TMDPDFs extractions with DNNs

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Outline

- Motivation
- A brief introduction to TMDs
- Sivers asymmetry from SIDIS
- DNN Approach for SU(3)_{flavor}
- DNN Method Testing
- DNN Fits & Results for Sivers function
- Unpolarized TMDs extraction
- Summary and Outlook

Motivation

- TMDs: so far, model-dependent extractions \rightarrow assumptions, limitations and biases
- Information from data: Has the full potential of the data been taken into account?

Introducing this 'novel' method of extracting TMDs with Deep Neural Networks (DNNs)

First time in TMD extractions (published in 2023)

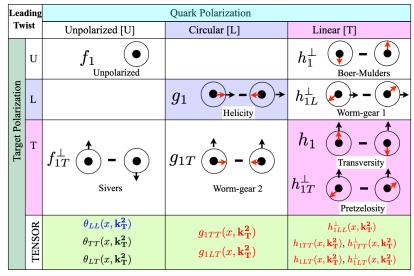
- Capacity to handle complex patterns, relationships in data with multi-D dependence.
- Data-driven
- Minimally biased → un-biased
- Uncertainty propagation (from data) using bootstrap method by generating 'replicas' (Statistical & Systematic uncertainties from the experimental data are combined in quadrature) Intended to be
- Recursive improvements to the DNN
- Systematic component can be quantified by a dedicated analysis

Exploratory & Instructional

TMDPDFs

$$\Phi(x, k_T; S) = \int \frac{d\xi^- d\xi_T}{(2\pi)^3} e^{ik.\xi} \langle P, S | \bar{\psi}(0) \mathcal{U}_{[0,\xi]} \psi(\xi) | P, S \rangle|_{\xi^+ = 0}$$

At leading-twist, the Quark correlator can be decomposed into 8 components (6 T - even and 2 T -odd terms)



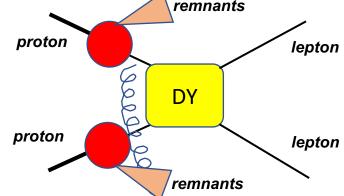
$$\Phi(x, k_T, P, S) = f_1(x, k_T^2) \frac{P}{2} + \frac{h_{1T}(x, k_T^2)}{4} \gamma_5 [\$_T, P] + \frac{S_L}{2} g_{1L}(x, k_T^2) \gamma_5 P + \frac{k_T \cdot S_T}{2M} g_{1T}(x, k_T^2) \gamma_5 P + \frac{S_L \cdot S_T}{2M} g_{1L}(x, k_T^2) \gamma_5 P + \frac{k_T \cdot S_T}{2M} g_{1L}(x, k_T^2) \gamma_5 P$$

T-even

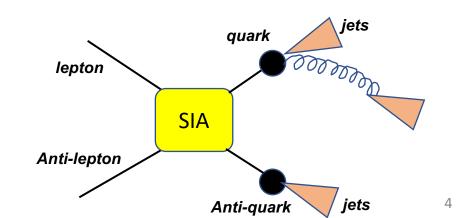
$$+ih_1^{\perp}(x,k_T^2)\frac{[k_T^{\prime},P^{\prime}]}{4M} - \frac{\epsilon_T^{k_TS_T}}{4M}f_{1T}^{\perp}(x,k_T^2)P^{\prime}$$

remnants

proton SIDIS Hadron + remnants

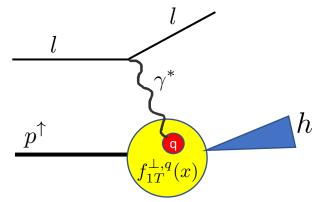


T-odd



TMDPDFs

Polarized Semi Inclusive DIS



$$\frac{d\sigma_{SIDIS}^{LO}}{dxdydzdp_T^2d\phi_hd\psi} = \left[\frac{\alpha}{xyQ^2}\frac{y^2}{2(1-\epsilon)}\left(1+\frac{y^2}{2x}\right)\right]$$

$$\times (F_{UU,T} + \epsilon F_{UU,L}) \left\{ 1 + \cos 2\phi_h \left(\epsilon A_{UU}^{\cos 2\phi_h} \right) \right\}$$

$$+S_T \left[\sin(\phi_h - \phi_s) \left(A_{UT}^{\sin(\phi_h - \phi_s)} \right) + \sin(\phi_h + \phi_s) \left(\epsilon A_{UT}^{\sin(\phi_h + \phi_s)} \right) + \sin(3\phi_h - \phi_s) \left(\epsilon A_{UT}^{\sin(3\phi_h - \phi_s)} \right) \right]$$

$$A_{UU}^{\cos 2\phi_h} \propto h_1^{\perp q} \circledast H_{1q}^{\perp h}$$

 $BM \circledast CF$

$$A_{UT}^{\sin(\phi_h - \phi_s)} \propto f_{1T}^{\perp q} \circledast D_{1q}^h$$

Sivers \circledast FF

$$A_{UT}^{\sin(\phi_h + \phi_s)} \propto h_1^q \circledast H_{1q}^{\perp h} \quad \text{Transv} \circledast \text{CF}$$

$$A_{UT}^{\sin(3\phi_h - \phi_s)} \propto h_{1T}^{\perp q} \circledast H_{1q}^{\perp h} \quad \text{Pretz} \circledast \text{CF}$$

$$\begin{aligned} h_1^{\perp q} \Big|_{SIDIS} &= -h_1^{\perp q} \Big|_{DY} \\ f_{1T}^{\perp q} \Big|_{SIDIS} &= -f_{1T}^{\perp q} \Big|_{DY} \end{aligned}$$

* For these two processes TMD factorization is proven

$$\frac{d\sigma^{LO}}{d\Omega} = \frac{\alpha_{em}^2}{Fq} F_v^1 \left\{ 1 + \cos^2\theta + \sin^2\theta \cos 2\phi_{CS} A_U^{\cos 2\phi_{CS}} \right\}$$

$$+S_T \left[\left(1 + \cos^2 \theta \right) \sin \phi_s A_T^{\sin \phi_s} + \sin^2 \theta \left(\sin(2\phi_{CS} + \phi_s) A_T^{\sin(2\phi_{CS} + \phi_s)} \right) \right]$$

$$+\sin(2\phi_{CS}-\phi_s)A_T^{\sin(2\phi_{CS}-\phi_s)}$$

$$A_T^{\cos 2\phi_{CS}} \propto h_1^{\perp q} \circledast h_1^{\perp q}$$

$$A_T^{\sin\phi_s}\propto f_1^q \circledast f_{1T}^{\perp q} \quad ext{PDF} \circledast ext{Sivers}$$

$$A_T^{\sin(2\phi_{CS}-\phi_S)} \propto h_1^{\perp q} \circledast h_1^q$$

$$A_T^{\sin(2\phi_{CS}+\phi_S)} \propto h_1^{\perp q} \circledast h_{1T}^{\perp q}$$

$$BM \circledast BM$$

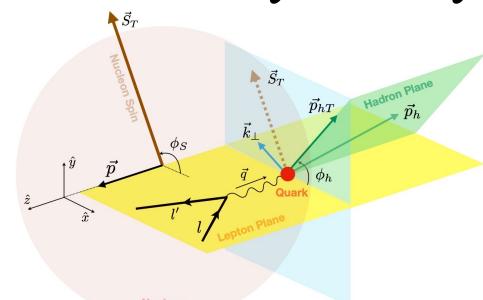
Polarized DY

$$BM \circledast Transv$$



Sivers Asymmetry from SIDIS $\frac{d^5\sigma^{lp\to lhX}}{dxdQ^2dzd^2p_\perp} = \sum_{q} e_q^2 \int d^2\mathbf{k}_\perp \, \left(\frac{2\pi\alpha^2}{x^2s^2}\frac{\hat{s}^2 + \hat{u}^2}{Q^4}\right)$

$$\frac{d^5 \sigma^{lp \to lhX}}{dx dQ^2 dz d^2 p_{\perp}} = \sum_{q} e_q^2 \int d^2 \mathbf{k}_{\perp} \left(\frac{2\pi\alpha^2}{x^2 s^2} \frac{\hat{s}^2 + \hat{u}^2}{Q^4} \right)$$
$$\times \hat{f}_{q/p\uparrow}(x, k_{\perp}) D_{h/q}(z, p_{\perp}) + \mathcal{O}(k_{\perp}/Q) = 0$$



$$\hat{f}_{q/p^{\uparrow}}(x,k_{\perp}) = f_{q/p}(x,k_{\perp}) + \frac{1}{2} \Delta^{N} f_{q/p^{\uparrow}}(x,k_{\perp}) \vec{S}_{T} \cdot (\hat{p} \times \hat{k}_{\perp})$$

$$\Delta^{N} f_{q/p\uparrow}(x, k_{\perp}) = 2\mathcal{N}_{q}(x)h(k_{\perp})f_{q/p}(x, k_{\perp})$$

Anselmino et al. (2017)

Single Spin Asymmetry (Sivers Asymmetry)

$$\langle p_{\perp}^2 \rangle = 0.12 \pm 0.01 \text{ GeV}^2$$

 $\langle k_{\perp}^2 \rangle = 0.57 \pm 0.08 \text{ GeV}^2$

$$A_{UT}^{\sin(\phi_h - \phi_S)}(x, y, z, p_{hT}) = \frac{d\sigma^{l \uparrow p \to hlX} - d\sigma^{l \downarrow p \to lhX}}{d\sigma^{l \uparrow p \to hlX} + d\sigma^{l \downarrow p \to hlX}} \equiv \frac{d\sigma \uparrow - d\sigma \downarrow}{d\sigma \uparrow + d\sigma \downarrow}$$

$$\mathcal{A}_{0}(z, p_{hT}, m_{1})$$

$$= \frac{\sqrt{2e}zp_{hT}}{m_{1}} \frac{[z^{2}\langle k_{\perp}^{2}\rangle + \langle p_{\perp}^{2}\rangle]\langle k_{S}^{2}\rangle^{2}}{[z^{2}\langle k_{S}^{2}\rangle + \langle p_{\perp}^{2}\rangle]^{2}\langle k_{\perp}^{2}\rangle}$$

$$\times \exp\left[-\frac{p_{hT}^{2}z^{2}\left(\langle k_{S}^{2}\rangle - \langle k_{\perp}^{2}\rangle\right)}{(z^{2}\langle k_{S}^{2}\rangle + \langle p_{\perp}^{2}\rangle)\left(z^{2}\langle k_{\perp}^{2}\rangle + \langle p_{\perp}^{2}\rangle\right)}\right]$$

$$\langle k_{S}^{2}\rangle = \frac{m_{1}\langle k_{\perp}^{2}\rangle}{m_{2}^{2} + \langle k_{\perp}^{2}\rangle}$$

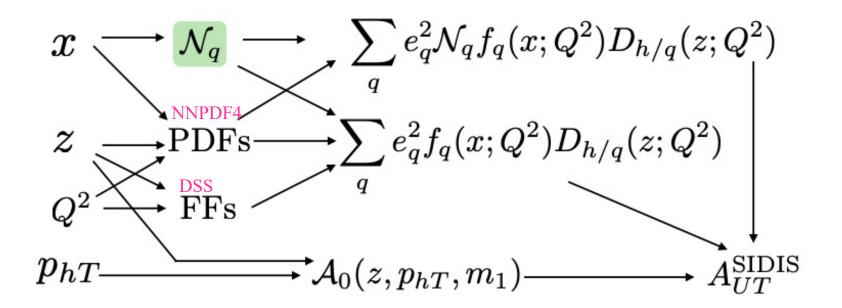
$$A_{UT}^{\sin(\phi_h - \phi_S)}(x, z, p_{hT}) = \mathcal{A}_0(z, p_{hT}, m_1) \left(\frac{\sum_q \mathcal{N}_q(x) e_q^2 f_q(x) D_{h/q}(z)}{\sum_q e_q^2 f_q(x) D_{h/q}(z)} \right)$$

$$N_q(x) = N_q x^{\alpha_q} (1 - x)^{\beta_q} \frac{(\alpha_q + \beta_q)^{(\alpha_q + \beta_q)}}{\alpha_q^{\alpha_q} \beta_q^{\beta_q}}$$

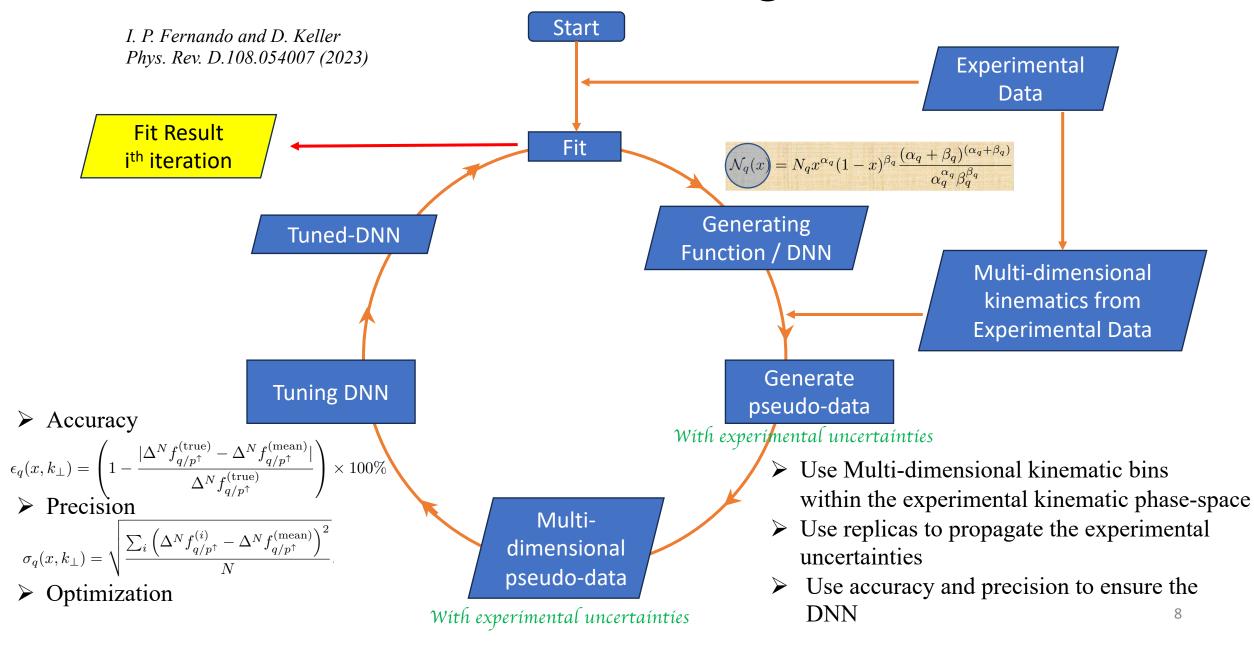
$$\widehat{\mathcal{N}_{\bar{q}}(x)} = N_{\bar{q}}$$

DNN Approach

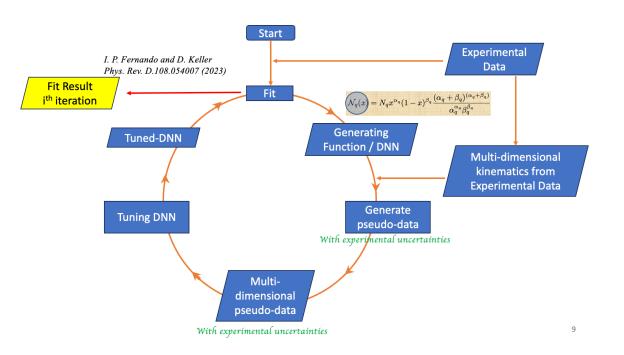
$$A_{UT}^{\sin(\phi_h - \phi_S)}(x, z, p_{hT}) = \mathcal{A}_0(z, p_{hT}, m_1) \left(\frac{\sum_q \mathcal{N}_q(x) e_q^2 f_q(x) D_{h/q}(z)}{\sum_q e_q^2 f_q(x) D_{h/q}(z)} \right)$$

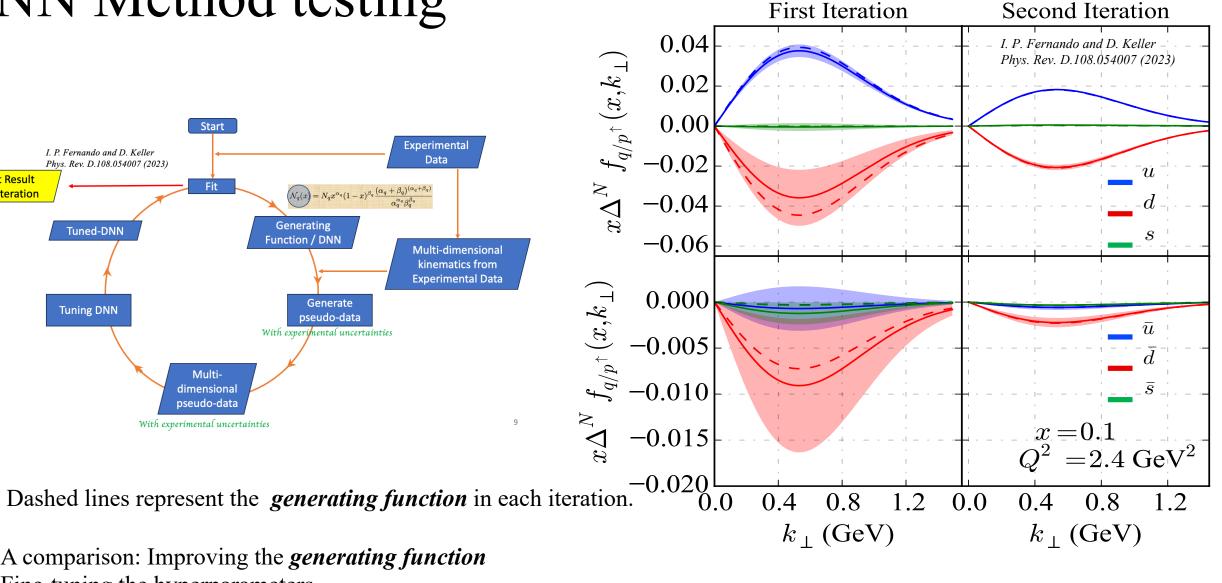


The DNN Method for extracting Sivers function



DNN Method testing





- A comparison: Improving the *generating function* Fine-tuning the hyperparameters
- Solid-lines and the band represent the mean and 68% CL with 1000 replicas of the DNN model.

Data Selection

Proton DNN
model

Deuteron DNN model

Projections from Deuteron DNN model

$$\Delta^{N} f_{q/p^{\uparrow}}(x, k_{\perp})\big|_{\text{SIDIS}} = -\Delta^{N} f_{q/p^{\uparrow}}(x, k_{\perp})\big|_{\text{DY}}$$

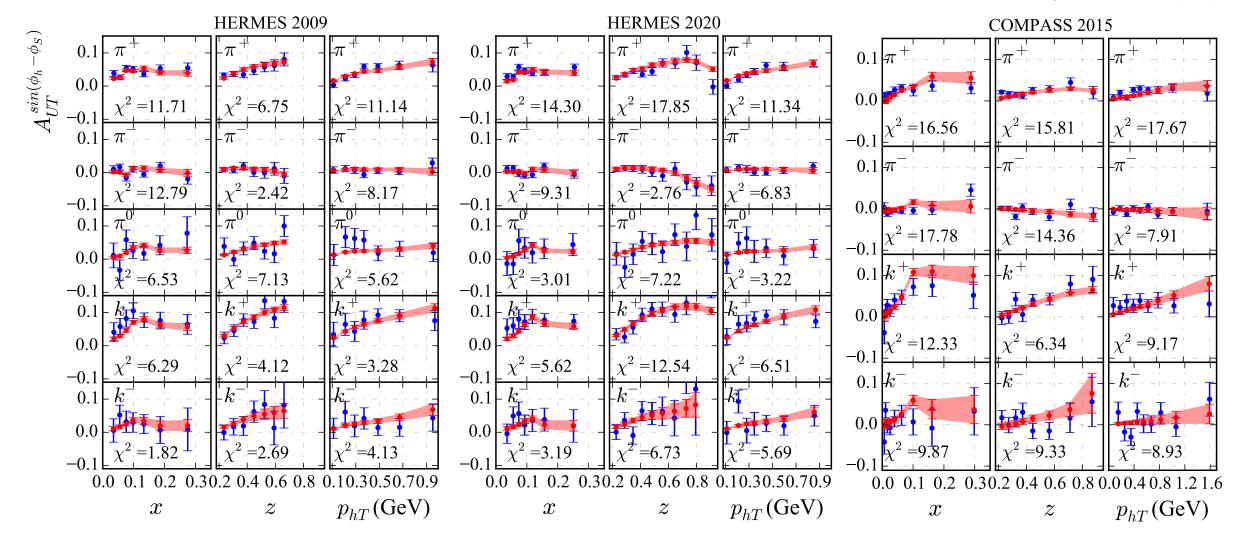
	Dataset	Kinematic	Reaction	Data	
•		coverage		points	
4	HERMES2009	0.023 < x < 0.4	$p^{\uparrow} + \gamma^* \to \pi^+$	21	
	(SIDIS)	0.2 < z < 0.7	$p^{\uparrow} + \gamma^* \to \pi^-$	21	
	[53]	$0.1 < p_{hT} < 0.9$	$p^{\uparrow} + \gamma^* \to \pi^0$	21	
		$Q^2 > 1 \text{ GeV}^2$	$p^{\uparrow} + \gamma^* \to K^+$	21	
			$p^{\uparrow} + \gamma^* \to K^-$	21	
	HERMES2020	0.023 < x < 0.6	$p^{\uparrow} + \gamma^* \to \pi^+$	27, 64	
	(SIDIS)	0.2 < z < 0.7	$p^{\uparrow} + \gamma^* \to \pi^-$	27, 64	
	[55]	$0.1 < p_{hT} < 0.9$	$p^{\uparrow} + \gamma^* \to \pi^0$	27	
		$Q^2 > 1 \text{ GeV}^2$	$p^{\uparrow} + \gamma^* \to K^+$	27, 64	
			$p^{\uparrow} + \gamma^* \to K^-$	27, 64	
	COMPASS2015	0.006 < x < 0.28	$p^{\uparrow} + \gamma^* \to \pi^+$	26	
	(SIDIS)	0.2 < z < 0.8	$p^{\uparrow} + \gamma^* \to \pi^-$	26	
	[54]	$0.15 < p_{hT} < 1.5$	$p^{\uparrow} + \gamma^* \to K^+$	26	
•		$Q^2 > 1 \text{ GeV}^2$	$p^{\uparrow} + \gamma^* \to K^-$	26	
1	COMPASS2009	0.006 < x < 0.28	$d^{\uparrow} + \gamma^* \to \pi^+$	26	
	(SIDIS)	0.2 < z < 0.8	$d^{\uparrow} + \gamma^* \to \pi^-$	26	
	[49]	$0.15 < p_{hT} < 1.5$	$d^{\uparrow} + \gamma^* \to K^+$	26	
		$Q^2 > 1 \text{ GeV}^2$	$d^{\uparrow} + \gamma^* \to K^-$	26	
	JLAB2011	0.156 < x < 0.396	$^{3}He^{\uparrow} + \gamma^{*} \rightarrow \pi^{+}$	4	
	(SIDIS) [52]	0.50 < z < 0.58	$^{3}He^{\uparrow} + \gamma^{*} \rightarrow \pi^{-}$	4	
	(/ []	$0.24 < p_{hT} < 0.43$	· •		
7		$1.3 < Q^2 < 2.7$			
	COMPASS2017	$0.1 < x_N < 0.25$	$p^{\uparrow} + \pi^- \to l^+ l^- X$	15	
	(DY) [50]	$0.3 < x_{\pi} < 0.7$			
		$4.3 < Q_M < 8.5$			
		$0.6 < q_T < 1.9$			

HERMES2020 3D binned data

Projections from Proton DNN model

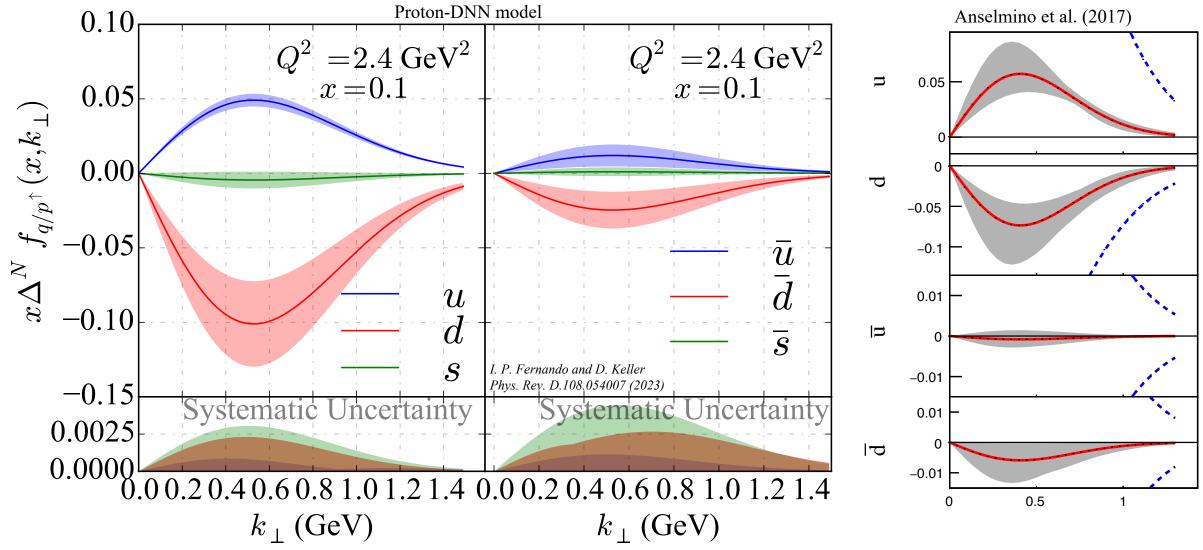
Proton DNN Fit Results

I. P. Fernando and D. Keller Phys. Rev. D.108.054007 (2023)



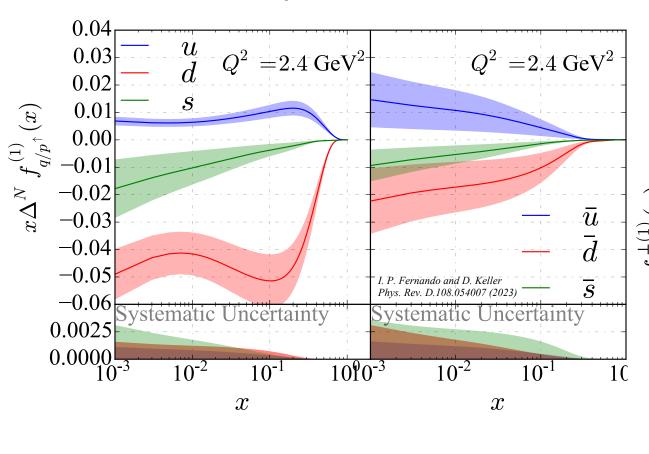
- ➤ All data points are well-described by the proton-DNN model.
- No kinematic cuts were implemented.

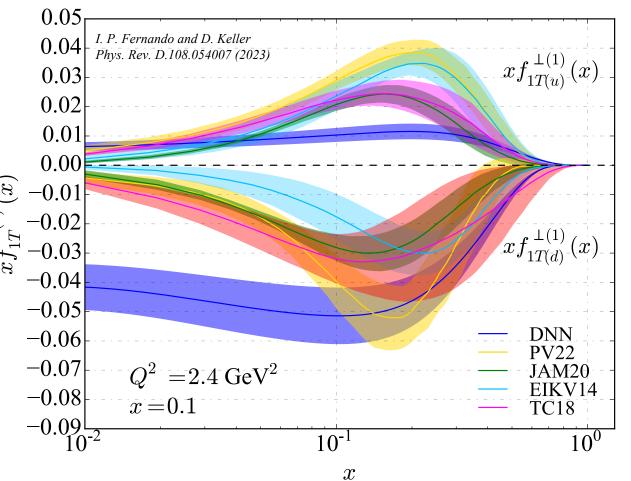
Sivers functions from the "Proton" DNN Model



Sivers 1st moments from the "Proton" Model

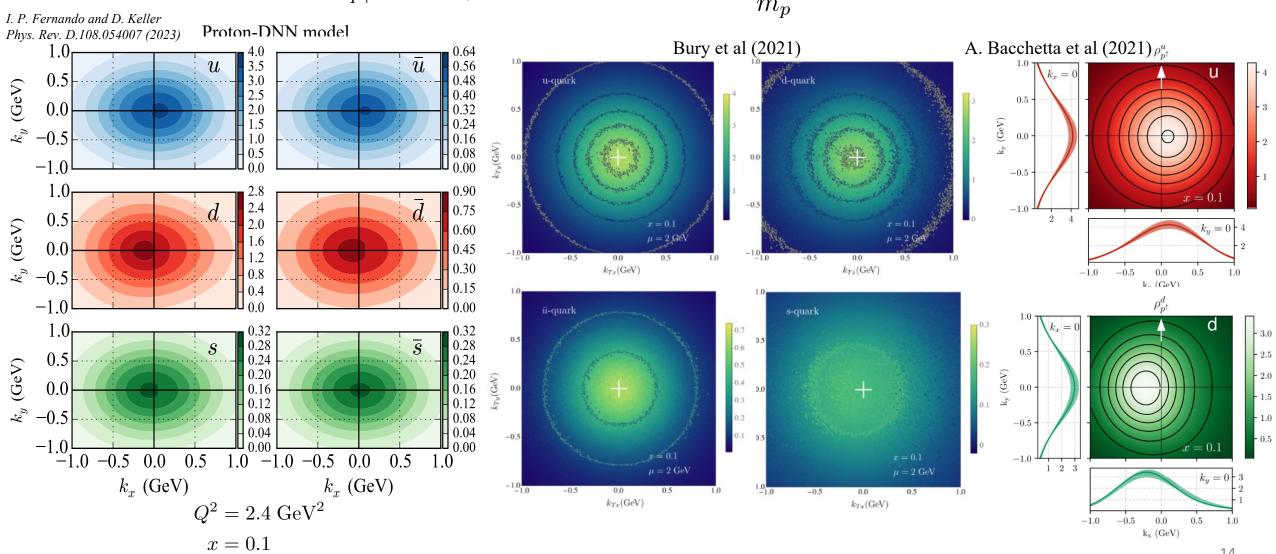
$$\Delta^{N} f_{q/p\uparrow}^{(1)}(x) = \int d^{2}k_{\perp} \frac{k_{\perp}}{4m_{p}} \Delta^{N} f_{q/p\uparrow}(x, k_{\perp}) = -f_{1T}^{\perp(1)q}(x)$$





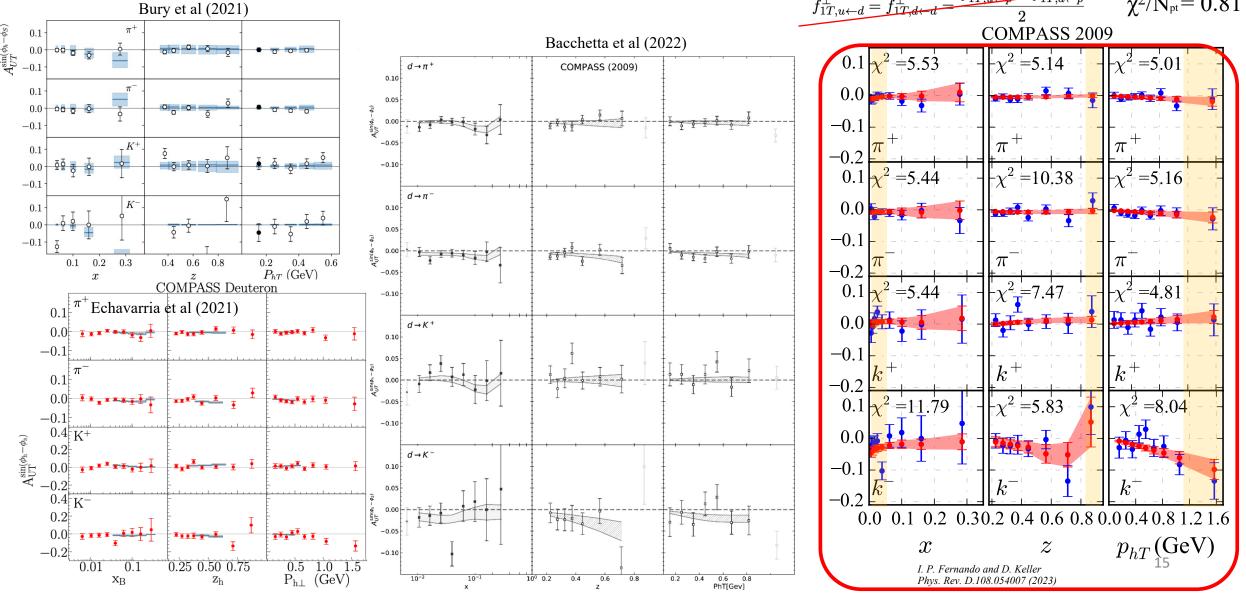
3D Tomography from the "Proton" DNN Model

$$\rho_{p\uparrow}^{a}(x, k_x, k_y; Q^2) = f_1^{a}(x, k_\perp^2; Q^2) - \frac{k_x}{m_p} f_{1T}^{\perp a}(x, k_\perp^2; Q^2)$$

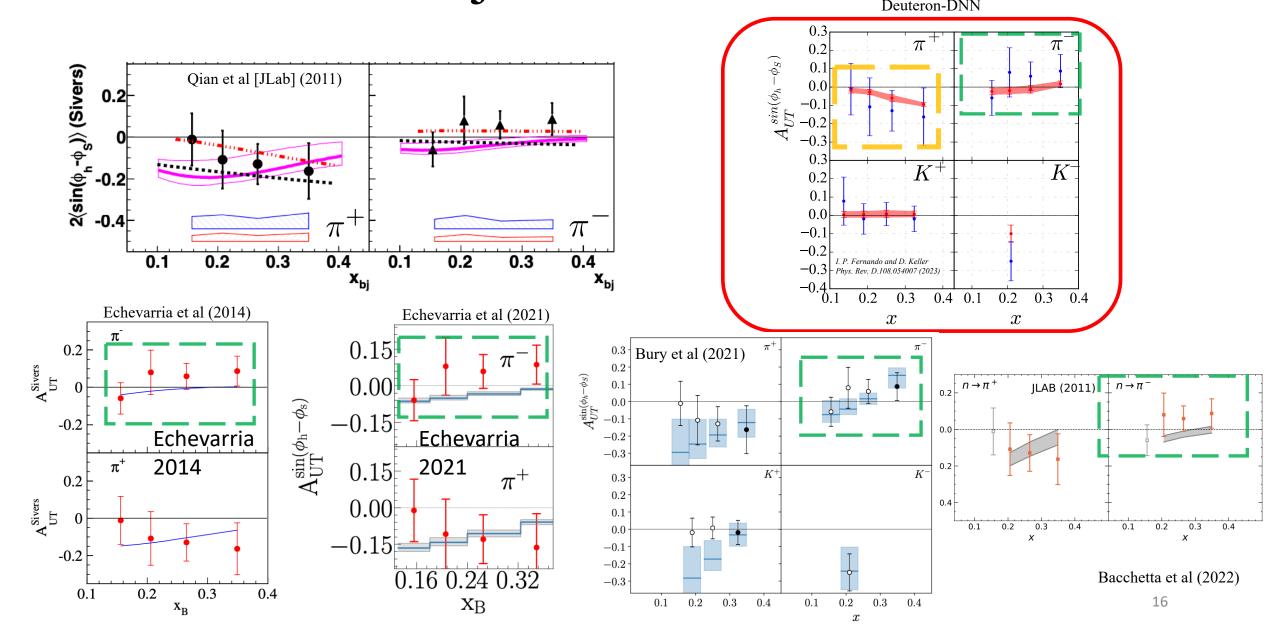


Deuteron DNN Fit Results

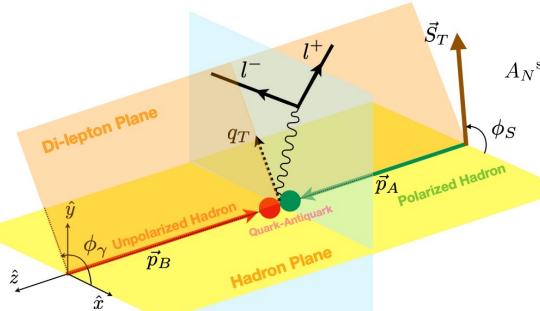
- No kinematic cuts are applied Deuteron-DNN model can describe data reasonably well
- No iso-spin symmetry conditions are applied Bury et al (2021)



Deuteron DNN Projections for JLab Kinematics



DNN Model Projections: DY



 $x_1 \longrightarrow \mathcal{N}_q \longrightarrow \sum_{q} \frac{e_q^2}{x_1 + x_2} \mathcal{N}_q(x_{1,2}) f_{q/A}(x_1) f_{\bar{q}/B}(x_2)$ $\sum_{q} \frac{e_q^2}{x_1 + x_2} f_{q/A}(x_1) f_{\bar{q}/B}(x_2)$ $Q_M \longrightarrow \mathcal{B}_0(q_T, m_1) \longrightarrow A_N^{\mathrm{DY}}$

Anselmino et al. (2017)

$$A_{N}^{\sin(\phi_{\gamma}-\phi_{S})}(x_{F}, M, q_{T}) = \mathcal{B}_{0}(q_{T}, m_{1}) \frac{\sum_{q} \frac{e_{q}^{2}}{x_{1}+x_{2}} \mathcal{N}_{q}(x_{1}) f_{q/A}(x_{1}) f_{\bar{q}/B}(x_{2})}{\sum_{q} \frac{e_{q}^{2}}{x_{1}+x_{2}} f_{q/A}(x_{1}) f_{\bar{q}/B}(x_{2})}$$

$$\phi_{S}$$

$$\mathcal{B}_{0}(q_{T}, m_{1}) = \frac{q_{T} \sqrt{2e}}{m_{1}} \frac{Y_{1}(q_{T}, k_{S}, k_{\perp 2})}{Y_{2}(q_{T}, k_{\perp 1}, k_{\perp 2})}$$

$$Y_{1}(q_{T}, k_{S}, k_{\perp 2}) = \left(\frac{\langle k_{S}^{2} \rangle^{2}}{\langle k_{\perp 2}^{2} \rangle \left(\langle k_{S}^{2} \rangle + \langle k_{\perp 2}^{2} \rangle\right)^{2}}\right) \times \exp\left(\frac{-q_{T}^{2}}{\langle k_{S}^{2} \rangle + \langle k_{\perp 2}^{2} \rangle}\right)$$

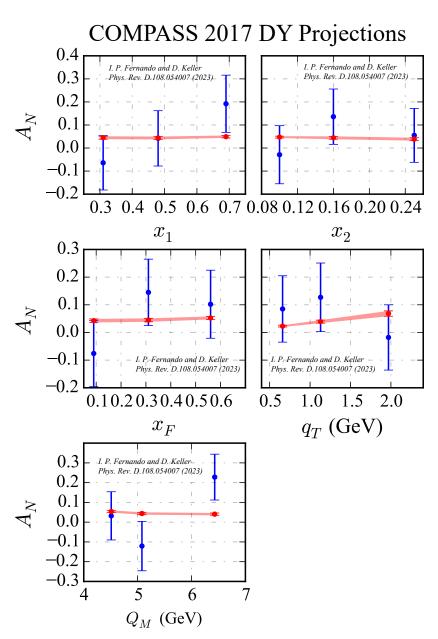
$$Y_{2}(q_{T}, k_{\perp 1}, k_{\perp 2}) = \left(\frac{1}{\langle k_{\perp 1}^{2} \rangle + \langle k_{\perp 2}^{2} \rangle}\right) \times \exp\left(\frac{-q_{T}^{2}}{\langle k_{\perp 1}^{2} \rangle + \langle k_{\perp 2}^{2} \rangle}\right)$$

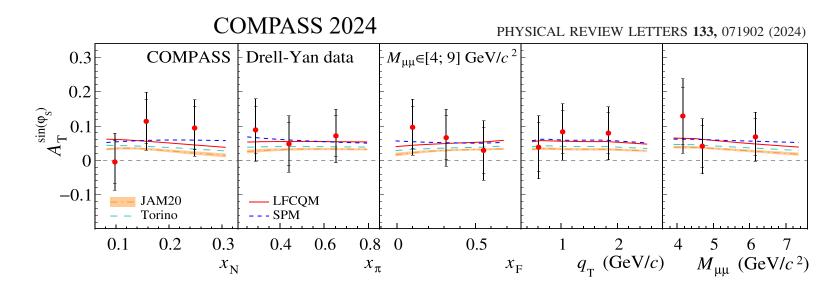
$$\frac{1}{\langle k_{S}^{2} \rangle} = \frac{1}{m_{1}^{2}} + \frac{1}{\langle k_{\perp 1}^{2} \rangle}$$

$$\Delta^{N} f_{q/p^{\uparrow}}(x, k_{\perp})\big|_{\text{SIDIS}} = -\Delta^{N} f_{q/p^{\uparrow}}(x, k_{\perp})\big|_{\text{DY}}$$

 $\langle k_{\perp 1}^2 \rangle = \langle k_{\perp 2}^2 \rangle = \langle k_{\perp}^2 \rangle = 0.25 \text{ GeV}^2$

DNN (Proton) Model Projections: DY

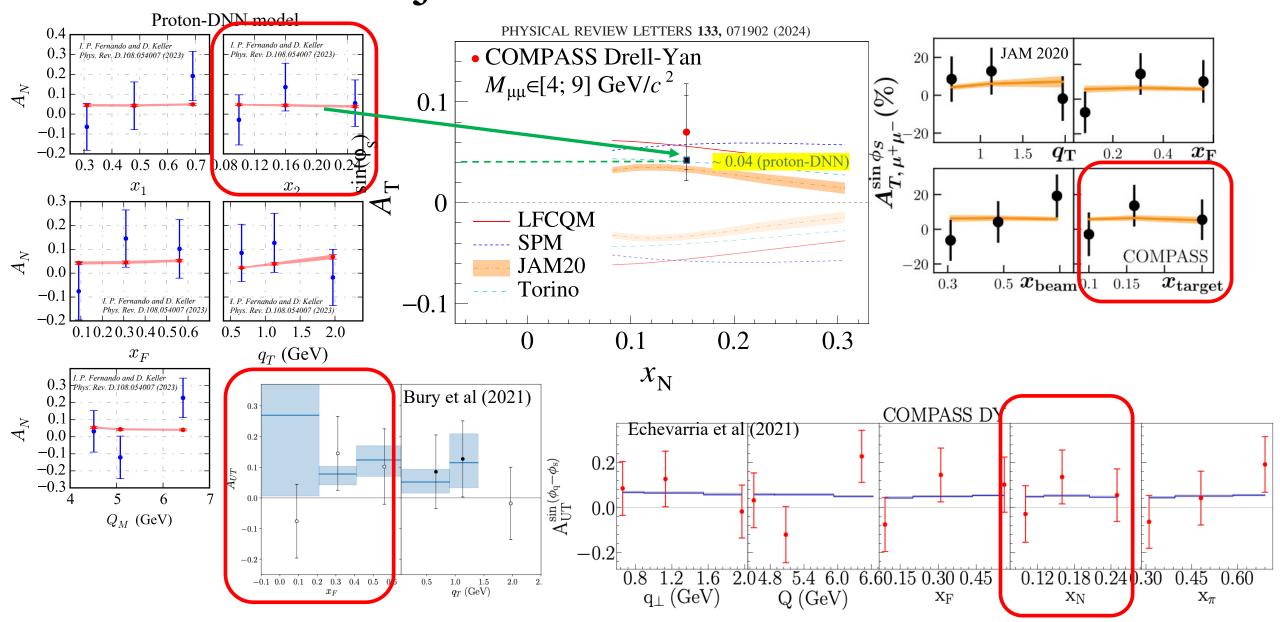




Note: These proton-DNN projections based on the assuming the sign-change of the Sivers functions

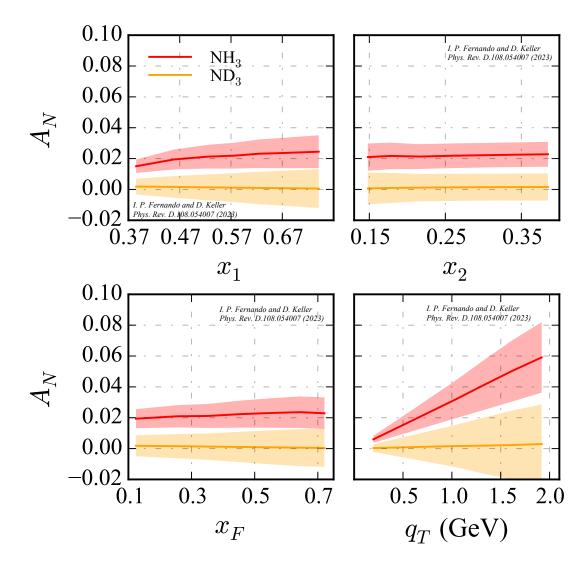
DNN Model Projections: DY

In Comparison with COMPASS 2024 Final



DNN Model Projections: DY @ SpinQuest

DNN Models



- ➤ SpinQuest (E1039) experiment at Fermilab is aiming to extract the Sivers function for the light-sea quarks.
- ➤ Unpolarized 120 GeV proton beam with polarized proton and deuteron targets (separately).
- ➤ Proton-DNN model predictions (Red)
 Deuteron-DNN model predictions
 (Orange)

Systematic Studies: data cuts

$$\begin{split} W^{\mu\nu} &= \sum_{f} |\mathcal{H}_{f}(Q^{2},\mu)|^{\mu\nu} \\ &\times \int d^{2}k_{\perp}d^{2}p_{\perp}\delta^{(2)}(z_{h}k_{\perp} + p_{\perp} - p_{hT}) \\ &\times F_{f/N^{\uparrow}}(x,z_{h}k_{\perp},S;\mu,\zeta_{F})D_{h/f}(z_{h},p_{\perp};\mu,\zeta_{D}) \\ &+ Y(p_{hT},Q^{2}), \end{split}$$

Examples:

- 1. Bury et al JHEP 05 (2021) 151 Q > 2 GeV $\delta = p_{hT}/zQ \le 0.3$
- 2. Echevarria et al JHEP 01 (2021) 126 $q_T/Q < 0.75$
- 3. JAM2020

$$Q^2 > 1.63 \text{ GeV}^2$$
, $0.2 < z < 0.6$, $0.2 < p_{hT} < 0.9 \text{ GeV}$

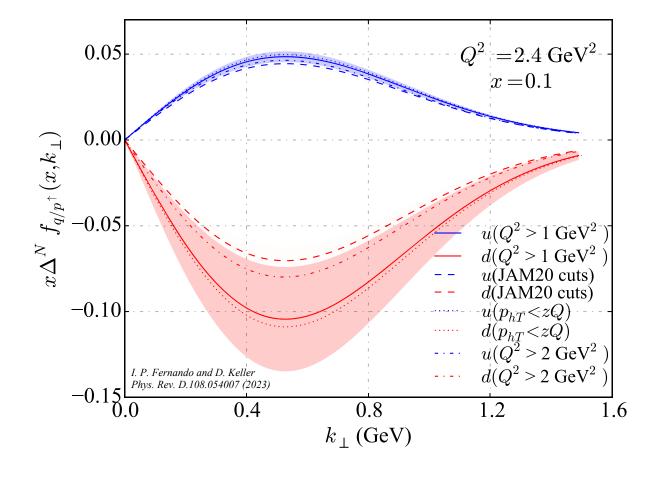
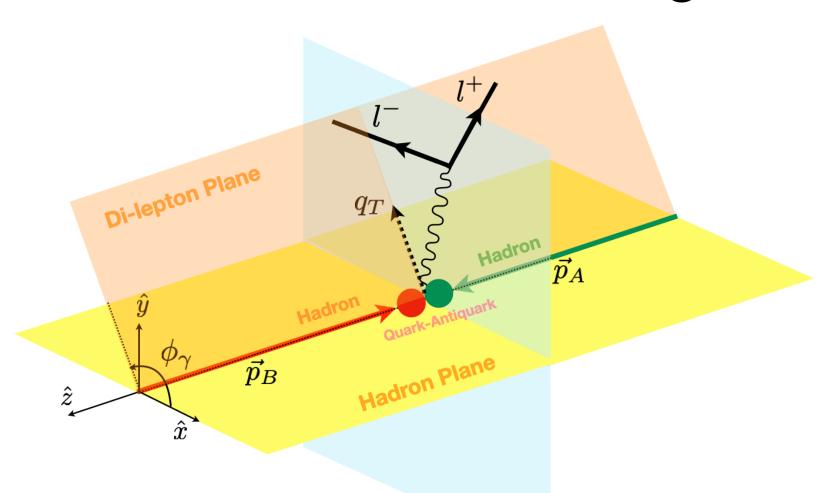


FIG. 17. Solid lines with light band represent the u (in blue), d (in red) Sivers functions using the cut $Q^2 > 1$ GeV². These resulting DNN models made from the cuts from all tests are also shown.

Unpolarized TMDs with fixed target DY data



Unpolarized TMDs: Motivation

• Can we obtain TMDs directly in k_T space instead of bT-space? We are interested in both. We think having a framework with k_T and complimentary to b_T space

$$\frac{d\sigma}{dq_T dQ_M dy} = \frac{16\pi^2 \alpha^2}{9Q_M^3} q_T \sum_q e_q^2 \mathcal{H}^{DY} \int d^2 \mathbf{k}_{aT} d^2 \mathbf{k}_{bT} f_{q/a}(x_a, \mathbf{k}_{aT}; \mu_Q, Q^2) f_{\bar{q}/b}(x_b, \mathbf{k}_{bT}; \mu_Q, Q^2) \delta^{(2)}(\mathbf{q}_T - \mathbf{k}_{aT} - \mathbf{k}_{bT}) + (x_a \leftrightarrow x_b)$$

- Number of data points vs cuts: Our goal is to use as much data as available while TMD factorization is still valid.
- Different types of parameterizations attempted in the literature: So, we want to explore a purely numerical approach utilizing DNNs

Unpolarized TMDs: Setup

Introducing (exploration of) a separable form

$$f_q(x, k_\perp, Q_M) = f_q(x, Q_M) s(x, k_\perp) \mathcal{B}(Q_M)$$

Colinear-PDFs with DGLAP Evolution

Transvers Momentum Kernal (DNN)

Transvers Momentum Evolution Kernal (DNN)

$$\hat{f}_{1}^{a}(x, \boldsymbol{b}_{T}^{2}; \mu_{f}, \zeta_{f}) = \hat{f}_{1}^{a}(x, \boldsymbol{b}_{T}^{2}; \mu_{i}, \zeta_{i}) \exp \left\{ \int_{\mu_{i}}^{\mu_{f}} \frac{d\mu}{\mu} \gamma(\mu, \zeta_{f}) \right\} \left(\frac{\zeta_{f}}{\zeta_{i}} \right)^{K(|\boldsymbol{b}_{T}|, \mu_{i})/2}$$

$$\frac{d\sigma}{dq_T dQ_M dy} = \frac{16\pi^2 \alpha^2}{9Q_M^3} q_T \sum_a e_q^2 \mathcal{H}^{DY} \int d^2 \mathbf{k}_{aT} d^2 \mathbf{k}_{bT} f_{q/a}(x_a, \mathbf{k}_{aT}; \mu_Q, Q^2) f_{\bar{q}/b}(x_b, \mathbf{k}_{bT}; \mu_Q, Q^2) \delta^{(2)}(\mathbf{q}_T - \mathbf{k}_{aT} - \mathbf{k}_{bT}) + (x_a \leftrightarrow x_b)$$

$$\frac{d\sigma}{dq_T dQ_M dy} = \frac{16\pi^2 \alpha^2}{9Q_M^3} q_T \sum_q e_q^2 x_a f_q(x_a, Q_M) x_b f_{\bar{q}}(x_b, Q_M) \mathcal{S}(q_T, x_a, x_b) \mathcal{B}^2(Q_M) + (x_a \leftrightarrow x_b)$$

Unpolarized TMDs: Setup

Introducing (exploration of) a separable form

$$f_q(x, k_\perp, Q_M) = f_q(x, Q_M) s(x, k_\perp) \mathcal{B}(Q_M)$$

Colinear-PDFs with DGLAP Evolution

Transvers Momentum Kernal

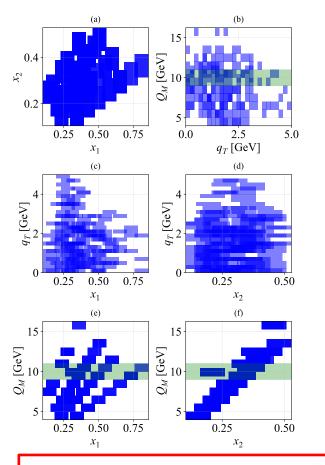
Transvers Momentum Evolution Kernal

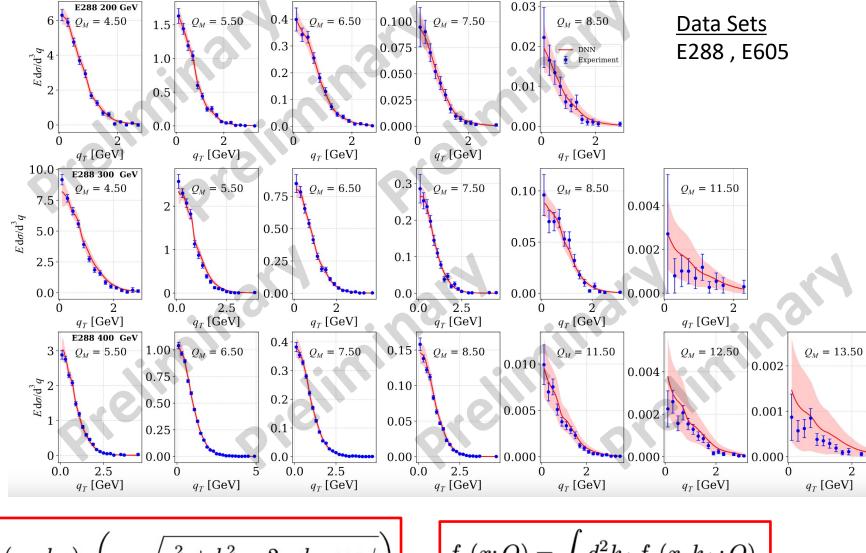
$$\frac{d\sigma}{dq_T dQ_M dy} = \frac{16\pi^2 \alpha^2}{9Q_M^3} q_T \sum_q e_q^2 \mathcal{H}^{DY} \int d^2 \mathbf{k}_{aT} d^2 \mathbf{k}_{bT} f_{q/a}(x_a, \mathbf{k}_{aT}; \mu_Q, Q^2) f_{\bar{q}/b}(x_b, \mathbf{k}_{bT}; \mu_Q, Q^2) \delta^{(2)}(\mathbf{q}_T - \mathbf{k}_{aT} - \mathbf{k}_{bT}) + (x_a \leftrightarrow x_b)$$

$$\frac{d\sigma}{dq_T dQ_M dy} = \frac{16\pi^2 \alpha^2}{9Q_M^3} q_T \sum_q e_q^2 x_a f_q(x_a, Q_M) x_b f_{\bar{q}}(x_b, Q_M) \mathcal{S}(q_T, x_a, x_b) \mathcal{B}^2(Q_M) + (x_a \leftrightarrow x_b)$$

$$S(q_T,x_a,x_b) = \int_0^\infty dk_\perp k_\perp \int_0^{2\pi} d\phi s(x_a,k_\perp) s\left(x_b,\sqrt{q_T^2+k_\perp^2-2q_Tk_\perp\cos\phi}
ight)$$

Current Status





$$S(q_T, x_a, x_b) = \int_0^\infty dk_\perp k_\perp \int_0^{2\pi} d\phi s(x_a, k_\perp) s\left(x_b, \sqrt{q_T^2 + k_\perp^2 - 2q_T k_\perp \cos \phi}\right)$$

$$f_q(x;Q) = \int d^2k_\perp f_q(x,k_\perp;Q)$$

Finalizing the simultaneous fits to obtain the TMDs (flavor dependent)

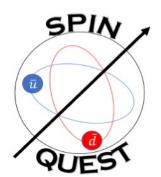
$$f_q(x, k_\perp, Q_M) = f_q(x, Q_M) s(x, k_\perp) \mathcal{B}(Q_M)$$

Summary & Outlook

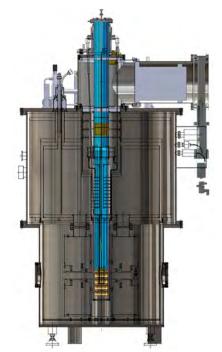
- We proposed a method for performing global fits to extract TMDs employing DNNs (first-ever application of DNNs in extracting TMDs).
- Extracting Sivers function was performed as an example of this method based on utilizing DNNs
- ➤ We have successfully tested our method with pseudo-data, also a dedicated systematic study.
- ➤ We projected SIDIS and DY Sivers asymmetries: for already completed experiments (as a validation check: COMPASS) and upcoming experiments (such as SpinQuest).
- Currently working on Unpolarized TMDPDF extraction...

Next:

- Applying the "DNN method" to extract other TMDs such as Transversity, Boer-Mulders function, as well as Spin-1 TMDs.
- ➤ Make projections to EIC kinematics



‡ Fermilab

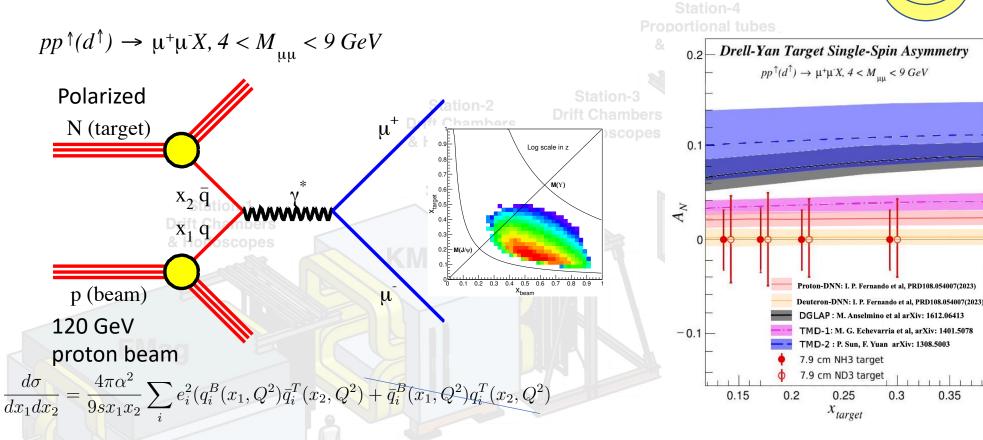


LANL-UVA
Polarized Target

https://spinquest.fnal.gov/ http://twist.phys.virginia.edu/E1039/

SpinQuest (E1039) Experiment at Fermilab

> Probing Sivers asymmetry from the (light) 'sea' quark contributions



Please Join The Effort

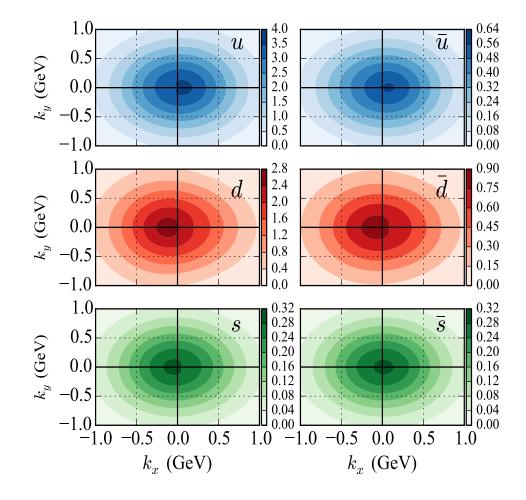
Dustin Keller (dustin@virginia.edu)[Spokesperson]
Kun Liu (liuk@lanl.gov)[Spokesperson])

Highest beam intensity on a polarized target ever!

Stay Tuned!

Thank you



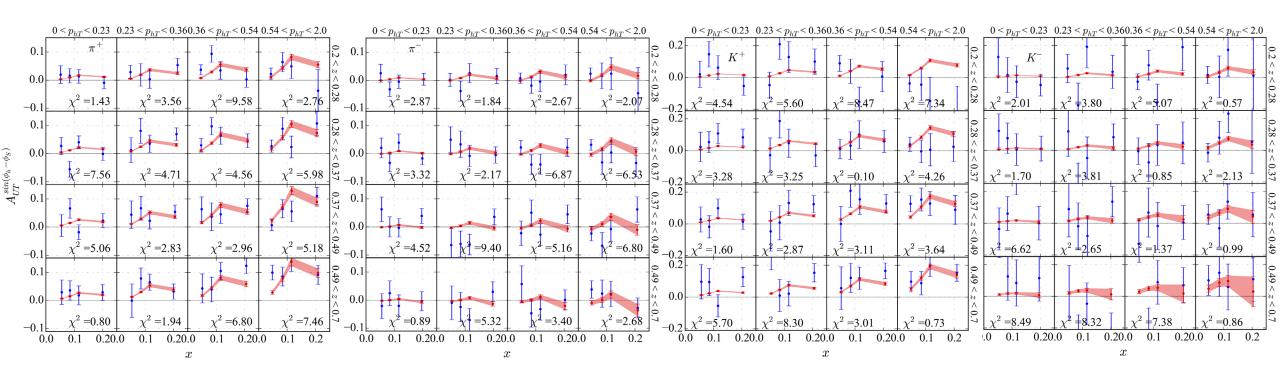




Office of Science

Backup Slides

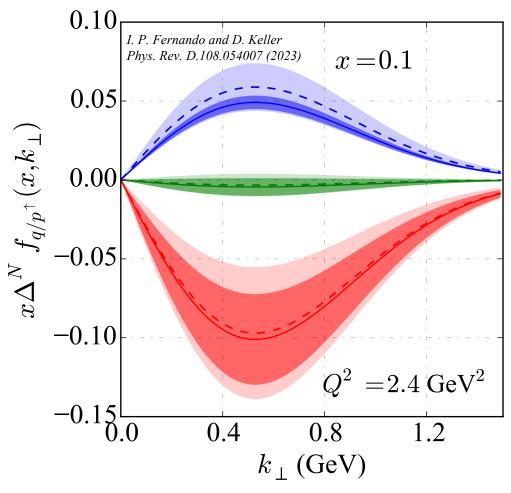
Projections from the "Proton" DNN Model



Projections of the of HERMES 2020 data for 3D kinematic bins, using the proton-DNN model including 68% C.L. error bands (in red) in comparison with the actual data points (in blue).

Backup

DNN Method: With Real data (Quality of the extraction)



The qualitative improvement of the extracted Sivers functions for u (blue), d (red), and s (green) quarks at x = 0.1 and $Q^2=2.4$ GeV² using the optimized proton-DNN model at the Second Iteration (solid-lines with dark-colored error bands with 68% CL), compared to the First Iteration (dashed-lines with light-colored error bands with 68% CL)

Systematic Studies: Choice of h(k)

$$\Delta^{N} f_{q/p^{\uparrow}}(x, k_{\perp}) = 2\mathcal{N}_{q}(x)h(k_{\perp})f_{q/p}(x, k_{\perp})$$

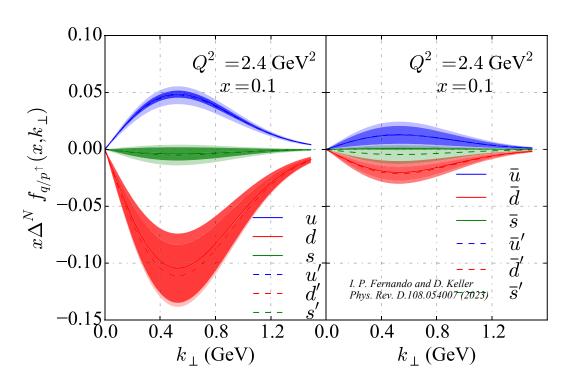


FIG. 19. Using two different $h(k_{\perp})$. Solid line with dark band represents the Sivers functions with $h(k_{\perp}) = \sqrt{2e} \frac{k_{\perp}}{m_1} e^{-k_{\perp}^2/m_1^2}$, whereas the dashed line with light band represents the Sivers functions with $h(k_{\perp}) = \frac{2k_{\perp}m_1}{m_1^2+k_1^2}$.

$$h(k_{\perp}) = \sqrt{2e} \frac{k_{\perp}}{m_1} e^{-k_{\perp}^2/m_1^2}$$

$$h(k_{\perp}) = \frac{2k_{\perp}m_1}{m_1^2 + k_{\perp}^2}$$

- ➤ It is clear that the DNN is capable of incorporating both types of h(k) without affecting the Sivers functions in the final model as well as the asymmetries (with deviation less than 1%).
- This is because DNN demonstrates that it maps to the h(k) such that the Sivers function is nearly unchanged.

Systematic Studies: data cuts

$$\begin{split} W^{\mu\nu} &= \sum_{f} |\mathcal{H}_{f}(Q^{2},\mu)|^{\mu\nu} \\ &\times \int d^{2}k_{\perp} d^{2}p_{\perp} \delta^{(2)}(z_{h}k_{\perp} + p_{\perp} - p_{hT}) \\ &\times F_{f/N^{\uparrow}}(x,z_{h}k_{\perp},S;\mu,\zeta_{F}) D_{h/f}(z_{h},p_{\perp};\mu,\zeta_{D}) \\ &+ Y(p_{hT},Q^{2}), \end{split}$$

In addition to the basic data cut $Q^2 > 1 \text{ GeV}^2$ we performed $Q^2 > 2 \text{ GeV}^2$ and, $p_{hT} < zQ$ cuts separately with the proton-DNN model to understand the impact on the extracted Sivers functions.

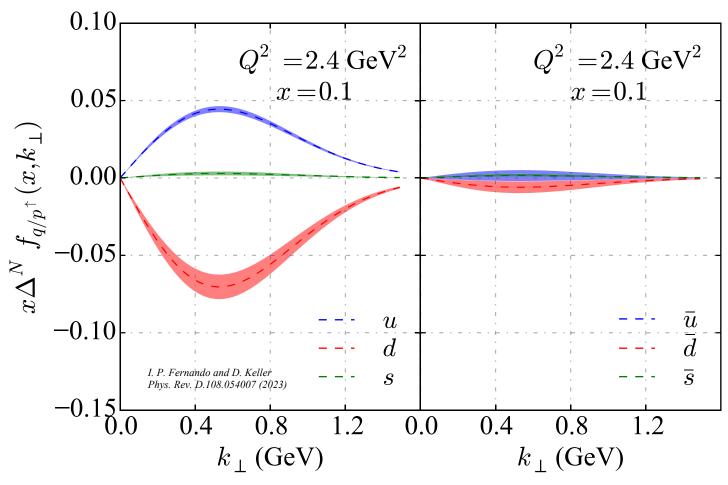


FIG. 18. Sivers functions from a retrained DNN model using the cuts [65] to the data demonstrating that being selective with the data can reduce the error bands of the fit but may also add an unintentional bias.

Systematic Studies: TMD Evolution

The solution of the TMD evolution equations

$$\mu^2 \frac{dF(x,b;\mu,\zeta)}{d\mu^2} = \frac{\gamma_F(\mu,\zeta)}{2} F(x,b;\mu,\zeta)$$

$$\zeta \frac{F(x,b;\mu,\zeta)}{d\zeta} = -\mathcal{D}(b,\mu)F(x,b;\mu,\zeta),$$

$$F(x,b;\mu,\zeta) = \left(\frac{\zeta}{\zeta_{\mu}(b)}\right)^{-\mathcal{D}(b,\mu)} F(x,b)$$

$$\mu \sim Q$$
, $\zeta_F \zeta_D \sim Q^4$, $\mu^2 = \zeta^2 = Q^2$

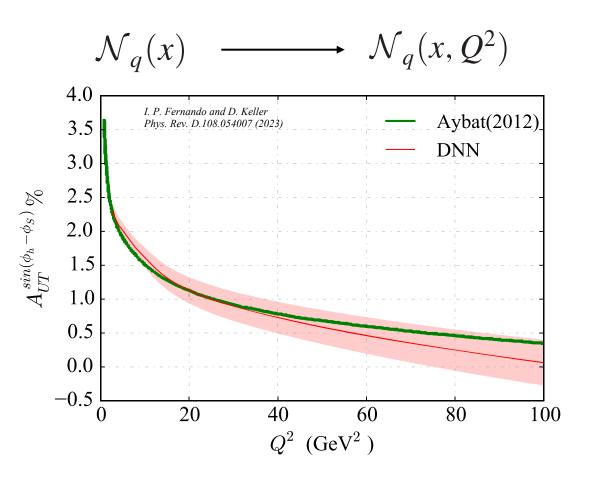
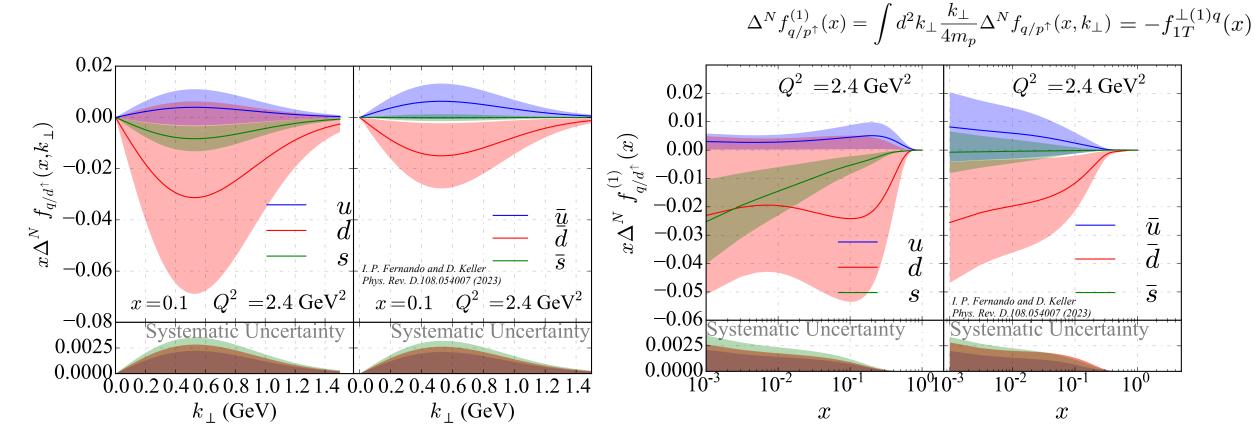


FIG. 21. The Sivers asymmetry evolution in Q^2 compared to the result from [6]. The red-colored solid line and the band represent the mean and standard deviation of the $A_{UT}^{\sin(\phi_h - \phi_S)}$ from 1000 replica models of the proton DNN at x = 0.12, z = 0.32, $p_{hT} = 0.14$ GeV.

DNN Method: Results from the "Deuteron" Model

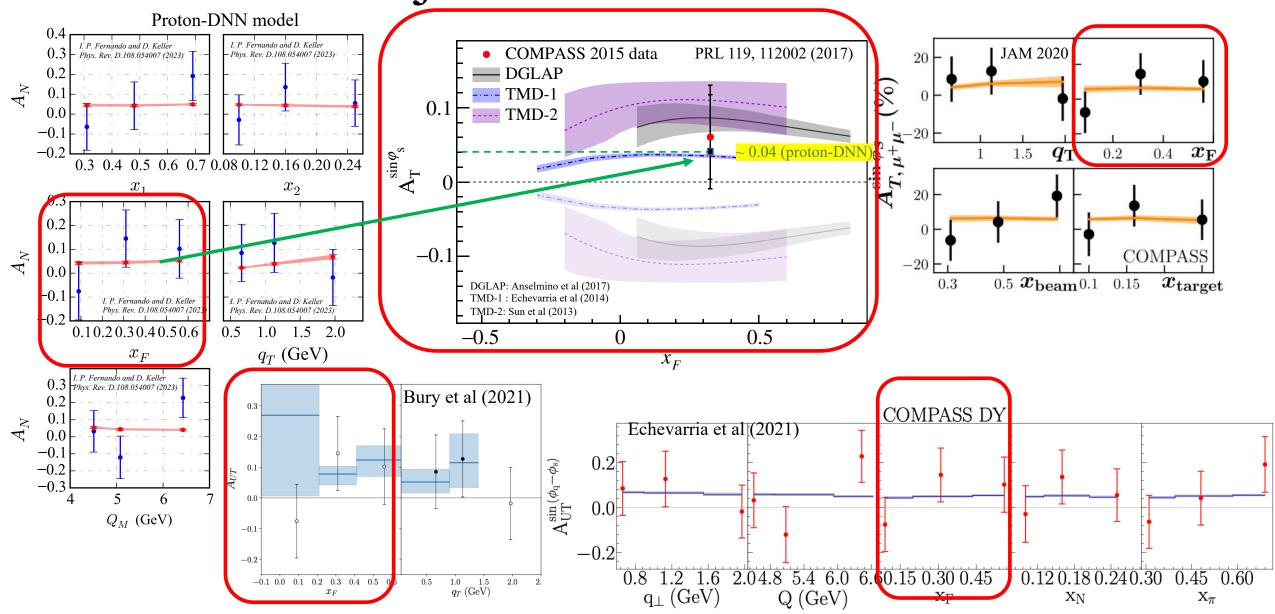
- ➤ Trained on COMPASS 2009 SIDIS data with Deuteron target.
- ➤ Did not imposed iso-spin symmetric conditions, or data cuts.

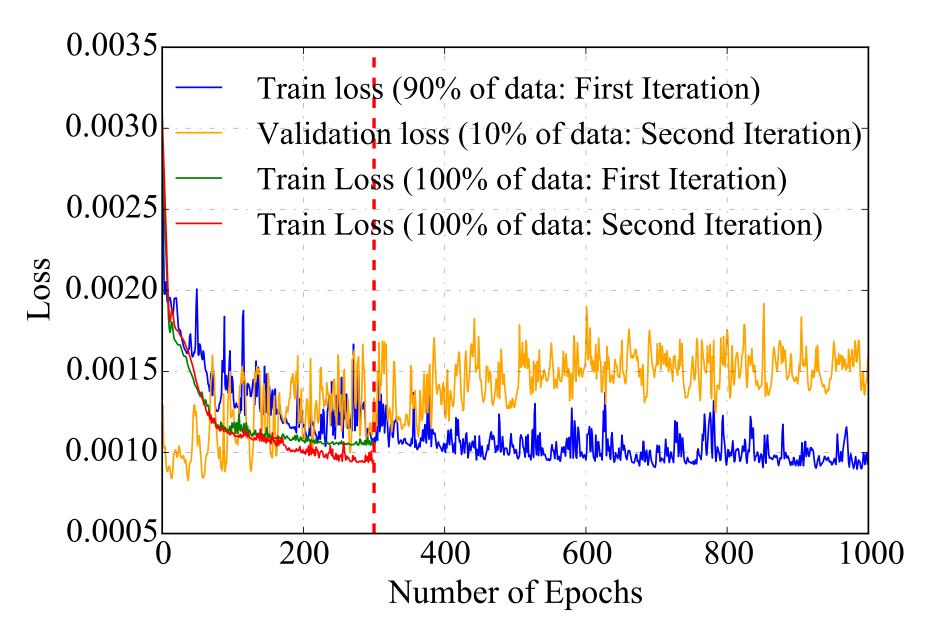
$$f_{1T,u \leftarrow d}^{\perp} = f_{1T,d \leftarrow d}^{\perp} = \frac{f_{1T,u \leftarrow p}^{\perp} + f_{1T,d \leftarrow p}^{\perp}}{2}$$



DNN Model Projections: DY

COMPASS 2017 DY Projections





Backup

TABLE III. The summary of the optimized sets of hyperparameters: The indications in the table are C_0^i and C_0^f for results from the pseudodata from the generating function, C_p^i , and C_p^f for results from SIDIS data from experiments associated with the polarized-proton target, and C_d^i and C_d^f for results from SIDIS data from experiments associated with the polarized-deuterium target, where i and f indicate the *First Iteration* and *Second Iteration* respectively. The initial learning rate is also listed (×10⁻⁴) as is the final training loss (×10⁻³). The accuracy and precision in each case are the maxima over the phase space.

Hyperparameter	\mathcal{C}_0^i	\mathcal{C}_0^f	${\cal C}^i_p$	${\mathcal C}_p^f$	\mathcal{C}_d^i	\mathcal{C}_d^f
Hidden layers	5	7	5	7	5	8
Nodes/layer	256	256	550	550	256	256
Learning rate	1	0.125	5	1	10	1
Batch size	200	256	300	300	100	20
Number of epochs	1000	1000	300	300	200	200
Training loss	0.6	0.05	1.5	1	2	1
$\varepsilon_u^{ ext{max}}$	95.67	99.27	55.21	94.04	56.80	93.02
$arepsilon_{ar{u}}^{\max}$	42.62	98.09	52.57	96.70	34.83	91.40
$arepsilon_d^{\max}$	80.46	98.89	55.69	93.13	52.44	89.27
$arepsilon_{ar{d}}^{\max}$	74.59	97.08	55.37	95.04	46.60	92.58
ε_s^{\max}	45.53	79.27	49.54	90.64	36.34	93.41
$\varepsilon_{\bar{s}}^{\max}$	59.27	91.13	33.89	82.51	65.57	91.45
σ_u^{\max}	3	0.1	5	2	2	0.4
$\sigma_{ar{u}}^{ ext{max}}$	2	0.2	6	2	8	2
σ_d^{\max}	10	1	20	6	2	1
$\sigma_{\bar{i}}^{\max}$	7	4	20	8	7	1
σ_{a}^{\max}	2	0.2	4	1	6	2
$\sigma_{ar{s}}^{ ext{max}}$	1	0.1	4	2	6	3