

Lattice Calculations of Nucleon Electromagnetic Form Factors at High Momentum Q^2

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(Lattice Hadron Physics collaboration)

JLab Hall A/C Collaboration meeting



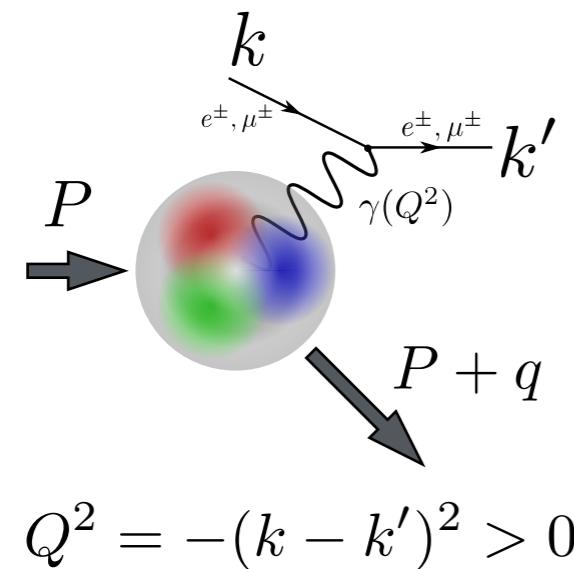
National Energy Research
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Outline

- Large- Q^2 Nucleon vector form factors on a lattice
 - *Motivation*
 - *Methodology*
 - *Challenges*
- Dominant (Quark-Connected) contributions
 - *Status*
 - *Examining systematic effects : excited states, discretization, pion-mass*
- Other systematics: Quark-Disconnected contributions
 - *Estimating impact on FFs and ratios*
- Summary & Outlook

Nucleon Elastic E&M Form Factors



Elastic e^-p amplitude

$$\langle P + q | \bar{q} \gamma^\mu q | P \rangle = \bar{U}_{P+q} \left[\begin{array}{l} \text{(Dirac)} \\ F_1(Q^2) \gamma^\mu + F_2(Q^2) \frac{i\sigma^{\mu\nu}q_\nu}{2M_N} \end{array} \right] U_P$$

Sachs Electric

$$G_E(Q^2) = F_1(Q^2) - \frac{Q^2}{4M^2} F_2(Q^2)$$

Magnetic

$$G_M(Q^2) = F_1(Q^2) + F_2(Q^2)$$

Elastic e^-p cross-section

- ➊ $G_{E,M}$ from ϵ -dep. at fixed $\tau(Q^2)$
("Rosenbluth separation")
- ➋ dominated by G_M at large Q^2
- ➌ 2γ corrections at $Q^2 \gtrsim 1 \text{ GeV}^2$

$$\frac{d\sigma}{d\Omega} = \frac{\sigma_{\text{Mott}}}{1 + \tau} \left[G_E^2 + \frac{\tau}{\epsilon} G_M^2 \right]$$

$$\tau = \frac{Q^2}{4M_N^2} \quad \epsilon = \left[1 + 2(1 + \tau) \tan^2 \frac{\theta}{2} \right]^{-1}$$

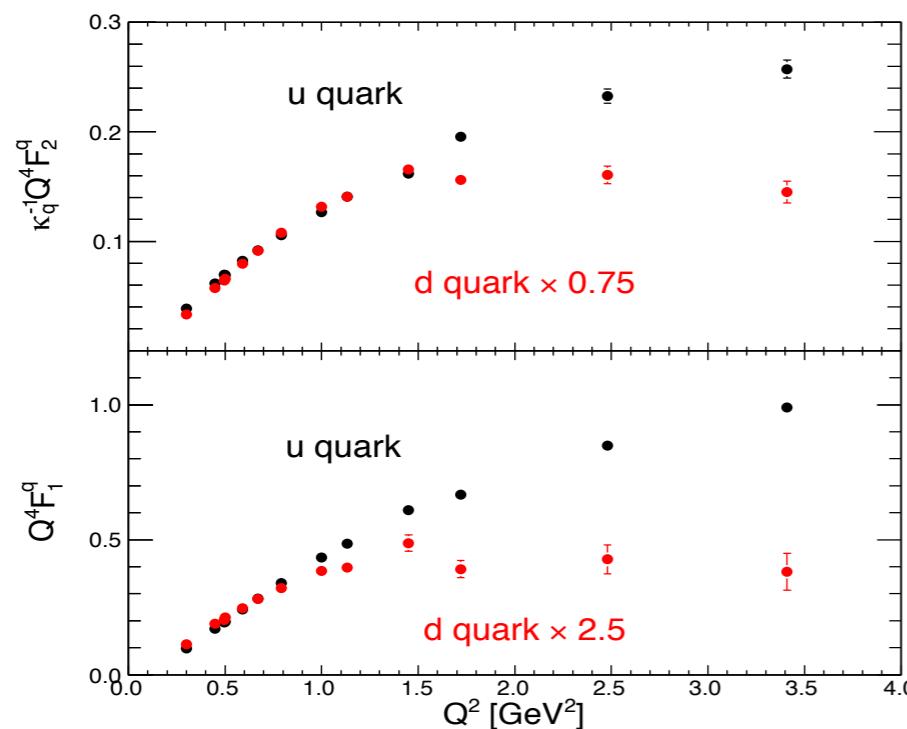
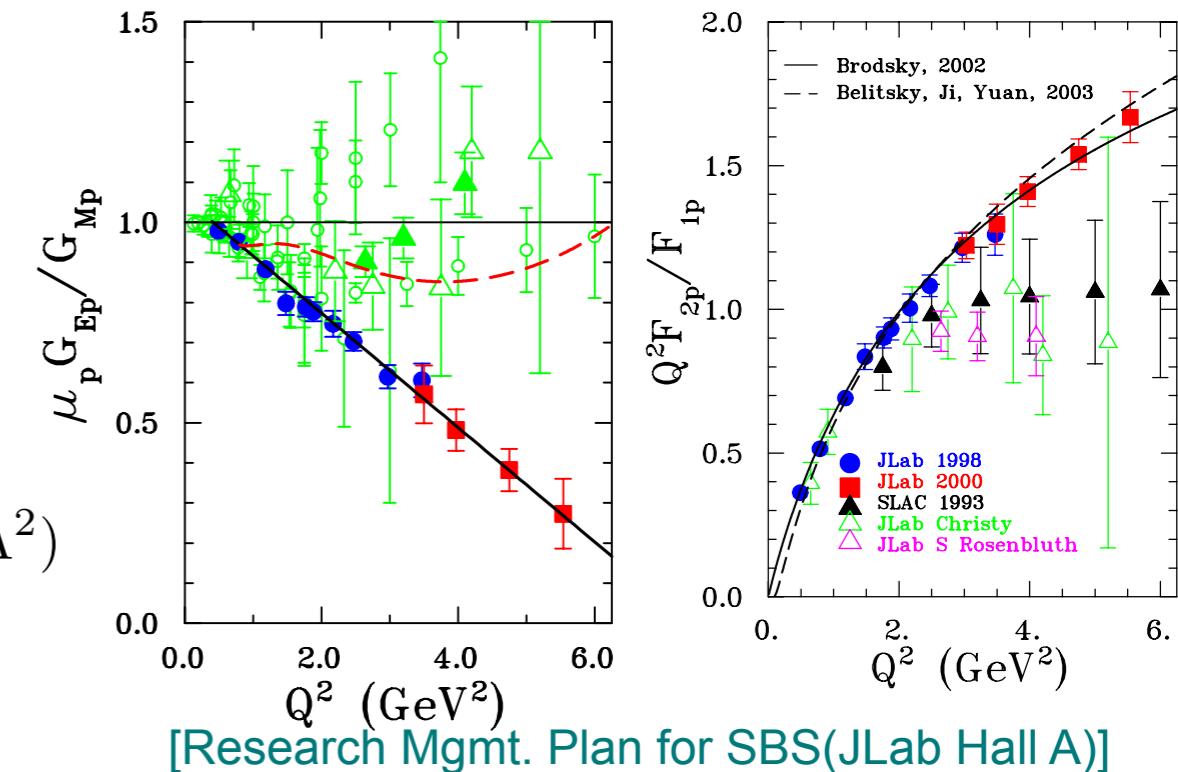
Polarization transfer: polarized e^- beam

- + detect polarization of recoil nucleon
- (alt.: transverse asymmetry on pol. target)
- ➊ G_E/G_M ratio (only small radiative corrections)

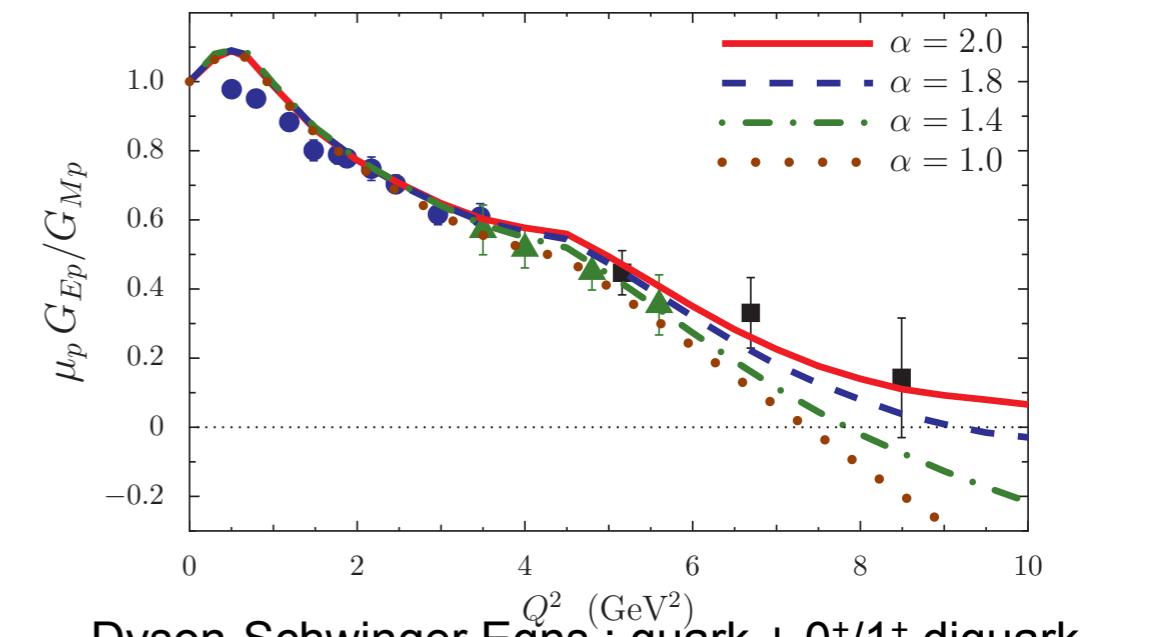
$$P_t/P_l \propto G_E/G_M$$

Why Large- Q^2 Nucleon Form Factors?

- Rosenbluth / Pol.transfer discrepancy in G_E/G_M ?
Input from lattice QCD may help
- Role of diquark correlations in elastic scattering?
Neutron & proton G_E/G_M at/above $Q^2 = 8 \text{ GeV}^2$
- Scale of transition to perturbative QCD?
 (F_2/F_1) scaling at large Q^2 : $Q^2 F_{2p}/F_{1p} \stackrel{?}{\propto} \log^2(Q^2/\Lambda^2)$
- What are contributions from u and d flavors?
Proton and neutron data needed in wide Q^2 range

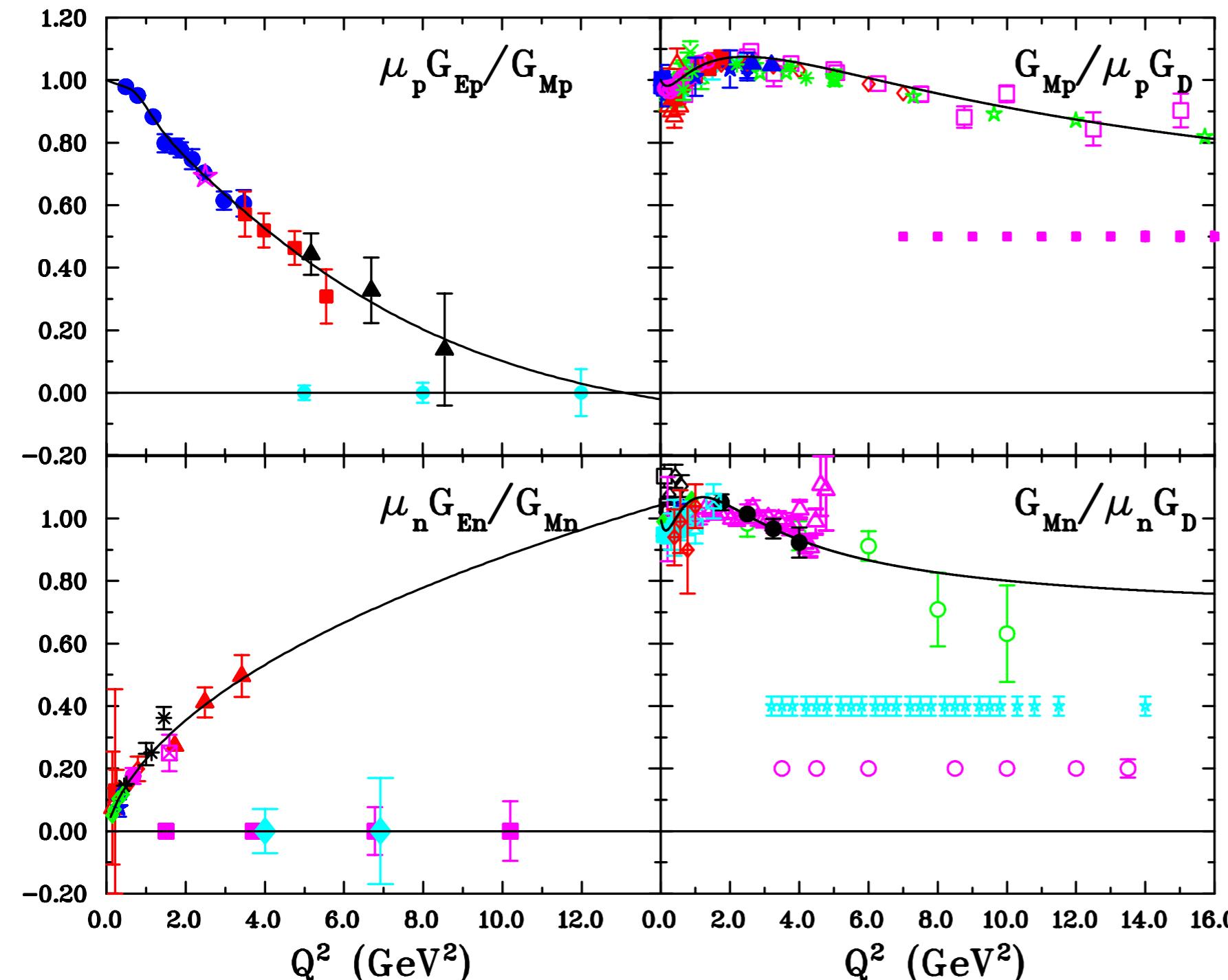


[G.D.Cates, C.W.de Jager, S.Riordan, B.Wojtsekhowski,
PRL106:252003, arXiv:1103.1808]



Dyson-Schwinger Eqns : quark + $0^+/1^+$ diquark
($\alpha \approx$ rate of transition const. quarks \rightarrow pQCD with Q^2)
[Cloet, Roberts, Prog.Part.Nucl.Phys 77:1 (2014)]

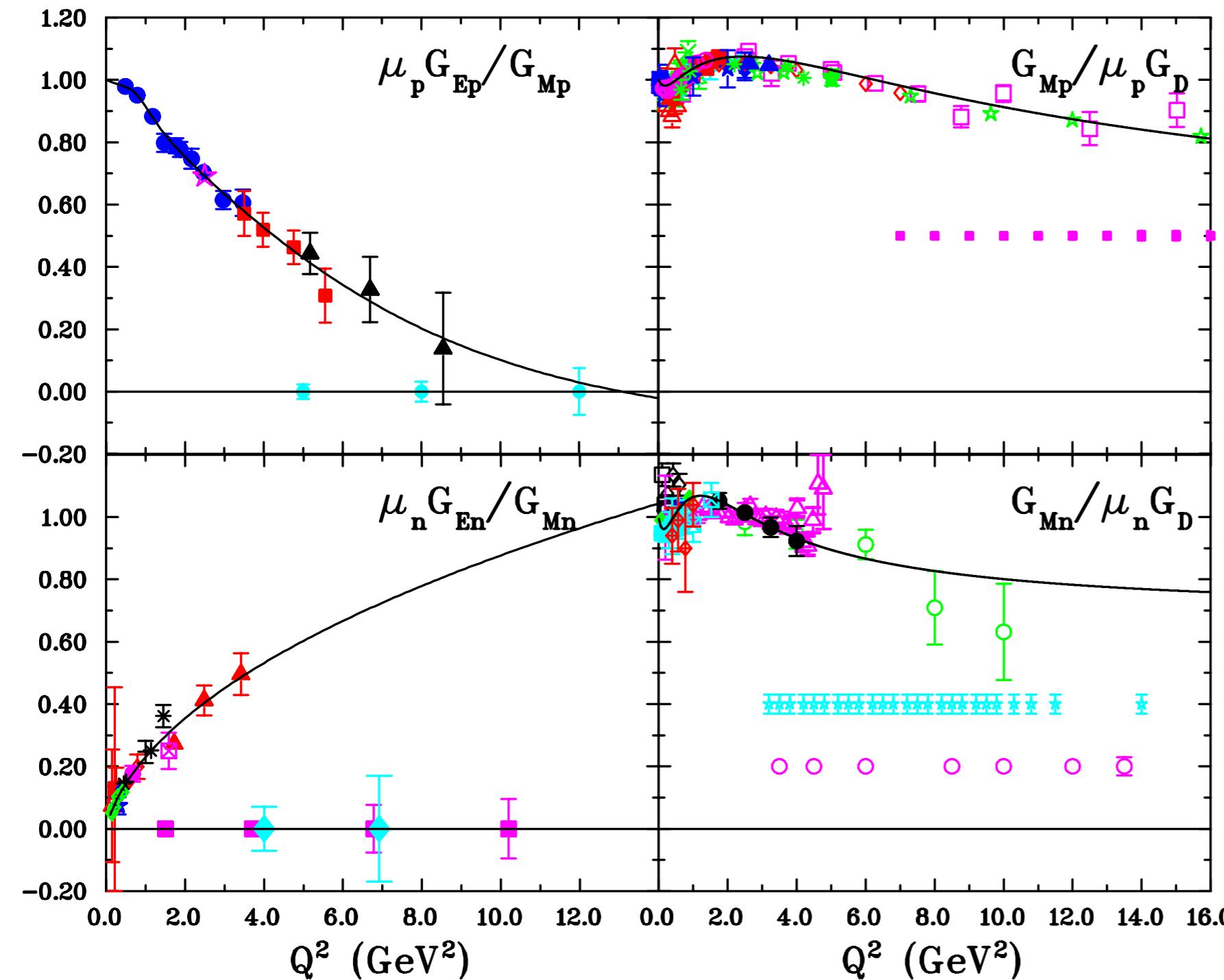
Recent Experiments @JLab : Projections



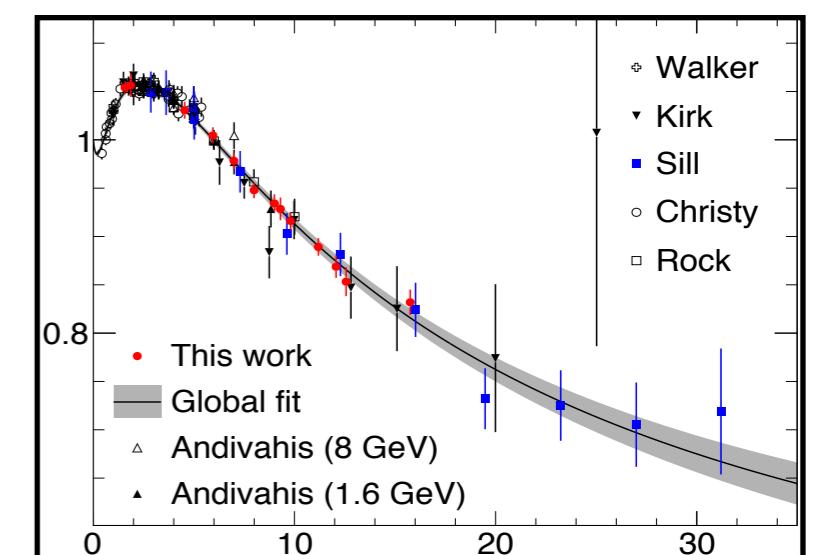
Experiments at JLab@12GeV

- Hall A (HRS, SBS):
 G_{Mp} @ $Q^2 \lesssim 17.5 \text{ GeV}^2$
 G_{Ep}/G_{Mp} @ $Q^2 \lesssim 15 \text{ GeV}^2$;
- Hall B (CLAS12):
 G_{Mn} @ $Q^2 \lesssim 14 \text{ GeV}^2$
- Hall C :
 G_{En}/G_{Mn} @ $Q^2 \lesssim 6.9 \text{ GeV}^2$

Recent Experiments: Projections + G_{Mp} Result



Projected new precision on proton & neutron form factors
 [V. Punjabi et al, EPJ A51: 79 (2015); arXiv: 1503.01452]



New G_{Mp} data from Hall A
 [Christy et al, PRL'22]

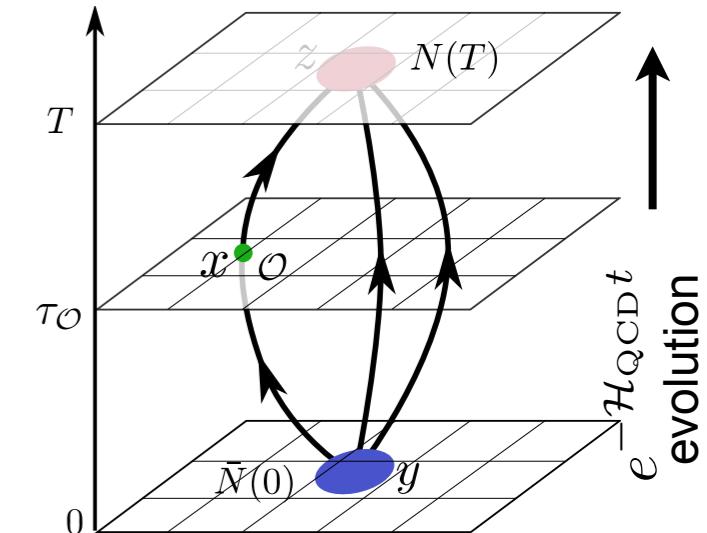
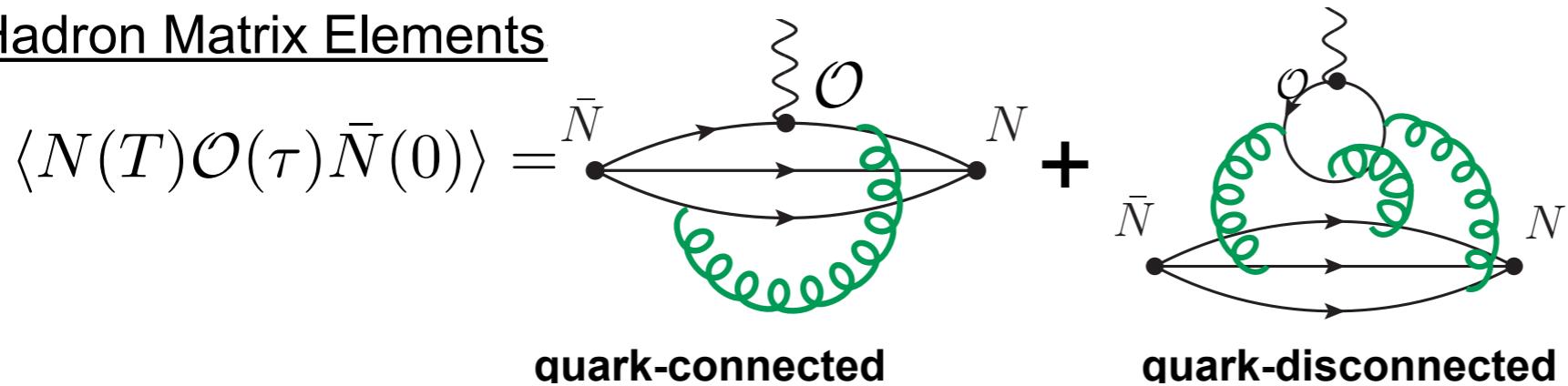
Nucleon Structure from Lattice QCD (in a Nutshell)

- Compute nucleon correlation functions

Hadron Spectrum:

$$\begin{aligned}\langle N(\vec{p}, T) \bar{N}(0) \rangle &= \sum_{\vec{y}} \langle N(\vec{y}, T) \bar{N}(0) \rangle \\ &\sim |Z_0|^2 e^{-E_0 T} [1 + e^{-\Delta E_{10} T} + \dots]\end{aligned}$$

Hadron Matrix Elements



$$\text{Quark lines} = (\not{D} + m)^{-1} \cdot \psi$$

- Extract ground state matrix elements

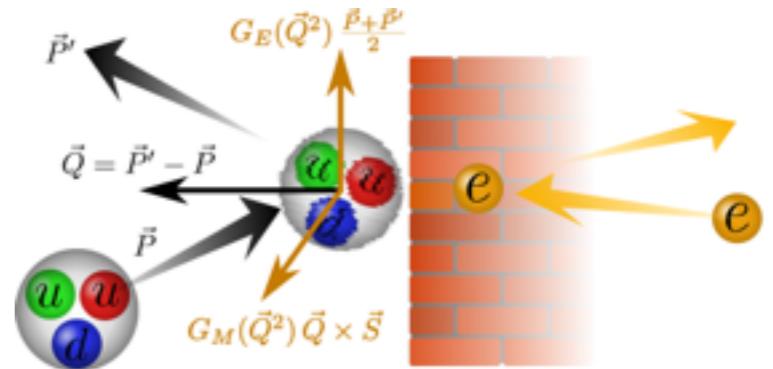
$$\begin{aligned}\langle N(T) \mathcal{O}(\tau) N(0) \rangle &= \sum_{n,m} Z_m e^{-E_n(T-\tau)} \langle n | \mathcal{O} | m \rangle e^{-E_m \tau} Z_n^* \\ T \rightarrow \infty & Z_0 Z_{0'}^* e^{-E_0 \tau - E_{0'}(T-\tau)} [\langle 0' | \mathcal{O} | 0 \rangle + \underbrace{O(e^{-\Delta E \{T, \tau, (T-\tau)\}})}_{\text{excited states}}] \\ &\quad \text{Ground state form factors} \qquad \qquad \qquad \text{Fit and discard} \\ &[\bar{u}' \gamma^\mu u] F_1 + [\bar{u}' \frac{i \sigma^{\mu\nu} q_\nu}{2m_N} u] F_2\end{aligned}$$

- Systematic uncertainties
- discretization effects
 - unphysical-heavy pion mass
 - finite volume
 - excited state contributions
 - quark-disconnected diagrams

Challenges at Large Q^2 on a Lattice

- Excited-state gaps shrink \Rightarrow **more contamination**

- in Breit frame : $p_{\text{in}} = p_{\text{out}} \approx 1.6 \text{ GeV}$ for $Q^2 = 10 \text{ GeV}^2$
- $E_1 - E_0 = \sqrt{M_1^2 + \vec{p}^2} - \sqrt{M_2^2 + \vec{p}^2} < M_1 - M_0$
- $N(\sim 1500) \Rightarrow \Delta E = 600 \rightarrow 360 \text{ MeV}$



- Stochastic noise grows faster with T [Lepage'89]:

Signal	$\langle N(T) \bar{N}(0) \rangle$	$\sim e^{-E_N T}$
Noise	$\langle N(T) \bar{N}(0) ^2 \rangle - \langle N(T) \bar{N}(0) \rangle ^2$	$\sim e^{-3m_\pi T}$
Signal/Noise		$\sim e^{-(E_N - \frac{3}{2}m_\pi)T}$

SNR / SNR[static] $\sim 1 / 100$
per $\Delta t = 1 \text{ fm/c}$ "time"
 $[Q^2 \approx 10 \text{ GeV}^2, \text{phys. quarks}]$

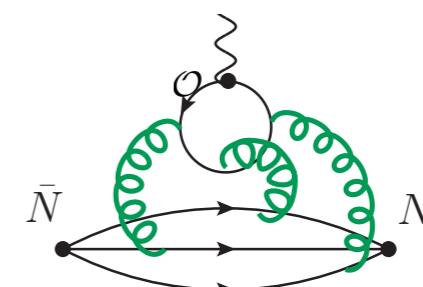
- Discretization effects:

O(a) Correction to current operator
expected to grow with momentum

$$(V_\mu)_I = [\bar{q} \gamma_\mu q] + c_V a \underbrace{\partial_\nu [\bar{q} i \sigma_{\mu\nu} q]}_{\propto Q}$$

- Quark-disconnected contributions:

negligible ($\lesssim 1\%$) at $Q^2 \leq 1 \text{ GeV}^2$, **unknown at large Q^2**



- No reliable EFT/ChPT for pion mass-, volume-extrapolation at large momentum p_N, Q

Many of the same challenges as for parton physics on a lattice
(LaMET, Quasi-/Pseudo- PDFs/DAs/TMDs)

Present QCD Calculation Parameters

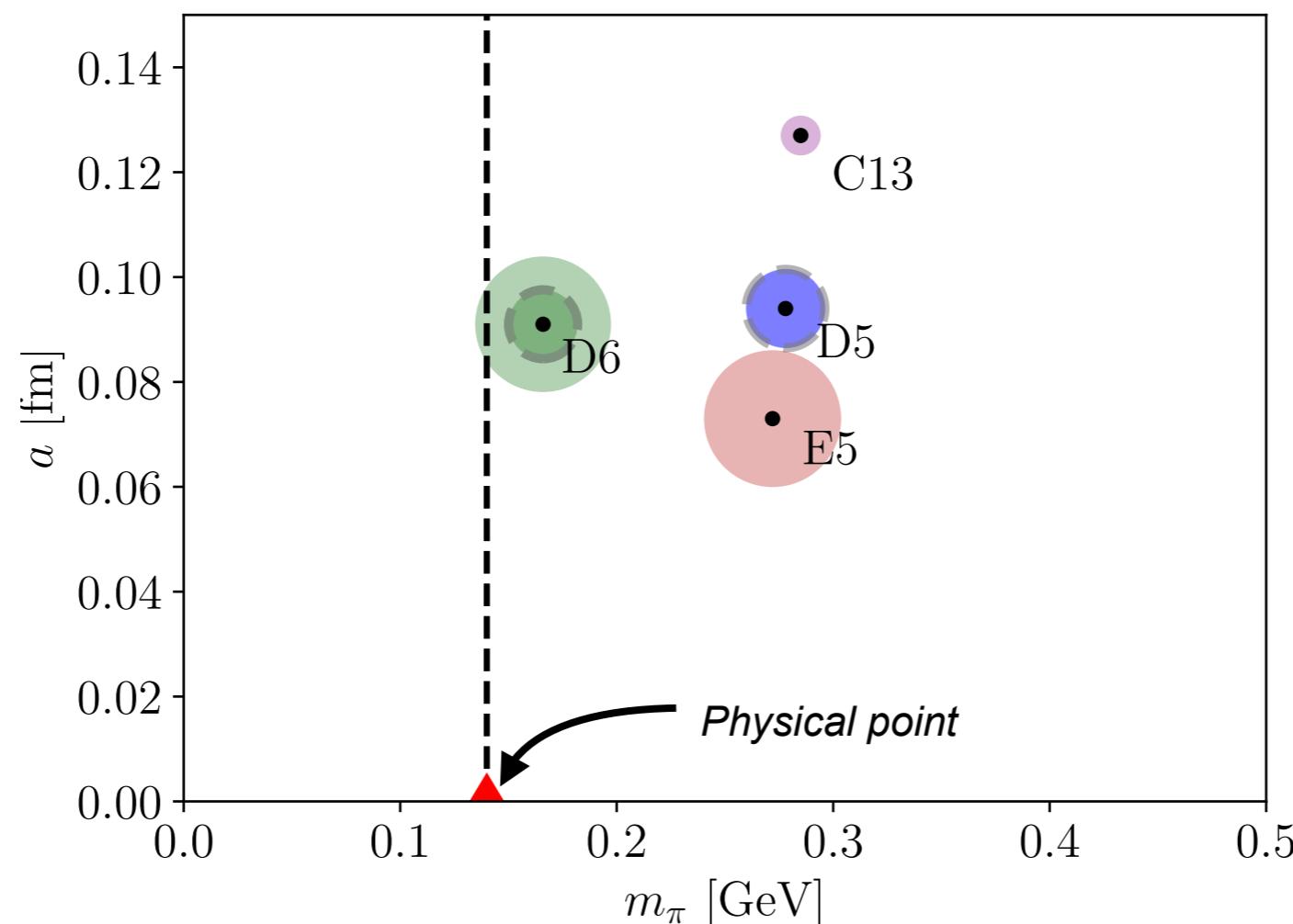
QCD action with $N_F = 2+1$ $O(a^2)$ -improved dynamic quarks

Many thanks to [JLab / W&M / LANL / MIT]

- Lattice spacing $a \approx 0.073 \div 0.091$ fm
- $SU(2)_f$ -symmetric + strange quarks
- Pion mass $m\pi = 170 \div 280$ MeV
- Physical volume $L \gtrsim 3.7 (m\pi)^{-1}$
- Euclidean time $t_{sep} = 0.5 \div 1.1$ fm

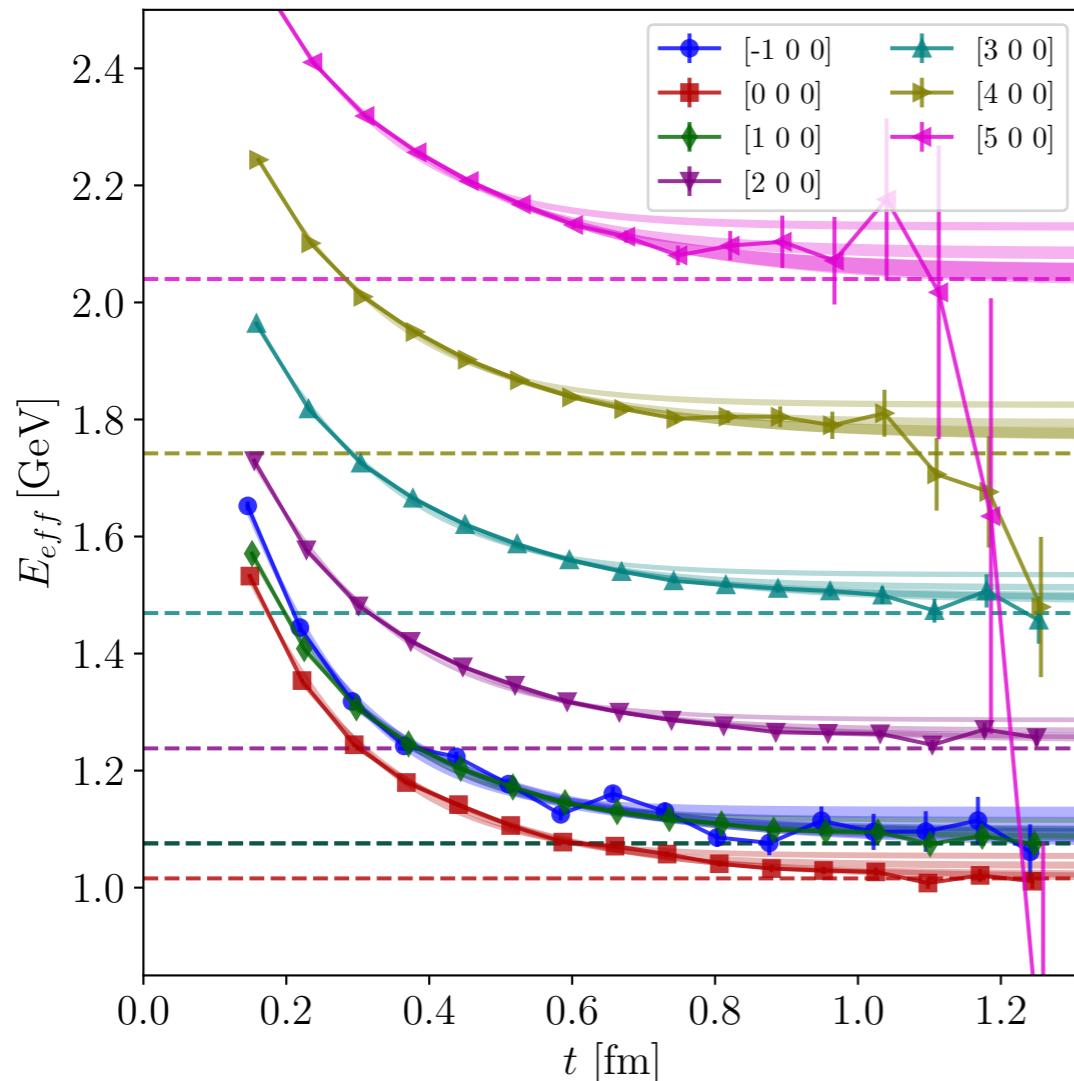
2022/24:

- MC Statistics $\gtrsim 250k$ on D6 ($48^3 \times 96$), E5 ($48^3 \times 128$)
- Disconnected contractions on D5, D6 (1200+ configs)



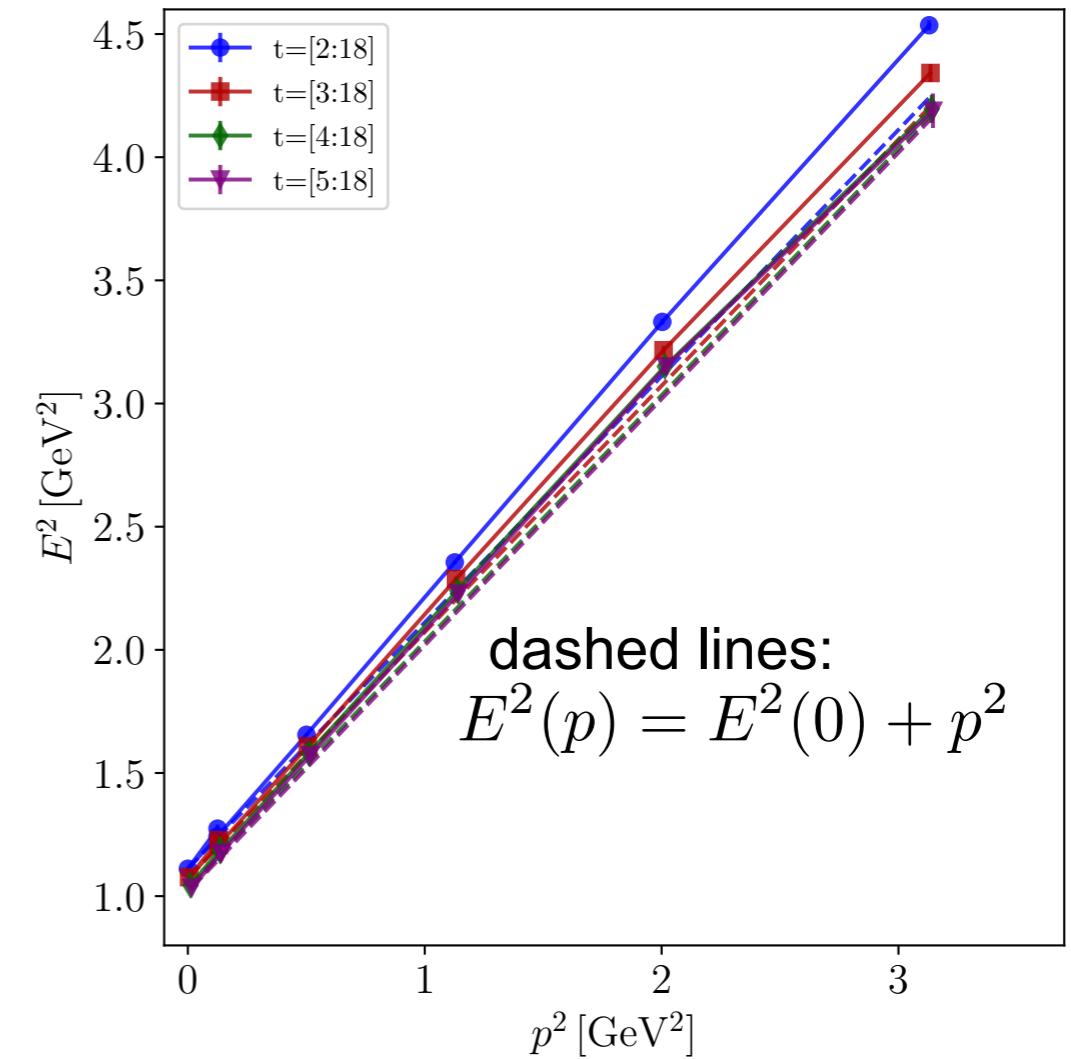
Lattice Nucleon Energy & Dispersion Relation (E5)

- E5 : $m\pi = 272 \text{ MeV}$, spacing $a = 0.073 \text{ fm}$, 266k MC samples
- Ground-state energy follows disp.relation for fits with $t_{sep} \gtrsim 0.3 \text{ fm}$



- 2-state fits compared to Effective energy

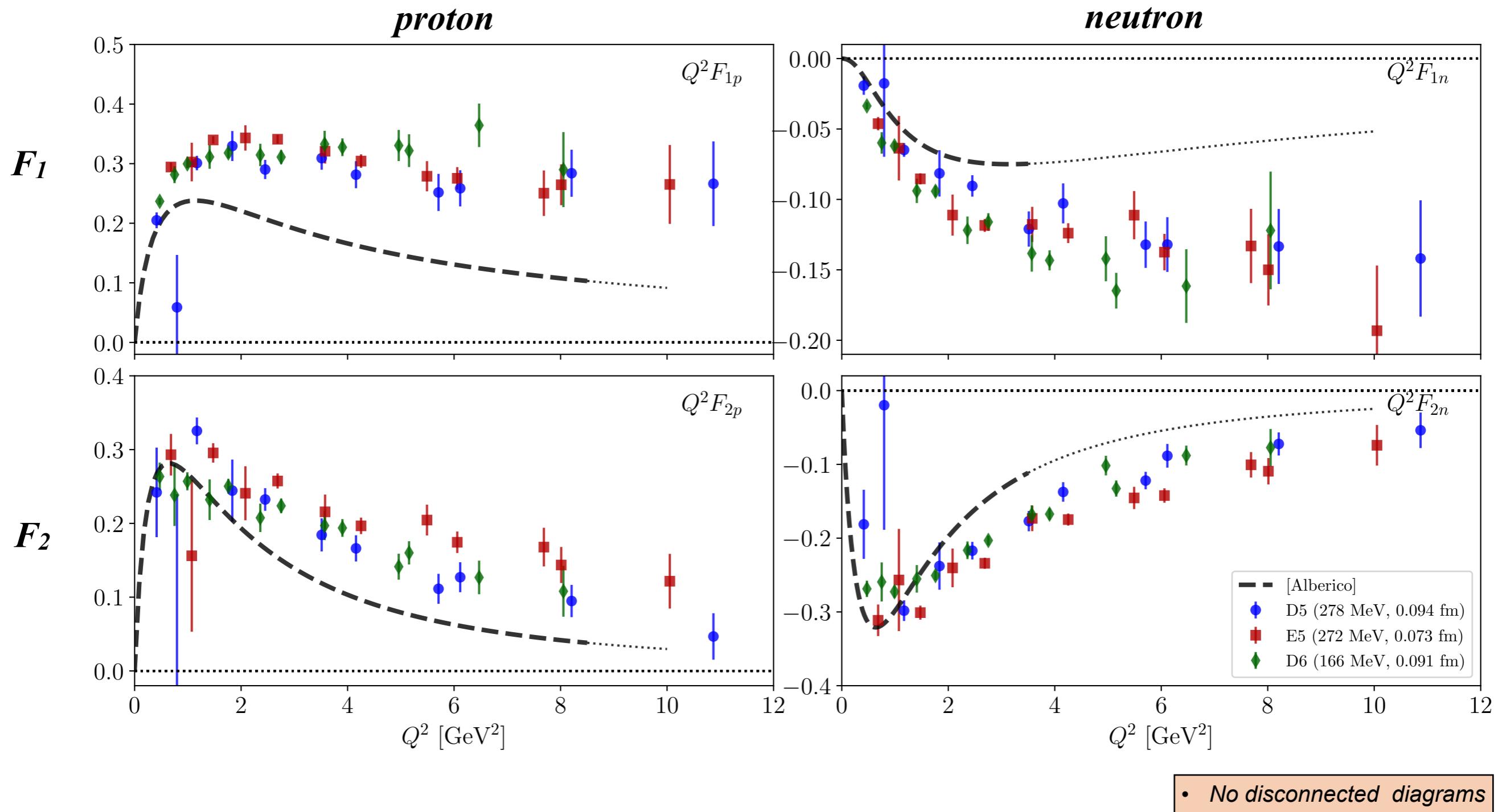
$$E_{eff} = \frac{1}{a} \log \frac{C_{N\bar{N}}(t)}{C_{N\bar{N}}(t+a)}$$



- Lattice energies vs. continuum dispersion relation (dashed)

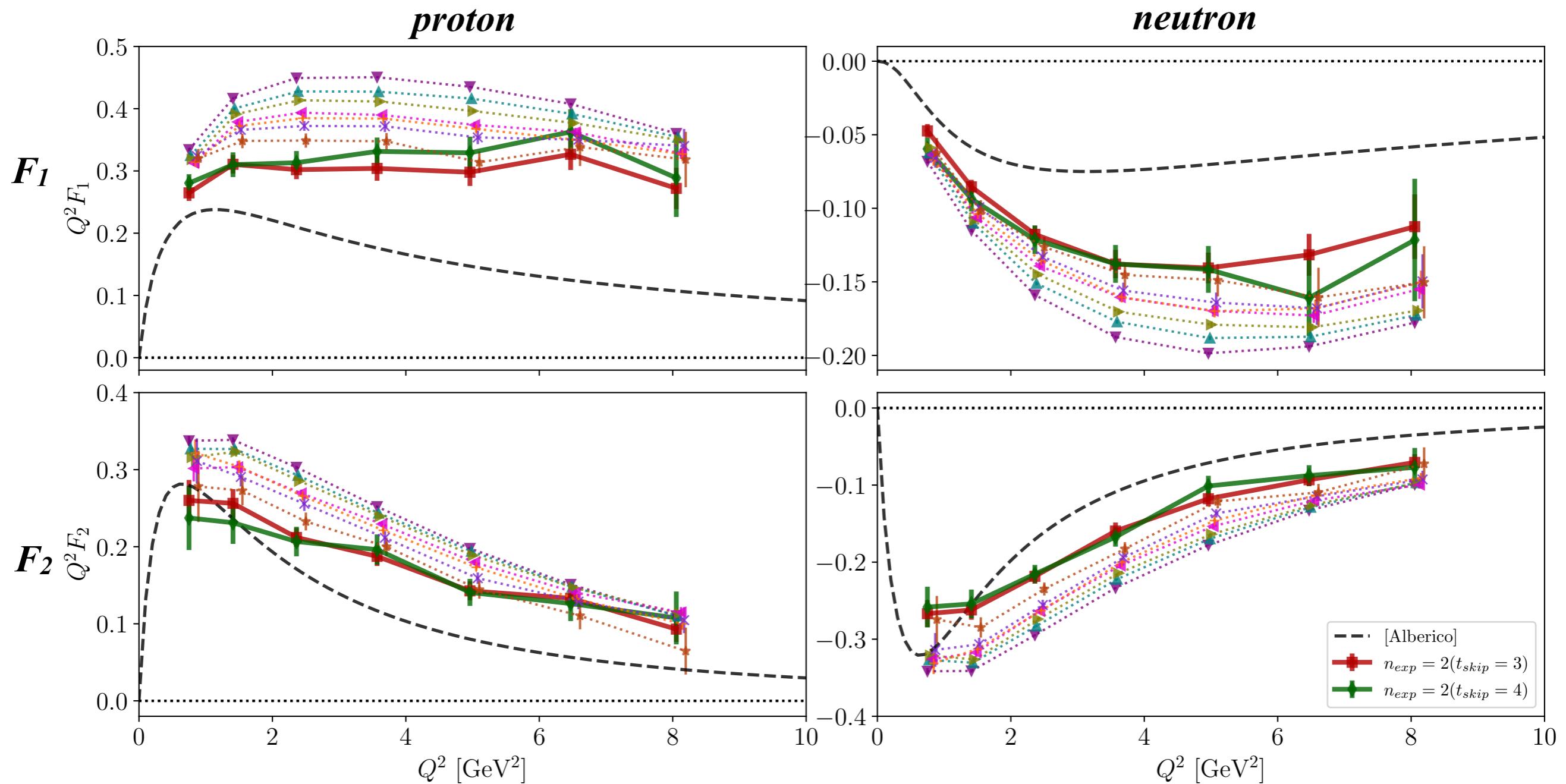
Nucleon Form Factors: Ensemble Comparison

- "Ground" state from **CONSERVATIVE** 2-state fits with $t_{sep} = 0.7 \div 1.1$ fm
- Comparison of 3 ensembles (D5 : 86k, E5 : 266k, D6 : 261k samples)
- Phenomenology (dashed) : [Alberico et al, PRC79:065204 (2008)]*



Nucleon Form Factors: 2-state fit vs. fixed $T_{sep}(D6)$

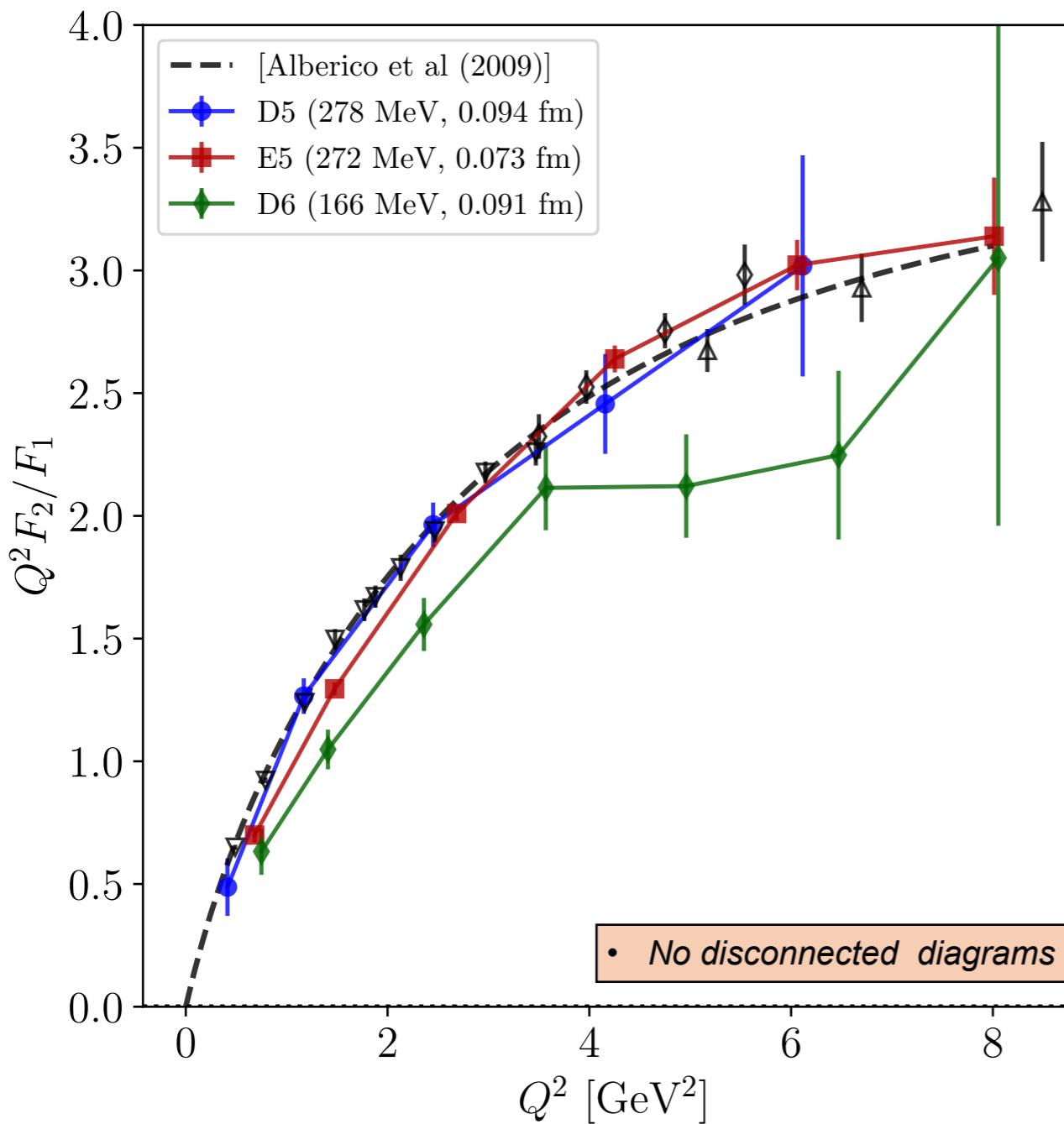
- D6 ensemble (260k samples) : fits with $t_{sep}=0.7\div1.1\text{ fm}$ and $t_{sep}=0.5\div1.1\text{ fm}$ range
- Dotted lines: comparison of 1-state approx. for various fixed t_{sep}
- Phenomenology (dashed)* : [Alberico et al, PRC79:065204 (2008)]



• No disconnected diagrams

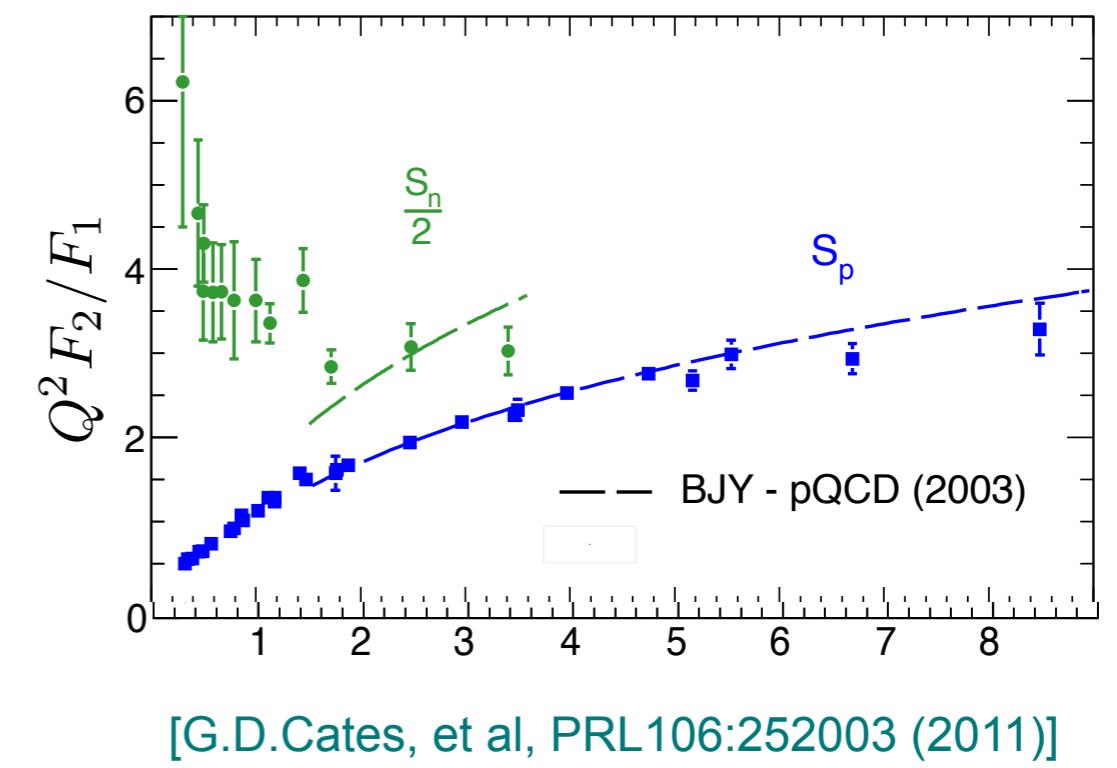
Proton F_2/F_1 Ratio

- Comparison of 3 ensembles (**D5** : 86k, **E5** : 266k, **D6** : 261k samples) ; fit $t_{sep} = 0.7 \div 1.1$ fm
- *Phenomenology (dashed)* : [Alberico et al, PRC79:065204 (2008)]
- Proton experimental data $Q^2 \leq 8.5$ GeV 2 (black points)



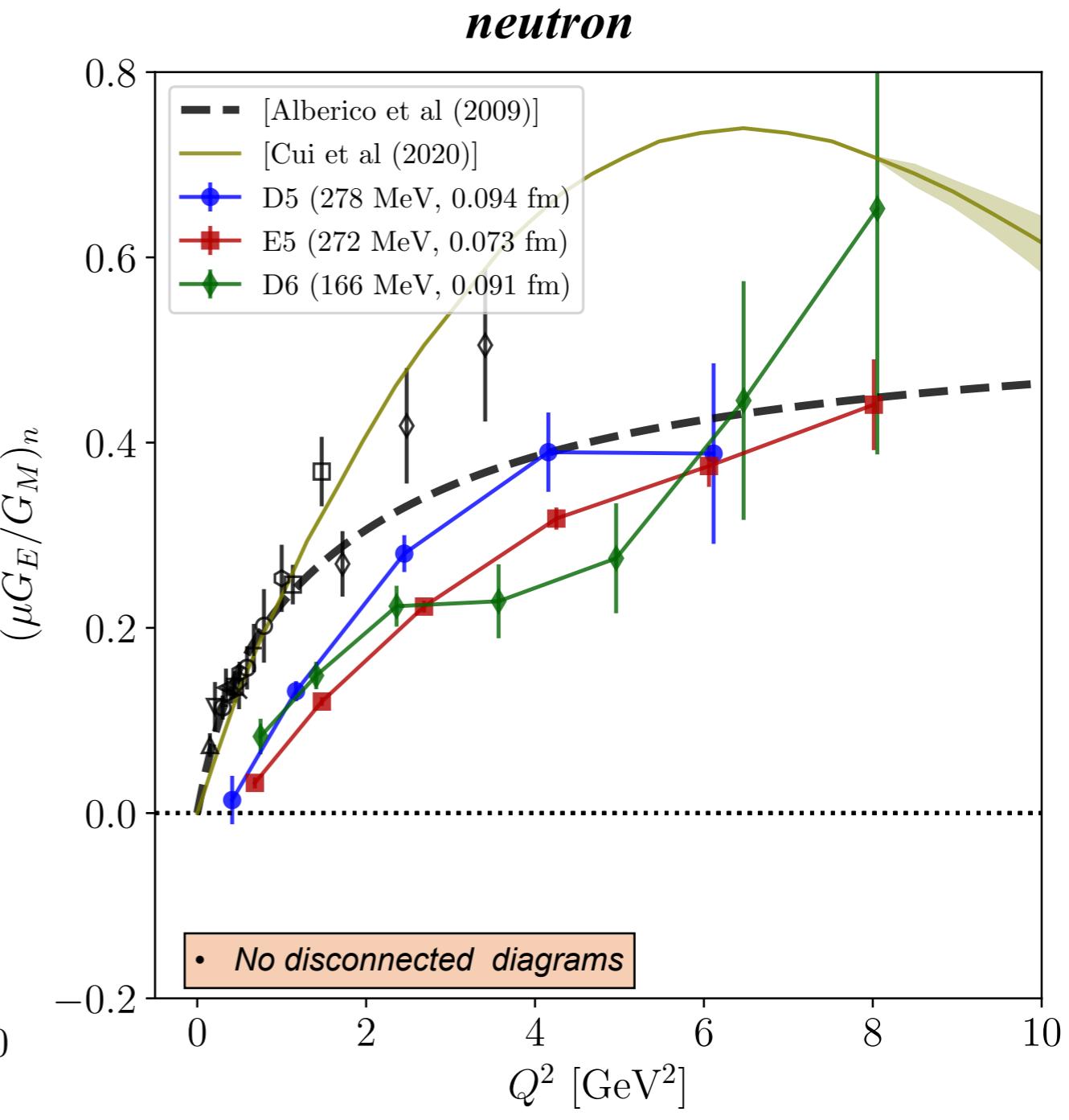
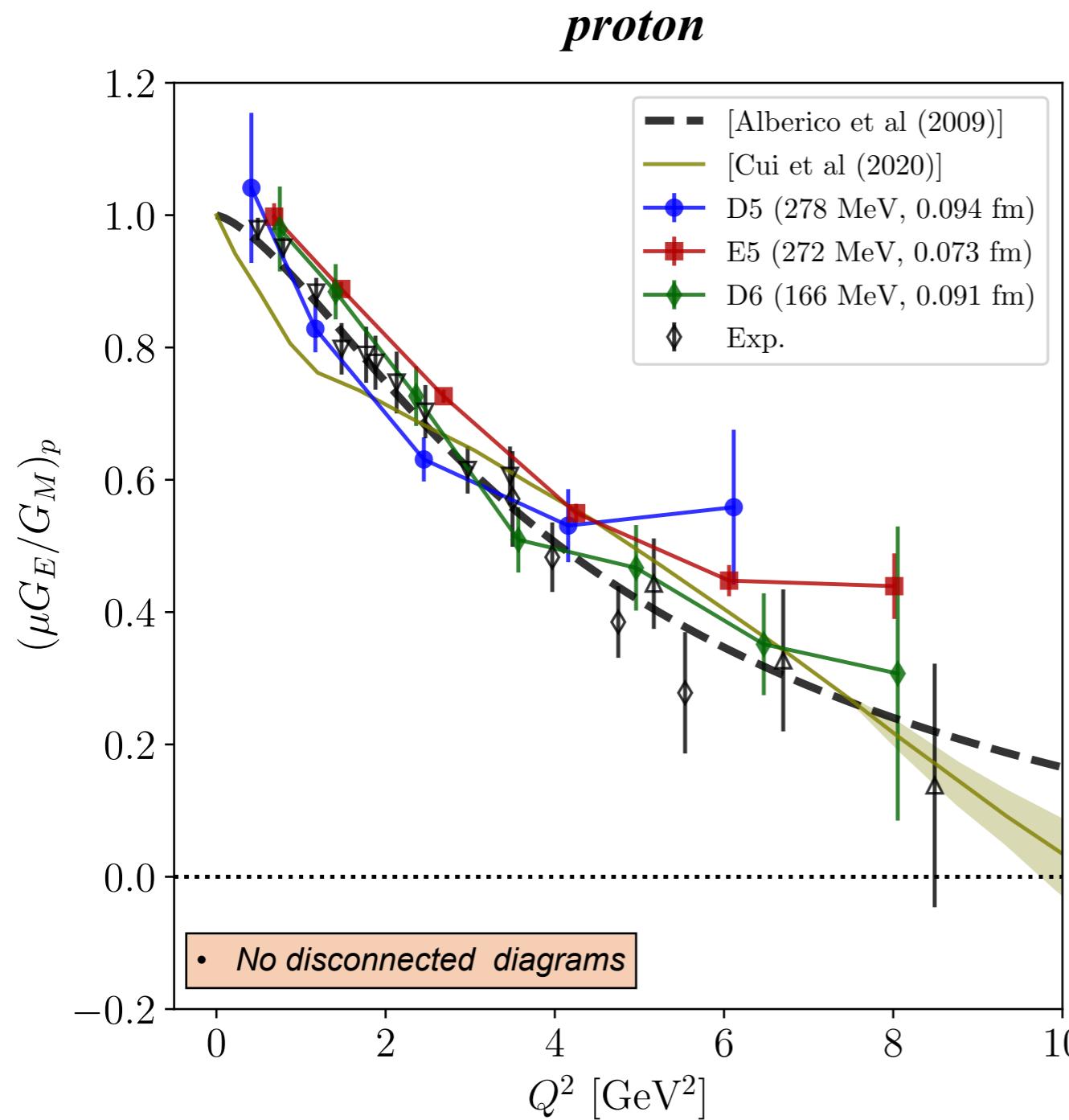
● Prediction from pQCD + quark OAM
[Balitsky, Ji, Yuan (2003)]

$$Q^2 F_{2p}/F_{1p} \stackrel{?}{\propto} \log^2(Q^2/\Lambda^2)$$



Proton & Neutron G_E/G_M Ratio (min. $t_{sep}= 0.5$ fm)

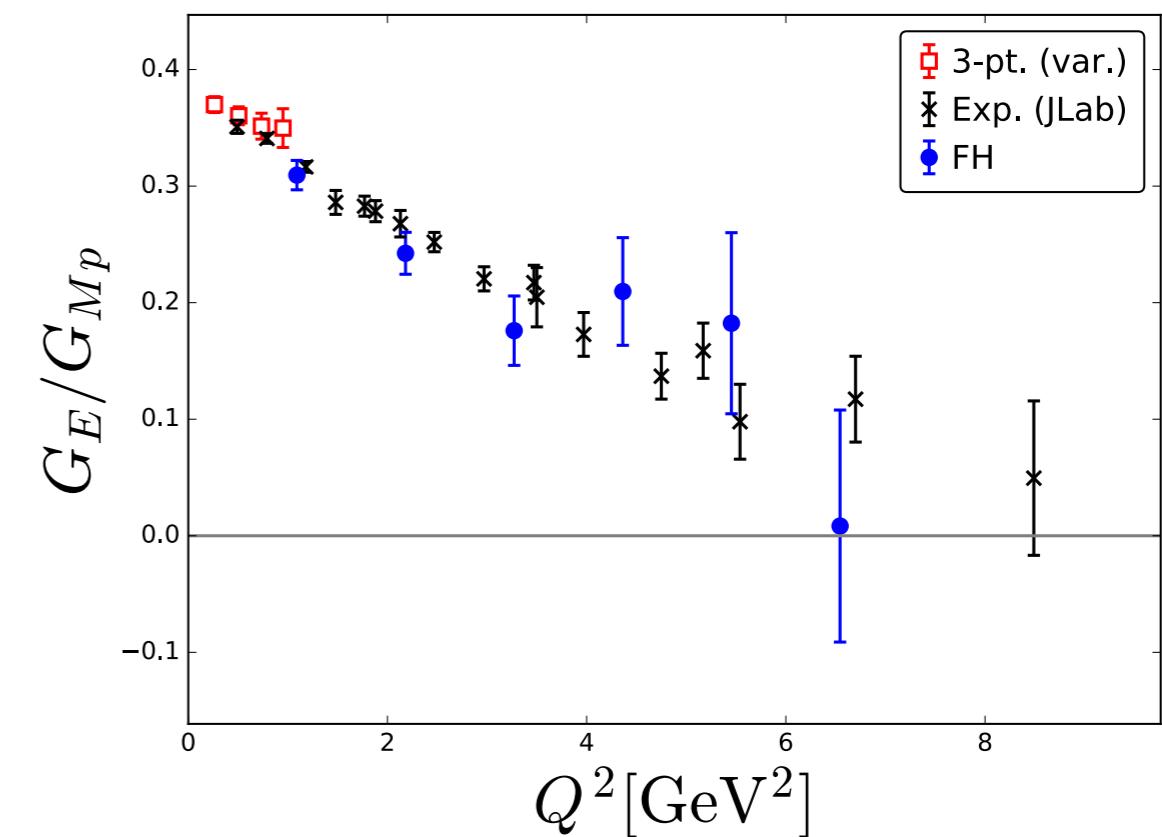
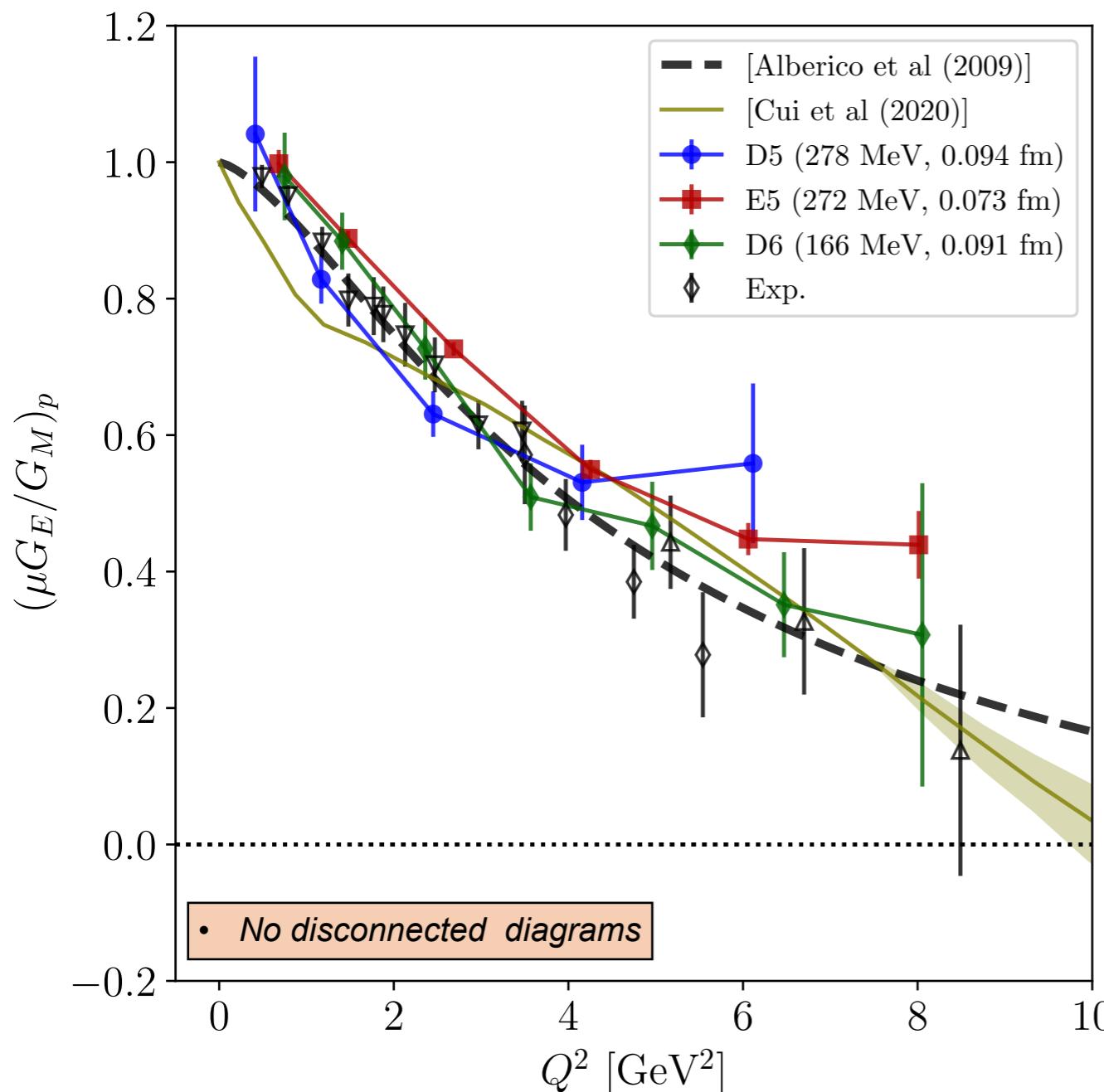
- Comparison of 3 ensembles (**D5** : 86k, **E5** : 266k, **D6** : 261k samples) ; fit $t_{sep}= 0.5 \div 1.1$ fm
- Phenomenology : [Alberico et al, PRC79:065204 (2008)] ;
- Experimental data (black points) $Q^2 \leq 8.5$ GeV 2 (proton) and $Q^2 \leq 3.4$ GeV 2 (neutron)



Proton & Neutron G_E/G_M Ratio (min. $t_{sep}= 0.5$ fm)

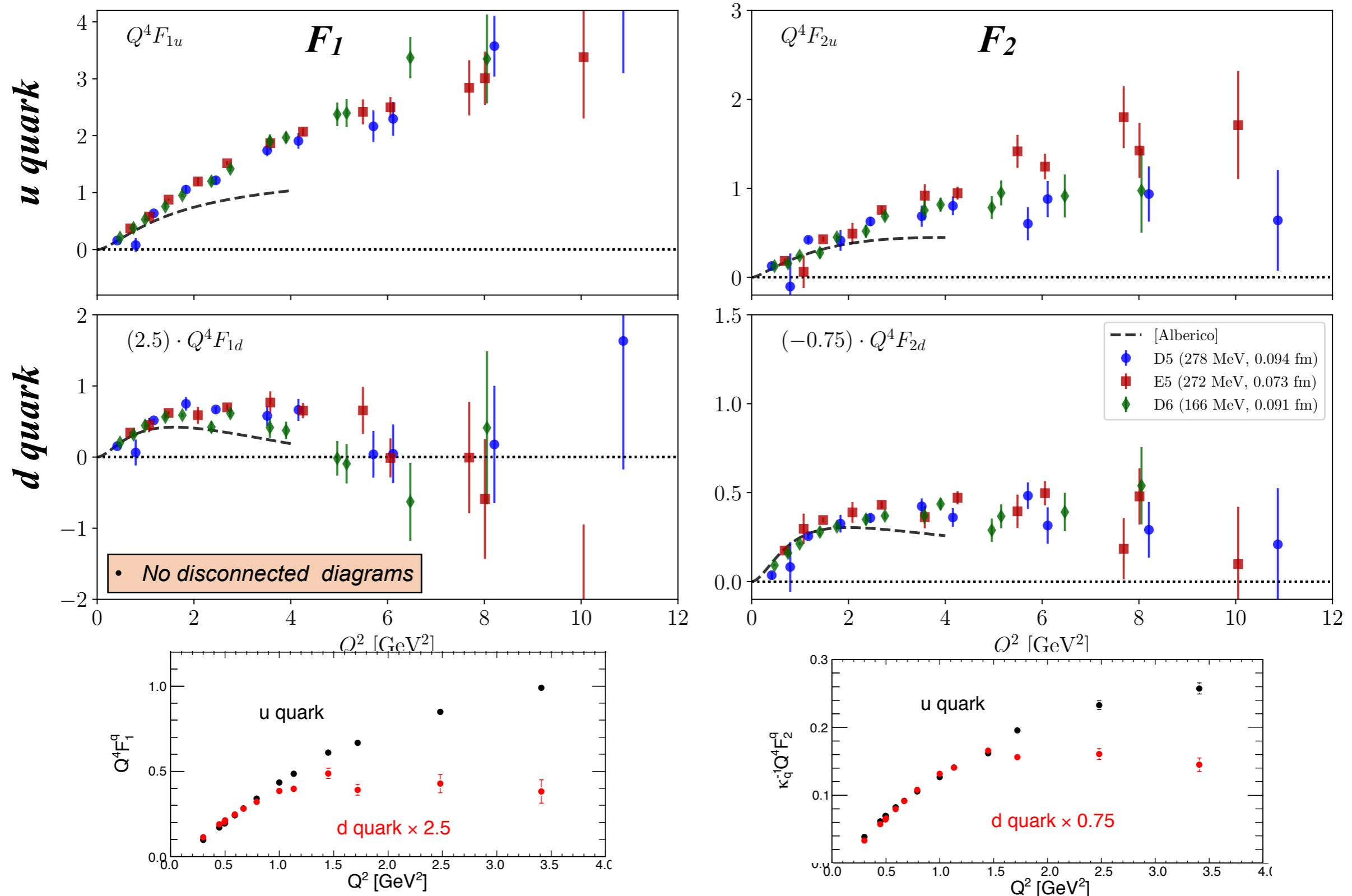
- Comparison of 3 ensembles (**D5** : 86k, **E5** : 266k, **D6** : 261k samples) ; fit $t_{sep}= 0.5 \div 1.1$ fm
- Phenomenology : [Alberico et al, PRC79:065204 (2008)] ;
- Experimental data (black points) $Q^2 \leq 8.5$ GeV 2 (proton) and $Q^2 \leq 3.4$ GeV 2 (neutron)

proton



Earlier calculation: ($a=0.074$ fm, $m_\pi=470$ MeV)
 Feynman-Hellman method
 [Chambers et al (CSSM), PRD96: 114509]

Light-Flavor Decomposition (Proton)



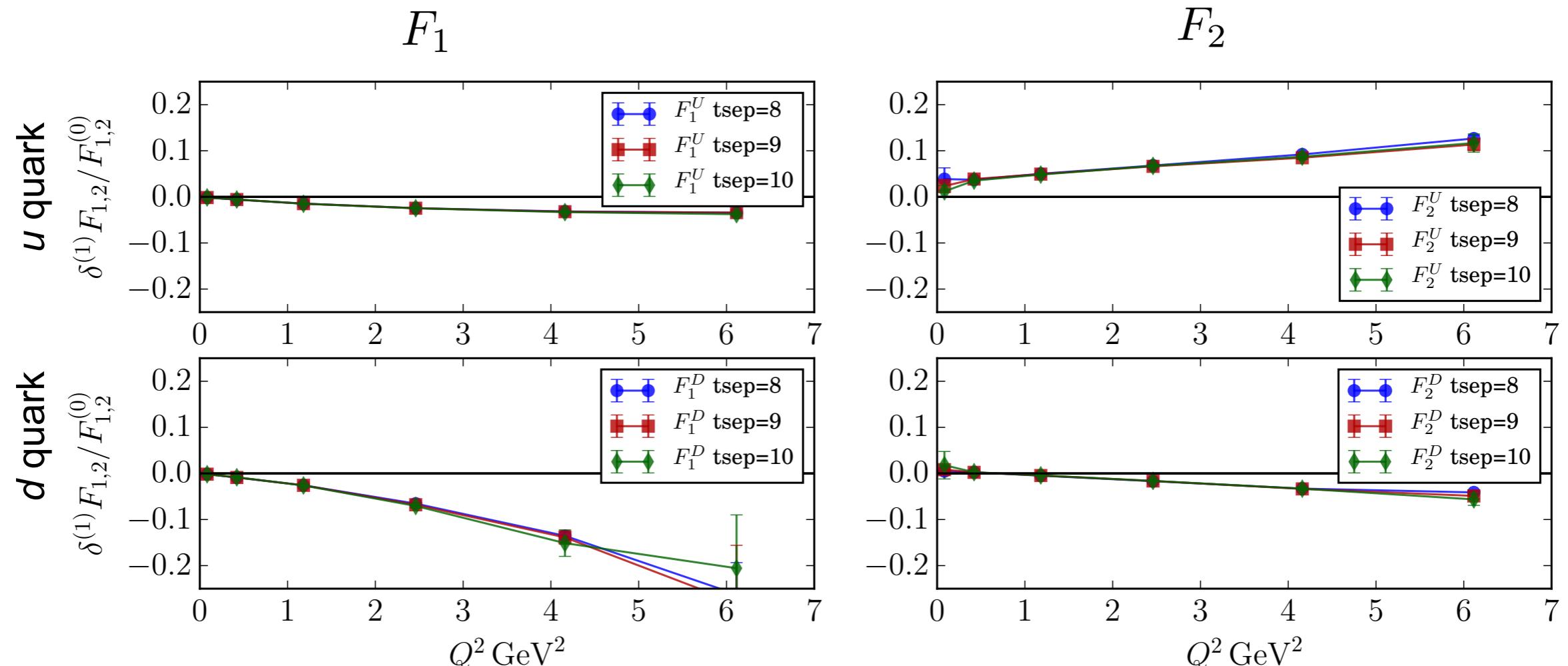
- Similar qual. features of flavor dependence [G.D.Cates, et al, PRL106:252003(2011)]

O(a) Vector Current Correction

- No disconnected diagrams

Improved vector current $(V_\mu)_I = \bar{q}\gamma_\mu q + c_V a\partial_\nu \bar{q}i\sigma_{\mu\nu}q$

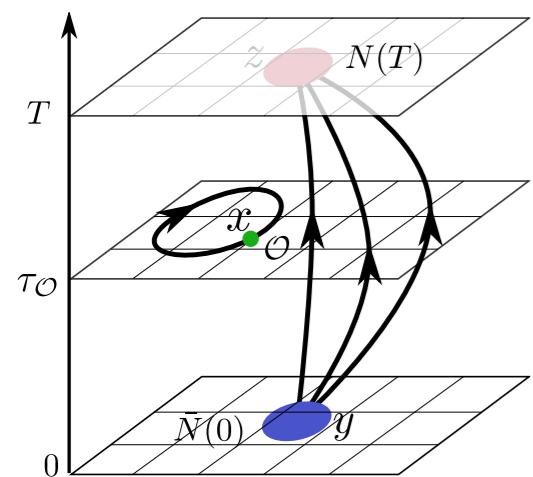
$O(a^1)$ correction : form factors of $a \langle N | \partial_\nu (\bar{q}i\sigma^{\mu\nu}q) | N \rangle$



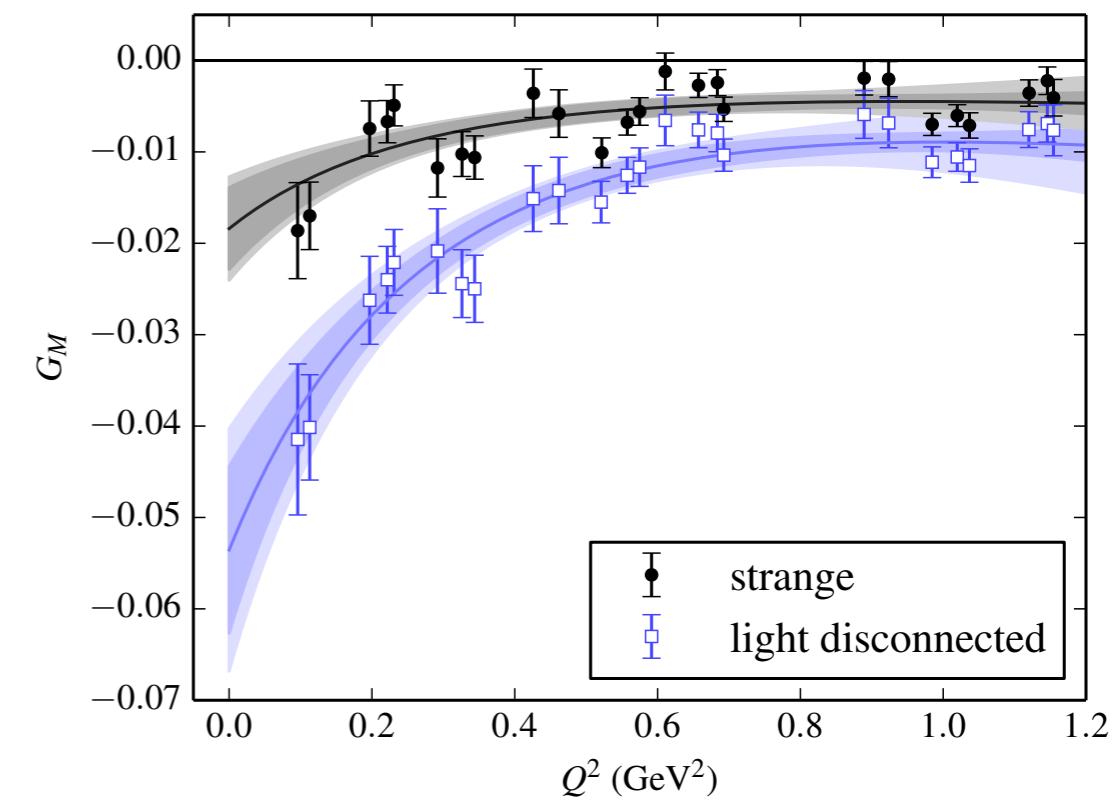
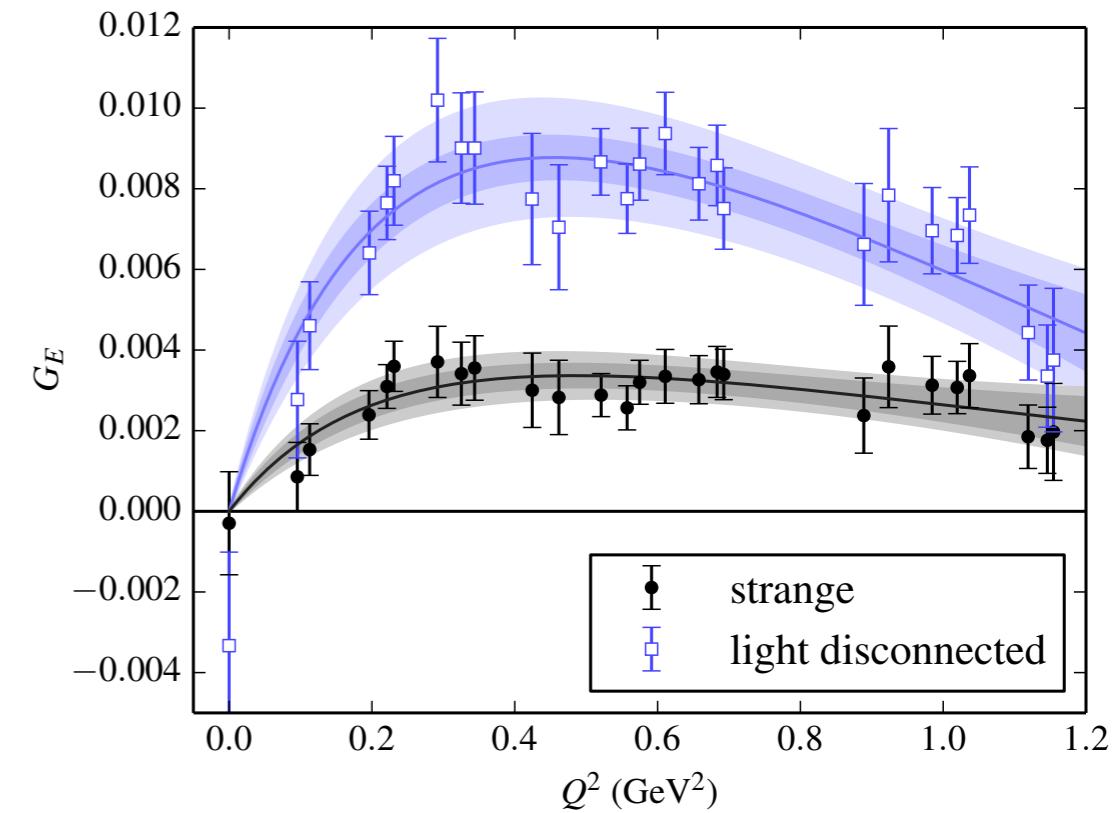
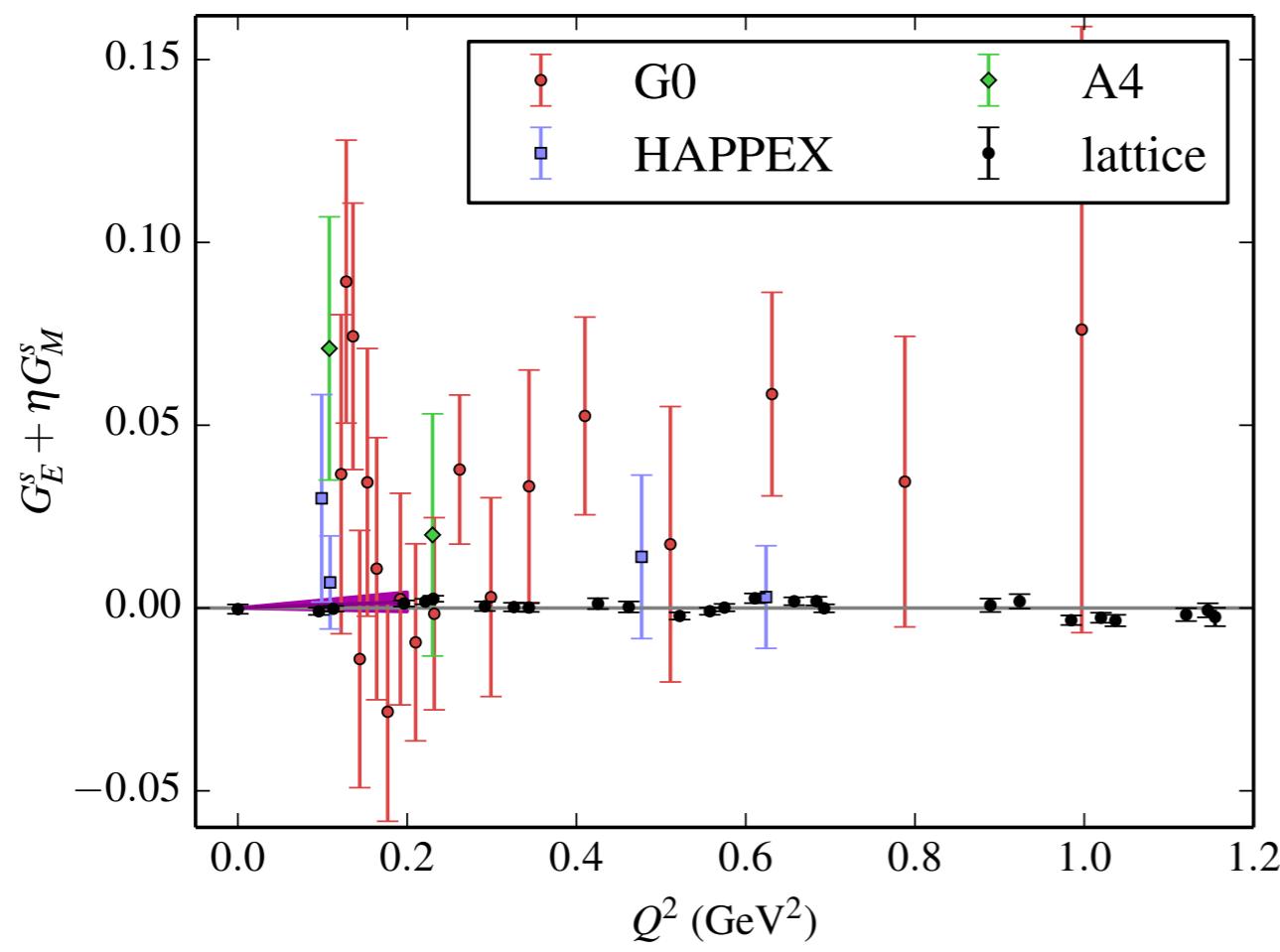
Relative magnitude of $O(a^1)$ effects : $\{O(a^1)\} / \{O(a^0)\}$ form factors
(assuming $c_V=0.05$)

- improvement coefficient c_V : must be computed on lattice from WI
- perturbation theory: $c_V \approx -0.01 C_F(g_0)^2$

Disconnected Contributions to Vector FFs?



Quark-Disconnected diagrams contribute $\lesssim 1\%$ to nucleon FFs
[J. Green, S. Meinel, S.S.
et al; PRD92:031501 (2015)]



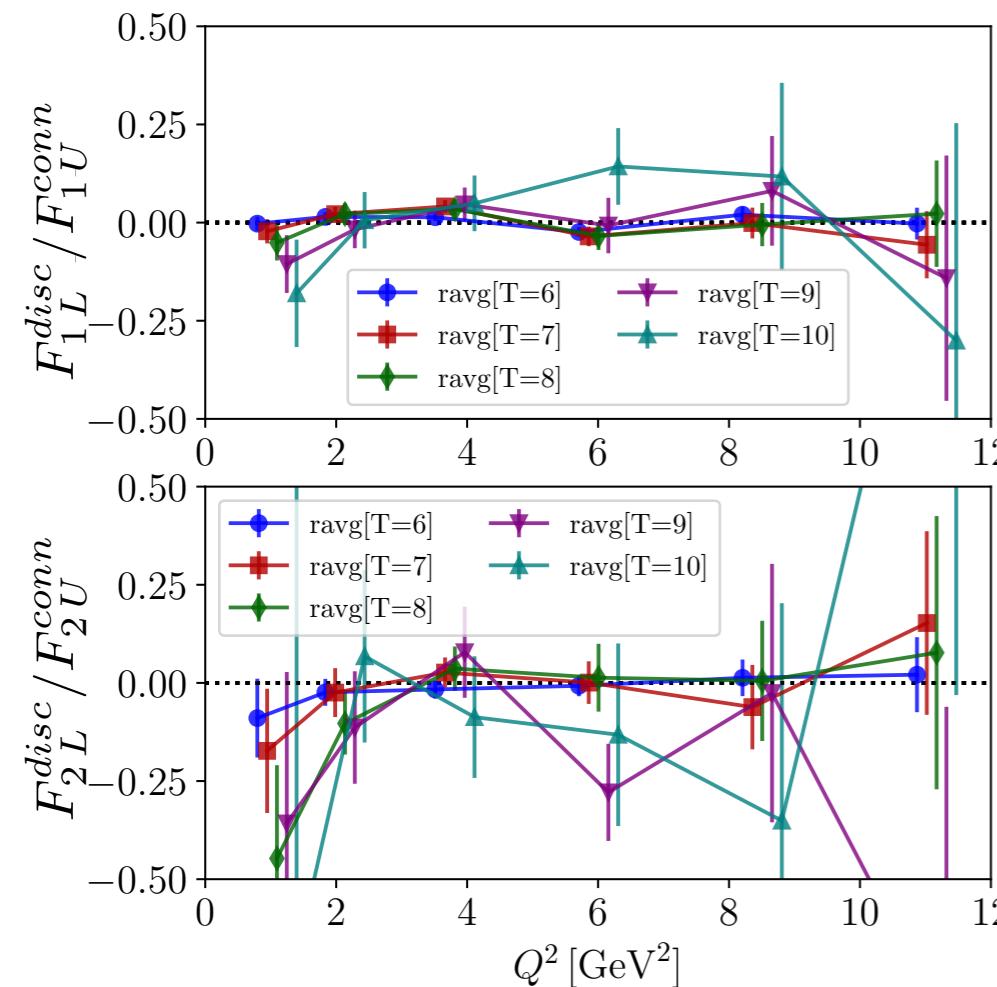
Disconnected Light & Strange vs. Connected (D5)

- relative correction $F_{1,2}^{\text{disc}} / F_{1,2}^{\text{conn}}$ from plateau averages $t_{\text{sep}} = 0.5 \div 0.9 \text{ fm}$, $Q^2 \lesssim 11 \text{ GeV}^2$
- D5 ensemble ($m_\pi = 280 \text{ MeV}$, $a = 0.094 \text{ fm}$), 1346 configs \otimes 64 samples of $\langle N\bar{N} \rangle$
- partial noise cancellation between $L=U/D$ and S in proton & neutron

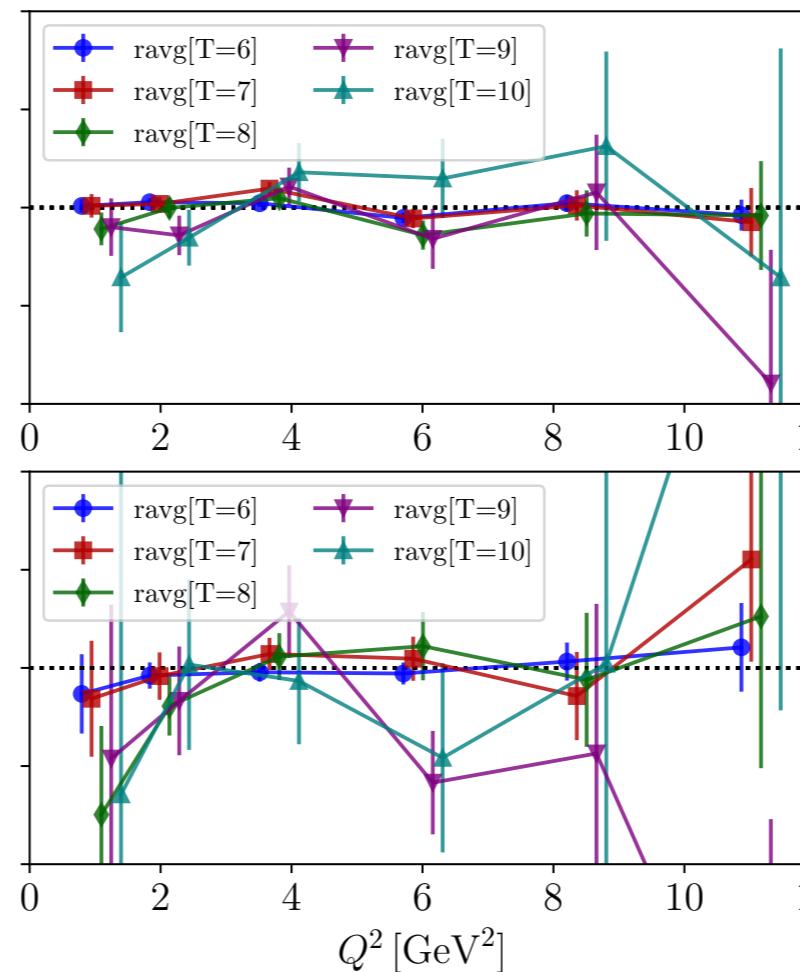
$$P = \frac{1}{3} [2U - D]_{\text{conn}} + \frac{1}{3} [L - S]_{\text{disc}}$$

$$N = \frac{1}{3} [2D - U]_{\text{conn}} + \frac{1}{3} [L - S]_{\text{disc}}$$

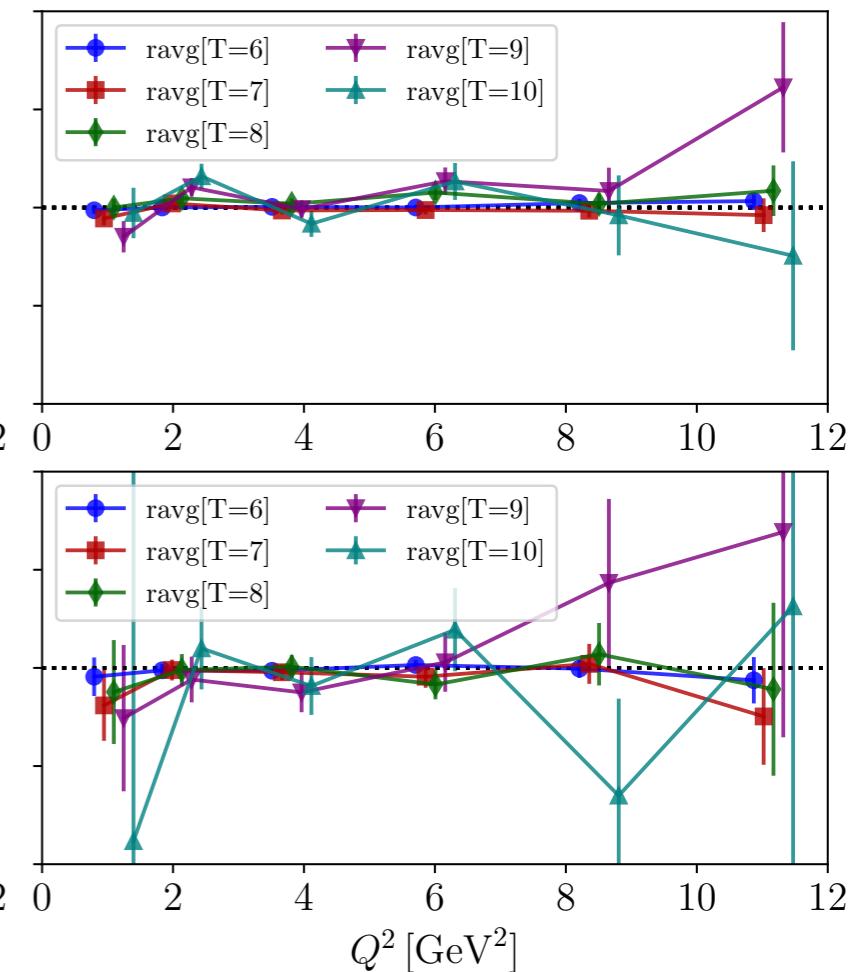
disconnected $L=U$ or D



disconnected S



disconnected $(L-S)$



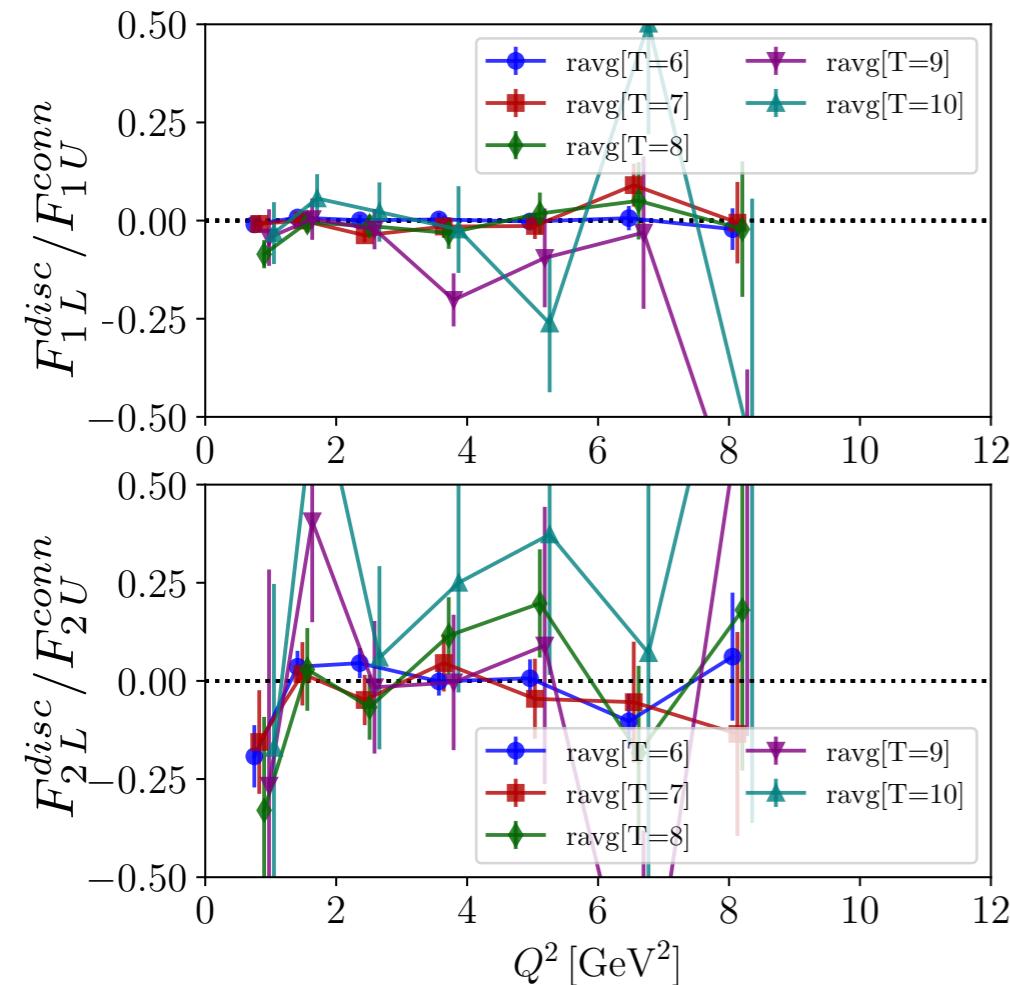
Disconnected Light & Strange vs. Connected (D6)

- relative correction $F_{1,2}^{\text{disc}} / F_{1,2}^{\text{conn}}$ from plateau averages $t_{\text{sep}} = 0.5 \div 0.74 \text{ fm}$, $Q^2 \leq 8 \text{ GeV}^2$
- D6 ensemble ($m_\pi = 170 \text{ MeV}$, $a = 0.092 \text{ fm}$), 727 configs \otimes 128 samples of $\langle N\bar{N} \rangle$
- partial noise cancellation between $L=U/D$ and S in proton & neutron

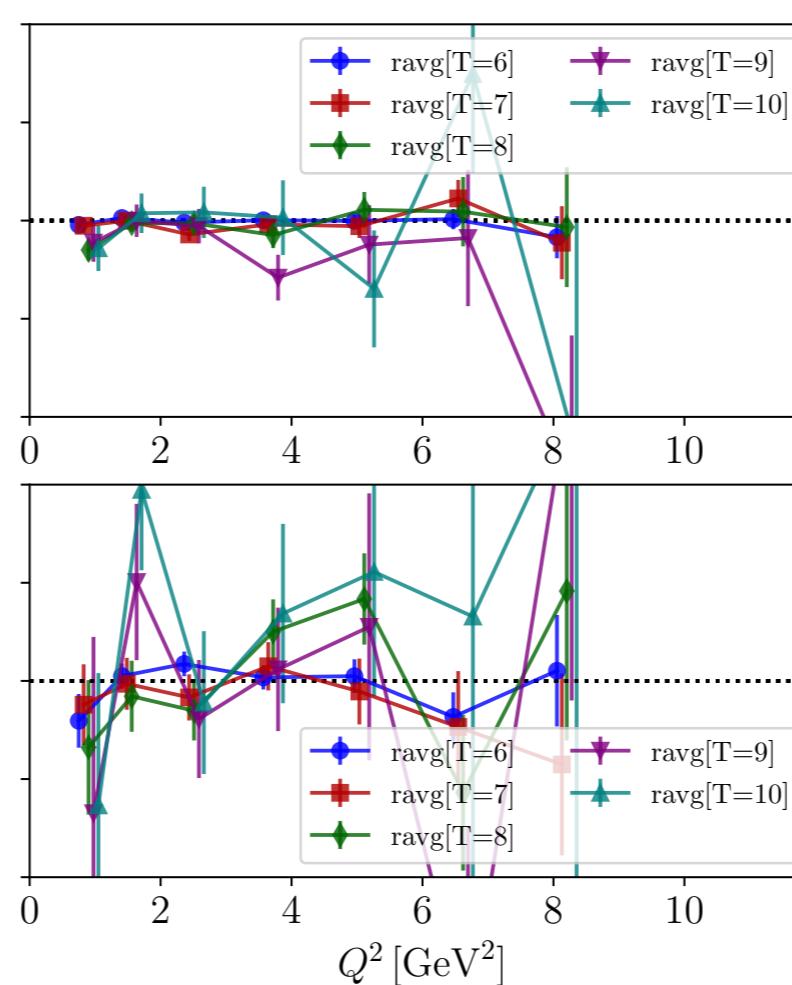
$$P = \frac{1}{3} [2U - D]_{\text{conn}} + \frac{1}{3} [L - S]_{\text{disc}}$$

$$N = \frac{1}{3} [2D - U]_{\text{conn}} + \frac{1}{3} [L - S]_{\text{disc}}$$

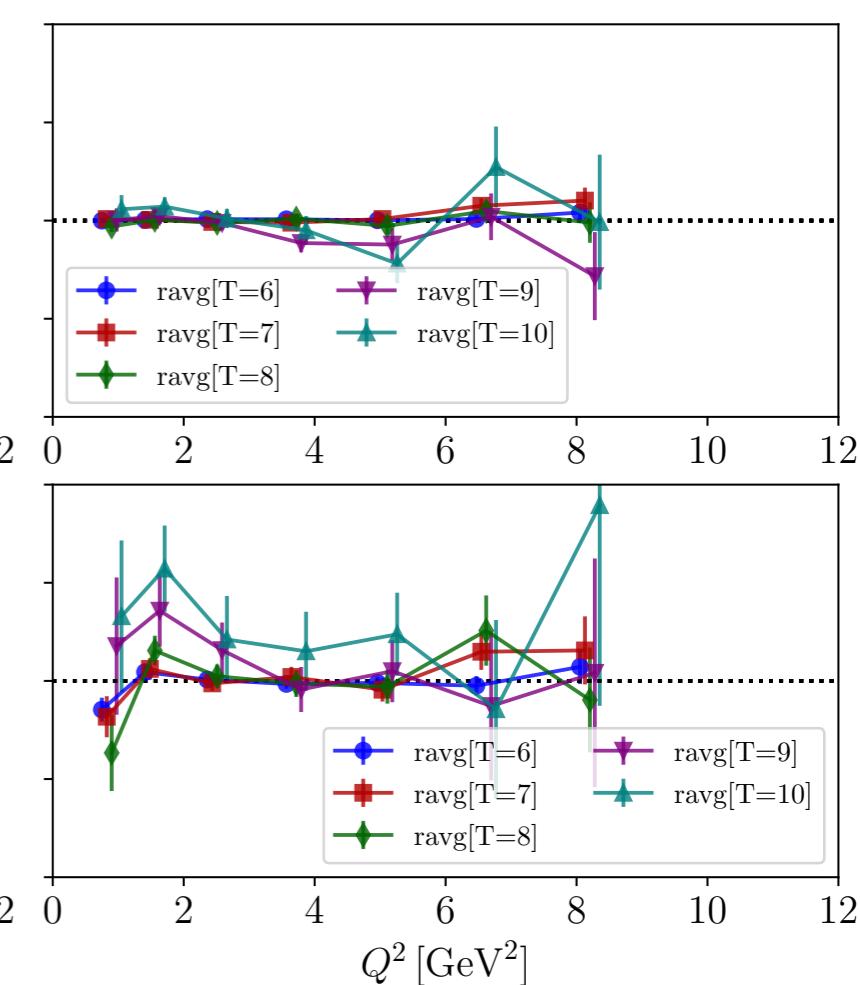
disconnected $L=U$ or D



disconnected S



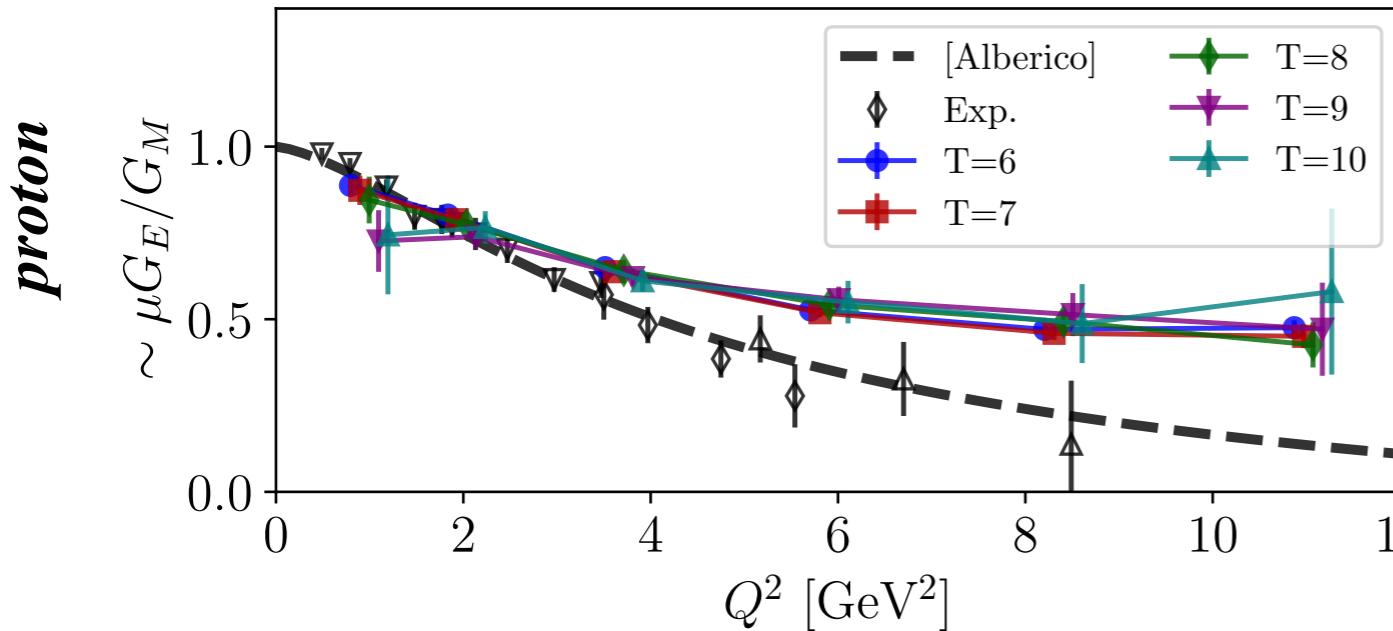
disconnected $(L-S)$



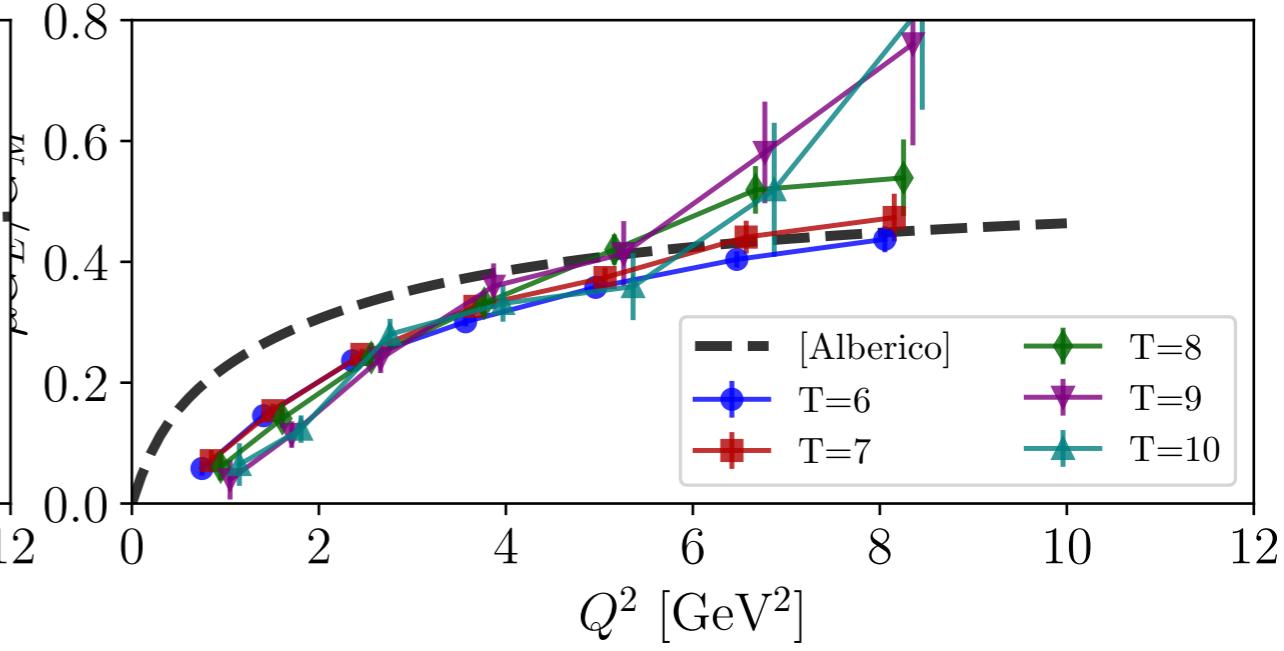
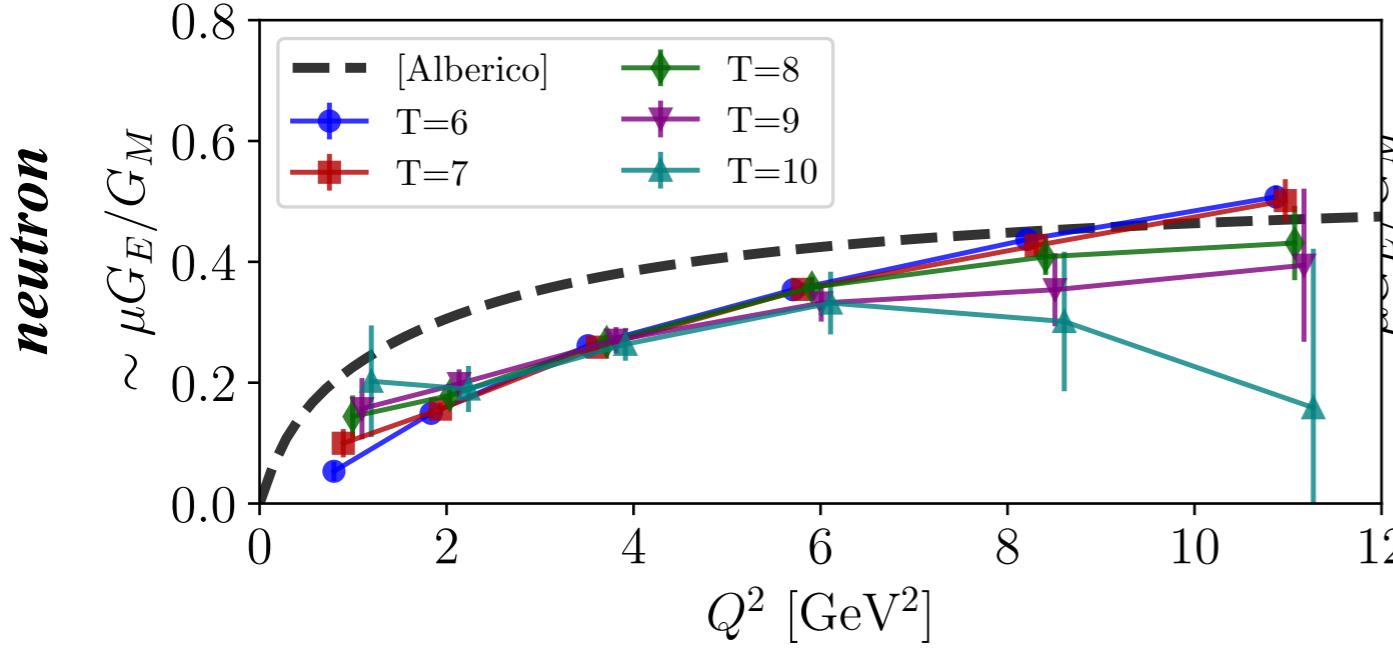
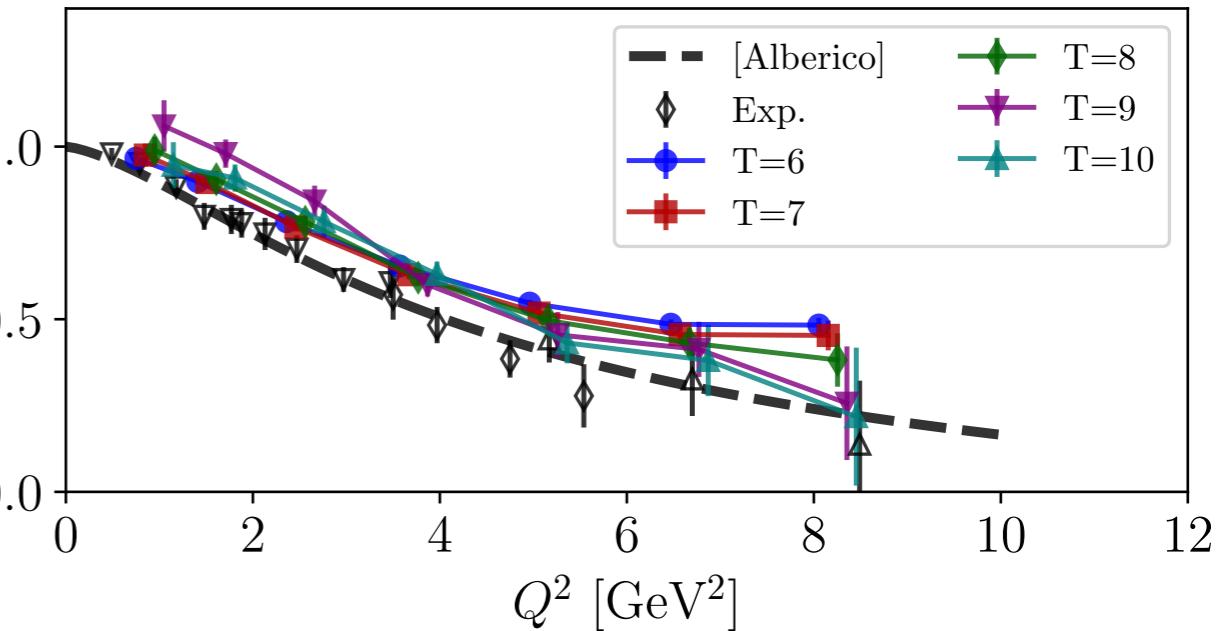
Proton&Neutron G_E/G_M : Connected-only

$$\left(\frac{\sinh \frac{\lambda' - \lambda}{2}}{\cosh \frac{\lambda' + \lambda}{2}} \right) \frac{\text{Re} \langle N_\uparrow(p'_x, T) J_t(T/2) \bar{N}_\uparrow(p_x, 0) \rangle}{\text{Re} \langle N_\uparrow(p'_x, T) J_y(T/2) \bar{N}_\uparrow(p_x, 0) \rangle} \xrightarrow{T \rightarrow \infty} G_E/G_M \quad \text{where} \quad \begin{pmatrix} p^{(\prime)} \\ E^{(\prime)} \end{pmatrix} = \begin{pmatrix} m_N \sinh \lambda^{(\prime)} \\ m_N \cosh \lambda^{(\prime)} \end{pmatrix}$$

D5($m\pi = 278$ MeV, $a = 0.094$ fm)



D6($m\pi = 166$ MeV, $a = 0.094$ fm)

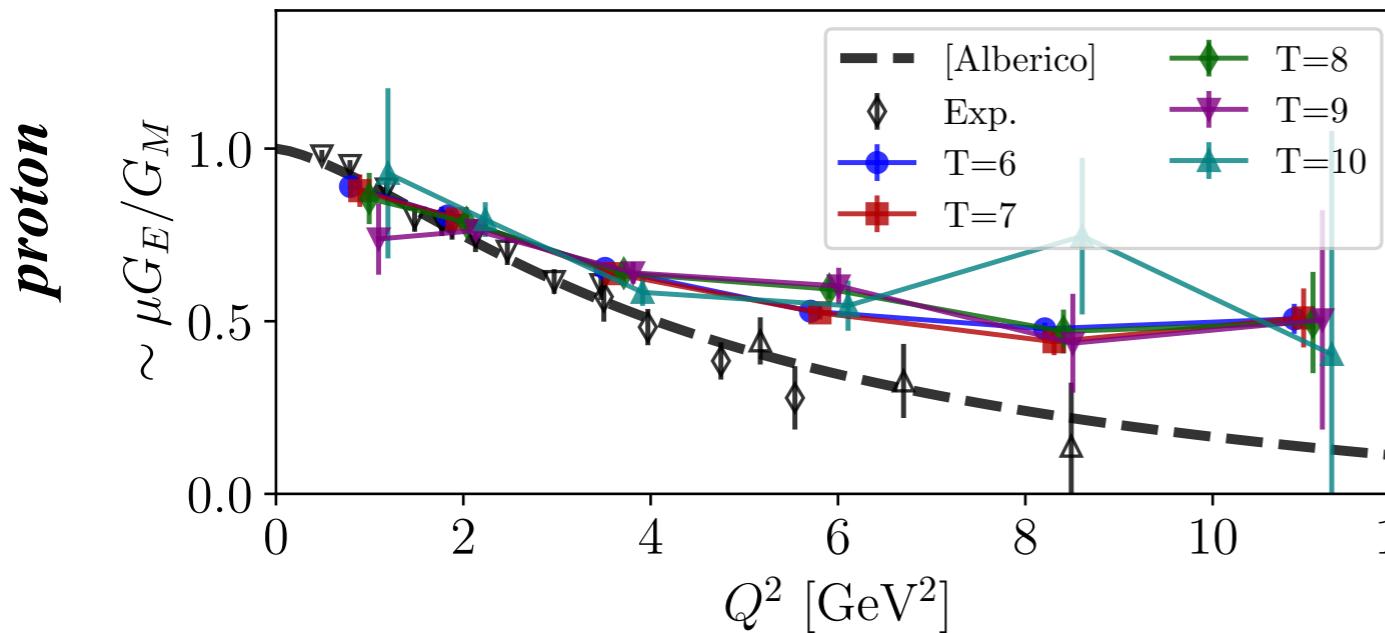


Proton&Neutron G_E/G_M : Connected+Disconnected

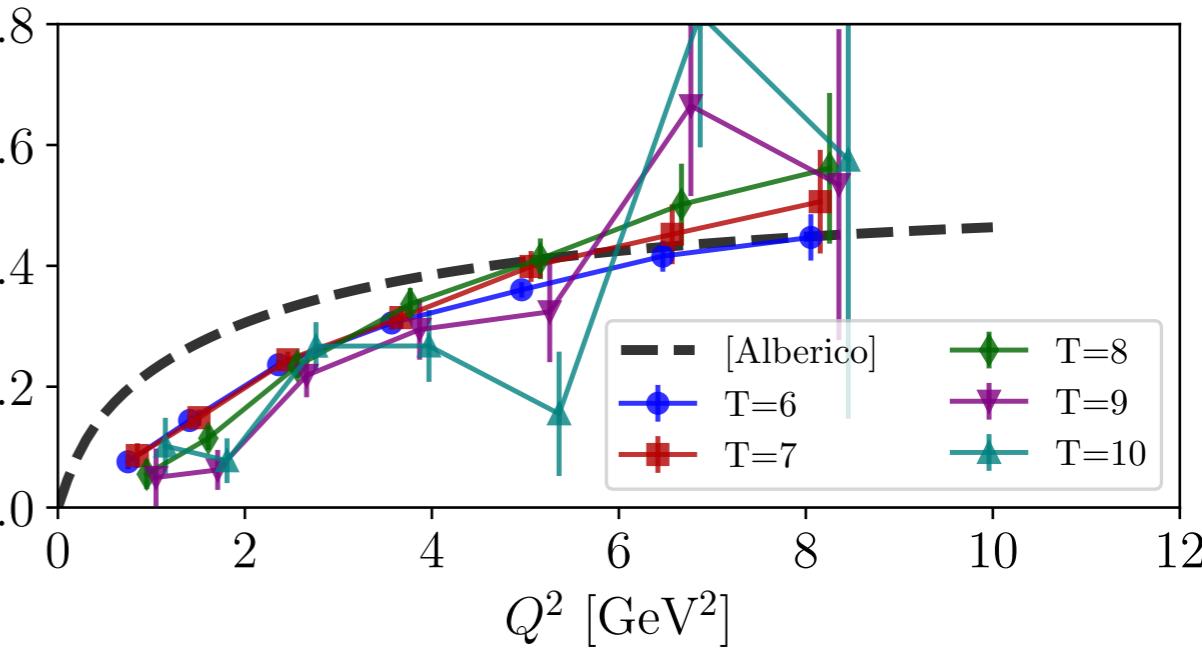
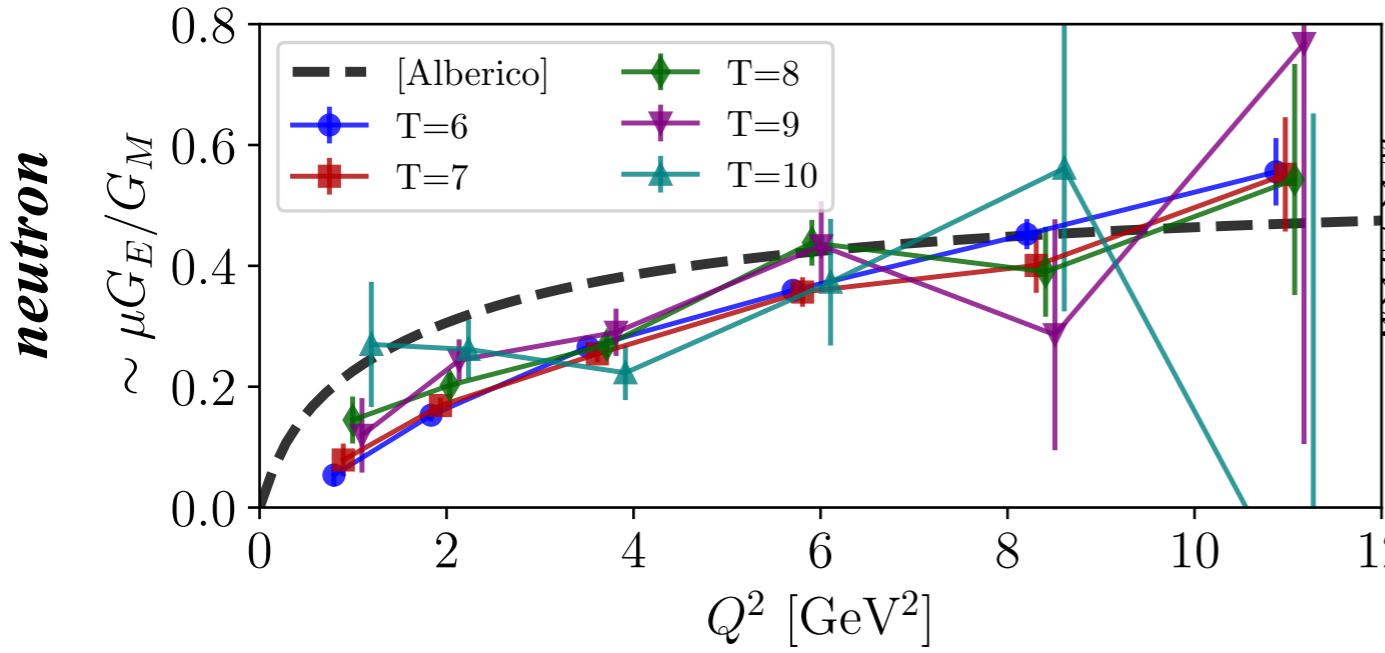
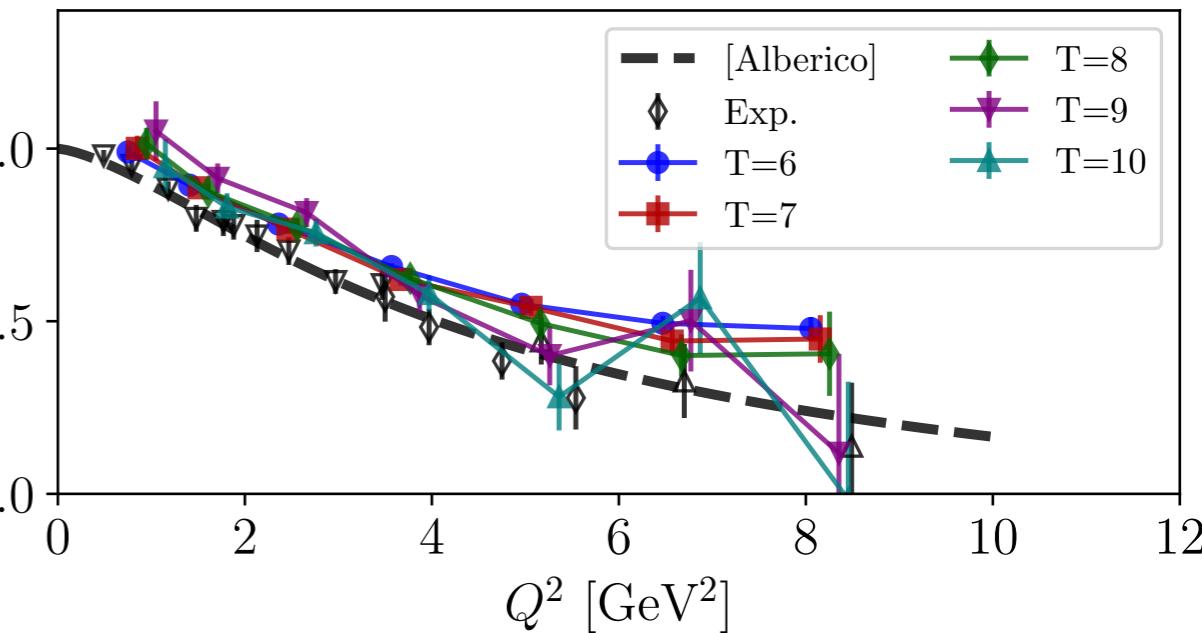
$$\left(\frac{\sinh \frac{\lambda' - \lambda}{2}}{\cosh \frac{\lambda' + \lambda}{2}} \right) \frac{\text{Re} \langle N_\uparrow(p'_x, T) J_t(T/2) \bar{N}_\uparrow(p_x, 0) \rangle}{\text{Re} \langle N_\uparrow(p'_x, T) J_y(T/2) \bar{N}_\uparrow(p_x, 0) \rangle} \underset{T \rightarrow \infty}{=} G_E/G_M$$

where $\begin{pmatrix} p^{(\prime)} \\ E^{(\prime)} \end{pmatrix} = \begin{pmatrix} m_N \sinh \lambda^{(\prime)} \\ m_N \cosh \lambda^{(\prime)} \end{pmatrix}$

D5($m\pi = 278$ MeV, $a = 0.094$ fm)



D6($m\pi = 166$ MeV, $a = 0.094$ fm)



Summary

- Preliminary results for high MC-statistics high-momentum form factors
 - up to $Q^2 \lesssim 10 \text{ GeV}^2$
 - two lattice spacings $a \gtrsim 0.07 \text{ fm}$
 - two pion masses $m\pi \gtrsim 170 \text{ MeV}$
- Quark-disconnected contributions (only $a \approx 0.09 \text{ fm}$)
 - consistent with zero but is a major source of uncertainty (esp. above $Q^2 \gtrsim 8 \text{ GeV}^2$)
 - little impact below $Q^2 \lesssim 6 \text{ GeV}^2$ (except in G_{Ep}/G_{Mp} and G_{En})
 - very unlikely to be significant, but longer t_{sep} , more stat. data needed
- Form factor results overshoot experimental data $x(2 \dots 2.5)$;
 F_2/F_1 , G_E/G_M ratios are in qualitative agreement with pheno. fits to experiment
 - Non-physical quarks masses?
 - Discretization? (less likely)
 - Excited states (most likely) — *large additional statistics will be needed for analysis*

*Important cross-check with experiments,
relevant for calculations of relativistic nucleon matrix elements
as well as TMDs, PDFs, DAs ...*

BACKUP