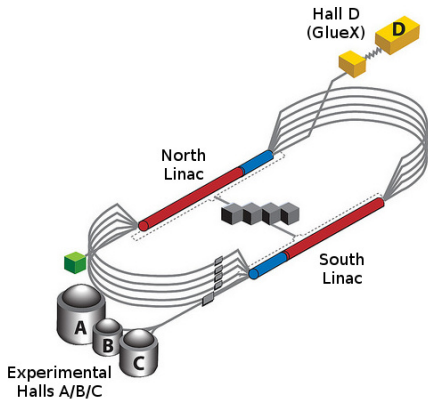
- Wide variety of intermediate states X accessible
- Photon polarization provides constraints on production process
- *Prerequisite:* understanding of production mechanism
- Existing photoproduction data very limited

Goal of GlueX

- Confirm π_1 and η_1
- Establish full light-quark hybrid spectrum

The GlueX Experiment at Jefferson Lab

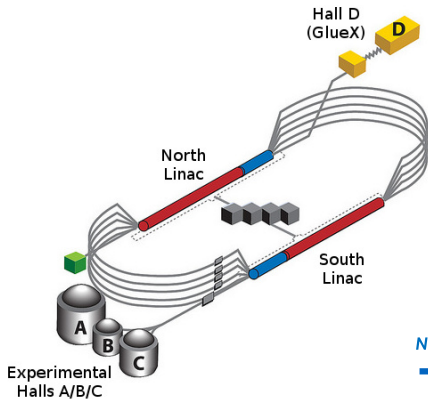
Adhikari *et al.*, NIMA **987** (2021) 164807



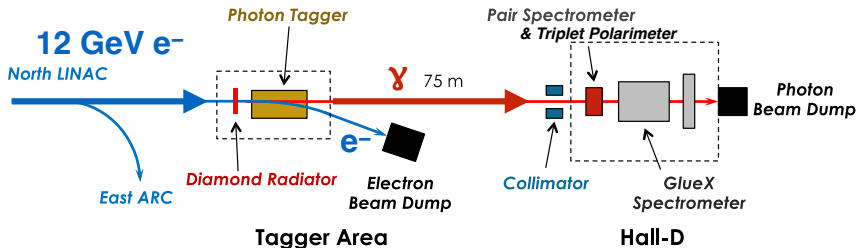
- 12 GeV e^- beam from CEBAF
- Diamond radiator \Rightarrow linearly polarized photon beam
- Photon energy tagged by scattered e^- ; $\sigma(E_\gamma) = 0.2\%$
- Coherent peak: 1 to 5×10^7 γ/s with $\approx 40\%$ polarization

The GlueX Experiment at Jefferson Lab

Adhikari *et al.*, NIMA **987** (2021) 164807

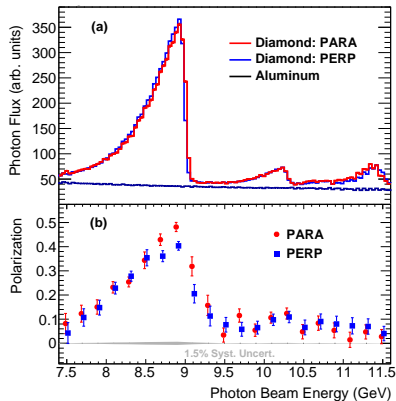


- 12 GeV e^- beam from CEBAF
- Diamond radiator \Rightarrow linearly polarized photon beam
- Photon energy tagged by scattered e^- ; $\sigma(E_\gamma) = 0.2\%$
- Coherent peak: 1 to 5×10^7 γ/s with $\approx 40\%$ polarization

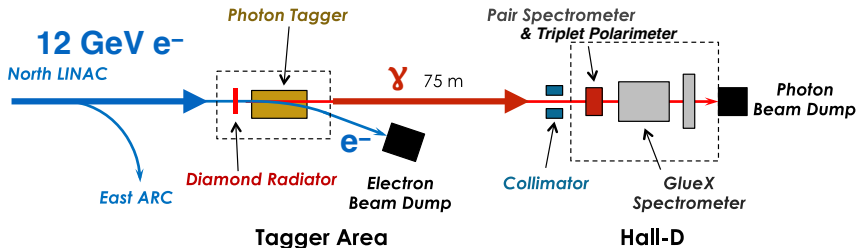


The GlueX Experiment at Jefferson Lab

Adhikari *et al.*, NIMA **987** (2021) 164807



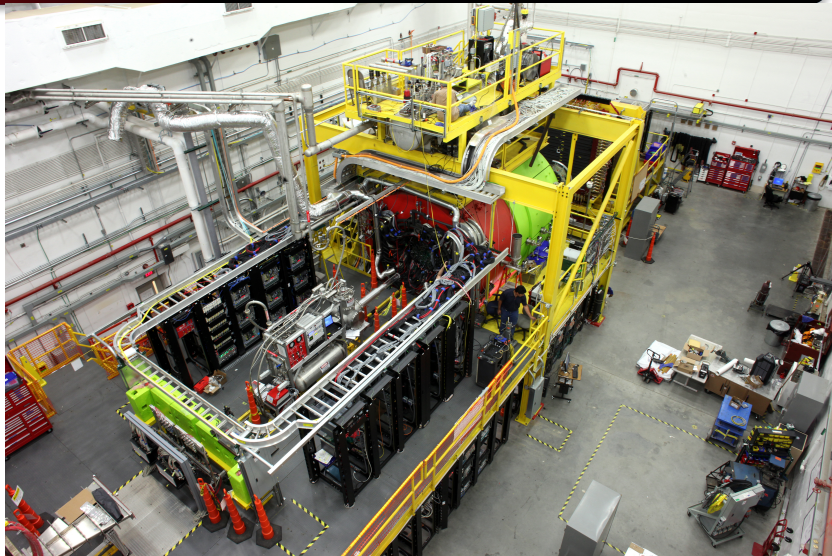
- 12 GeV e^- beam from CEBAF
- Diamond radiator \Rightarrow linearly polarized photon beam
- Photon energy tagged by scattered e^- ; $\sigma(E_\gamma) = 0.2\%$
- Coherent peak: 1 to 5×10^7 γ /s with $\approx 40\%$ polarization



The GlueX Experiment at Jefferson Lab

Adhikari *et al.*, NIMA **987** (2021) 164807

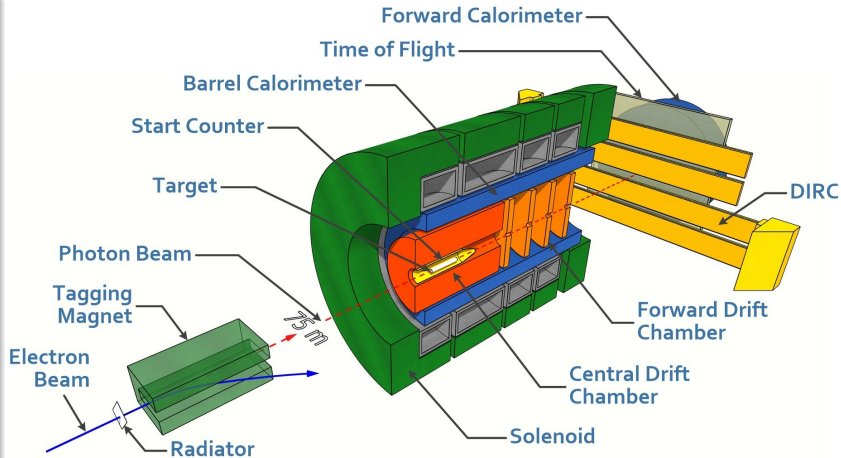
- Large-acceptance spectrometer
- Optimized for light-meson spectroscopy
- GlueX Phase I
 - 2017 to 2018
 - $\int \mathcal{L} = 125 \text{ pb}^{-1}$ in coherent peak
- GlueX Phase II
 - 2020 to 2025?
 - Added DIRC for PID
 - So far $\int \mathcal{L} = 190 \text{ pb}^{-1}$ in coherent peak
 - Expect 3 to 4 \times Phase I



The GlueX Experiment at Jefferson Lab

Adhikari *et al.*, NIMA **987** (2021) 164807

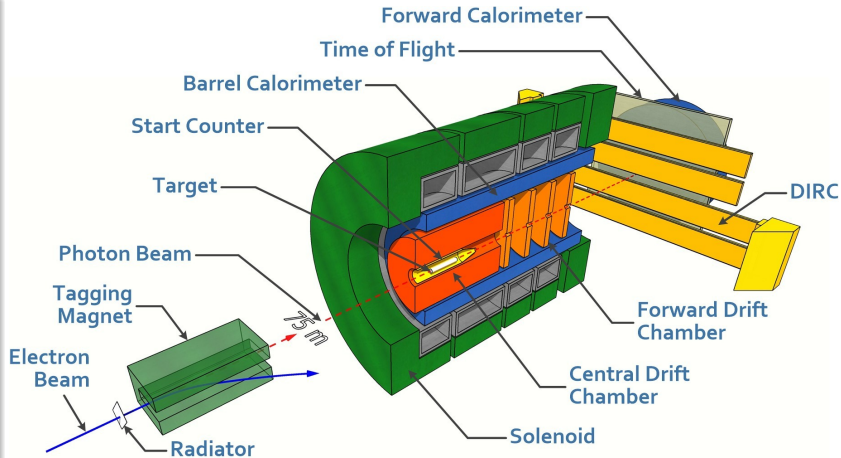
- Large-acceptance spectrometer
- Optimized for light-meson spectroscopy
- GlueX Phase I
 - 2017 to 2018
 - $\int \mathcal{L} = 125 \text{ pb}^{-1}$ in coherent peak
- GlueX Phase II
 - 2020 to 2025?
 - Added DIRC for PID
 - So far $\int \mathcal{L} = 190 \text{ pb}^{-1}$ in coherent peak
 - Expect 3 to 4 \times Phase I



The GlueX Experiment at Jefferson Lab

Adhikari *et al.*, NIMA **987** (2021) 164807

- Large-acceptance spectrometer
- Optimized for light-meson spectroscopy
- GlueX Phase I
 - 2017 to 2018
 - $\int \mathcal{L} = 125 \text{ pb}^{-1}$ in coherent peak
- GlueX Phase II
 - 2020 to 2025?
 - Added DIRC for PID
 - So far $\int \mathcal{L} = 190 \text{ pb}^{-1}$ in coherent peak
 - Expect 3 to 4 \times Phase I



The GlueX Experiment at Jefferson Lab

Adhikari *et al.*, NIMA **987** (2021) 164807

- Large-acceptance spectrometer
- Optimized for light-meson spectroscopy
- GlueX Phase I
 - 2017 to 2018
 - $\int \mathcal{L} = 125 \text{ pb}^{-1}$ in coherent peak
- GlueX Phase II
 - 2020 to 2025?
 - Added DIRC for PID
 - So far $\int \mathcal{L} = 190 \text{ pb}^{-1}$ in coherent peak
 - Expect 3 to 4 \times Phase I

