

Forward Heavy-Ion Physics at the LHC



Daniel Tapia Takaki
The University of Kansas

High Energy Nuclear Physics with Spectator Tagging workshop
Old Dominion University, March 10, 2015

Plan of this talk

Introduction: low-x physics

Ultra-Peripheral (pp, pA and AA) Collisions

What are UPCs

Why at LHC

Recent results

Future directions

Recent workshop on Forward HI physics

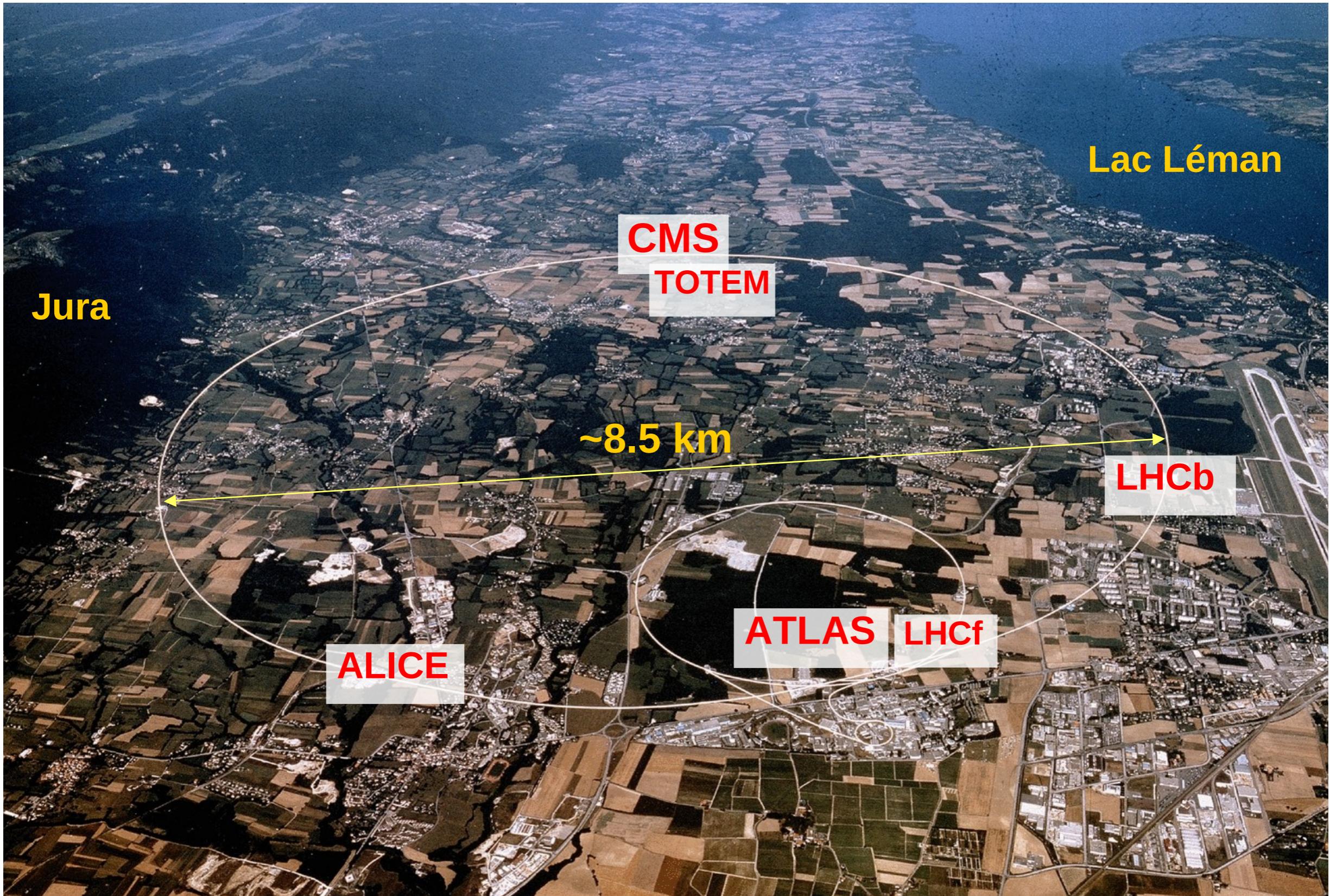
**LHC WG meeting on difraction and forward
physics & Future directions heavy-ion
physics in the forward region**

<http://cern.ch/lawrence2014>

September 3-6, 2014

Lawrence and Kansas City

***CERN Yellow Report on Forward physics,
in preparation***



Forward HI physics

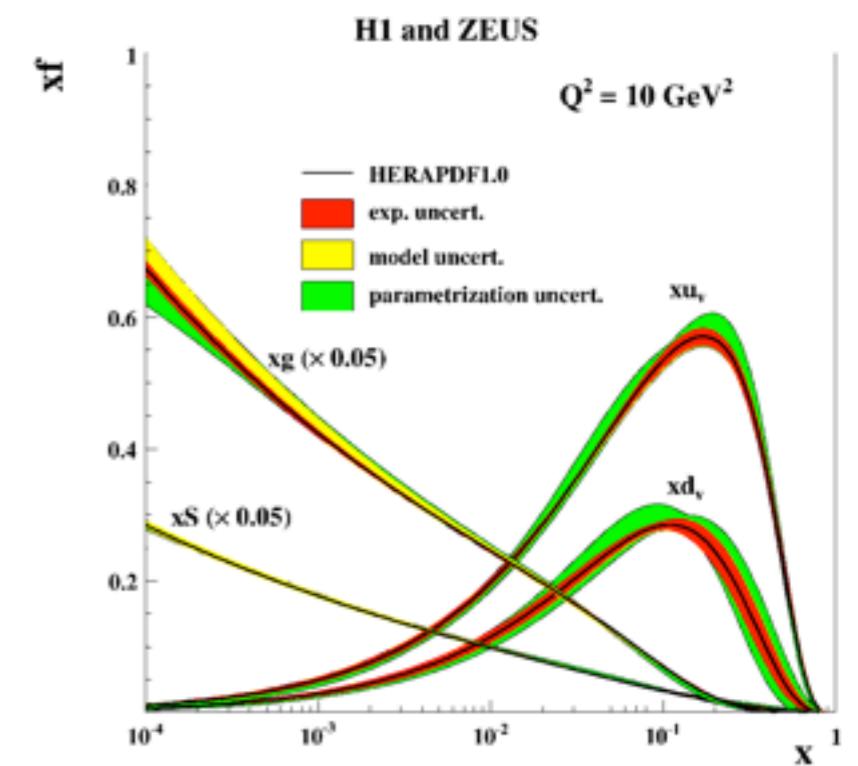
Large rapidity gap between small- x partons and most forward particles

$$x \approx p_{parton}/p_{proton}$$
$$Q^2 = -(p'_e - p_e)^2$$

***the nature of the initial state
is one of the most important
questions in relativistic
heavy-ion physics.***

UPCs are cleaner probes of nPDFs

Low- x regime dominated by gluons

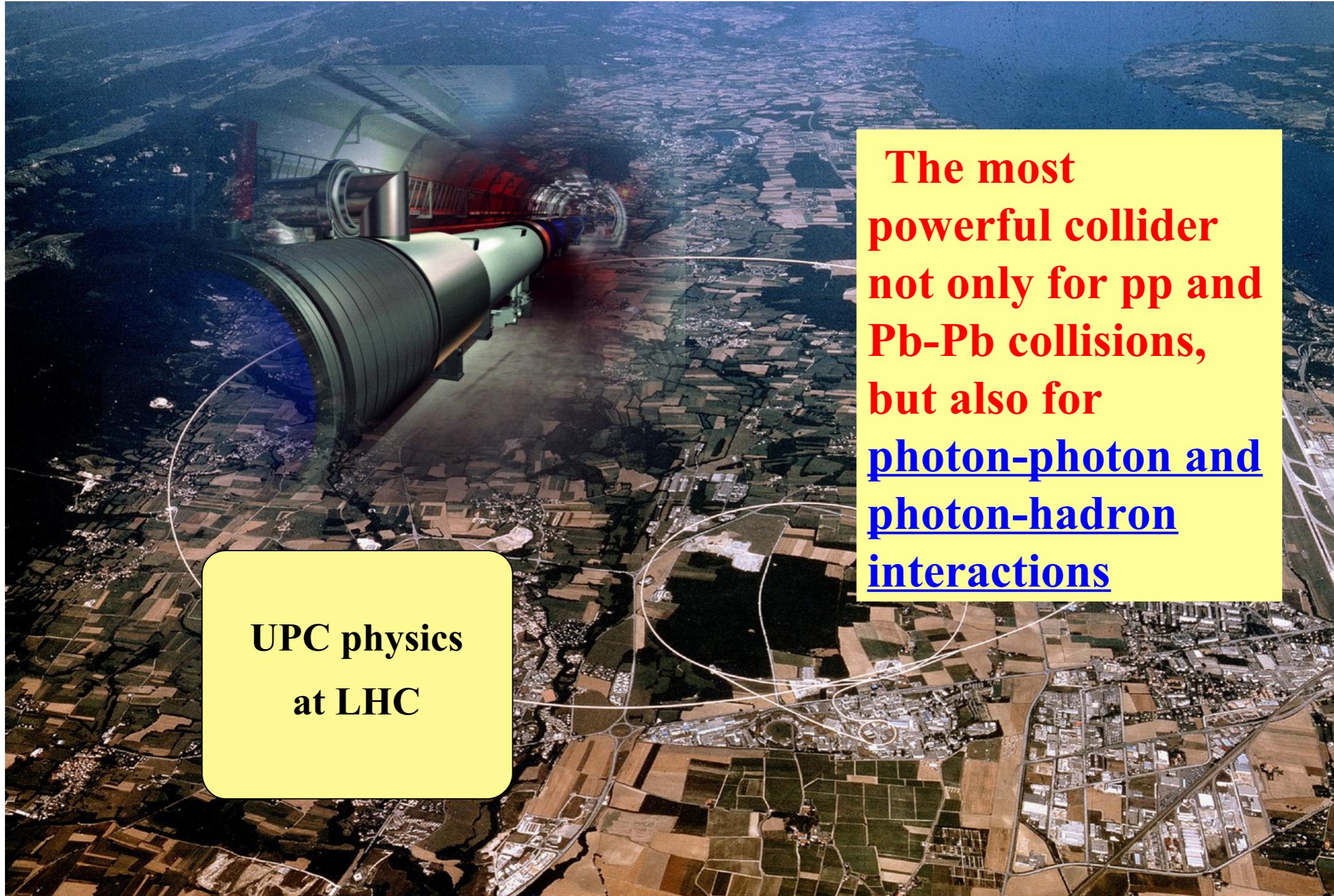


UPC at LHC can be seen as the precursor of part of the EIC physics

UPC described in the two recent White Papers:
EIC white paper: arXiv:1212.1701 [nucl-ex]

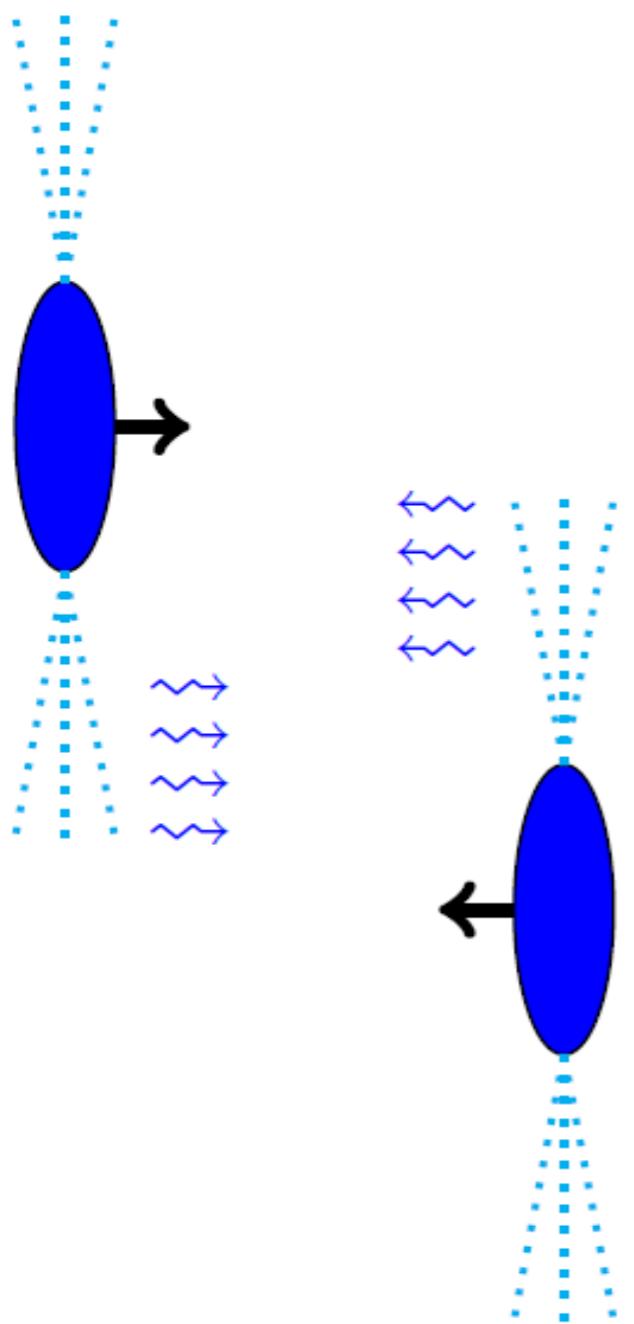
HI White paper: arXiv:1502.02730 [nucl-ex]

Using the LHC as a $\gamma\gamma$, γPb , γp collider



UPCs in Pb-Pb

Why Ultra-Peripheral collisions

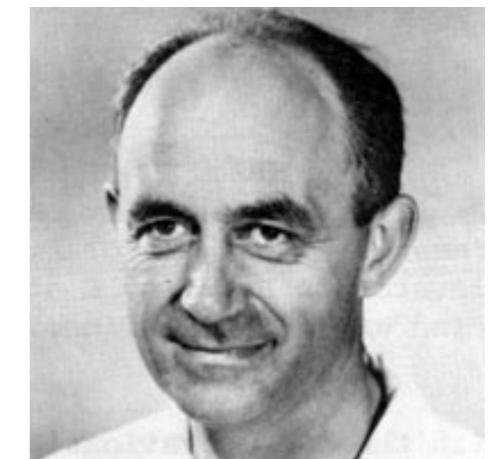


Nuovo Cim.,2:143-158,1925

<http://arxiv.org/abs/hep-th/0205086>

Therefore, we consider that when a charged particle passes near a point, it produces, at that point, a variable electric field. If we decompose this field, via a Fourier transform, into its harmonic components we find that it is equivalent to the electric field at the same point if it were struck by light with an appropriate continuous distribution of frequencies.

High photon flux $\sim Z^2$
→ well described by the Weizsäcker-Williams approximation



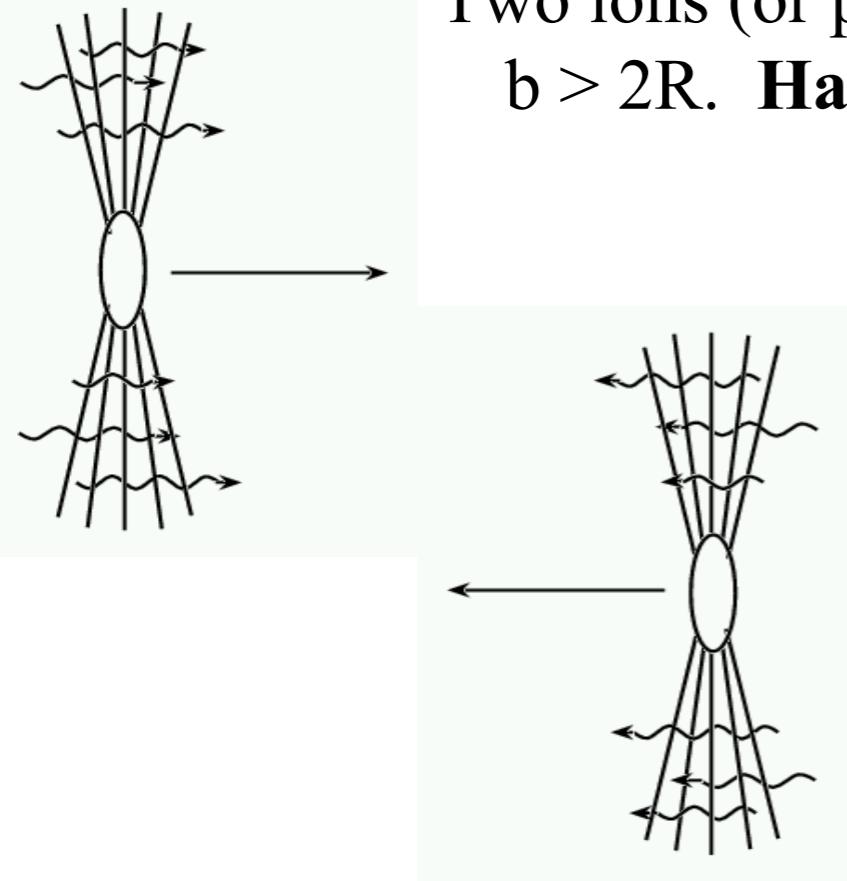
Enrico FERMI

The electromagnetic field surrounding these protons/ions can be treated as a beam of quasi real photons

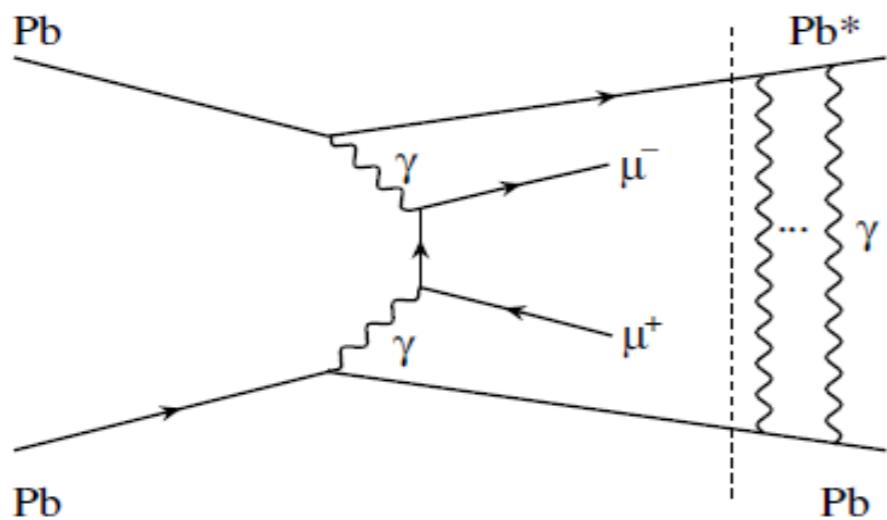
Two ions (or protons) pass by each other with impact parameters $b > 2R$. **Hadronic interactions are strongly suppressed**

Why ultra-peripheral heavy-ion collisions

Two ions (or protons) pass by each other with impact parameters
 $b > 2R$. Hadronic interactions are strongly suppressed

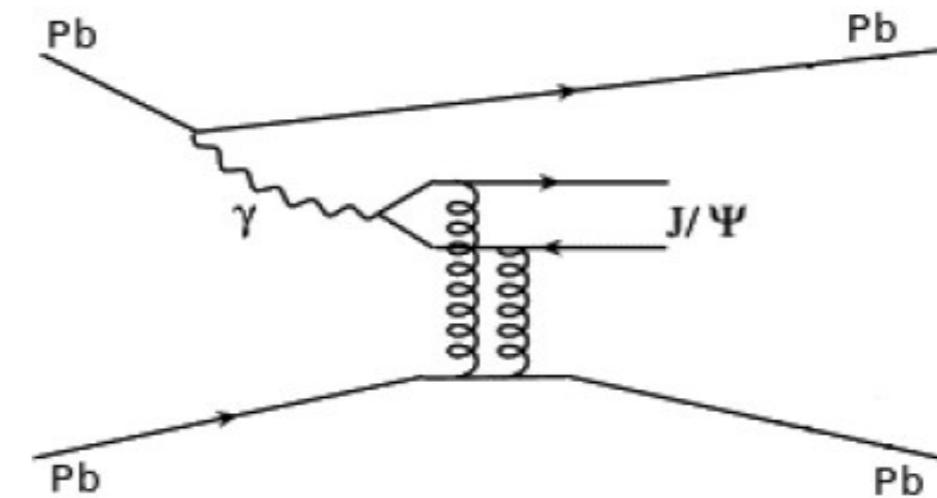


Number of photons scales like Z^2 for a single source \Rightarrow
exclusive particle production in heavy-ion collisions
dominated by electromagnetic interactions.
The virtuality of the photons $\rightarrow 1/R \sim 30 \text{ MeV}/c$



Two-photon production

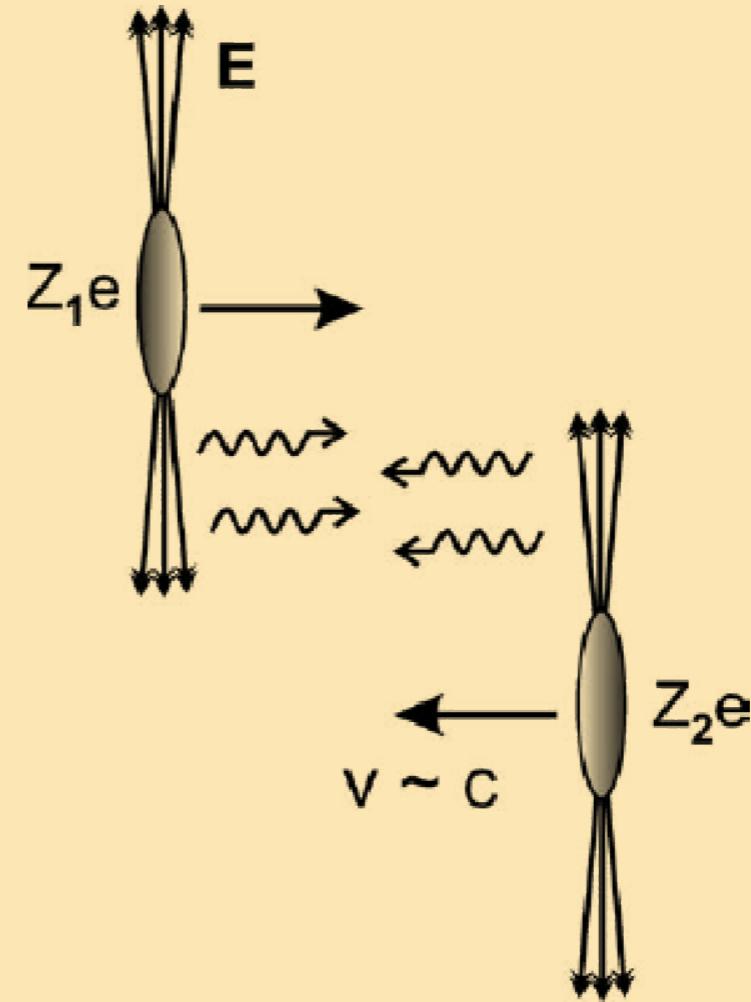
Photon-induced reactions



$\gamma + p \rightarrow J/\Psi + p$
modelled in pQCD: exchange of two
gluons with no net-colour transfer

What is an UPC?

- ▶ Range of strong interaction
 $\sim 1/m_\pi \sim 1\text{fm}$
 - ▶ Range of electromagnetic force ∞
 - ▶ Impact parameter $b > 2R_A$
 - ➡ **photon-nucleus** collision
 - ➡ Just like DIS, only $Q^2 = 0$
 - ▶ Look for **exclusive** events: A intact
 - ▶ Signature: “two muons in an otherwise empty detector”
-
- ▶ In DIS: Q^2 provides hard scale ➡ QCD perturbation theory
 - ▶ In UPC: hard scale e.g. from heavy quark ➡ quarkonia



T. Lappi
IS conference, Dec 2014

Photon-induced interactions in eA vs. in pA or AA

- Energy reach very favorable in UPC:

LHC: $W_{\gamma N} \leq 500$ GeV for γA (Pb-Pb collisions)

$W_{\gamma N} \leq 1500$ GeV for γp (p-Pb collisions)

MEIC: $W_{\gamma N} \sim 15\text{-}70$ GeV

eRHIC: $W_{\gamma N} \sim 50\text{-}100$ GeV

LHeC: $W_{\gamma N} \sim 1300$ GeV (γp); 800 GeV (γA) [$E_e = 60$ GeV]

- UPC restricted to photoproduction ($Q^2 \approx 0$) because of the Form Factor.

Why J/ ψ photo-production at LHC

Total J/ ψ cross section: 23 mb (STARLIGHT) vs 10.3 mb Rebyakova, Strikman and Zhalov

Models differ by the way photo-nuclear interaction is treated...

STARLIGHT

<http://starlight.hepforge.org>

Adeluyi and Bertulani (AB)

Phys. Rev. C 85 (2012) 044904

Goncalves and Machado (GM)

Phys. Rev C 84 (2011) 011902

Cisek, Szczerba, Schafer (CSC)

Phys. Rev. C 86 (2012) 014905

Rebyakova, Strikman and Zhalov (RSZ)

Phys. Lett. B 710 (2012) 252

Five model predictions available
- published in the last two years-

$$\left. \frac{d\sigma}{dt} \right|_{t=0} = \frac{\alpha_s^2 \Gamma_{ee}}{3\alpha M_V^5} 16\pi^3 \left[xg(x, \frac{M_V^2}{4}) \right]^2 \quad \text{Ryskin 1993}$$

$$\left. \frac{\frac{d\sigma(\gamma A \rightarrow VA)}{dt}}{\frac{d\sigma(\gamma N \rightarrow VN)}{dt}} \right|_{t=0} = \left[\frac{G_A(x, M_V^2/4)}{G_N(x, M_V^2/4)} \right]^2$$

$$\left. \frac{d\sigma_{\gamma A \rightarrow J/\psi A}}{dt} \right|_{t=0} = \frac{M_{J/\psi}^3 \Gamma_{ee} \pi^3 \alpha_s^2(Q^2)}{48\alpha_{em} Q^8} \left[xG_A(x, Q^2) \right]^2$$

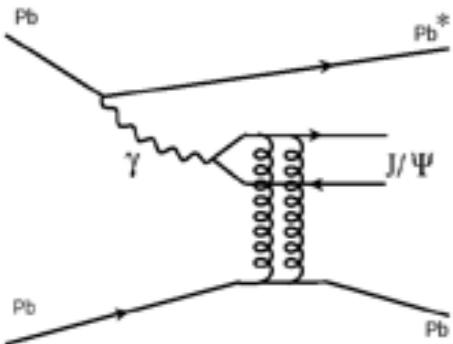
Mass of J/ ψ serves as a hard scale: $Q^2 \sim \frac{M_{J/\psi}^2}{4} \sim 2.5 \text{ GeV}^2$

Bjorken $x \sim 10^{-2} - 10^{-5}$ accessible at LHC: $x = \frac{M_{J/\psi}^2}{W_{\gamma p}^2}$

Also a more recent calculation

T. Lappi, H. Mäntysaari
<http://arxiv.org/abs/1301.4095>

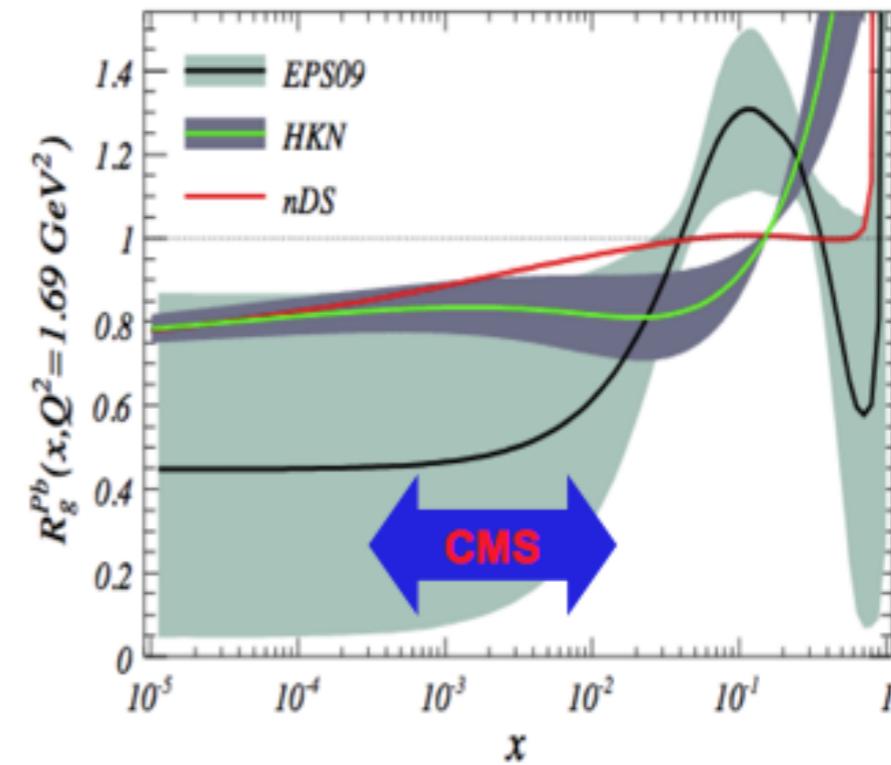
UPC Quarkonia Probe Nuclear Glue



$$\frac{d\sigma_{\gamma A \rightarrow J/\Psi A}}{dt} \Big|_{t=0} = \xi_{J/\Psi} \left(\frac{16\pi^3 \alpha_s^2 \Gamma_{l+l^-}}{3\alpha M_{J/\Psi}^5} \right) [x G_A(x, \mu^2)]^2$$

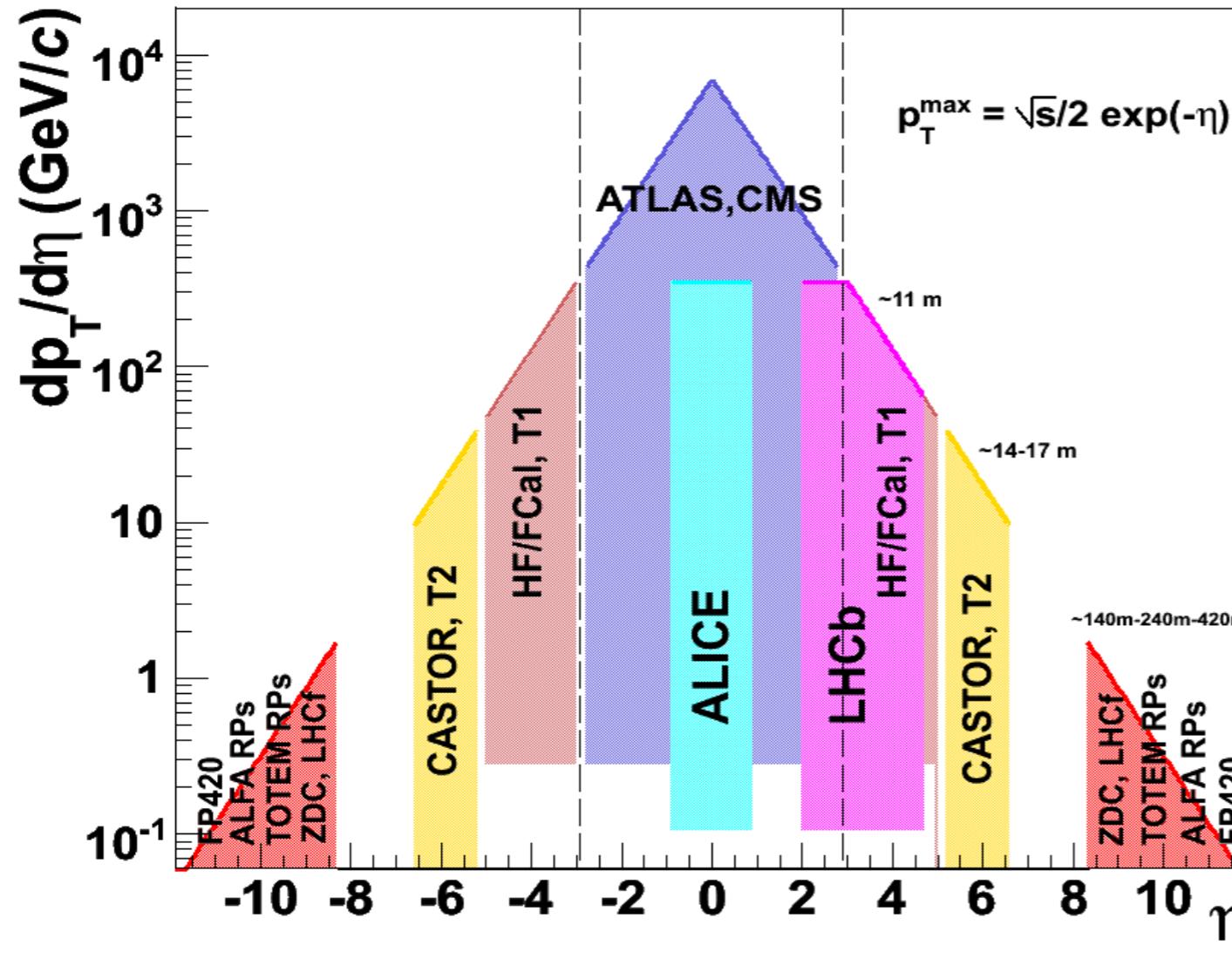
The ultra-peripheral
coherent J/ψ
photoproduction cross
section depends on the
nuclear gluon density
squared

C. A. Salgado et al 2012 J. Phys. G : Nucl. Part. Phys. 39 015010

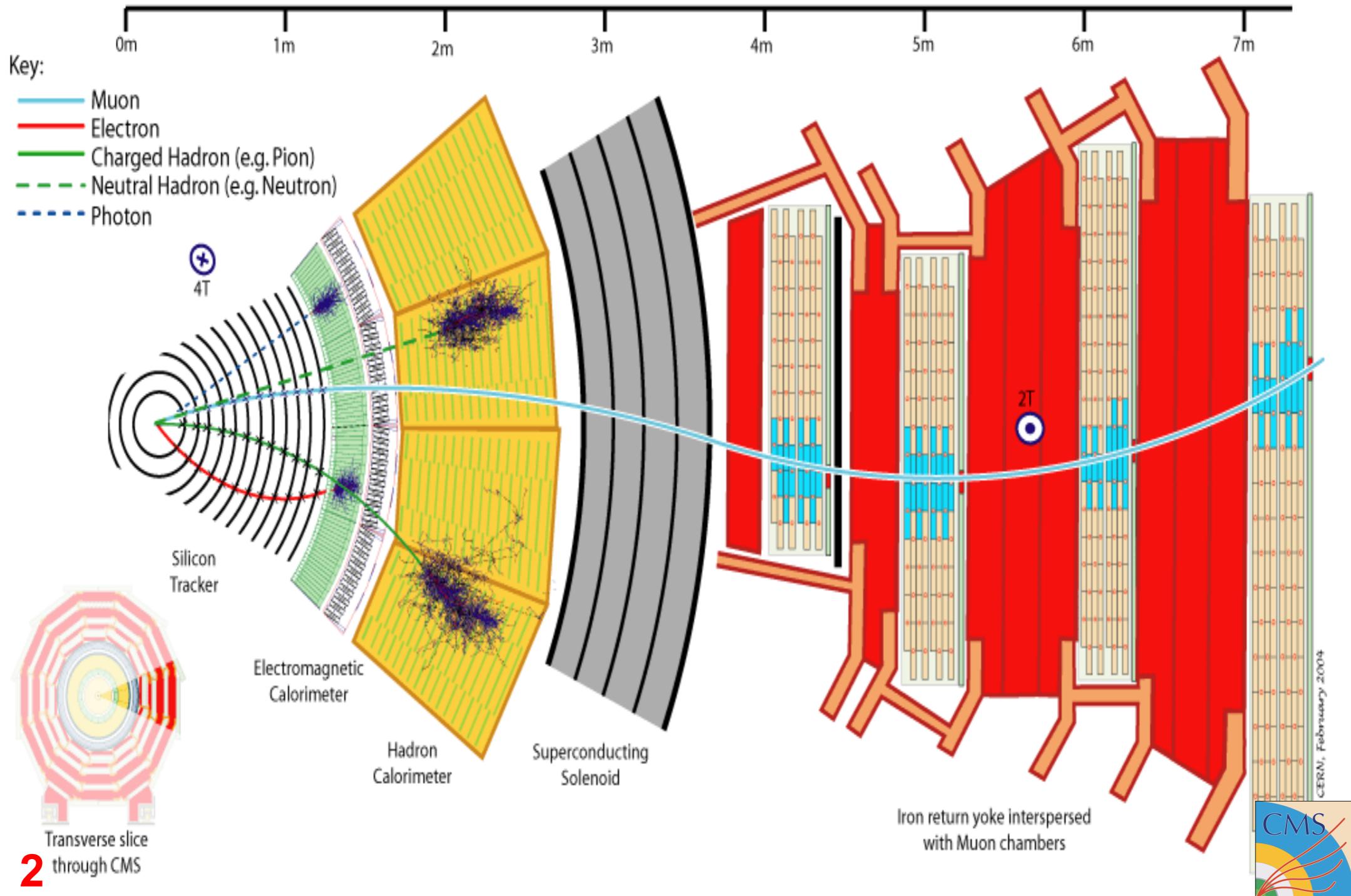


Forward physics at LHC

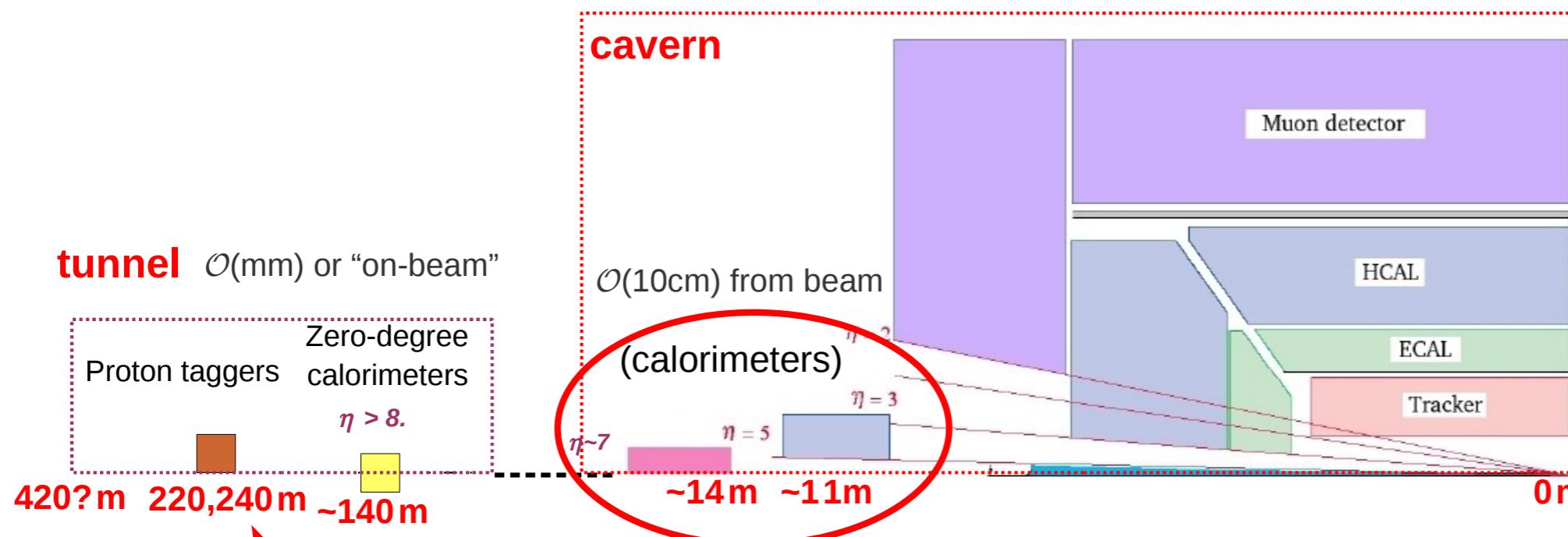
DdE, arXiv:0708.0551



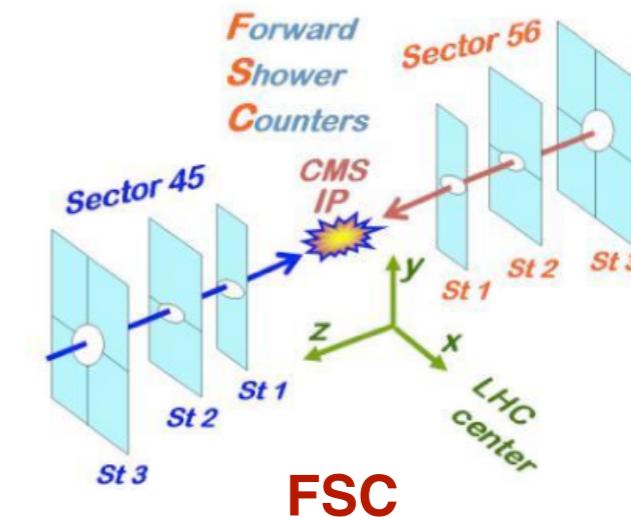
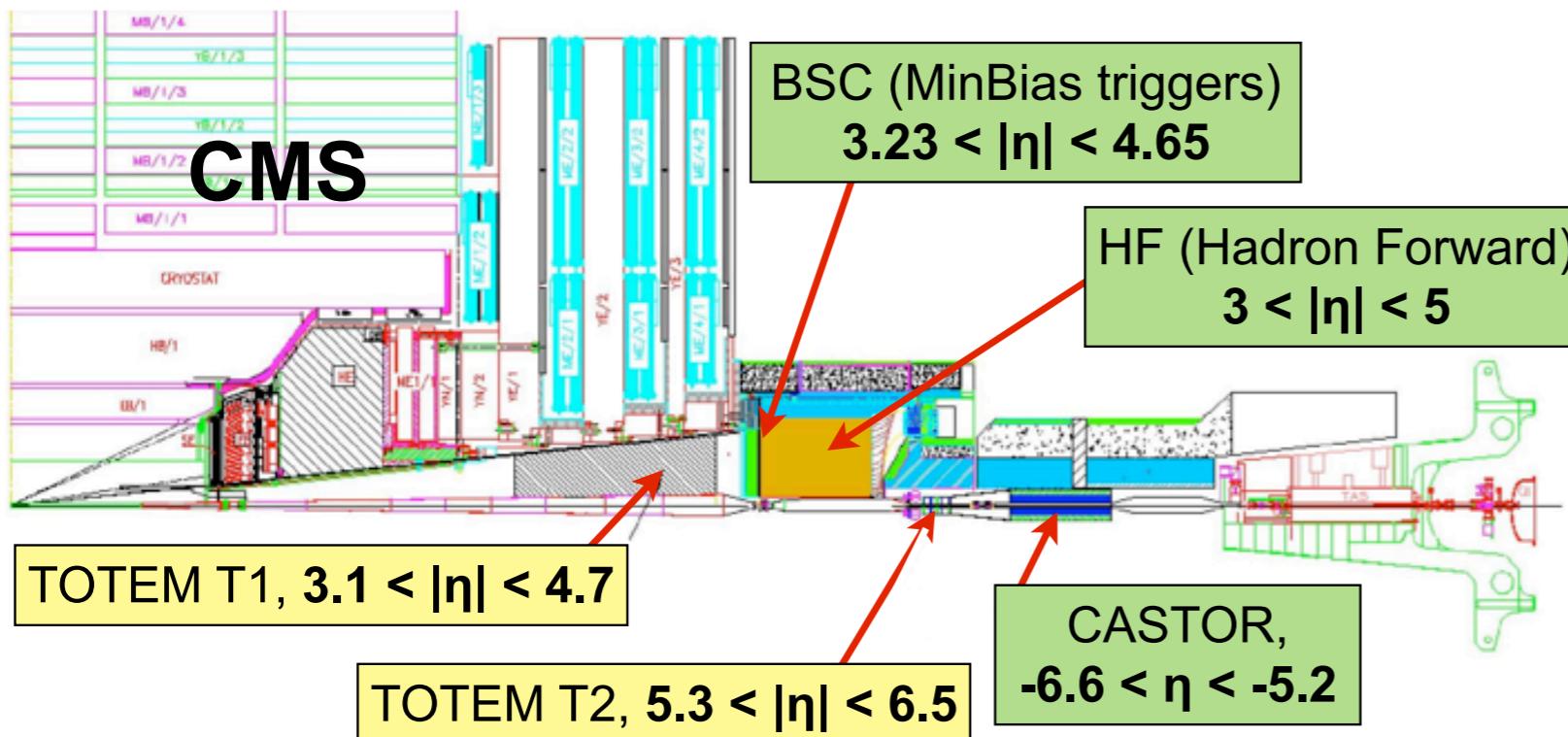
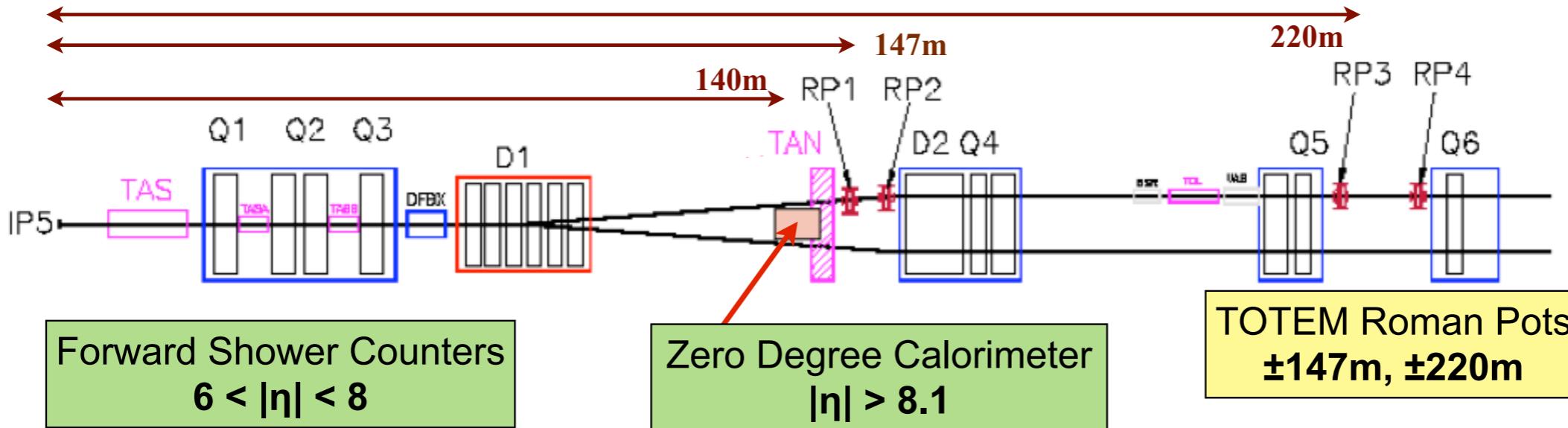
A transverse slice through CMS



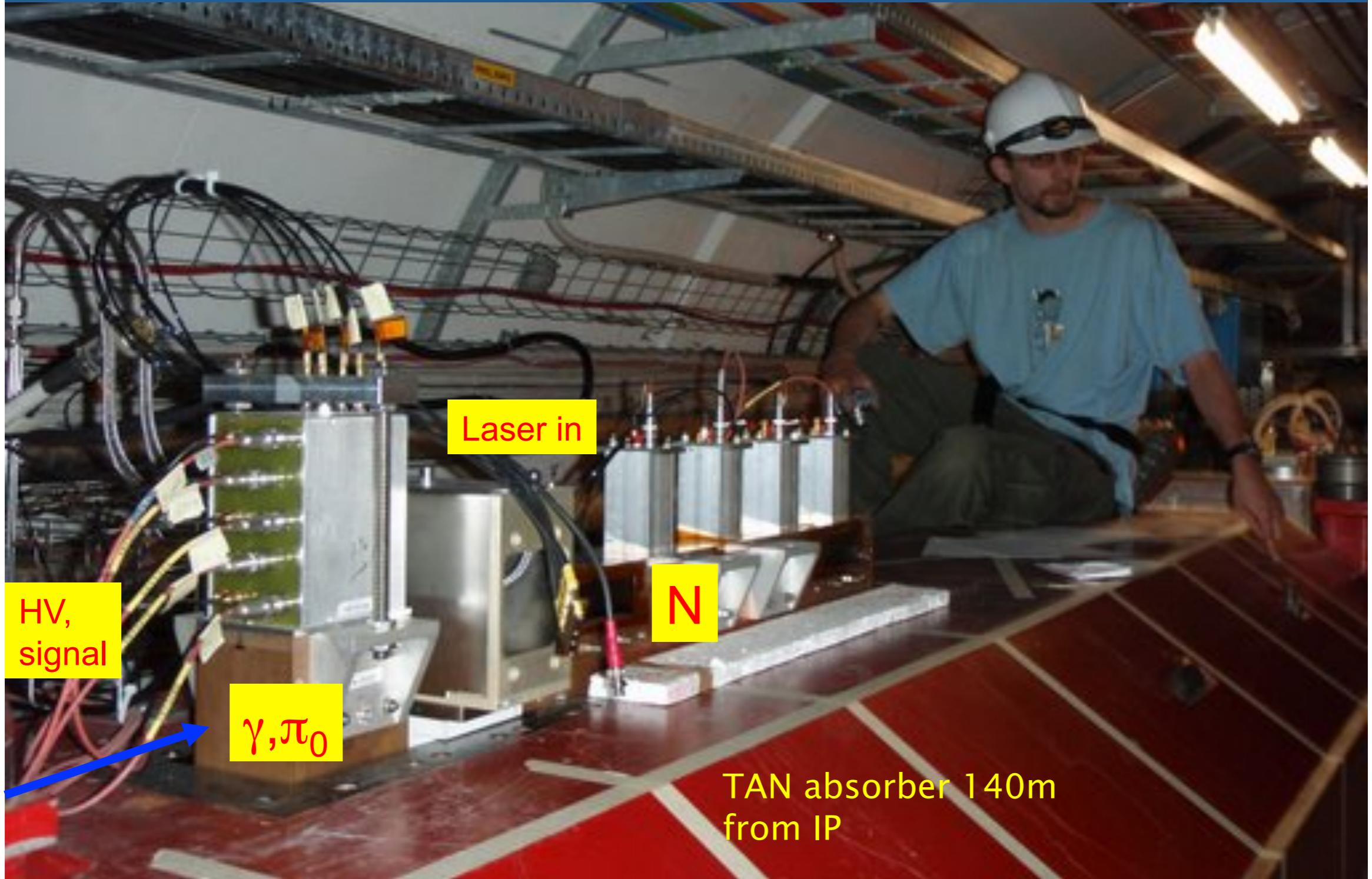
Forward physics at CMS



Forward detectors at CMS

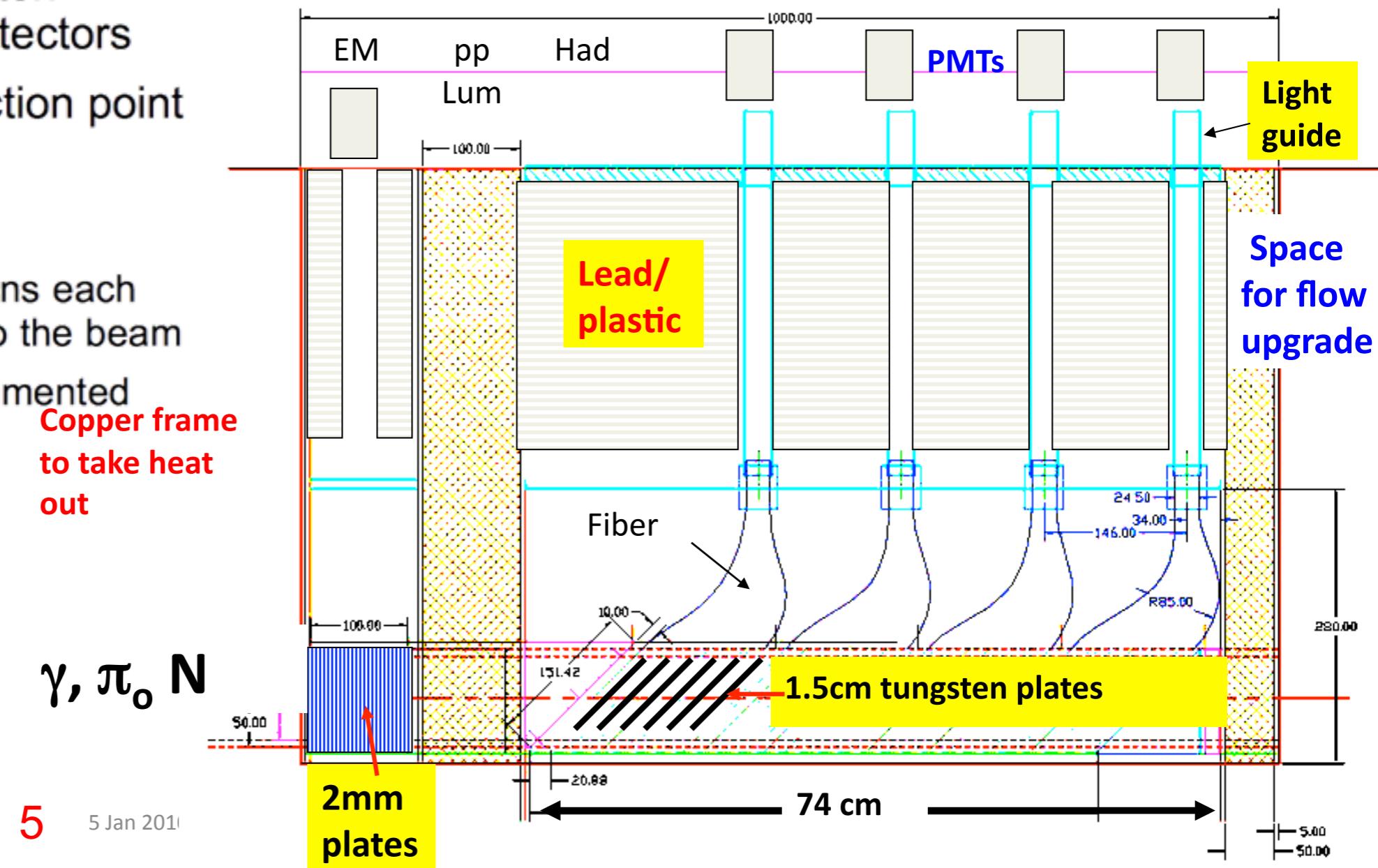


Zero Degree Calorimeters

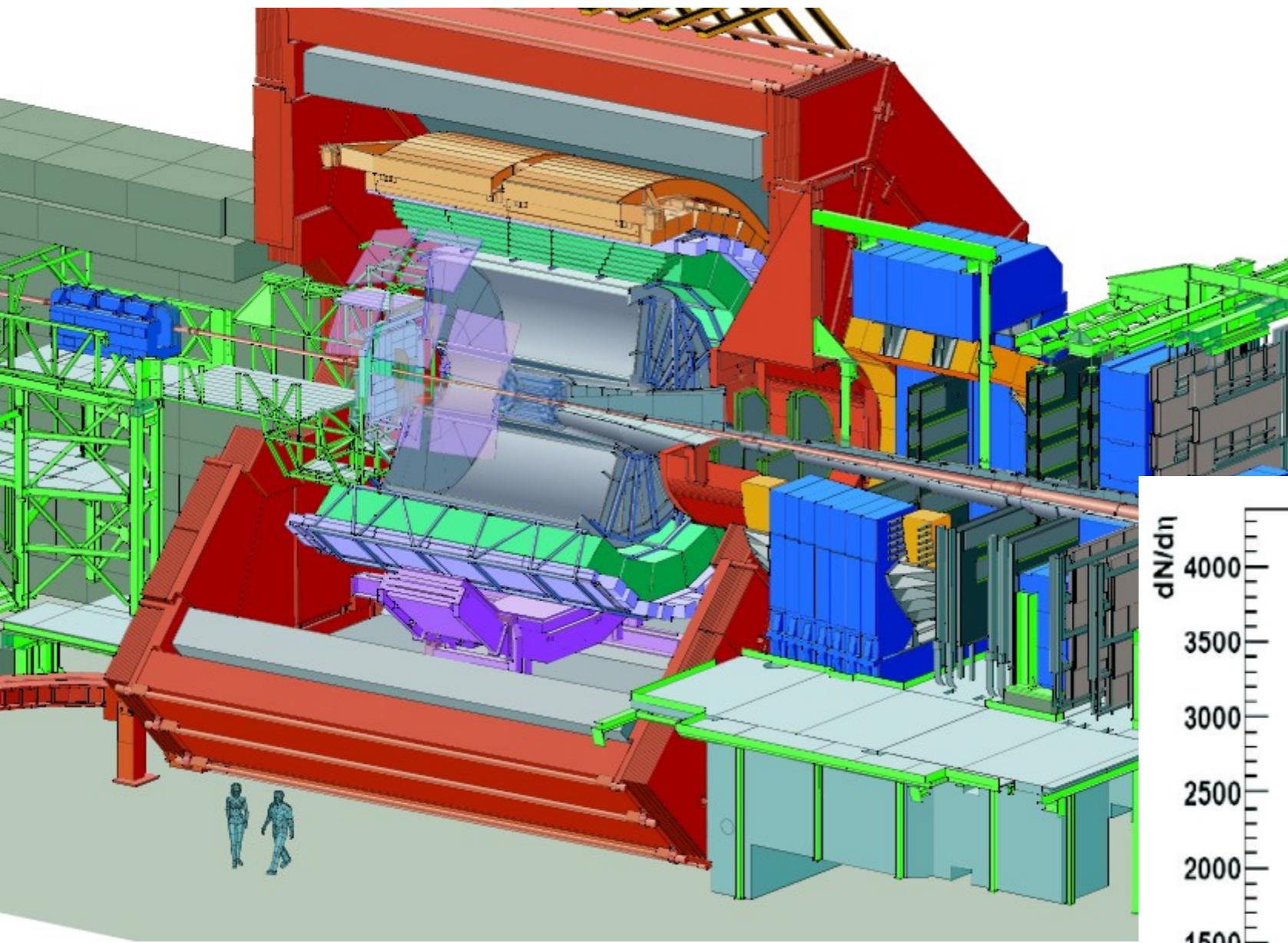


Side View of ZDC

- Quartz fibers and tungsten absorber Cherenkov detectors
- 140 meters from interaction point on either side
- Total of 18 channels
 - 5 electromagnetic sections each segmented transverse to the beam
 - 4 hadronic sections segmented longitudinally



The ALICE experiment at LHC



Dedicated triggers for UPC, using
VZERO forward detectors for vetoing
And MUON, TOF and SPD

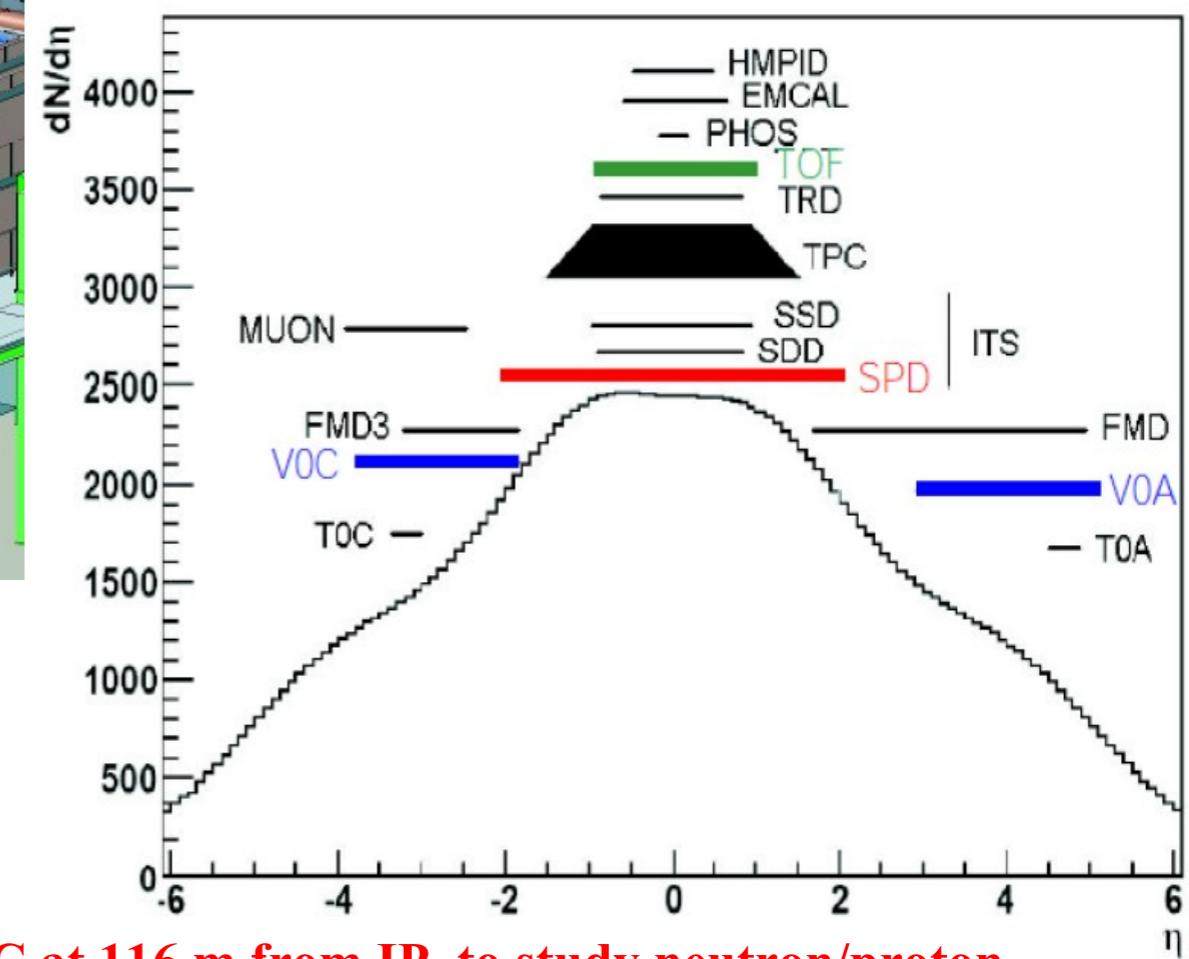
ALICE can measure J/ ψ
mesons down to zero p_T

Central rapidity

Inner Tracking (ITS), Time
Projection Chamber (TPC),
Time-of-Flight, TRD, EMCAL
 $|\eta| < 0.9$

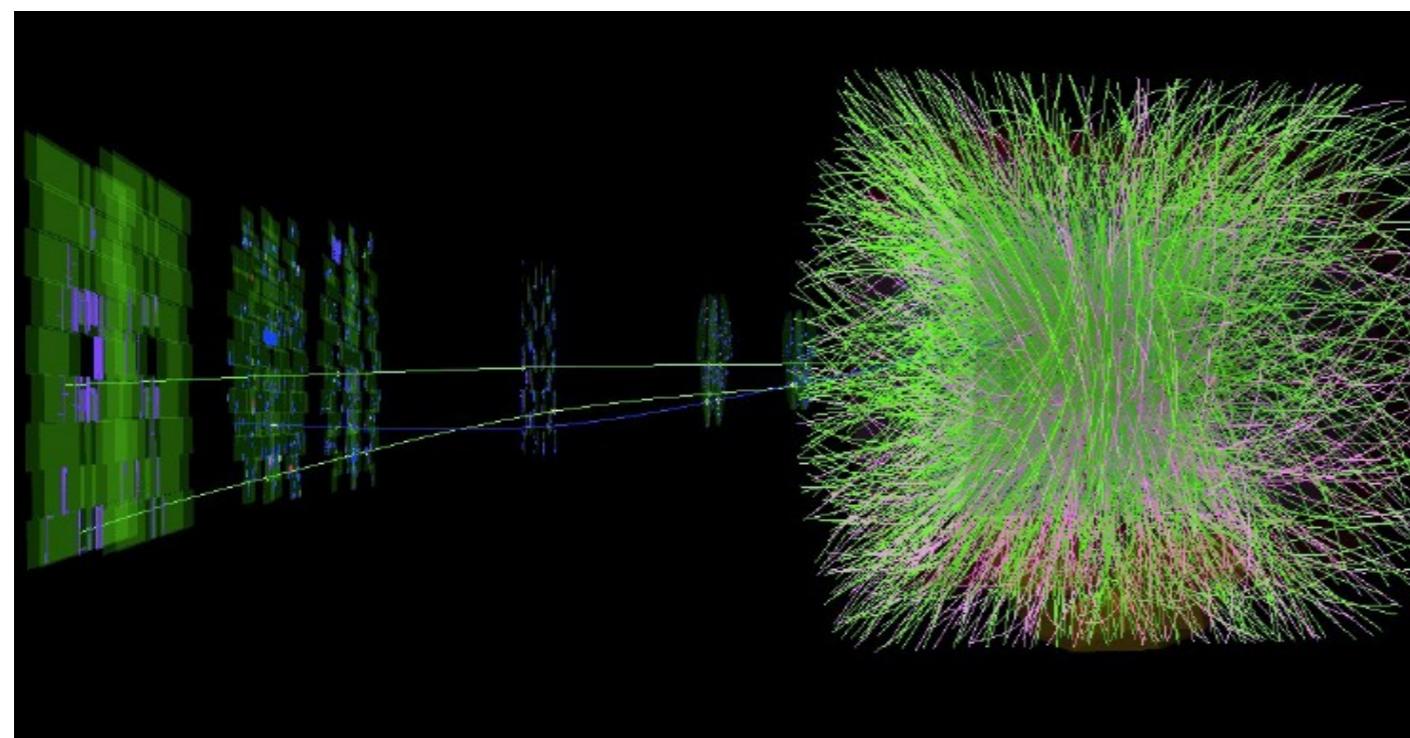
Forward rapidity

Muon Spectrometer
 $-4 < \eta < -2.5$



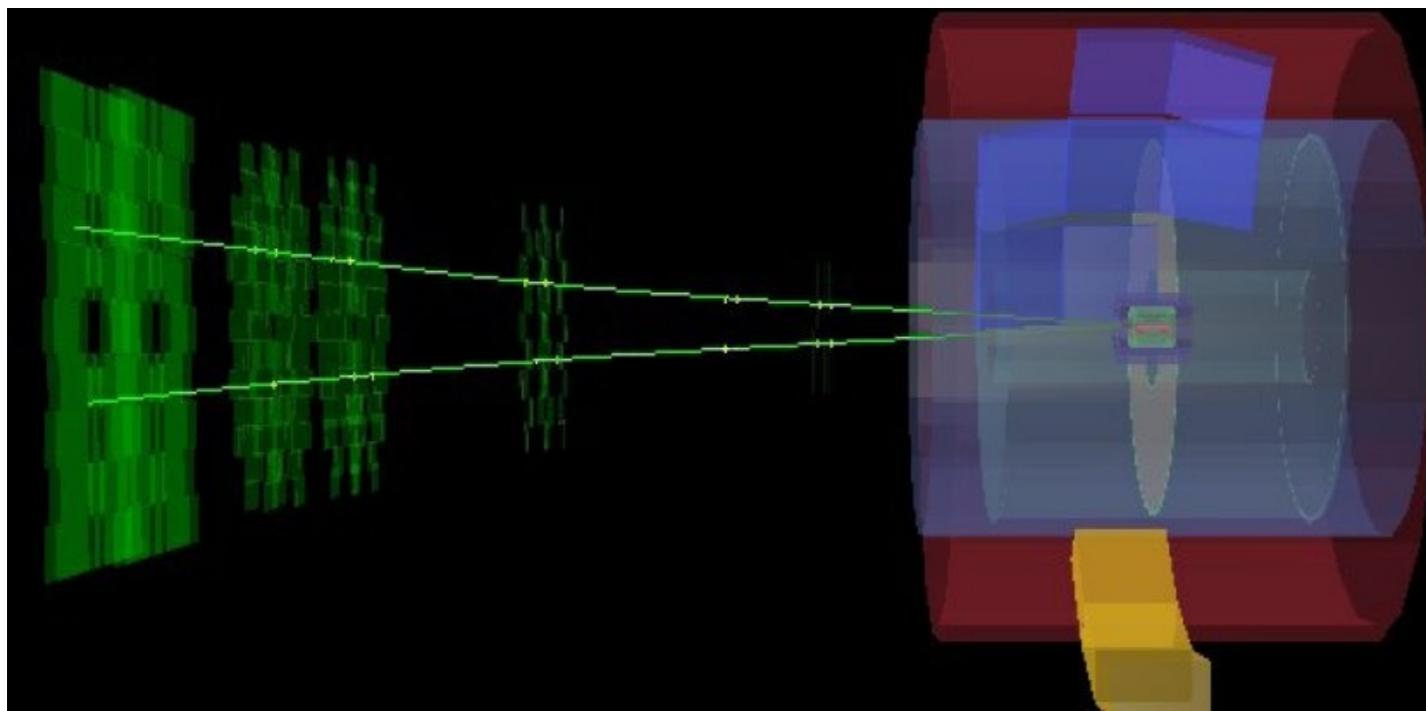
ZDC at 116 m from IP, to study neutron/proton
emitted at the very forward region

Exclusive J/ ψ analysis at forward rapidity



From a typical inclusive
J/ ψ candidate in
Pb-Pb collisions...

....to an exclusive J/ ψ
candidate

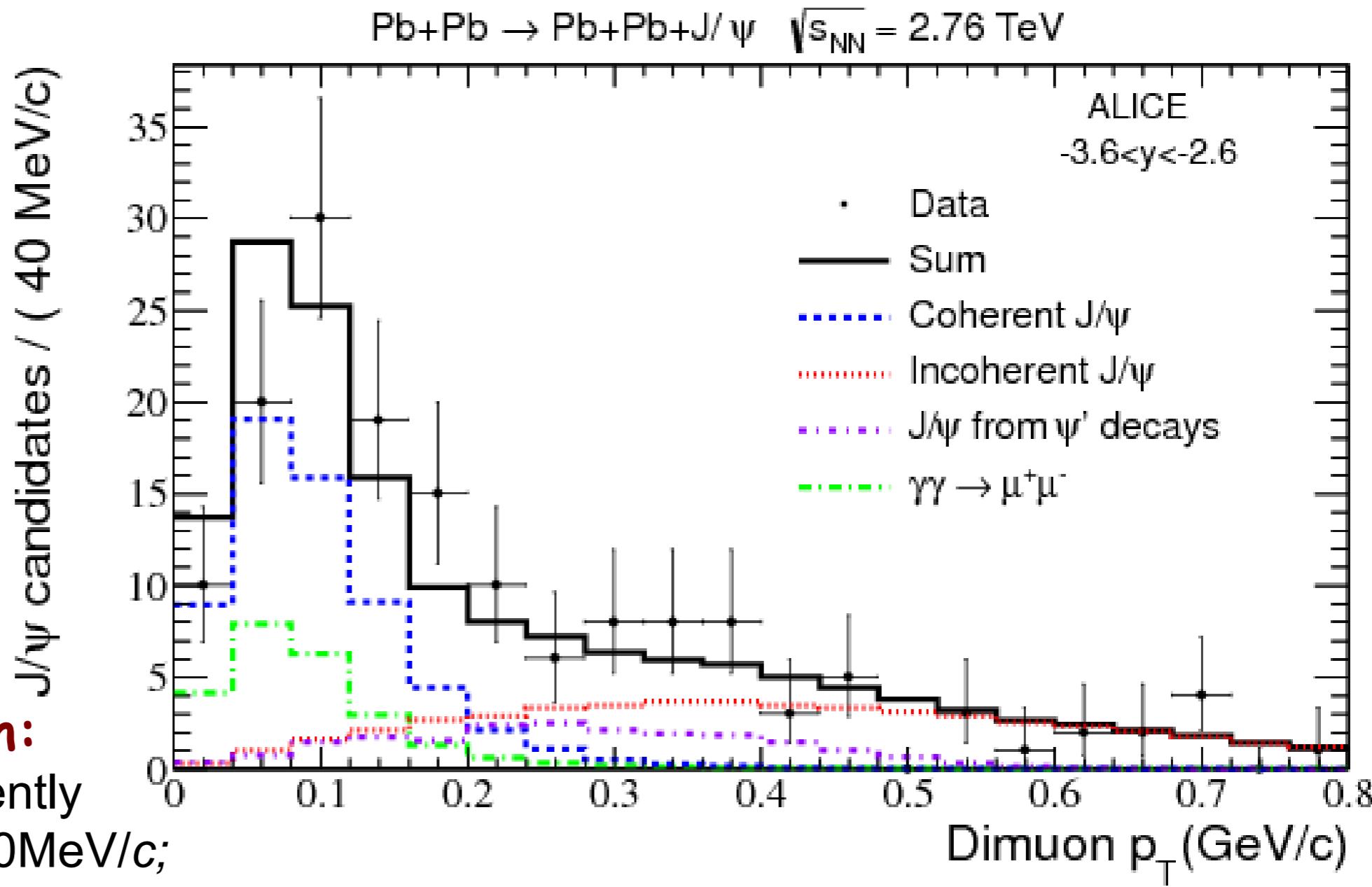


Two UPC publications
by ALICE

Phys.Lett. B718 (2013) 1273-1283

Eur. J. Phys. C73, 2617 (2013)

p_T distribution for J/ ψ candidates

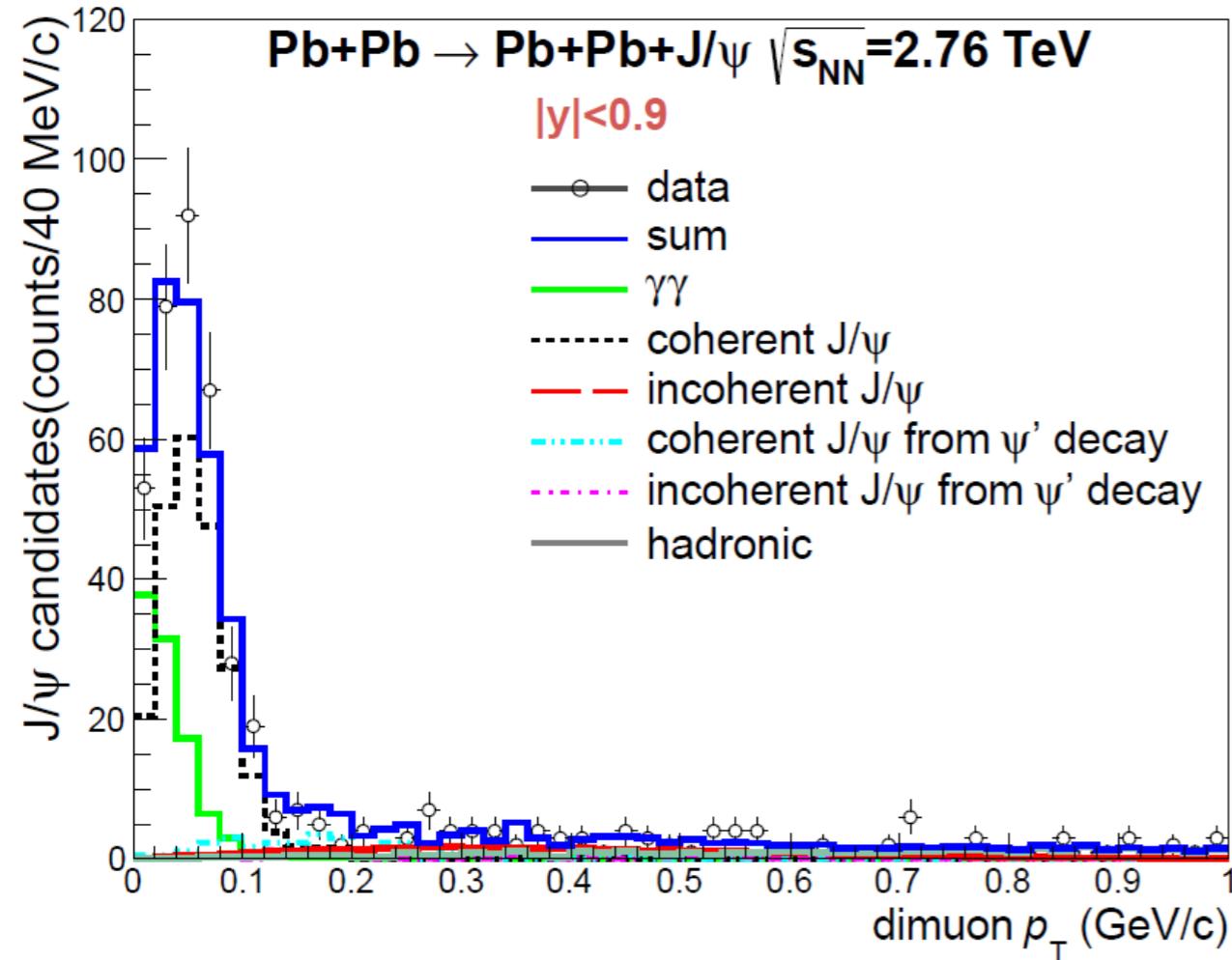


Four physics processes:

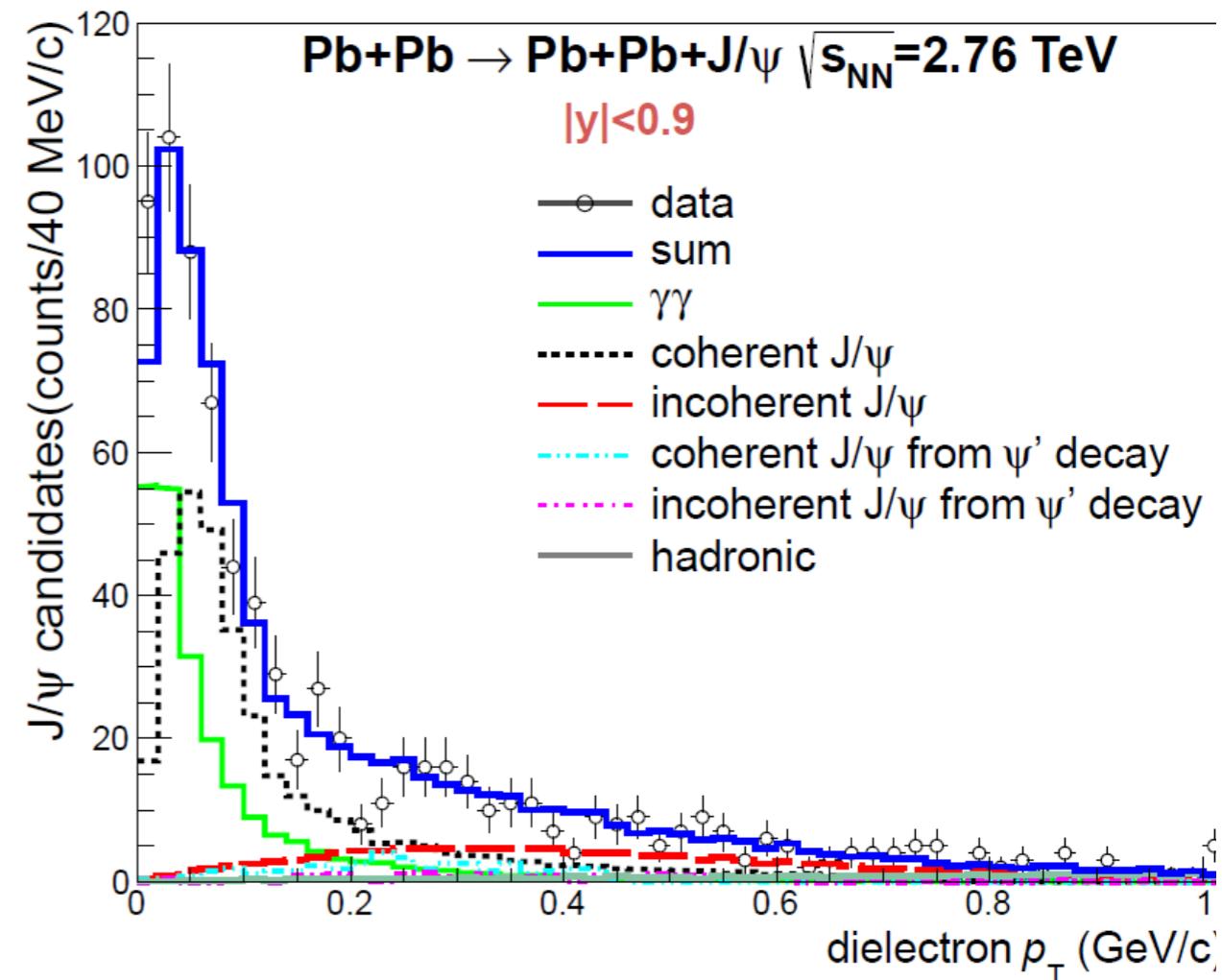
- Coherent J/ ψ
- Incoherent J/ ψ
- J/ ψ from ψ' decays
- $\gamma\gamma \rightarrow \mu^+\mu^-$

Central barrel measurements in UPC

J/ ψ in the dimuon channel



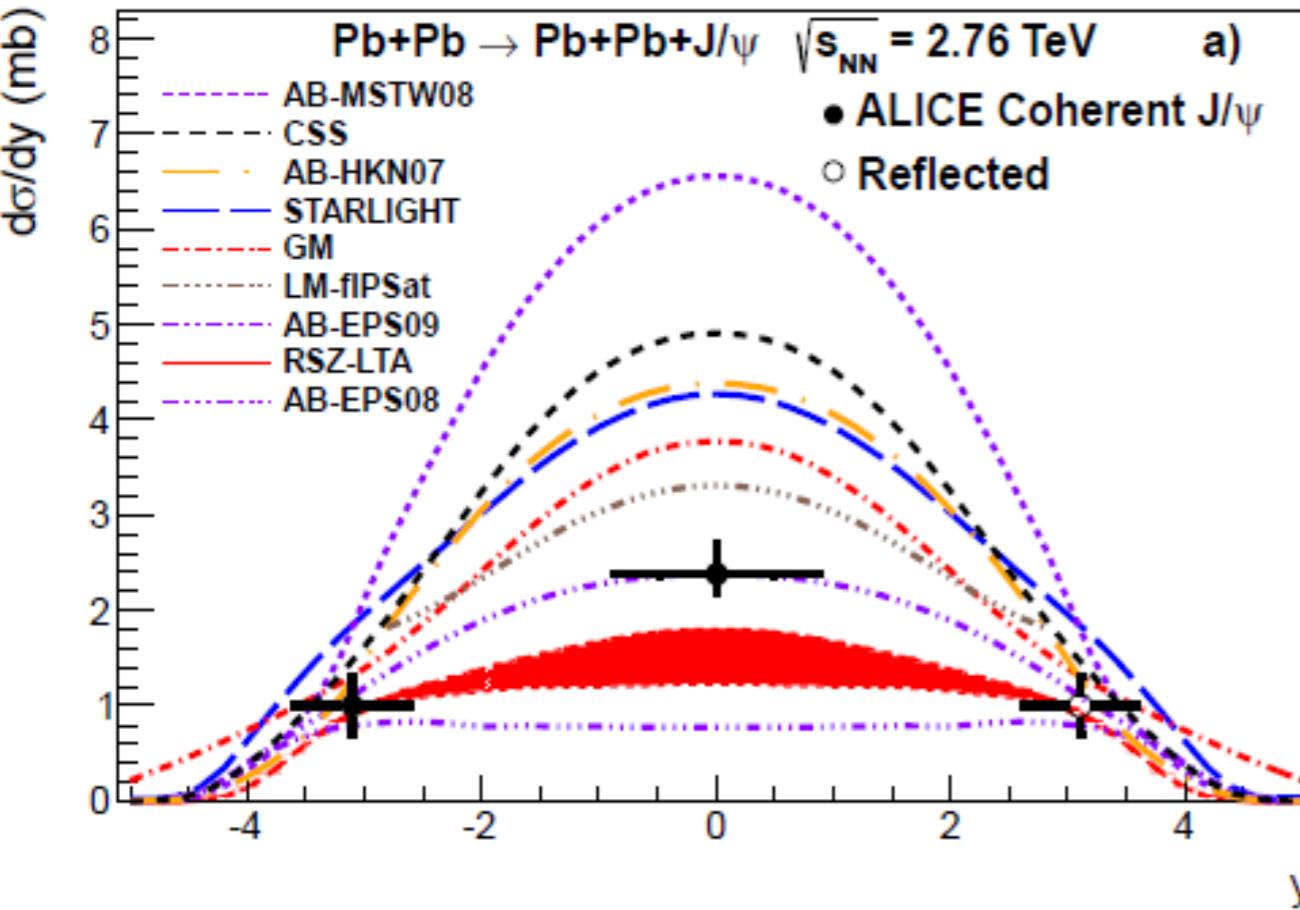
J/ ψ in the dielectron channel



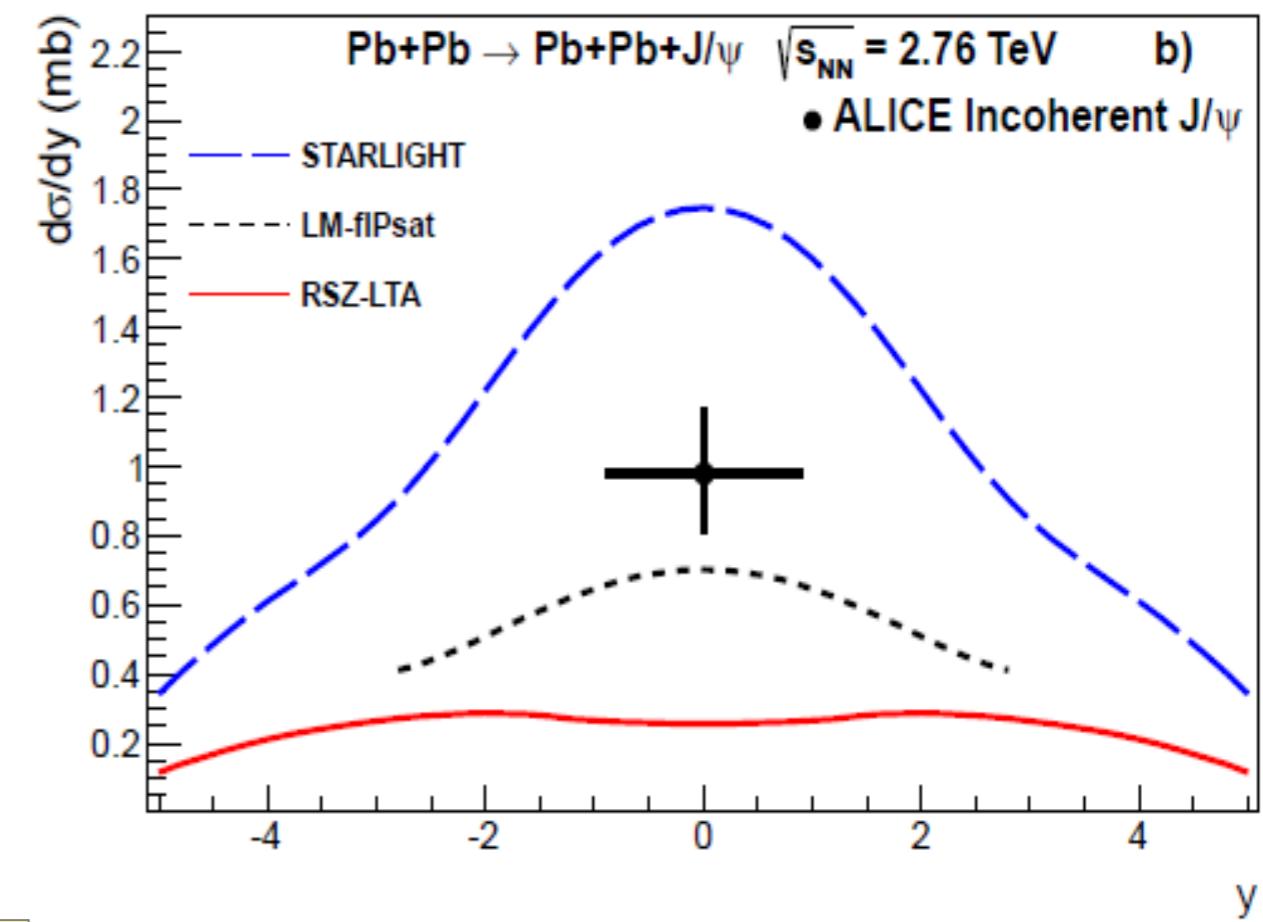
Data is well described by signals/backgrounds expected in UPC

Data and theoretical predictions

Coherent J/ ψ



Incoherent J/ ψ



Direct evidence of nuclear gluon shadowing

At mid-rapidity, Bjorken- $x \sim 10^{-3}$

ALICE shows that the distribution in $x \approx 10^{-2} - 10^{-3}$ range is consistent with the EPS09 parameterization

Two UPC publications by ALICE

Phys.Lett. B718 (2013) 1273-1283

Eur. J. Phys. C73, 2617 (2013)

Nuclear gluon shadowing from ALICE data

V. Guzei, E. Kryshen, M. Strikman, M. Zhalov. Phys. Lett. B726 (2013) 290

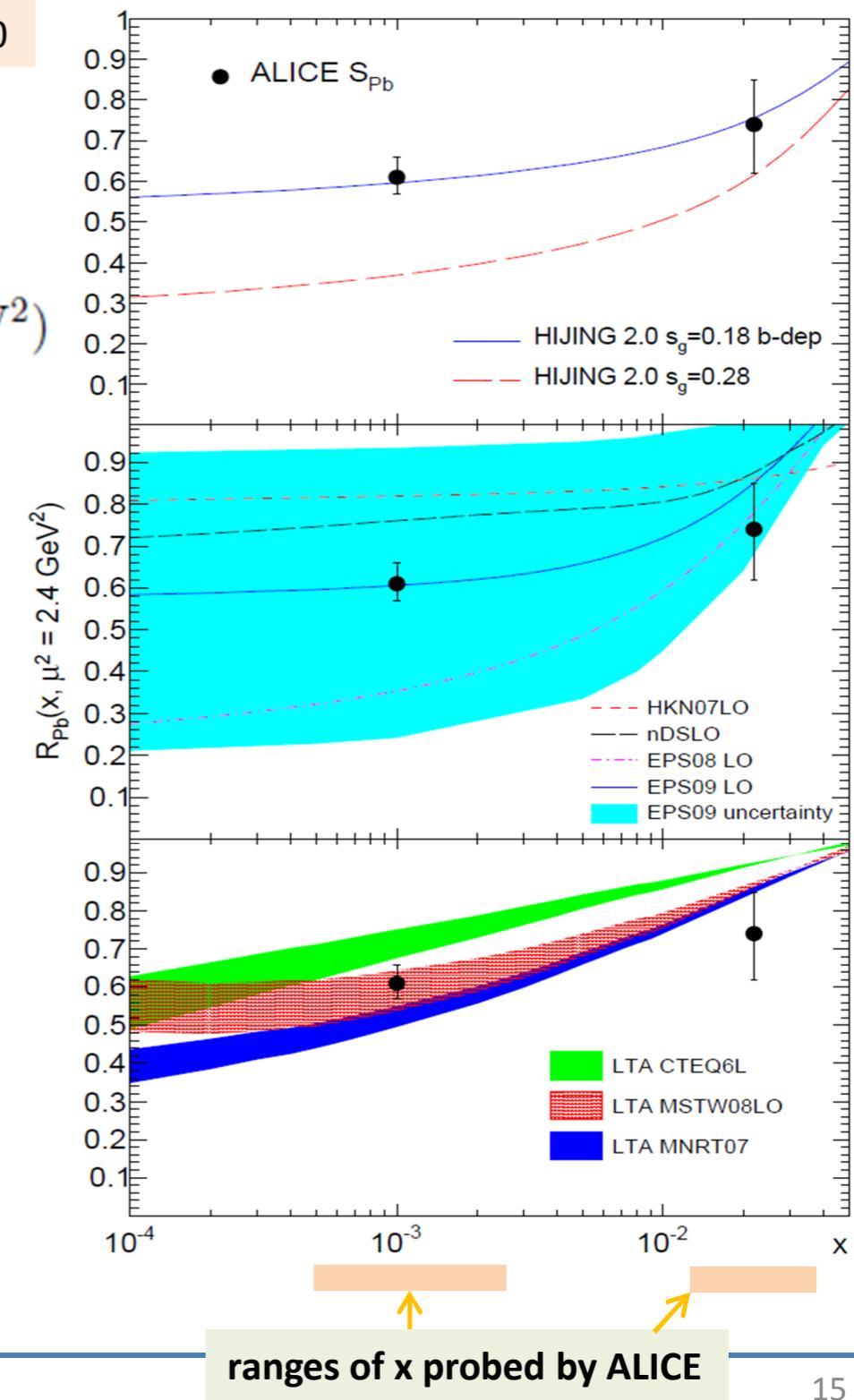
Nuclear suppression factor in J/ψ photoproduction:

ALICE data corrected for photon flux

$$S(W_{\gamma p}) \equiv \left[\frac{\sigma_{\gamma \text{Pb} \rightarrow \text{J}/\psi \text{Pb}}^{\text{exp}}(W_{\gamma p})}{\sigma_{\gamma \text{Pb} \rightarrow \text{J}/\psi \text{Pb}}^{\text{IA}}(W_{\gamma p})} \right]^{1/2} \rightarrow R(x, \mu^2 = 2.4 \text{ GeV}^2)$$

Impulse Approximation: J/ψ photoproduction cross section from HERA corrected for the integral over squared Pb form-factor

- **Hijing:** scale-independent gluon shadowing, characterized by parameter s_g
- **Shadowing parametrizations (EPS,nDS,HKN07)** use DIS and Drell-Yan data + π^0 data from RHIC (EPS) – gluon shadowing essentially unconstrained at low x
- **Leading twist approximation:** propagation of color dipoles in nuclei via intermediate diffractive states (Gribov-Glauber shadowing theory). Incorporates diffractive parton distributions in proton (from HERA)



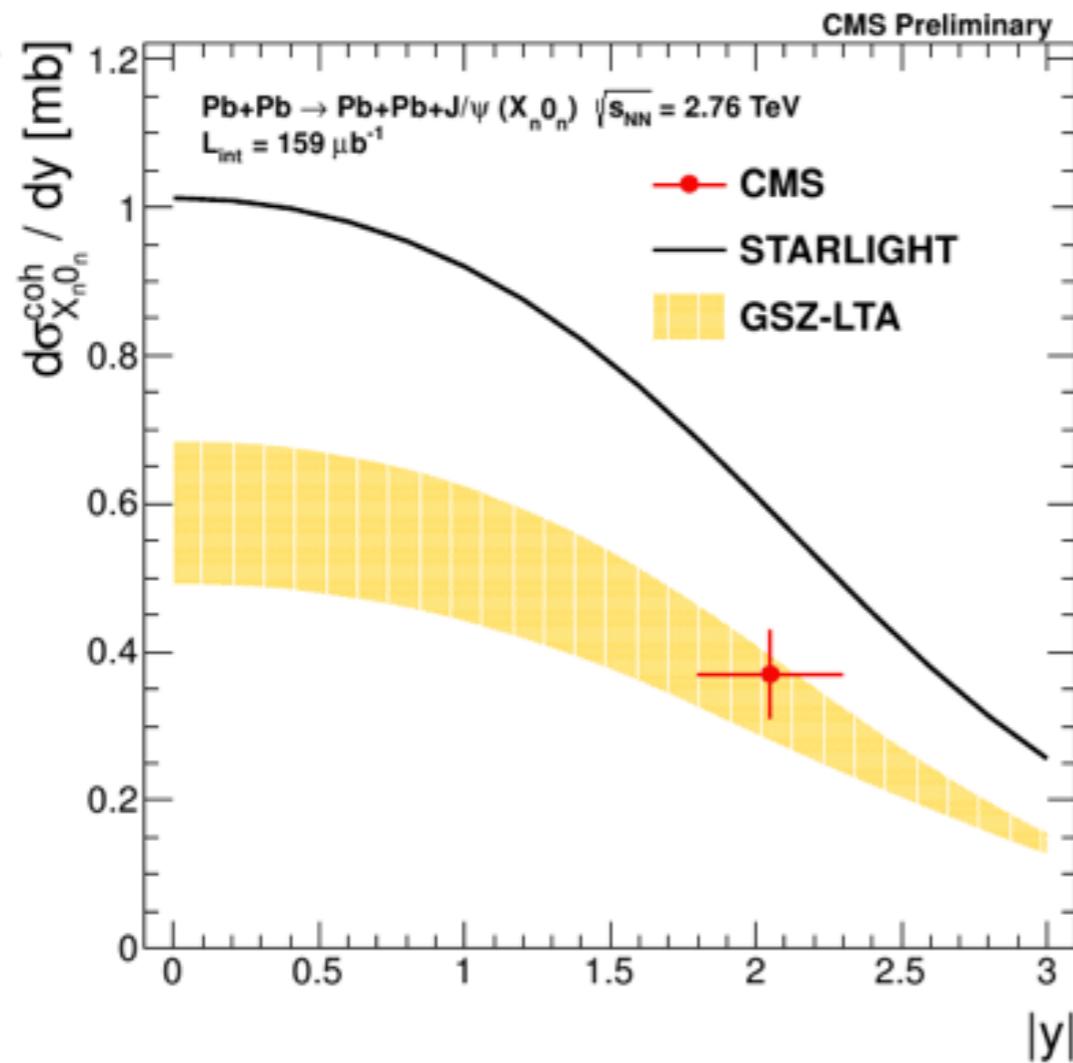
Evgeny Kryshen

ranges of x probed by ALICE

15

Coherent J/y Cross Section in PbPb UPC

NEW



- Cross section measured for events with single sided neutron emission, the $X_n 0_n$ break-up mode
- $X_n 0_n$ is the largest cross section available given ZDC trigger requirement

$$ds/dy(\text{coh}/X_n 0_n) = 0.37 \pm 0.04(\text{stat}) \pm 0.04(\text{syst}) \text{ mb}$$

P. Kenny
IS conference. Dec 2014



Break-up Modes Ratios

J/ ψ with $p_T < 0.15 \text{ GeV}/c$	$X_n X_n / X_n 0_n$	$1_n 0_n / X_n 0_n$	$1_n 1_n / X_n 0_n$
Data	0.36 ± 0.04	0.26 ± 0.03	0.03 ± 0.01
STARLIGHT	0.37	N/A	0.02
GSZ	0.32	0.30	0.02

First measurement
of break-up ratios
for UPC J/ ψ

$X_n 0_n$ single-sided neutron emission with any number of neutrons

$X_n X_n$ double-sided neutron emission with any number of neutrons

$1_n 0_n$ single-sided neutron emission with only one neutron

$1_n 1_n$ double-sided neutron emission with only one neutron on each side

The multiple photon-exchange model of nuclear break-up in coherent interactions describes the data reasonably well

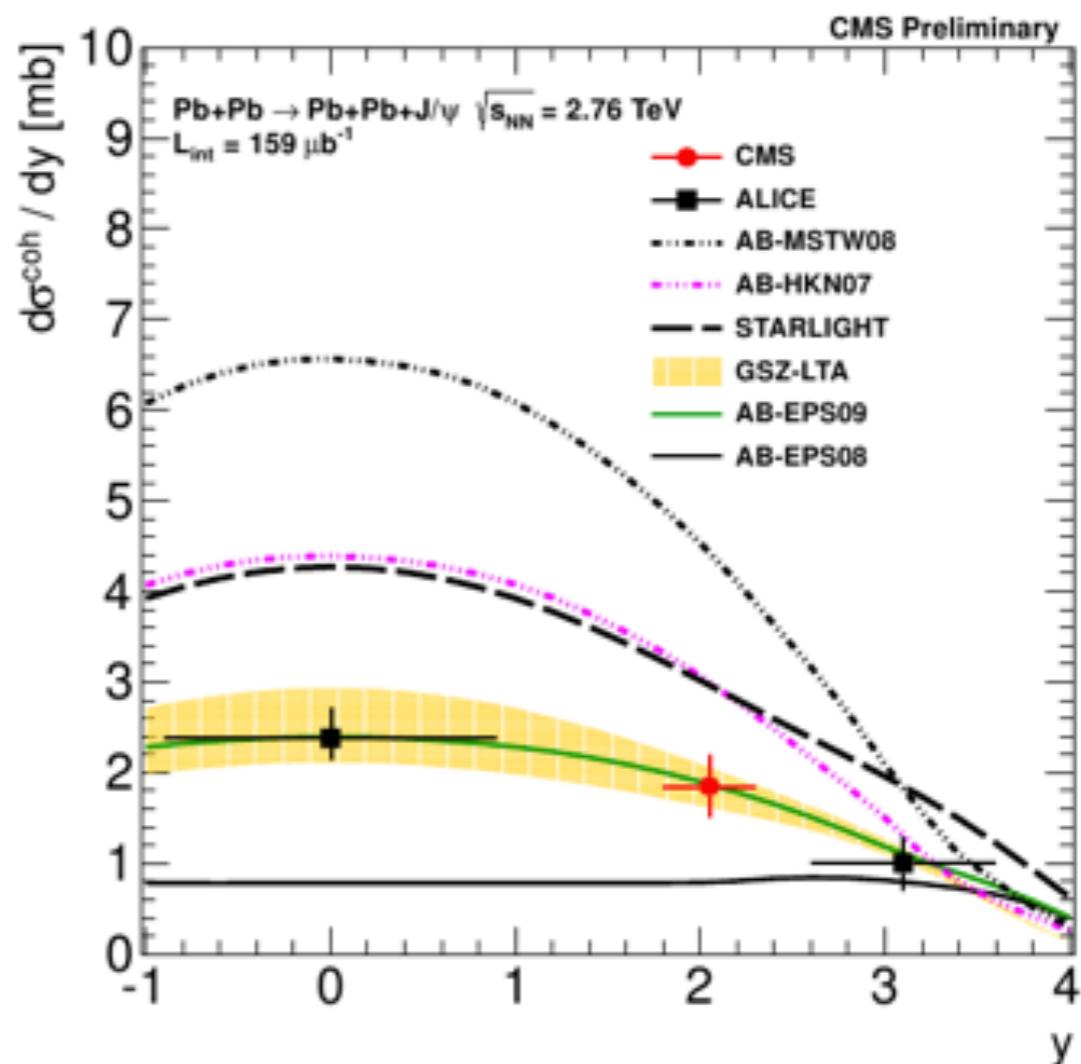
P. Kenny
IS conference. Dec 2014

Coherent J/y Cross Section in PbPb UPC

NEW

- Cross section for X_{n_n} is scaled up to the total cross section using STARLIGHT
- CMS and ALICE results favor the same theoretical models
- ALICE and CMS measurements favor models containing moderate gluon shadowing

CMS: HIN-12-009: <http://cds.cern.ch/record/1971267>



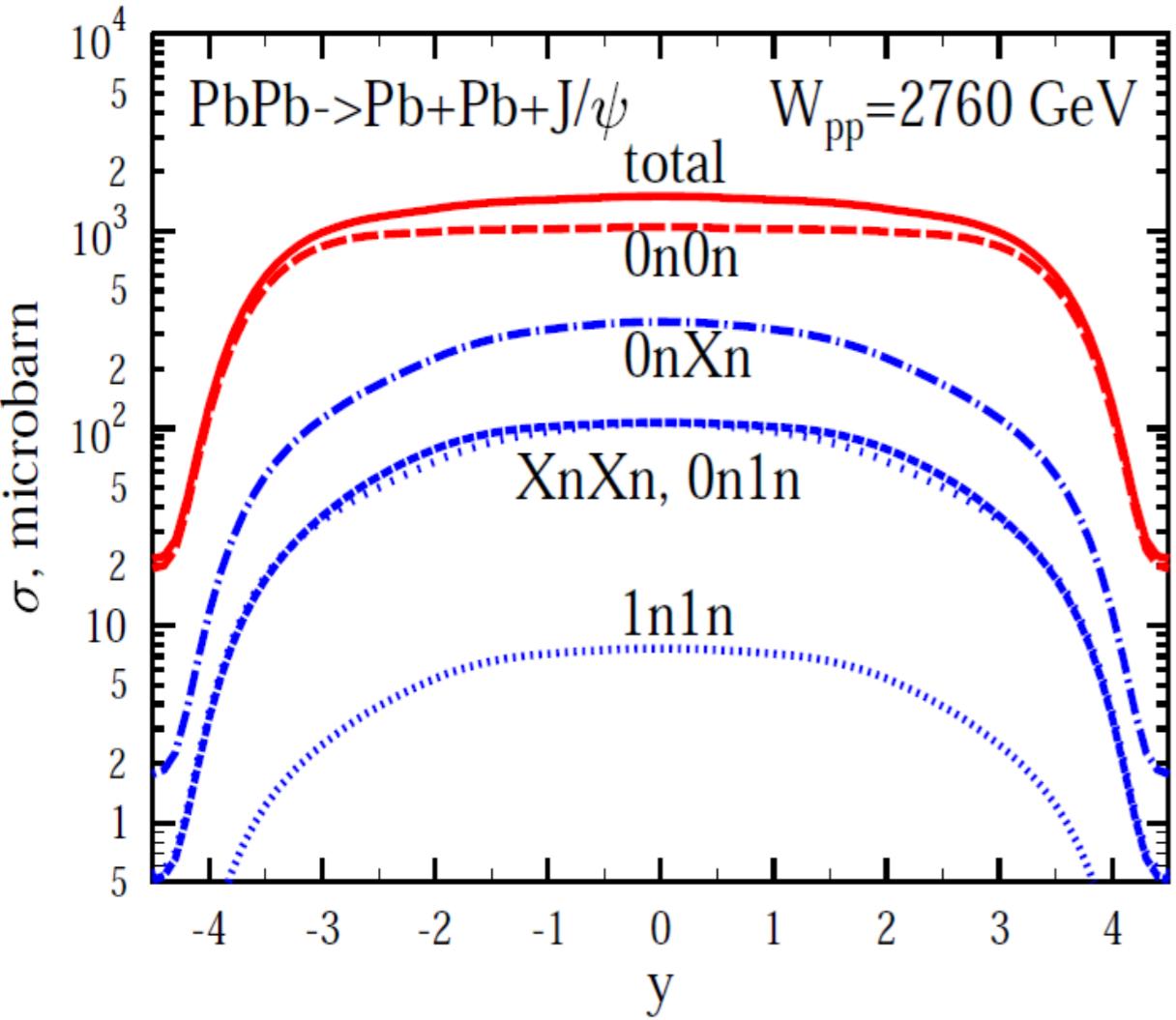
ALICE: [Eur.Phys.J. C73 \(2013\) 2617](http://cds.cern.ch/record/1971267)

[Phys.Lett. B 718 \(2013\) 1273-1283](http://cds.cern.ch/record/1971267)

P. Kenny
IS conference, Dec 2014

New at the LHC: Dependence on neutron emission

Using Zero Degree Calorimeters (ZDC) it is possible to select coherent production with ion excitation, where neutrons are emitted from at least one of the nuclei



Different configurations:

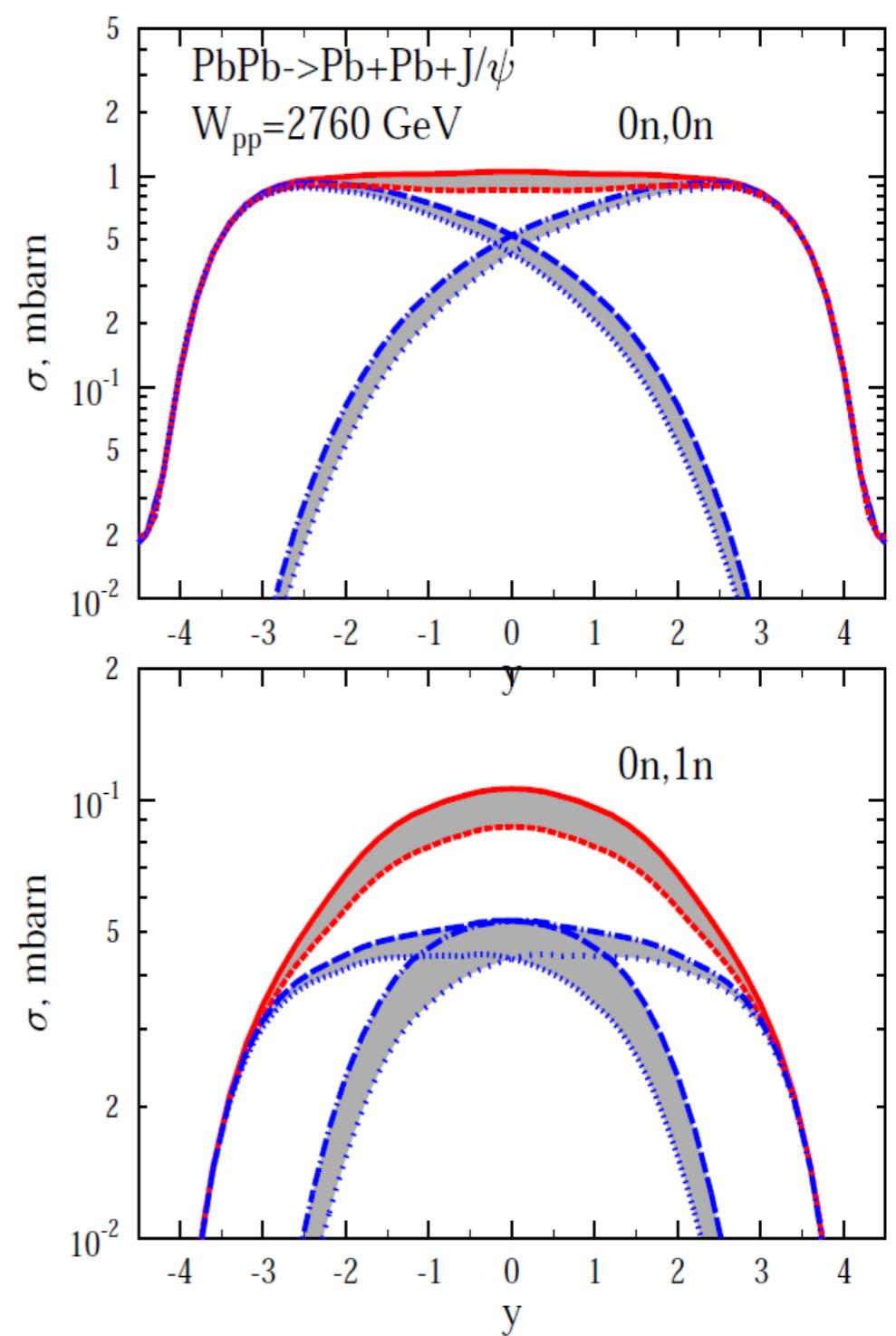
1n1n: one neutron emission by each ion;

XnXn: emission of several neutrons;

0n1n and 0nXn: excitation and decay of one of the ions, and

0n0n: no neutron emission

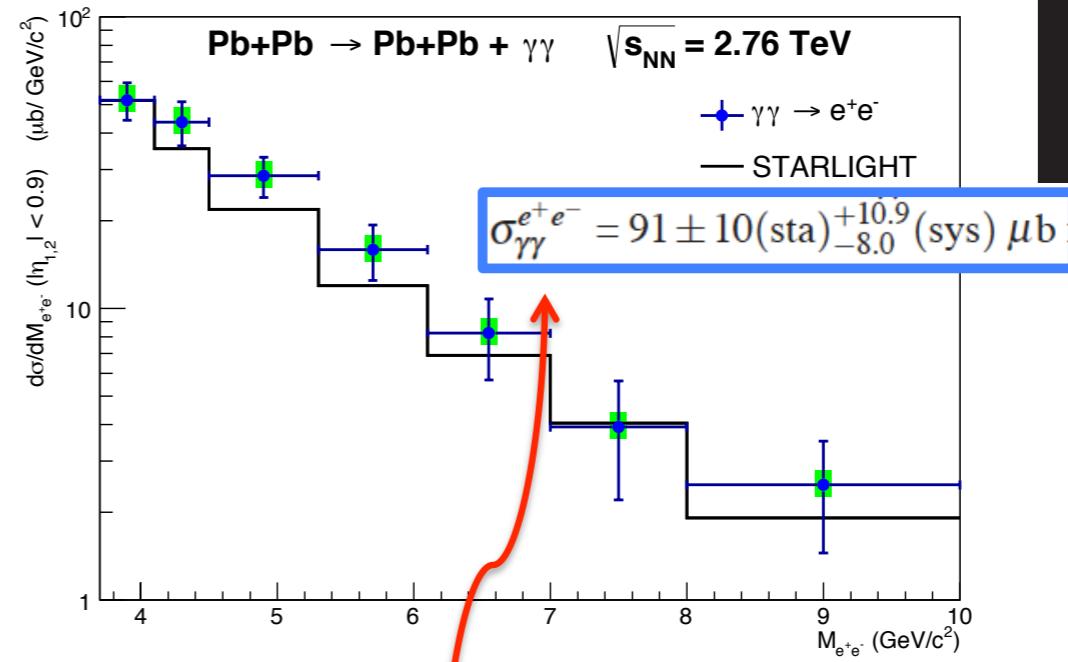
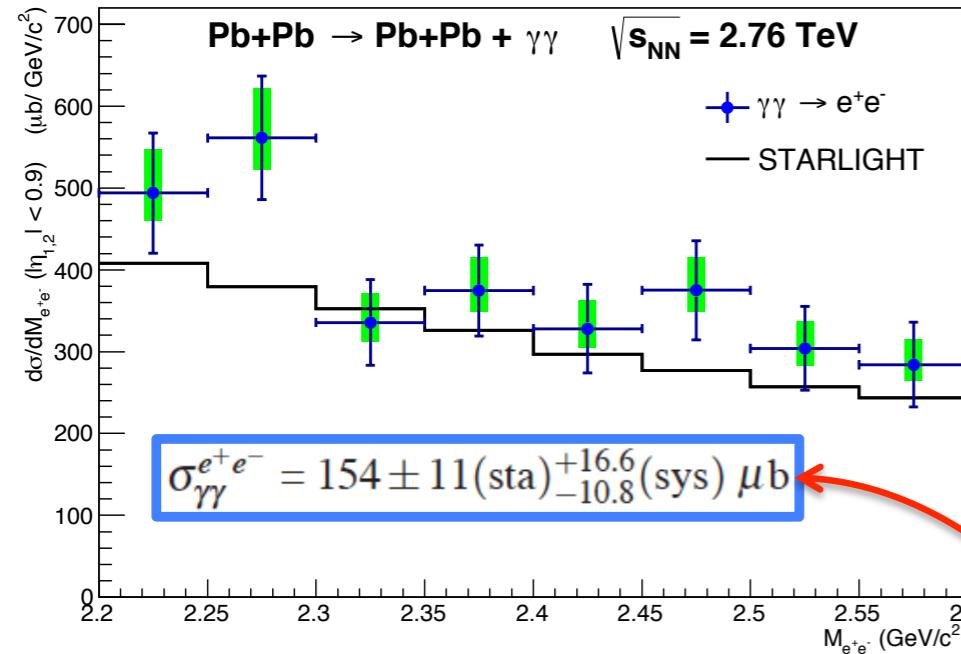
V. Rebyakova, M. Strikman and M. Zhalov
ArXiv:1109.0737, Sept 2011



Shaded area: Uncertainty on nuclear gluon shadowing

$\gamma + \gamma \rightarrow e^+ + e^-$ production in Pb-Pb (Central Barrel)

Eur. Phys. J. C (2013) 73:2617



- ✓ QED process ... but uncertainties due to
 - Higher order corrections because the coupling is enhanced by a factor of Z
 - Nuclear form factor and the minimum momentum transfer in the interaction
- Different models predict a **reduction** of the LO cross section up to 30%
- (see for example: A. J. Baltz, Phys. Rev. C 80 (2009) 034901; Phys. Rev. Lett. 100 (2008) 062302)

- ✓ Measurement in two different mass ranges: [2.2, 2.6] and [3.7, 10] GeV/c^2
- ✓ Precision of 12% and 16% respectively
- ✓ Data slightly **above** STARLIGHT, a **LO prediction**

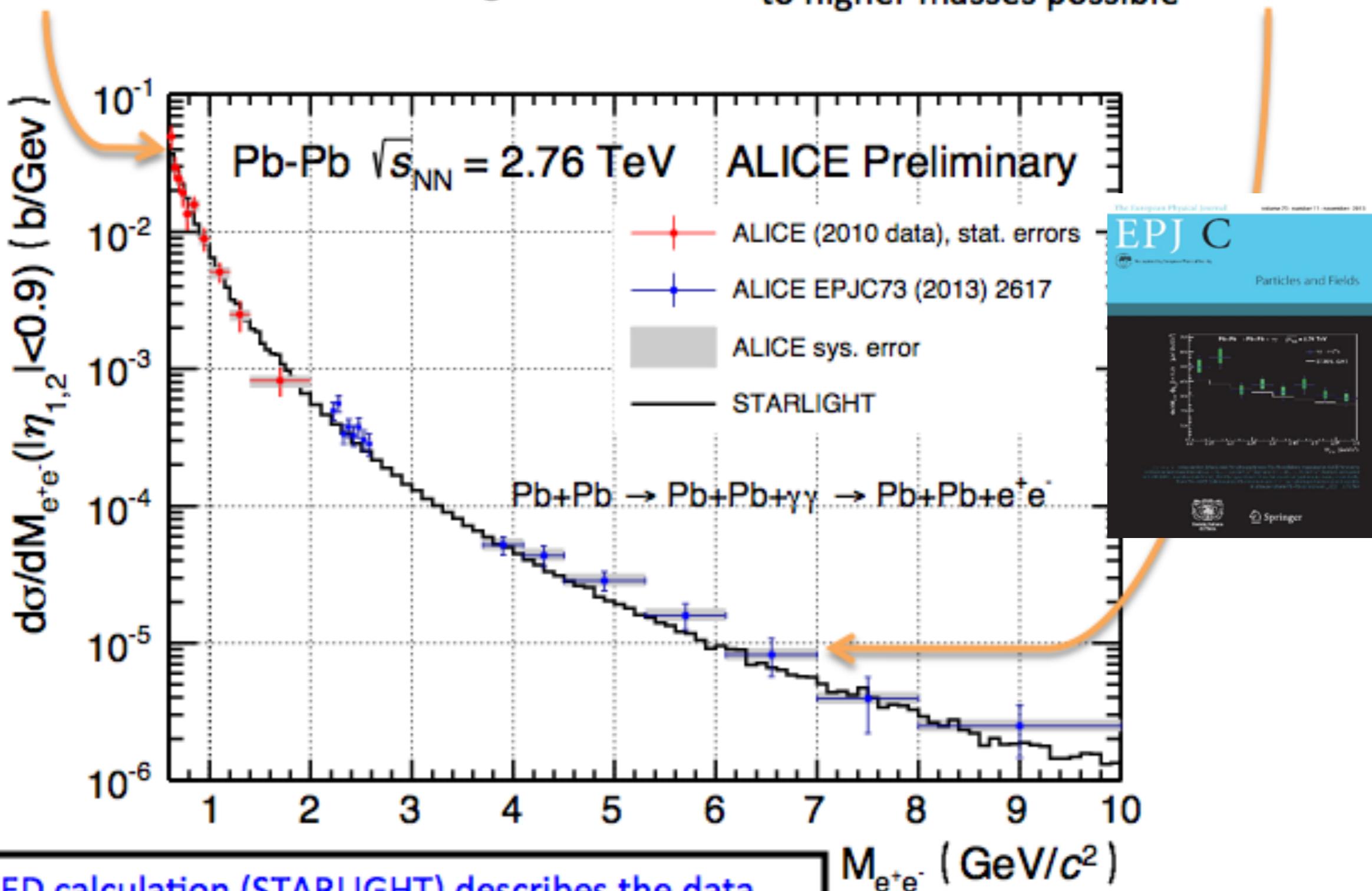
ALICE data sets stringent limits on the contribution from high order terms

Eur. J. Phys. C73, 2617 (2013)



2010: Low luminosity, but trigger allows to cover the low mass region

2011: High luminosity, measurement to higher masses possible



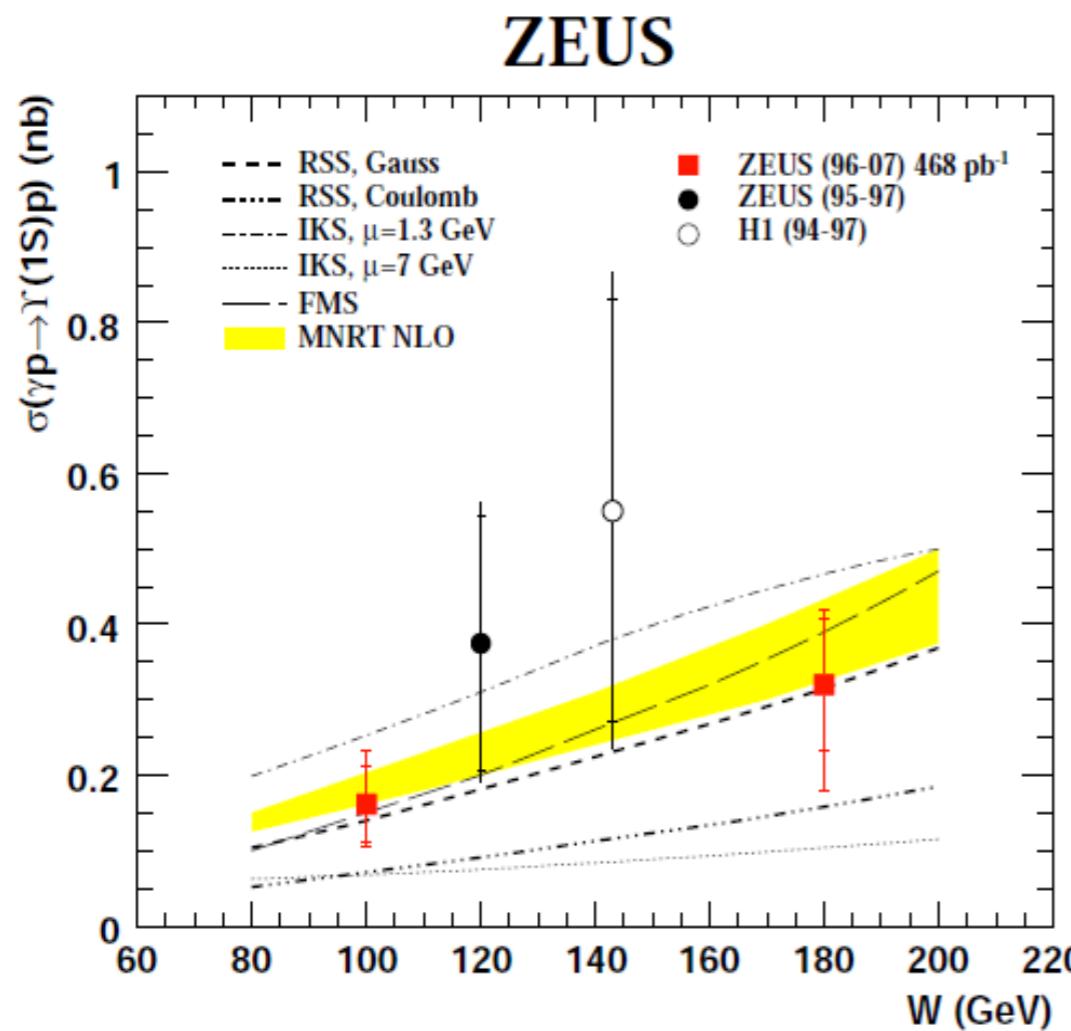
A LO QED calculation (STARLIGHT) describes the data
Strong constraint on NLO contributions

Upsilon photoproduction

$\gamma + p \rightarrow Y + p$: possible thanks to strong photon flux of the proton hitting the Pb nuclues

Very limited statistics from HERA (H1 and ZEUS) ~ 100 candidates

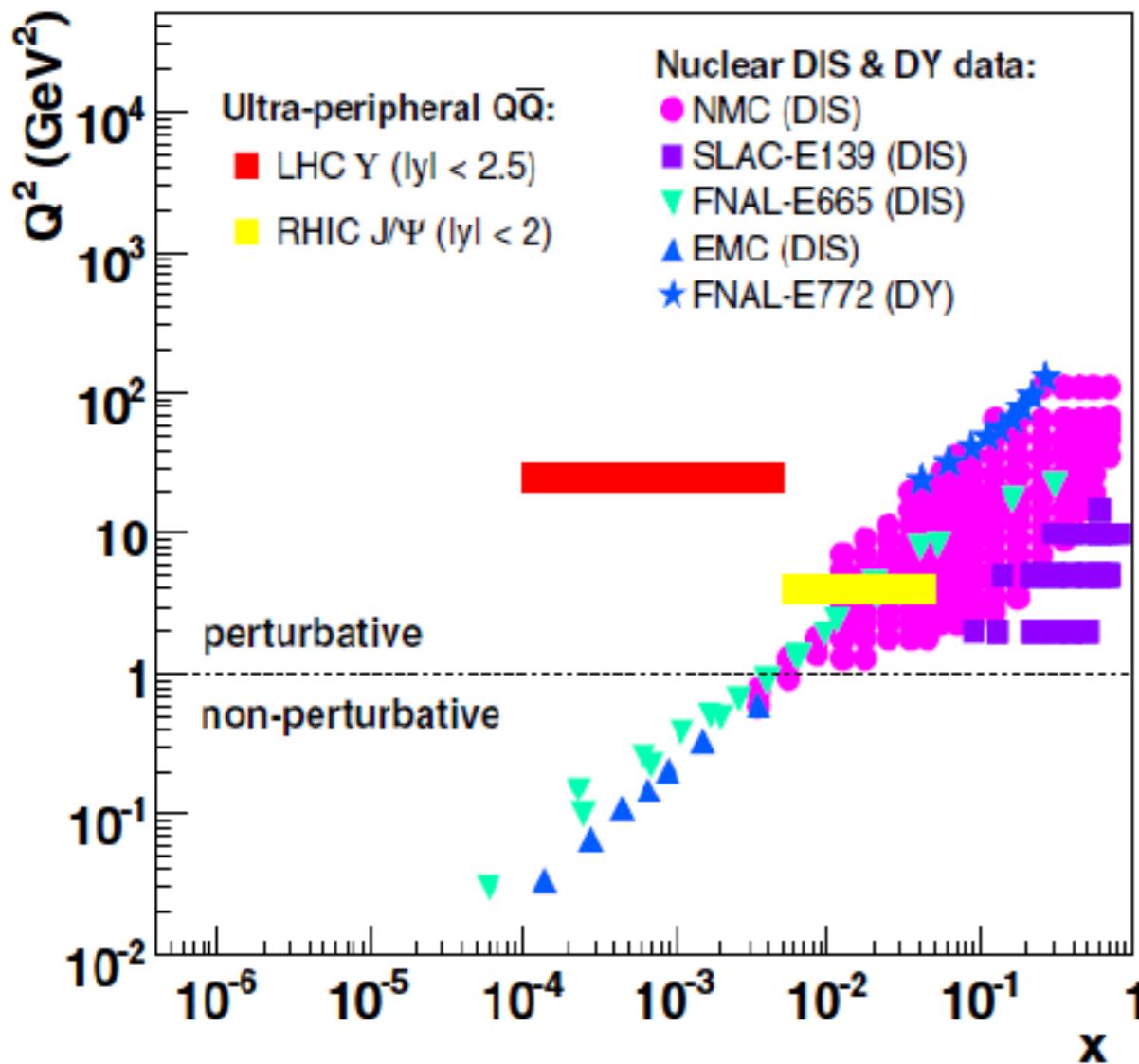
Uncertainty in measured cross section larger than a factor 3



Ideal way to measure this process at LHC

Needed to have a baseline for
 $\gamma + \text{Pb} \rightarrow Y + \text{Pb}$

Here CMS is very competitive as Upsilon acceptance is down to zero transverse momenta

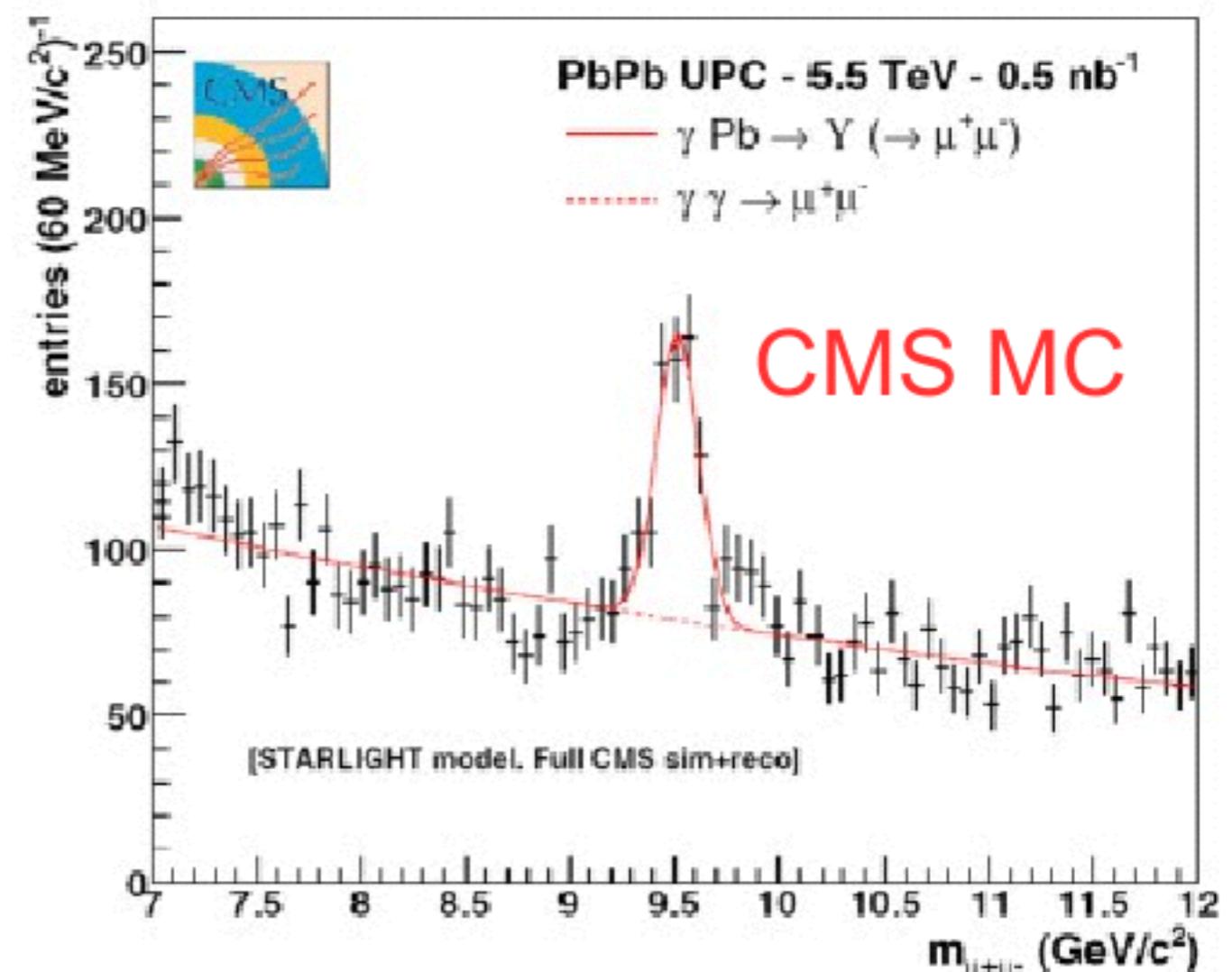


2011: 150 mb^{-1}

at $\sqrt{s}_{NN} = 2.76 \text{ TeV}$

2015: $0.5\text{-}1.5 \text{ nb}^{-1}$

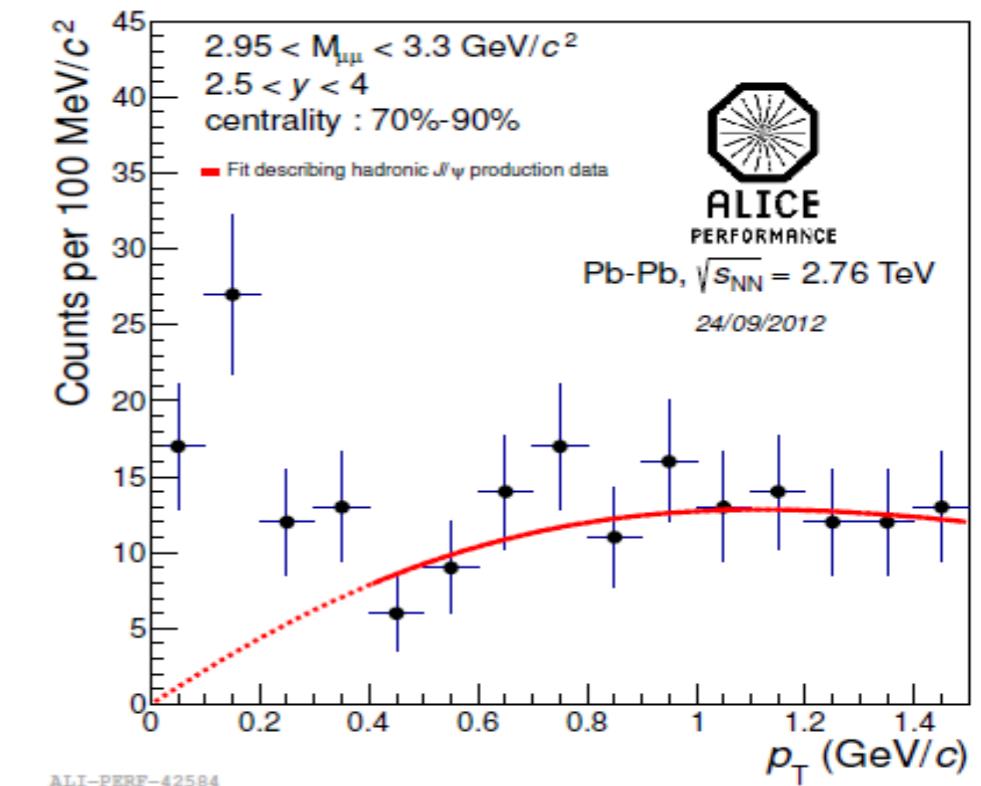
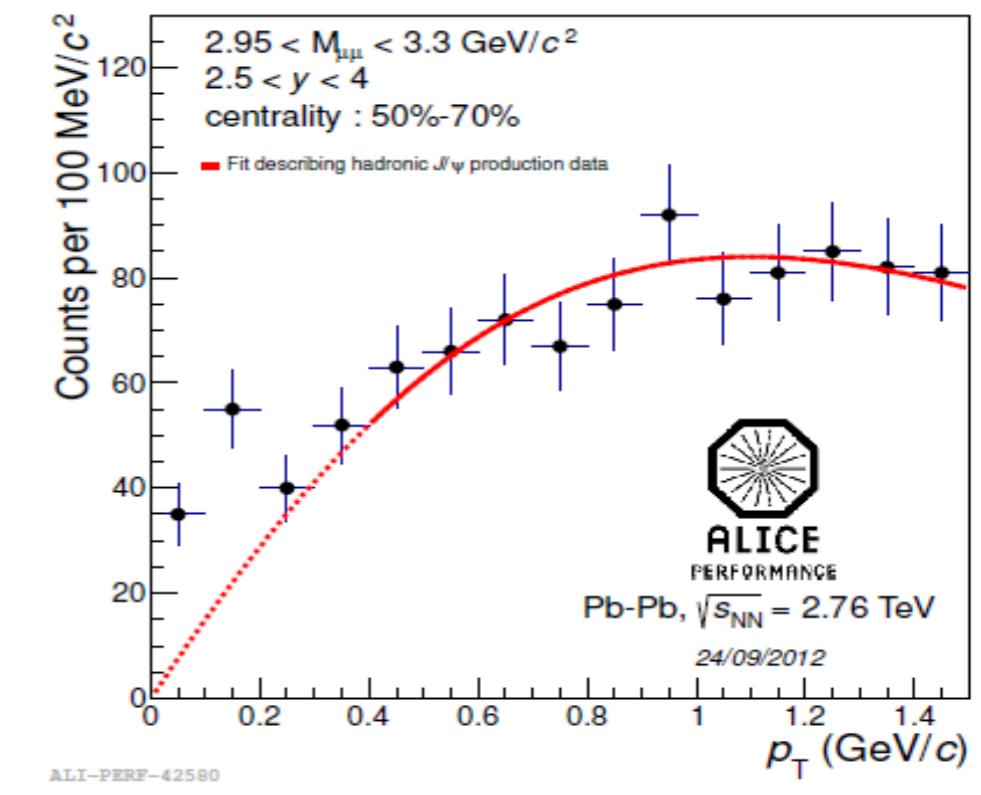
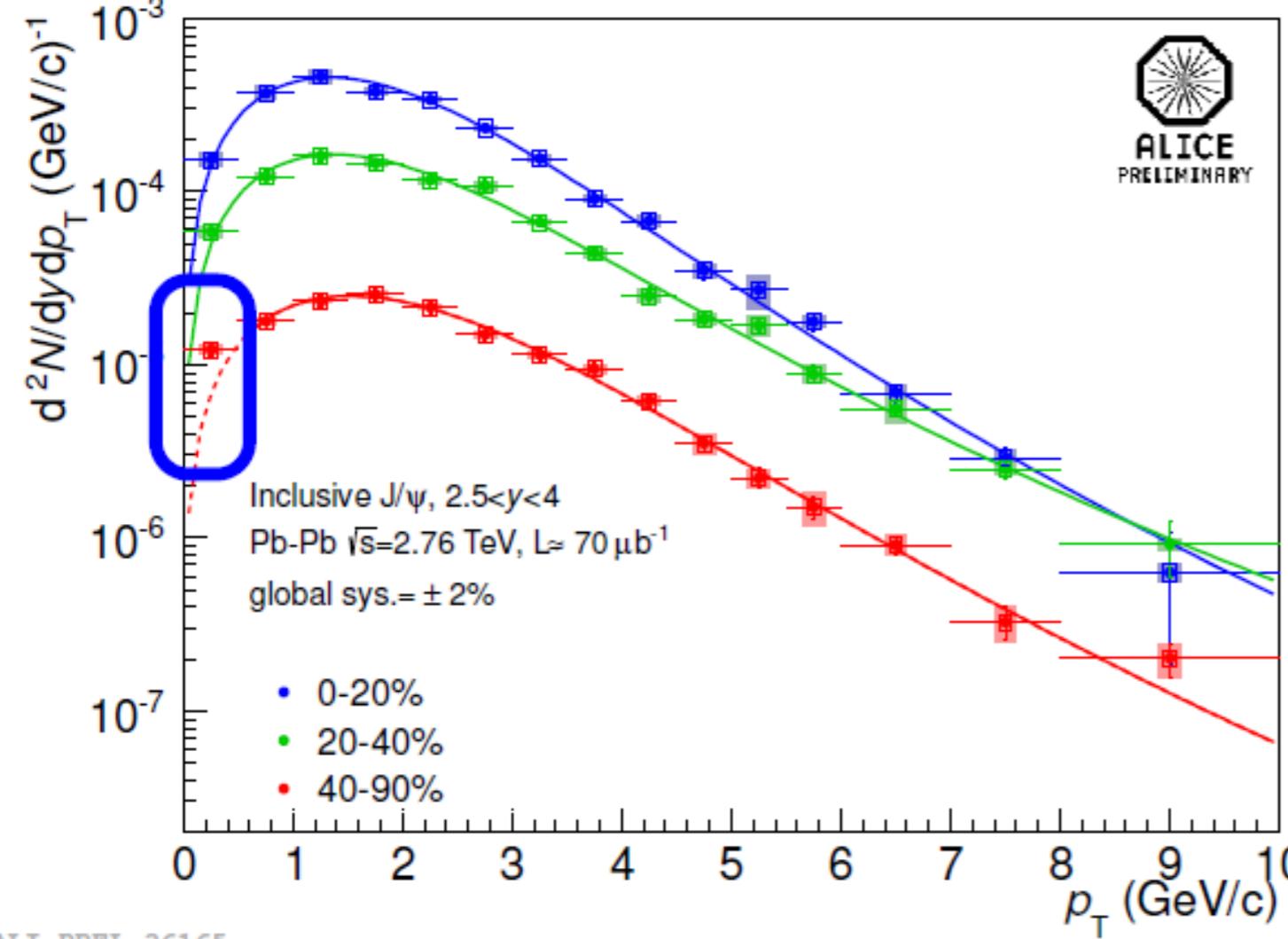
at $\sqrt{s}_{NN} = 5.1 \text{ TeV}$



2015 Pb-Pb run with CMS

- Between 200-1000 Upsilon candidates expected

UPC in inclusive peripheral Pb-Pb at forward rapidity?!

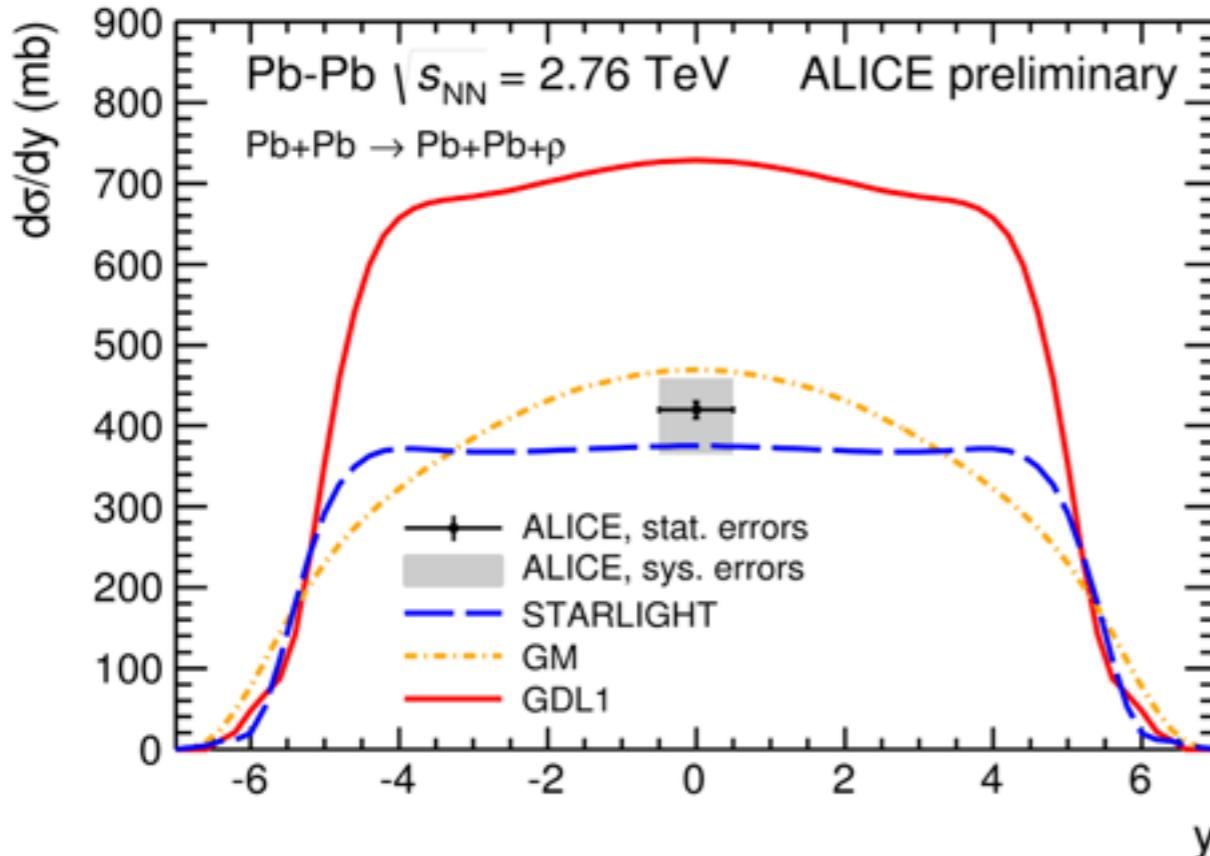




Coherent rho⁰ photoproduction in Pb-

D1

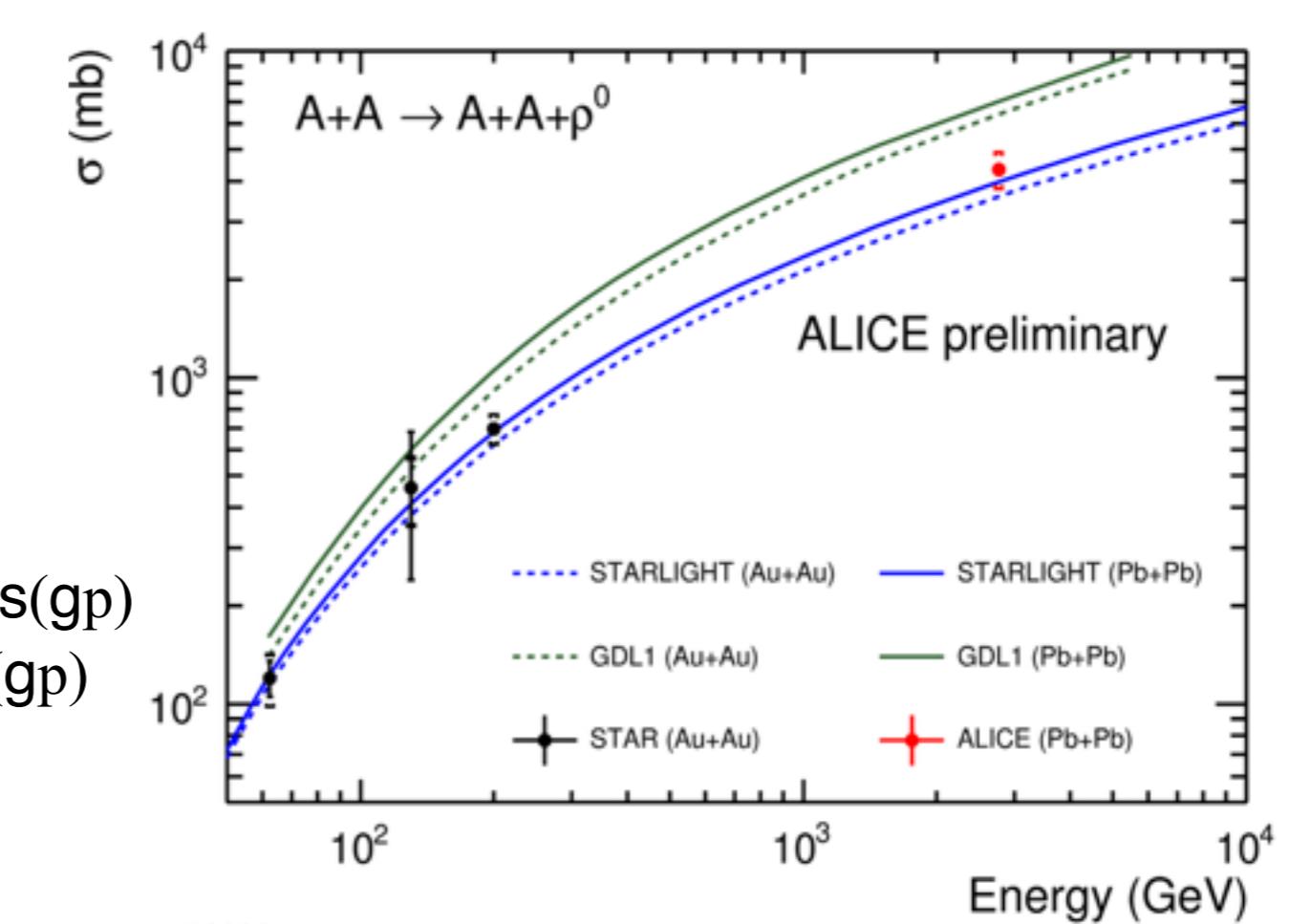
$$\left. \frac{d\sigma(\rho^0)^{\text{coh.}}}{dy} \right|_{y=0} = (420 \pm 10(\text{stat.})^{+39}_{-55}(\text{sys.})) \text{ mb}$$



ALI-PREL-73819

GDL: proper QM Glauber calculation for scaling $s(gp) \rightarrow s(gA)$, uses Donnachie-Landshoff model for $s(gp)$

GM: Based on color dipole model with saturation implementation by the CGC formalism



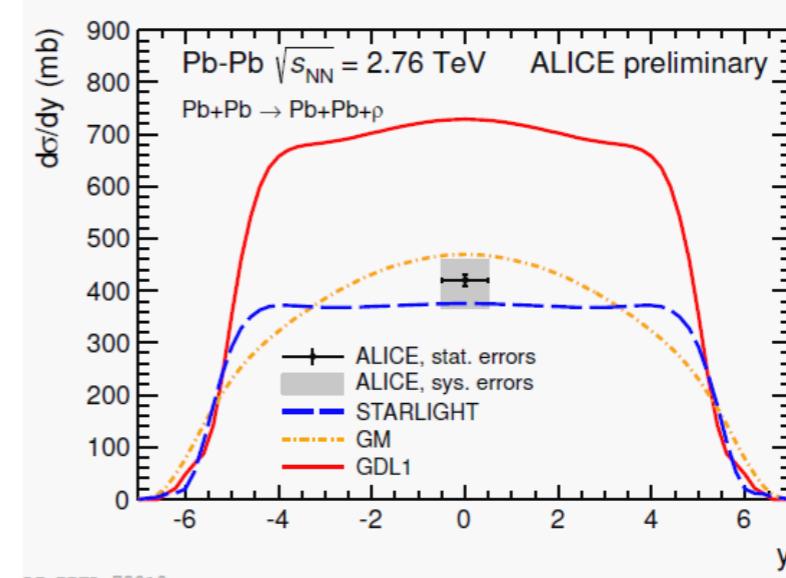
ALI-PREL-73823

STARLIGHT: scales the experimentally measured gp cross section using a Glauber model, neglecting the elastic nuclear cross section

$$\sigma(\rho^0)^{\text{coh.}} = (4.3 \pm 0.1(\text{stat.})^{+0.6}_{-0.5}(\text{sys.})) \text{ b}$$

Cross-sections below Quantum Glauber?

- $\sigma(\gamma A \rightarrow pA)$ should be calculable by a quantum Glauber calculation, with input from $\sigma(\gamma A \rightarrow pA)$
 - ◆ $\sigma(\gamma A \rightarrow pA)$ fixed (or checked) by HERA
- Both ALICE & 62 & 200 GeV STAR measurements find σ 's $\sim 40\%$ lower
- Quantum Glauber approach should be straightforward
 - ◆ Works OK at lower (fixed target, $k \sim 10$ GeV) energies
- Evidence that nuclei do not behave like individual nucleons?
 - ◆ Is pQCD applicable at low Q^2
 - ☞ Shadowing??



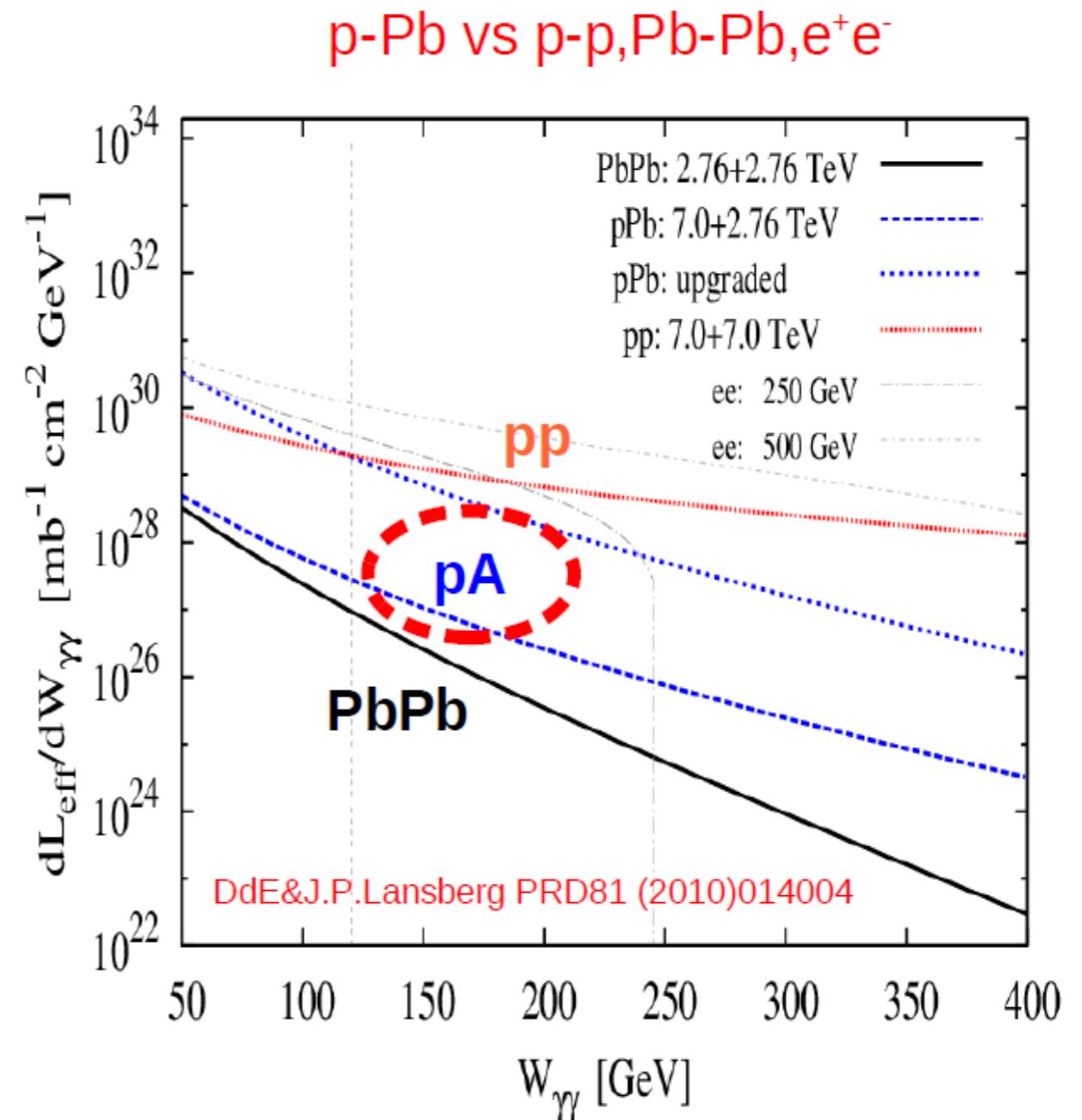
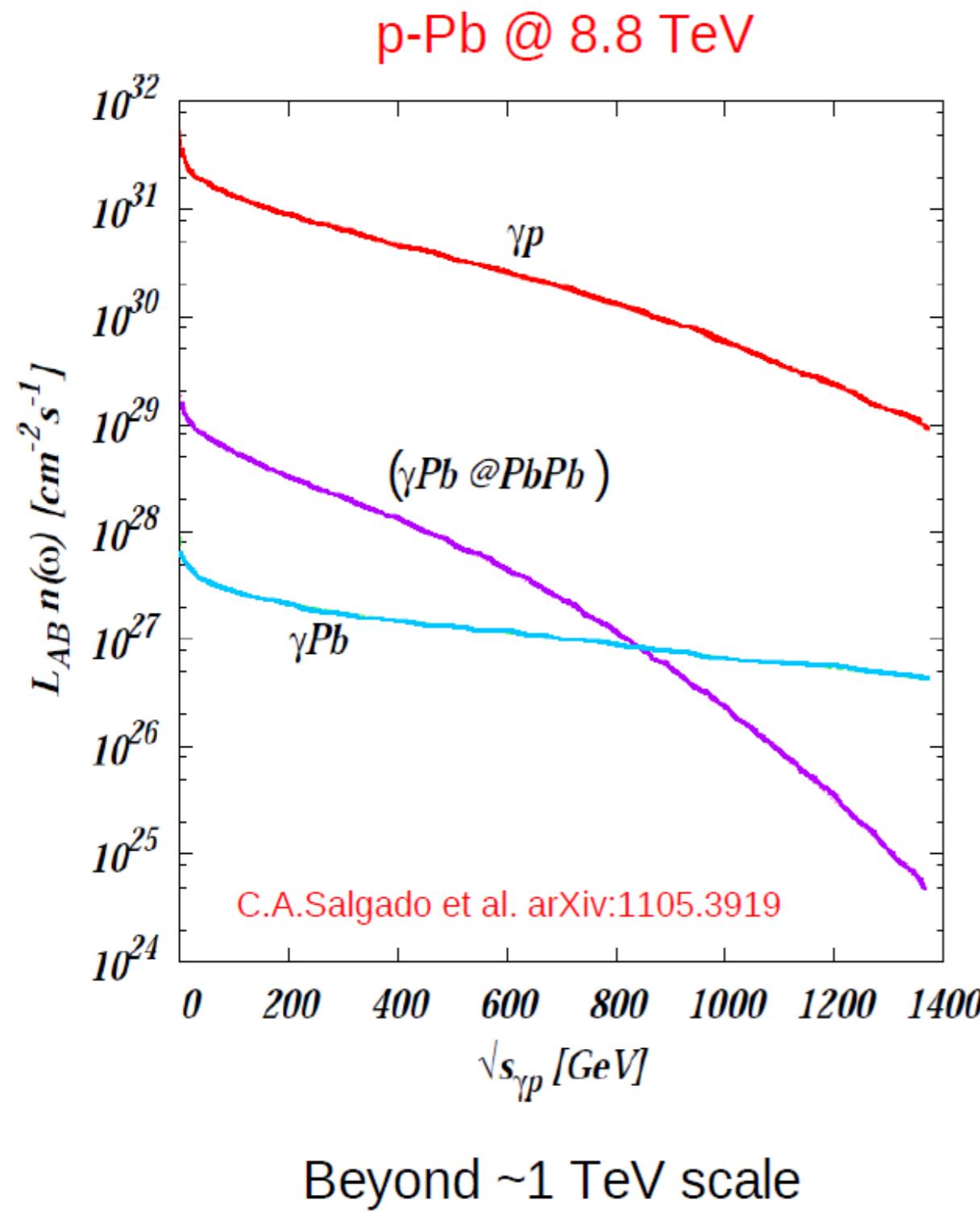
ALICE data presented by C. Mayer
at the Wkshp. on Photon induced
collisions at the LHC (2014).

S. Klein
IS conference. Dec 2014

p

UPCs in Pb-Pb

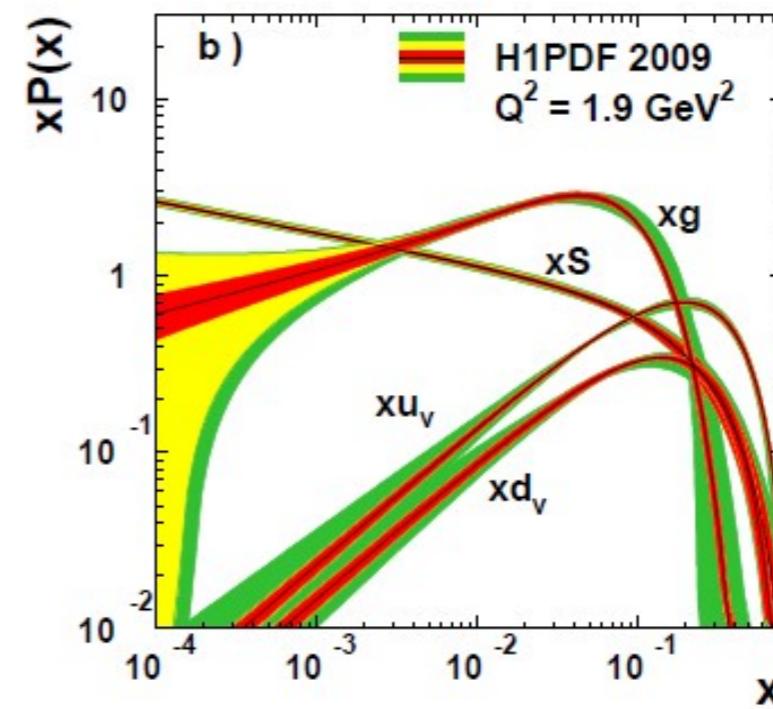
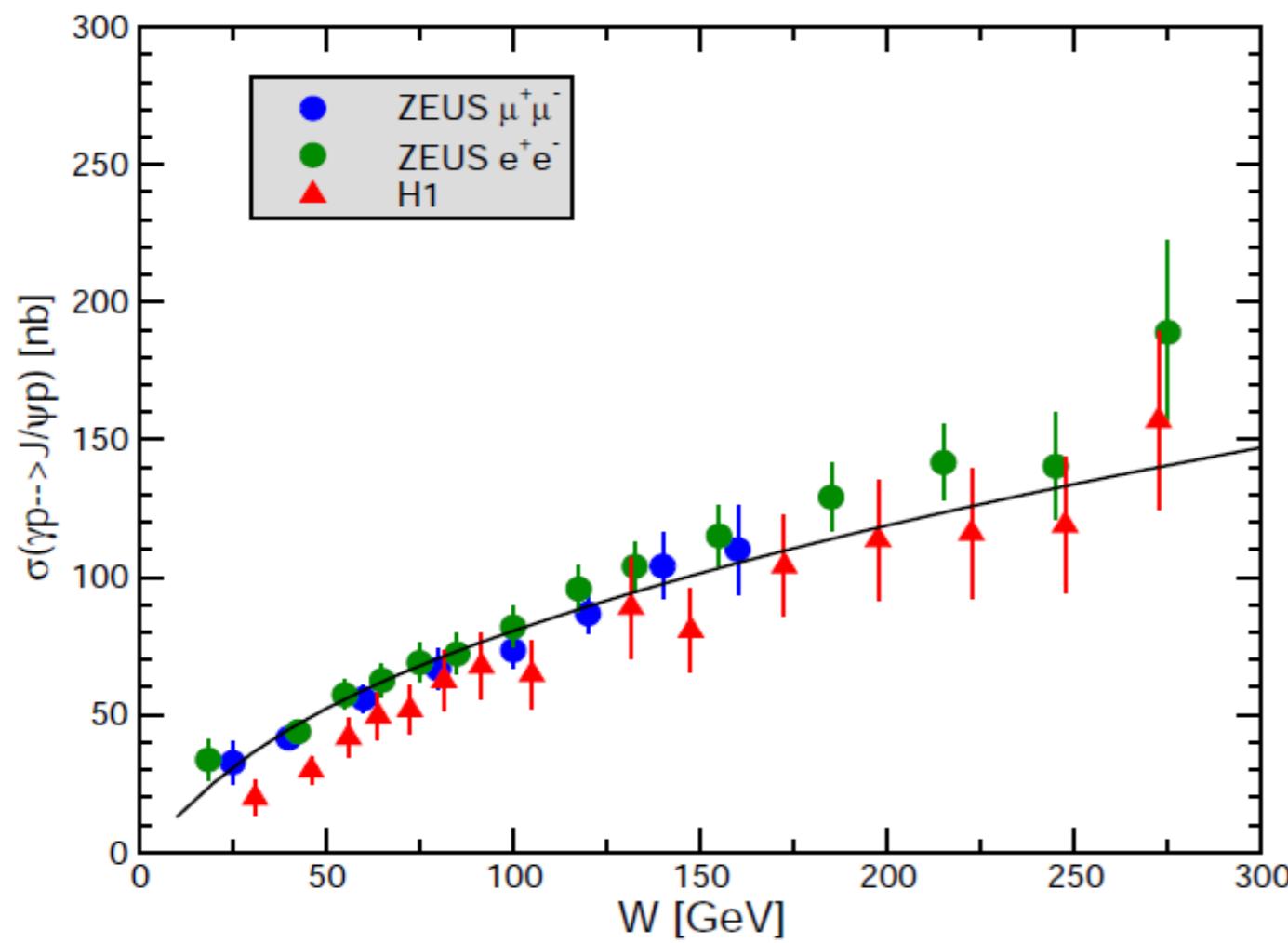
Effective luminosities in UPC



$$\mathcal{L}_{\gamma\gamma}^{\text{eff}} \equiv \mathcal{L}_{AB} d\mathcal{L}_{\gamma\gamma}/dW_{\gamma\gamma}$$

J/ ψ photoproduction in γp

$$\frac{d\sigma_{\gamma p \rightarrow p J/\psi}}{dt} = \frac{\Gamma_{ee} M_{J/\psi}^3 \pi^3}{48 \alpha_{em}} \cdot \frac{\alpha_S^2(\bar{Q}^2)}{\bar{Q}^8} \left[x g_N(x, \bar{Q}^2) \right]^2 \exp[B_{J/\psi}(s)t]$$

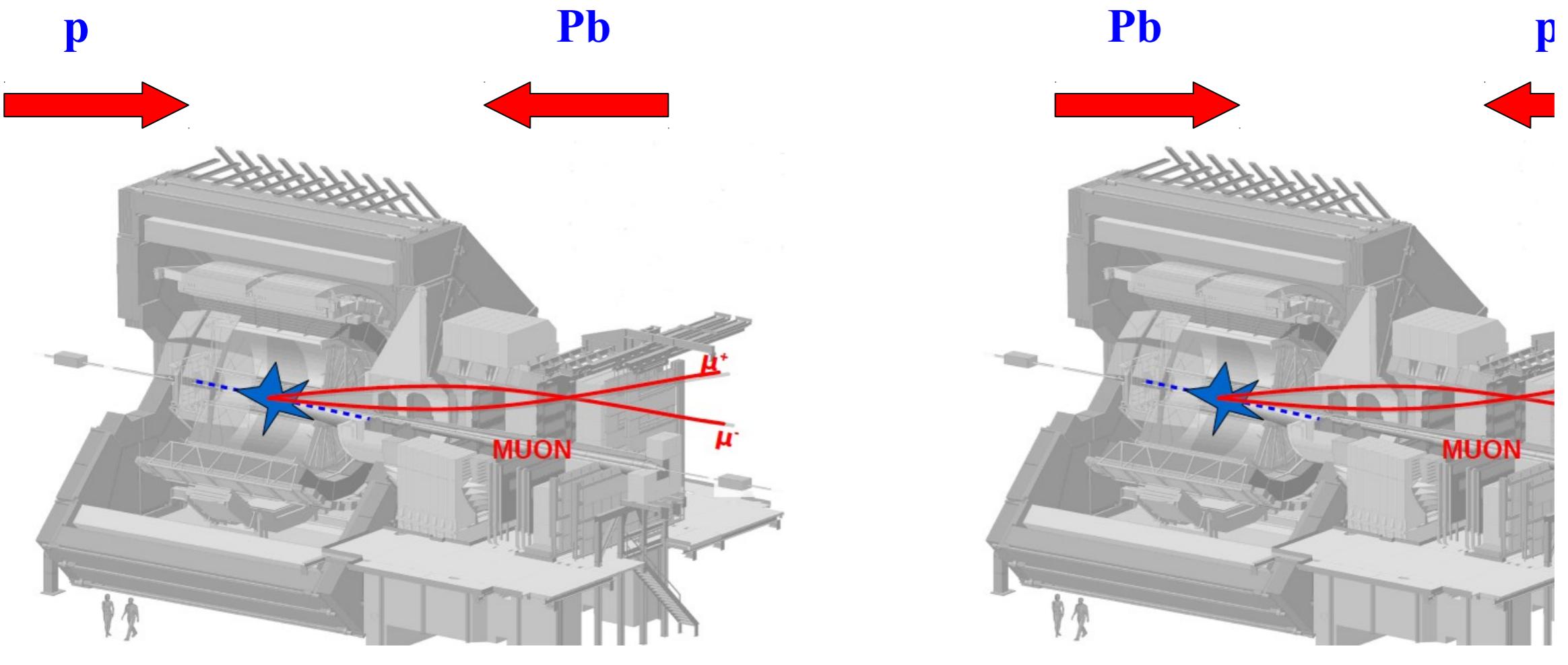


The Pb nucleus acts as photon emitter (enhanced flux by factor $Z^2 \approx 7000$ compared to the photon flux from the proton

At LHC Bjorken-x down to 10^{-5}

γp centre-of-mass energies at the 1 TeV energy scale

H1: A. Aktas *et al.* Eur.Phys. J.C46:585-603,2006
 ZEUS:S. Chekanov *et al.*, Nucl. Phys. B695 (2004) 3.
 A. Martin *et al.* Phys.Lett. B 662:252-258, 2008



$21 < W_{\gamma p} < 45 \text{ GeV}$

$549 < W_{\gamma p} < 1163 \text{ GeV}$

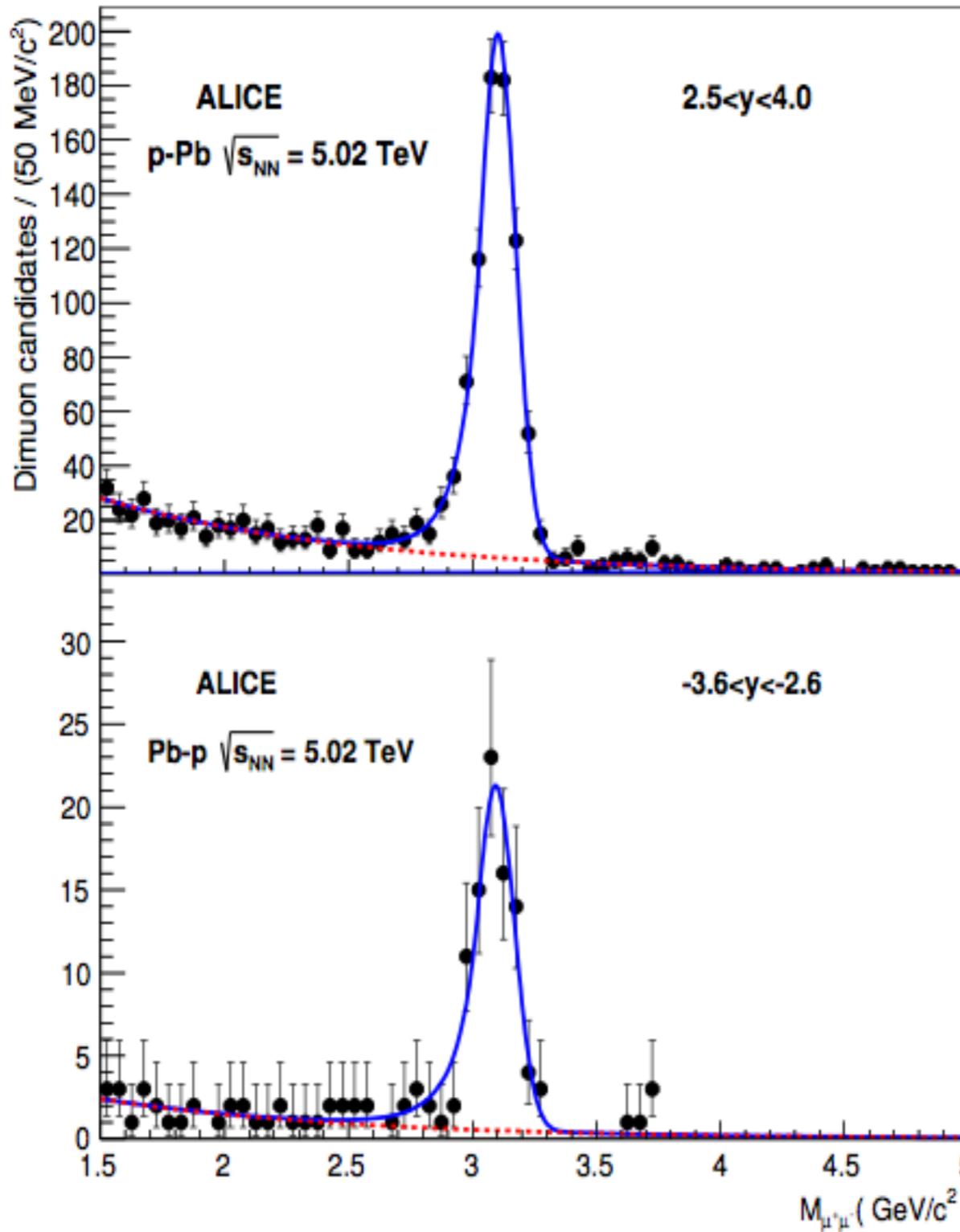
- The fact that the Pb nucleus is the dominant photon emitter allows us to separate the two $W_{\gamma p}$ regimes unambiguously.
 - “p-Pb” (*) corresponds to the **lower** energy range
 - “Pb-p” corresponds to the **higher** energy range.

$$x = \left(M_{J/\psi} / \sqrt{s_{NN}} \right) \exp(\pm y)$$



Exclusive J/psi in p-Pb

ALICE



$\langle W_{gp} \rangle \sim 30 \text{ GeV}$

$\langle W_{gp} \rangle \sim 700 \text{ GeV}$

[Phys.Rev.Lett. 113 \(2014\) 23, 232504](#)

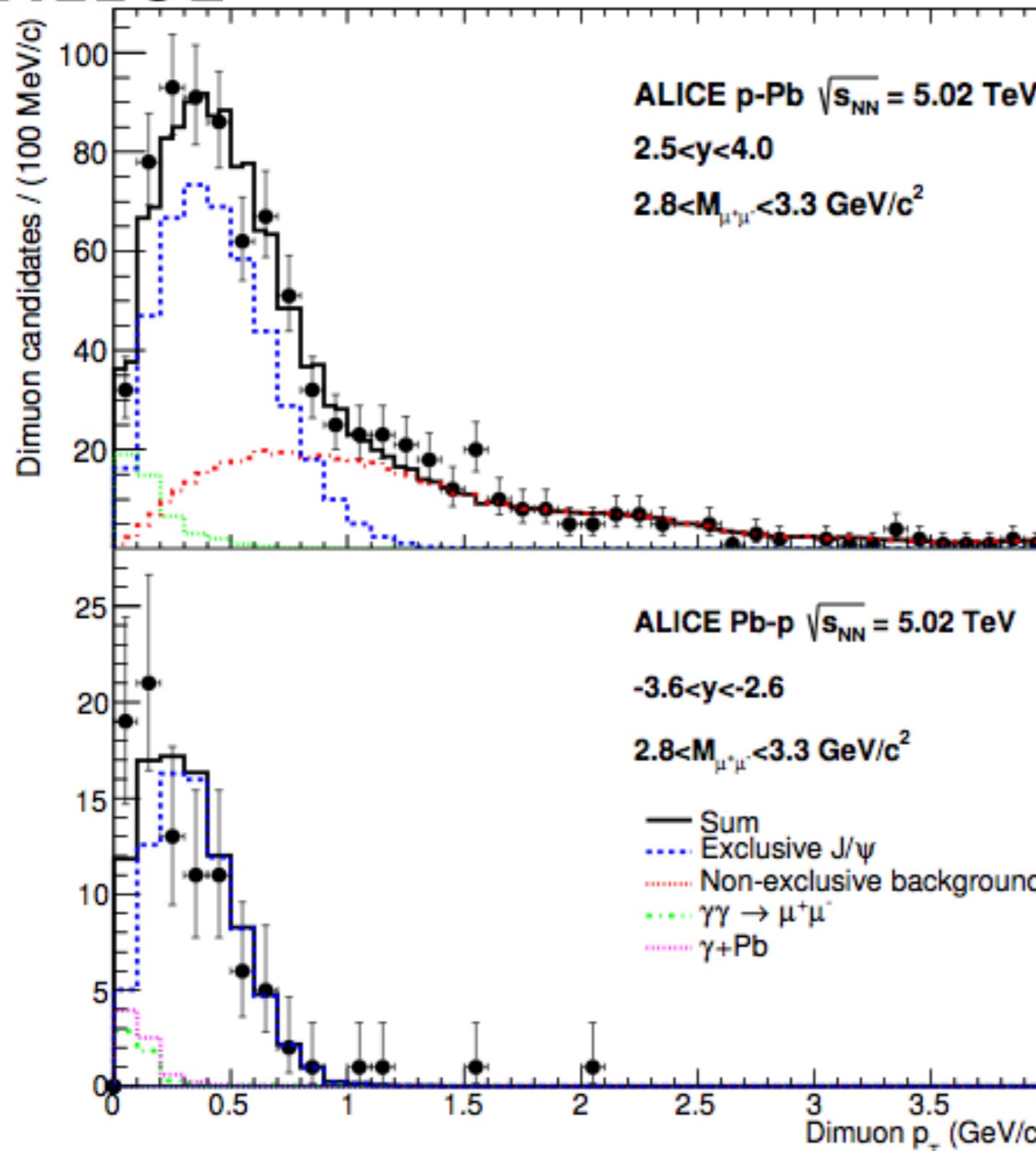
DTT
ICHEP 2014



Exclusive J/psi in p-Pb

Phys.Rev.Lett. 113 (2014) 23, 232504

ALICE



Data well described by
templates

**Energy dependence
is clearly visible**

Low W_{gp} energy point
 $\langle W_{gp} \rangle \sim 30$ GeV

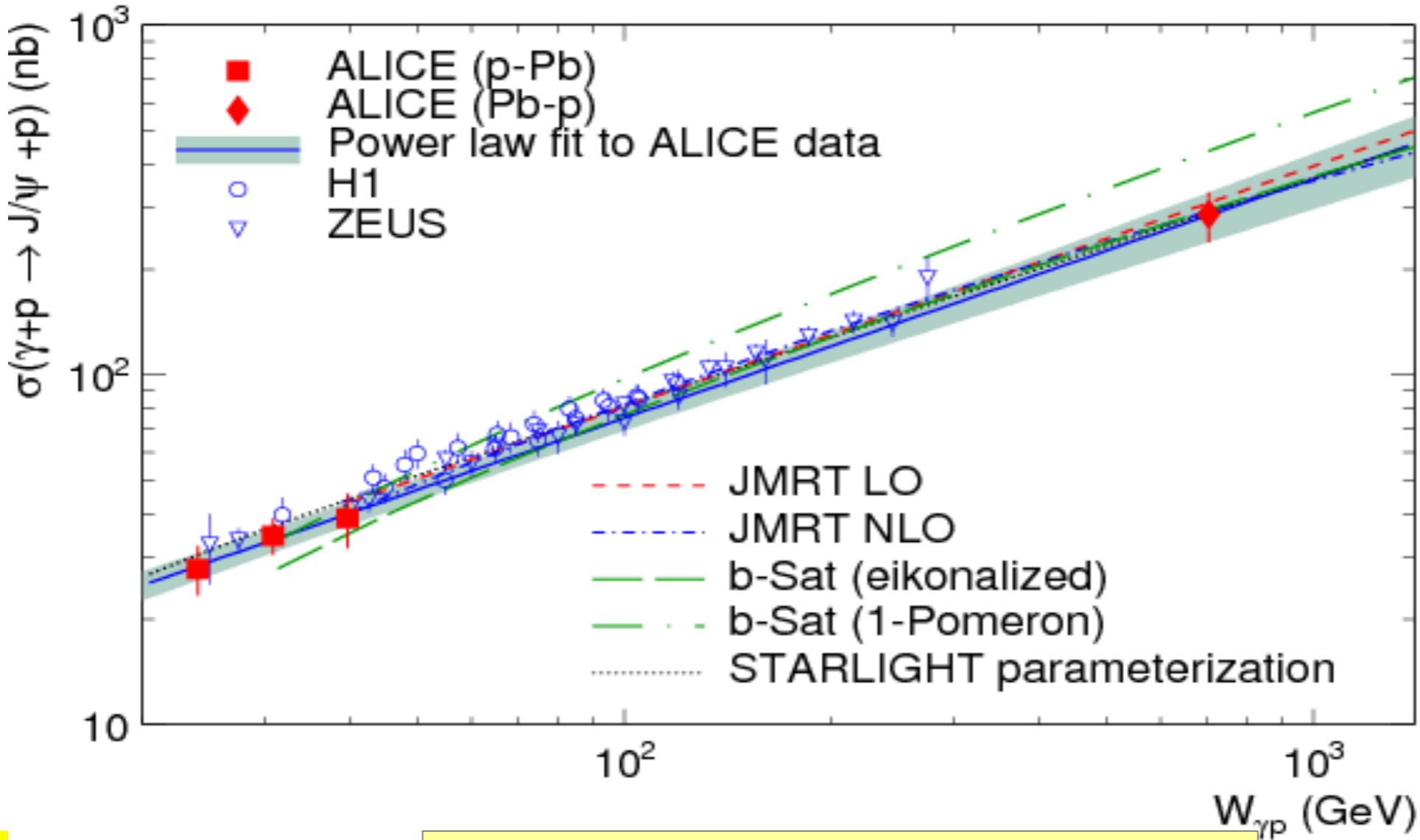
High W_{gp} energy point
 $\langle W_{gp} \rangle \sim 700$ GeV



Exclusive J/psi in p-Pb

ALICE

Phys.Rev.Lett. 113 (2014) 23, 232504

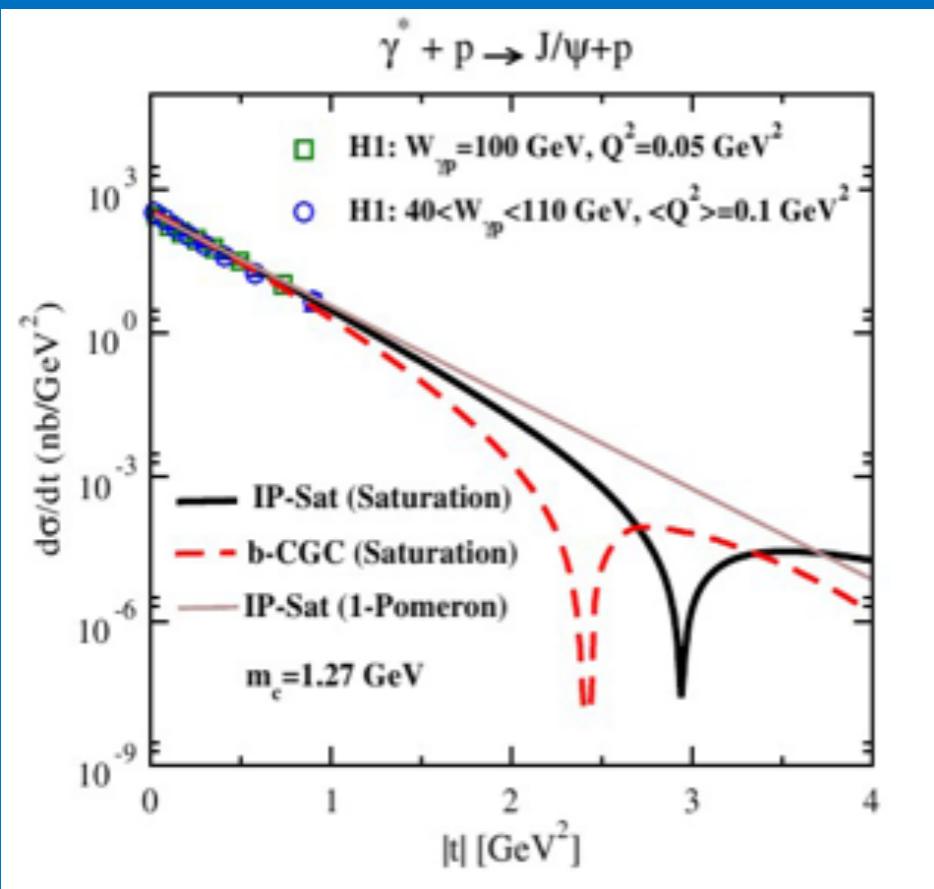


HERA
Measurements
H1
 $\delta=0.67\pm0.03$
ZEUS
 $\delta=0.69\pm0.02$

No change on the proton gluon density between HERA and LHC energies

DTT
ICHEP 2014

- Our knowledge of the photon emitter allows us to solve for $\sigma(W_{\gamma p})$ using the measured $d\sigma/dy$
- A power law fit ($\sigma(W) \sim W^\delta$) to ALICE data points gives $\delta = 0.68 \pm 0.06$.



J/Ψ with high statistics

N. Armesto, A.H. Rezaeian, Phys. Rev. D 90 (2014) 054003:

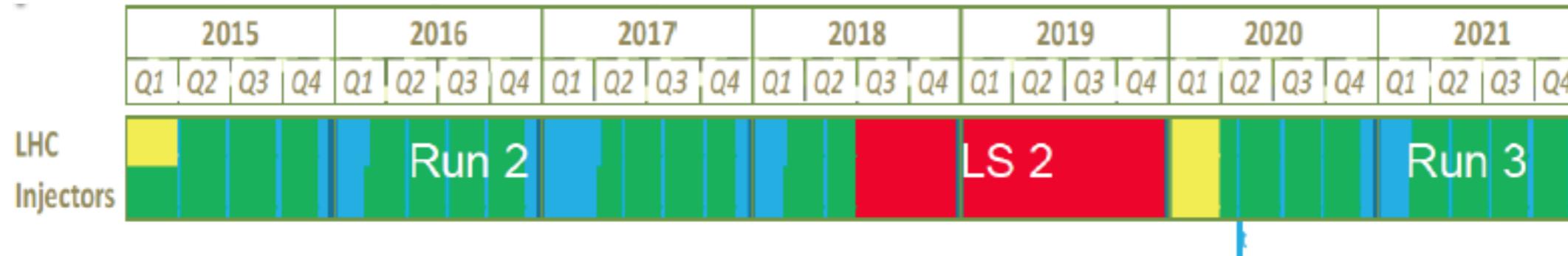
t ($=p_{T2}$) distribution of differential cross sections of photo-production of vector mesons may discriminate among saturation and non-saturation models.
 \rightarrow dip (or multiple dips) in the t distribution of diffractive photoproduction of vector mesons”

J/Ψ rapidity tag

- At forward rapidity we have contributions both from $x \sim 10^{-2}$ (95%) and 10^{-4} (5%), depending on the Pb nucleus emitting the photon.
 \rightarrow tagging by using ZDC activity, see [arXiv:1109.0737](https://arxiv.org/abs/1109.0737) \rightarrow Gluon shadowing at 10^{-4} feasible !
- \rightarrow Run2: 2,500 J/Ψ^* $5\% * 30\% \sim 40$ tagged J/Ψ at $x \sim 10^{-4}$ (1 nb^{-1})

E. Scapparone
IS conference. Dec 2014

Run 2 and beyond



5.1 TeV for the 2015 Pb-Pb run

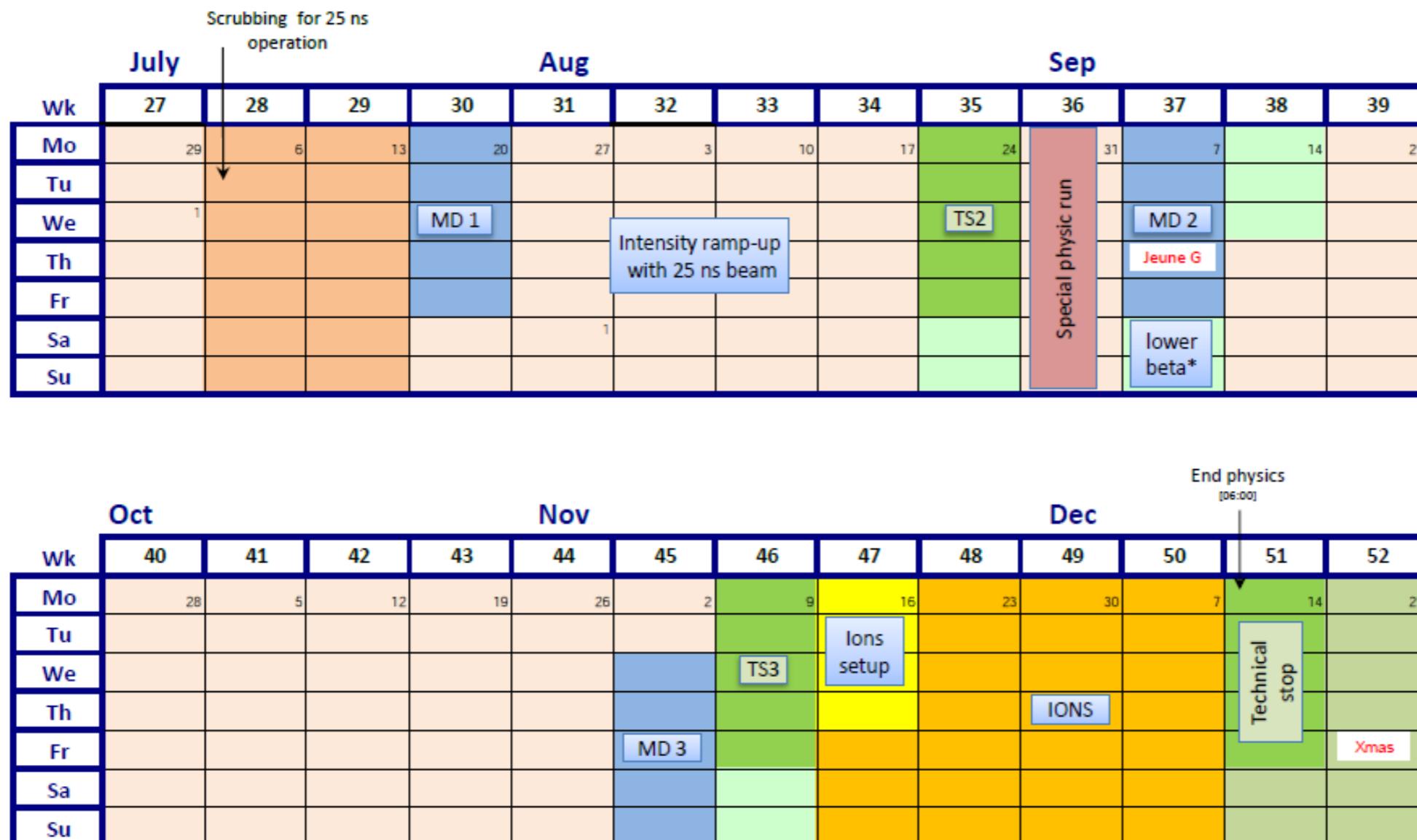
1 nb-1 for CMS/ATLAS

Next heavy-ion run in 2016 and 2017

Heavy Ion Preparations for Run 2

- The 2015 Ion Period will be the first high luminosity Pb run in the LHC Ion program
 - Peak Lumi: $3.7 \times 10^{27} \text{ cm}^{-2}\text{s}^{-1}$, interaction rate $\sim 30\text{kHz}$
 - 8 times higher than the 2011 PbPb interaction rate, and 4 times higher than the LHC design value!

HI and special pp runs in 2015



- 4 day floating MD removed
- Otherwise as was



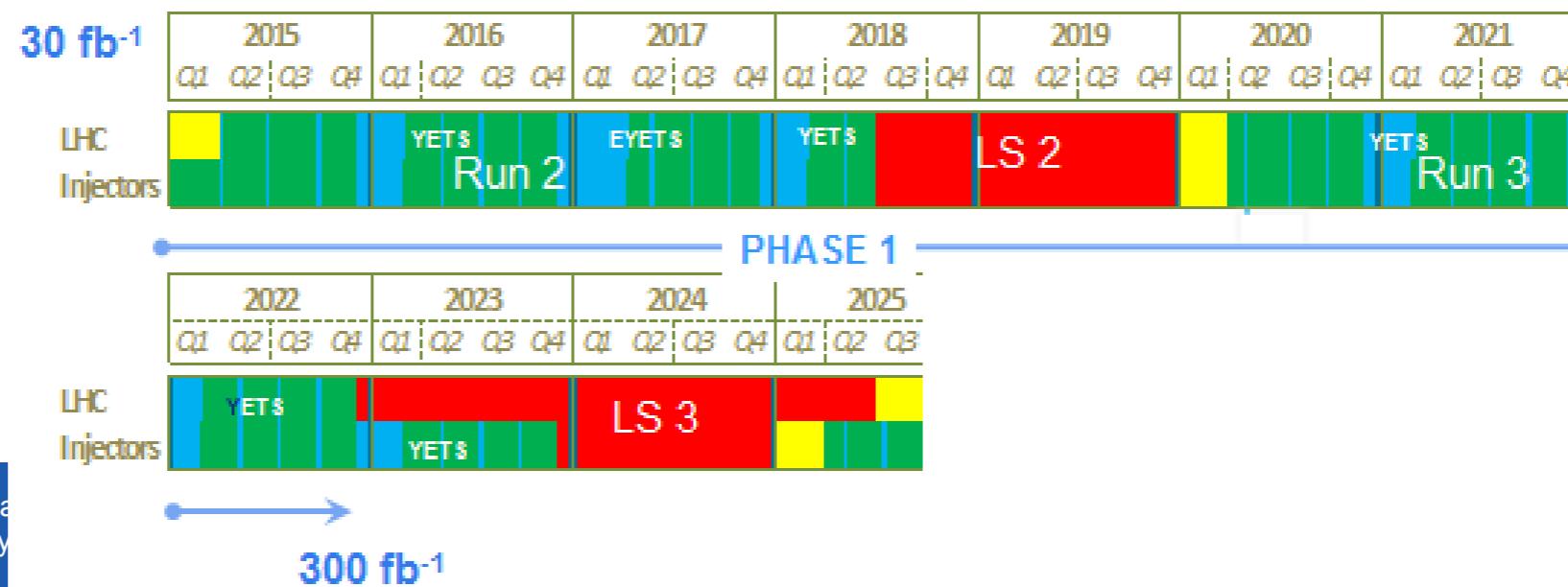
LHC schedule

LHC goal for 2015 and for Run 2 and 3

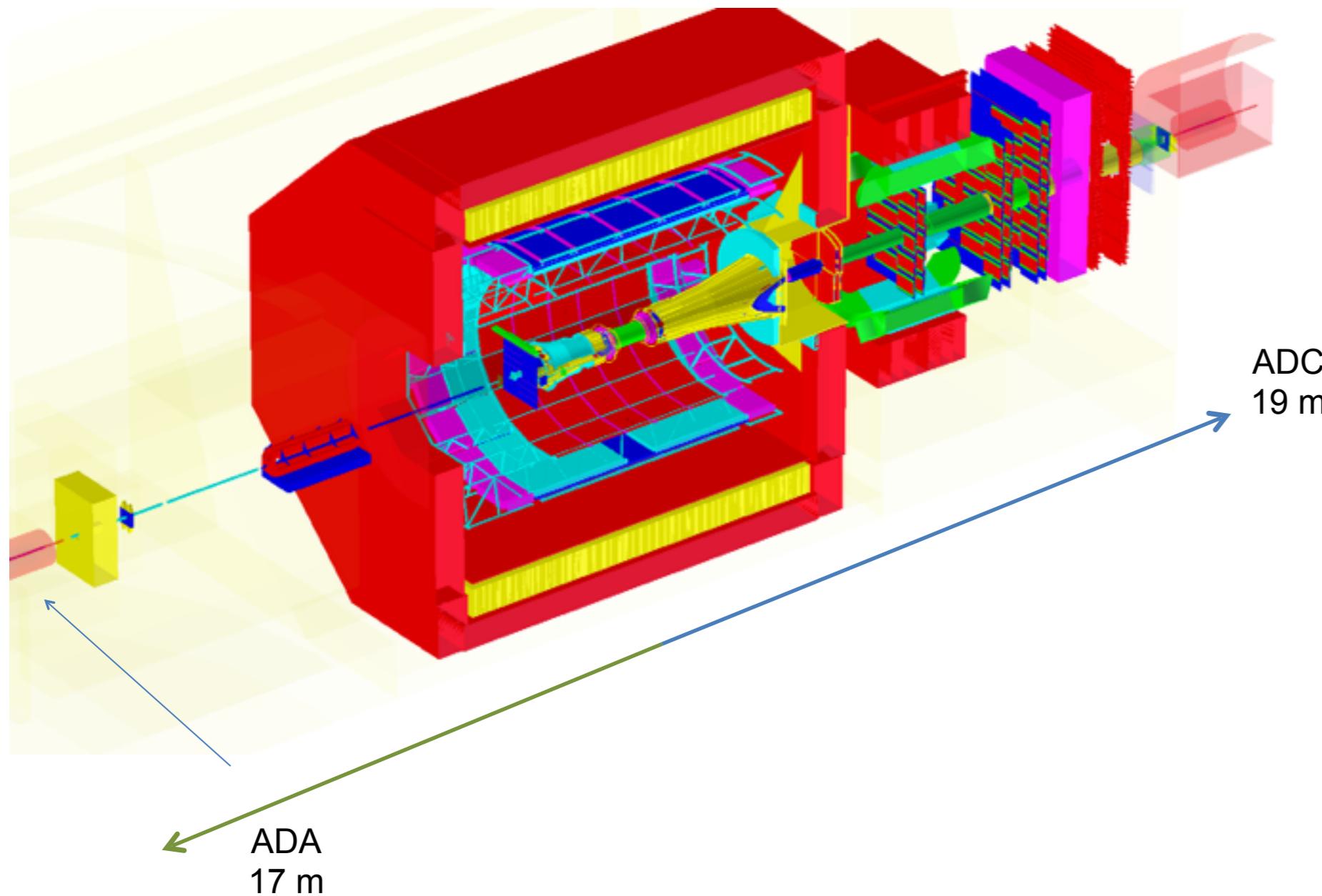
Integrated luminosity goal:
2015 : 10 fb^{-1}

Run2: $\sim 100\text{-}120 \text{ fb}^{-1}$
(better estimation by end of 2015)

300 fb^{-1} before LS3

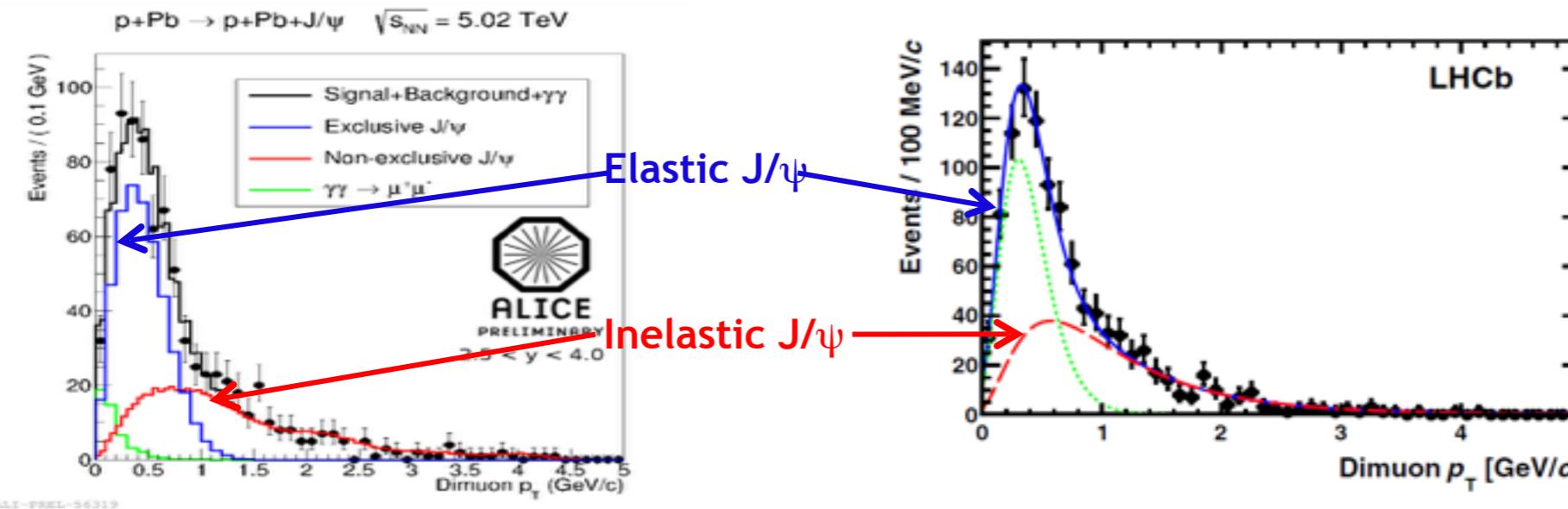


Ongoing projects at ALICE: ADA/ADC



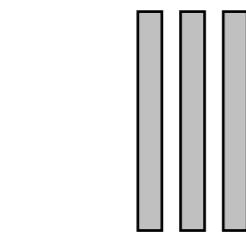
J/ ψ in ultra-peripheral pp and p-Pb collisions

- Exclusive J/ ψ photoproduction measured by ALICE in p-Pb (Pb as a source of photons) and by LHCb in pp. ALICE can also measure J/ ψ photoproduction in pp in run2 ($\sqrt{s} = 13$ TeV)
- Experimentally events with only two muons are selected by applying vetos on central and forward detector activity
- Inelastic J/ ψ are characterized by broader pt distribution, but separation of inelastic and elastic contributions is a delicate task
- **ADA and ADC detectors will help to suppress inelastic contribution**



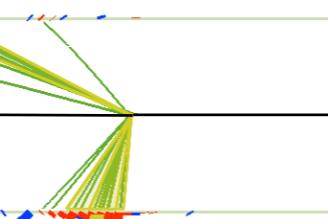
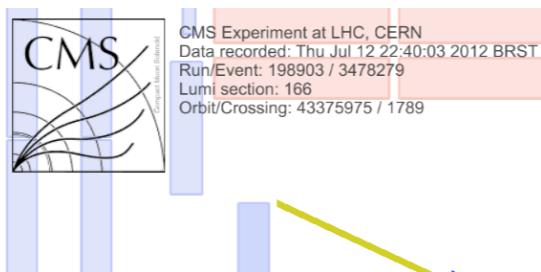
Dijet: CMS+TOTEM

Forward Shower
Counters

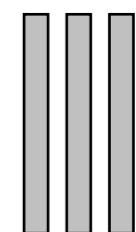


59 - 114 m

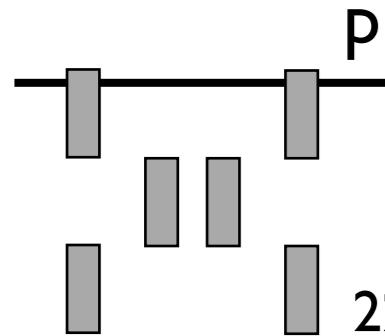
P



Forward Shower
Counters



P



TOTEM
Roman Pots

TOTEM T2



TOTEM T2

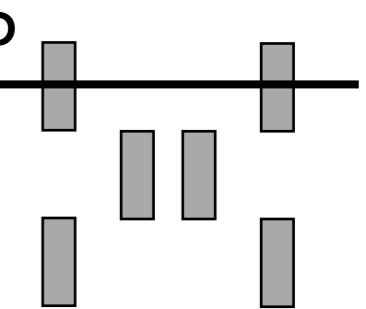
CMS: $|\eta| < 5$

T2: $5.3 < |\eta| < 6.5$

FSC: $6 < |\eta| < 8$

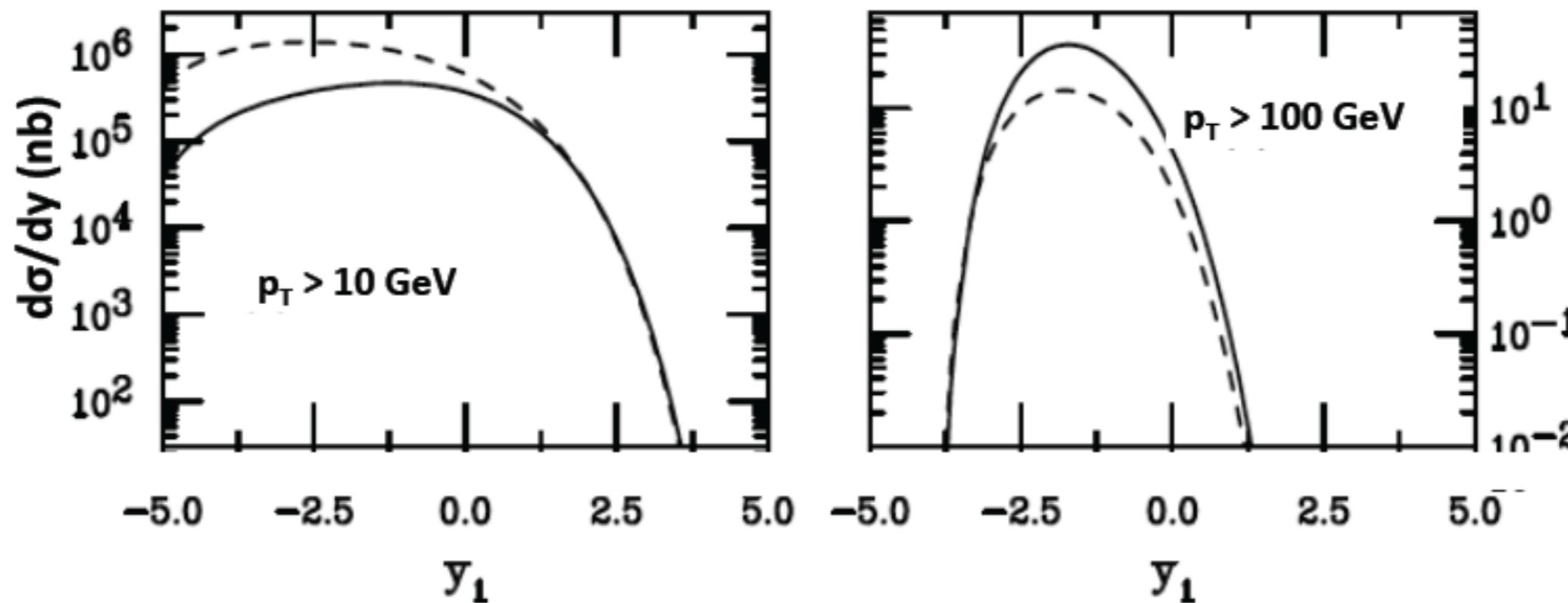
TOTEM RP

very large rapidity coverage !



TOTEM
Roman Pots

Dijet production in UPC



The photon is coming from the left and its direction can be resolved by the correlation with neutrons in the ZDCs.

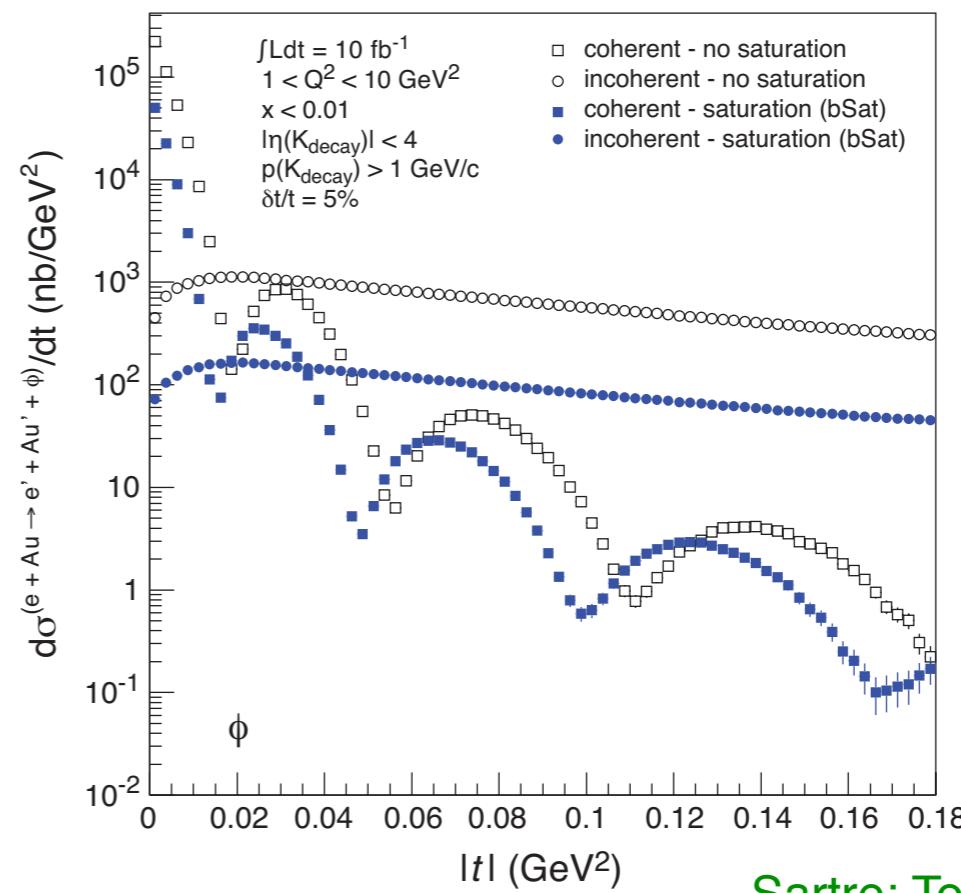
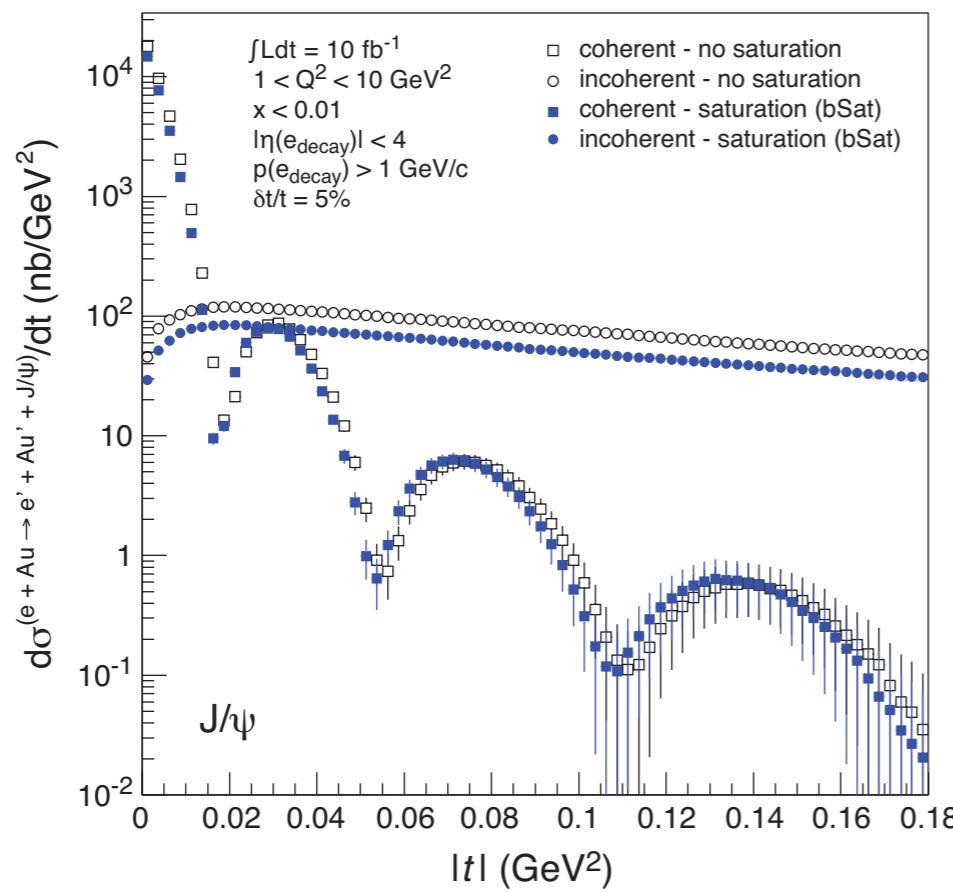
In the direct process (solid), the entire photon energy contributes to the hard process while in the resolved case (dashed) part does.

UPCs and the EIC

- UPCs are a ‘limiting case’ of EIC physics
 - ◆ Almost real photons
 - ◆ No tagging
- UPC data for many channels are available now
 - ◆ No need to wait
 - ◆ Other channels can be studied with improved triggering
- At the very least, this is a good testbed for EIC physics

S. Klein
IS conference. Dec 2014

Exclusive Vector Meson Production in e+A



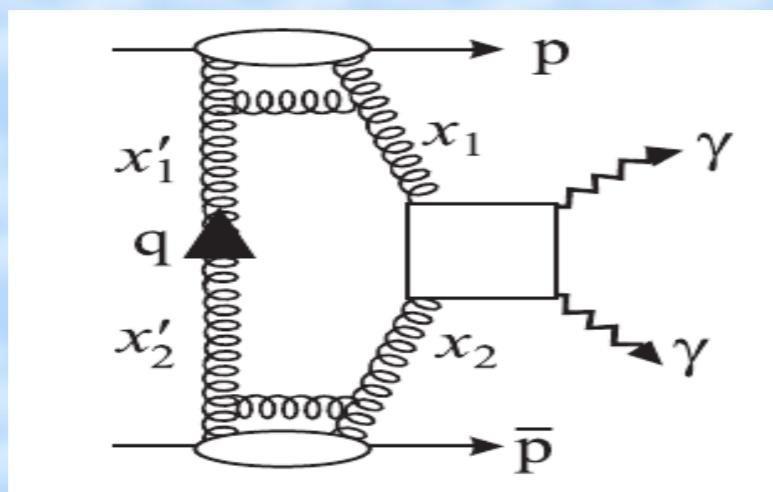
Sartre: Toll, Ullrich,
Phys.Rev. C87, 024913 (2013)

- Low-t: coherent diffraction dominates - gluon density
- High-t: incoherent diffraction dominates - gluon correlations
 - Need good breakup detection efficiency to discriminate between the two scenarios
 - unlike protons, forward spectrometer won't work for heavy ions
 - measure emitted neutrons in a ZDC
 - rapidity gap with absence of break-up fragments sufficient to identify coherent events

Exclusive production in pp vs. AA

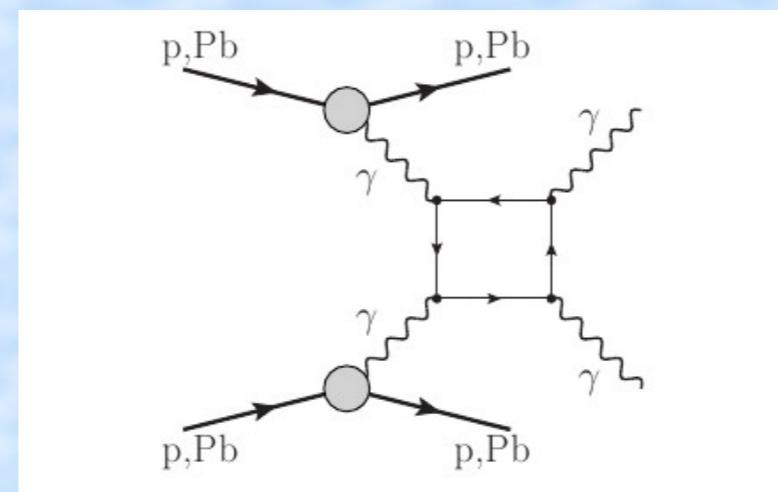
Different production mechanisms may dominate. Consider exclusive $\gamma\gamma$ (or Higgs) production:

p-p



V. A. Khoze, A.D. Martin,
M.G. Ryskin, W.J. Stirling,
Eur. Phys. J C 38 (2005) 475.

Pb-Pb



D. d'Enterria, G.G. Silveira,
PRL 111 (2013) 080405.

In p-p collisions, 3 (or more) gluon exchange dominate, whereas for heavy-ion collisions, $\gamma\gamma \rightarrow \gamma\gamma$ dominate.

Summary

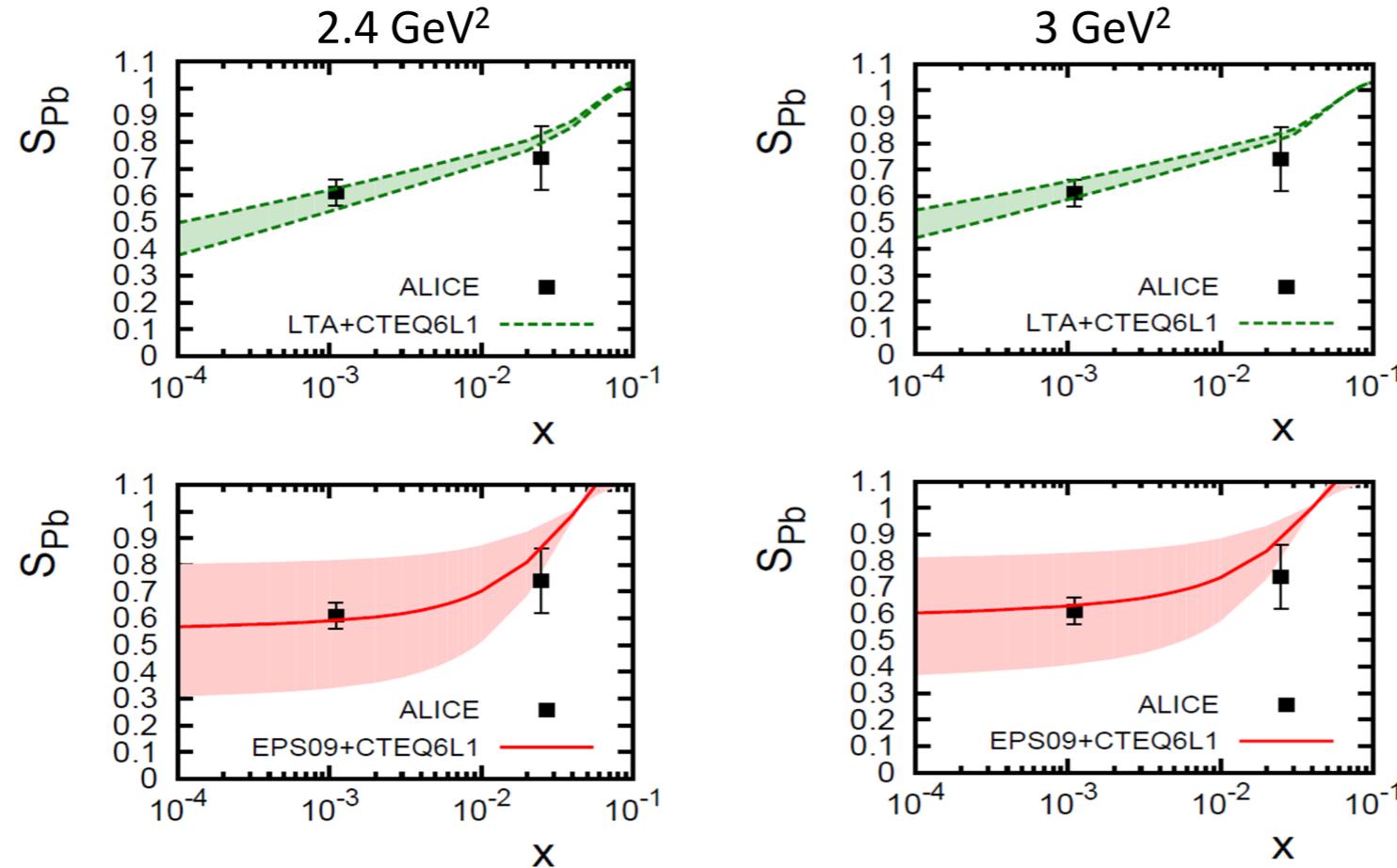
Forward Heavy-Ion Physics at LHC is exploring QDC phenomena at novel x-values

Many new interesting topics still there to study QCD and New Physics at high energies/luminosities: nuclear shadowing, saturation, excited states of vector mesons, Higgs production ...

UPC physics can be seen as the precursor for the Electron-Ion Collider

Scale dependence

- Studied in detail in Guzey, Zhalov: JHEP 1310 (2013) 207.
- Scale of 3 GeV^2 found to be most appropriate for the description of J/ψ photoproduction data



EPS09, variation of scale by factor 4:

$$R(x=0.011, Q^2 = 2.4 \text{ GeV}^2) = 0.569$$

$$R(x=0.011, Q^2 = 9.6 \text{ GeV}^2) = 0.671$$

Future measurements of heavier vector mesons (ψ' , Υ) will further elucidate the importance of the scale

pp collisions at LHC

- Hadrons are extended composite objects: Even at asymptotically large c.m. energies, ~40% of hadronic interactions are not “point-like”:

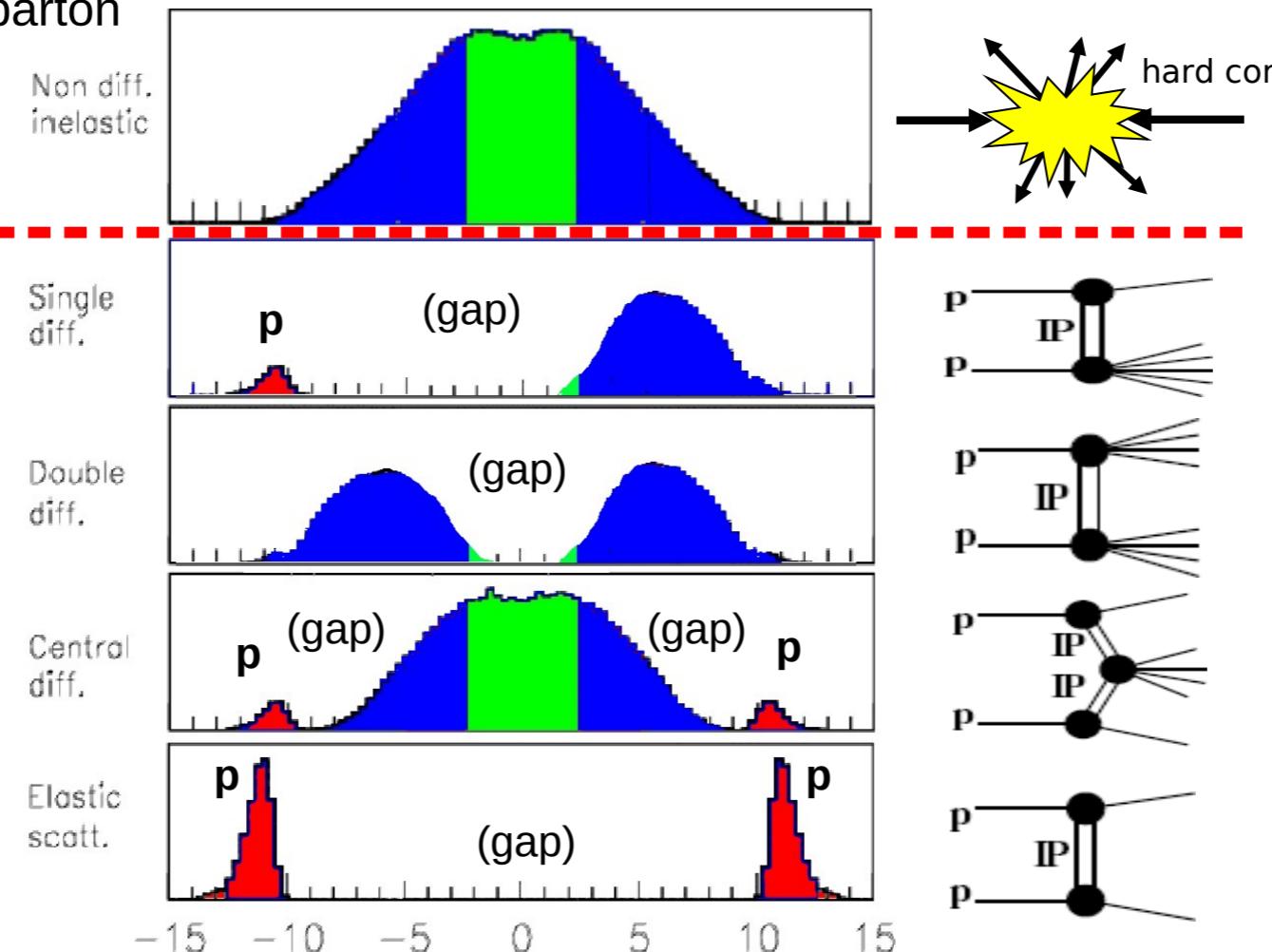
- (1) Perturbative parton-parton collisions

~60%

- (2) Diffractive, elastic

~40%

- 1 or 2 protons “intact” + 1 or 2 rapidity gaps:
- No colour flux.
- Colourless exchange with vacuum $J^{PC}=0^{++}$ quantum-numbers:
|Pomeron = 2-gluons in colour-singlet state.



- pQCD (~60 mb) + elastic (~25 mb) + diffractive (~15mb) ~ 100 mb at the LHC.

Forward physics at LHC

