Lattice QCD Calculation of the Pion Distribution Amplitude with Domain Wall Fermions at Physical Pion Mass

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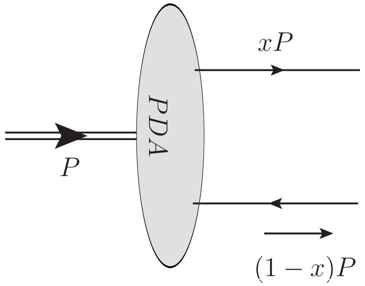
Outline

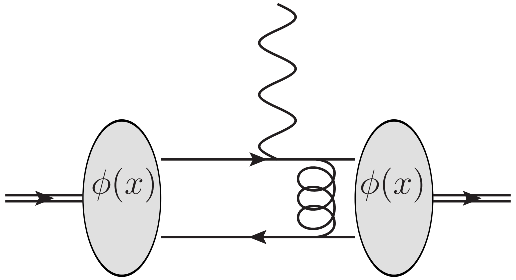
- Introduction to pion distribution amplitude
- Lattice calculation of pion DA
- Resummation in quasi-DA matching
- Conclusion and Outlook

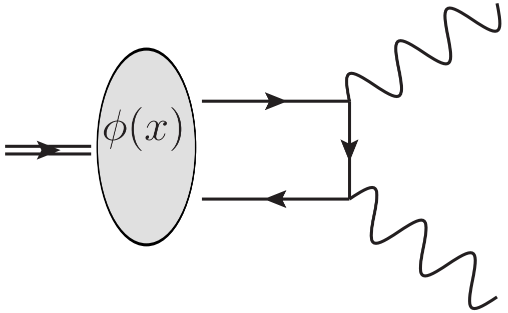
Outline

Introduction to pion distribution amplitude

Pion Distribution Amplitude

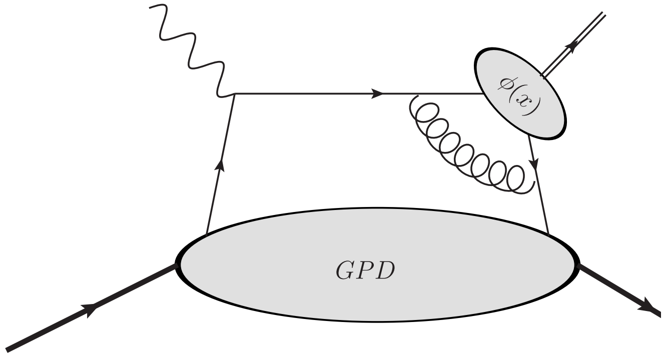


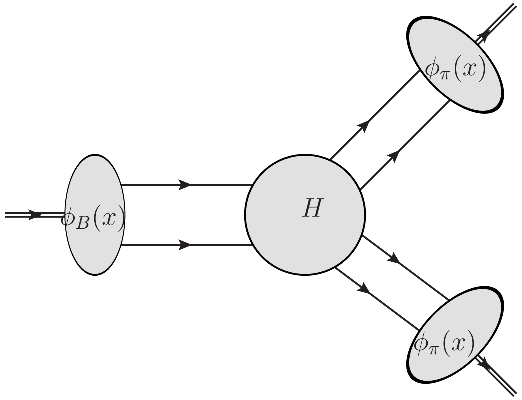


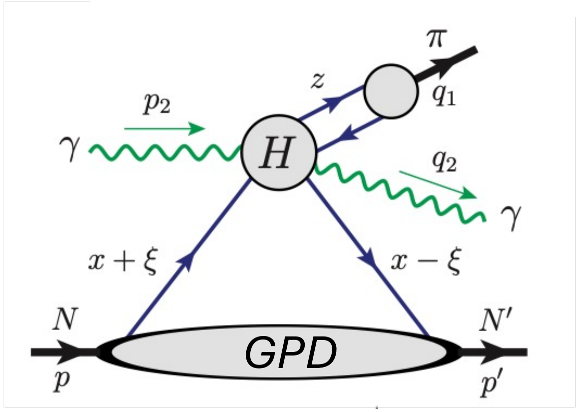


Universal inputs to various hard exclusive processes at large momentum transfer Q^2

- $\pi \to \gamma\gamma^*$ transition form factor
- Pion electromagnetic form factor
- Deeply virtual meson production Brodsky, et.al, PRD (1994)
- Heavy meson decay Beneke, et.al, PRL (1999)
- Exclusive Photoproduction Z.Yu & J.Qiu, PRL (2024)
- ...







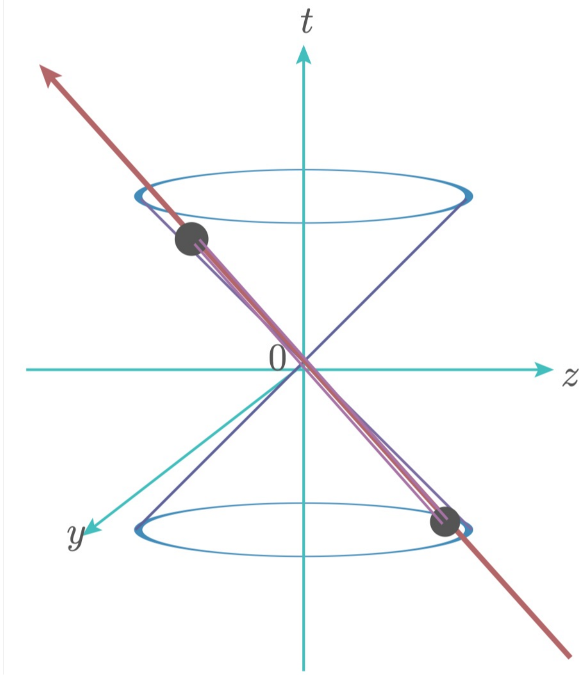
Weakly constrained by experiments!

Direct calculation from lattice QCD?

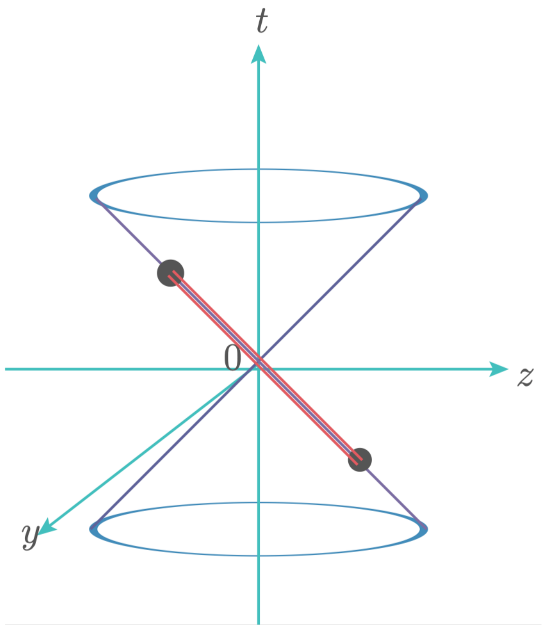
Large Momentum Effective Theory (LaMET)

Ji, PRL (2013)
Ji, SCPMA(2014)

Momentum distribution



Large P_z Expansion



$$+\mathcal{O}\left(\frac{1}{P_z^n}\right)$$

Quasi-DA: $\tilde{\phi}(x, P_z) =$

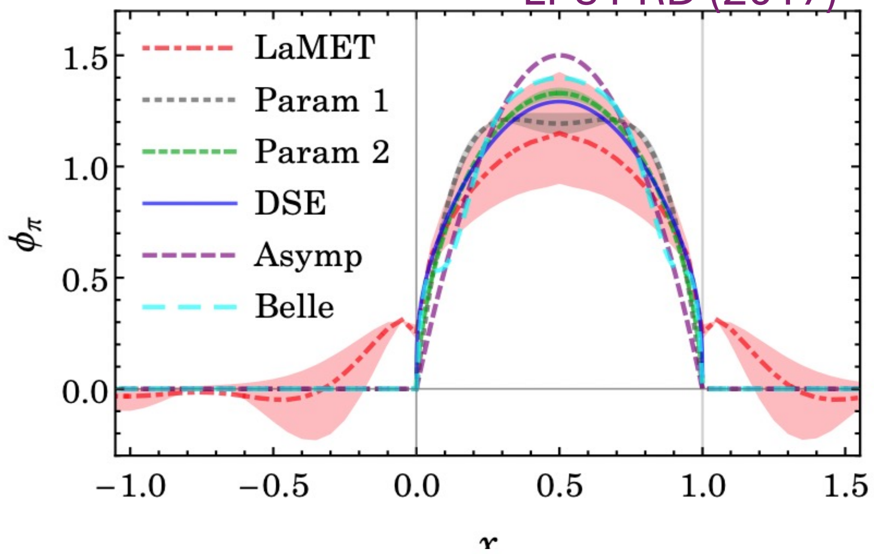
$$\int \frac{dz}{2\pi} e^{i\left(\frac{1}{2}-x\right)zP_z} \langle 0|\bar{q}\left(-\frac{z}{2}\right)\gamma_z\gamma_5 U(0,z) q\left(\frac{z}{2}\right)|\pi\rangle$$

$$C\,(x, y, \mu, P_z) \otimes \phi(y, \mu) \qquad +\mathcal{O}\left(\frac{\Lambda_{QCD}^2}{x^2 P_z^2}, \frac{\Lambda_{QCD}^2}{\bar{x}^2 P_z^2}\right)$$

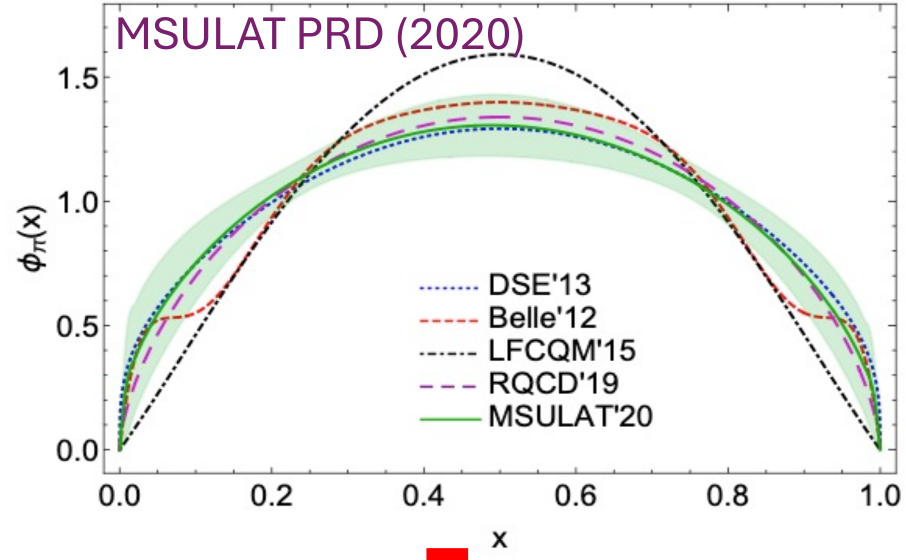
$$= \frac{1}{if_\pi} \int \frac{d\eta^-}{2\pi} e^{i\left(\frac{1}{2}-x\right)\eta^- p^+} \langle 0|\bar{q}\left(\frac{\eta^-}{2}\right)\gamma_+\gamma_5 U\left(\frac{\eta^-}{2}, -\frac{\eta^-}{2}\right) q\left(-\frac{\eta^-}{2}\right)|\pi(p)\rangle$$

Progress in x-dependent DA calculations

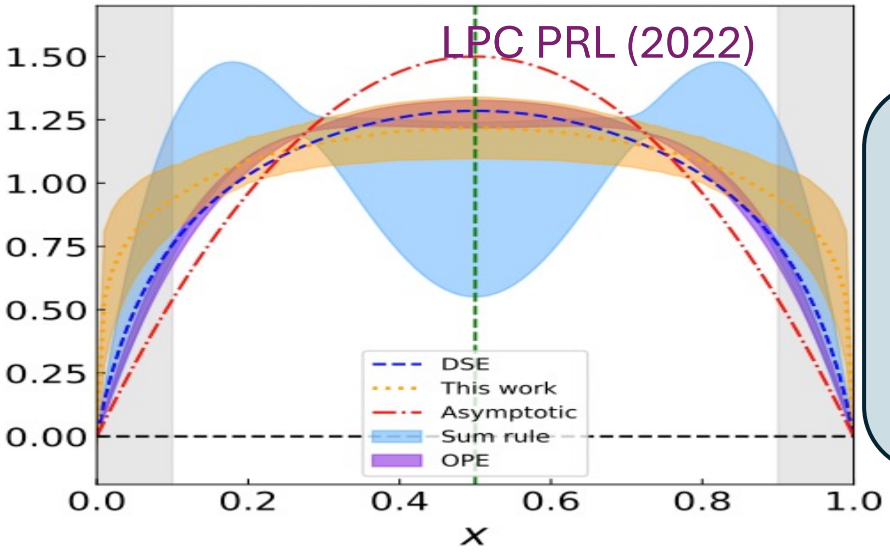
LP3 PRD (2017)



$a \to 0$

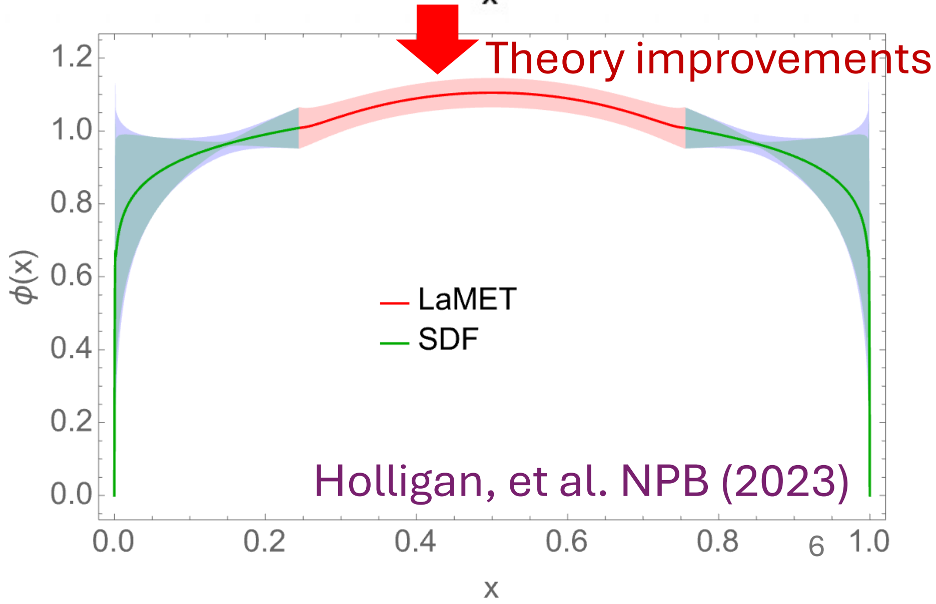


$m_\pi \to 130\ MeV$



This work:
Chiral Symmetry

Large Logarithm
Resummation



Outline

Lattice calculation of pion DA

Lattice Setup

- Physical pion mass
- Chiral symmetric Fermion action – domain wall fermions
- Momentum smeared quark source

Lattice Spacing-a	Pion Mass	Lattice Volume	$m_\pi L$	Fermion Action
0.0836 fm	137 MeV	$64^3 \times 128 \times 12$	3.73	2+1f DW

Momentum Smearing	Pion Momentum	Samples	Sources	Effective Statistics
$k = \{0, 1.4\}$ GeV	$P_z = [0, 1.85]$ GeV	55	{32, 128}	Up to 28,160

Lattice raw data and fitting

$$C_{\pi\pi}(t) = \langle O_\pi(0)|O_\pi(t)\rangle,$$

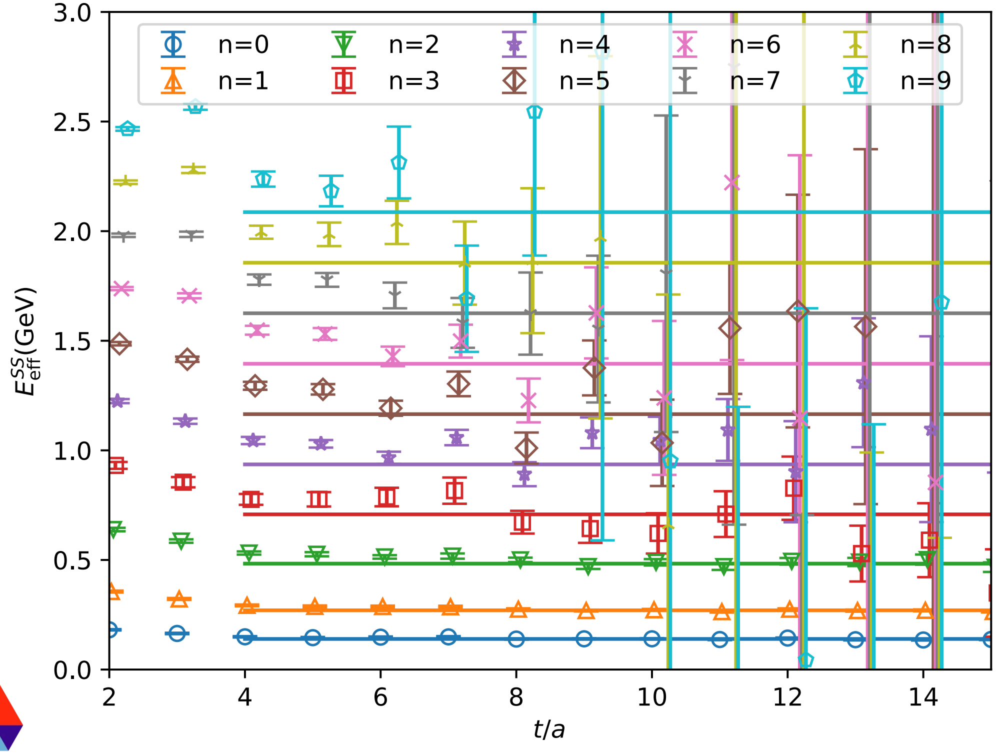
$$C_{\pi O_0}(t,z) = \langle O_\pi(0)|\bar{\psi}(-\frac{z}{2},t)\gamma_t\gamma_5 W(-\frac{z}{2},\frac{z}{2})\psi(\frac{z}{2},t)|\Omega\rangle,$$

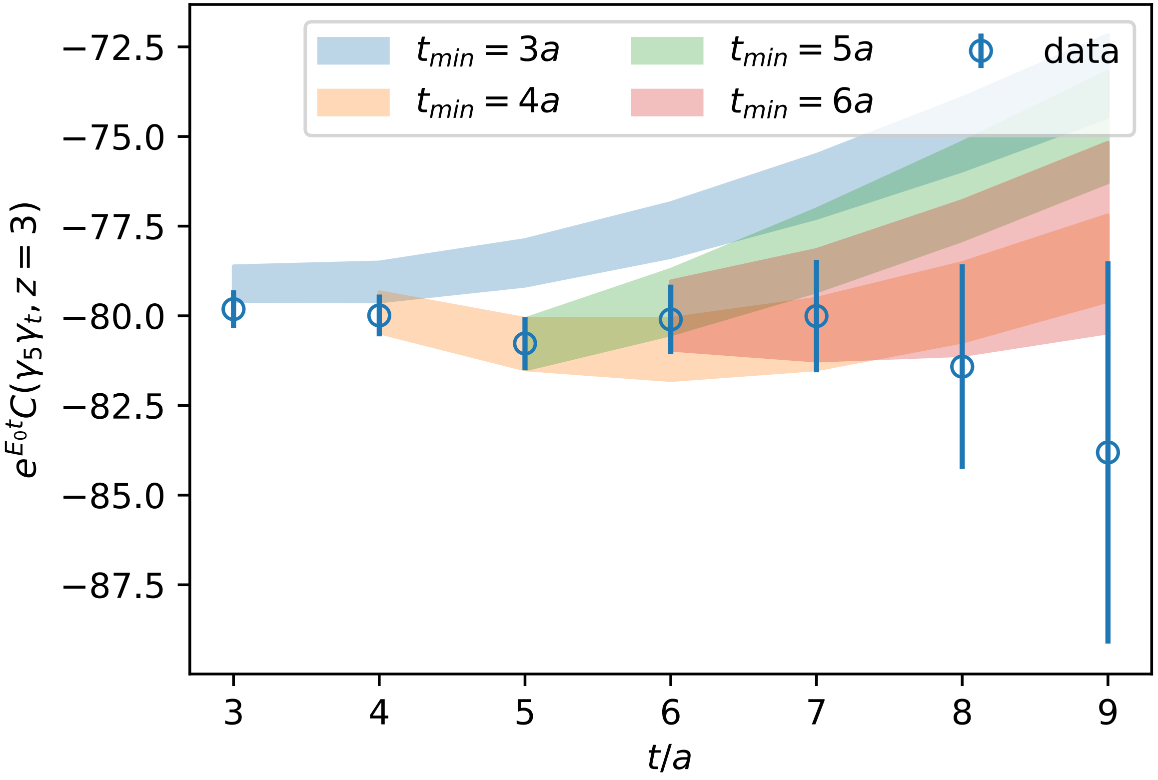
$$C_{\pi O_3}(t,z) = \langle O_\pi(0)|\bar{\psi}(-\frac{z}{2},t)\gamma_z\gamma_5 W(-\frac{z}{2},\frac{z}{2})\psi(\frac{z}{2},t)|\Omega\rangle$$

$$C_{\pi\pi}(t) = \sum A_i^\pi(e^{-E_i t} + e^{-E_i(N_t - t)}),$$

$$C_{\pi O_0}(t,z) = \sum A_i^{O_0}(z)(e^{-E_i t} + e^{-E_i(N_t - t)}),$$

$$C_{\pi O_3}(t,z) = \sum A_i^{O_3}(z)(e^{-E_i t} + e^{-E_i(N_t - t)}),$$





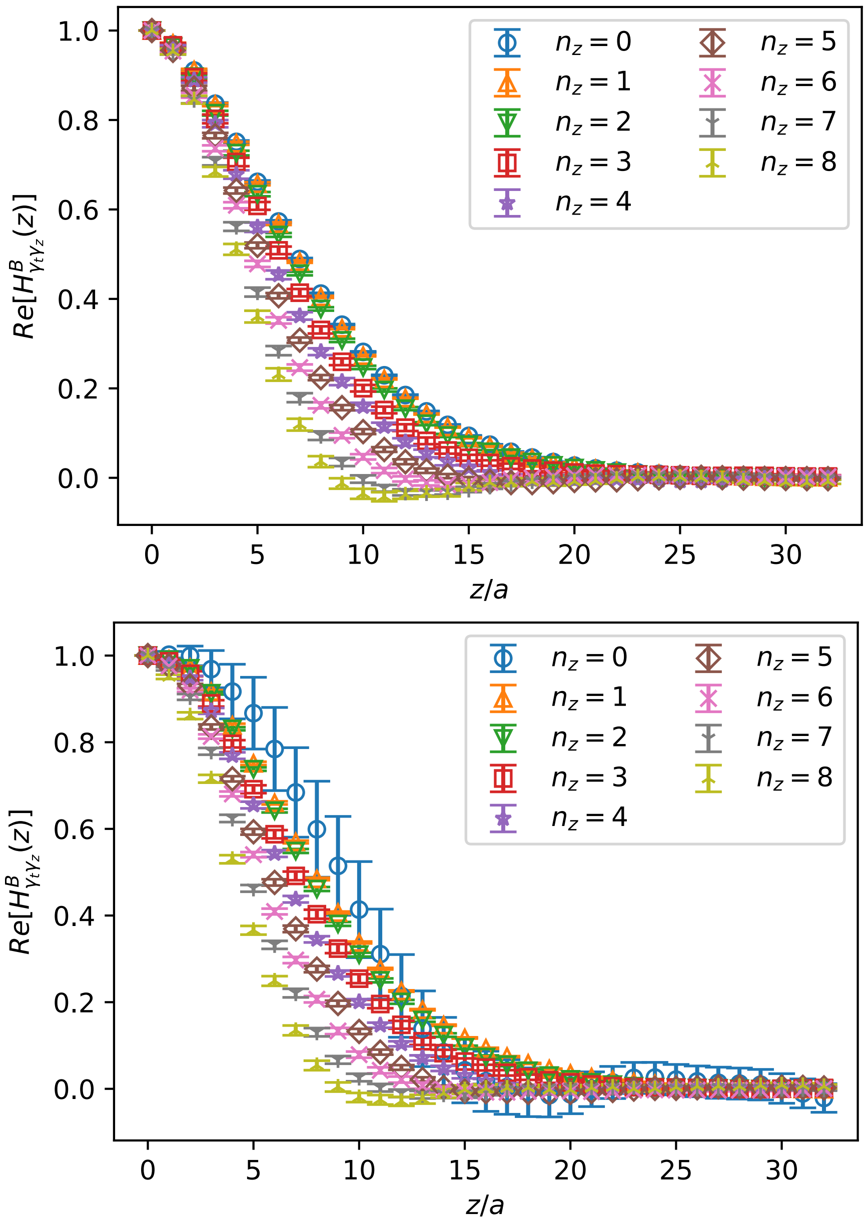
Bare matrix elements

$$A_0^{\pi} = \frac{|\langle O_\pi|\pi\rangle|^2}{2E_0},$$
$$A_0^{O_0}(z) = \frac{\langle O_\pi|\pi\rangle}{2E_0} f_\pi H_{\gamma_t\gamma_5}(z) E_0,$$
$$A_0^{O_3}(z) = \frac{\langle O_\pi|\pi\rangle}{2E_0} i f_\pi H_{\gamma_z\gamma_5}(z) P_z,$$

Zero OP mixing w/ DWF

Pion DA is symmetric (vanishing imaginary part)

The lattice data decays exponentially with the Wilson link length.
The bare results contains both logarithmic and linear divergence in lattice spacing a



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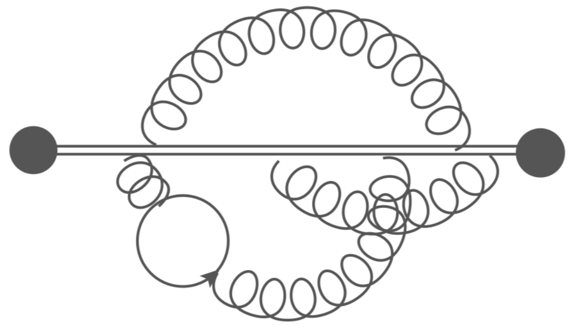
Renormalizing linear divergence



- Linearly divergence in Wilson line: $U(0, z)$
  - $h^B(\mathrm{z}) \sim e^{-\delta m(a)\cdot z}$ Ji, et.al, PRL (2017)
- Renormalon ambiguity in $\Delta\big(\delta m(a)\big) \sim \Lambda_{QCD}$ Beneke, PLB (1995)
  - Renormalon also in the matching kernel Braun, et al., PRD (2018)

- $h^R(\mathrm{z}) \sim h^B(\mathrm{z}) e^{\delta m\cdot z}$ uncertain up to $e^{\mathcal{O}(\mathrm{z}\Lambda_{\mathrm{QCD}})} \to \mathcal{O}\left(\frac{\Lambda_{\mathrm{QCD}}}{\mathrm{x}\, P_z}\right)$ in $\tilde{q}$

Achieving power accuracy: Zhang, et al., PLB (2023)

- Extracting δm with Leading Renormalon Resummation
- Using LRR-improved matching

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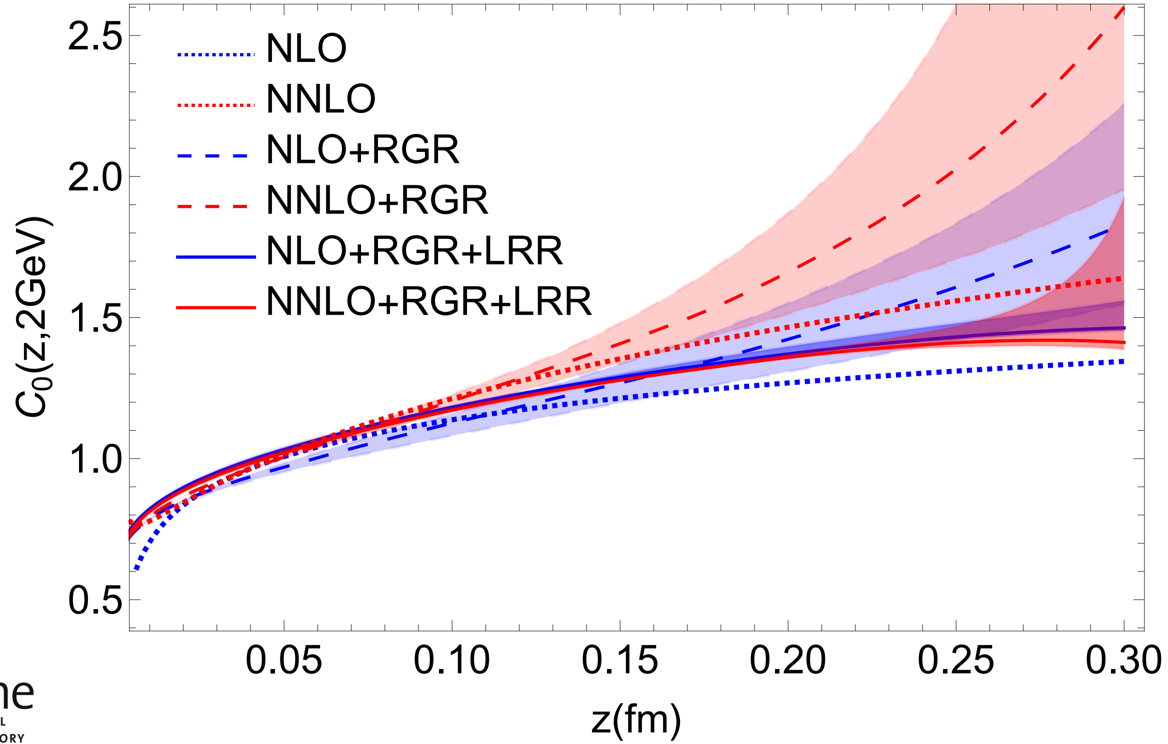
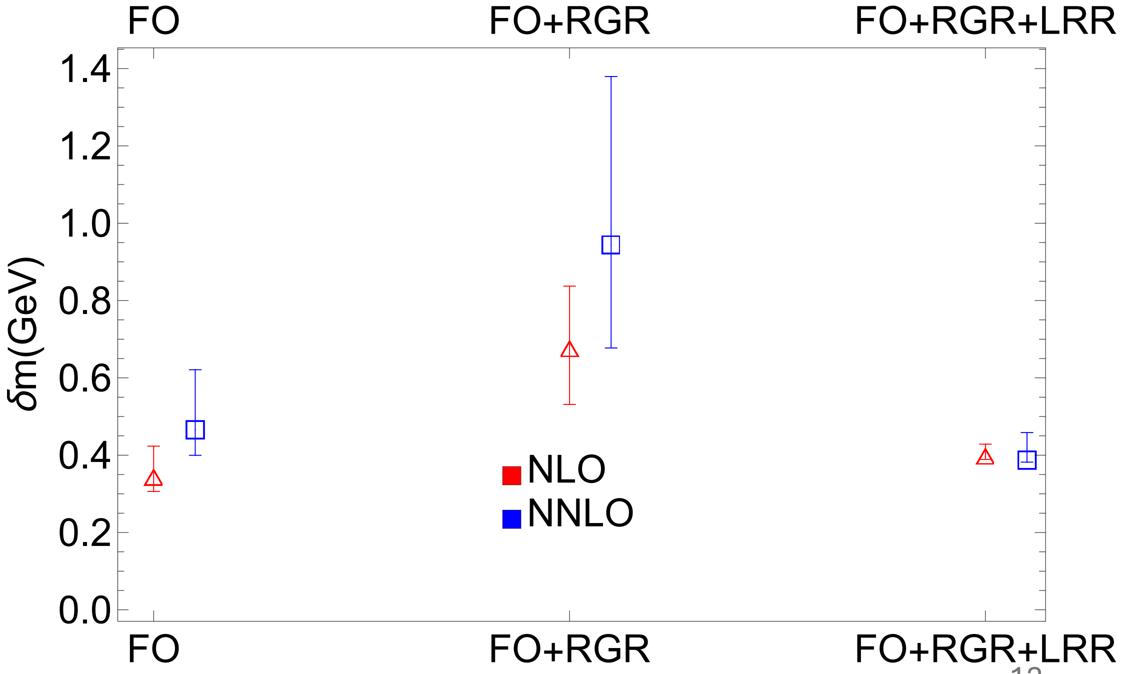
δm extraction with LRR

Zhang, et al., PLB (2023)

$$\ln\left(\frac{C_0(z, z^{-1})\exp(-I(z))}{h^B(z,0)}\right) = \delta m\,|z| + b$$





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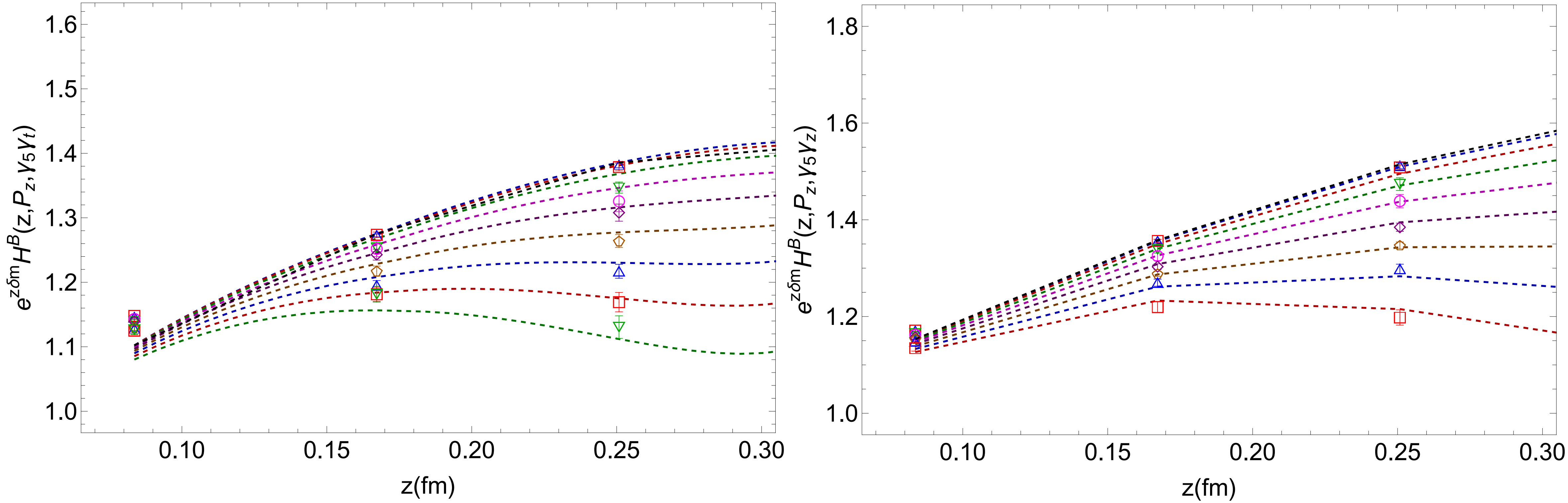
Consistency with OPE

$$H^R(z, P_z, \mu) = \sum_{m=0}^{\infty}\sum_{n=0}^{m}\frac{1}{m!}\left(\frac{i\,zP_z}{2}\right)^m C_{mn}(z,\mu)\langle\xi^n\rangle(\mu)$$

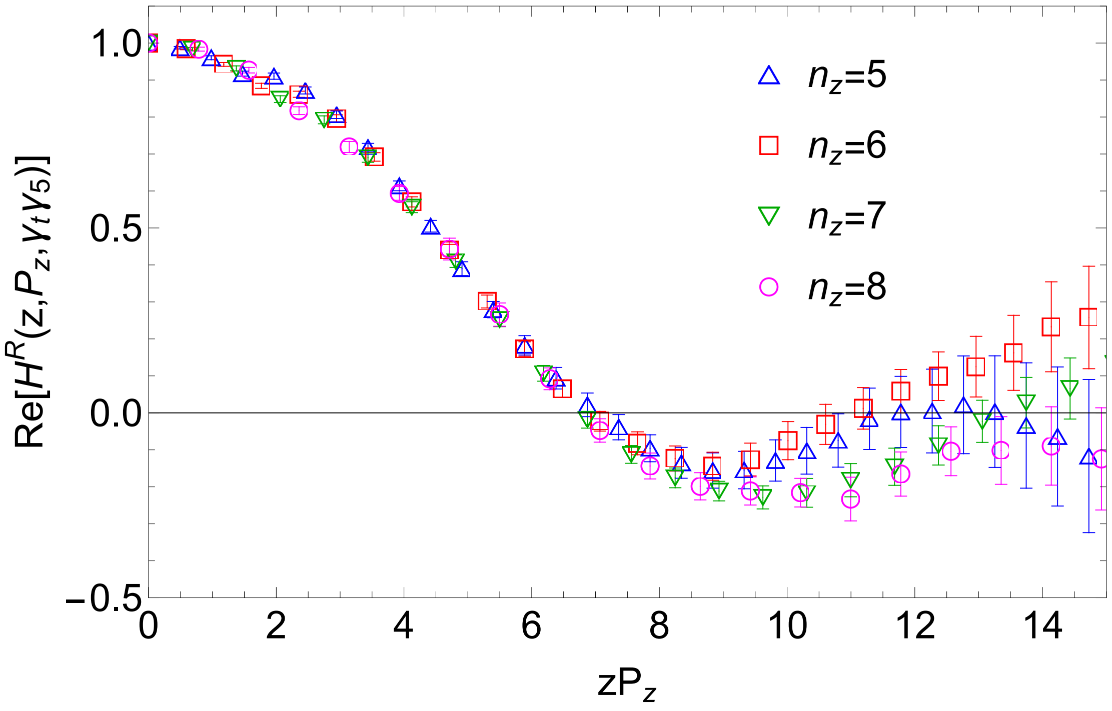


Renormalization in hybrid scheme

Ji, et al., NPB (2020)

$$h^R(z, P_z) = \frac{h^B(z,P_z)}{Z_h(z)}$$

$$Z_h(z) = \begin{cases} h^B(z,0), & |z| < z_s \\ e^{\delta m|z-z_s|} h^B(z_s,0), & |z| > z_s \end{cases}$$



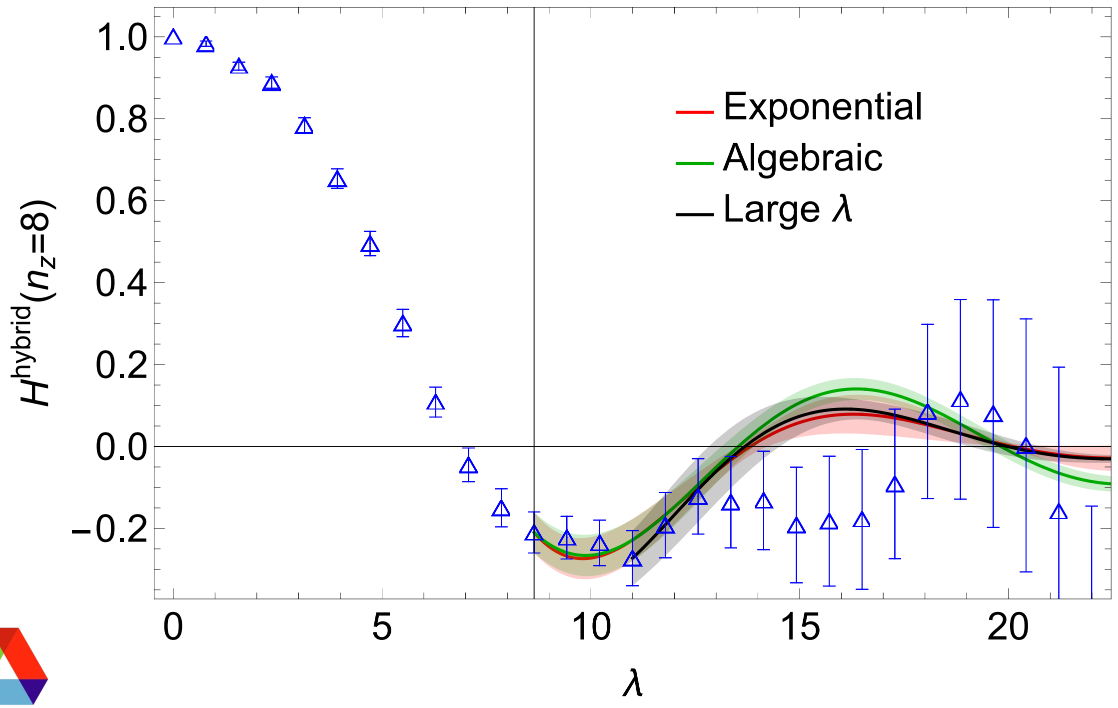
Longtail extrapolation ($\lambda = zP_z \to \infty$)

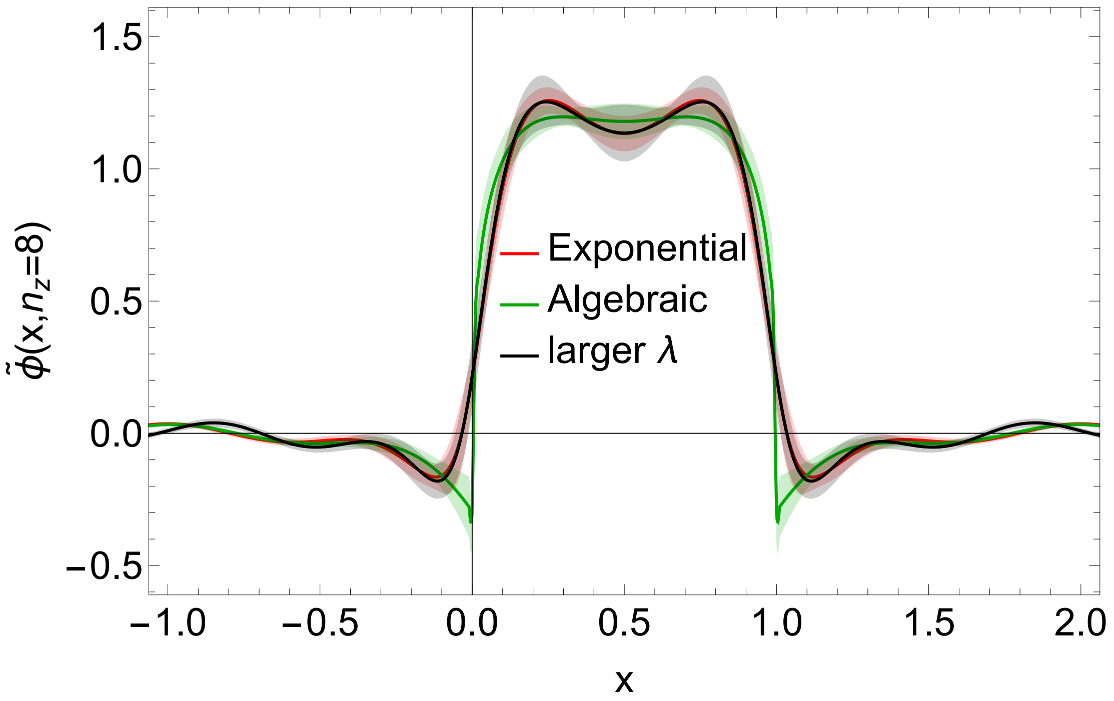
Ji, et al., NPB (2020)

Quasi-DA matrix elements have finite correlation length:

Inferred from Regge behavior

$$h^R(\lambda \to \infty) = e^{-\frac{\lambda}{\lambda_0}}\left(e^{-\frac{i\lambda}{2}}\frac{c_1}{(-i\lambda)^{d_1}} + e^{\frac{i\lambda}{2}}\frac{c_1}{(i\lambda)^{d_1}}\right)$$





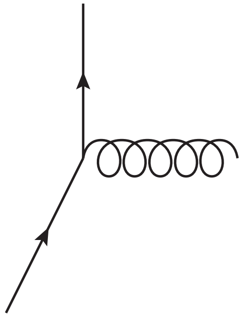
Outline

Resummation in quasi-DA matching

Logarithms in the Matching Kernel

$$C^{\gamma_t\gamma_5}(x,y,\mu,P_z) = \delta(x-y) + \frac{\alpha_s(\mu)C_F}{2\pi}\left[\begin{cases} \frac{1+x-y}{y-x}\frac{\bar{x}}{\bar{y}}\ln\frac{(y-x)}{\bar{x}} + \frac{1+y-x}{y-x}\frac{x}{y}\ln\frac{(y-x)}{-x} & x<0 \\ \frac{1+y-x}{y-x}\frac{x}{y}\ln\frac{4x(y-x)P_z^2}{\mu^2} + \frac{1+x-y}{y-x}\left(\frac{\bar{x}}{\bar{y}}\ln\frac{y-x}{\bar{x}} - \frac{x}{y}\right) & 0<x<y<1 \\ \frac{1+x-y}{x-y}\frac{\bar{x}}{\bar{y}}\ln\frac{4\bar{x}(x-y)P_z^2}{\mu^2} + \frac{1+y-x}{x-y}\left(\frac{x}{y}\ln\frac{x-y}{x} - \frac{\bar{x}}{\bar{y}}\right) & 0<y<x<1 \\ \frac{1+y-x}{x-y}\frac{x}{y}\ln\frac{(x-y)}{x} + \frac{1+x-y}{x-y}\frac{\bar{x}}{\bar{y}}\ln\frac{(x-y)}{-\bar{x}} & 1<x \end{cases}\right.$$

- Efremov-Radyushkin-Brodsky-Lepage logarithm
  - Physical scale of the system
    - Quark momentum logarithm $L = \ln x$
    - Anti-quark momentum logarithm $L = \ln \bar{x}$
- Threshold logarithm
  - Gluon momentum $L = \ln|x-y|$
- Only one RG equation (ERBL evolution): How to resum?



Both become important only in the threshold limit $x \to y$

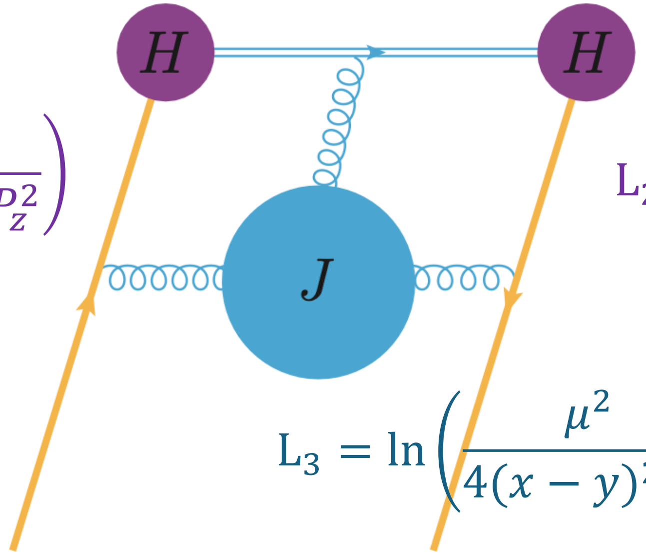
Factorizing Hard and "Soft" scales

Becher, Neubert & Pecjak JHEP(2007)

- All three logarithms are important only in the threshold limit
  - $x - y \to 0$, soft gluon emission
- Integrate out hard modes
  - Sudakov factor H
    - Quark component
    - Anti-quark component
- Integrate out hard collinear modes
  - Jet function J

$$\mathcal{C}(x, y, \mu, P_z) \xrightarrow{x \to y} H(xP_z, \bar{x}P_z, \mu) \otimes J(|x-y|P_z, \mu)$$

$$\mathrm{L}_1 = \ln\left(\frac{\mu^2}{4x^2 P_z^2}\right)$$



$$\mathrm{L}_2 = \ln\left(\frac{\mu^2}{4\bar{x}^2 P_z^2}\right)$$

$$\mathrm{L}_3 = \ln\left(\frac{\mu^2}{4(x-y)^2 P_z^2}\right)$$

Threshold log: soft gluon

Ji, Liu & Su JHEP (2023)

Resumming individual pieces

- $\frac{\partial \ln H^{\pm}}{\partial \ln \mu} = \frac{1}{2}\Gamma_{cusp} \ln\left(\frac{4x^2P^2}{\mu^2}\right) + \gamma_H(\alpha) \pm i\pi\Gamma_{cusp}$
- $\frac{\partial \ln J}{\partial \ln \mu} = \Gamma_{cusp} \ln\left(\frac{z^2\mu^2 e^{2\gamma_E}}{4}\right) + \gamma_J(\alpha)$
- Γ_{cusp} is the universal cusp anomalous dimension
- $\frac{\partial \ln H(xP)}{\partial \ln \mu} + \frac{\partial \ln H(\bar{x}P)}{\partial \ln \mu} + \frac{\partial \ln J}{\partial \ln \mu} = \left(\frac{\partial \ln C}{\partial \ln \mu}\right)_{\mathrm{x} \to \mathrm{y}}$

- Solving the three RG equations independently

Correcting the matching kernel

- Resummed Sudakov factor: $H = |H|e^{i\hat{A}}$

$$|H(\mu)| = |H(\mu_1,\mu_2)|e^{S(\mu_1,\mu)+S(\mu_2,\mu)-a_c(\mu_1,\mu)-a_c(\mu_2,\mu)} \times \left(\frac{2xP_z}{\mu_1}\right)^{-a_\Gamma(\mu_1,\mu)}\left(\frac{2\bar{x}P_z}{\mu_2}\right)^{-a_\Gamma(\mu_2,\mu)}$$

$$\hat{A}^{\rm RGR}(xP_z,\bar{x}P_z,\mu_1,\mu_2) = \pi{\rm sign}(z)\left[\frac{\alpha_s(\mu_1)C_F}{2\pi}\left(1-\ln\frac{4x^2P_z^2}{\mu_1^2}\right)-\frac{\alpha_s(\mu_2)C_F}{2\pi}\left(1-\ln\frac{4\bar{x}^2P_z^2}{\mu_2^2}\right)+2\int_{\mu_1}^{\mu_2}\frac{\Gamma_{\rm cusp}}{\mu}d\mu\right]$$

- Resummed Jet function:

$$J(\Delta,\mu) = e^{[-2S(\mu_i,\mu)+a_J(\mu_i,\mu)]}\tilde{J}_z(l_z=-2\partial_\eta,\alpha_s(\mu_i))\left[\frac{\sin(\eta\pi/2)}{|\Delta|}\left(\frac{2|\Delta|}{\mu_i}\right)^\eta\right]_*\frac{\Gamma(1-\eta)e^{-\eta\gamma_E}}{\pi}\bigg|_{\eta=2a_\Gamma(\mu_i,\mu)}$$

- $C_{TR} = (H\otimes J)_{TR}\otimes(H\otimes J)_{NLO}^{-1}\otimes C_{NLO}$
- Inverse matching:

$$C_{TR}^{-1} = C_{NLO}^{-1}\otimes(H\otimes J)_{NLO}\otimes(H\otimes J)_{TR}^{-1}$$

What are the scale choices of $\mu_{1,2}$ and μ_i?

Scale choices of resummation

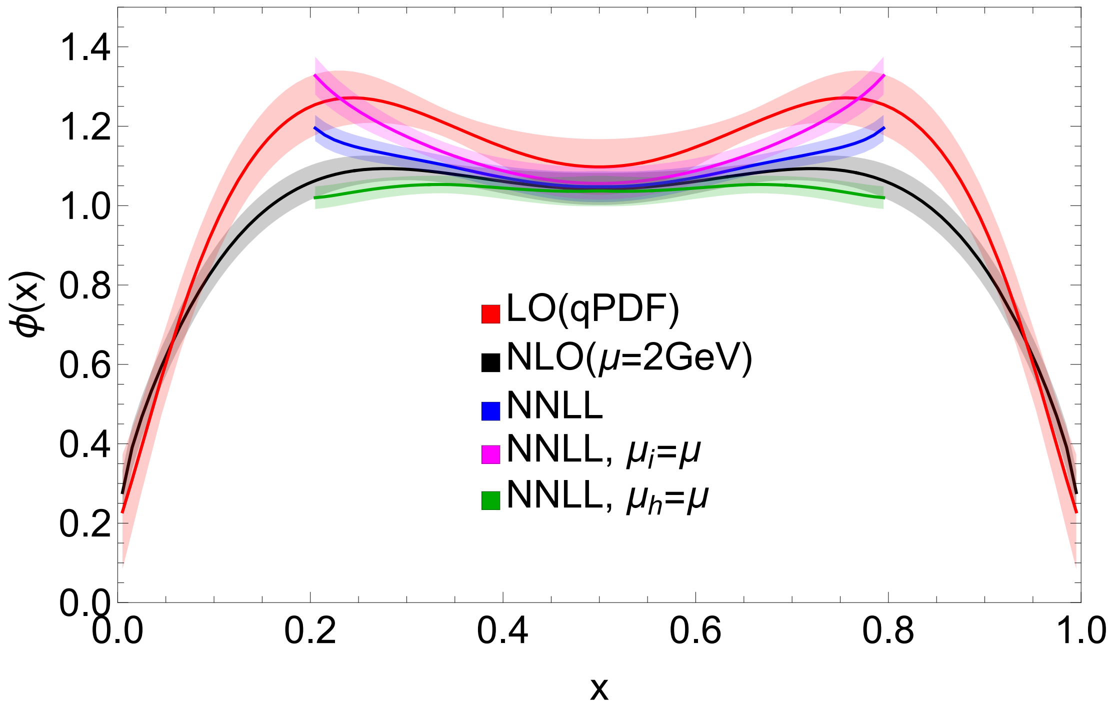
- Hard scale:
  - $H(xP,\mu)$: quark momentum $\mu_{h_1} = 2xP$
  - $H(\bar{x}P,\mu)$: anti-quark momentum $\mu_{h_2} = 2\bar{x}P$
- Semi-hard scale:
  - $J(|y-x|P,\mu)$: gluon momentum $\mu_i = 2|y-x|P$?
  - This scale choice is not applicable because $\mu_i \to 0$ hits the Landau Pole for any given x!
- Actual x-dependent semi-hard scale found in $\int dy\, J(|x-y|)\phi(y)$
  - $2xP$ when $x \to 0$
  - $2\bar{x}P$ when $x \to 1$
  - Choosing $\mu_i = 2\min(x,\bar{x})\,P$

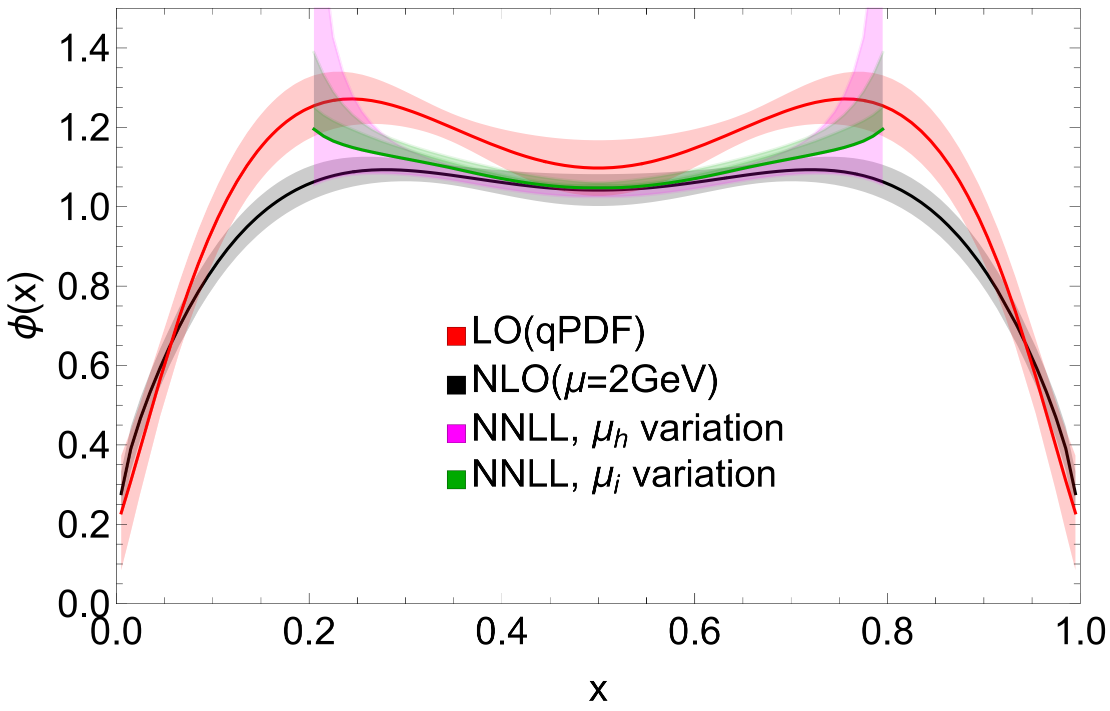
Becher, Neubert & Pecjak JHEP(2007)

Matching with Resummed Kernel

Scale variation: $\mu_i \to c * \mu_i,\ c = [\frac{1}{\sqrt{2}}, \sqrt{2}]$

When scale variation becomes large, perturbation theory is no longer reliable

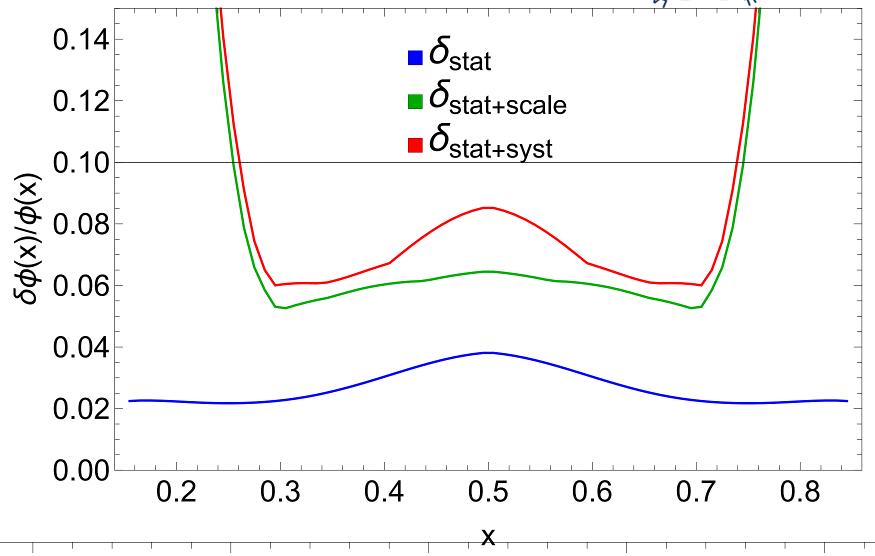


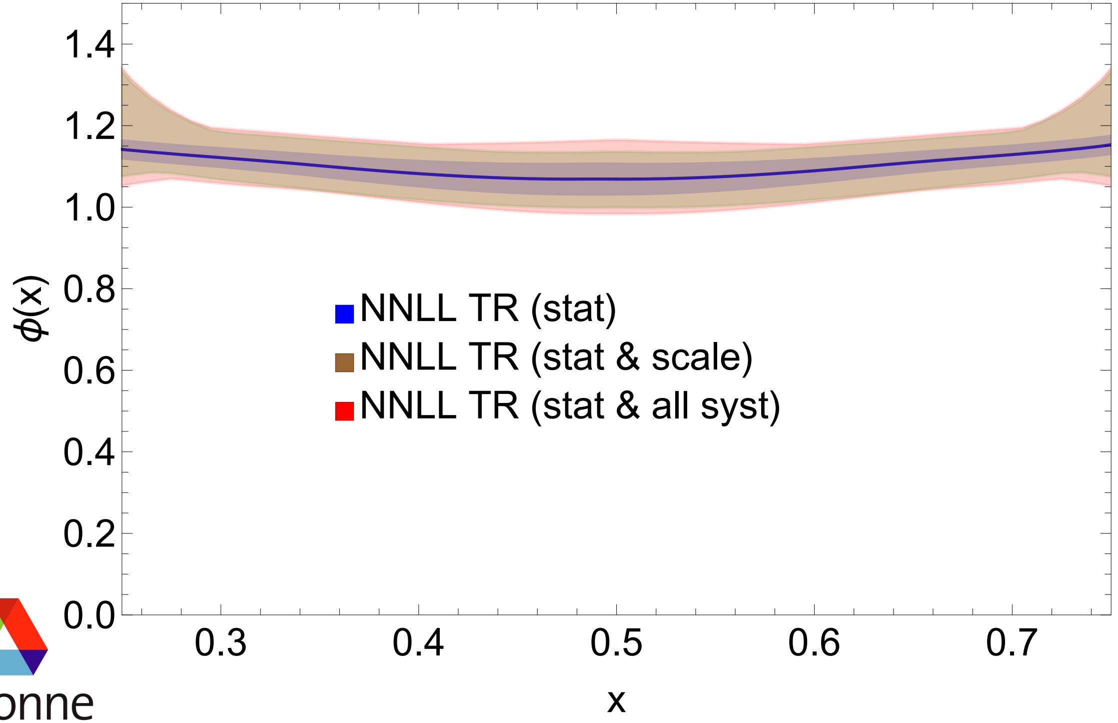


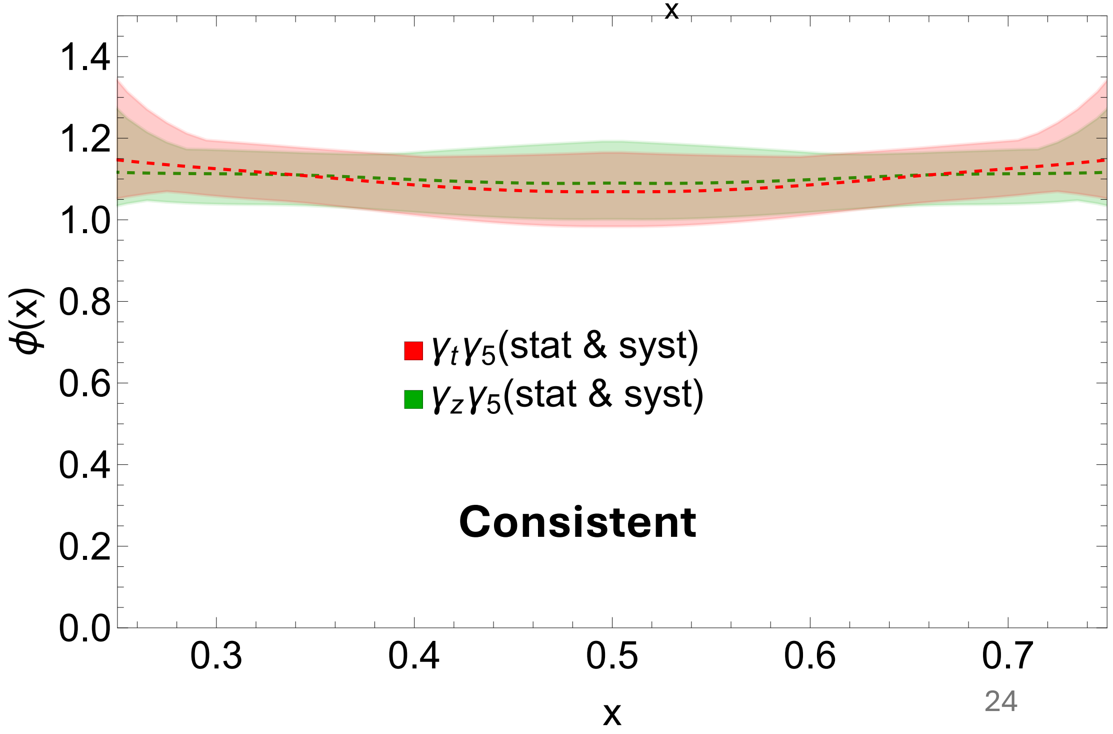
Including all systematics

- Different z_s
- Different extrapolations
- Scale Variation

$x \in [0.25, 0.75]$
$\delta_{total} \approx 10\%$

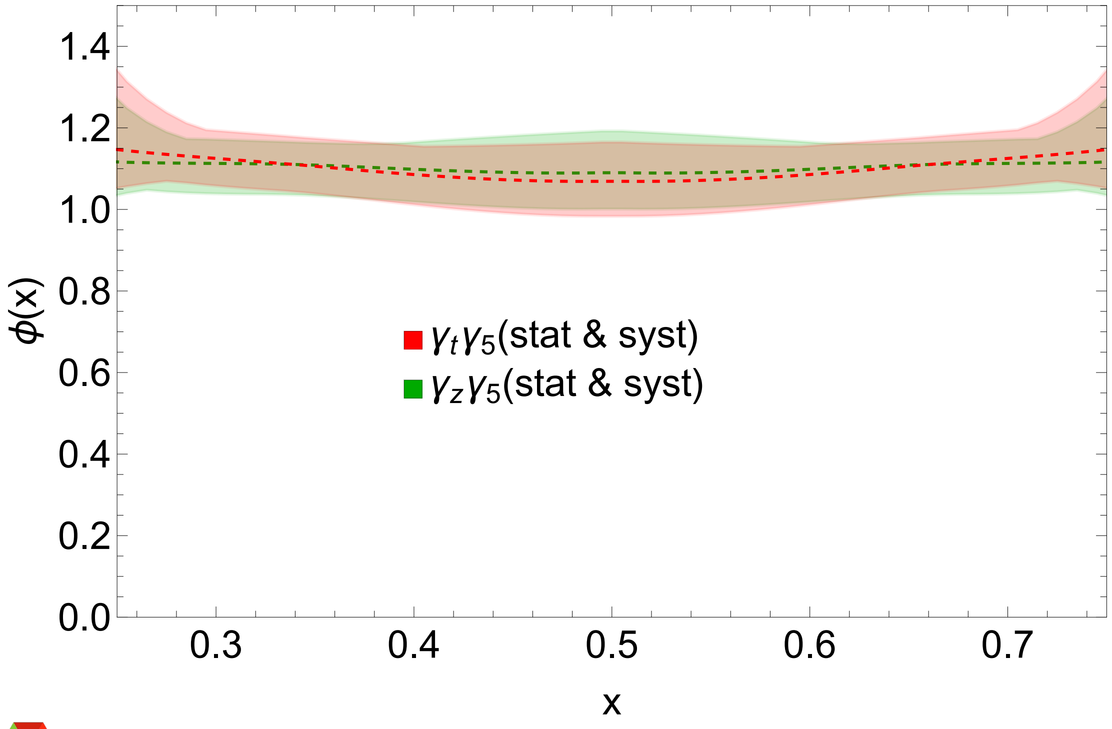




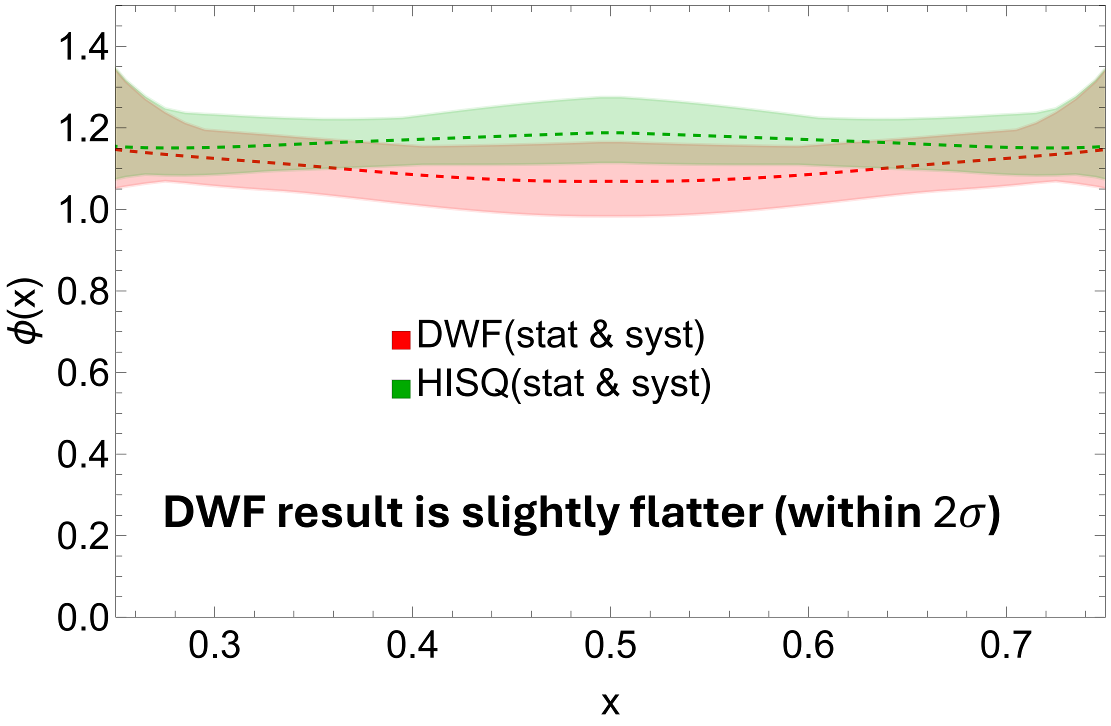


Comparison of Final Results

- Different operators



- Different fermion actions on similar lattice



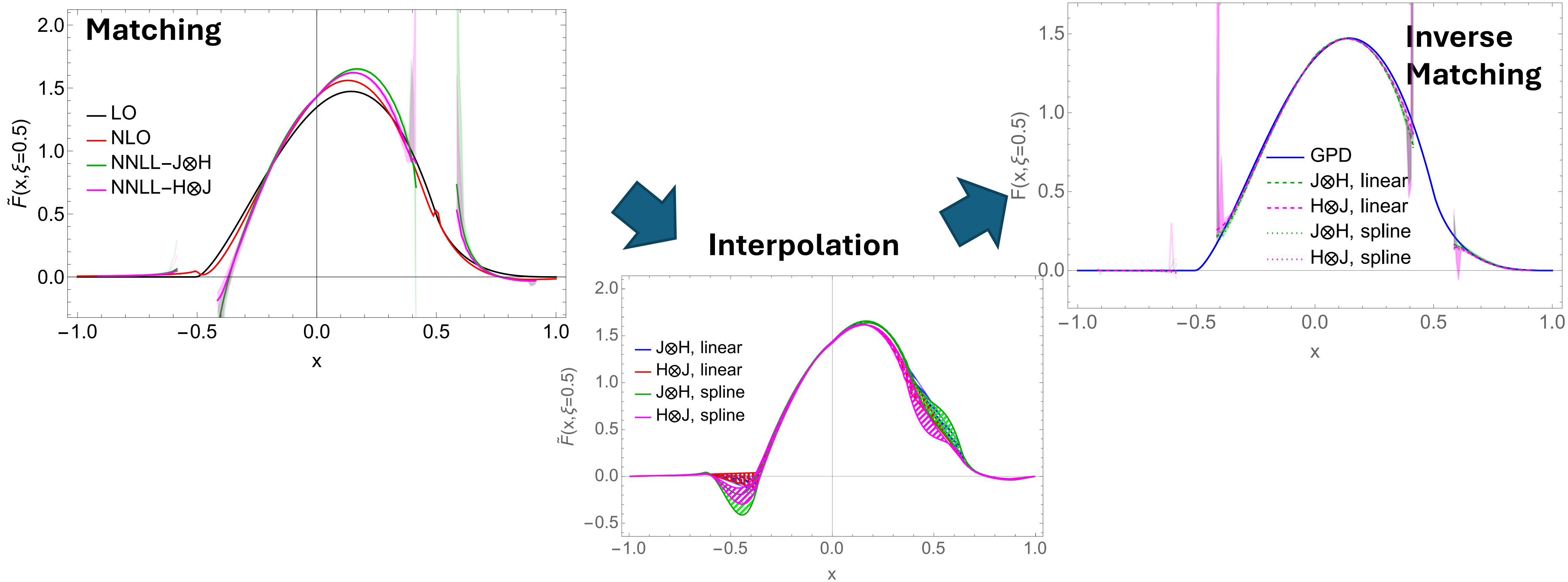
Conclusion and Outlook

- We present a pion DA calculation using gauge ensembles with domain wall fermions;
- We propose and develop a more robust method to resum the small-momentum logarithms in the perturbative matching kernel of DA, the first implementation of threshold resummation in the LaMET DA calculation;
- We observe a slightly flatter distribution for domain wall fermions.

- ❑ Continuum limit is needed for a more conclusive comparison
- ❑ Larger pion momentum is needed to extend the x range of calculation
- ❑ More precise measurement of DA longtail is needed (Coulomb Gauge?)
- ❑ Generalize the method to GPD matching

Resummation in GPD at non-zero skewness

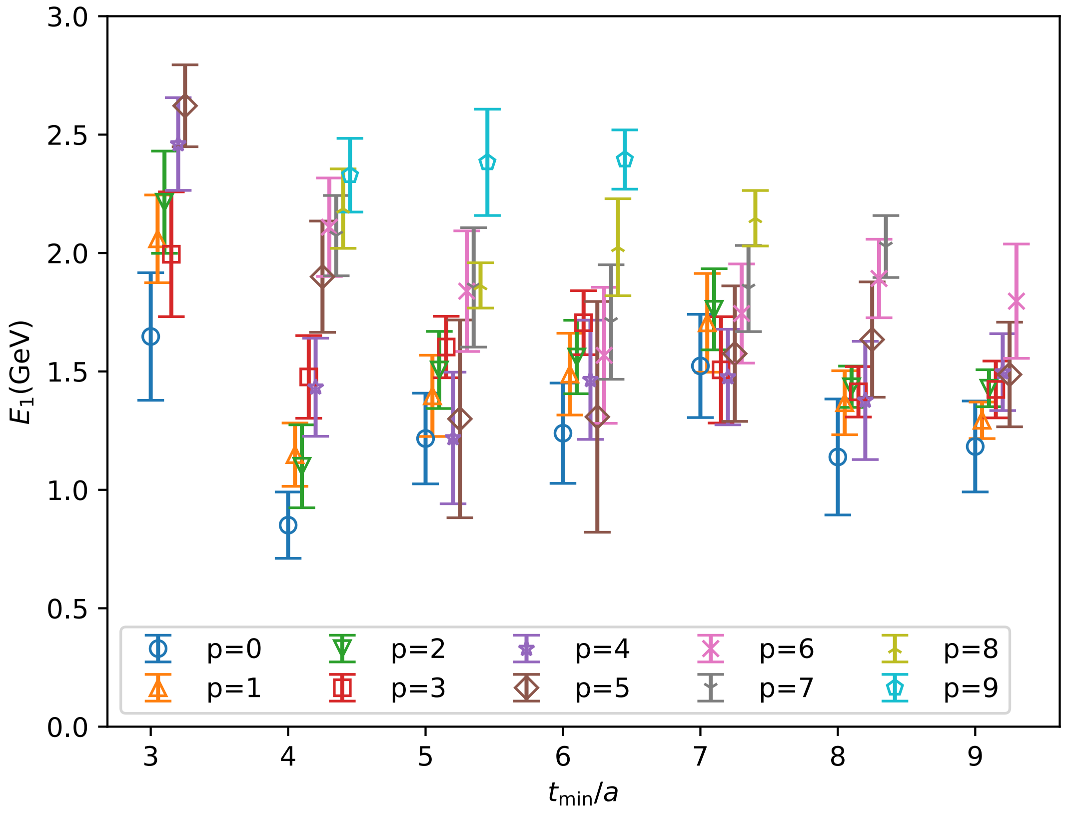
LaMET (inverse) matching is highly local and not spreading out higher-power or non-perturbative effects.

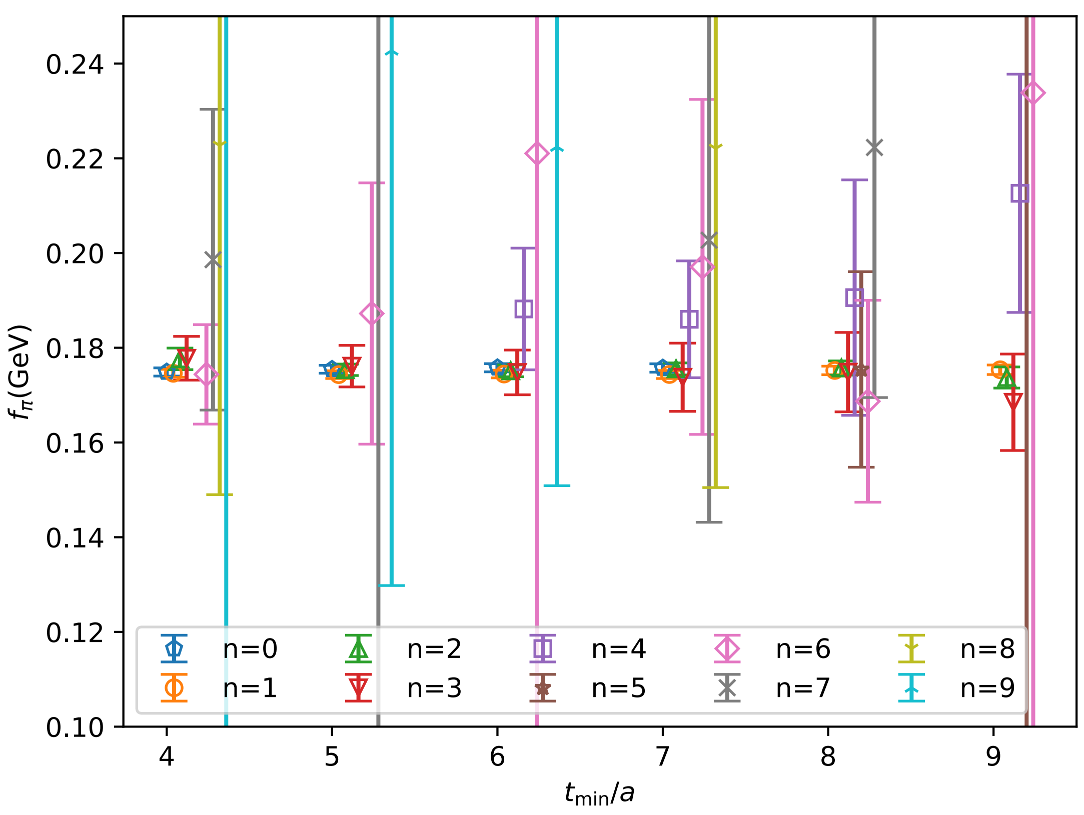


Thank you for listening!

Backup Slides

Fits of energy and extraction of f_π





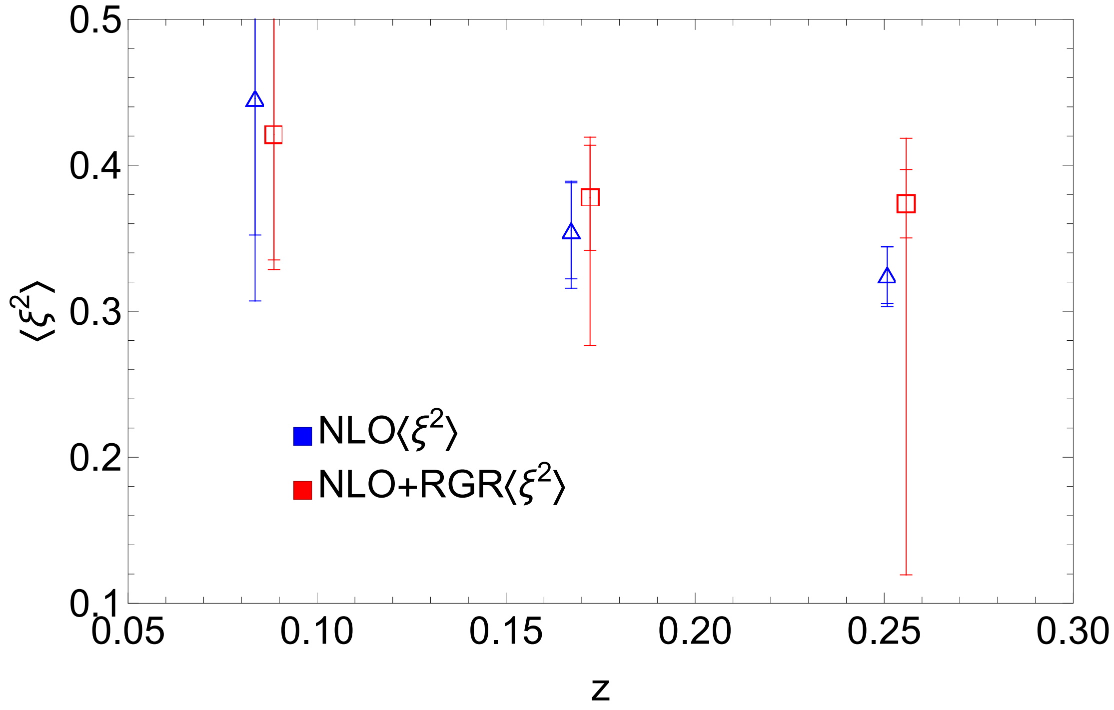
Extracting Moments from OPE

- RG-invariant ratio:

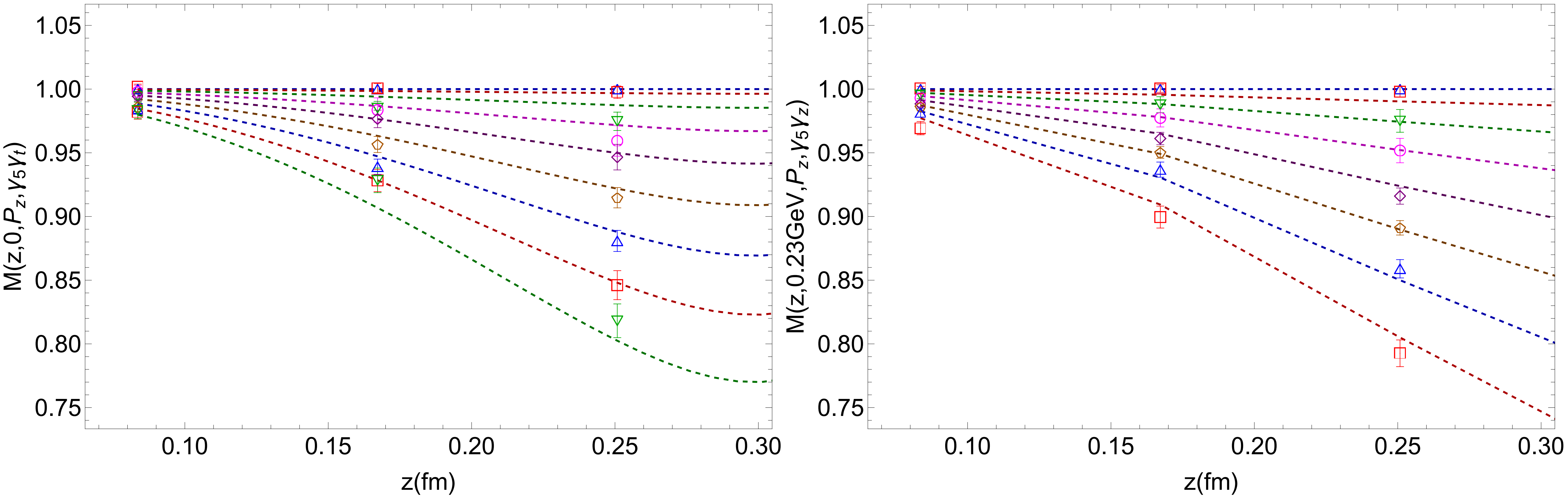
$$\mathcal{M}(z,P_1,P_2) = \lim_{a\to 0} \frac{H^B(z,P_2,a)}{H^B(z,P_1,a)} = \frac{H^R(z,P_2)}{H^R(z,P_1)}$$

- Fit to OPE

$$\mathcal{M}(z,P_1,P_2) \approx \frac{\sum_{n=0}^{\infty}\sum_{m=0}^{n} \frac{(\frac{-izP_2}{2})^n}{n!} C_{nm}(z,\mu)\langle\xi^m\rangle}{\sum_{n}^{\infty}\sum_{m=0}^{n} \frac{(\frac{-izP_1}{2})^n}{n!} C_{nm}(z,\mu)\langle\xi^m\rangle}$$

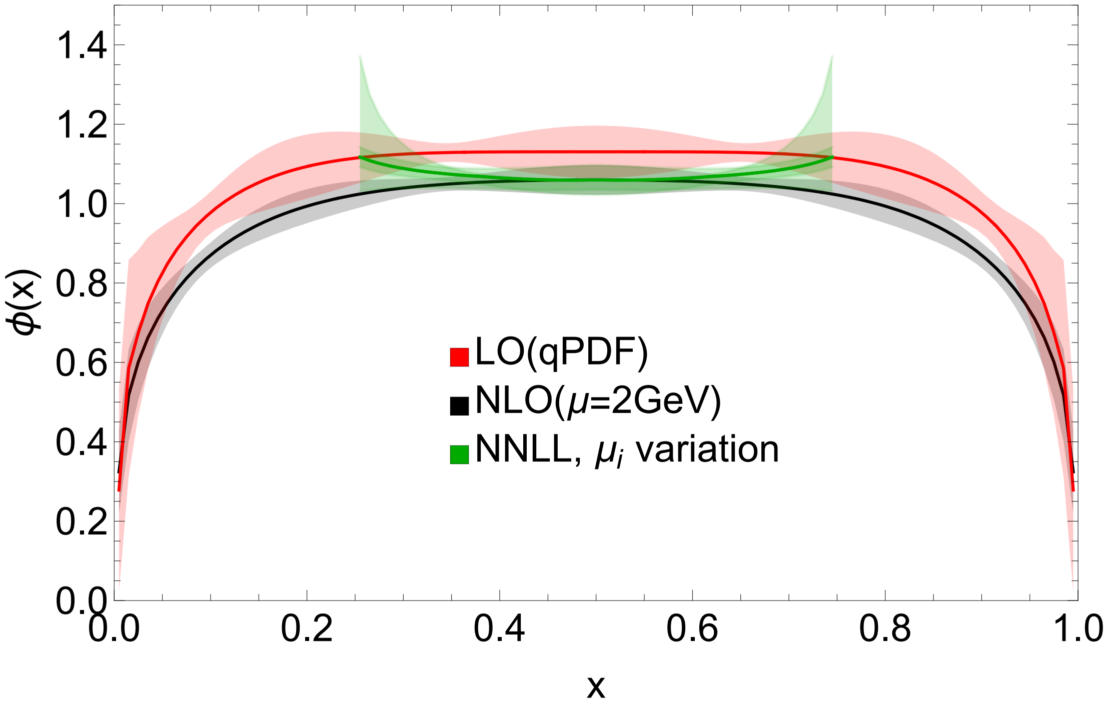


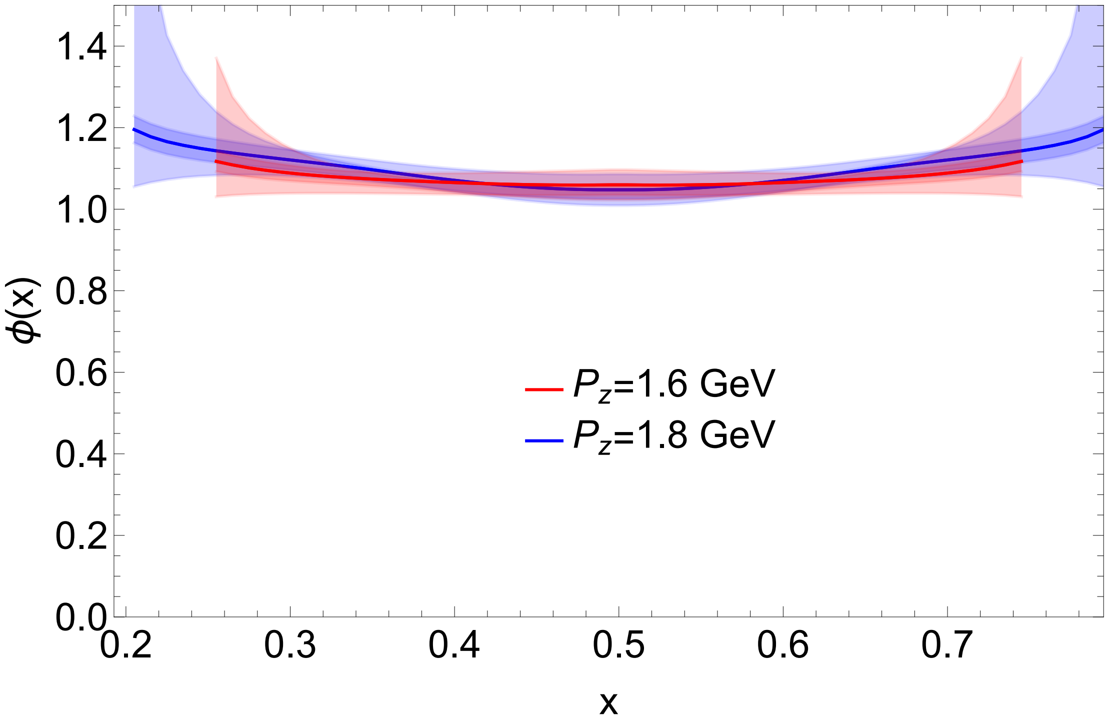
Consistency with OPE (Ratio)



Momentum dependence of calculable range

- Compare with $P_z = 1.6GeV$
- The range of calculation increases with momentum





Pion and Kaon DA on HISQ ensembles

arxiv:2407.00206