

DVCS on a polarised proton target at CLAS12

Samy Polcher Rafael

Introduction

- Outstanding question : How does the nucleon's mass and spin arise from partons ?
- First step is to map out the nucleon
- **GPDs** describe the longitudinal momentum and transverse position of partons in the nucleon





 $\frac{1}{2} = \sum J^q + J^g$

- Access to the spin decomposition of the nucleon
- Access to its mechanical properties

$$=\sum_{q}^{1} \frac{1}{2} \int_{-1}^{1} \mathrm{d}x \, x(H^{q}(x,\xi,0) + E^{q}(x,\xi,0)) + \frac{1}{2} \int_{-1}^{1} \mathrm{d}x \, H^{g}(x,\xi,0) + E^{g}(x,\xi,0)$$



[arXiv:hep-ph/0504030] [arXiv:1807.07620]

Deeply Virtual Compton Scattering

- DVCS offers the most straightforward access to GPDs
- DVCS can be factorised into :
 - → Hard part y*q scattering computed in perturbative QCD
- Two indistinguishable processes, DVCS and Bethe-Heitler

$$|T|^{2} = |T_{\rm DVCS}|^{2} + |T_{\rm BH}|^{2} + \underbrace{T_{\rm DVCS}T_{\rm BH}^{*} + T_{\rm DVCS}^{*}T_{\rm BH}}_{\rm I}$$

Amplitude is expressed as a function of FFs and CFFs which are functions of GPDs

$$\mathcal{F} = \int_{-1}^{1} dx F(\mp x, \xi, t) \left[\frac{1}{x - \xi + i\epsilon} \pm \frac{1}{x + \xi - i\epsilon} \right]$$

[d'Hose2016]

Observables

- Asymmetries in the DVCS cross section are sensitive to CFFs
- Beam spin asymmetry (BSA), polarised electron and unpolarised proton

$$A_{\rm LU}(\phi) \sim \frac{s_{1,\rm unp}^{\mathcal{I}} \sin \phi}{c_{0,\rm unp}^{\rm BH} + (c_{1,\rm unp}^{\rm BH} + c_{1,\rm unp}^{\mathcal{I}} + ...) \cos \phi...}$$
$$s_{1,\rm unp}^{\mathcal{I}} \propto \Im [F_1 \mathcal{H} + \xi (F_1 + F_2) \widetilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E}].$$

• Target spin asymmetry (TSA), unpolarised electron and polarised proton

$$A_{\rm UL}(\phi) \sim \frac{s_{1,\rm LP}^{\mathcal{I}} \sin \phi}{c_{0,\rm unp}^{\rm BH} + (c_{1,\rm unp}^{\rm BH} + c_{1,\rm unp}^{\mathcal{I}} + ...) \cos \phi + ...}$$
$$s_{1,\rm LP}^{\mathcal{I}} \propto \Im m[F_1 \widetilde{\mathcal{H}} + \xi(F_1 + F_2)(\mathcal{H} + \frac{x_b}{2}\mathcal{E}) - \xi(\frac{x_b}{2}F_1 + \frac{t}{4M^2}F_2)\widetilde{\mathcal{E}}]$$

- Both are sensitive to the imaginary part of the H and \tilde{H} CFFs, measuring both is essential to separate the two contributions.

The RGC experiment

- New polarised target APOLLO, cryogenic solid target
- Dynamic polarisation of hydrogen or deuterium in NH3 or ND3 cells
- Took data from June 2022 to March 2023 in three run periods
 - Summer 22, FTON 28 PAC days (pass 1 completed)
 - Fall 22, FTOUT 30 PAC days (pass 1 ongoing)
 - Winter 23, FTON 22 PAC days
- Only Summer22 is in the current analysis (~50% of the DVCS stat collected)

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Previous CLAS6 measurement

- In 2009 data taken at CLAS on a polarised proton target with JLab 6GeV
- Measurement of the BSA, TSA and DSA in JLab6 kinematics
- RGC can expand the phase space probed with the Jlab 12GeV and CLAS12 upgrades

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DVCS Event selection

- The event must have at least one electron one proton and one photon
- Apply particle identification and fiducial cuts
- Exclusive process → exclusivity variables
- Nuclear background due to the unpolarised nitrogen in the target
 - Data taken on Carbon target to estimate the background
- Apply exclusivity cuts to remove as much of the nuclear background as possible

Dilution factor

$$A_{UL} = \frac{1}{D_f} \frac{N^{++} + N^{-+} - N^{+-} - N^{--}}{P_t^{-}(N^{++} + N^{-+}) + P_t^{+}(N^{+-} + N^{--})}$$

• To take into account the remaining nuclear background, the TSA and DSA are scaled by the dilution factor

$$n_{NH3}=f_c[(L-l_{NH3})
ho_{He}\Delta\sigma_{He}+l_{NH3}
ho_{NH3}(rac{7}{6}\Delta\sigma_C+3\Delta\sigma_H)]$$

- RGC took data with multiple targets:
 - Empty target \rightarrow Foil contribution
 - He bath \rightarrow He contribution
 - Carbon \rightarrow Nitrogen contribution
 - CH2 \rightarrow Hydrogen contribution

e - NH3 beadsHelium Bath

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[S.E. Kuhn]

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Dilution factor

$$Df = rac{3 l
ho_{NH3} \Delta \sigma_{H}}{(L-l)
ho_{He} \Delta \sigma_{He} + l
ho_{NH3} (rac{7}{6} \Delta \sigma_{C} + 3 \Delta \sigma_{H})}$$

• The dilution factor is stable as a function of all kinematic variable so we use a single value for all bins

 $D_f = 94 \pm 1\%$

• Target packing fraction : ~57%

۵ طَ 0.9 0.9 0.8 0.8 0.7 0.7 $N^{NH3} = 61885, N^{C} = 562$ $N^{NH3} = 61885, N^{C} = 562$ 0.6 0.6F D₄=0.938 ± 0.0033 D_i=0.940 ± 0.0031 0.5 0.5 350 ¢ [°] 250 0.2 0.6 0.8 1 1.2 -t [GeV²] 50 200 300 0.4 100 150 ے ۵Ť 0.9F 0.9 0.8 0.8 0.7 0.7 $N^{NH3} = 61885, N^{C} = 562$ $N^{NH3} = 61885, N^{C} = 562$ 0.6 0.6 D₄=0.940 ± 0.0031 D₄=0.940 ± 0.0031 0.5<u>L.</u> 0.5 0.5 0.€ xb [GeV²] Q² [GeV²] 0.1 0.2 0.3 0.4 2 3 4

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Target polarisation

$$A_{UL} = \frac{1}{D_f} \frac{N^{++} + N^{-+} - N^{+-} - N^{--}}{P_t^{-}(N^{++} + N^{-+}) + P_t^{+}(N^{+-} + N^{--})}$$

• The $ep \rightarrow e'p'$ elastic double spin asymmetry is well known

$$A_{th} = \frac{A_{exp}}{P_b P_t} \qquad A_{exp} = \frac{N^+ - N^-}{D_f (N^+ + N^-)}$$

• We can extract the target polarisation by comparing it to the measured elastic asymmetry

$$P_t^+ = 87 \pm 3 \%$$

 $P_t^- = 80 \pm 3 \%$

N. Pilleux

Yields and π^0 background

$$A_{UL} = \frac{1}{D_f} \frac{N^{++} + N^{-+} - N^{+-} - N^{--}}{P_t^{-}(N^{++} + N^{-+}) + P_t^{+}(N^{+-} + N^{--})}$$

$$N^{bt} = \frac{Y^{bt}}{\mathrm{FC}^{bt}} (1 - R^{bt})$$

- Y^{bt} event count in the bin with beam polarisation b and target polarisation t
- FC^{bt} Beam charge in the spin configuration
- $R^{bt} \pi^0$ contamination fraction

- A π^0 decay $\pi^0 \rightarrow \gamma \gamma$ can pass as a DVCS is one of the photons carries most of the momentum
- Significant background contribution that needs to be subtracted

π^0 subtraction method

fDVCS rec

 π^{0} events that pass the DVCS selection \rightarrow What we need to estimate

π^{0} rec

Wide π^0 sample selected from data with loose exclusivity cuts

π^0 decays

- Each data π^0 is randomly decayed 1000x
- Decays are passed through the detector simulation
- π^{0} and DVCS event selections are applied

π^0 data and decays

• The π^0 decay distributions matches data

Method validation

- Apply the method to a π^0 simulation
- All events reconstructed as DVCS are background contamination
 - The false DVCS distribution estimated match the mock data well

Method validation

π^0 contamination in data

Binning

- 2 bins in Q^2 , t, x_b so that there is the same number of events per bin
- N bins in phi with at least 500events per bin and at least 15° wide.

Beam spin asymmetry

Only statistical errors are included .

 $0.5 = -0.12, < Q^2 > = 1.72, < x_b > = 0.12, < t > = -0.58$

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Target spin asymmetry

Only statistical errors are included

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[⊥]/_↓ <Q²> = 1.94, <x_↓> = 0.20, <t> = -0.71

 $_{4} = \langle Q^{2} \rangle = 1.72, \langle x \rangle = 0.12, \langle t \rangle = -0.58$

CIQS

Work in progress: Mahalanobis distance cut

- 1D cuts on 8 different variables are difficult to adjust systematically
- We know the value of each variable for an ideal exclusive event
- Define a distance between each event to that ideal reference
- Mahalanobis distance: distance weighted by the covariance matrix between all variables
- The covariance matrix is taken from a GEMC dvcs simulation

$$d(ec{x},ec{y}) = \sqrt{(ec{x}-ec{y})^T \Sigma^{-1} (ec{x}-ec{y})}$$

Work in progress: Mahalanobis distance cut

- Cut at distance<4 to have a cut comparable to the 1D cuts
- **10% increase in statistics**, more π^0 background is removed but the dilution factor goes from 94% to 92%

Work in progress: Mahalanobis distance cut

Summary & Outlook

- RGC is the first polarised target experiment with the CLAS12 detector
- The pDVCS analysis is ongoing and shows promising results
- → Tools are in place to take into account nuclear and π^0 backgrounds
- The main remaining steps are:
 - → Refined momentum correction
 - Systematic error estimates
- The rest of the data collected is being processed and will be available in the coming month. It is expected to double the proton DVCS statistics for this analysis.

π^0 method comparison

- Full Monte-Carlo method, the 1y to 2y acceptance ratio is computed from a ٠ simulation and scaled to data
- Genepi π^0 simulation on NH3 target •

^IV Data

Work in progress: Momentum corrections

- Electron and proton corrections from Monte-Carlo
- Photon in FT: RGA pass2 correction (Asli Acar)

Work in progress: Photon in the FD correction

- Energy correction as a function of momentum from electron (trackin ECAL energy deposit)
- sector/low E correction with pi0 mass, fit of the pi0 mass for events with $|\theta g1 \theta g2| < 5^{\circ}$

Calibration

p [GeV]

Potential FCup issues

31/10/2024

Fiducial cuts, DC and FTCAL

- DC: edge > 4cm on region 2
- FTCAL: 8.25 < r < 15.75 cm

Fiducial cuts PCAL

• RGA common analysis note medium cuts u,v,w > 14cm. ~3 scintillator bars

Beam polarisation

$$A_{\rm LU} = \frac{P_t^{-}(N^{++} - N^{-+}) + P_t^{+}(N^{+-} - N^{--})}{P_b(P_t^{-}(N^{++} + N^{-+}) + P_t^{+}(N^{+-} + N^{--}))}$$
$$A_{UL} = \frac{1}{D_f} \frac{N^{++} + N^{-+} - N^{+-} - N^{--}}{P_t^{-}(N^{++} + N^{-+}) + P_t^{+}(N^{+-} + N^{--})}$$

P_b: Beam polarisation

Measured with a Moller polarimeter regularly all along the experimental run

$$P_{b} = 82.6 \pm 0.2 \%$$

π^0 event selection

- Wide selection to include all potential contamination including from some SIDIS $\pi^{\scriptscriptstyle 0}$
- There is a significant contribution from nuclear background π^0

Data and decays agreement

- We apply the π^0 cuts to the simulated decays
- The decays need to be weighted so that one data event has the same weight as all the decays coming from that event

$$W_i = \frac{1}{N_i^{\text{lev2}}} \left(1 - \frac{N_i^{\text{badElec}}}{N^{\text{decay}}}\right)$$

Double spin asymmetry

