

First Neural Network extraction of unpolarized TMDs



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Map Collaboration

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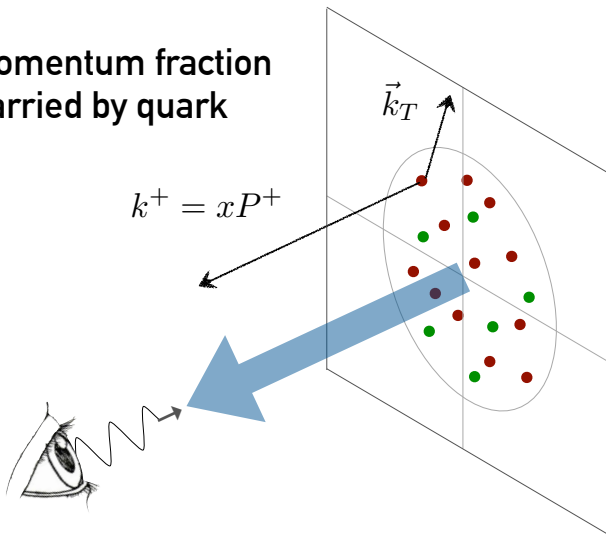
based on [arXiv:2502.04166](https://arxiv.org/abs/2502.04166)

TMDs: 3D maps

in momentum space

x : momentum fraction
carried by quark

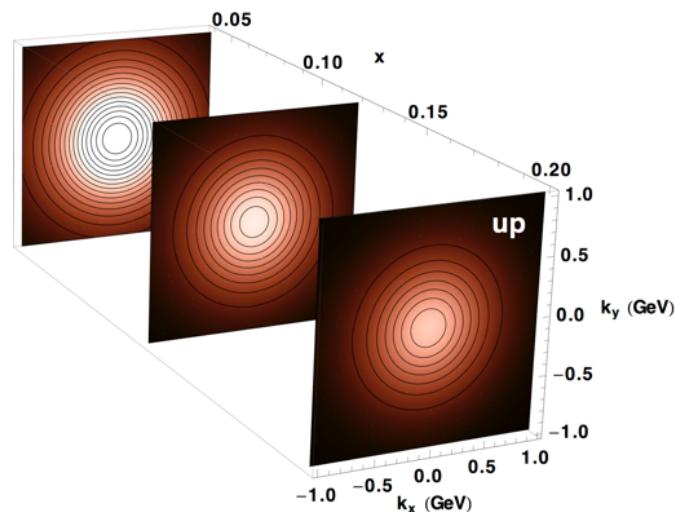
$$k^+ = xP^+$$



Transverse Momentum
Distributions

TMDs

$$f^q(x, \mathbf{k}_T)$$



Transverse Momentum Distributions

the distribution of quarks sharply depends on the orientation of their spins

Leading Quark TMDPDFs  Nucleon Spin  Quark Spin

this work

		Quark Polarization		
		Un-Polarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 = \text{Unpolarized}$		$h_1^\perp = \text{Boer-Mulders}$
	L		$g_1 = \text{Helicity}$	$h_{1L}^\perp = \text{Worm-gear}$
	T	$f_{1T}^\perp = \text{Sivers}$	$g_{1T}^\perp = \text{Worm-gear}$	$h_1 = \text{Transversity}$ $h_{1T}^\perp = \text{Pretzelosity}$

from the TMD handbook
[arXiv:2304.03302](https://arxiv.org/abs/2304.03302)

TMDs — *unpolarized quark TMD PDF*

matching to the collinear region

factorizes as **hard** and longitudinal non-perturbative

$b_T \ll 1/\Lambda_{\text{QCD}}$

$$F_{f/P}(x, \mathbf{b}_T; \mu, \zeta) = \sum_j C_{f/j}(x, b_*; \mu_b, \zeta_F) \otimes f_{j/P}(x, \mu_b)$$

CS and RGE evolution

$$\times \exp \left\{ K(b_*; \mu_b) \ln \frac{\sqrt{\zeta_F}}{\mu_b} + \int_{\mu_b}^{\mu} \frac{d\mu'}{\mu'} \left[\gamma_F - \gamma_K \ln \frac{\sqrt{\zeta_F}}{\mu'} \right] \right\}$$

non-perturbative transverse content

$$\times \exp \left\{ g_{j/P}(x, b_T) + g_K(b_T) \ln \frac{\sqrt{\zeta_F}}{\sqrt{\zeta_{F,0}}} \right\}$$

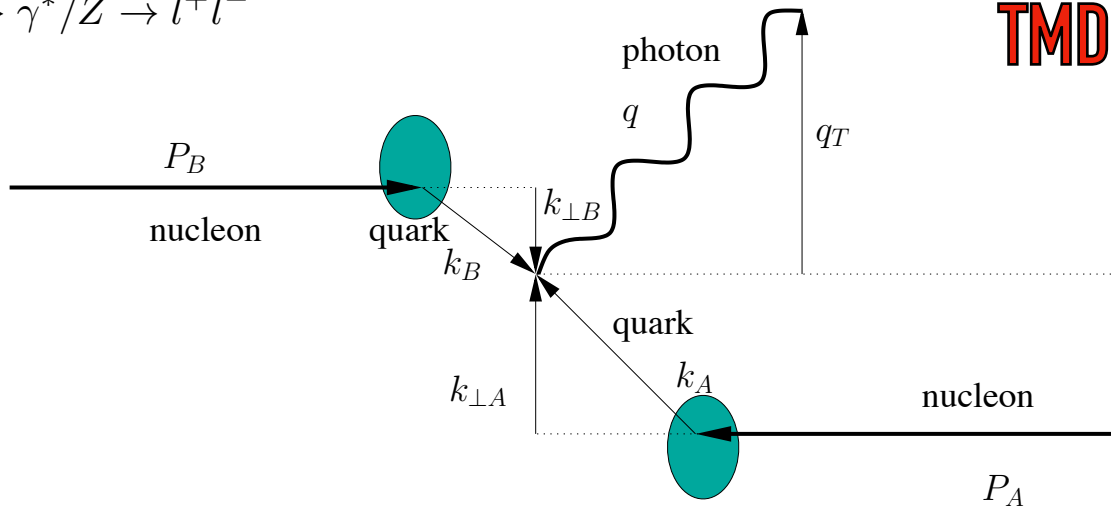
parametrized and fitted to data

Drell-Yan

$$N(P_A) + N(P_B) \rightarrow \gamma^*/Z \rightarrow l^+ l^-$$

for $q_T \ll Q$

TMD factorization



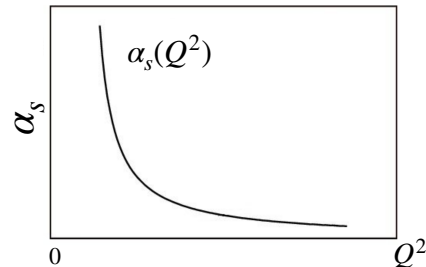
$$\left(\frac{d\sigma}{dq_T} \right)_{\text{res.}} \propto H(Q, \mu) \int \frac{d^2 \mathbf{b}}{4\pi} e^{i\mathbf{b} \cdot \mathbf{q}_T} x_1 f_1^q(x_1, \mathbf{b}; \mu, \zeta_1) x_2 f_1^{\bar{q}}(x_2, \mathbf{b}; \mu, \zeta_2)$$

b^* prescription

$$\frac{d\sigma}{dq_T} \propto \int \frac{d^2\mathbf{b}}{4\pi} e^{i\mathbf{b}\cdot\mathbf{q}_T} f_1^q(x_1, \mathbf{b}) f_1^{\bar{q}}(x_2, \mathbf{b})$$

integration up to infinity

cross section

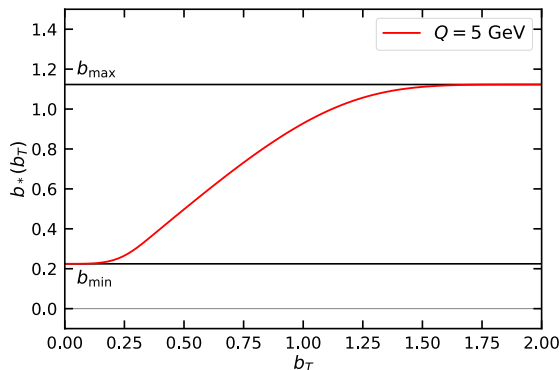


when b_T becomes large

$$\alpha_s(\mu_b) = \alpha \left(\frac{2e^{-\gamma_E}}{b} \right) \gg 1$$

invalidates perturbative calculations

$\Rightarrow b_{\max}$



$$b_{\max} = 2e^{-\gamma_E}$$

$$b_{\min} = 2e^{-\gamma_E}/Q$$

necessary to introduce a prescription

$$b_*(b) = b_{\max} \left(\frac{1 - \exp\left(-\frac{b^4}{b_{\max}^4}\right)}{1 - \exp\left(-\frac{b^4}{b_{\min}^4}\right)} \right)^{\frac{1}{4}}$$

b^* prescription

and definition of f_{NP}

└─▶ perturbative

$$f(x, b; \mu, \zeta) = \left[\frac{f(x, b; \mu, \zeta)}{f(x, b_*(b); \mu, \zeta)} \right] f(x, b_*(b); \mu, \zeta)$$

non perturbative



$f_{\text{NP}}(x, b, \zeta)$

→ fit to data

Non perturbative function depends on
the choice of b^* -prescription

Motivation

What would we like to accomplish?

goal: fit for the first time TMDs with NN

proof of concept

* do NN offer an advantage in fitting (real) data?

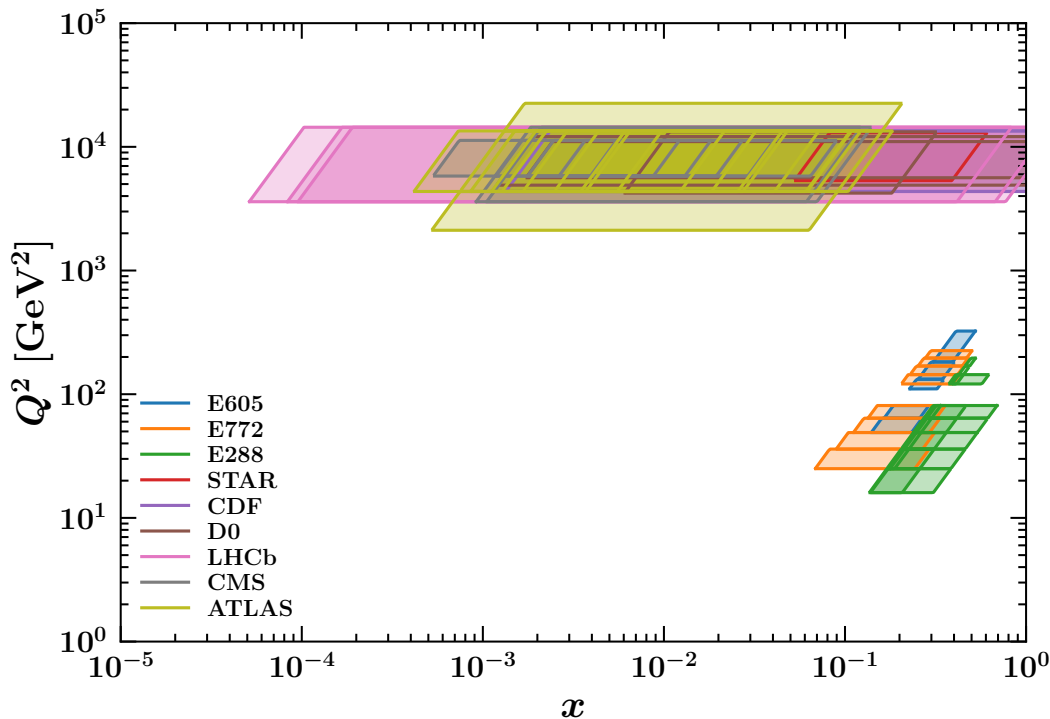
* we want to **test the methodology**



with **closure tests**

Data kinematical coverage

high and low energy Drell-Yan data



TMD region

$$q_T/Q < 0.2$$

total number of data: 482

included all available DY data with the exception
of PHENIX (2 data points)

Parametrization

of the non-perturbative part of $f_1(x, k_T)$

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[arXiv:2502.04166](https://arxiv.org/abs/2502.04166)

parametrization

$$f_{\text{NP}}(x, b_T; \zeta) = \frac{\text{NN}(x, b_T, \{p_i\})}{\text{NN}(x, 0, \{p_i\})} \exp\left[-g_2^2 b_T^2 \ln\left(\frac{\zeta}{Q_0^2}\right)\right]$$

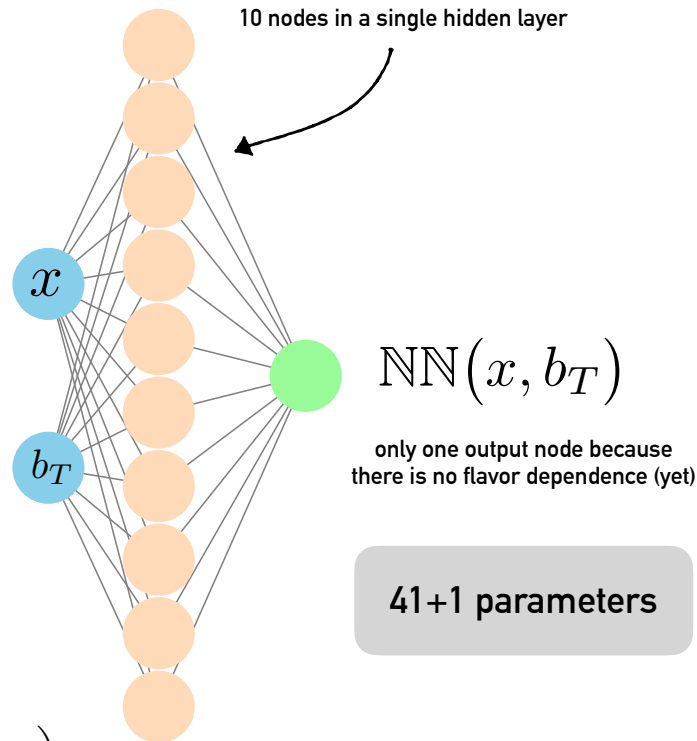
- physically required constraints

$$f_{\text{NP}} \rightarrow 1 \quad \text{for} \quad b_T \rightarrow 0$$

with activation function

$$\sigma(z) = \frac{1}{2} \left(1 + \frac{z}{1 + |z|} \right)$$

Neural Network



Parameterization

of the non-perturbative part of TMDs

"analytic", "with a functional form"

TMD PDF

$$f_{1\text{NP}}(x, b_T^2) \propto \text{F.T. of } \left(e^{-\frac{k_\perp^2}{g_{1A}}} + \lambda_B k_\perp^2 e^{-\frac{k_\perp^2}{g_{1B}}} + \lambda_C e^{-\frac{k_\perp^2}{g_{1C}}} \right)$$

Gaussians

weighted Gaussian

$$g_1(x) = N_1 \frac{(1-x)^\alpha x^\sigma}{(1-\hat{x})^\alpha \hat{x}^\sigma}$$

NP evolution

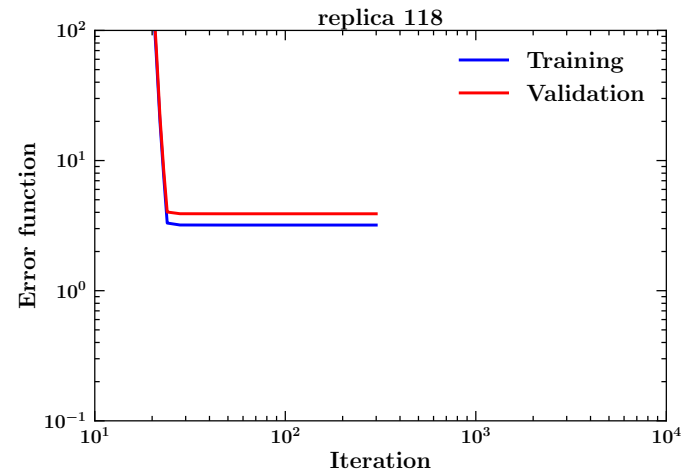
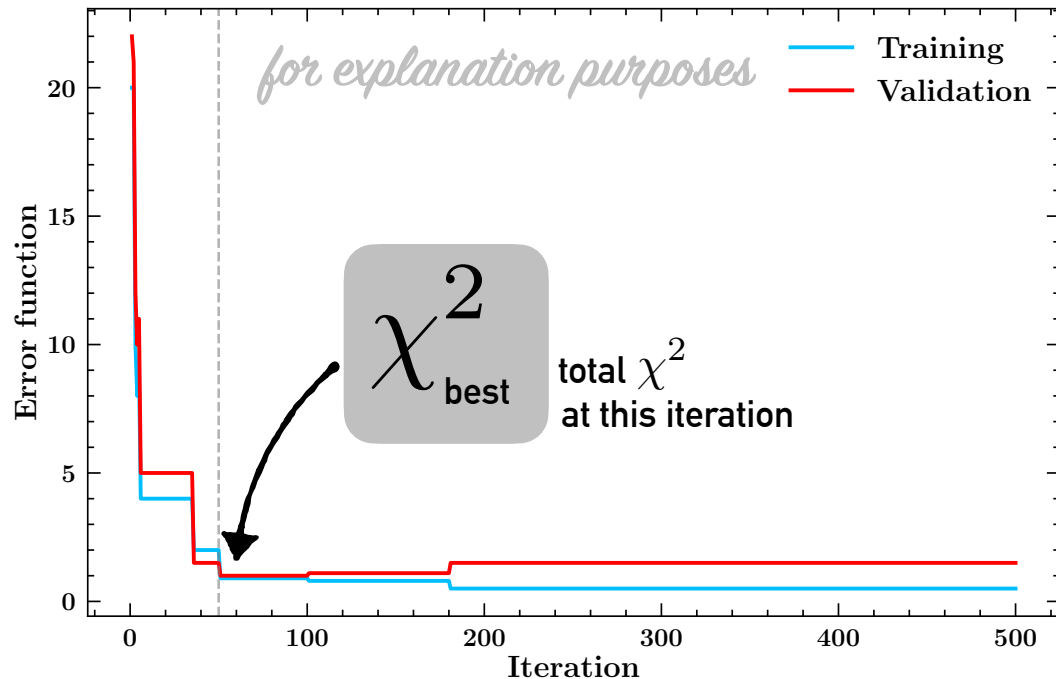
$$g_K(b_T^2) = -g_2^2 \frac{b_T^2}{4}$$

MAP22

12 parameters

avoid overfitting

split data into training and validation set



random splitting

50% - 50%

dataset by dataset

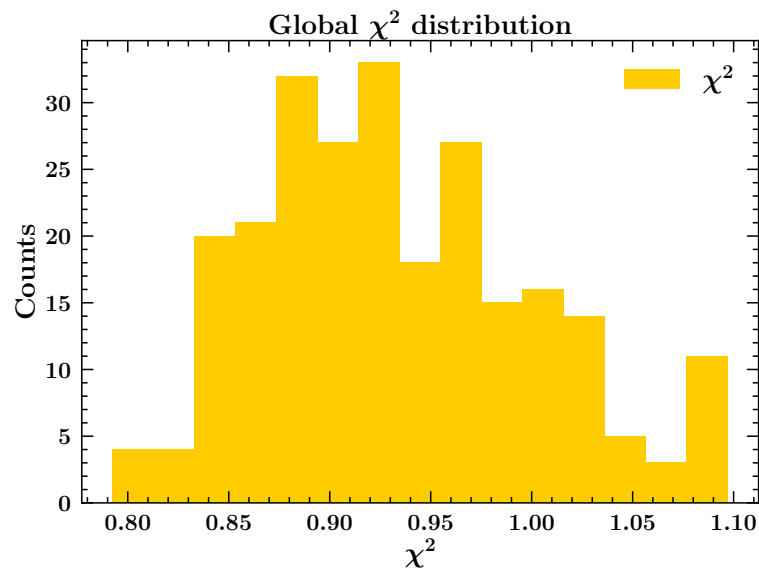
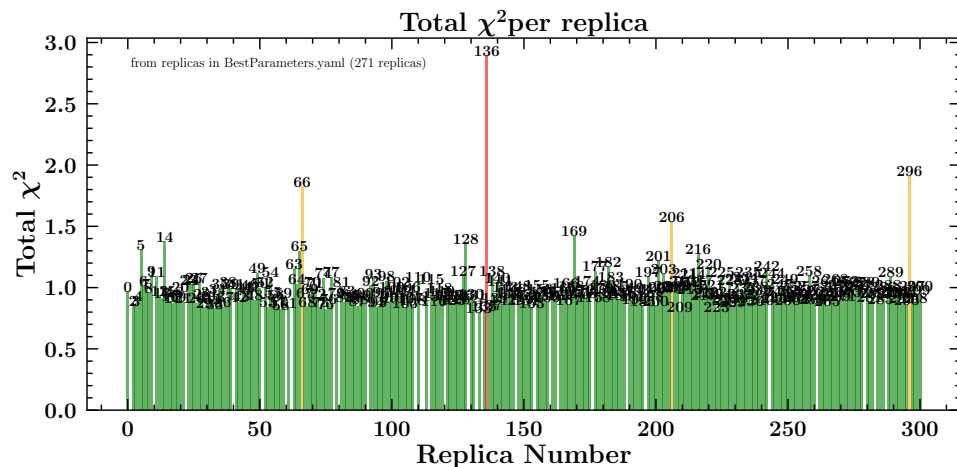
Monte Carlo replicas

each replica is an independent fit

a fit runs in ~15 mins

nearly all replicas get to the end of the fit

$$\chi_0^2 = 0.97$$

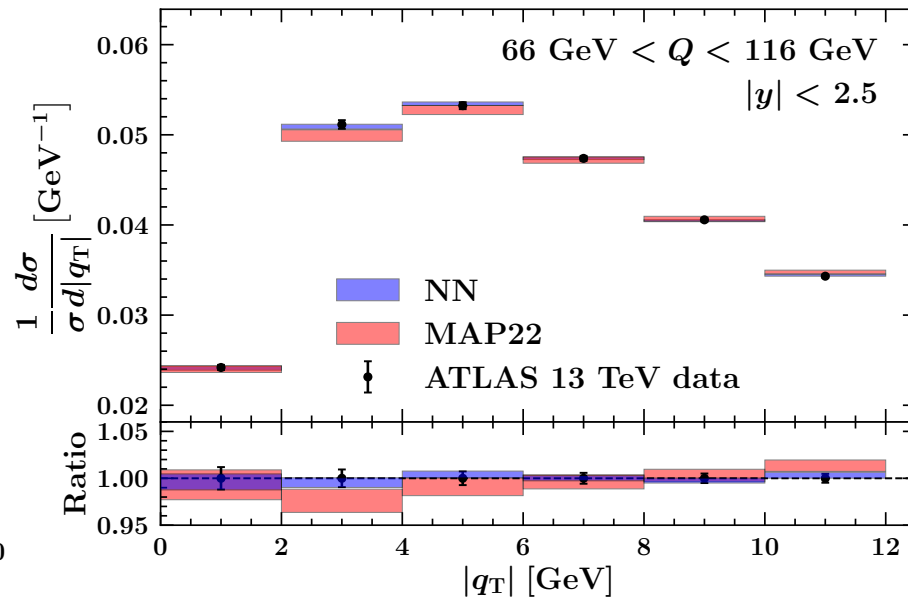
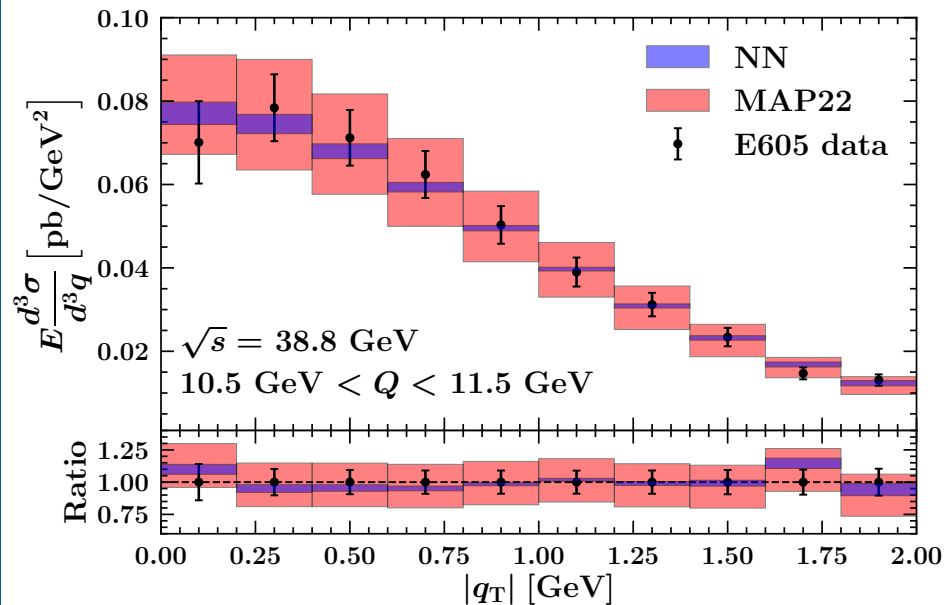


Comparison with data

logarithmic accuracy: N3LL

$$\chi^2 = 0.97$$

great description of data



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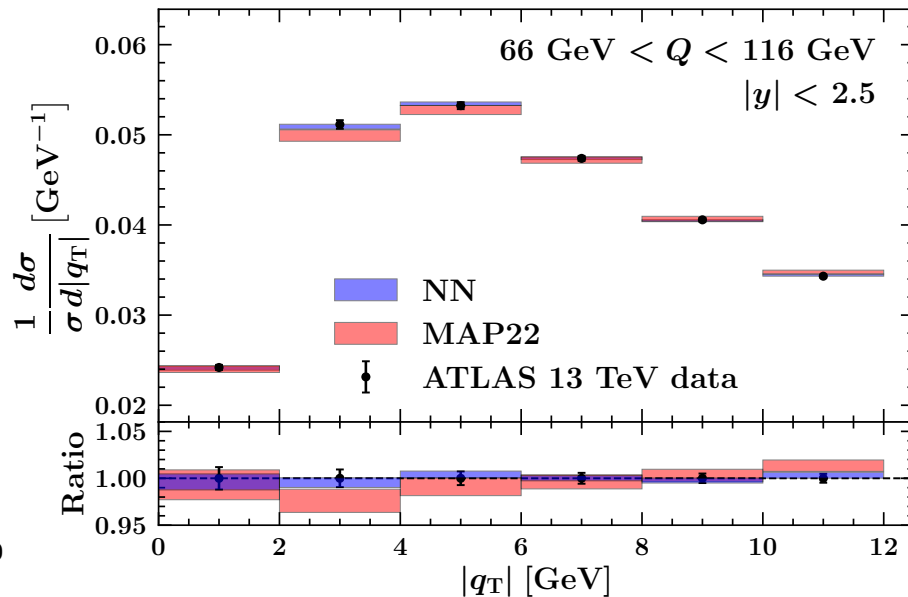
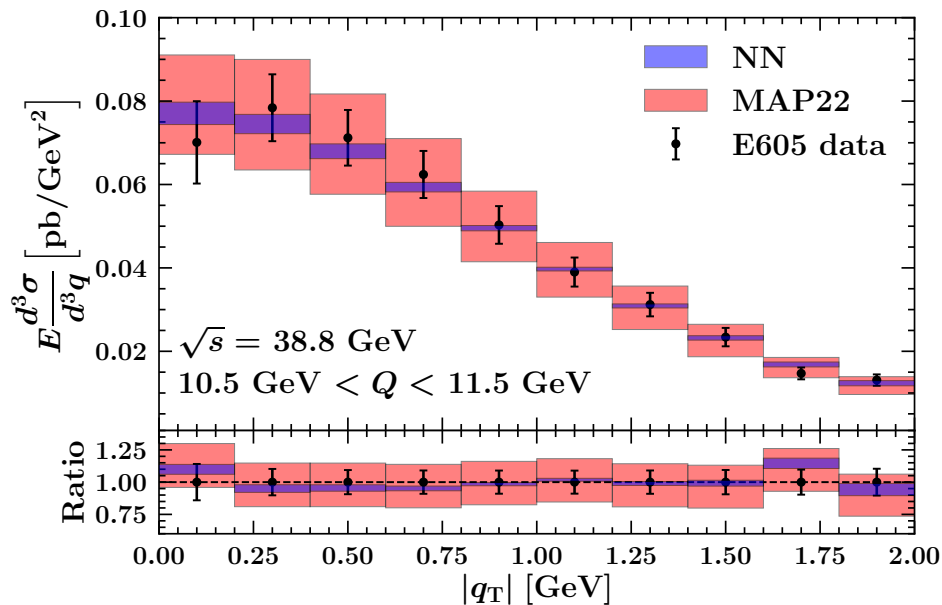
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Comparison with data

logarithmic accuracy: N3LL

$$\chi^2 = 0.97$$

great description of data



why are **NN bands** so much smaller than **MAP22's**?

Shifted predictions

& decomposition of the χ^2

experimental data

$$m_i \pm \sigma_{i,\text{unc}} \pm \sigma_{i,\text{corr}}^{(1)} \pm \dots \pm \sigma_{i,\text{corr}}^{(k)}$$

uncorrelated

correlated

$$\chi^2 = \sum_{i=1}^n \left(\frac{m_i - \bar{t}_i}{s_i} \right)^2 + \sum_{\alpha=1}^k \lambda_{\alpha}^2$$

uncorrelated contribution

penalty term

nuisance parameters $\frac{\partial \chi^2}{\partial \lambda_{\alpha}} = 0$

$$\bar{t}_i = t_i + d_i$$

shifted prediction

shift

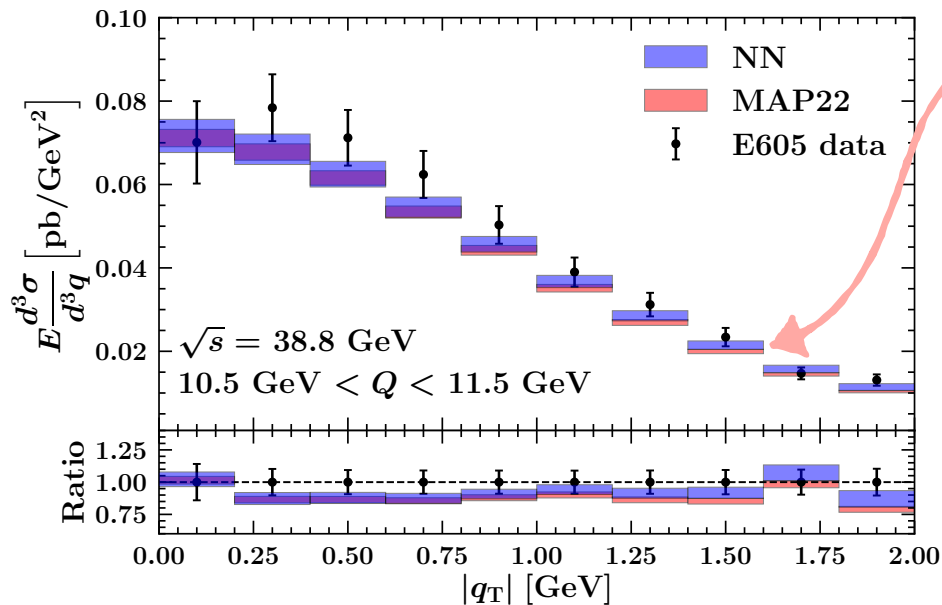
$$d_i = \sum_{\alpha=1}^k \lambda_{\alpha} \sigma_{i,\text{corr}}^{(\alpha)}$$

Comparison with data

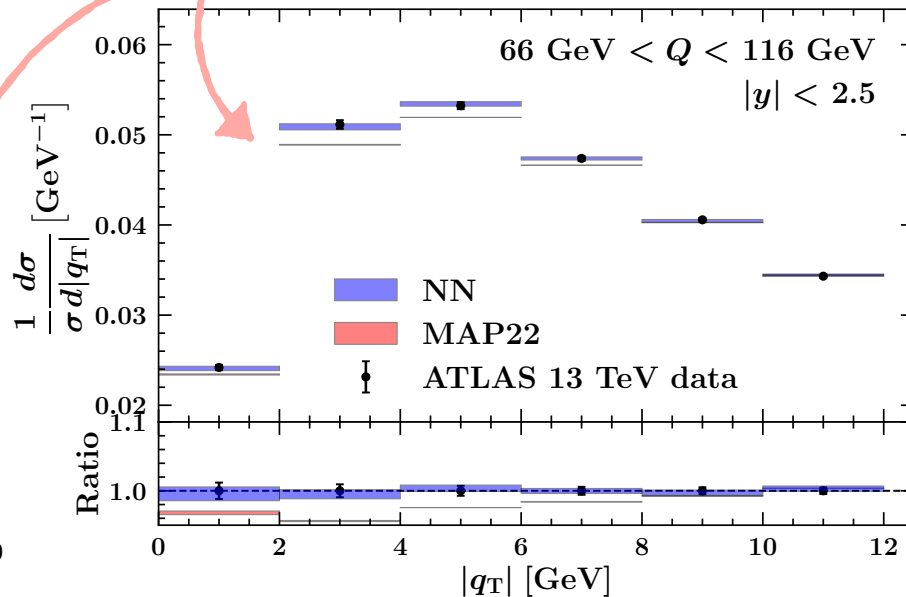
without including systematic shifts

$$\chi^2 = 0.97$$

$$\chi^2 = 1.28$$



significant reduction
of the bands for MAP22



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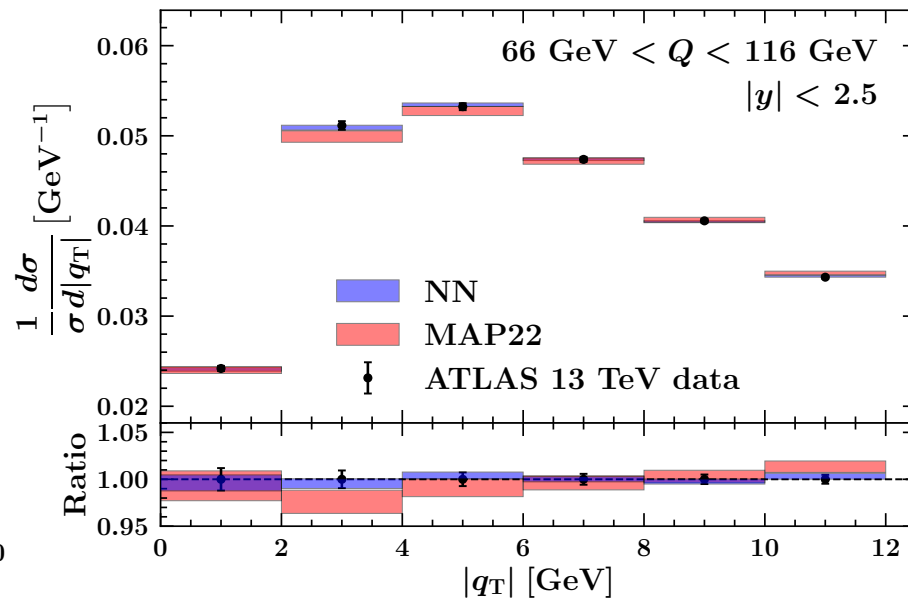
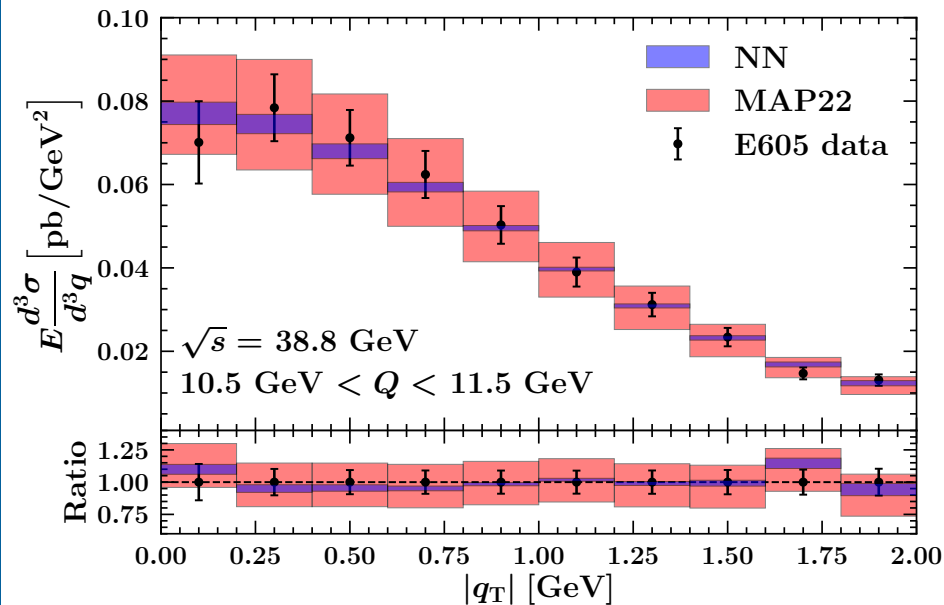
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... just to compare again

N3LL

$$\chi^2 = 0.97$$

$$\chi^2 = 1.28$$



Description of data

numerical breakdown

Experiment	N_{dat}	$\bar{\chi}^2 (\bar{\chi}_D^2 + \bar{\chi}_\lambda^2)$	
		NN	MAP22
Fixed-target	233	1.08 (0.98 + 0.10)	0.91 (0.70 + 0.21)
RHIC	7	1.11 (1.03 + 0.07)	1.45 (1.37 + 0.08)
Tevatron	71	0.80 (0.73 + 0.06)	1.20 (1.17 + 0.04)
LHCb	21	0.98 (0.88 + 0.10)	1.25 (1.05 + 0.20)
CMS	78	0.40 (0.38 + 0.02)	0.41 (0.35 + 0.06)
ATLAS	72	1.38 (1.09 + 0.29)	3.51 (3.03 + 0.49)
Total	482	0.97 (0.86 + 0.11)	1.28 (1.09 + 0.20)

$$\chi^2 = \sum_{i=1}^n \left(\frac{m_i - \bar{t}_i}{s_i} \right)^2 + \sum_{\alpha=1}^k \lambda_\alpha^2$$

uncorrelated
contribution

penalty term



rule of thumb: it's best
when it's small

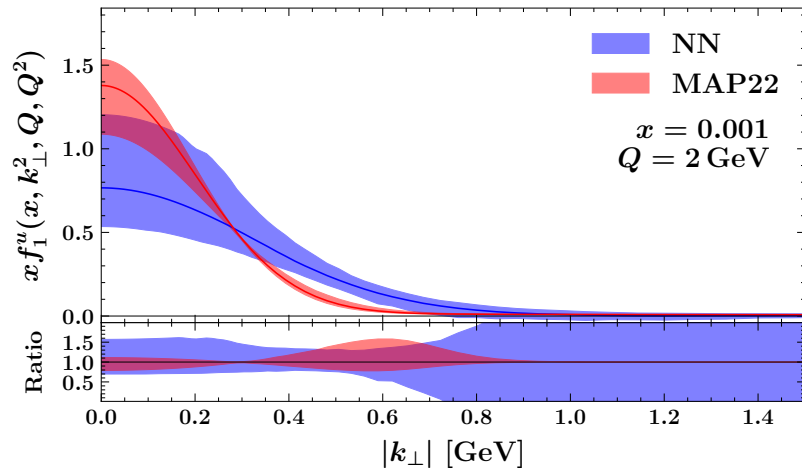
great improvement

MAP22 relies more on the shifts with respect to the NN fit

NN TMDs

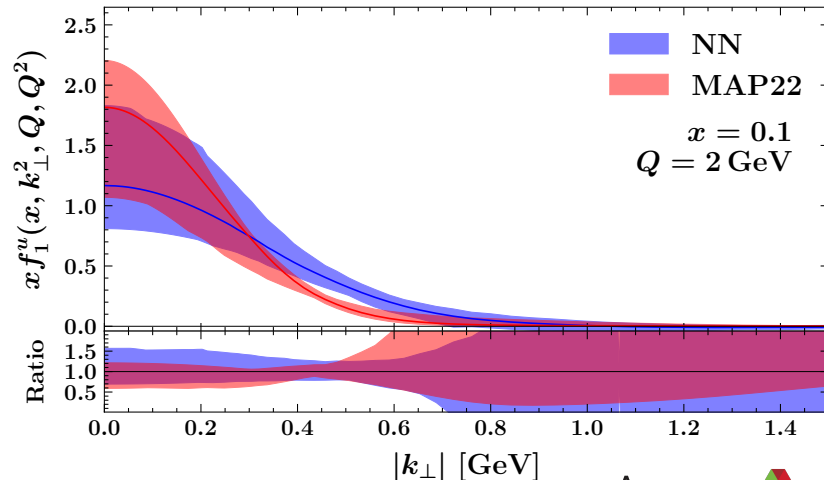
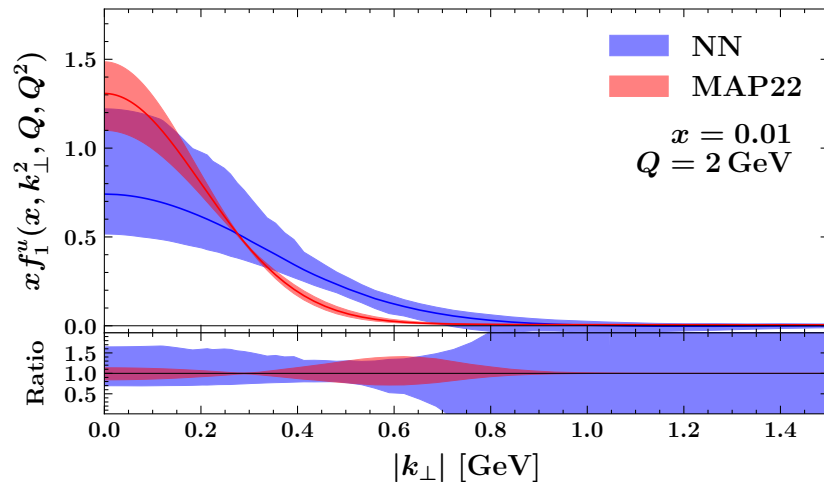
extraction of $f_1(x, k_T)$ from DY data

MAP Collaboration
arXiv:2502.04166



NN fit has larger bands

more reliable estimation uncertainties



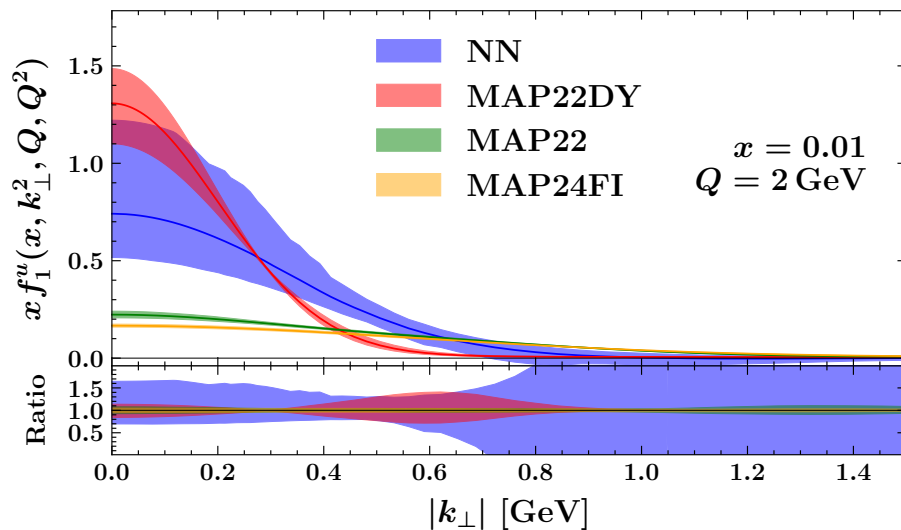
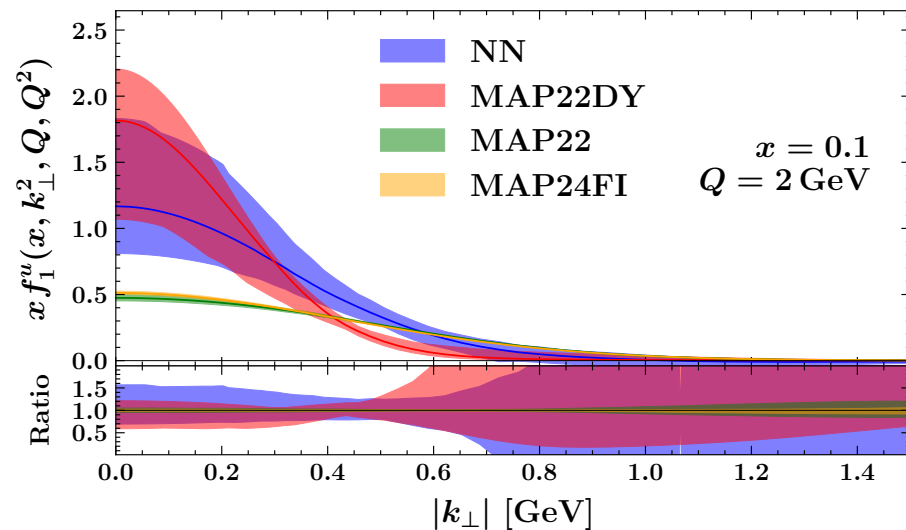
TMDs

comparison with other MAP fits

MAP22 and MAP24

full global fits, including SIDIS data

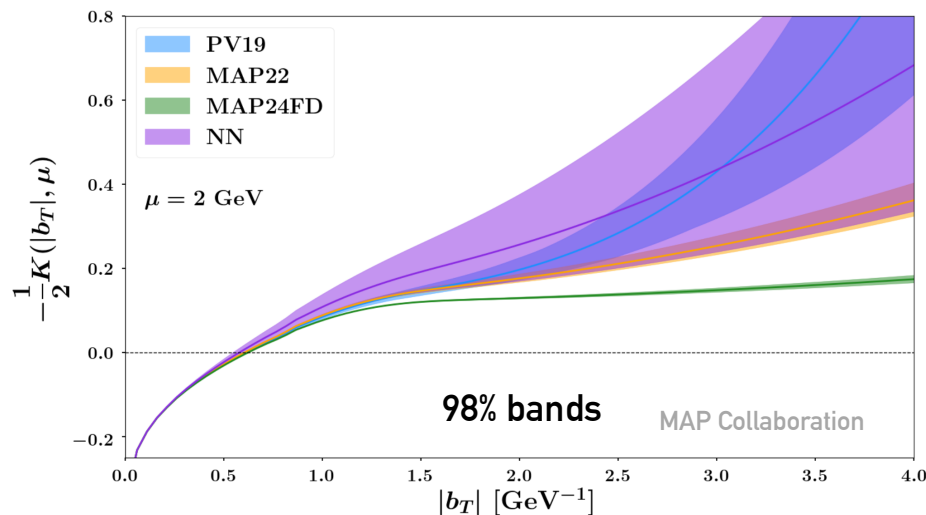
- narrower uncertainty bands
- wider TMDs



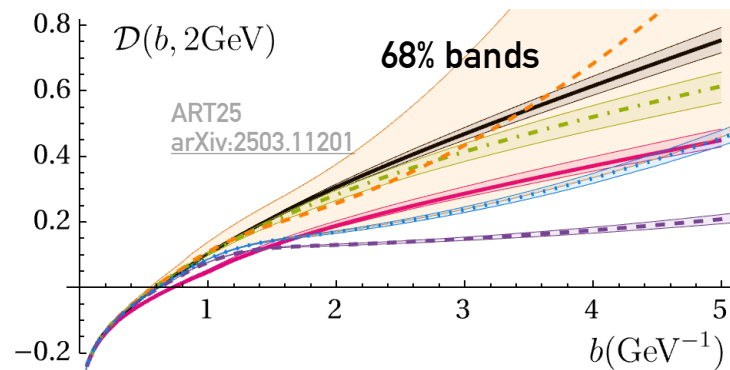
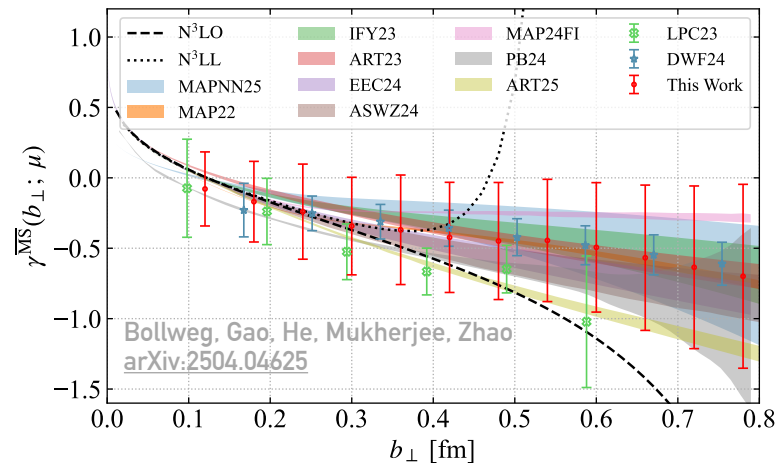
Collins-Soper kernel

“strong” universality

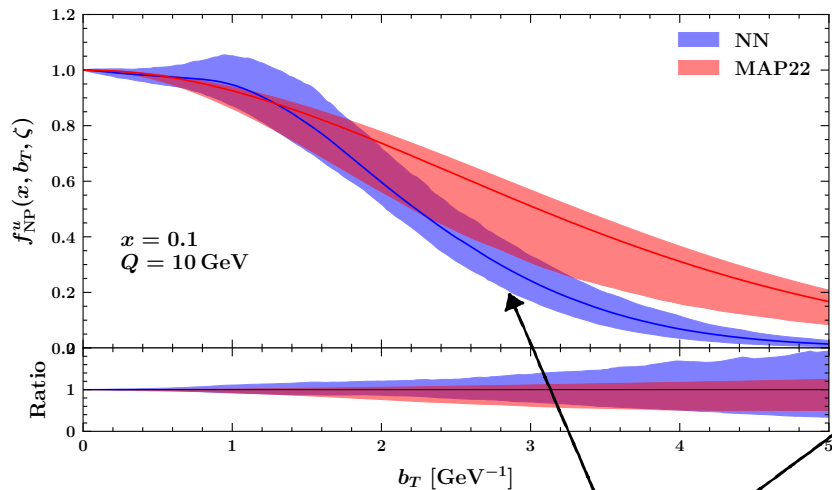
process independent,
insensitive to the types of external hadrons involved,
not dependent on polarization, on the flavors of the quarks, and on the scale Q



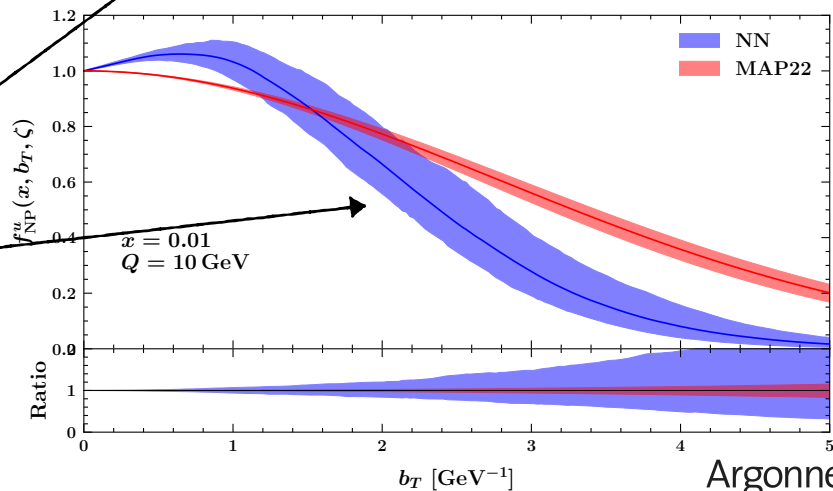
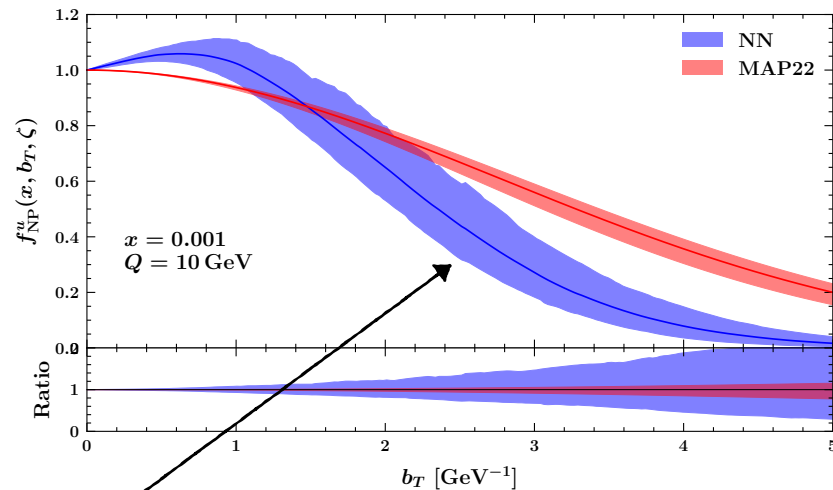
good agreement between the extractions



comparison with MAP22 only DY



NN falls sharply with b_T



Closure tests *only Drell-Yan data*

to validate the methodology

closure test - level 0

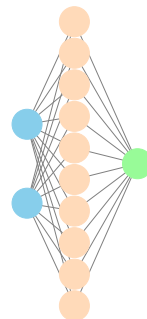
Generate **pseudodata** based
on a model that we know

- * central value: computed with **DWS** parameters
- * uncertainties: real DY data

and for MAP22?

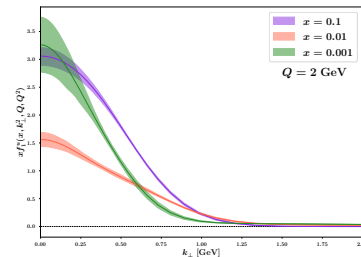
...work in progress

fit with NN



$$g_1 = 0.304$$
$$g_2 = 0.028$$

$$f_{\text{NP}}^{\text{DWS}}(b_T, \zeta) = \exp \left[-\frac{1}{2} \left(g_1 + g_2 \ln \left(\frac{\zeta}{2Q_0^2} \right) \right) b_T^2 \right]$$



can we recover the
DWS result?

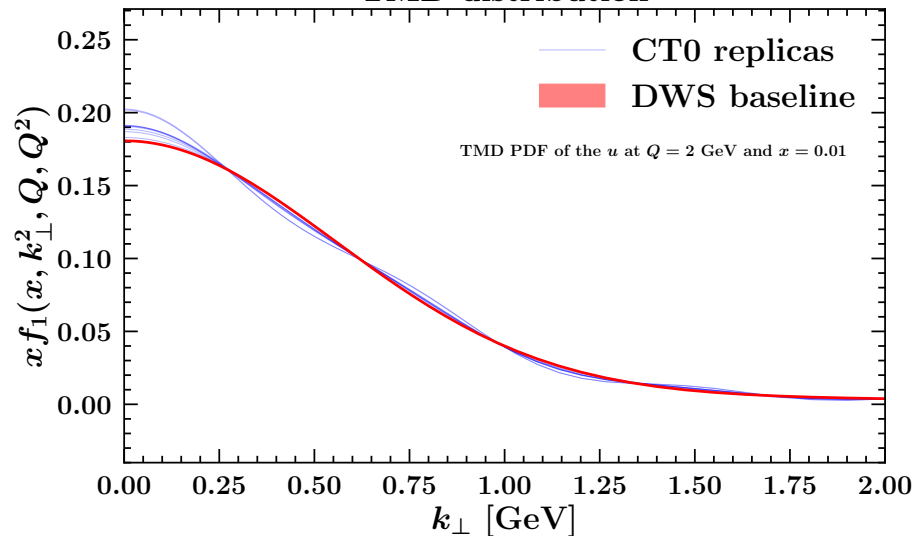
yes!

Closure tests

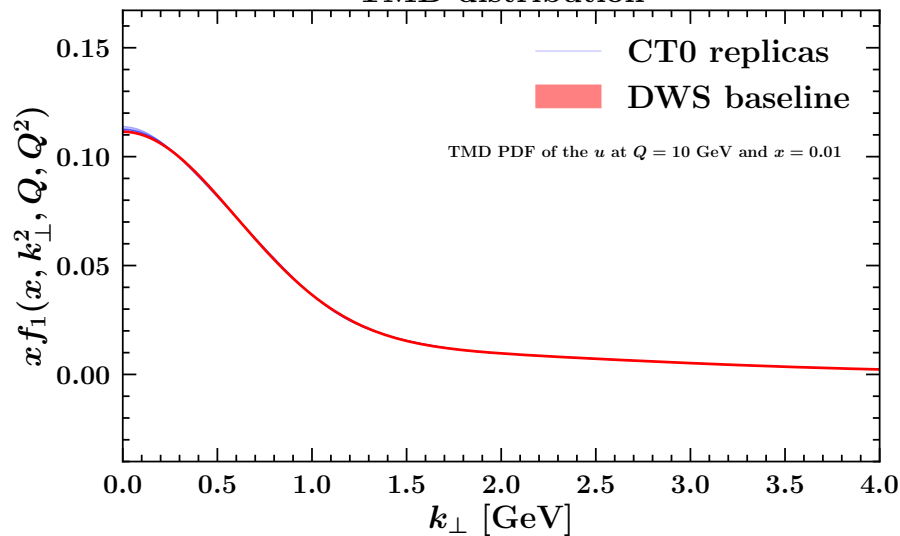
closure test - level 0

$$f_{\text{NP}}^{\text{DWS}}(b_T, \zeta) = \exp \left[-\frac{1}{2} \left(g_1 + g_2 \ln \left(\frac{\zeta}{2Q_0^2} \right) \right) b_T^2 \right]$$

TMD distribution



TMD distribution



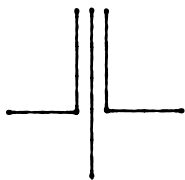
$$\chi^2 \sim 10^{-6}$$

Conclusions

◆ first Neural Networks extraction

of the unpolarized quark TMD PDFs in the proton

at N3LL

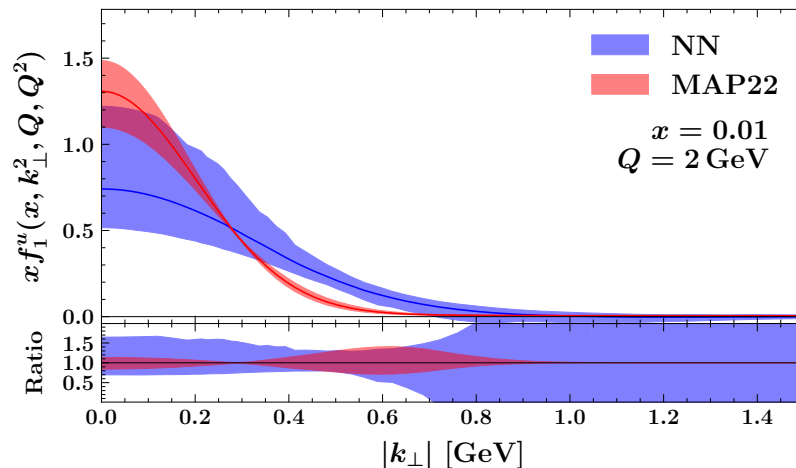


Drell-Yan

data, 482 points

great description of data

$$\chi^2 = 0.97$$



more reliable uncertainties

OROS

the genesis of a new fitting framework

MAP Collaboration



ὄρος

— more flexible, modular and versatile

- suited for fitting all kinds distributions, PDFs, TMDs, GPDs
- very fast
- written in C++ but with a python wrapper

to do next:

- ◆ Study b^* prescription - is there an impact on the extracted TMDs?
- ◆ Perform closure tests - validate the use of NN
- ◆ perform NN TMD fit for DY + SIDIS



TMD

every TMD has the same general structure

many subtleties involved in TMD analyses

b* prescription, ζ -prescription,
logarithmic accuracy

matching to the
collinear region

collinear PDFs

$$f_1^q(x, b; \mu, \zeta) = \sum_j (C_j \otimes f^j)(x, b_*; \mu_b) e^{R(b_*; \mu_b, \mu)} f_{\text{NP}}(x, b)$$

double scale dependence

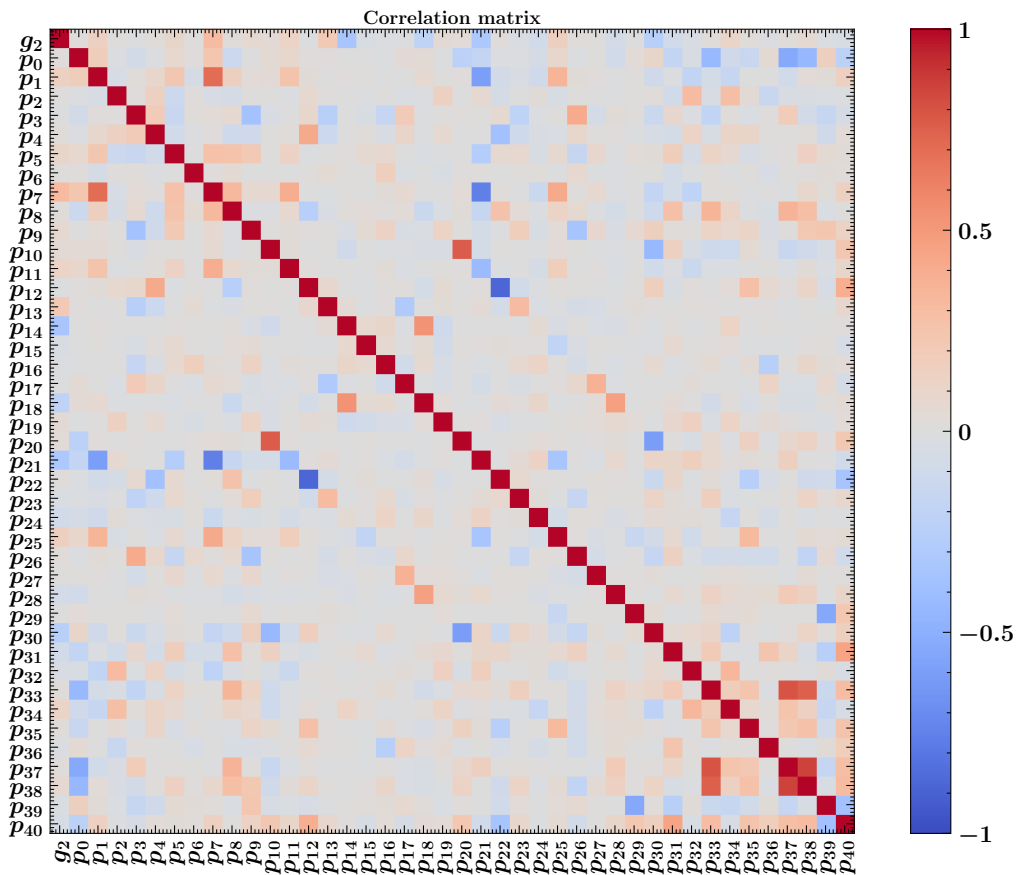
perturbative
evolution

non perturbative
transverse content

parametrized
and fitted to data

correlation matrix

41 parameters



Closure tests

to validate the methodology

* Level 0

- _____ central pseudo-data is given by **central predictions** of the known model
- _____ **no Monte Carlo noise** is added on top of the central data
— each replica is fitting the same set of data

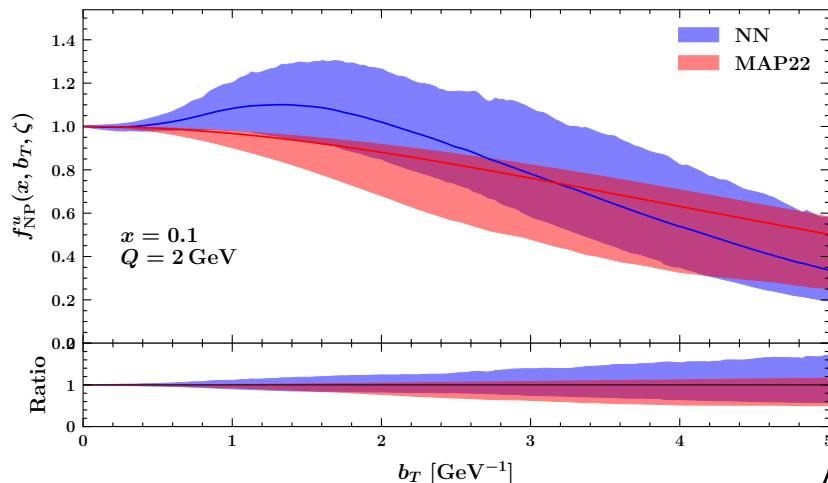
* Level 1

- _____ central pseudo-data is shifted by some noise η
drawn from the experimental covariance matrix
- _____ no MC noise is added — each replica fits a subset of the same shifted data

* Level 2

- _____ central pseudo-data is shifted by level 1 noise η
- _____ MC noise is added on top of the level 1 shift

comparison with MAP22 only DY



NN has a bump at very small b_T
it reaches 1 from above

