

Kinematic power corrections in TMD factorization

Based on S.Piloñeta and A.Vladimirov JHEP12(2024)059 and work in progress

QCD Evolution 2025

Sara Piloñeta. May 20th 2025



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Introduction

The necessity for power corrections

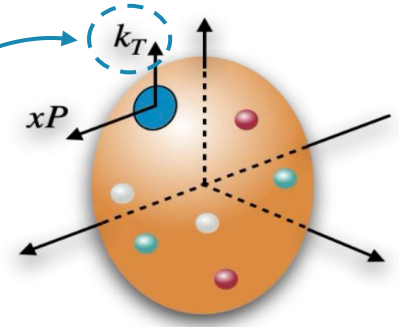
- What do we do? We are interested in **studying** the **structure** of **nucleons**

TMD factorization theorem

Transverse Momentum Dependent Distributions (TMDs)

Very powerful tool but has some **limitations**

Parton's transverse momentum



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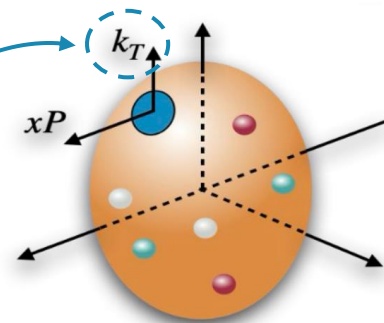
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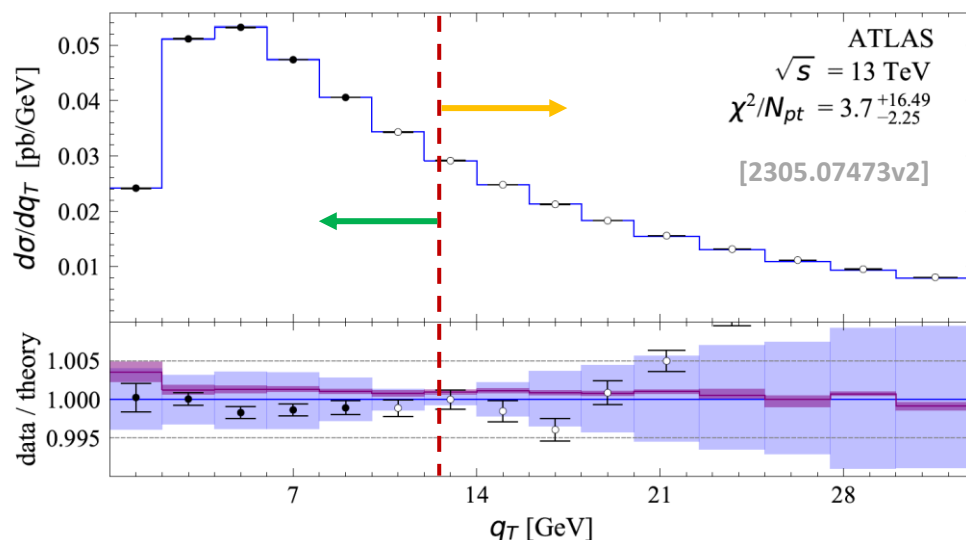
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Very powerful tool but has some **limitations**

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Z-boson production ATLAS measurement comparison



Very **precise description** at **low q_T**

Prediction systematically **lower** than the data at **larger q_T**

Signal of the **necessity** for **power corrections**

Introduction

The necessity for power corrections

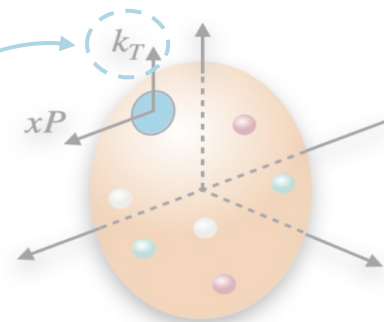
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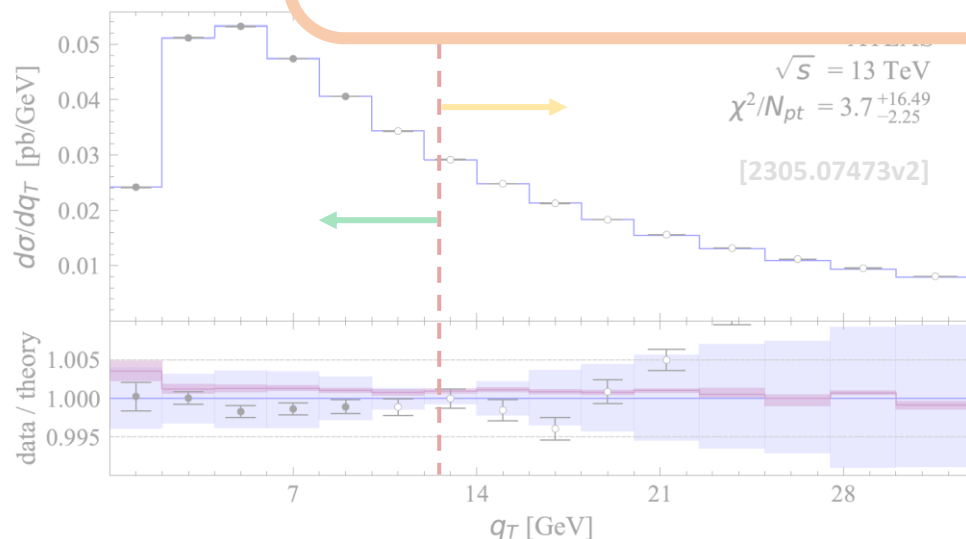
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Parton's transverse momentum



This has motivated us to explore the structure of the TMD factorization theorem deeper



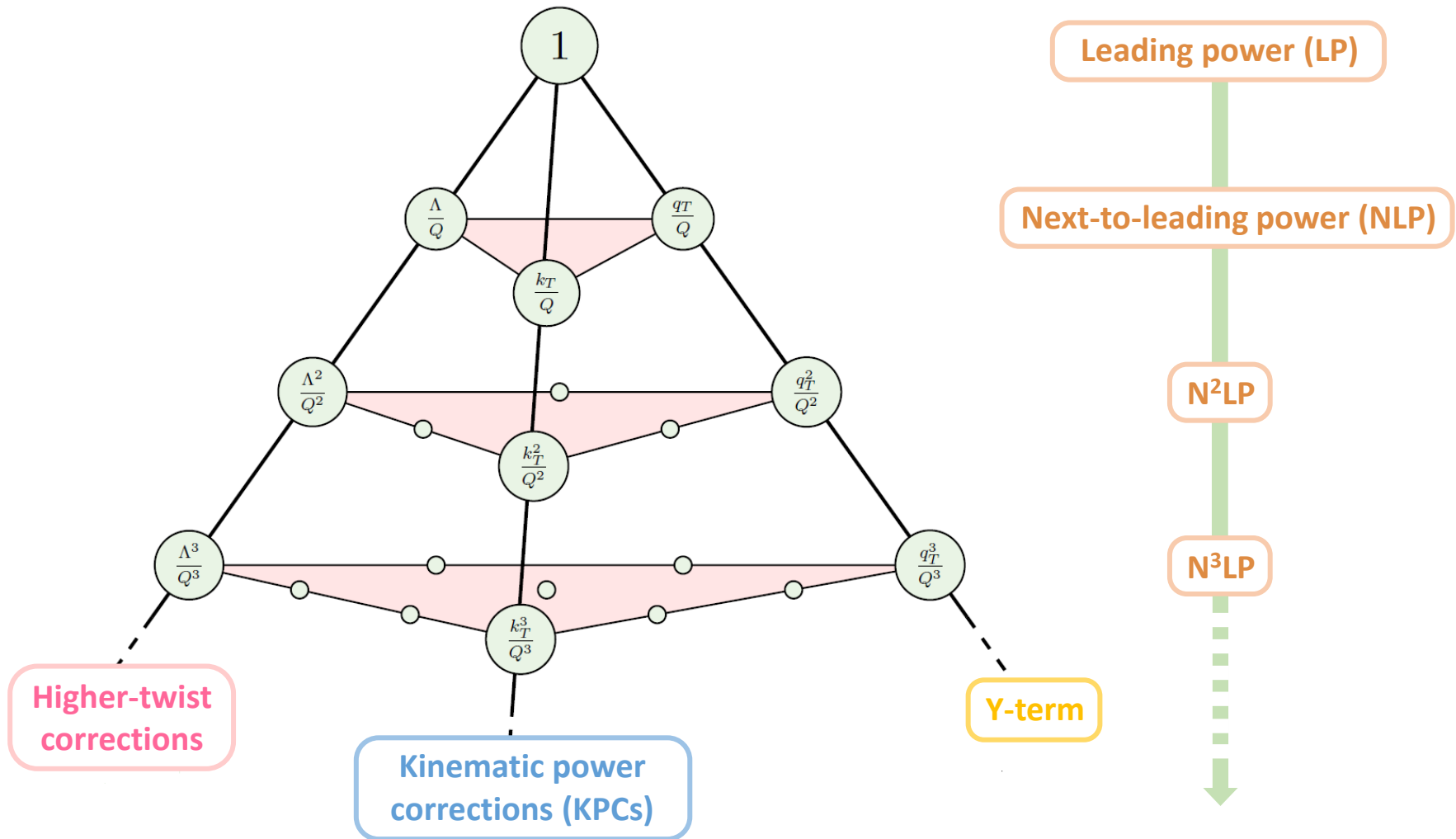
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The hierarchy of the TMD factorization theorem

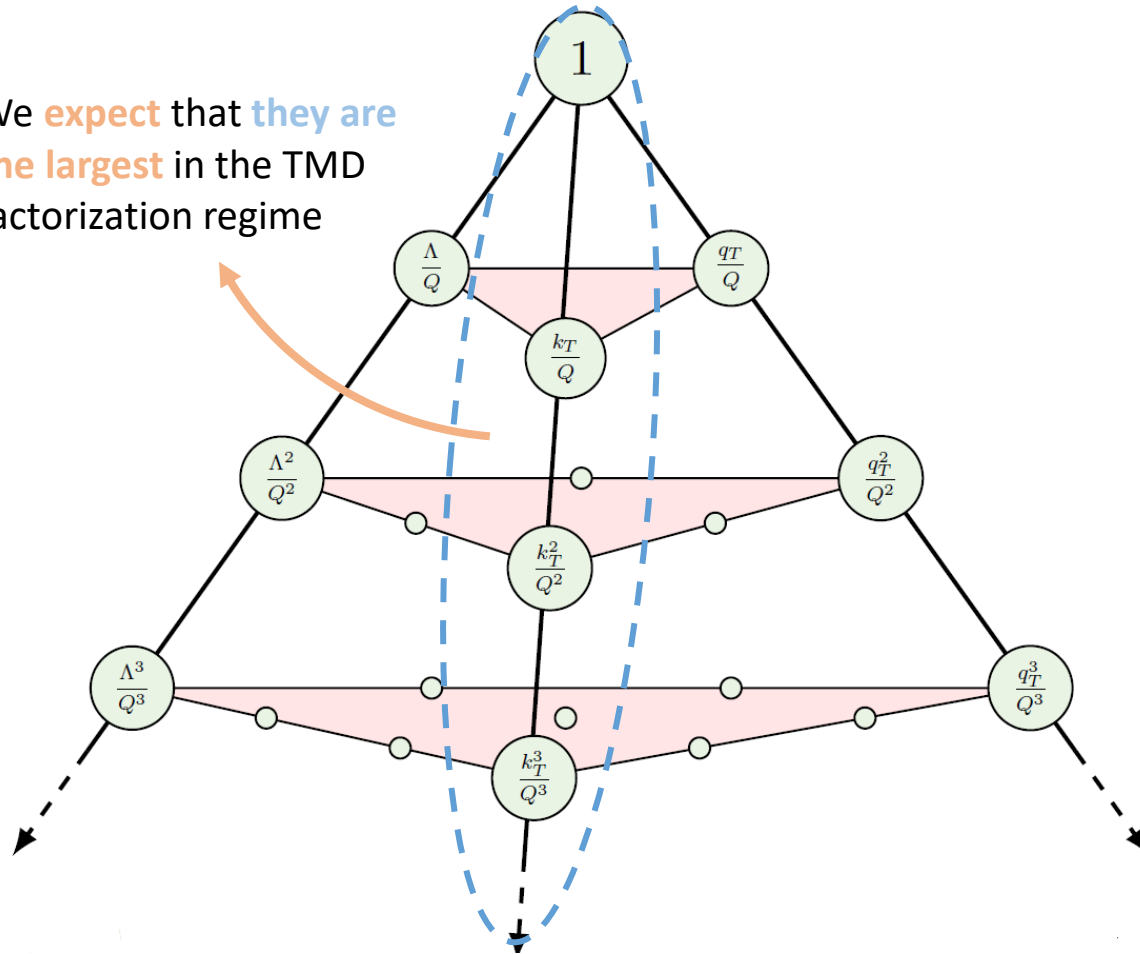
- There is a whole “pyramid” of power corrections



The TMD-with-KPCs factorization theorem

- We focus on the Kinematic power corrections (KPCs) that follow the LP term

We expect that they are the largest in the TMD factorization regime

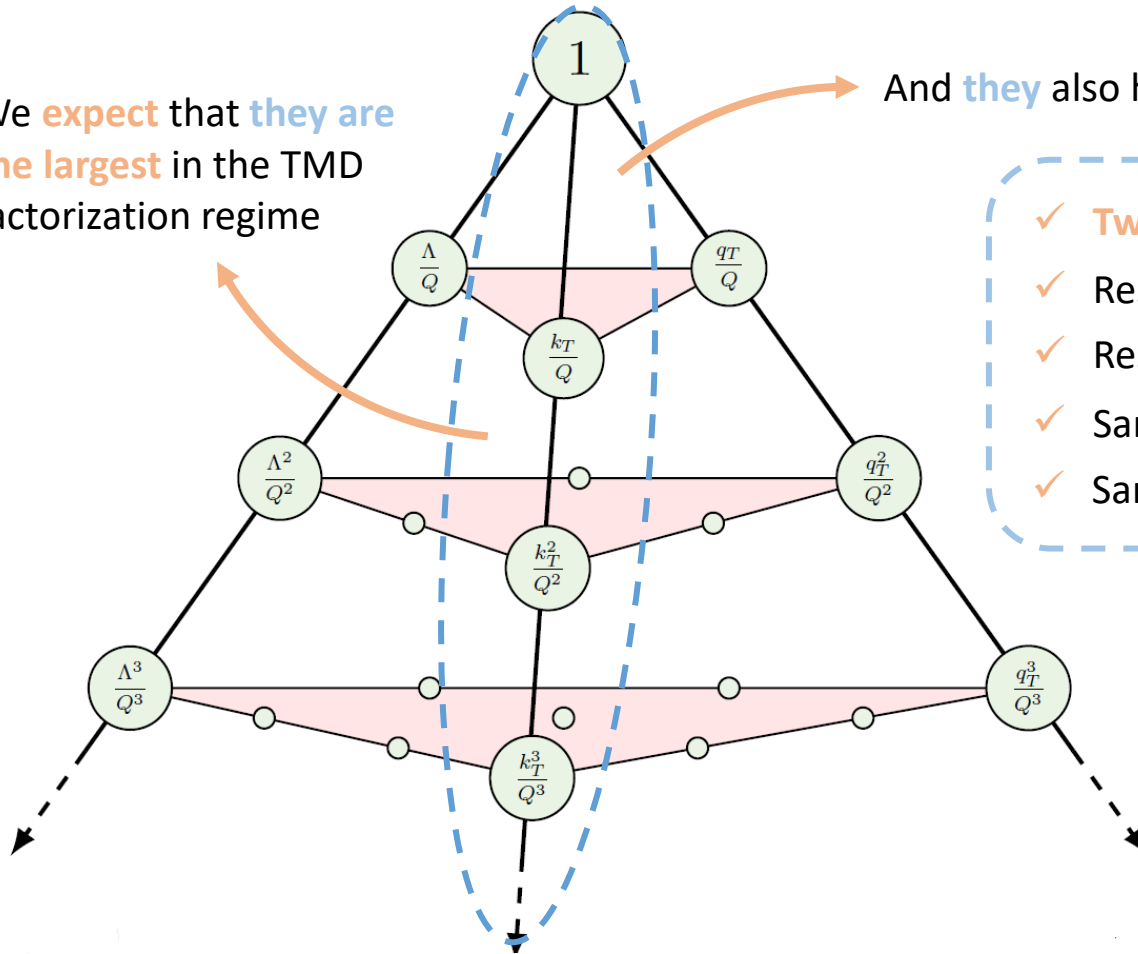


The TMD-with-KPCs factorization theorem

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And **they** also have **nice properties!**



- ✓ **Twist-2** terms at $q_T \ll Q$
- ✓ Restore **EM-gauge invariance**
- ✓ Restore **frame-invariance**
- ✓ Same TMD **distributions** as **LP**
- ✓ Same **coefficient function** as **LP**

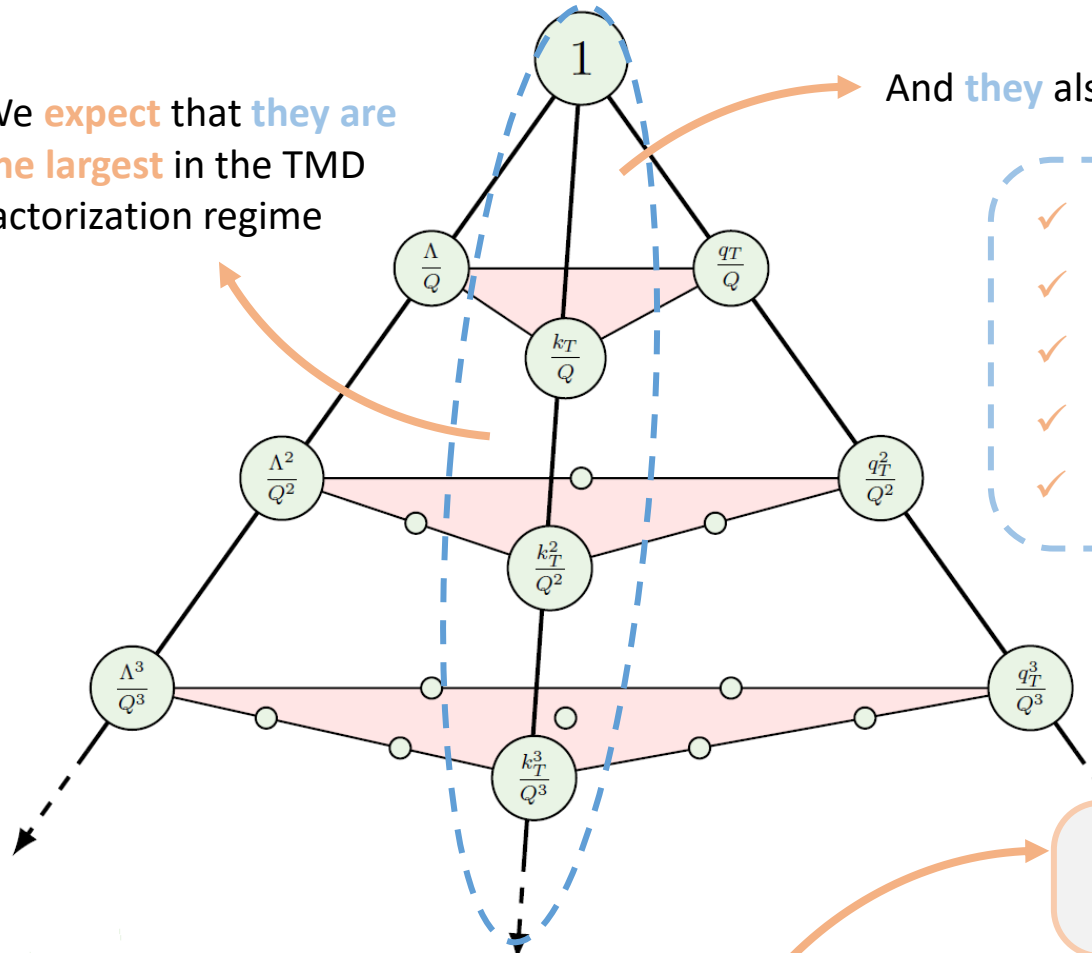
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Theoretical **framework** including **them**

**TMD-with-KPCs
factorization theorem**

For details, see [AV, 2307.13054v2]

Angular distributions of Drell-Yan leptons

- **Feasibility** of using the **TMD-with-KPCs factorization** → unpolarized Drell-Yan as a starting point

→ **Angular distributions of Drell-Yan leptons in the TMD factorization approach**

→ First practical application of this framework [S. Piloñeta, AV, 2407.06277]

Angular decomposition of the cross-section

$$\frac{d\sigma}{dp_T dy dQ d\Omega} = \frac{3}{16\pi} \left[\frac{d\sigma^{U+L}}{dp_T dy dQ} \right] \left[(1 + \cos^2 \theta) + \frac{1 - 3 \cos^2 \theta}{2} A_0 + \sin 2\theta \cos \phi A_1 + \frac{\sin^2 \theta \cos 2\phi}{2} A_2 + \sin \theta \cos \phi A_3 + \cos \theta A_4 + \sin^2 \theta \sin 2\phi A_5 + \sin 2\theta \sin \phi A_6 + \sin \theta \sin \phi A_7 \right]$$

A_U → **Unpolarized angle-integrated cross-section**

$A_{0,\dots,7}$ → Only **four** at **LP** $\{A_U, A_2, A_4, A_5\}$
 → The **rest** are **power-suppressed**

Angular distributions of Drell-Yan leptons

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Angular decomposition of the cross-section

$$\frac{d\sigma}{dp_T dy dQ d\Omega}$$

A_U

Unpolarized angle-integrated cross-section

When we incorporate the resummed KPCs we can describe them all!

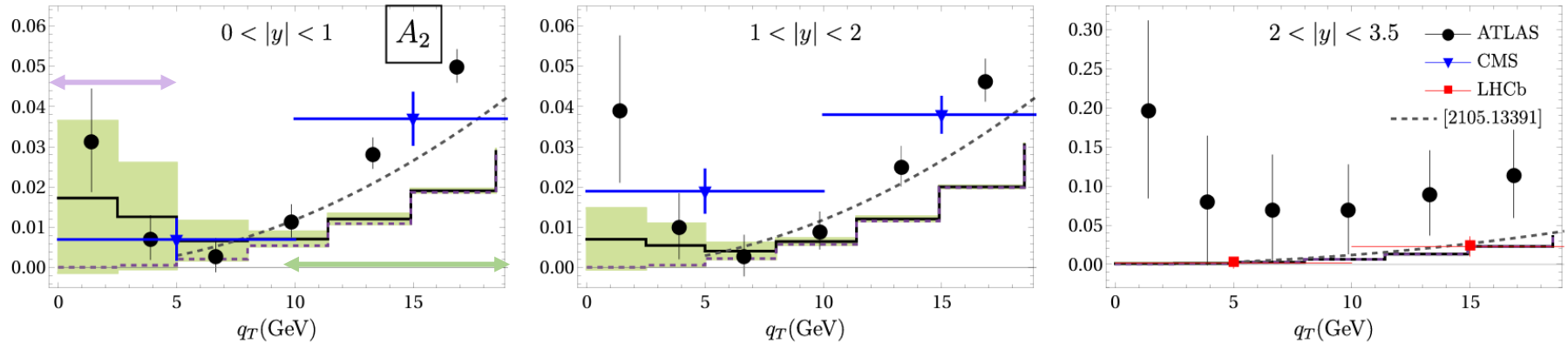
$$\begin{aligned} & \frac{\sin^2 \theta \cos 2\phi}{2} A_2 + \sin \theta \cos \phi A_3 + \cos \theta A_4 \\ & \sin^2 \theta \sin 2\phi A_5 + \sin 2\theta \sin \phi A_6 + \sin \theta \sin \phi A_7 \end{aligned}$$

$A_{0,\dots,7}$

Only **four** at **LP** $\{A_U, A_2, A_4, A_5\}$

The **rest** are **power-suppressed**

Angular distribution A_2



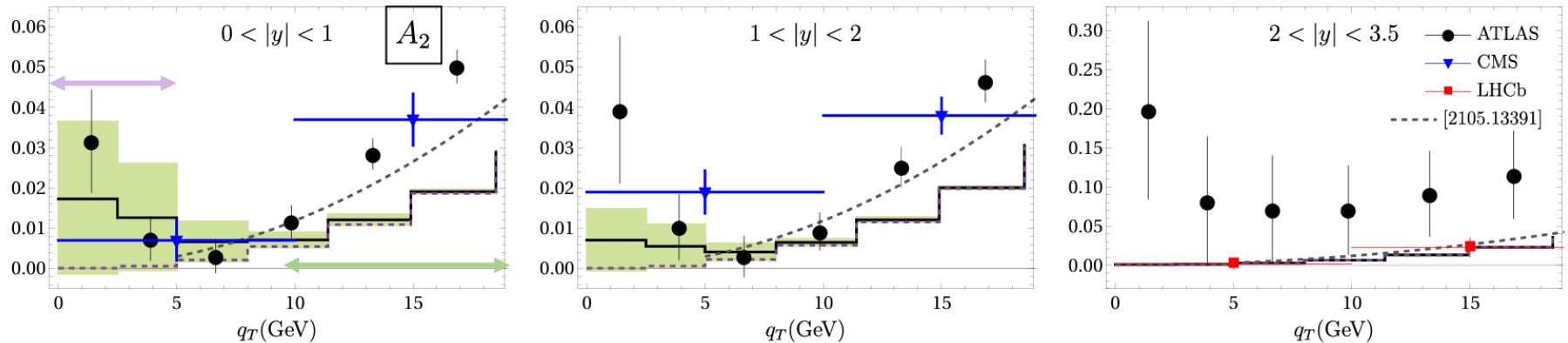
$$A_2 \sim \left[(v^2 - a^2)_q \frac{h_1^\perp h_1^\perp}{M^2} + (v^2 + a^2)_q \frac{f_1 f_1}{Q^2} \right]$$

Dominant at $q_T \rightarrow 0$ (pointing to the first term)

Dominant at larger q_T (pointing to the second term)

- Contains both **Boer-Mulders** and **unpolarized** distributions

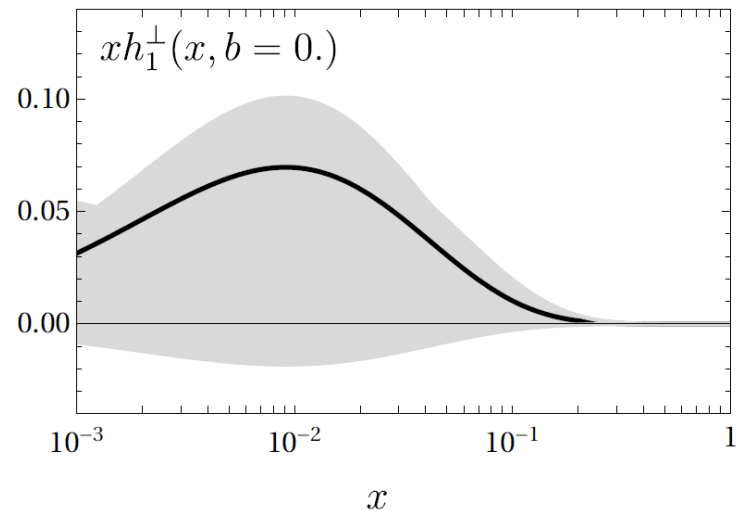
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- Boer-Mulders distribution (extraction from A_2)
 - Very primitive model
 - New fit is required
 - First evidence (?) of twist-3 effects at LHC

$$x|E(-x, 0, x)|$$



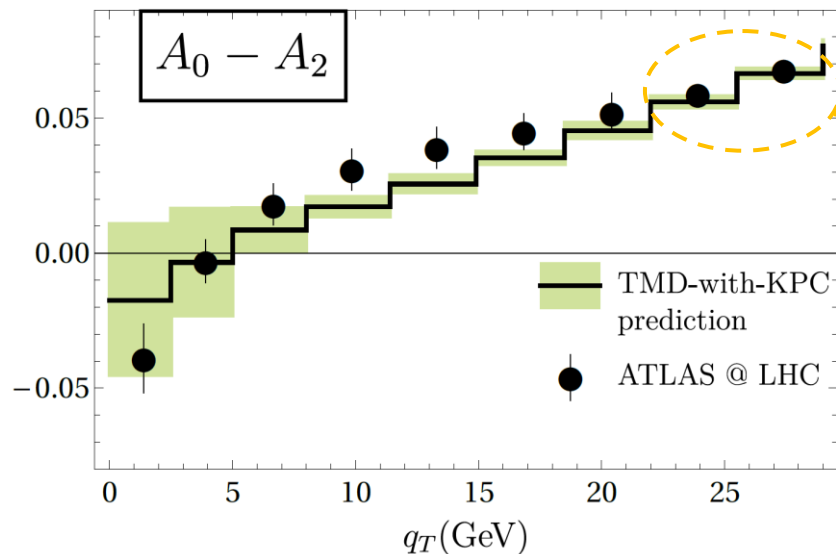
Lam-Tung relation ($A_0 - A_2$) description

- If we use the **TMD factorization** theorem the **relation does not hold at LP**

$$\Sigma_{LT} \sim k^2/M^2 h_1^\perp h_1^\perp$$

The double **Boer-Mulders** is **negligible!**

- Using the A_2 and A_0 computed **including KPCs** the theoretical **expression** also **contains f_1**



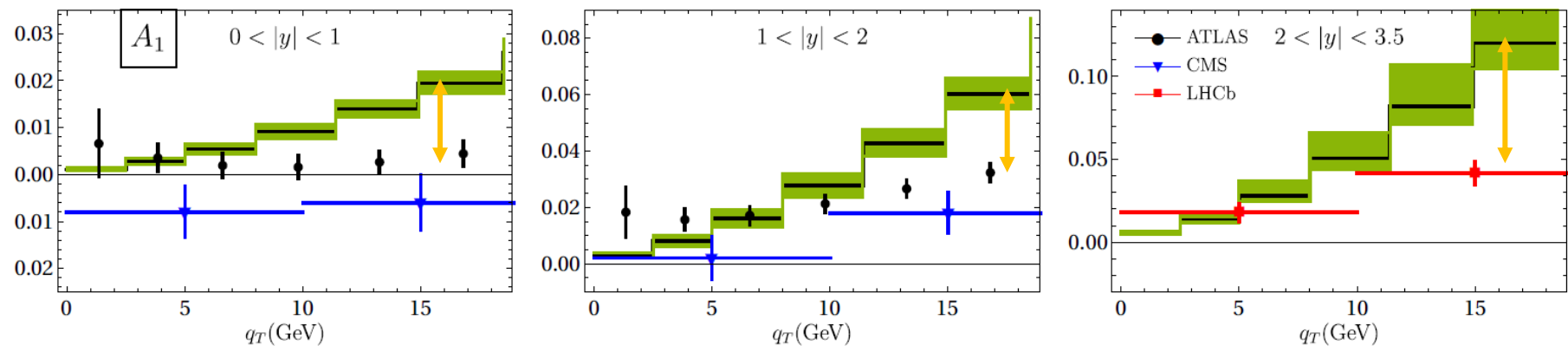
Y-term suppressed by an extra factor

Smaller q_T/Q corrections

Good agreement in contrast to A_0 and A_2 independently

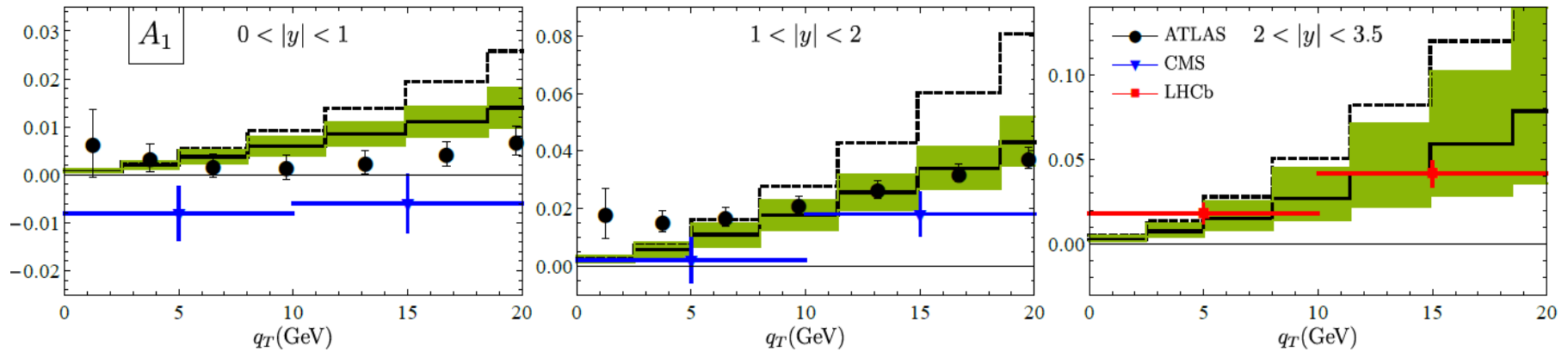
Good **description** only possible **due to KPCs inclusion!**

Angular distribution A_1

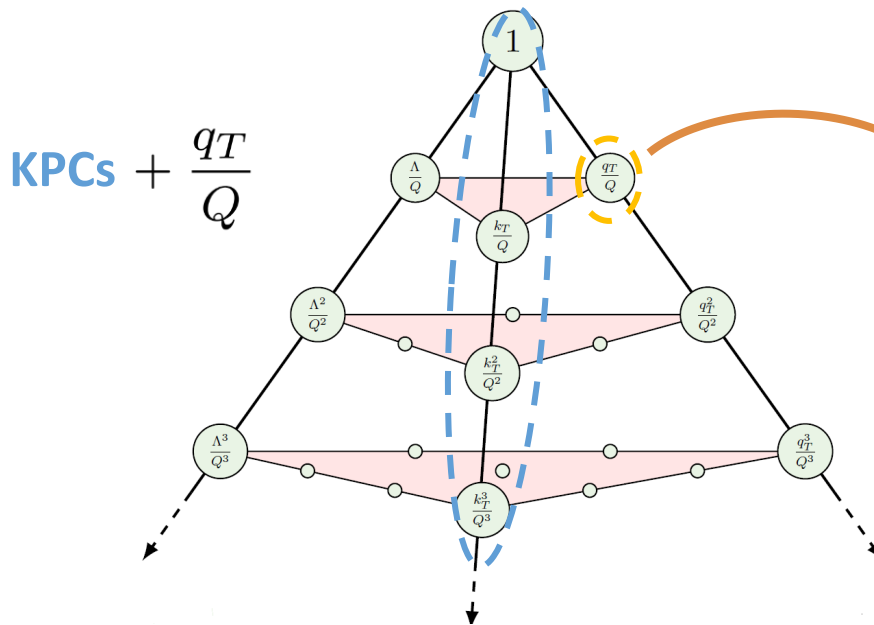


- Problems with the A_1 data description at larger values of q_T

Angular distribution A_1



- Problems with the A_1 data description at larger values of q_T



Including the computation of the leading q_T/Q correction fixes the jump!

[A.Arroyo, I.Scimemi, AV, 2503.24336]

Moving to SIDIS: structure functions computation using KPCs

- Now, let's shift our focus to **SIDIS** → essential process for probing hadron structure

Cross-section decomposition in terms of structure functions

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h d\mathbf{p}_\perp^2} = \frac{\alpha_{\text{em}}^2}{xy Q^2} \frac{y^2}{2(1-\varepsilon)} \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos \phi_h F_{UU}^{\cos \phi_h} \right. \\ \left. + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin \phi_h F_{LU}^{\sin \phi_h} + \dots \right.$$

→ We want to compute them including KPCs

$$\left\{ F_{UU,T} + \dots \right\} = \frac{x}{4z} \frac{1-\varepsilon}{Q^2} L_{\mu\nu} W^{\mu\nu}$$

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$$\left\{ F_{UU,T} + \dots \right\} = \frac{x}{4z} \frac{1-\varepsilon}{Q^2} \underbrace{L_{\mu\nu}}_{\textcircled{1}} W^{\mu\nu}$$

1 Lepton tensor conveniently decomposed via the tensors P^μ , q^μ , $p_\perp^\mu$

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$$\left\{ F_{UU,T} + \dots \right\} = \frac{x}{4z} \frac{1-\varepsilon}{Q^2} L_{\mu\nu} \underbrace{W^{\mu\nu}}_{(2)}$$

(2) Hadron tensor computed using the TMD-with-KPCs factorization theorem

→ Same coefficient function as LP $(q_\mu W^{\mu\nu} = 0)$
→ Main difference with LP → Convolution integral

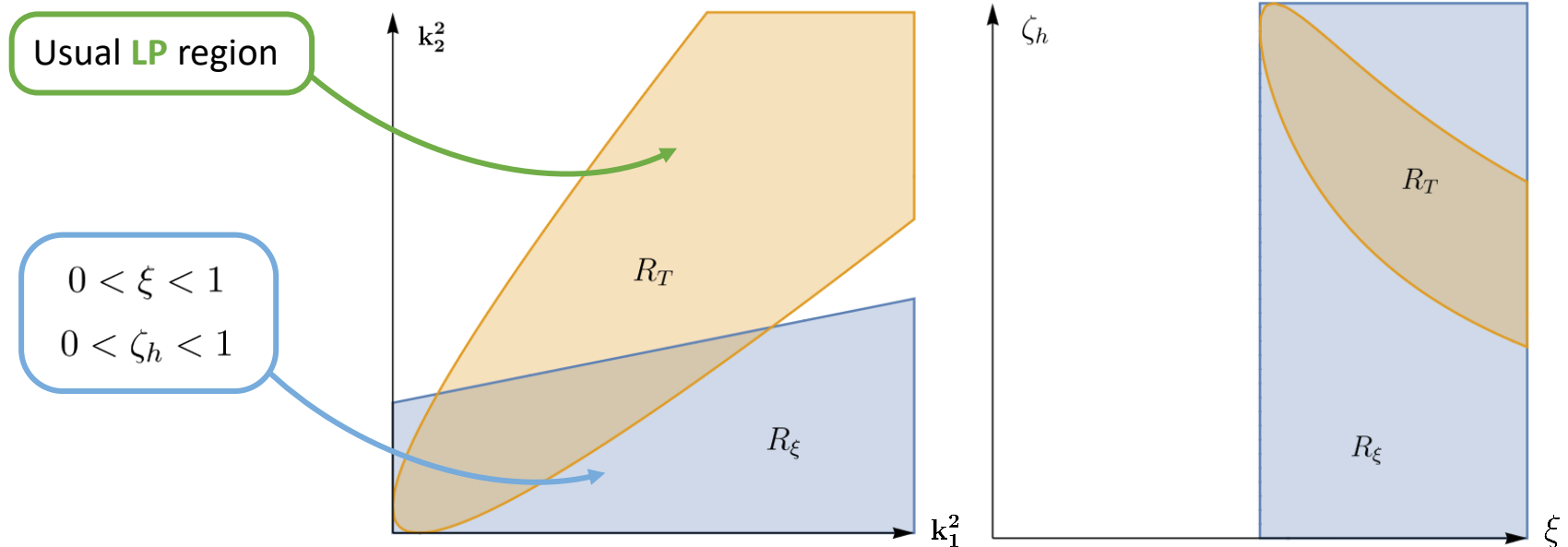
A closer look at the convolution integral

- The **convolution integral** is **more complicated** now

(LP case) $\mathcal{C}_{LP}[A, f_1, D_1] \sim \int d^2\mathbf{k}_1 d^2\mathbf{k}_2 \delta(q_T + \mathbf{k}_1 - \mathbf{k}_2) f_1(x_1, \mathbf{k}_1^2) D_1(z_1, \mathbf{k}_2^2)$

Additional dependance
coming from extra δ -functions

(KPCs case) $\mathcal{C}_{KPC}[A, f_1, D_1] \sim \int d^2\mathbf{k}_1 d^2\mathbf{k}_2 \delta(q_T + \mathbf{k}_1 - \mathbf{k}_2) f_1(\xi(x_1, \mathbf{k}_{1,2}^2), \mathbf{k}_1^2) D_1(\zeta_h(z_1, \mathbf{k}_{1,2}^2), \mathbf{k}_2^2)$



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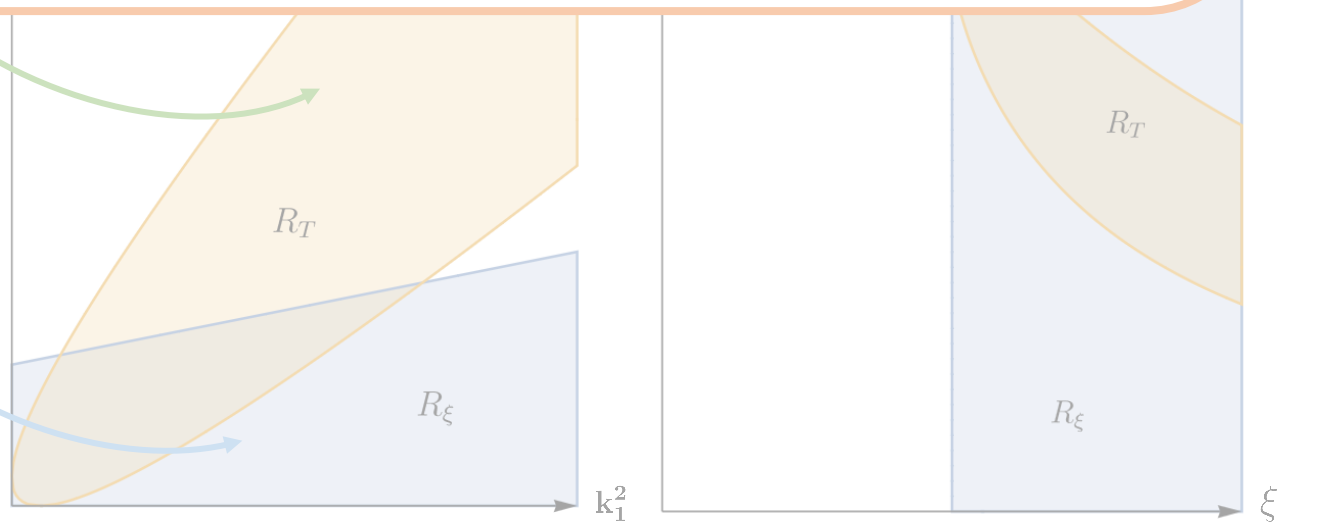
The convolution integral and the KPCs have been implemented in
arTeMiDe, although a few improvements are still needed

github.com/VladimirovAlexey/artemide-development

Usual

$$0 < \xi < 1$$

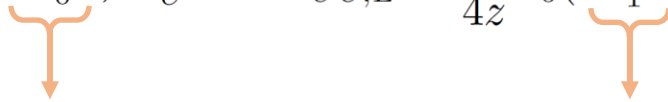
$$0 < \zeta_h < 1$$



Focusing on $F_{UU,T}$ and $F_{UU,L}$

- I have obtained the **theoretical expressions** for all of them using the **TMD-with-KPCs factorization**

↳ But in **this talk** I am going to **focus on** $F_{UU,T}$ and $F_{UU,L}$

$$F_{UU,T} = \frac{x}{4z} F_0 (\mathcal{S}_1^{\mu\nu} - \mathcal{S}_0^{\mu\nu}) W_U^{\mu\nu} \quad F_{UU,L} = \frac{x}{4z} F_0 (2\mathcal{S}_1^{\mu\nu}) W_U^{\mu\nu}$$

$$\mathcal{S}_0^{\mu\nu} = g^{\mu\nu} - \frac{q^\mu q^\nu}{Q^2} \quad \mathcal{S}_1^{\mu\nu} = \frac{(2xP^\mu + q^\mu)(2xP^\nu + q^\nu)}{(1 + \gamma^2)Q^2}$$

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The theoretical **expression** for $F_{UU,T}$ **including KPCs** is

$$F_{UU,T} = \frac{x}{2z} Q^2 \mathcal{C}_{KPC}[1, f_1 D_1] + \frac{1}{2} F_{UU,L}$$

$$\mathcal{C}_{LP}[1, f_1 D_1] + \varepsilon F_{UU,L} \quad \xrightarrow{\text{Cross-section gets modified}} \quad \mathcal{C}_{KPC}[1, f_1 D_1] + \left(\frac{1}{2} + \varepsilon \right) F_{UU,L}$$

The **inclusion** of the **KPCs** implies that $F_{UU,T}$ is **not exactly 1** as it was **in the LP case!**

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So far, I have **studied** the **KPCs impact** on

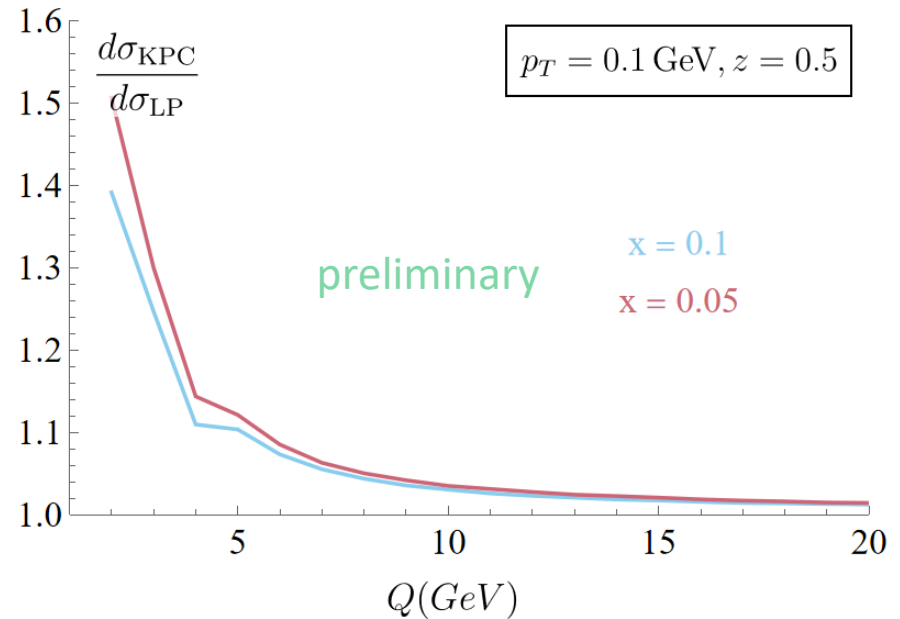
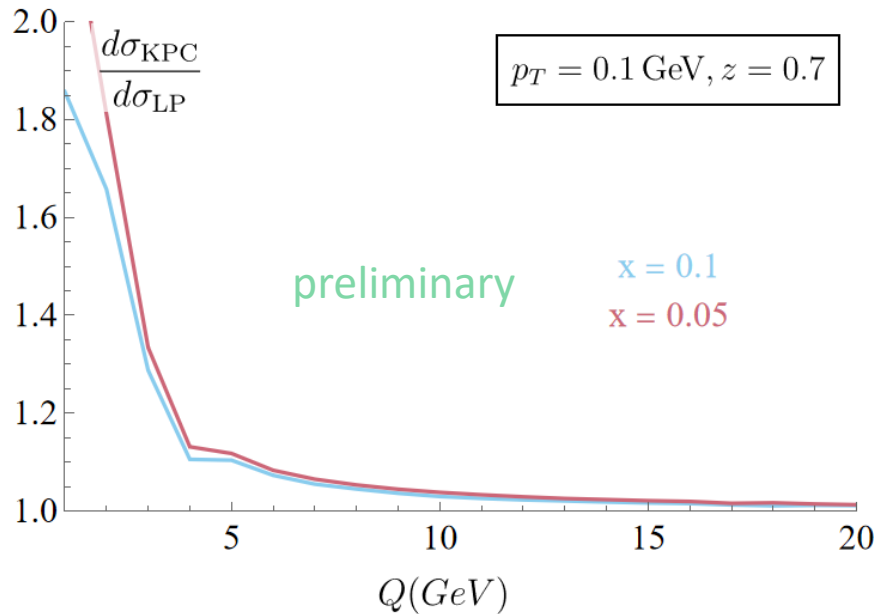
The **SIDIS cross-section**

The **structure functions** $F_{UU,T}$ and $F_{UU,L}$

Work is still in progress!

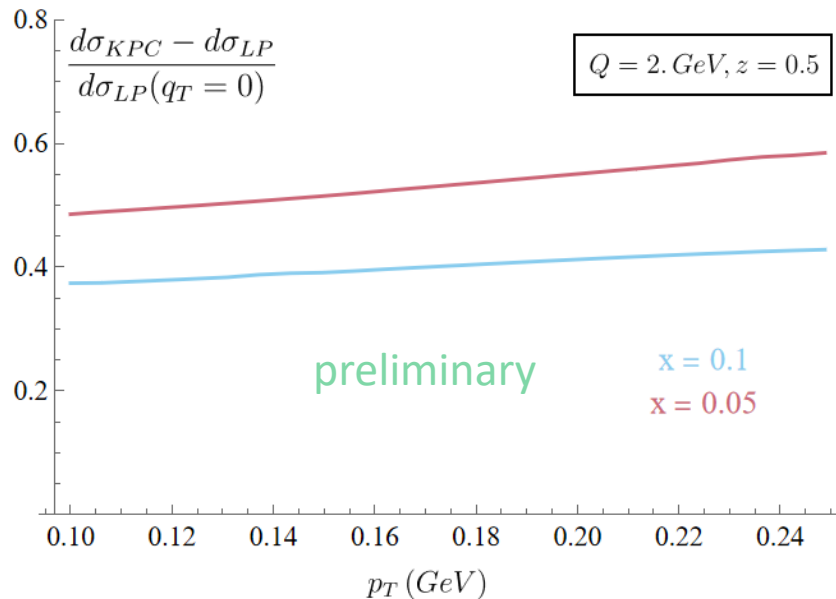
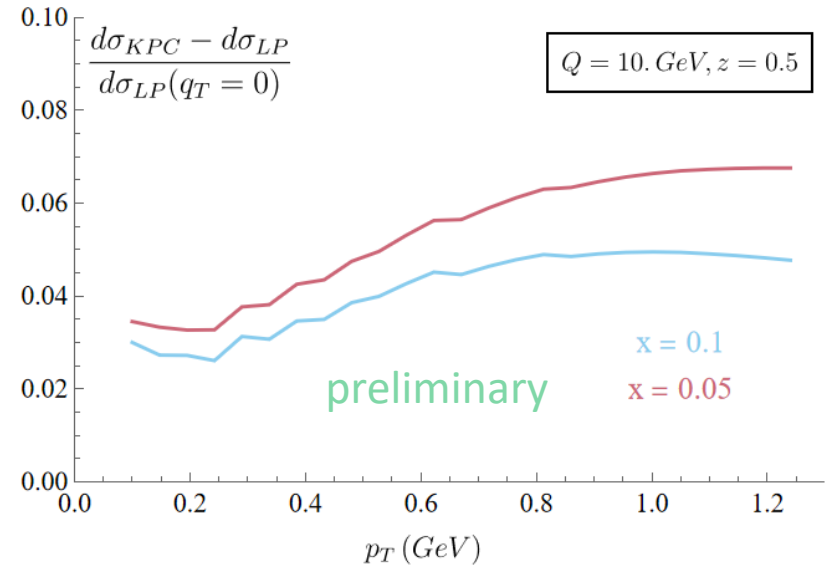
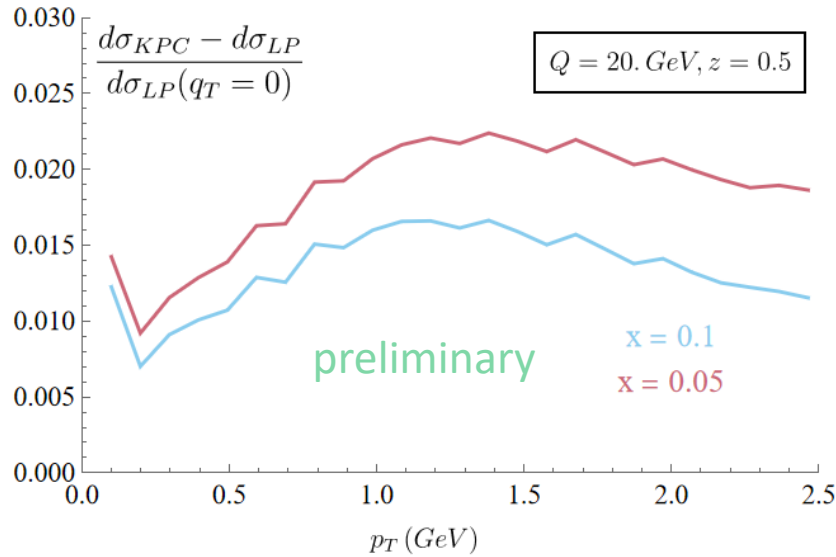
Ratio of KPC-summed to LP cross-sections

- Studying the size of $\mathcal{C}_{KPC}[1, f_1 D_1] + \frac{1}{2} F_{UU,L}$ in comparison with $\mathcal{C}_{LP}[1, f_1 D_1]$



- ✓ Cross-section with KPCs is bigger
- Specially at low energies
- Satisfies that $\lim_{Q \rightarrow \infty} d\sigma_{KPC} = d\sigma_{LP}$

Cross-sections difference relative to the $q_T = 0$ LP cross-section



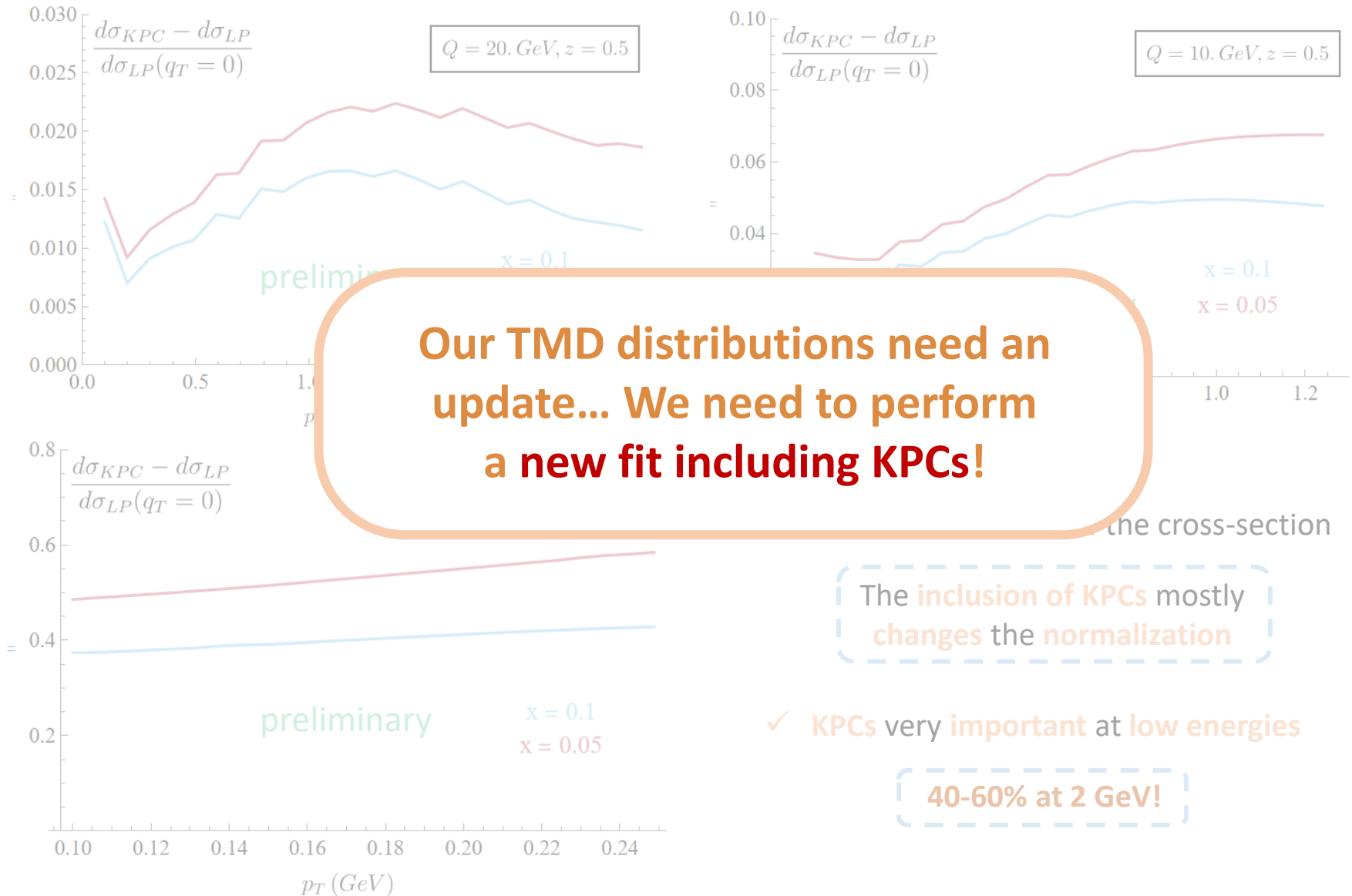
✓ Almost flat increase of the cross-section

The **inclusion of KPCs** mostly **changes the normalization**

✓ **KPCs** very **important at low energies**

40-60% at 2 GeV!

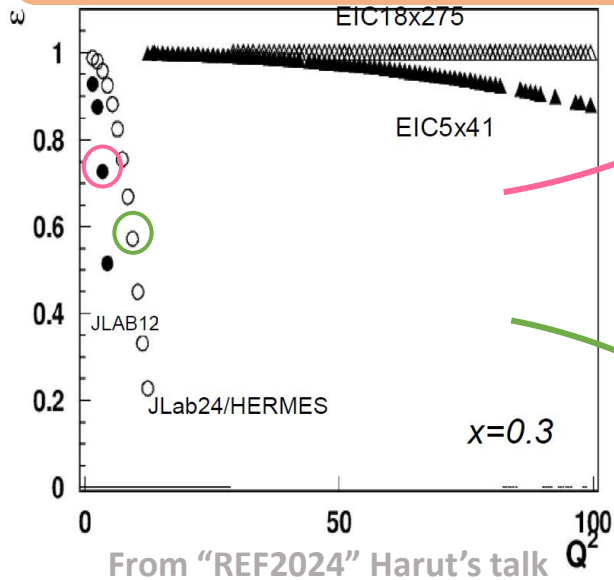
Cross-sections difference relative to the $q_T = 0$ LP cross-section



The contribution of longitudinal photons

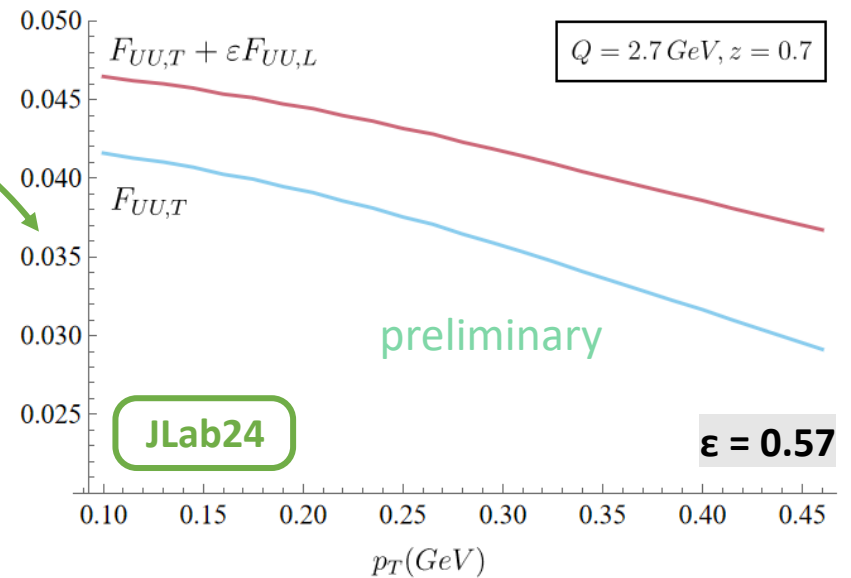
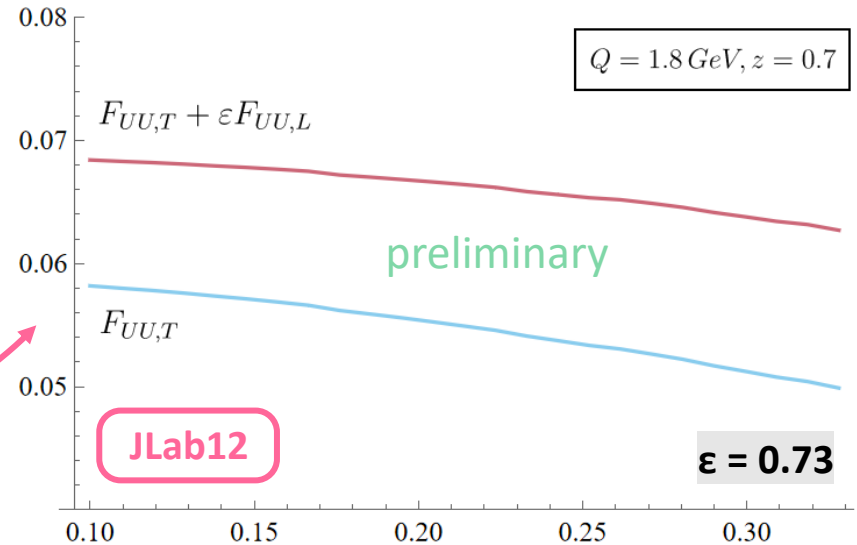
What about the contribution of
longitudinal photons?

ϵ values for different experiments



From "REF2024" Harut's talk

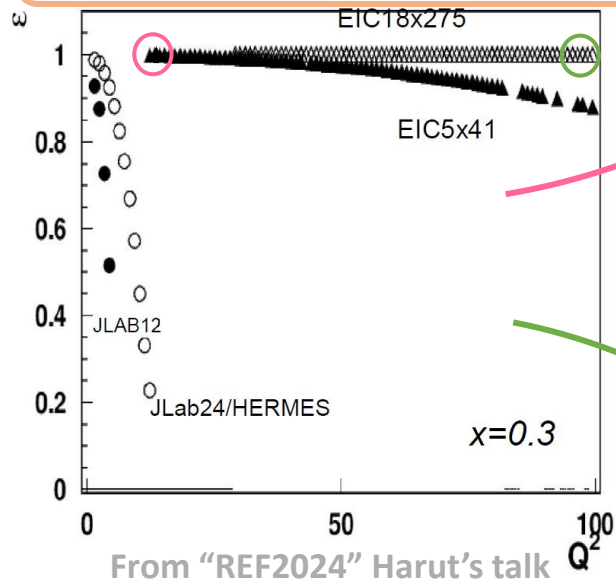
At low energies, the contribution of the $F_{UU,L}$ structure function is not negligible!



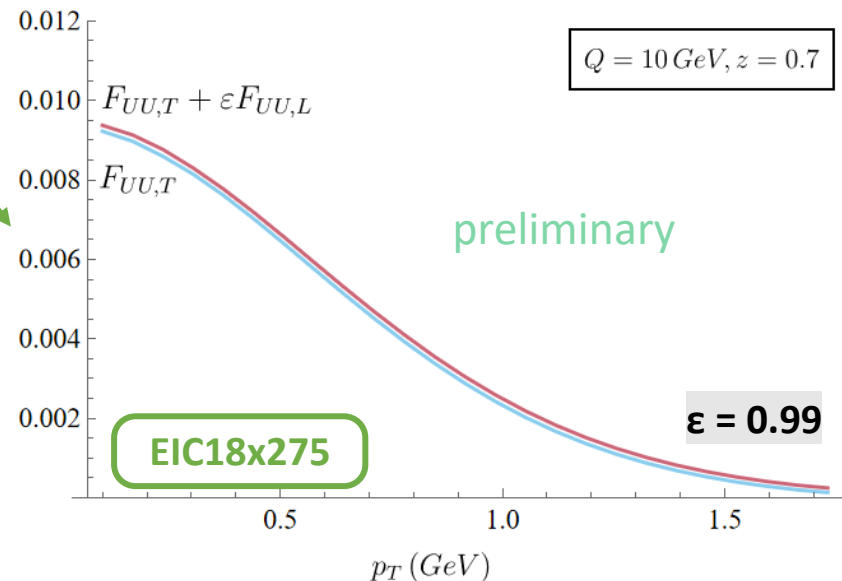
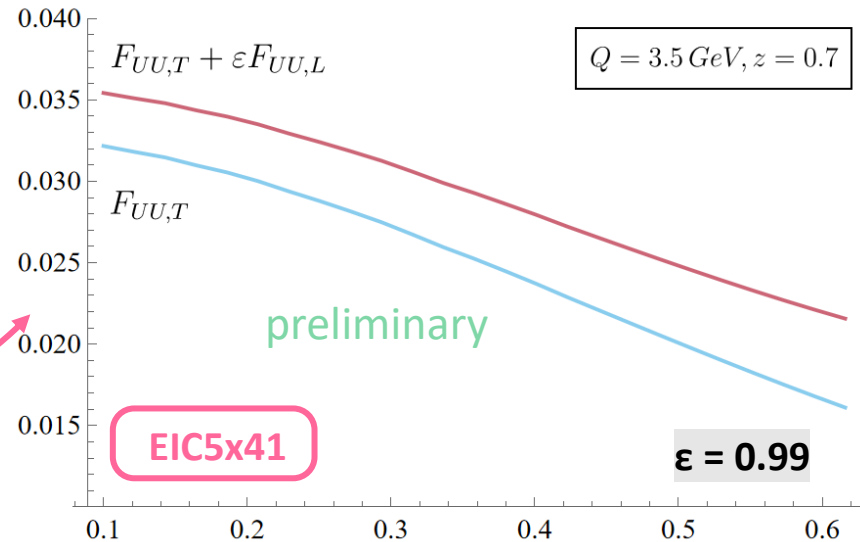
The contribution of longitudinal photons

What about the contribution of
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As Q grows, the contribution of $F_{UU,L}$ becomes less important, even if ϵ is very close to 1



Conclusions

- The **TMD-with-KPCs factorization theorem** [AV, 2307.13054v2] has been **tested**
 - ❑ The **angular distributions** of **Drell-Yan** leptons can be satisfactorily **described**
 - ❑ It gives a nice **description** of the **Lam-Tung relation**
- The **subleading** $F_{UU,T}$ and $F_{UU,L}$ SIDIS structure functions have been **computed** using it
 - ❑ The **SIDIS cross-section grows** when including KPCs
 - ❑ The $F_{UU,L}$ contribution is **not negligible** at **low energies** like **2 – 5 GeV**

Our TMD distributions need an update... We need to perform a new fit including KPCs!

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Thank you for your attention! 😊