

# MINERvA's Axial FF Measurement

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CHARTERED 1693

# Overview

- MINERvA finished a 10 year run at Fermilab in 2019 with an unprecedentedly large dataset.
- One of the the initial goals of MINERvA was to measure the axial vector form factor, only accessible in weak interactions.
- While  $F_A$  is important it it's own right, neutrino oscillation experiments have to reconstruct the neutrino energy from neutrino nucleus collisions.
  - Messy, messy.
  - A solid understanding of neutrino nucleon scattering is a necessary input
- Also, neutrino experiments would really benefit from a standard candle cross-section
  - Well,  $\nu e^- \rightarrow \nu e^-$  &  $\nu_\mu e^- \rightarrow \nu_e \mu^-$  but the cross-sections are small and you can't get  $E_\nu$  easily.

# Formalism - $d\sigma/dQ^2$

$$\frac{d\sigma}{dQ^2} \begin{pmatrix} \nu n \rightarrow l^- p \\ \bar{\nu} p \rightarrow l^+ n \end{pmatrix} = \frac{M^2 G_F^2 \cos^2 \theta_c}{8\pi E_\nu^2} \left[ A(Q^2) \mp B(Q^2) \frac{(s-u)}{M^2} + C(Q^2) \frac{(s-u)^2}{M^4} \right] \quad (4)$$

$$A(Q^2) = \frac{m^2 + Q^2}{4M^2} \left[ \left( 4 + \frac{Q^2}{M^2} \right) |F_A|^2 - \left( 4 - \frac{Q^2}{M^2} \right) |F_V^1|^2 + \frac{Q^2}{M^2} \left( 1 - \frac{Q^2}{4M^2} \right) |\xi F_V^2|^2 + \frac{4Q^2}{M^2} \text{Re} F_V^{1*} \xi F_V^2 + \mathcal{O} \left( \frac{m^2}{M^2} \right) \right],$$

$$B(Q^2) = \frac{Q^2}{M^2} \text{Re} F_A^* (F_V^1 + \xi F_V^2),$$

$$C(Q^2) = \frac{1}{4} \left( |F_A|^2 + |F_V^1|^2 + \frac{Q^2}{4M^2} |\xi F_V^2|^2 \right) \quad (5)$$

- Vector form factors  $F_V^1(Q^2)$  and  $F_V^2(Q^2)$  measured and parameterized by you all.
- $F_A^1(Q^2)$  only appears in weak interactions which give neutrinos a chance to measure it.

# Measurement concept

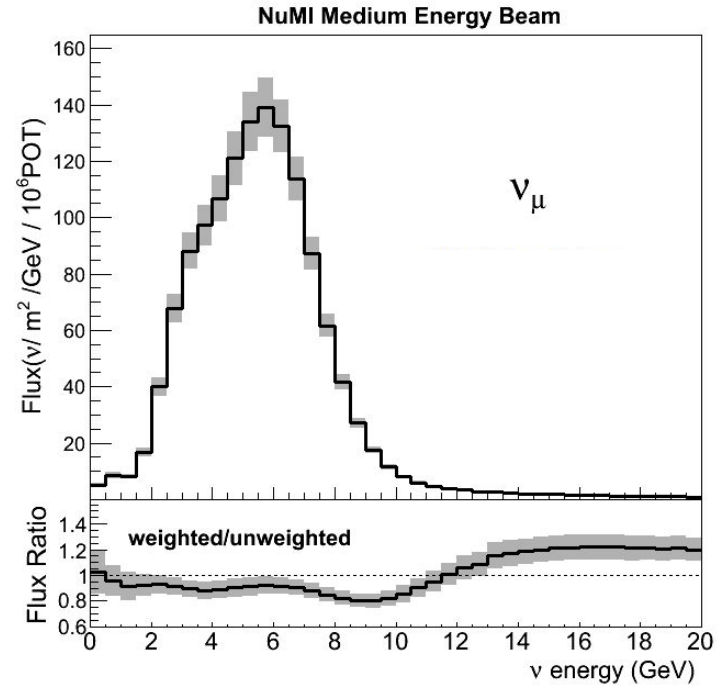
- In MINERvA neutrinos interact on plastic scintillator, CH
- The nucleons in carbon are bound and are therefore moving targets (Fermi motion, etc).
- The free protons in hydrogen are at rest.
- Elastic scattering  $\bar{\nu}_\mu p \rightarrow \mu^+ n$  is a 2-body problem.
- We don't know the energy of each incoming neutrino but we do know its direction and we can measure  $(p_\mu, \theta_\mu)$ .
- This is enough to completely constrain the system
  - namely, momentum must be conserved in the reaction plane and perpendicular to it.
  - not true for interactions on C or inelastic interactions.
- We'll use this to select a sample enriched in elastic events with sidebands that can be used to constrain backgrounds in the signal region.

## Main Reference

T. Cai et al., The MINERvA Collaboration, Nature 614 (2023) 7946, 48-53

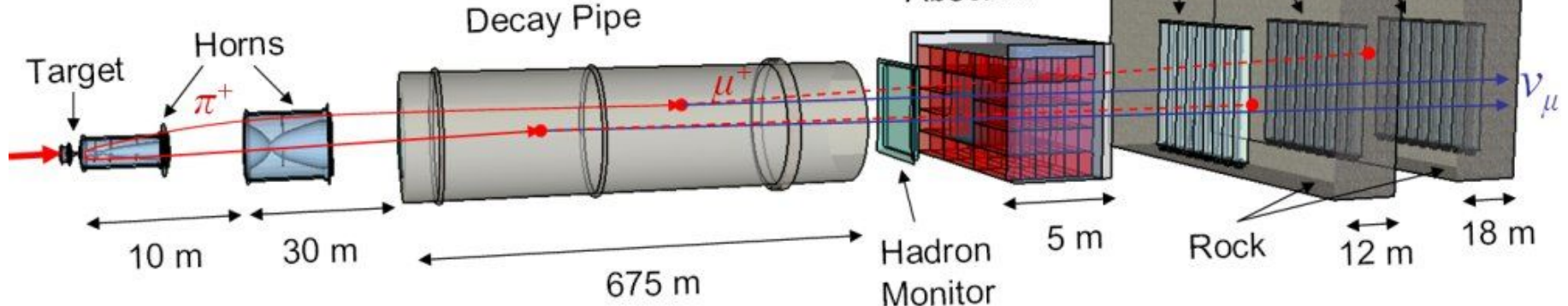
# What was MINERvA?

- MINERvA was a dedicated neutrino interaction experiment that ran in the NuMI beam at Fermilab from 2009-2019.
- Two beam configurations:
  - Low Energy  $\sim <3$  GeV
  - Medium Energy  $\sim <6$  GeV
- Goal was to measure inclusive and exclusive scattering off of a wide range of nuclei (Pb/Fe/C/O/He)
- Collected  $\sim 10^7$  neutrino interactions



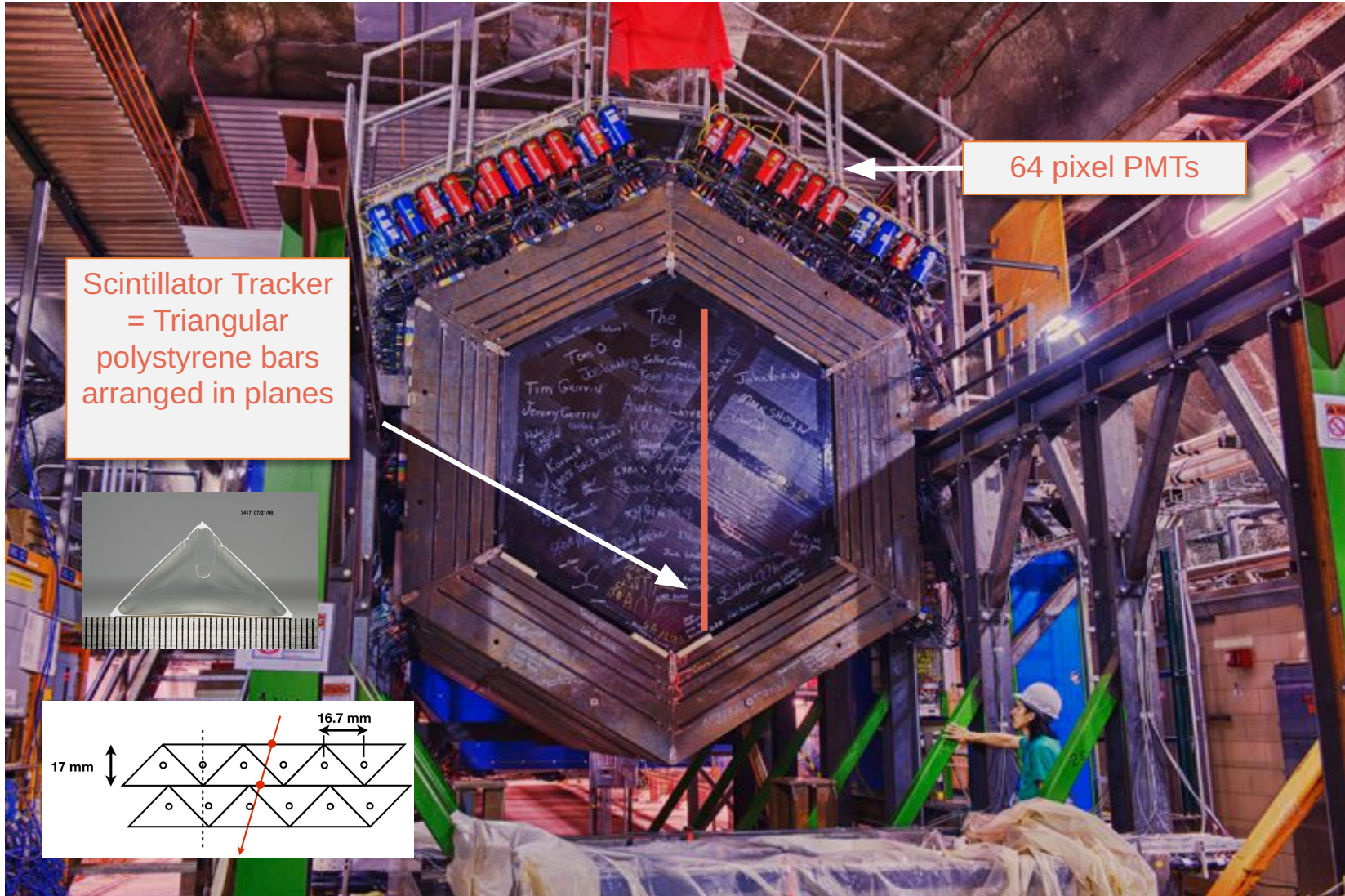
flux correction

## The NuMI Beam



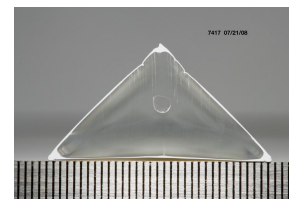


# The MINERvA Experiment



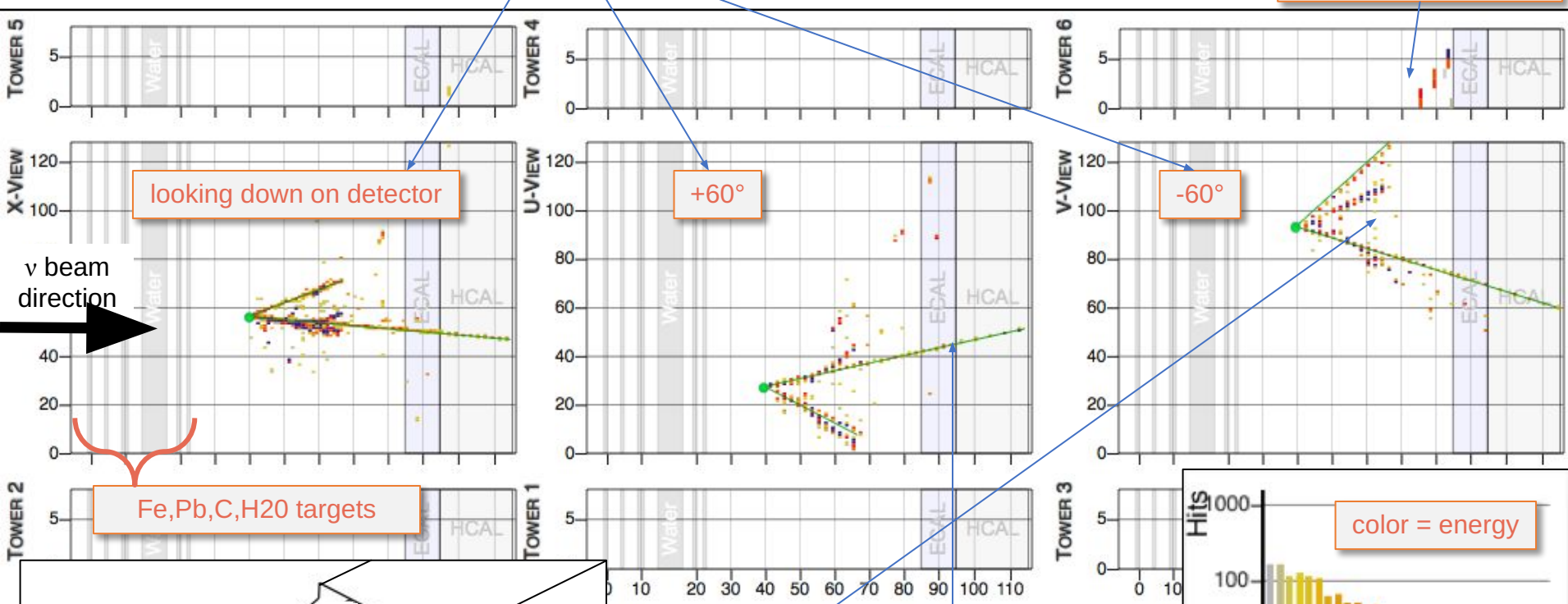
# An event in MINERvA's scintillator tracker

Triangular polystyrene bars



Particle leaves the inner detector, stops in outer iron calorimeter

3 stereo views, X-U-V, shown separately



looking down on detector

+60°

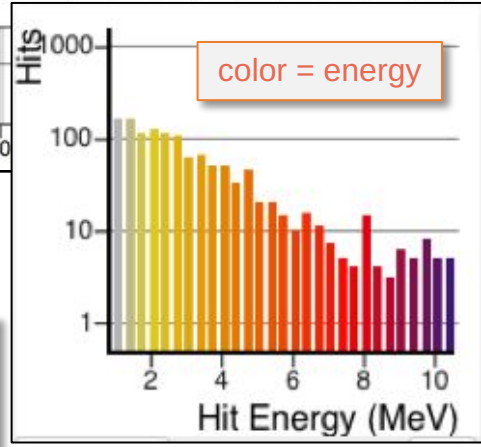
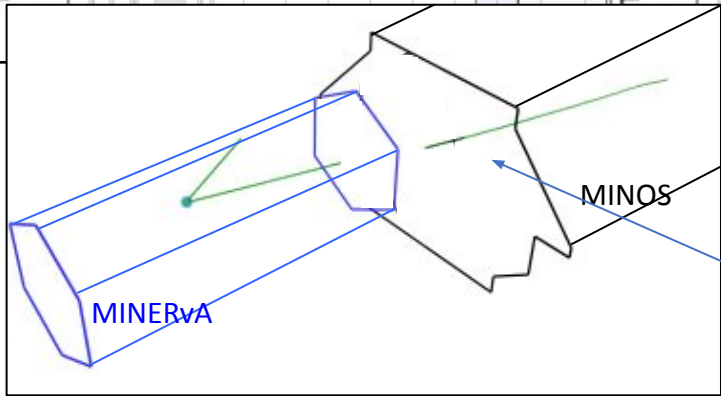
-60°

Fe,Pb,C,H2O targets

color = energy

Stops in Scintillator, best hadron particle ID

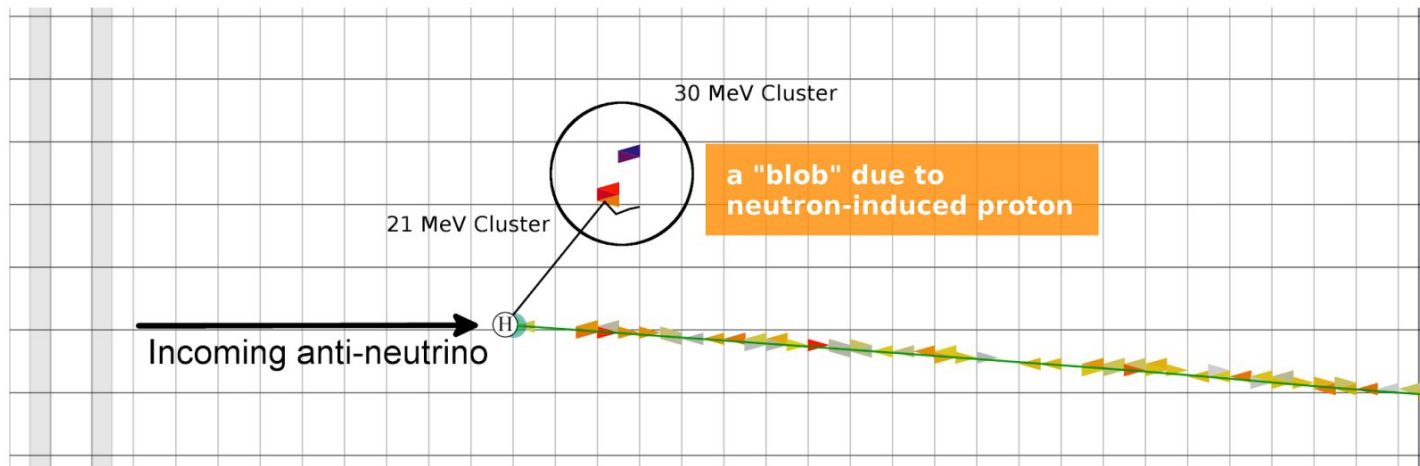
Muon leaves the back of the detector headed toward magnetized spectrometer (a.k.a. MINOS Near Detector)



# Event selection

- We are looking for  $\bar{\nu}_\mu p \rightarrow \mu^+ n$  (“elastic scattering”)
- Backgrounds:
  - Primary: same process but on a proton bound in C (“quasi-elastic”)
  - Multi-nucleon emission & inelastic (i.e., pion production) events that produce isolated blobs but not much other activity
- $E_\nu$  and  $Q^2$  are reconstructed from the muon arm assuming a target at rest:

$$E_{\bar{\nu}} = \frac{M_n^2 - M_p^2 - m_\mu^2 + 2M_p E_\mu}{2(M_p - E_\mu + p_\mu \cos \theta_\mu)},$$
$$Q_{QE}^2 = 2E_{\bar{\nu}}(E_\mu - p_\mu \cos \theta_\mu) - m_\mu^2,$$

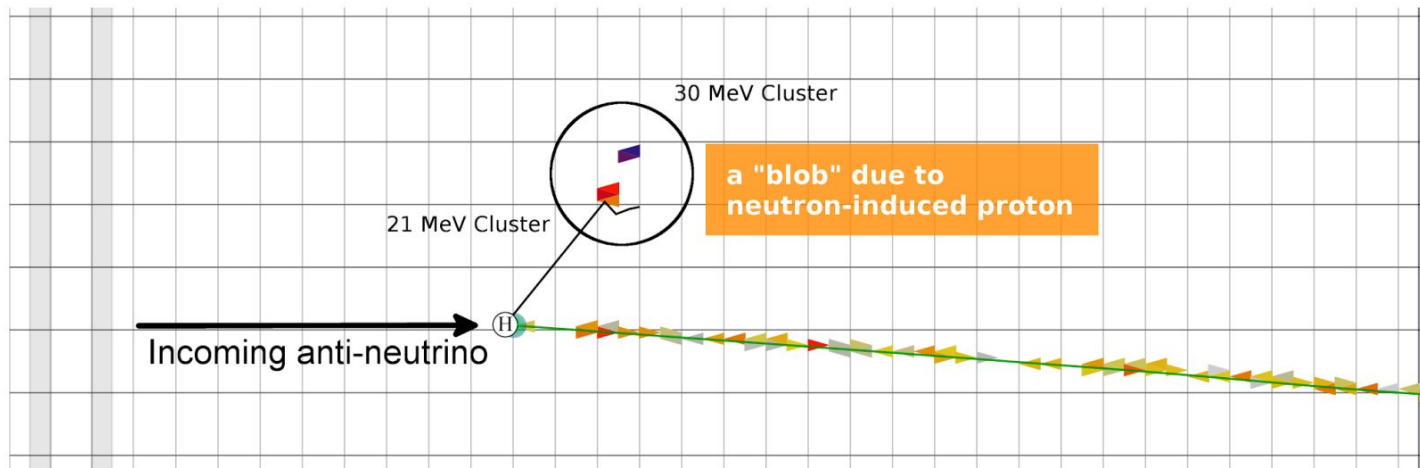




# Event selection

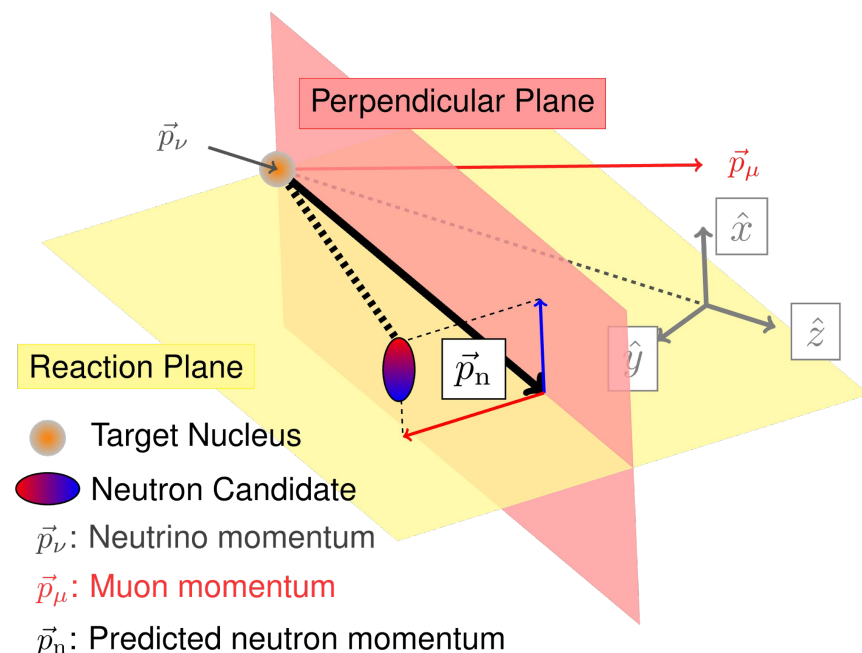
- Require:
  - $\mu^+$  track starting in the scintillator tracker and reconstructed in MINOS with  $\Theta_\mu < 20^\circ$  and  $1.5 < p_\mu < 20 \text{ GeV}/c$
  - No other tracks from protons or pions at the vertex
  - Isolated blob or proton-like track 20 cm or more from the vertex. No other significant blobs are allowed.
  - Because recoil energy  $= q_0 = Q^2/2m_n$  we require that the energy not associated with the muon is less than a  $Q^2$  dependent threshold

Condition	$E_{\text{max}}^{\text{recoil}}$ (GeV)	
$Q_{\text{QE}}^2 < 0.3 \text{ (GeV}/c)^2$	$0.04 + 0.43Q_{\text{QE}}^2/(\text{GeV}/c)^2$	→ 170 MeV
$Q_{\text{QE}}^2 < 1.4 \text{ (GeV}/c)^2$	$0.08 + 0.3Q_{\text{QE}}^2/(\text{GeV}/c)^2$	→ 500 MeV
$Q_{\text{QE}}^2 > 1.4 \text{ (GeV}/c)^2$	0.50	

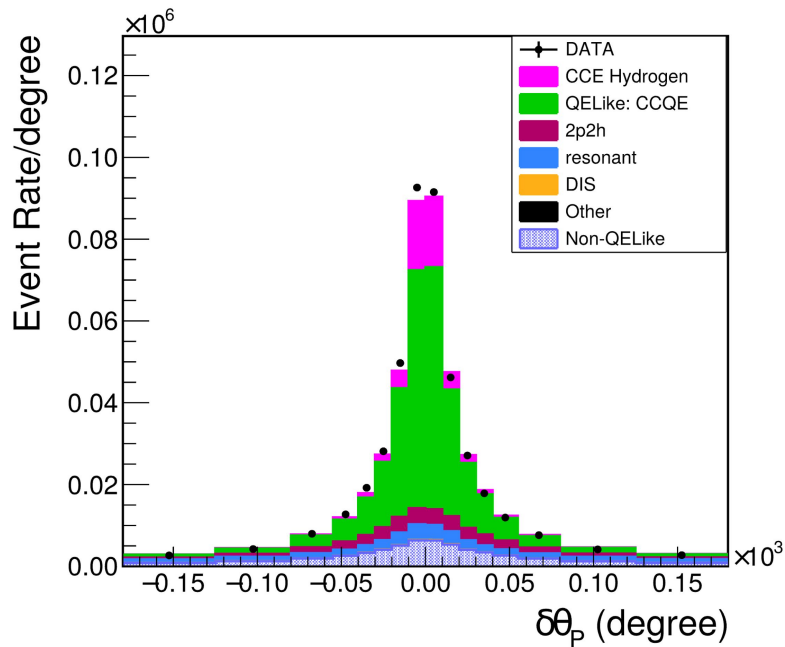


# Signal vs Background

- When scattering on a free proton the direction of the neutron is predicted from the muon kinematics.
- Not the case for bound protons due to initial state interactions (e.g., Fermi motion) and final state interactions (neutron rescattering).
- Other backgrounds are from multi-nucleon scattering (2p2h) and inelastic processes.
- We separate signal from background using the difference between the observed neutron angle and what would be expected for scattering off a free proton:  $\delta\theta_R$  and  $\delta\theta_P$



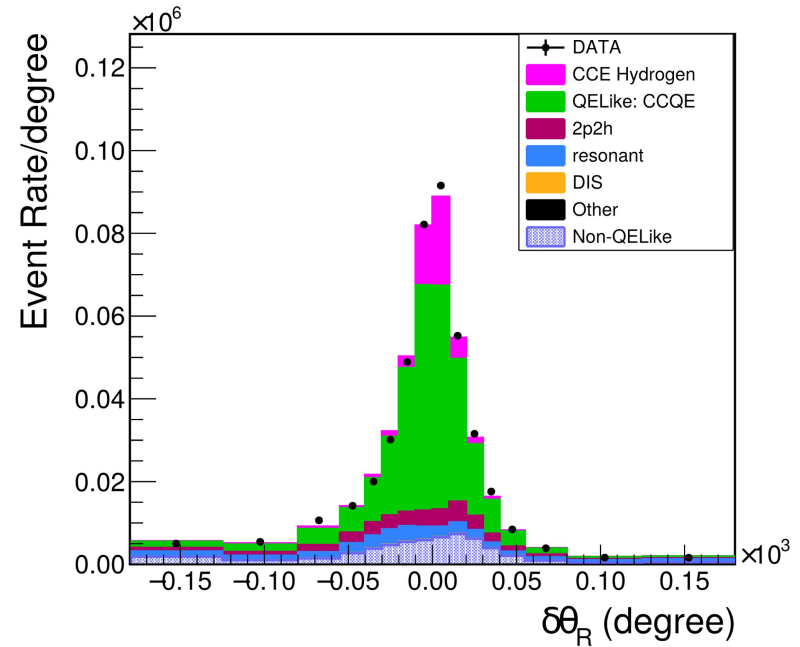
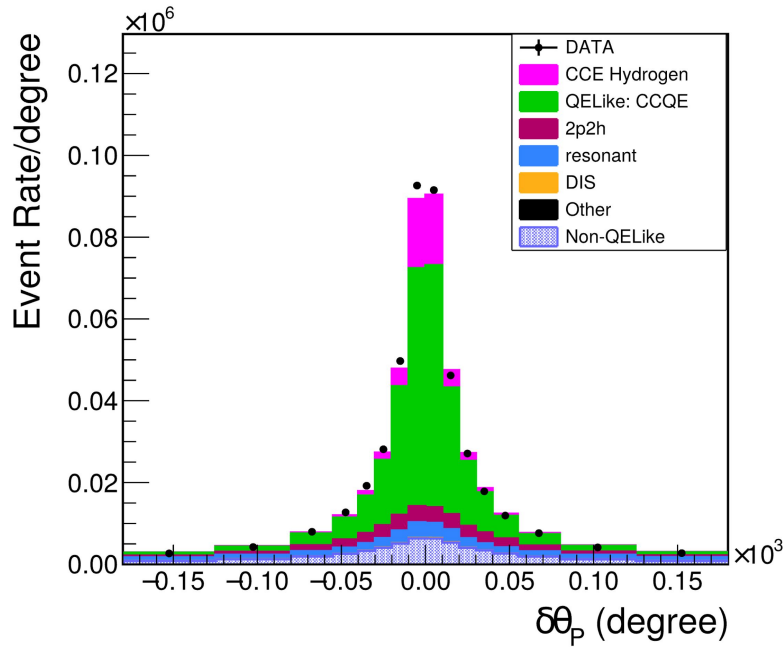
# $\delta\theta_P$ and $\delta\theta_R$ distributions



## Event categories

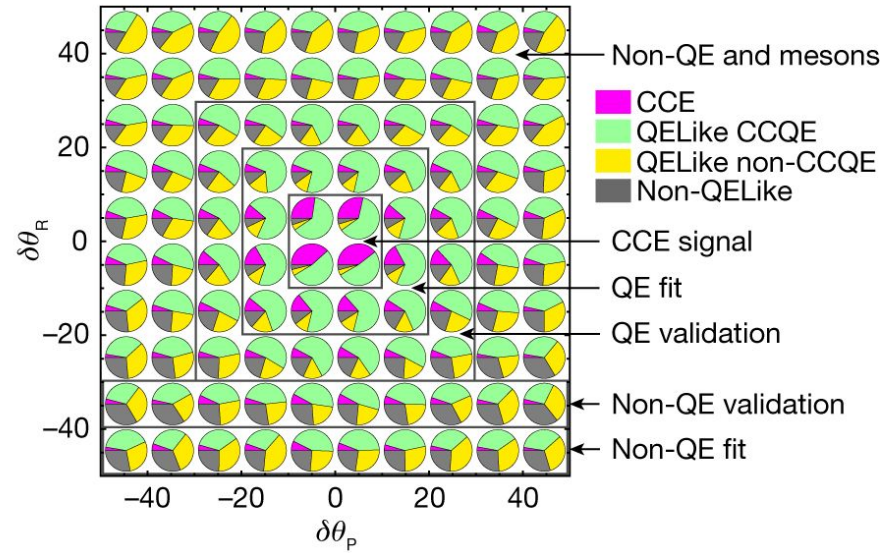
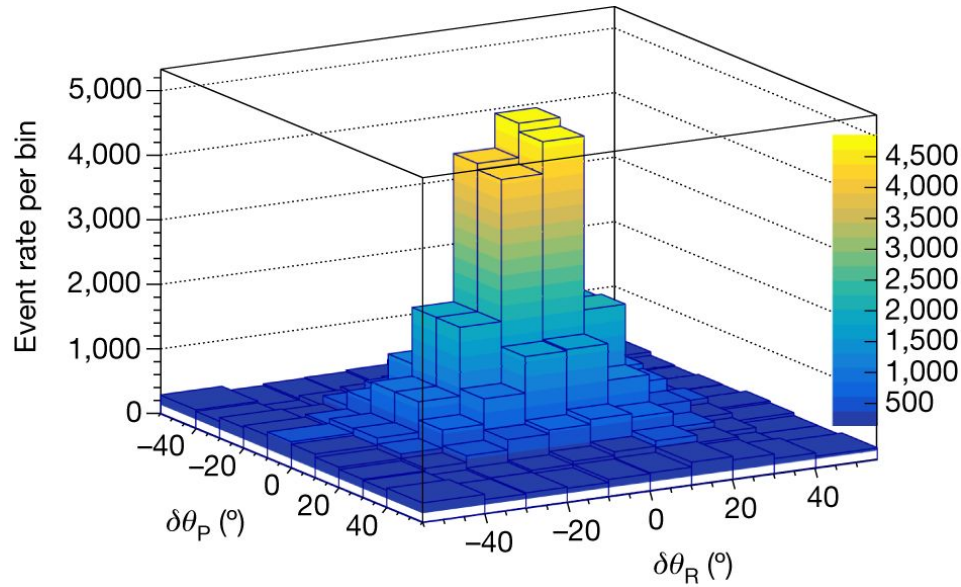
- Based (mostly) on particles in the final state according to MC truth
- CCE=signal  $\bar{\nu}_\mu p \rightarrow \mu^+ n$
- QELike:
  - only  $\mu^+$  and nucleons in final state
  - reconstructed proton tracks must have  $T_p < 150$  MeV
- QELike backgrounds:
  - CCQE=quasi-elastic
  - 2p2h=multi-nucleon
  - resonant+DIS=inelastic
- Non-QELike:
  - events with other particles in the final state or those with  $T_p > 150$  MeV

# $\delta\theta_P$ and $\delta\theta_R$ distributions



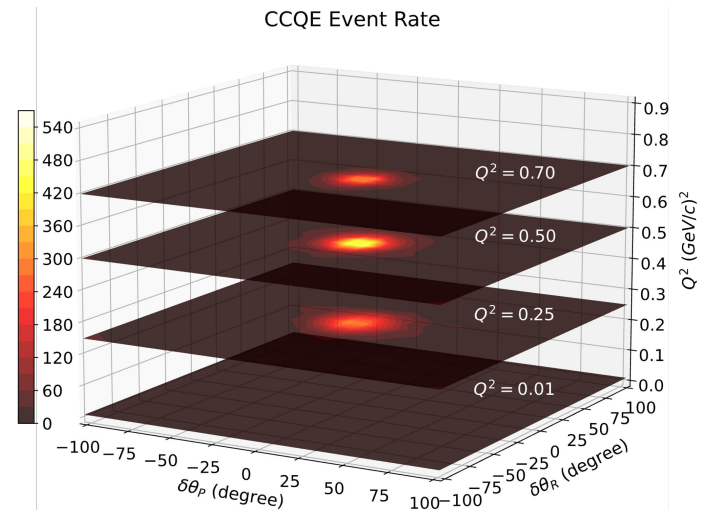
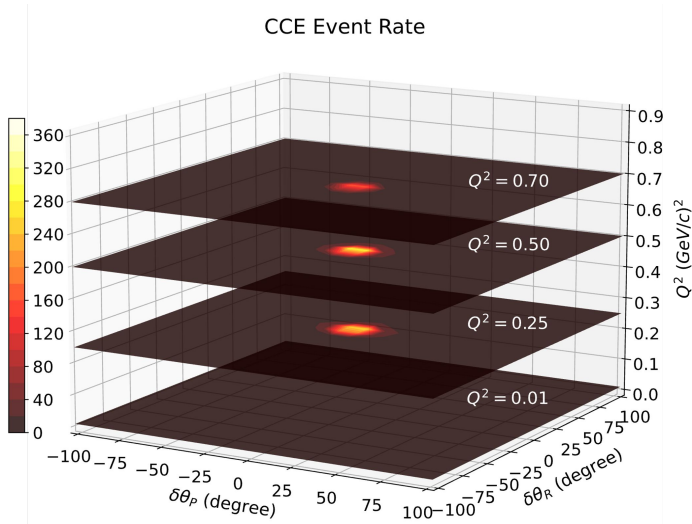
- Signal concentrated around  $\delta\theta_P$  &  $\delta\theta_R = 0^\circ$
- Symmetry in  $\delta\theta_P$
- Asymmetry for backgrounds in  $\delta\theta_R$

# $\delta\theta_P$ and $\delta\theta_R$ distributions: in two dimensions

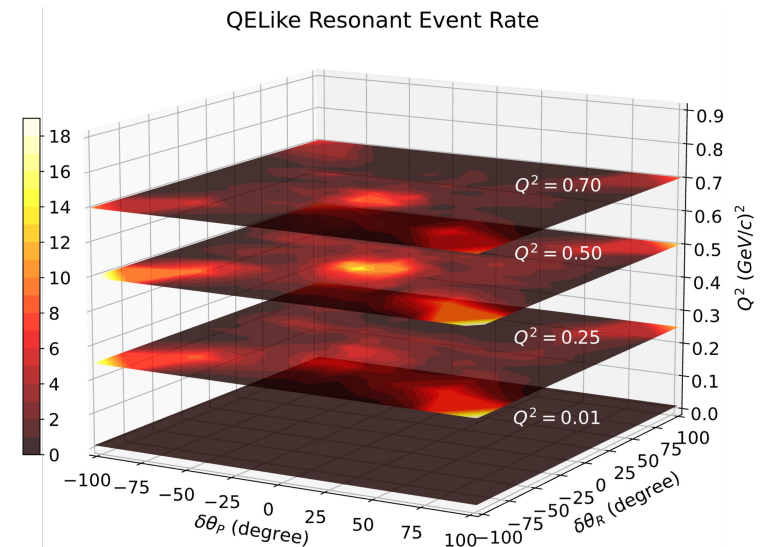
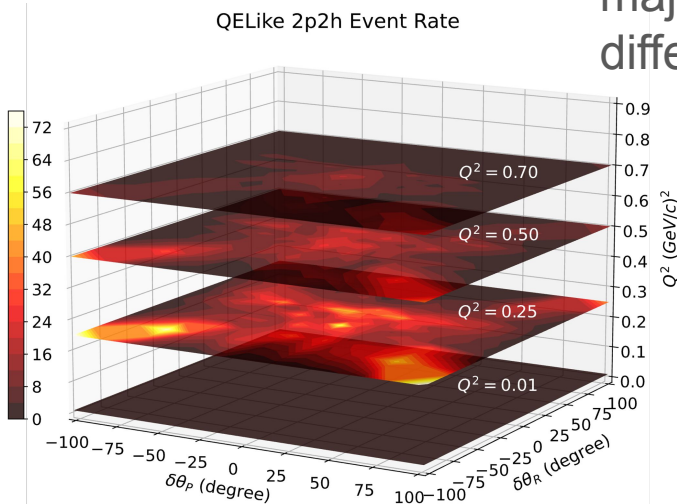




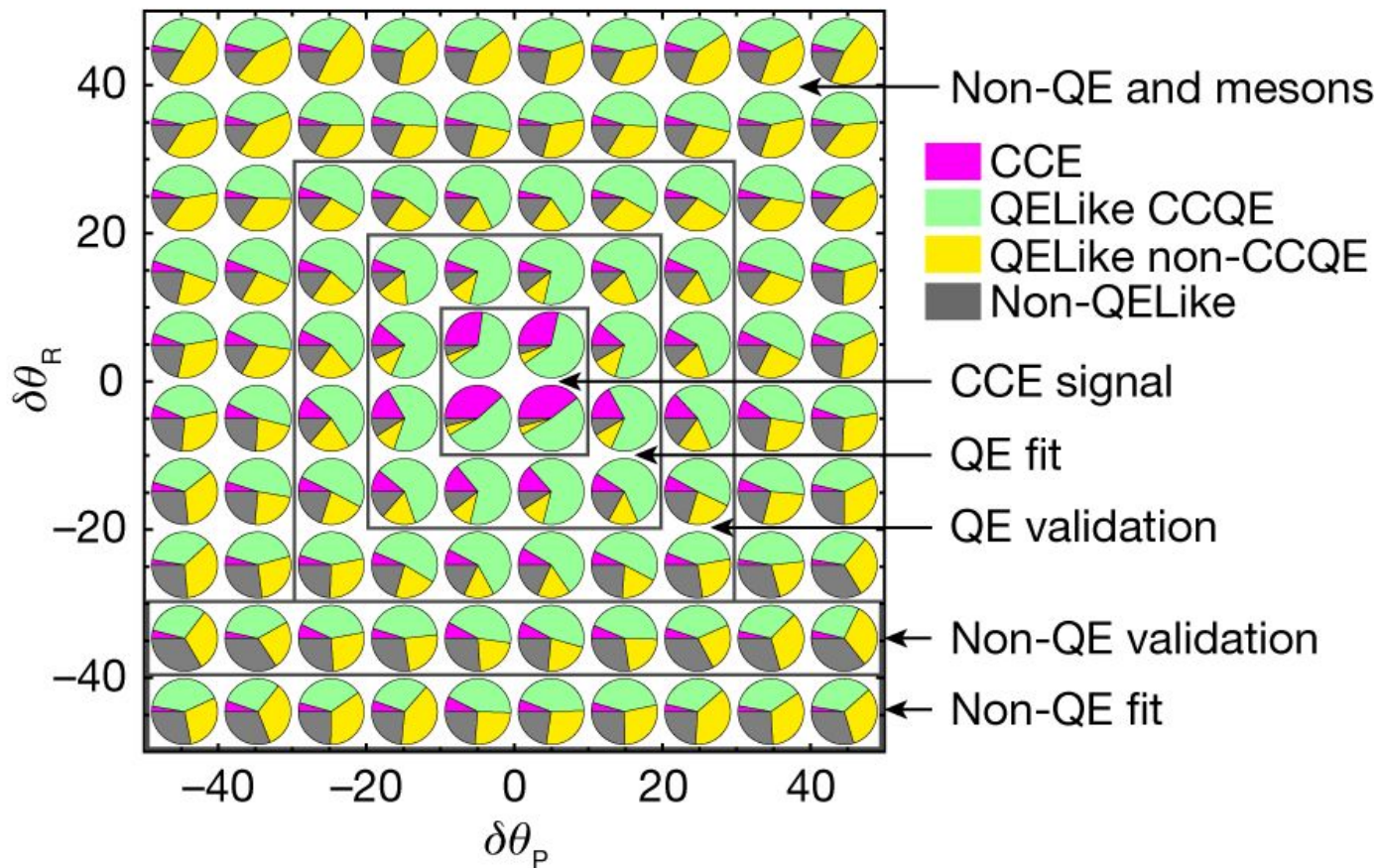
# $\delta\theta_P$ and $\delta\theta_R$ distributions: in three dimensions



Heat maps of signal and major backgrounds at different  $Q^2$

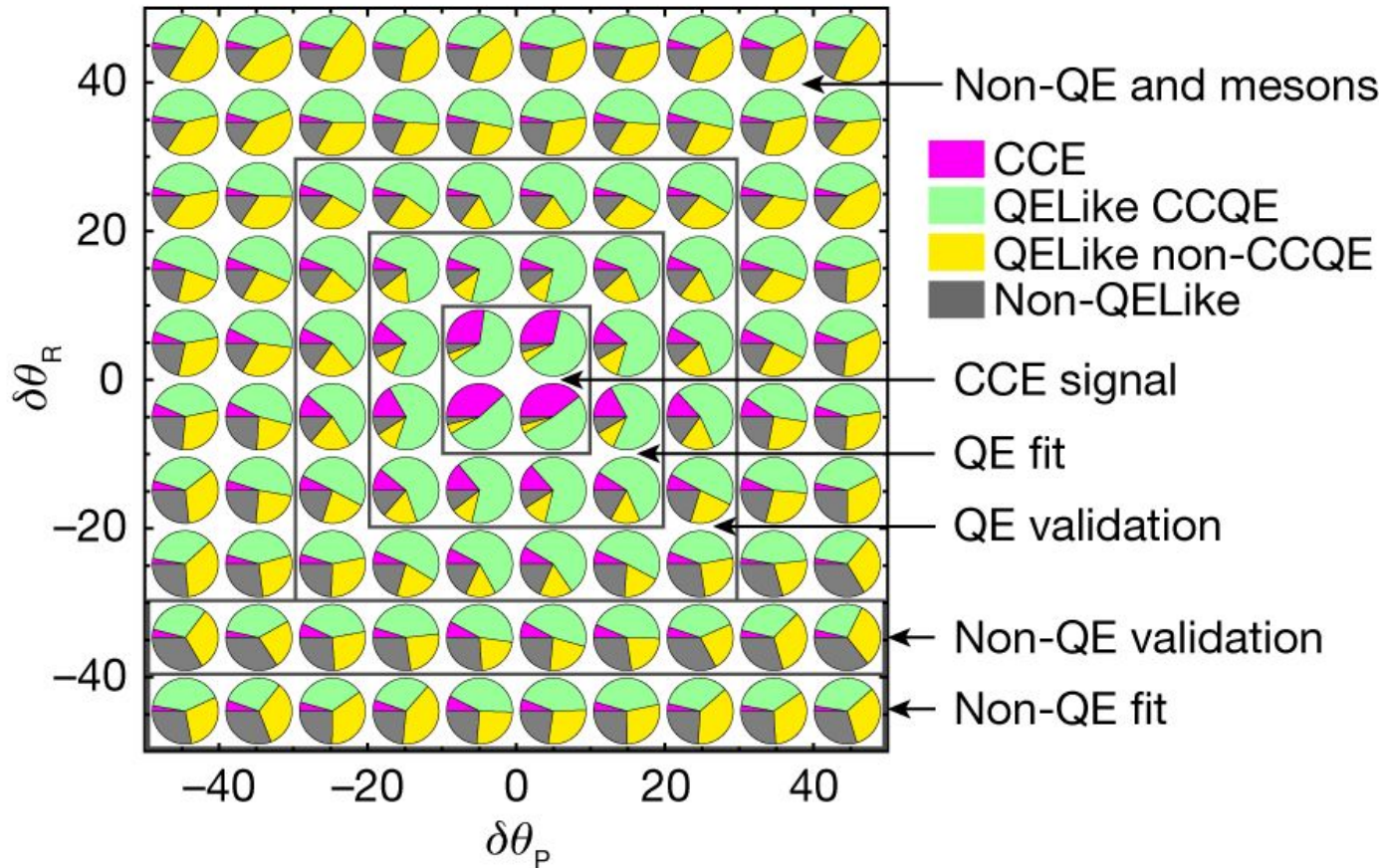


# Background constraint via sidebands



- The events are divided into two samples:
  - those that pass the full event selection (“single-blob”)
  - those that fail because they have more than one significant blob (“multi-blob”)
  - This latter sample is enriched in Non-QELike events

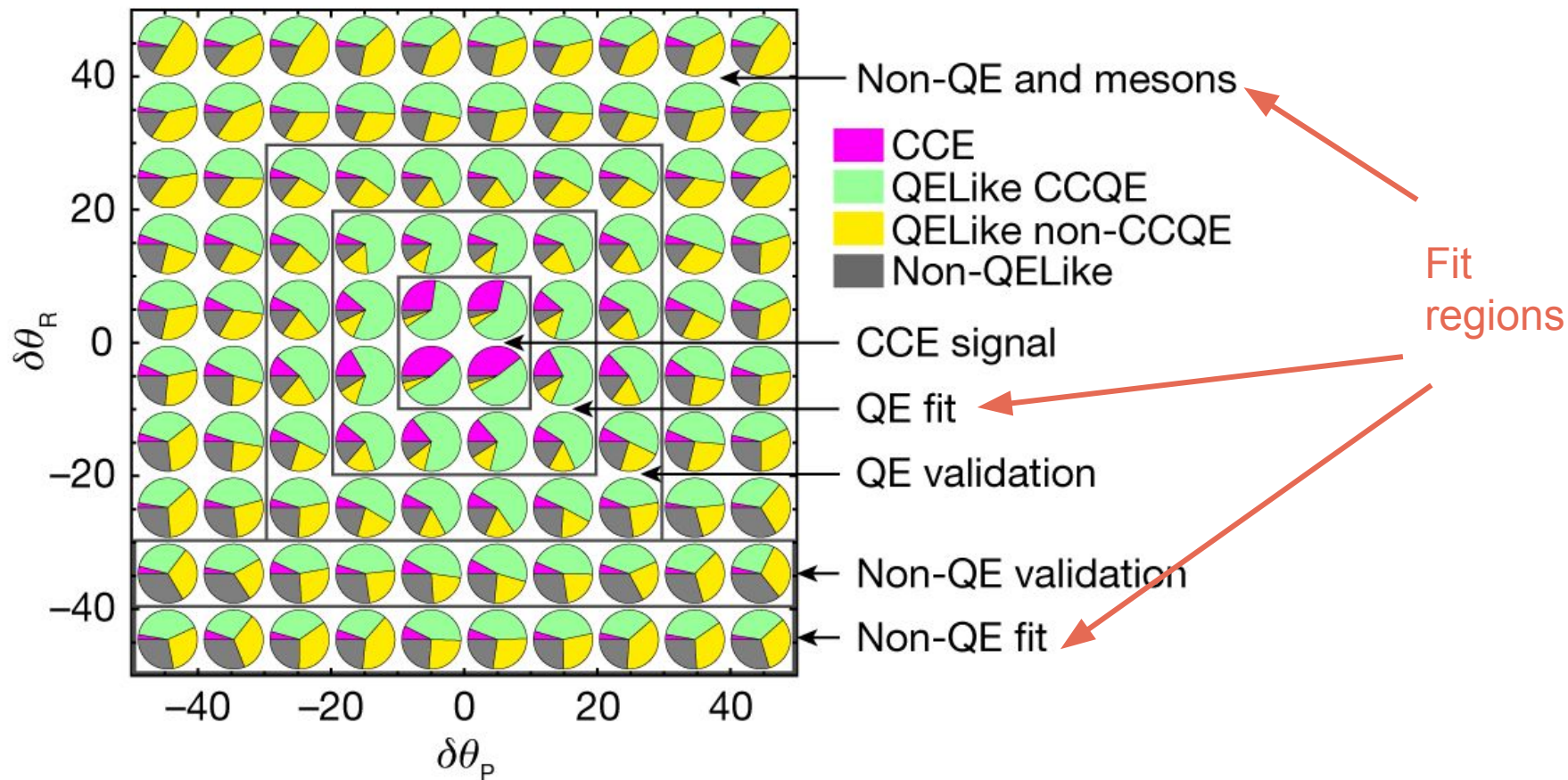
# Background constraint via sidebands



- QE-like CCQE and non-QE are fit in different regions of the  $(\delta\theta_p, \delta\theta_R)$  plane
- The results are tested in the validation regions.
- The fit is done in bins of  $Q^2$  with each bin weighted up or down. A regularization term is added to the  $\chi^2$  to ensure smoothness.



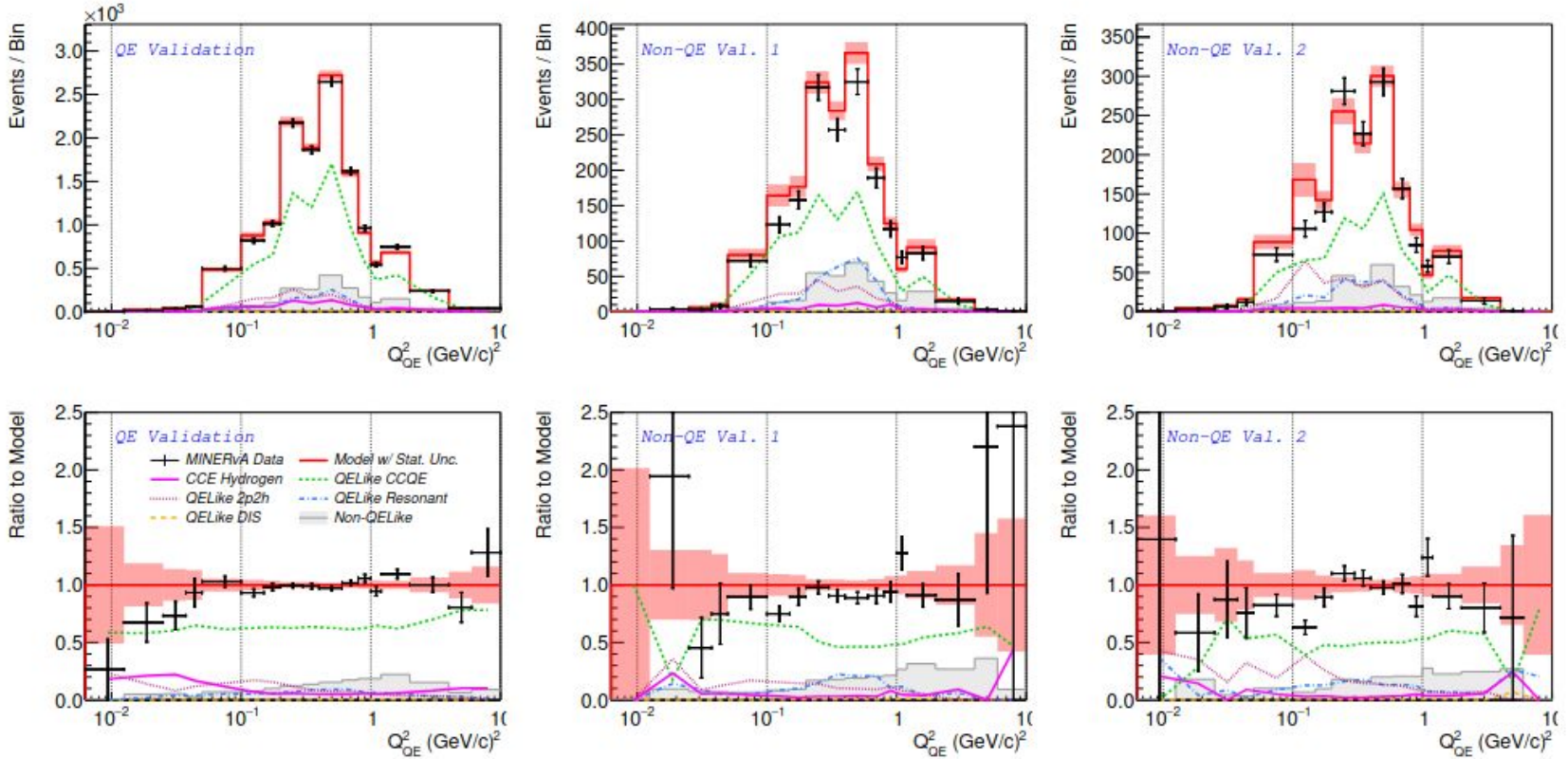
# Background constraint via sidebands



- In each Q2 bin a weight is fit for each of these categories:
  - CCQE, 2p2h, “resonant” ( $\pi$  absorbed in target nucleus)
  - Non-QELike
- The fit is done using both the single and multi-blob samples separately taking the events from the fit regions above

# Fit validation

The results of the fit in the validation regions



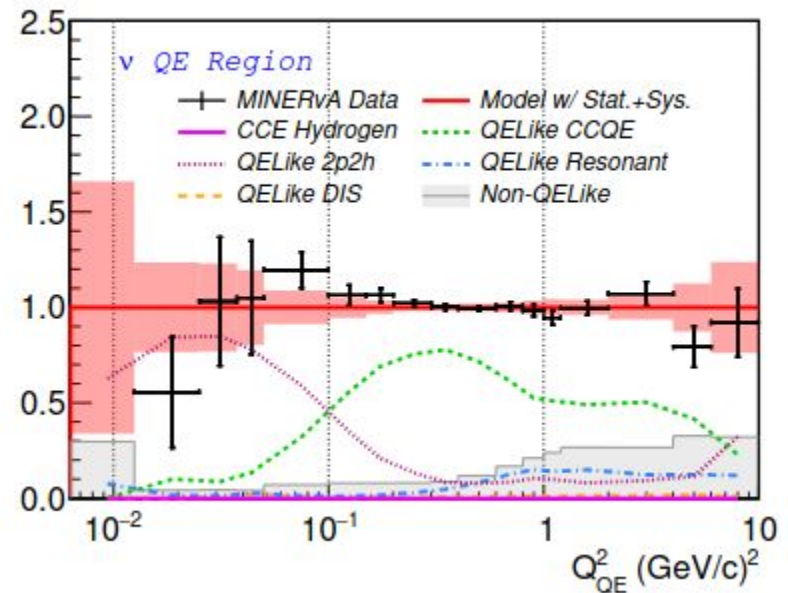
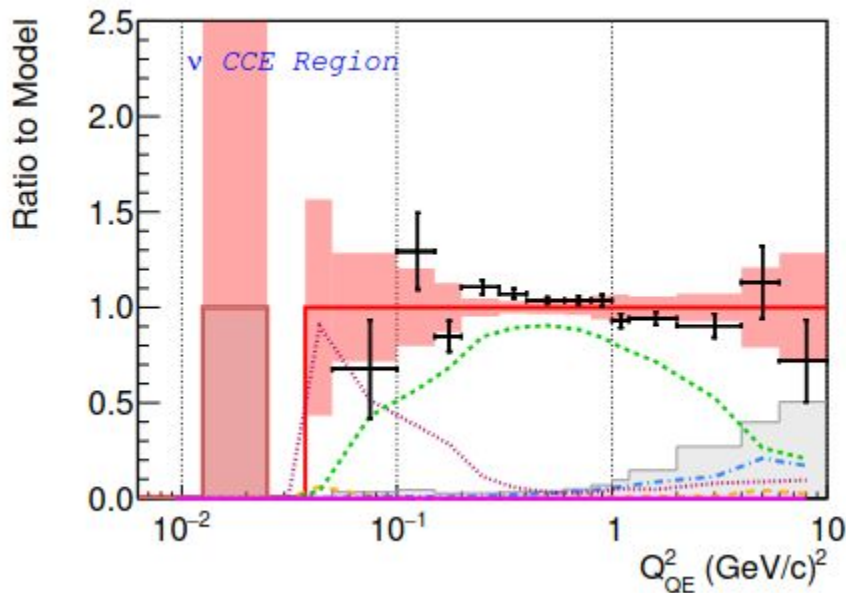
**Fig. 8** Post-fit event rate and ratio in the validation regions. The “Non-QE validation” region is shown separated into two sub-regions. “Non-QE Val. 1” spans  $|\delta\theta_P| < 20^\circ$ , “Non-QE Val. 2” occupies  $20^\circ < |\delta\theta_P| < 55^\circ$ . The vertical error bars around the data points and the error band around the model prediction account for 1 standard deviation due to statistical uncertainty. The CCE signal and the regions used in the background fits are shown in Fig. 2 and 3, respectively.



# Validation using the $\nu_\mu$ sample

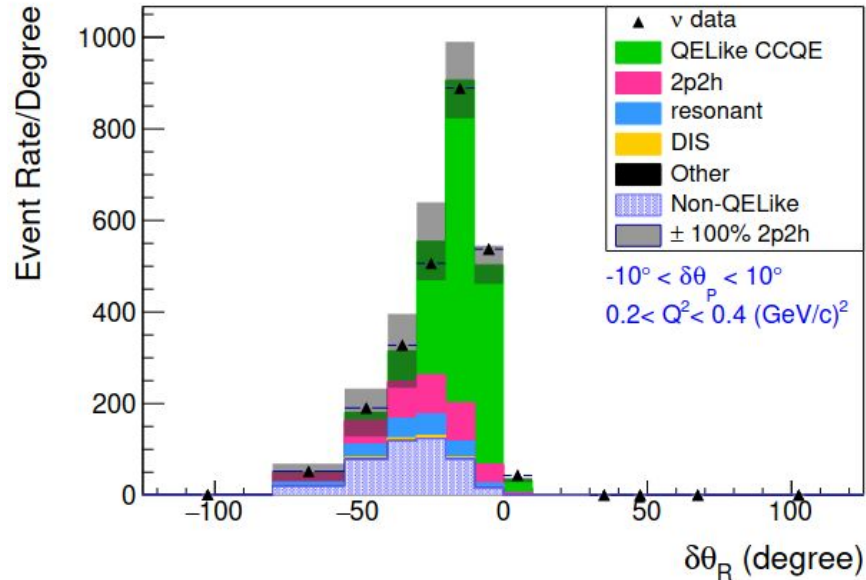
- The data here were from an anti-neutrino dominated beam
- We also have data from a neutrino dominated beam and the analogous process with a  $\mu^-$  and a proton in the f.s.
  - But, there is no elastic signal!
  - So we can look to see how well the signal region in the  $(\delta\theta_P, \delta\theta_R)$  plane is modeled!
- We use events with a tracked proton from the vertex to predict where the neutron would have gone.
- With this we can repeat the background constraint exercise.

# Validation using the $\nu_\mu$ sample



- Generally good agreement.
- Small discrepancy between 0.2 and 0.4 GeV<sup>2</sup>
- This can be accounted for by a 100% increase in the 2p2h rate

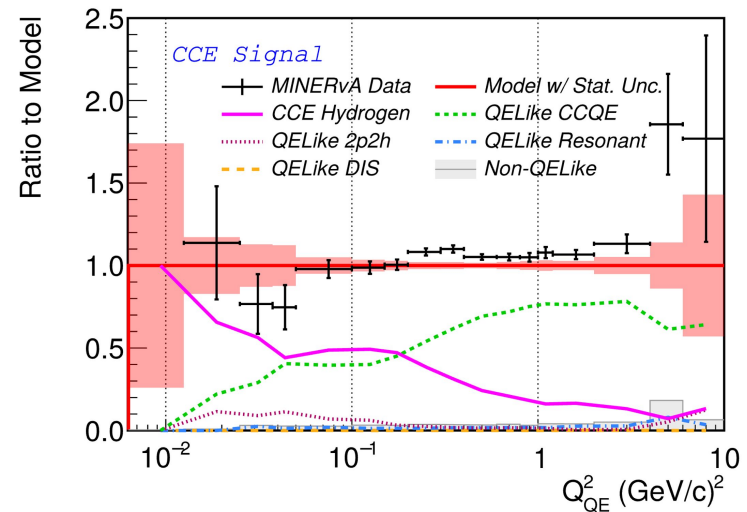
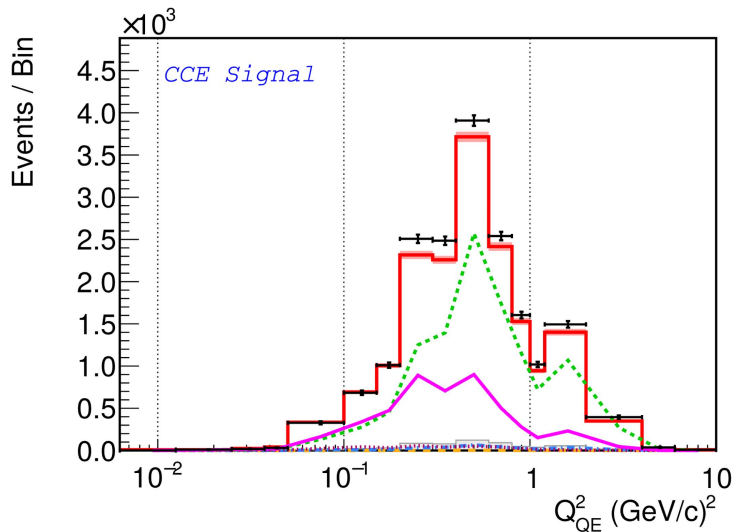
# Validation using the $\nu_\mu$ sample



- The elastic signal region is just the central two bins.
- The gray band indicates +/- 100% 2p2h (pink)

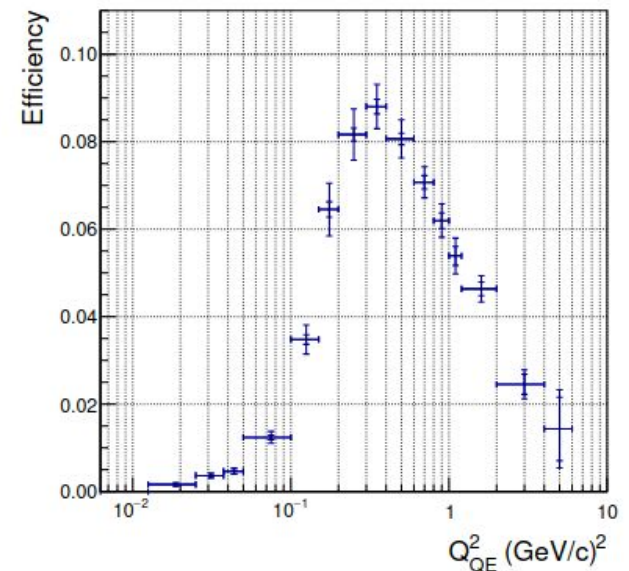
# Getting to a cross-section

- The observable is events per  $Q^2$  bin in the signal region ( $\pm |\delta\theta_{P,R}| < 10^\circ$ )
- After constraint the backgrounds are subtracted from the event sample in the signal region



# Getting to a cross-section

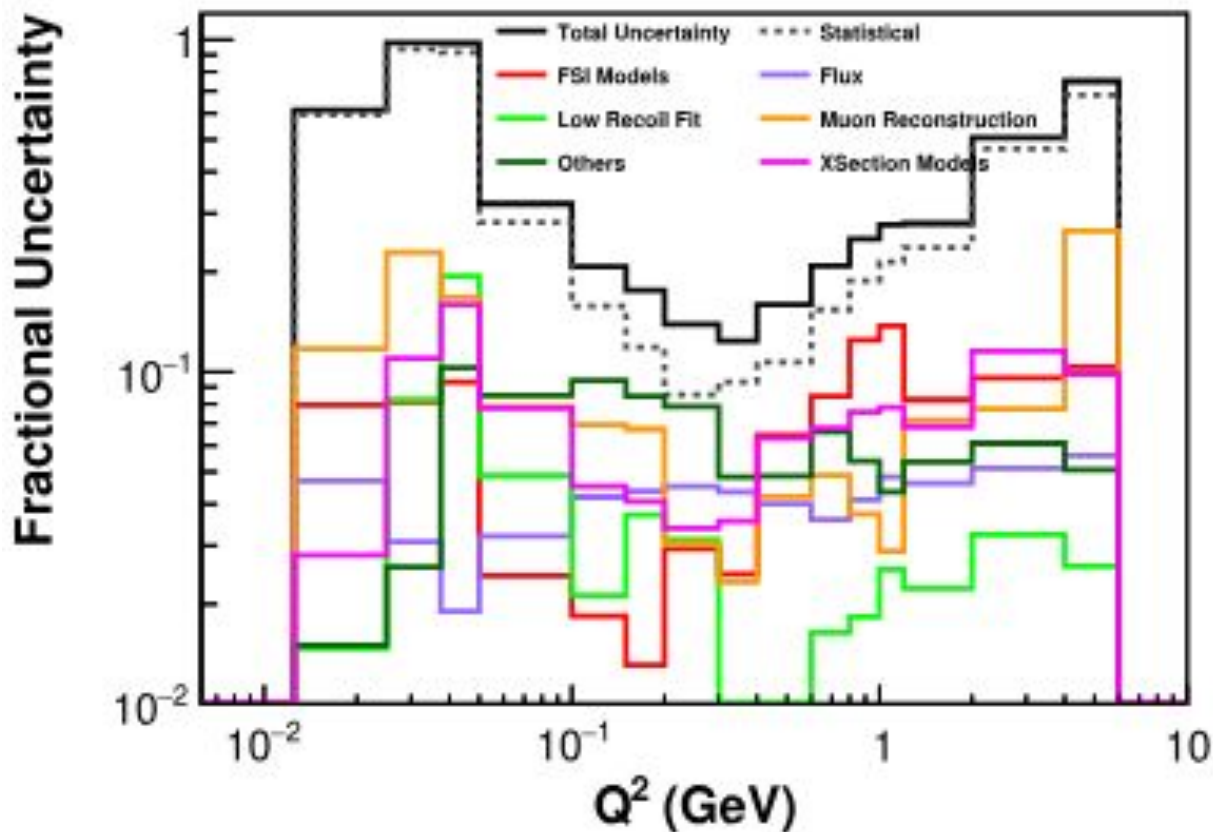
- The observable is events per  $Q^2$  bin in the signal region ( $\pm |\delta\theta_{P,R}| < 10^\circ$ )
- After constraint the backgrounds are subtracted from the event sample in the signal region
- The distribution is then unfolded using the d'Agostini method
- An efficiency correction is applied
- The predicted neutrino flux and target mass are incorporated to produce a cross-section





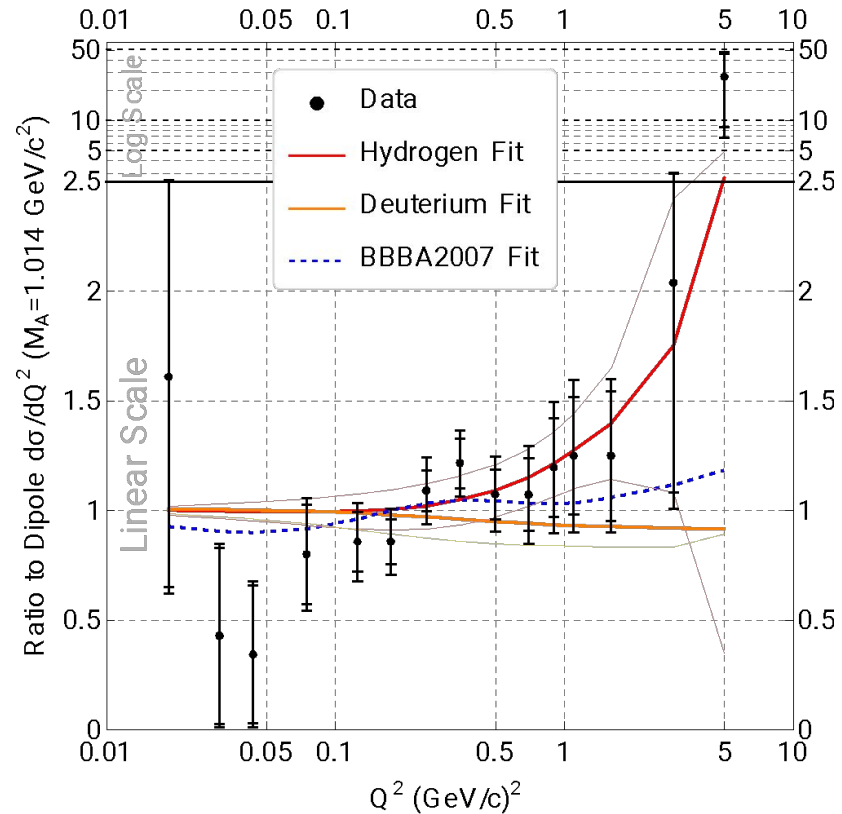
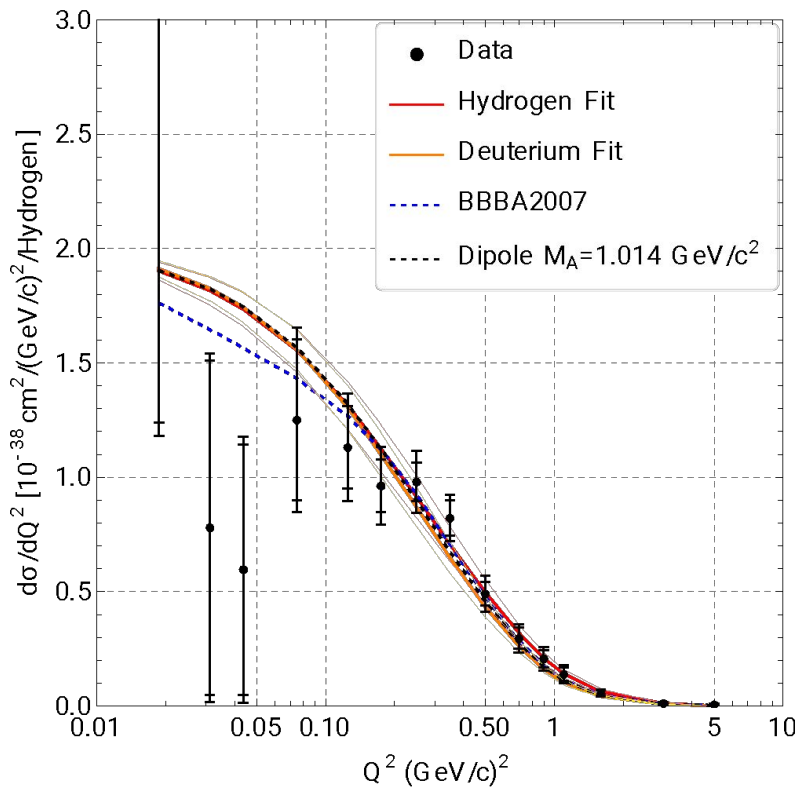
# Systematics

- Systematics are incorporated by redoing the entire analysis with one or more uncertain parameters in the MC shifted. There are several 10s of these.
- This produces a new result.
- After doing this for all uncertain parameters the variance in the result gives the uncertainty.



- Secondary neutron interactions: 4.8%
- CCQE normalization: 4.5%
- $\mu$  energy scale: 4.2% (MINOS), 3.1% (MINERvA)
- Flux: 3.9%
- neutron FSI: 3%
- 2p2h: 2.3%

# Result: $d\sigma/dQ^2$



- The H (this work) and D fits use the z-expansion formalism as in Meyer et al, Phys Rev D 93 (2016) 11, 113015
- BBBA is R. Bradford et al., Nuclear Physics B, Proceedings Supplements 159 (2006) 127–132, doi:<https://doi.org/10.1016/j.nuclphysbps.2006.08.028>.

# Result: $F_A(Q^2)$

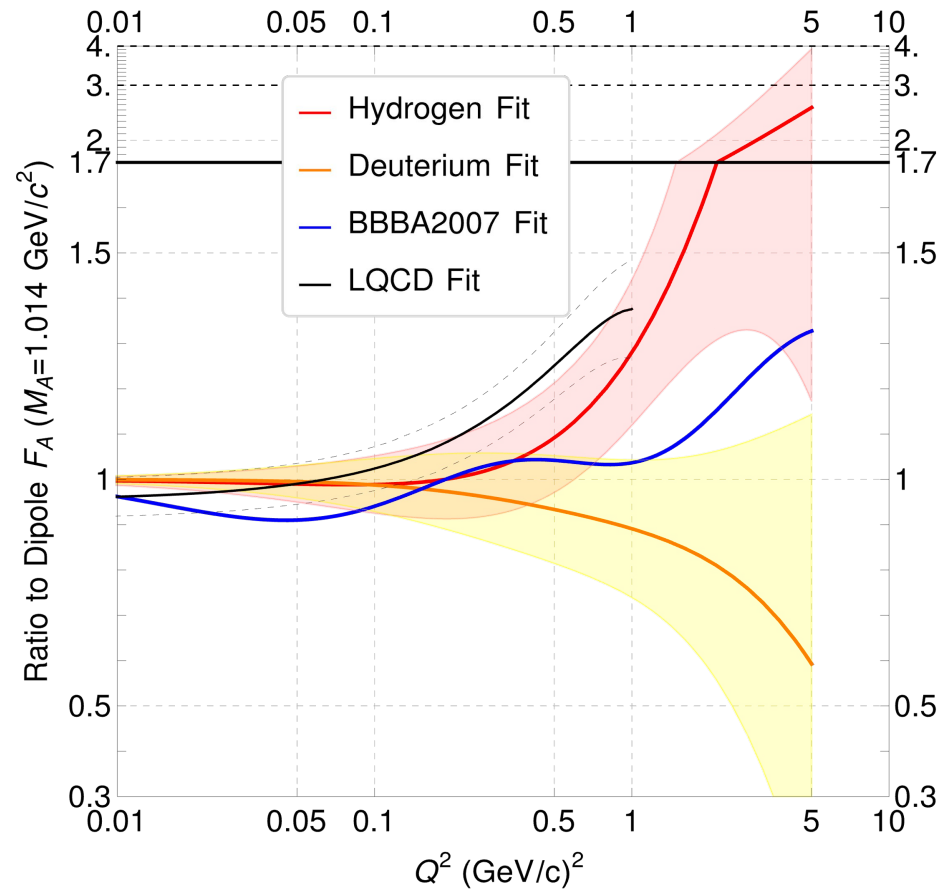
## z-expansion

$$z(Q^2, t_{\text{cut}}, t_0) = \frac{\sqrt{t_{\text{cut}} + Q^2} - \sqrt{t_{\text{cut}} - t_0}}{\sqrt{t_{\text{cut}} + Q^2} + \sqrt{t_{\text{cut}} - t_0}}$$

$$F_A(Q^2) = \sum_{k=0}^{k_{\text{max}}} a_k z^k$$

- $k_{\text{max}}=8$  here and  $t_0=-750$  MeV
- $t_{\text{cut}} = 9m_{\pi}^2$
- The chi-square includes a regularization term that penalizes higher order  $a_k$  terms to prevent overfitting
- There is a sum-rule that makes only  $k_{\text{max}}-4$  of the  $a_k$  free
- $F_A(0)=-1.2723$  from beta decay

$$F_A(Q^2) = F_A(0) \left( 1 - \frac{\langle r_A^2 \rangle}{3!} Q^2 + \frac{\langle r_A^4 \rangle}{5!} Q^4 + \dots \right) \longrightarrow r_A \equiv \sqrt{\langle r_A^2 \rangle} = 0.73(17) \text{ fm}$$



# Summary

- A major goal of MINERvA was to measure  $F_A(Q^2)$
- Thanks to the large dataset enabled by the powerful NuMI beam we were eventually able to collect 5580 signal events with a data constrained background of 12500 events.
- Roughly 10% uncertainties on the ratio to dipole in the range  $0.1 < Q^2 < 1.0 \text{ GeV}^2$  dominated by statistics but systematics starting to become noticeable.
- No new data for a while. The DUNE near detector complex's SAND detector may be able to do a similar measurement (i.e., kinematics selection) and there are proposals for H and D targets or bubble chambers or gas TPCS with (some) H in them.
- Data release at <https://minerva.fnal.gov/data-release-page/>

Questions?



# Some colors

Spirit Gold

Patina

Colonial Yellow

Vine

College Sky

Weathered Brick

Moss

College Woods

Slate

Griffin Green

Wren Twilight