First Measurement of the Proton Generalized Polarizabilities with a polarized electron beam in Virtual Compton Scattering (VCS-IIIp)

A Proposal to Jefferson Lab PAC-53 Speaker: Michael Paolone (New Mexico State University) On behalf of the VCS-IIIp collaboration and its spokespeople: N. Sparveris*, H. Atac, S. Lee, M.K. Jones and M. Paolone

*contact person

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The proposal follows the Letter-of-Intent LOI12-23-001

PAC51 Recommendation:

"The physics case presented in the proposal is robust, and the document is very well-written. Feasibility concerns have not been identified, and the PAC encourages the proponents to proceed and submit a full proposal ..."

The proposal adds a unique component to an ongoing experimental program at JLab

First measurement of the proton GPs utilizing polarization measurements (BSAs) as an alternative to the unpolarized VCS cross section measurements



If we want to understand the characteristics of the proton as a building block of the universe... we **need** to understand the dynamics of the proton's constituents and how they contribute to those emergent characteristics.

- **Polarizability** is in an important characteristic of the proton:
 - How rigid is the proton in the presence of an EM field?
 - A fundamental property of the proton!
 - Sensitive to the excitation of the proton.
 - Can be accessed by Compton scattering (the photon acting as an induced EM field)
- Generalized Polarizabilities (GPs)
 - Accessed via virtual photon interaction.
 - Probe length (Q2) provides information on proton constituents in relation to structure of the proton

The proton is the only known stable composite particle!







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$$\overrightarrow{m} = \beta_M \overrightarrow{B}$$

Magnetic Polarizability

Paramagnetic Contribution from direct alignment









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Magnetic Polarizability

 $\overrightarrow{m} = \beta_M B$

Paramagnetic Contribution from direct alignment

> **Diamagnetic contribution** from meson cloud partially cancels with paramagnetic



















Static Polarizabilities



N BARYONS (S = 0, I = 1/2) p, N^+ = $u \ u \ d; \ n, \ N^0$ = $u \ d \ d$	PDGID:S016 JSON (beta) INSPIRE Q	
p $I(J^P) = 1/2(1/2^+)$		
p MASS (atomic mass units u)	$1.007276466621 \pm 0.00000000053$ ~	/
p MASS (MeV)	938.27208816 \pm 0.00000029 MeV \sim	·
$m_p - m_{\overline{p}} /m_p$	$< 7 imes 10^{-10}$ CL=90.0% \sim	·
$rac{q_{\overline{p}}}{m_{\overline{p}}}ig ig/ig(rac{q_p}{m_p}ig)$	$1.0000000003 \pm 0.000000016$ ~	/
$ rac{q_{ar p}}{m_{ar p}} -rac{q_p}{m_p})/rac{q_p}{m_p}$	$(0.1\pm 6.9) imes 10^{-12}$ ~	·
$ q_p+q_{\overline{p}} /e $	$< 7 imes 10^{-10}$, L=90.0% \sim	/ D
$ q_p+q_e /e$	$< 1 imes 10^{-21}$ \sim	·
MAGNETIC MOMENT	$2.7928473446 \pm 0.000000008 \mu_N$ ~	·
	$-2.792847344\pm 0.00000004\mu_N$ ~	·
$(\mu_p + \mu_{\overline{p}}) \ / \ \mu_p$	$(2\pm4) imes10^{-9}$ \sim	· -
p ELECTRIC DIPOLE MOMENT	$< 2.1 imes 10^{-25} e ext{cm}$ $igsaclassical{eq}$ $igsaclassical{eq}$	·
Electric polarizability $lpha$	$0.00112\pm0.00004\mathrm{fm}^3$ \sim	
Magnetic polarizability $oldsymbol{eta}$	$(2.5\pm0.4) imes10^{-4}$ fm 3 (S = 1.2) Normalized values of the second secon	
Charge radius	0.8409 ± 0.0004 fm \sim	
Magnetic radius	$0.851\pm0.026~{ m fm}$ \sim	·
Mean life $ au$	$> 9 imes 10^{29}$ years CL=90.0% \sim	·
D MEAN LIFE	\sim	·

Decay Modes

Expand all decays

Listed in the PDG as a fundamental property:



Generalized Polarizabilities

Accessed via virtual Compton scattering (VCS) as a function of Q2 (As $Q^2 \rightarrow 0$, the RCS static polarizabilities are found)

Can provide:

- spatial distribution of the polarization densities
- electric & magnetic polarizability radii
- manifestation of dynamical mechanisms for the EM polarizability effects

N BARYONS (S = 0, I = 1/2) p, N ⁺ = u u d; n, N ⁰ = u d d	PDGID:S016 JSON (beta) INS	Spire Q
${oldsymbol{p}}$ $I(J^P)$ = $1/2(1/2^+)$		
p MASS (atomic mass units u)	$1.007276466621 \pm 0.00000000053$ u	~
p MASS (MeV)	$938.27208816 \pm 0.00000029$ MeV	~
$ m_p\!\!-\!\!m_{\overline{p}} /m_p$	$< 7 imes 10^{-10}$ CL=90.0%	~
$ rac{q_{ar{p}}}{m_{ar{p}}} /(rac{q_p}{m_p})$	$1.00000000003 \pm 0.00000000016$	~
$(\left rac{q_{ar{p}}}{m_{ar{p}}} ight -rac{q_p}{m_p})/rac{q_p}{m_p}$	$(0.1\pm 6.9) imes 10^{-11}$	~
$ q_p+q_{ar p} /e$	$< 7 imes 10^{-10}$ CL=90.0%	~
$ q_p+q_e /e$	$< 1 imes 10^{-21}$	~
<i>p</i> MAGNETIC MOMENT	$2.7928473446 \pm 0.000000008\mu_N$	~
\overline{p} MAGNETIC MOMENT	$-2.792847344\pm 0.000000004\mu_N$	~
$(\mu_p + \mu_{\overline{p}}) \ / \ \mu_p$	$(2\pm4) imes10^{-9}$	~
<i>p</i> ELECTRIC DIPOLE MOMENT	$< 2.1 imes 10^{-25}~e$ cm	~
Electric polarizability $lpha$	$0.00112\pm0.00004 ext{fm}^3$	~
Magnetic polarizability $meta$	$(2.5\pm0.4) imes10^{-4}$ fm 3 (S = 1.2)	~
Charge radius	0.8409 ± 0.0004 fm	~
Magnetic radius	0.851 ± 0.026 fm	~
Mean life $ au$	$> 9 imes 10^{29}$ years <code>CL=90.0%</code>	~
\overline{p} MEAN LIFE		~



Decay Modes





Virtual Compton Scattering



SCATTERING PLANE



DR

valid below & above Pion threshold



- The **DR** (dispersion relations) method:
 - Available above and below the pion threshold
 - The scalar polarizabilities enter as free parameters to be fit.
- The **LEX** (Low energy expansion) method
 - Valid only below the pion threshold

Dispersive integrals for Non Born amplitudes

Spin GPs are fixed

Scalar GPs have an unconstrained part

Fit to the experimental cross sections at each Q^2



valid only below **Pion threshold**

LEX





scalar GPs α_F and β_M

MIT-Bates @ Q²=0.06 GeV²



MAMI-A1 @ $Q^2=0.33 \text{ GeV}^2$



Early Experiments



Jlab-Hall A @ Q²=0.9 & 1.8 GeV²



Initial investigations showed that the proton generally increases in stiffness as Q2 increases.



MIT-Bates @ Q²=0.06 GeV²



MAMI-A1 @ $Q^2=0.33 \ GeV^2$



Jlab-Hall A @ Q²=0.9 & 1.8 GeV²



Early Experiments



Initial investigations showed that the proton generally increases in stiffness as Q2 increases.

Early Q2 = 0.33 GeV2 measurement at MAMI seemed to buck the trend of a monatomic decrease. A second measurement at MAMI had similar results. Phys. Rev. Lett 85, 708 (2000)

Eur. Phys. J. A37, 1-8 (2008)



Theoretical Predictions





A. Yu. Korchin and O. Scholten A. Metz and D. Drechsel

Phys. Rev. D 62, 014013 (2000) Phys. Rev. C 63, 025205 (2001) Phys. Rev. C 58, 1098 (1998) Z. Phys. A 356, 351 (1996)



Theoretical Predictions





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Theoretical Predictions



Recent Experiments





Current Landscape and Questions

• Electric Polarizability:

• Is the observed structure coincidental?







Current Landscape and Questions

- - If so: More precise measurements will help inform
 - If not: Strong tension exists in the world data. Additional measurements can help pinpoint possible sources of





Magnetic Polarizability

- Large uncertainties and discrepancies exist in the world data.
- High precision data is needed to disentangle diamagnetic and paramagnetic contributions in the nucleon.

Current Landscape and Questions

• Electric Polarizability:

- Is the observed structure coincidental?
 - If so: More precise measurements will help inform theory.
 - If not: Strong tension exists in the world data. Additional measurements can help pinpoint possible sources of tension







- Magnetic Polarizability
 - Large uncertainties ar the world data.
 - extractions (this proposal) High precision data is diamagnetic and paramagnetic contributions in the nucleon.

Current Landscape and Questions

• Electric Polarizability:

• Is the observed structure coincidental?

• If so: More precise measurements will help inform theory.





A new path to measure GPs: Induced asymmetries

One can set-up and measure beam-charge or beam-spin asymmetries as an independent path to extracting the generalized polariabilities.

Virtual Compton scattering at low energies with a positron \mathbf{beam}

Barbara Pasquini^{a,1,2}, Marc Vanderhaeghen^{b,3}

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²Istituto Nazionale di Fisica Nucleare, Sezione di Pavia, 27100 Pavia, Italy

³Institut für Kernphysik and PRISMA⁺ Cluster of Excellence, Johannes Gutenberg Universität, D-55099 Mainz, Germany



This proposal:

Beam-Spin Asymmetry (BSA):

$$A^e_{LU} = \frac{d\sigma^e_+ - d\sigma^e_-}{d\sigma^e_+ + d\sigma^e_-}$$

Eur. Phys. J. A 57 (2021) 11, 316

(a): The beam-charge asymmetry as a function of the photon scattering angle at Q2 = 0.43 GeV2.

(b) & (c): The electron and positron beam-spin asymmetry as a function of the photon scattering angle for out-of-plane kinematics.

Future proposal (positron-beam): Beam-Charge Asymmetry (BCA): $A_{UU}^{C} = \frac{(d\sigma_{+}^{+} + d\sigma_{-}^{+}) - (d\sigma_{+}^{-} + d\sigma_{-}^{-})}{d\sigma_{+}^{+} + d\sigma_{-}^{+} + d\sigma_{-}^{-} + d\sigma_{-}^{-}}$





Extensions: Spatial Dependence of induced polarizations

Light-front quark charge densities \rightarrow Deformation under EM field



Phys. Rev. Lett. 104, 112001 (2010)

M. Gorchtein, C. Lorce, B. Pasquini, M. Vanderhaeghen



Extensions: Spatial Dependence of induced polarizations

Light-front quark charge densities \rightarrow Deformation under EM field



x-y defines the transverse plane with the z-axis being the direction of the fast-moving proton

Extensions: Polarizability Radii $\langle r_{\alpha_E}^2 \rangle = \frac{-6}{\alpha_E(0)} \cdot \frac{d}{dQ^2} \alpha_E(Q^2) \Big|_{Q^2 = 0}$





$$\langle r_{\beta_M}^2 \rangle = \frac{-6}{\beta_M(0)} \cdot \frac{d}{dQ^2} \beta_M(Q^2) \Big|_{Q^2 = 0}$$







Extensions: Polarizability Radii $\langle r_{\alpha_E}^2 \rangle = \frac{-6}{\alpha_E(0)} \cdot \frac{d}{dQ^2} \alpha_E(Q^2) \Big|_{Q^2=0}$



$$36 \pm 0.29 \text{ fm}^2$$
Atac 2021
Atac 2021
Atac 2021
Atac 2021
Atac 2020
Atac

$$\langle r_{\alpha_E}^2 \rangle = 1.36 \pm 0.29 \text{ fm}^2$$

$$\langle r_{\beta_M}^2 \rangle = \frac{-6}{\beta_M(0)} \cdot \frac{d}{dQ^2} \beta_M(Q^2) \Big|_{Q^2 = 0}$$







Experimental Setup



$$A^e_{LU} = \frac{d\sigma^e_+ - d\sigma^e_-}{d\sigma^e_+ + d\sigma^e_-}$$

Effectively, the experiment uses the same experimental setup as in the VCS-II with the addition of a polarized electron beam

Hall C: SHMS, HMS (in standard configuration) E=2.2 GeV I=70µA 85% polarization LH2 10 cm







Experience running, extracting, and publishing with similar setups in Hall-C

Settings are within normal operating conditions

Singles rates are low





Studies to optimize the experiment kinematics

This proposal focuses on extracting one Q2 point with high precision



$Q^2 = 0.35 \,\mathrm{GeV^2}$

Better sensitivity at 2.2 GeV

Sensitivity grows near pole-mass of $\Delta(1230)$ resonance



Studies to optimize the experiment kinematics

This proposal focuses on extracting one Q2 point with high precision





Better sensitivity at 2.2 GeV

Decreased sensitivity: $\theta_{\gamma^*\gamma} < 100 \text{ deg}$

Sensitivity grows near pole-mass of Δ (1230) resonance



Simulation Studies

Simulation weighted with the DR model for VCS (Pasquini et al)

Systematics:

Acceptance / bin migration	5%
Bin centering	6%
Radiative effects	2%
Beam polarization	3%
Background subtraction	1%
Target wall	1%
Pion contamination	1%

System. uncertainties ~ 50% compared to the statistical



Projected BSA measurements



 $\Lambda_{\alpha} \& \Lambda_{\beta}$ fitted by an χ^2 minimization which compares the DR asymmetries to the measured ones and the GPs are determined

GPs extracted through a DR fit to the BSA measurements.

Experiment Kinematics & Projected Measurement

Beam Energy (GeV)	Beam Current (uA)	Beam Requirements	Target	$ heta_e^{\circ}$	$P_e^{\prime} \ (MeV/c)$	$ heta_p^\circ$	$P_p^{\prime} \ (MeV/c)$	Beamtime (days)	
2.2	70	Polarized	LH2 (10cm)	17.72	1676.41	60.71	723.69	2	
2.2	70	Polarized	LH2 (10cm)	17.72	1676.41	56.21	808.93	6	Total beam time : 20.
2.2	70	Polarized	LH2 (10cm)	17.72	1676.41	51.12	874.74	6	
2.2	70	Polarized	LH2 (10cm)	17.72	1676.41	47.10	908.37	6	
Total Time								20	+ 0.5 (dummy & check outs)





Experiment Kinematics & Projected Measurement

	Beam	Beam	\mathbf{Beam}	Target	θ_e°	P'_{e}
	Energy	$\mathbf{Current}$	Requirements			(MeV)
	(GeV)	(uA)				
	2.2	70	Polarized	LH2 (10cm)	17.72	1676.4
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	2.2	70	Polarized	LH2 (10cm)	17.72	1676.4
_	2.2	70	Polarized	LH2 (10cm)	17.72	1676.4
1	Total Time					





Total beam time (with VCS-II) : 16.5 days



Experiment Kinematics & Projected Measurement

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	Energy	Current	Requirements			(MeV)
_	(GeV)	(uA)				
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_	2.2	70	Polarized	LH2 (10cm)	17.72	1676.4







Summary

• Focus on the study of fundamental proton properties

• Gain insight to the response of the proton constituents to an external EM field, to the interplay of the polarizability mechanisms & to the tomography of the polarizability effects in the proton

• First measurement of the polarizabilities using a new experimental method

- Unique independent cross check to the world data
- Employing alternative experimental methods is instrumental in the scientific process, frequently revealing surprises and new knowledge (e.g. proton EM form factors, proton charge radius)
- Particularly valuable in addressing the theoretical challenges with the electric GP
- unpolarized VCS measurements and is well tested and understood
- The proposal complements an ongoing experimental program at Jlab with a unique & novel approach • The experiment uses a standard experimental setup in Hall C that is already employed for the
 - E=2.2 GeV, I=70µA, 85% polarization, LH2 10 cm
 - $\theta_{\rm SHMS}$ ~ 18 deg, $p_{\rm SHMS}$ ~ 1.7 GeV, $\theta_{\rm HMS}$ ~ 47 to 60 deg, $p_{\rm HMS}$ ~ 0.7 to 0.9 GeV

• The collaboration is well prepared & ready, considering the upcoming VCS-II experiment

• Matching the kinematics of the two experiments reduces the beam-time request by 4 days

