

Selected

Recent advances in extracting x-dependent GPDs from lattice QCD

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POETIC 2025

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OUTLINE

A. Methods to access GPDs from lattice QCD

B. Selected results for the proton:

- twist-2 GPDs
- twist-3 GPDs

$$f_i = f_i^{(0)} + \frac{f_i^{(1)}}{Q} + \frac{f_i^{(2)}}{Q^2} \dots$$

C. Synergy with phenomenology

D. Concluding remarks

Twist-2 ($f_i^{(0)}$)

Quark \ Nucleon	U (γ^+)	L ($\gamma^+\gamma^5$)	T (σ^{+j})
U	$H(x, \xi, t)$ $E(x, \xi, t)$ unpolarized		
L		$\widetilde{H}(x, \xi, t)$ $\widetilde{E}(x, \xi, t)$ helicity	
T			H_T, E_T $\widetilde{H}_T, \widetilde{E}_T$ transversity

(Selected) Twist-3 ($f_i^{(1)}$)

Quark \ Nucleon	\mathcal{O}	γ^j	$\gamma^j \gamma^5$	σ^{jk}
U		G_1, G_2 G_3, G_4		
L			$\widetilde{G}_1, \widetilde{G}_2$ $\widetilde{G}_3, \widetilde{G}_4$	
T				$H'_2(x, \xi, t)$ $E'_2(x, \xi, t)$

OUTLINE

Theoretical/Technical
slide warning



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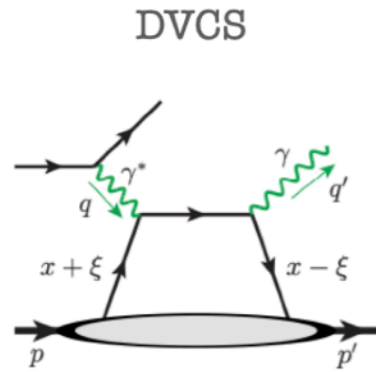
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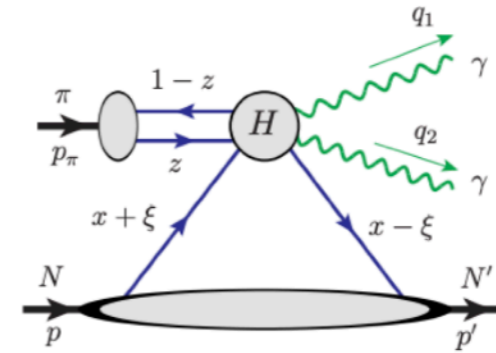
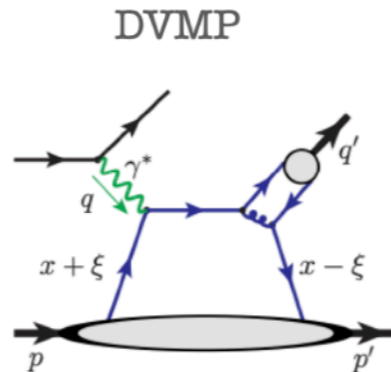
Generalized Parton Distributions

★ GPDs may be accessed via exclusive reactions (DVCS, DVMP)



[X.-D. Ji, PRD 55, 7114 (1997)]

★ exclusive pion-nucleon diffractive production of a γ pair of high p_{\perp}



[J. Qiu et al, arXiv:2205.07846]

Talk by J.W. Qiu

★ GPDs are not well-constrained experimentally:

- **x-dependence extraction is not direct. DVCS amplitude:** $\mathcal{H} = \int_{-1}^{+1} \frac{H(x, \xi, t)}{x - \xi + i\epsilon} dx$

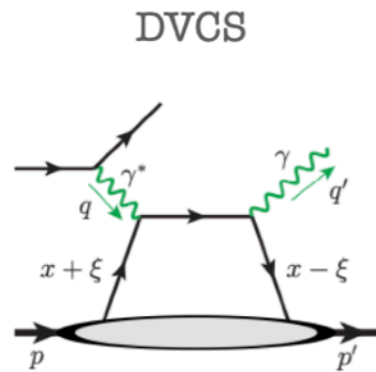
(SDHEP [J. Qiu et al, arXiv:2205.07846] gives access to x)

- independent measurements to disentangle GPDs
- GPDs phenomenology more complicated than PDFs (multi-dimensionality)
- and more challenges ...

Theoretical issues discussed by A. Freese

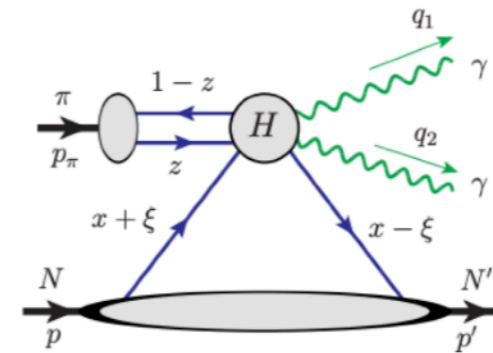
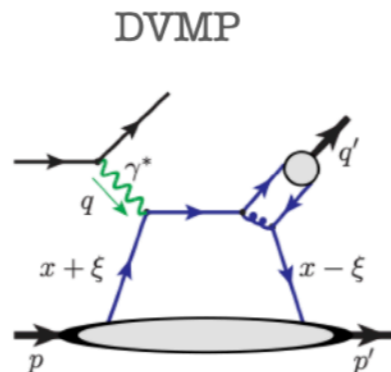
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Theoretical issues discussed by A. Freese

★ Essential to complement the knowledge on GPD from lattice QCD

★ Lattice data may be incorporated in global analysis of experimental data and may influence parametrization of t and ξ dependence

Accessing information on PDFs/GPDs



★PDFs parameterized via matrix elements of nonlocal light-cone operators

$$f(x) = \frac{1}{4\pi} \int dy^- e^{-ixP^+y^-} \langle P, S | \bar{\psi}_f \gamma^+ \mathcal{W} \psi_f | P, S \rangle$$

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★ **Mellin moments**
(local OPE expansion)

$$\bar{q}\left(-\frac{1}{2}z\right) \gamma^\sigma W\left[-\frac{1}{2}z, \frac{1}{2}z\right] q\left(\frac{1}{2}z\right) = \sum_{n=0}^{\infty} \frac{1}{n!} z_{\alpha_1} \cdots z_{\alpha_n} \left[\bar{q} \gamma^\sigma \overleftrightarrow{D}^{\alpha_1} \cdots \overleftrightarrow{D}^{\alpha_n} q \right]$$

local operators

$$\langle N(P') | \mathcal{O}_V^{\mu_1 \cdots \mu_{n-1}} | N(P) \rangle \sim \sum_{\substack{i=0 \\ \text{even}}}^{n-1} \left\{ \gamma^{\mu} \Delta^{\mu_1} \cdots \Delta^{\mu_i} \bar{P}^{\mu_{i+1}} \cdots \bar{P}^{\mu_{n-1}} \left[A_{n,i}(t) - i \frac{\Delta_\alpha \sigma^{\alpha\mu}}{2m_N} \Delta^{\mu_1} \cdots \Delta^{\mu_i} \bar{P}^{\mu_{i+1}} \cdots \bar{P}^{\mu_{n-1}} B_{n,i}(t) \right] + \frac{\Delta^\mu \Delta^{\mu_1} \cdots \Delta^{\mu_{n-1}}}{m_N} C_{n,0}(\Delta^2) \Big|_{n \text{ even}} \right\} + \frac{\Delta^\mu \Delta^{\mu_1} \cdots \Delta^{\mu_{n-1}}}{m_N} C_{n,0}(\Delta^2) \Big|_{n \text{ even}} U(P)$$

Reconstruction of PDFs/GPDs very challenging

Accessing information on PDFs/GPDs



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Reconstruction of PDFs/GPDs very challenging

★ Matrix elements of nonlocal operators
(quasi-GPDs, pseudo-GPDs, ...)

$$\langle N(P_f) | \bar{\Psi}(z) \Gamma \mathcal{W}(z,0) \Psi(0) | N(P_i) \rangle_\mu$$

Nonlocal operator with Wilson line

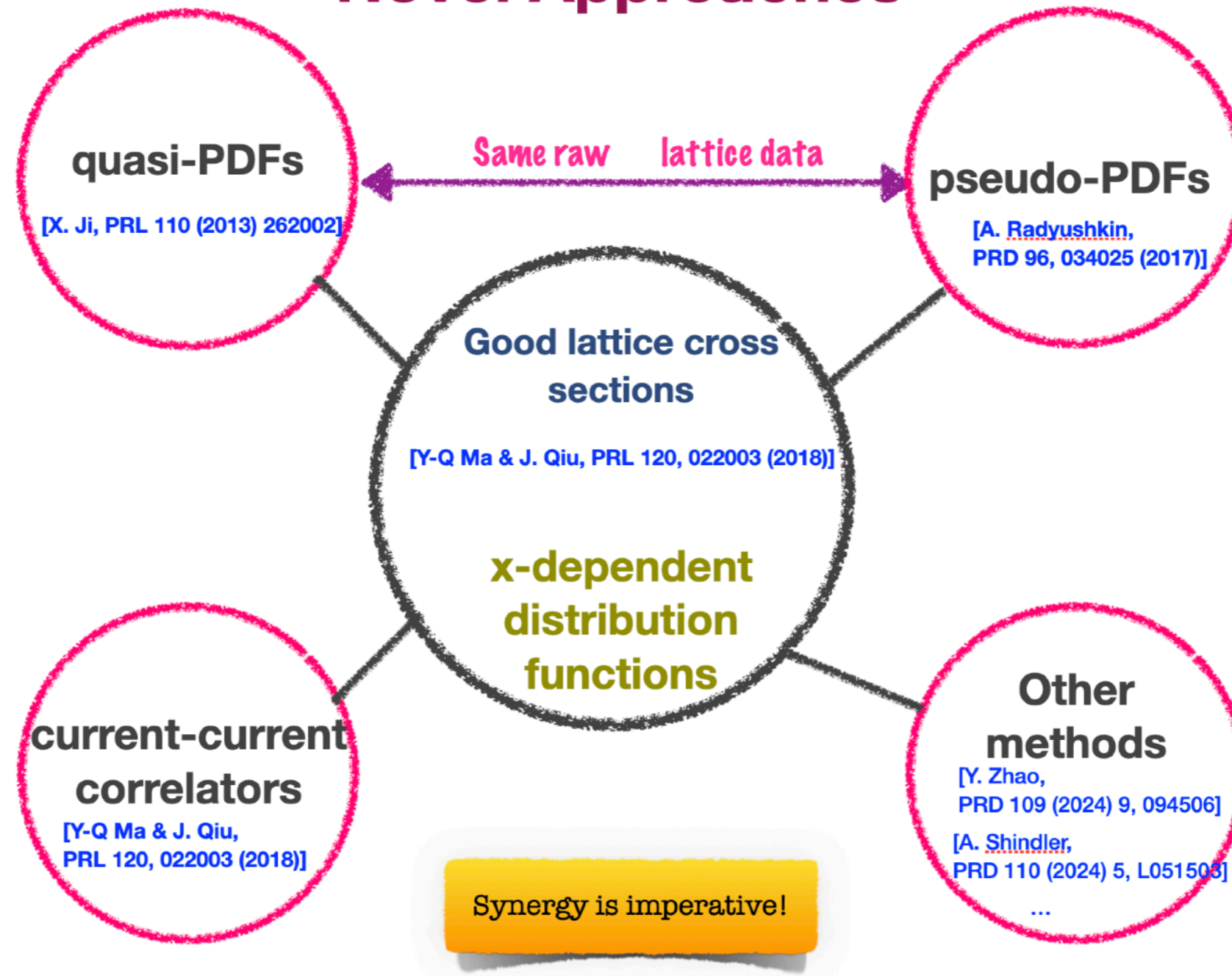
$$\langle N(P') | \mathcal{O}_V^\mu(x) | N(P) \rangle = \bar{U}(P') \left\{ \gamma^\mu H(x, \xi, t) + \frac{i\sigma^{\mu\nu} \Delta_\nu}{2m_N} E(x, \xi, t) \right\} U(P) + \text{ht},$$

$$\langle N(P') | \mathcal{O}_A^\mu(x) | N(P) \rangle = \bar{U}(P') \left\{ \gamma^\mu \gamma_5 \tilde{H}(x, \xi, t) + \frac{\gamma_5 \Delta^\mu}{2m_N} \tilde{E}(x, \xi, t) \right\} U(P) + \text{ht},$$

$$\langle N(P') | \mathcal{O}_T^{\mu\nu}(x) | N(P) \rangle = \bar{U}(P') \left\{ i\sigma^{\mu\nu} H_T(x, \xi, t) + \frac{\gamma^{[\mu} \Delta^{\nu]}}{2m_N} E_T(x, \xi, t) + \frac{\bar{P}^{[\mu} \Delta^{\nu]}}{m_N^2} \tilde{H}_T(x, \xi, t) + \frac{\gamma^{[\mu} \bar{P}^{\nu]}}{m_N} \tilde{E}_T(x, \xi, t) \right\} U(P) + \text{ht}$$

This talk

Novel Approaches



Reviews of methods and applications

- *A guide to light-cone PDFs from Lattice QCD: an overview of approaches, techniques and results*
K. Cichy & M. Constantinou (invited review) Advances in HEP 2019, 3036904, arXiv:1811.07248
- *Large Momentum Effective Theory*
X. Ji, Y.-S. Liu, Y. Liu, J.-H. Zhang, and Y. Zhao (2020), 2004.03543
- *The x-dependence of hadronic parton distributions: A review on the progress of lattice QCD*
M. Constantinou (invited review) Eur. Phys. J. A 57 (2021) 2, 77, arXiv:2010.02445

Well-studied “novel” methods for PDFs/GPDs in LQCD



Matrix elements of non-local operators (space-like separated fields)
with **boosted hadrons**

$$\mathcal{M}(P_f, P_i, z) = \langle N(P_f) | \bar{\Psi}(z) \Gamma \mathcal{W}(z, 0) \Psi(0) | N(P_i) \rangle_\mu$$

Calculation very taxing!

- length of the Wilson line (z)
 - nucleon momentum boost (P_3)
 - momentum transfer (t)
 - skewness (ξ)
- } PDFs, GPDs
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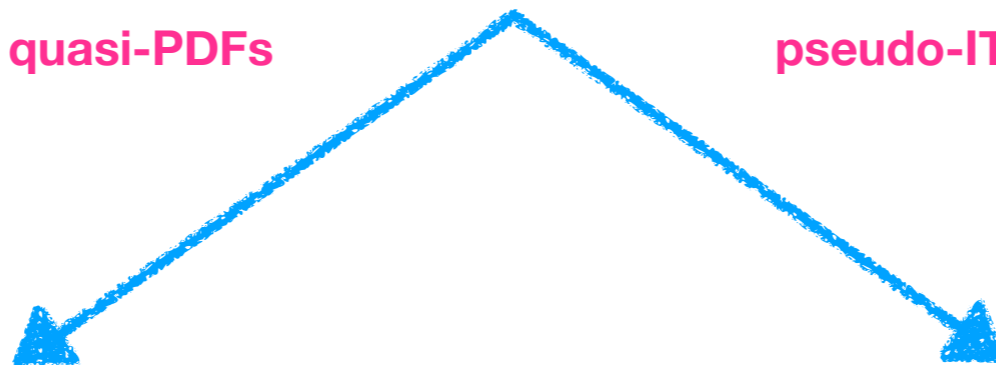
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[X. Ji, Phys. Rev. Lett. 110 (2013) 262002]
[X. Ji, Sci. China Phys. M.A. 57 (2014) 1407]

quasi-PDFs

pseudo-ITD

[A. Radyushkin, PRD 96, 034025 (2017)]



$$\tilde{q}_\Gamma^{\text{GPD}}(x, t, \xi, P_3, \mu) = \int \frac{dz}{4\pi} e^{-i x P_3 z} \mathcal{M}(P_f, P_i, z)$$

$$\mathfrak{M}(\nu, \xi, t; z_3^2) \equiv \frac{\mathcal{M}(\nu, \xi, t; z_3^2)}{\mathcal{M}(0, 0, 0; z_3^2)} \quad (\nu = z \cdot p)$$

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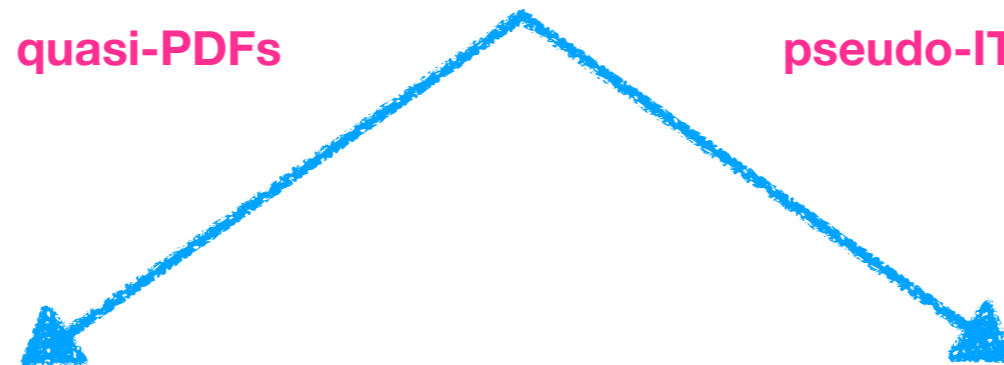
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Matching in momentum space
(Large Momentum
Effective Theory)

Matching in ν space

$$Q(\nu, \mu^2) = \int_{-1}^1 dx e^{i\nu x} q(x, \mu^2)$$

Light-cone PDFs & GPDs

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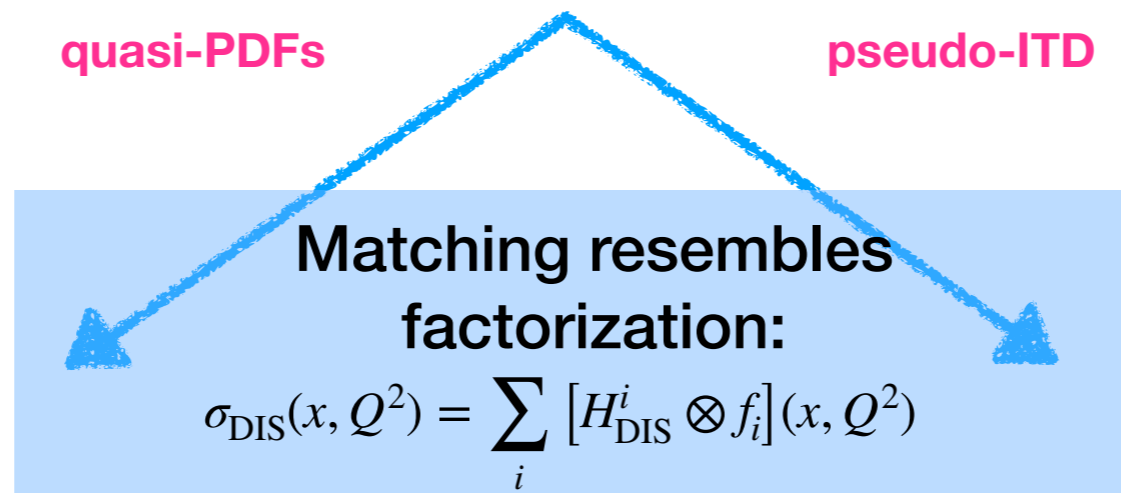
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Matching resembles factorization:

$$\sigma_{\text{DIS}}(x, Q^2) = \sum_i [H_{\text{DIS}}^i \otimes f_i](x, Q^2)$$

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GPDs on the lattice



- ★ Off-forward matrix elements of non-local light-cone operators

$$F^{[\gamma^+]}(x, \Delta; \lambda, \lambda') = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ik \cdot z} \langle p'; \lambda' | \bar{\psi}(-\frac{z}{2}) \gamma^+ \mathcal{W}(-\frac{z}{2}, \frac{z}{2}) \psi(\frac{z}{2}) | p; \lambda \rangle \Big|_{z^+=0, \vec{z}_\perp=\vec{0}_\perp}$$

- ★ Parametrization in two leading twist GPDs

$$F^{[\gamma^+]}(x, \Delta; \lambda, \lambda') = \frac{1}{2P^+} \bar{u}(p', \lambda') \left[\gamma^+ H(x, \xi, t) + \frac{i\sigma^{+\mu} \Delta_\mu}{2M} E(x, \xi, t) \right] u(p, \lambda)$$

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How can one define GPDs on a Euclidean lattice?

- ★ Potential parametrization (γ^+ inspired)

$$F^{[\gamma^3]}(x, \Delta; \lambda, \lambda'; P^3) = \frac{1}{2P^0} \bar{u}(p', \lambda') \left[\gamma^3 H_{Q(0)}(x, \xi, t; P^3) + \frac{i\sigma^{3\mu} \Delta_\mu}{2M} E_{Q(0)}(x, \xi, t; P^3) \right] u(p, \lambda)$$

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[Constantinou & Panagopoulos (2017)]

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γ^0 ideal for PDFs

GPDs on the lattice



- ★ Off-forward matrix elements of non-local light-cone operators

$$F^{[\gamma^+]}(x, \Delta; \lambda, \lambda') = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ik \cdot z} \langle p'; \lambda' | \bar{\psi}(-\frac{z}{2}) \gamma^+ \mathcal{W}(-\frac{z}{2}, \frac{z}{2}) \psi(\frac{z}{2}) | p; \lambda \rangle \Big|_{z^+=0, \vec{z}_\perp=\vec{0}_\perp}$$

- ★ Parametrization in two leading twist GPDs

$$F^{[\gamma^+]}(x, \Delta; \lambda, \lambda') = \frac{1}{2P^+} \bar{u}(p', \lambda') \left[\gamma^+ H(x, \xi, t) + \frac{i\sigma^{+\mu} \Delta_\mu}{2M} E(x, \xi, t) \right] u(p, \lambda)$$

How can one define GPDs on a Euclidean lattice?

- ★ Potential parametrization (γ^+ inspired)

$$F^{[\gamma^3]}(x, \Delta; \lambda, \lambda'; P^3) = \frac{1}{2P^0} \bar{u}(p', \lambda') \left[\text{no hand icon} t; P^3 + \frac{i\sigma^{3\mu} \Delta_\mu}{2M} E_{Q(0)}(x, \xi, t; P^3) \right] u(p, \lambda) \rightarrow \text{finite mixing with scalar [Constantinou \& Panagopoulos (2017)]}$$

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γ^0 ideal for PDFs

γ^0 parametrization is prohibitively expensive

GPDs on the lattice



★ Lorentz invariant parametrization

$$F_{\lambda,\lambda'}^\mu = \bar{u}(p', \lambda') \left[\frac{P^\mu}{M} A_1 + z^\mu M A_2 + \frac{\Delta^\mu}{M} A_3 + i\sigma^{\mu z} M A_4 + \frac{i\sigma^{\mu\Delta}}{M} A_5 + \frac{P^\mu i\sigma^{z\Delta}}{M} A_6 + \frac{z^\mu i\sigma^{z\Delta}}{M} A_7 + \frac{\Delta^\mu i\sigma^{z\Delta}}{M} A_8 \right] u(p, \lambda)$$

Goals

- ★ Extraction of standard GPDs using A_i obtained from any frame
- ★ quasi-GPDs may be redefined (Lorentz covariant) inspired by light-cone

GPDs on the lattice



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Light-cone GPDs using lattice correlators in non-symmetric frames

→ Proof-of-concept calculation ($\xi = 0$):

- symmetric frame: $\vec{p}_f^s = \vec{P} + \frac{\vec{Q}}{2}, \quad \vec{p}_i^s = \vec{P} - \frac{\vec{Q}}{2} \quad -t^s = \vec{Q}^2 = 0.69 \text{ GeV}^2$

- asymmetric frame: $\vec{p}_f^a = \vec{P}, \quad \vec{p}_i^a = \vec{P} - \vec{Q} \quad t^a = -\vec{Q}^2 + (E_f - E_i)^2 = 0.65 \text{ GeV}^2$

GPDs on the lattice



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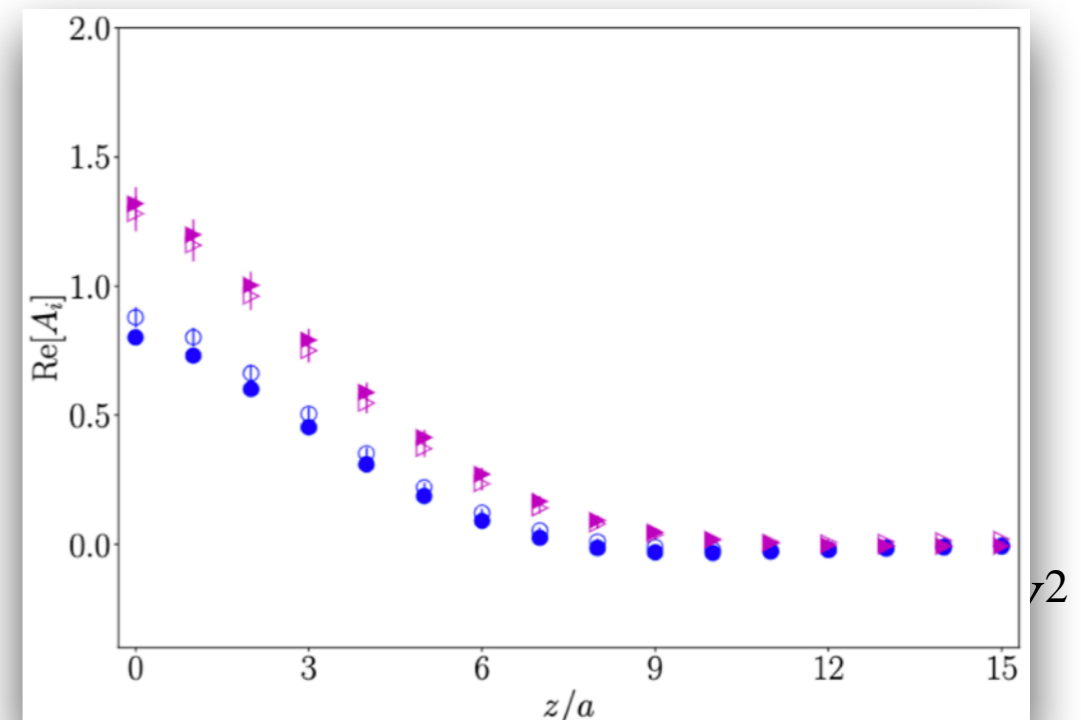
Light-cone GPDs using lattice correlators in non-symmetric frames

[S. Bhattacharya et al., PRD 106 (2022) 11, 114512]

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GPDs on the lattice



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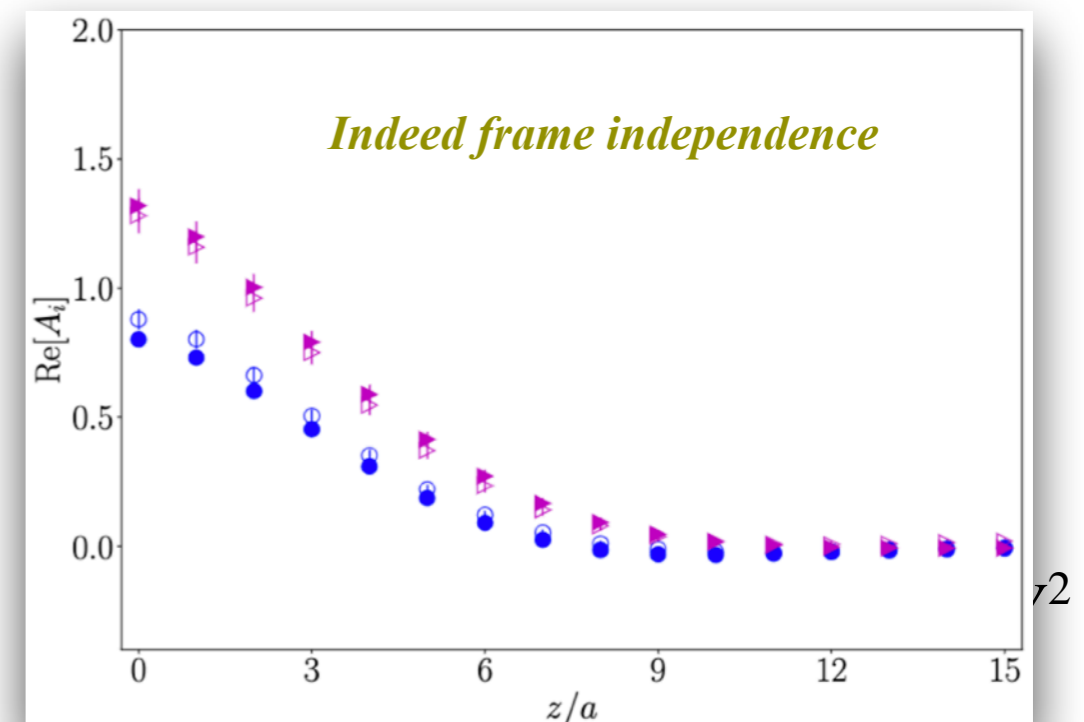
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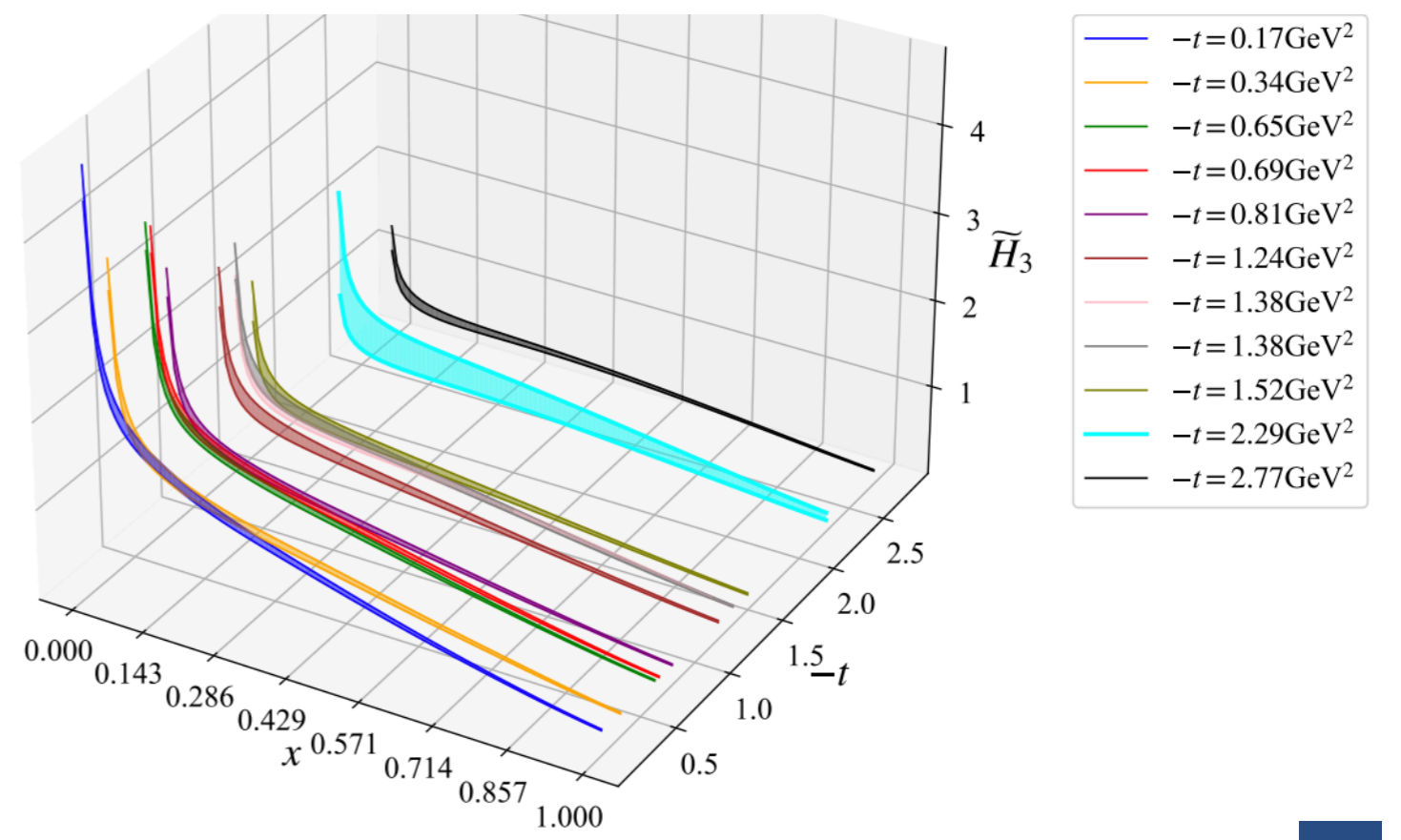
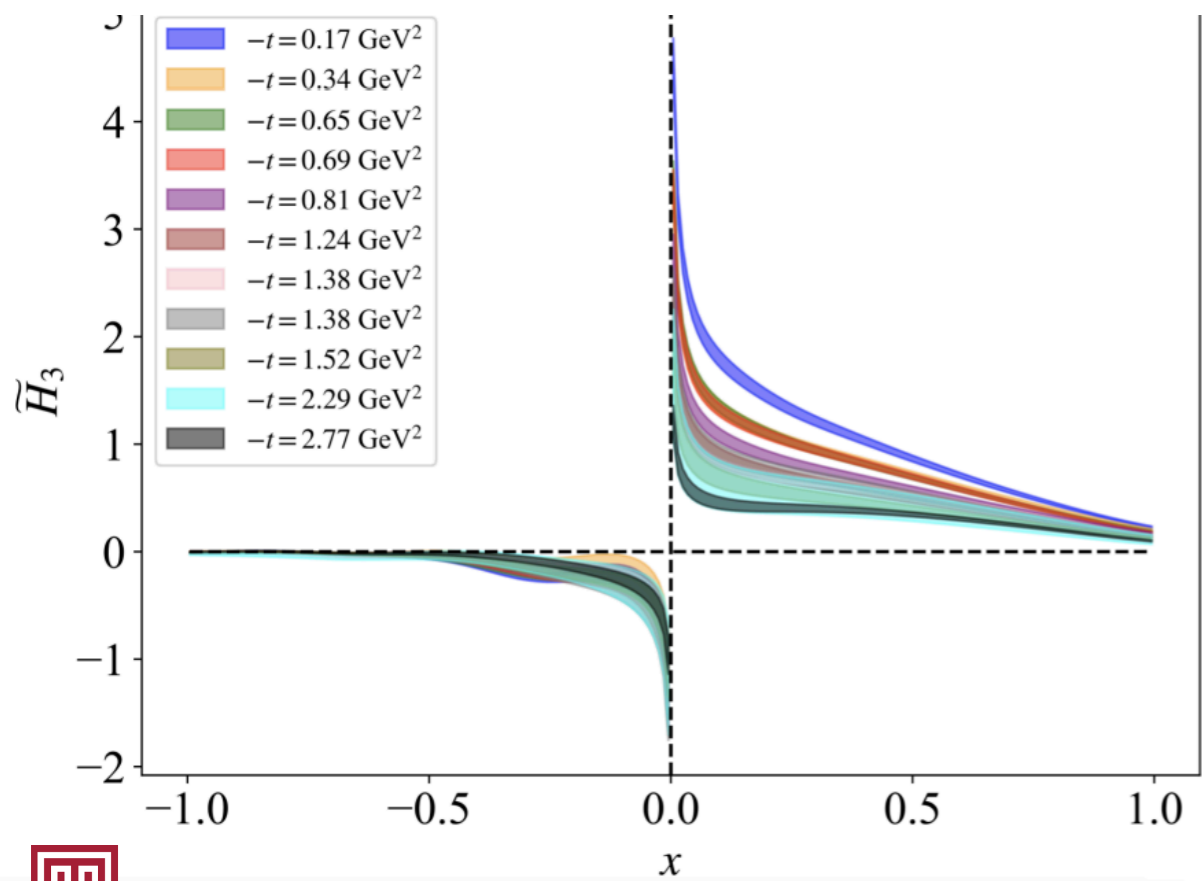
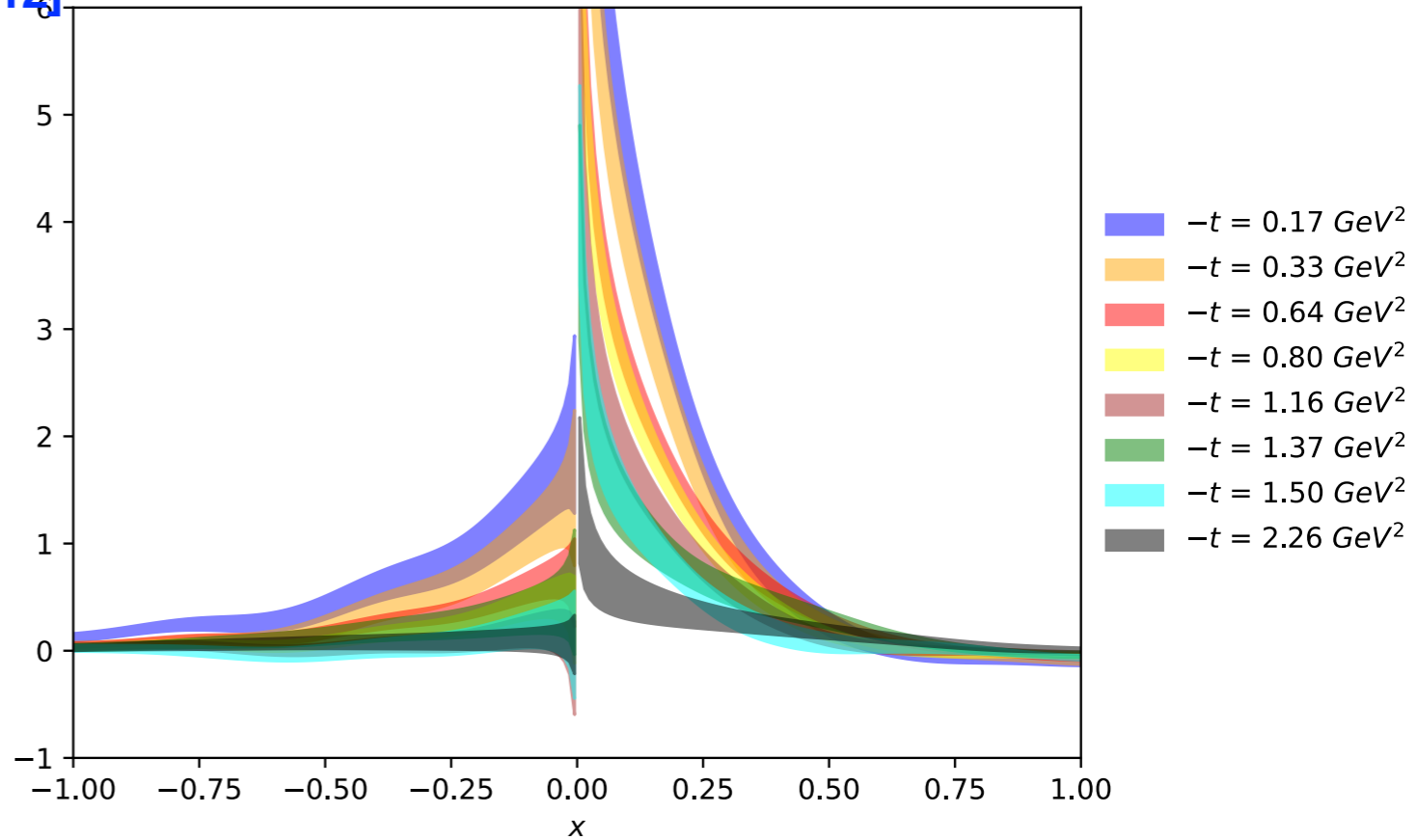
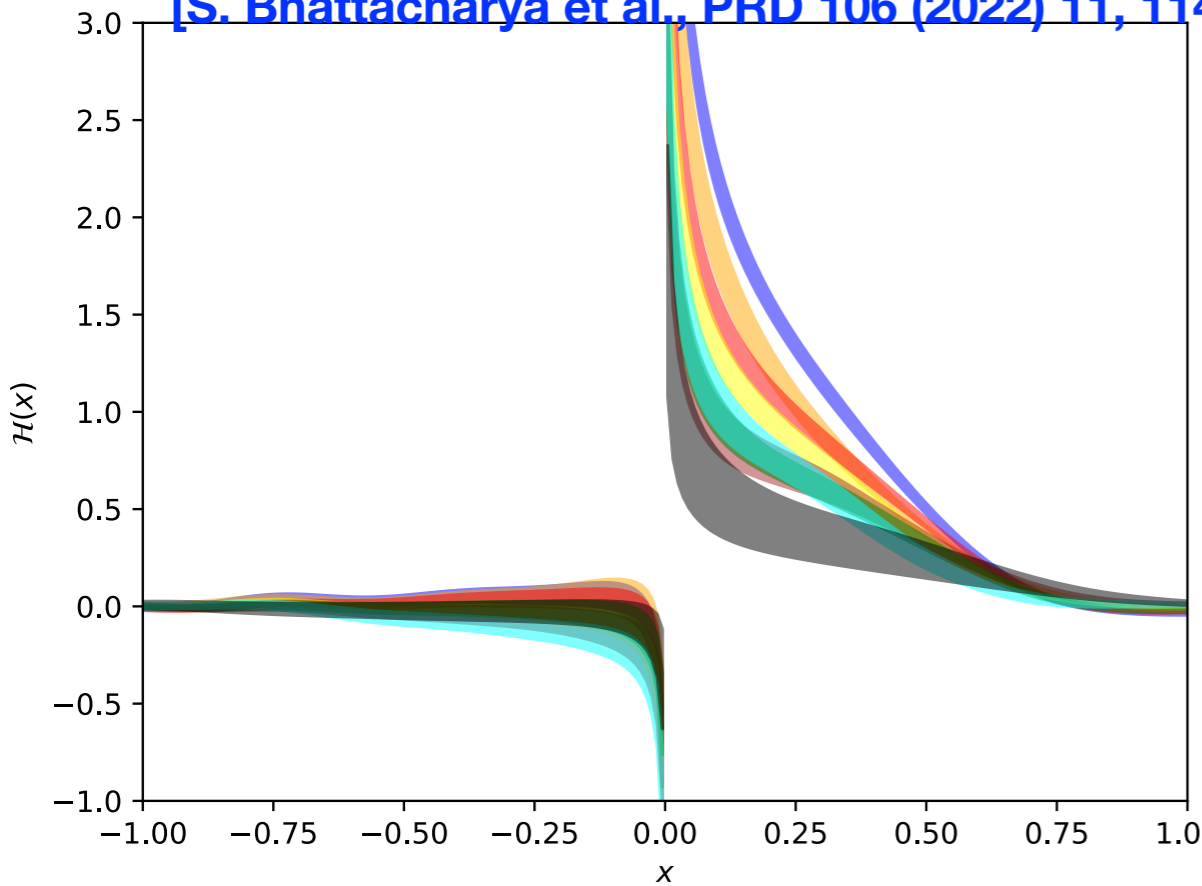
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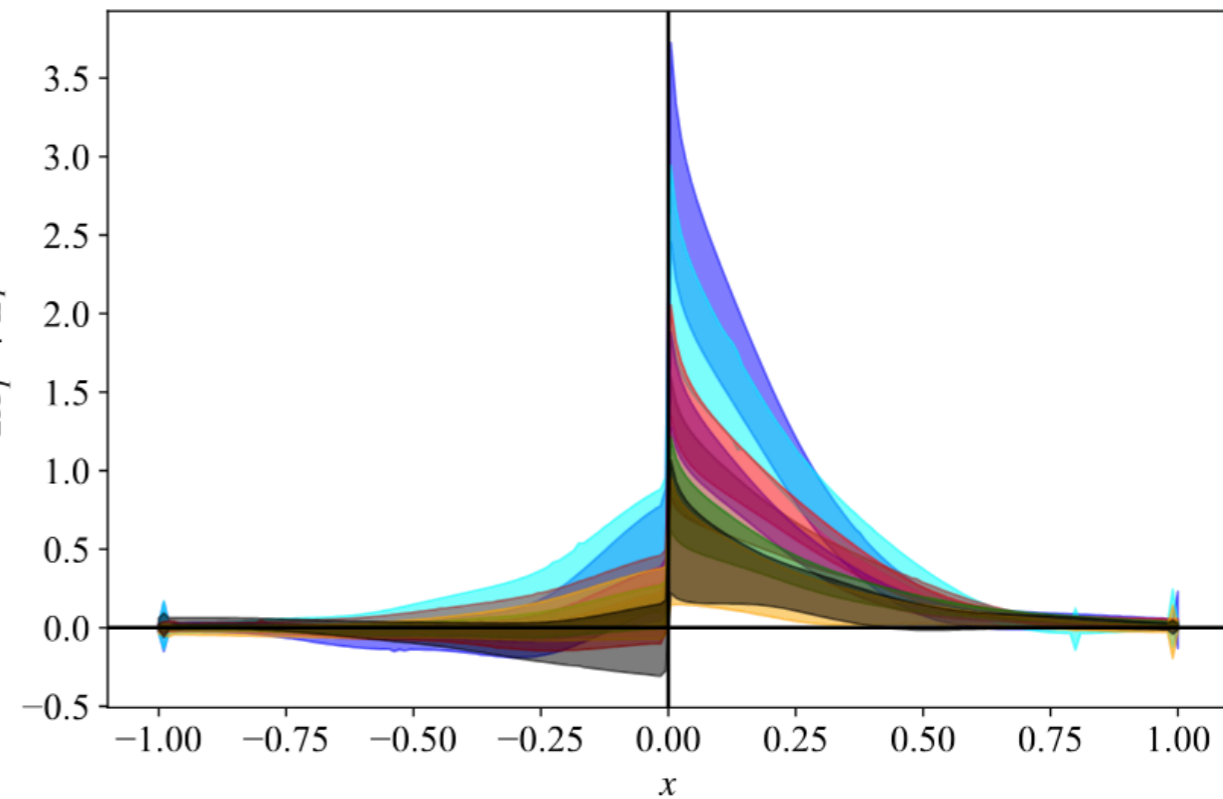
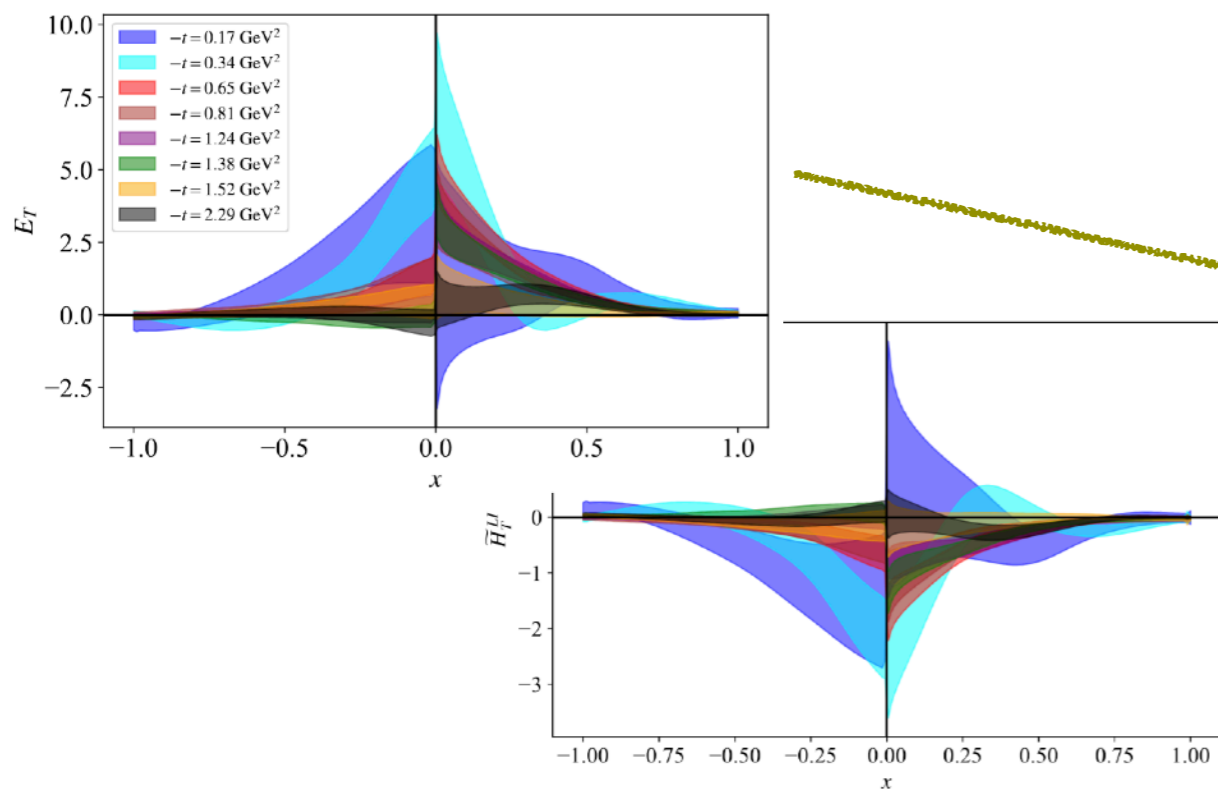
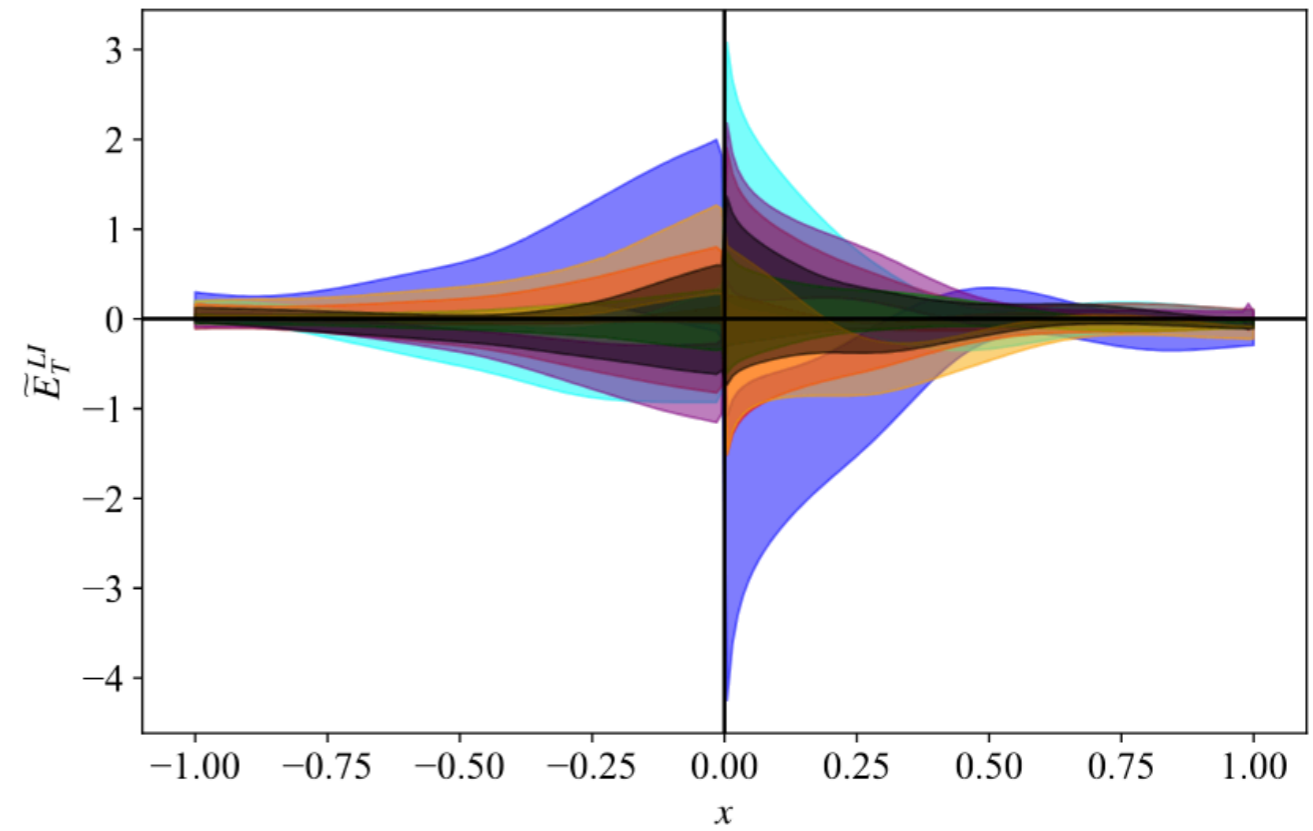
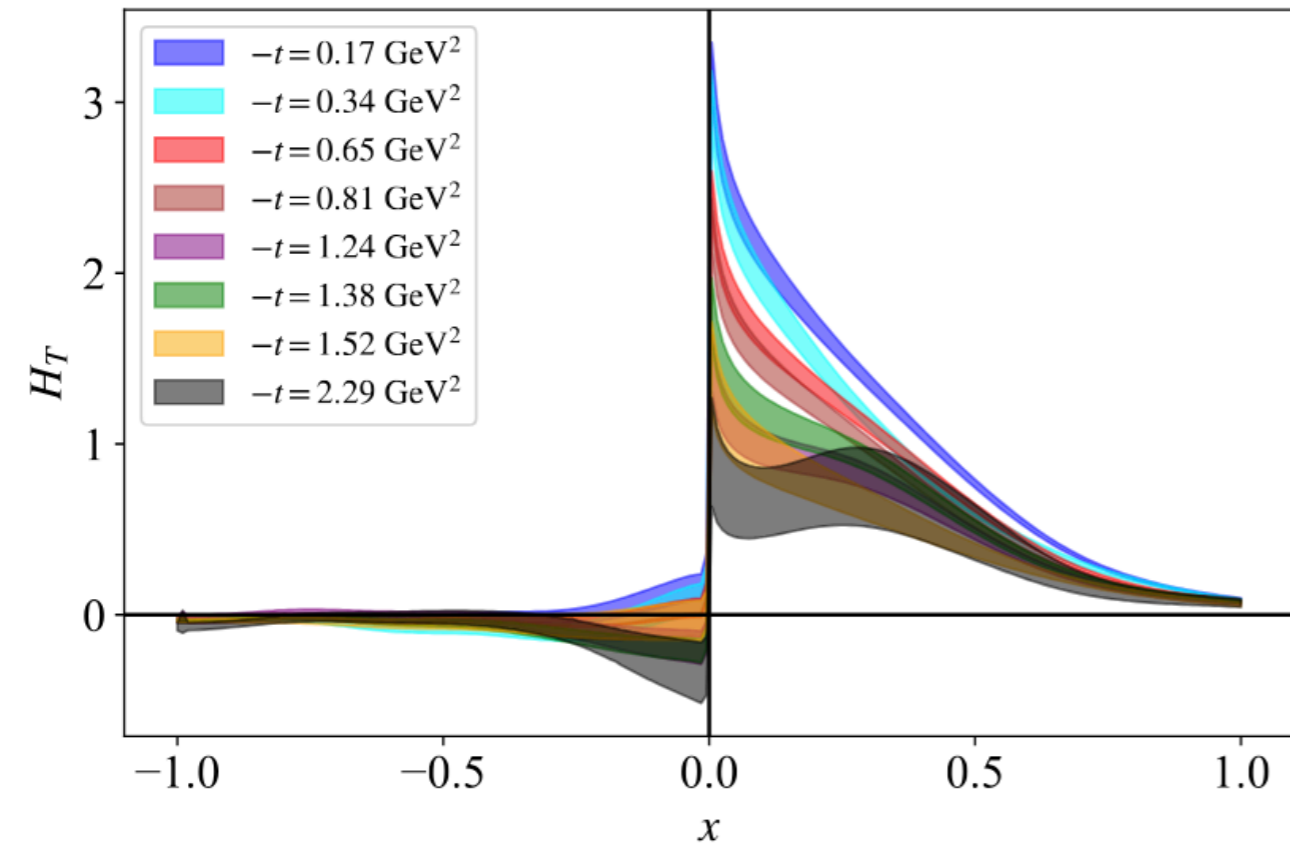


Results: unpolarized, helicity

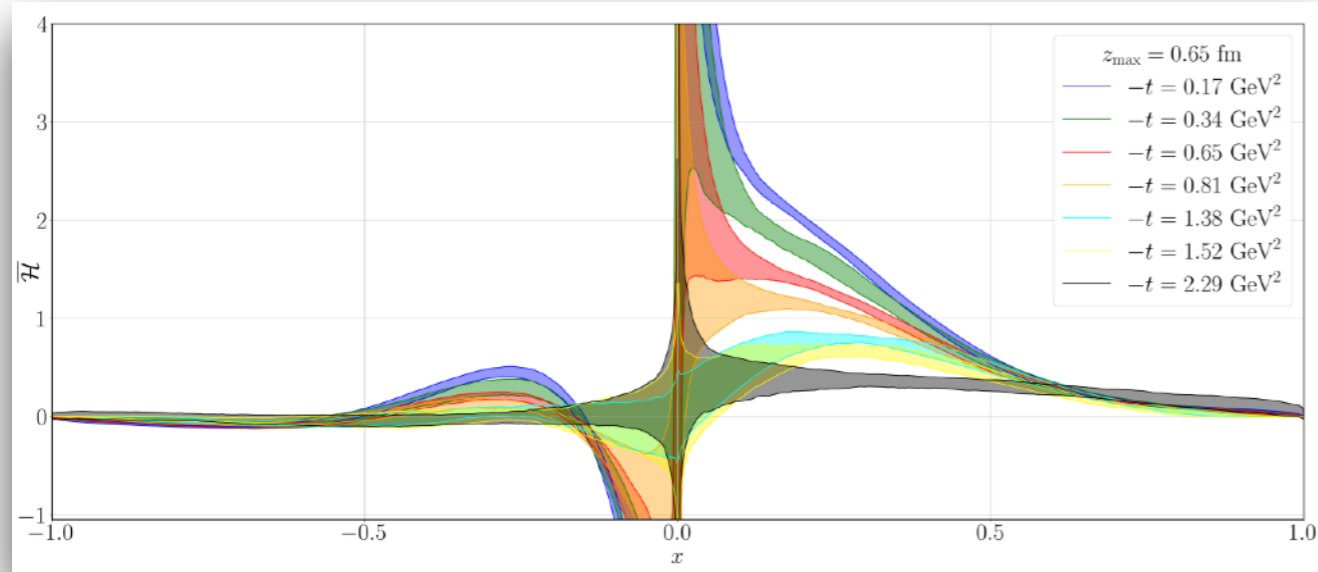
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Results: Transversity GPDs



Alternative approach: pseudo-ITD



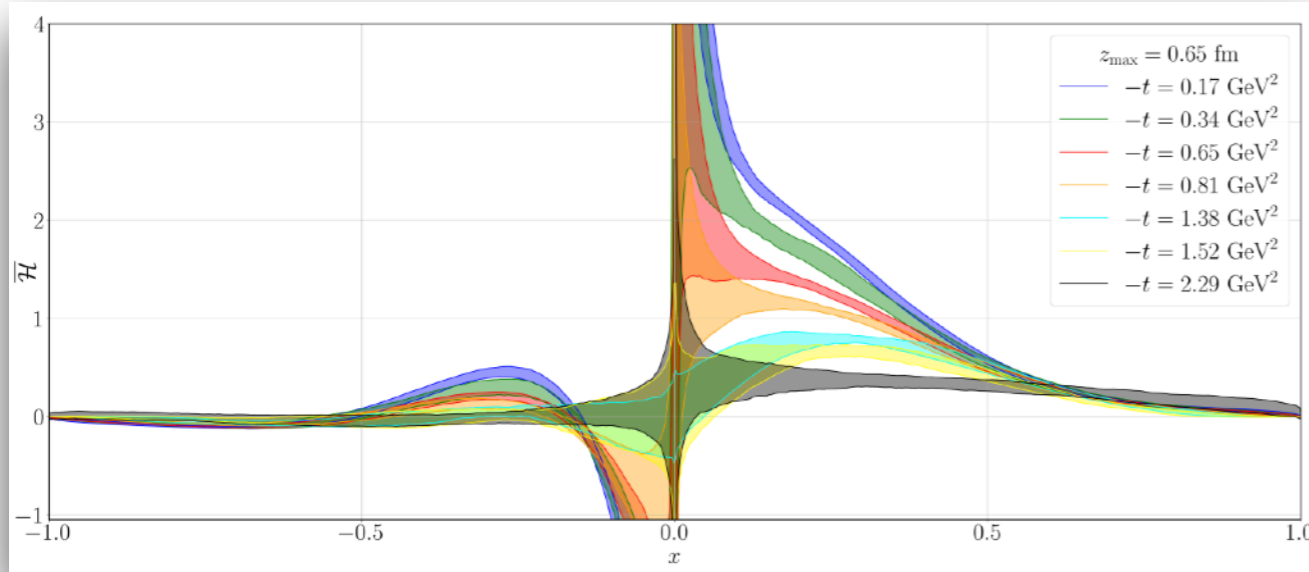
[Battacharya et al., PRD 110 (2024) 5, 054502]

Different steps between approaches:

- *renormalization*
- *x-dependence reconstruction*
- *matching formalism*

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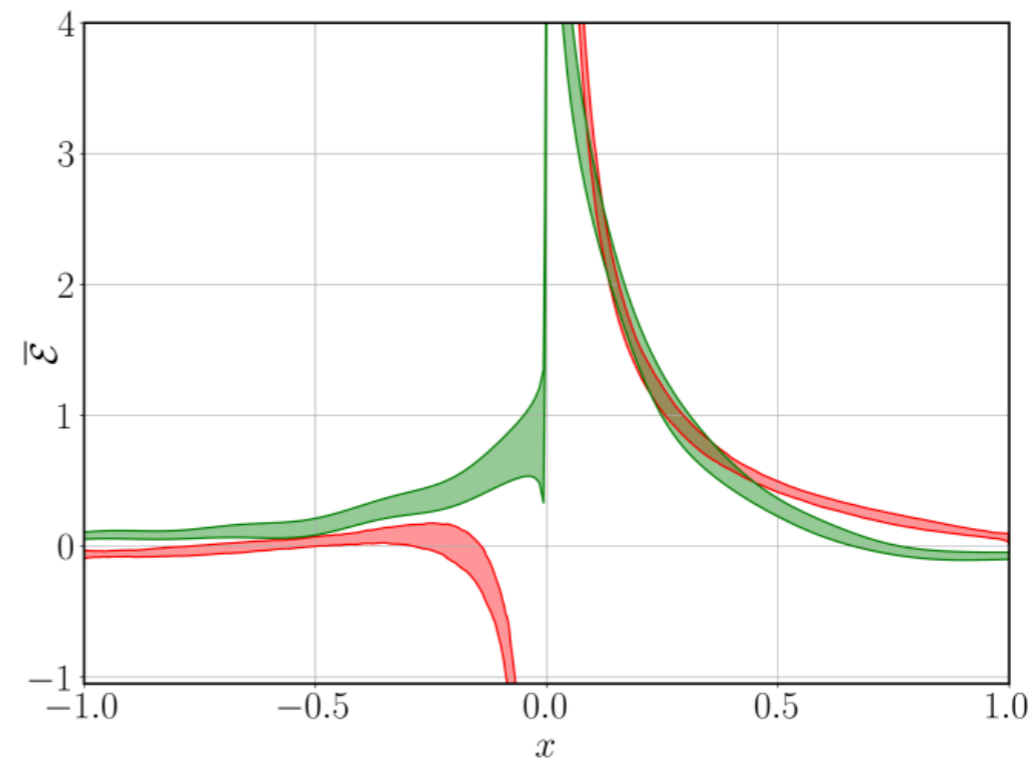
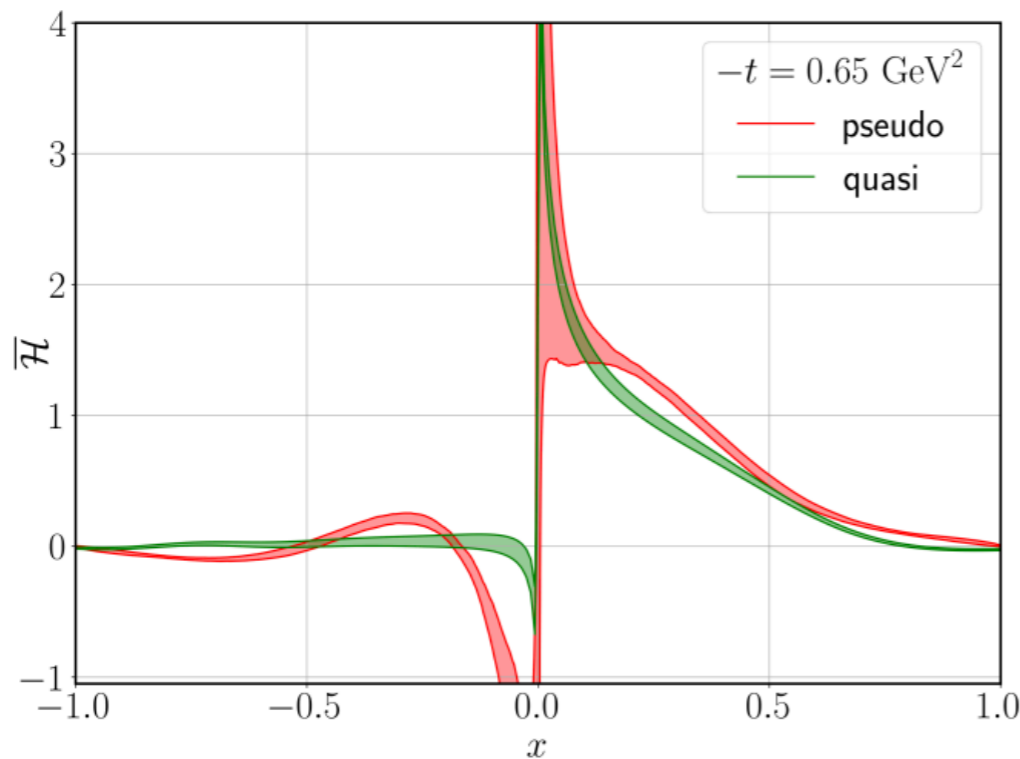
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Different steps between approaches:

- renormalization
- x -dependence reconstruction
- matching formalism

★ Comparison between methods helps assess systematic effects



- ★ $x < 0$ and small- x regions susceptible to systematic effects
- ★ Comparison only includes systematic uncertainties

Mellin moments from non-local operators

[S. Bhattacharya et al., PRD 108 (2023) 1, 014507; arXiv:2410.03539]

★ Leading-twist factorization formula

$$\mathcal{M}(z, P, \Delta) \equiv \frac{\mathcal{F}(z, P, \Delta)}{\mathcal{F}(z, P=0, \Delta=0)} = \sum_{n=0} \frac{(-izP)^n}{n!} \frac{C_n^{\overline{\text{MS}}}(\mu^2 z^2)}{C_0^{\overline{\text{MS}}}(\mu^2 z^2)} \langle x^n \rangle + \mathcal{O}(\Lambda_{\text{QCD}}^2 z^2)$$



★ Avoid power-divergent mixing of multi-derivative operators

★ Wilson coefficients known to NLO (or NNLO)

★ Both isovector and isoscalar (ignores disconnected; found to be tiny)

[C. Alexandrou et al., PRD 104 (2021) 5, 054503]

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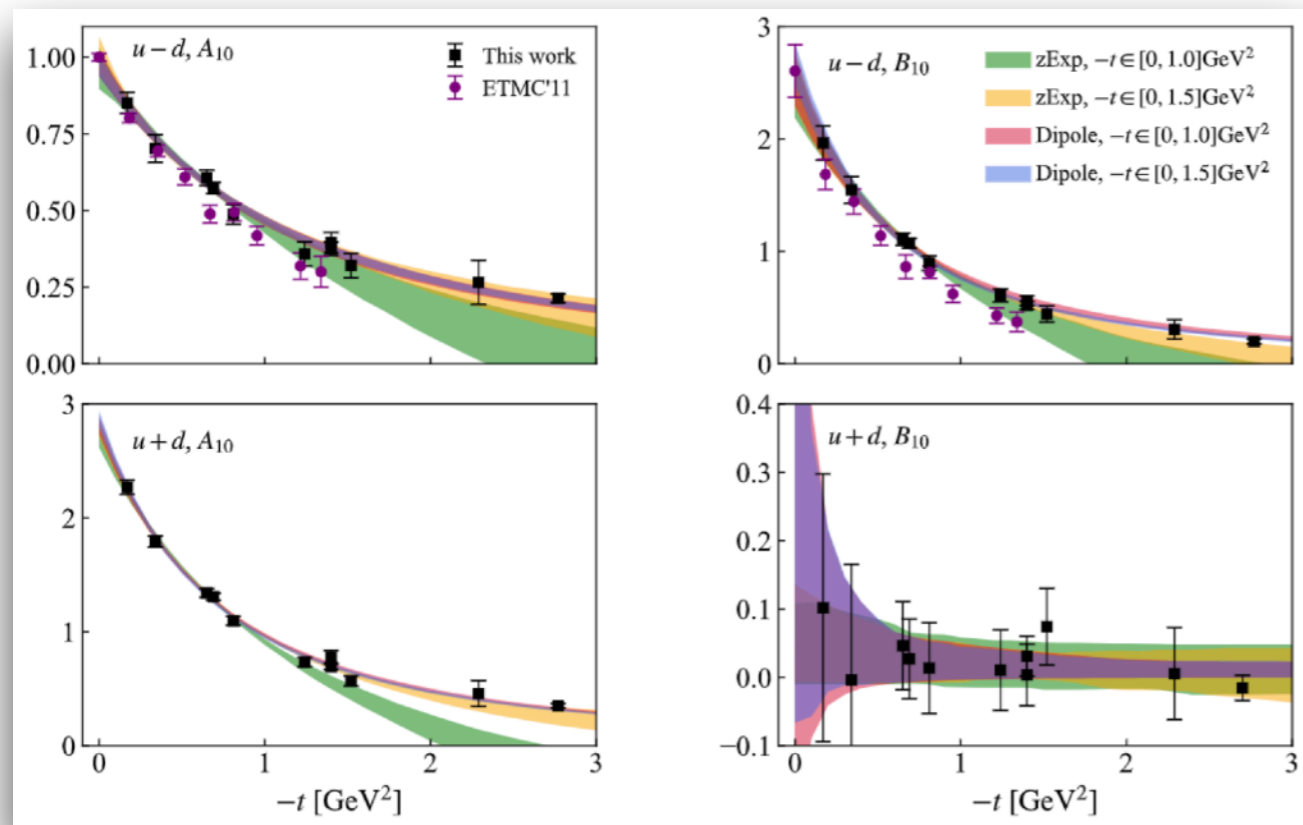
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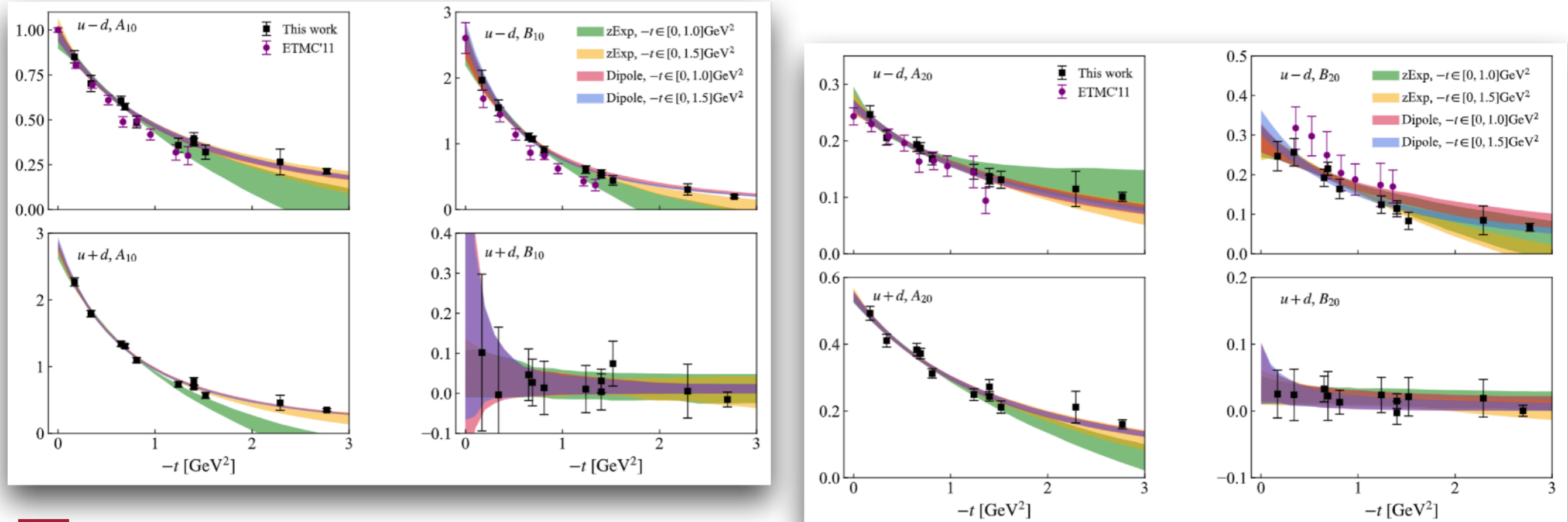
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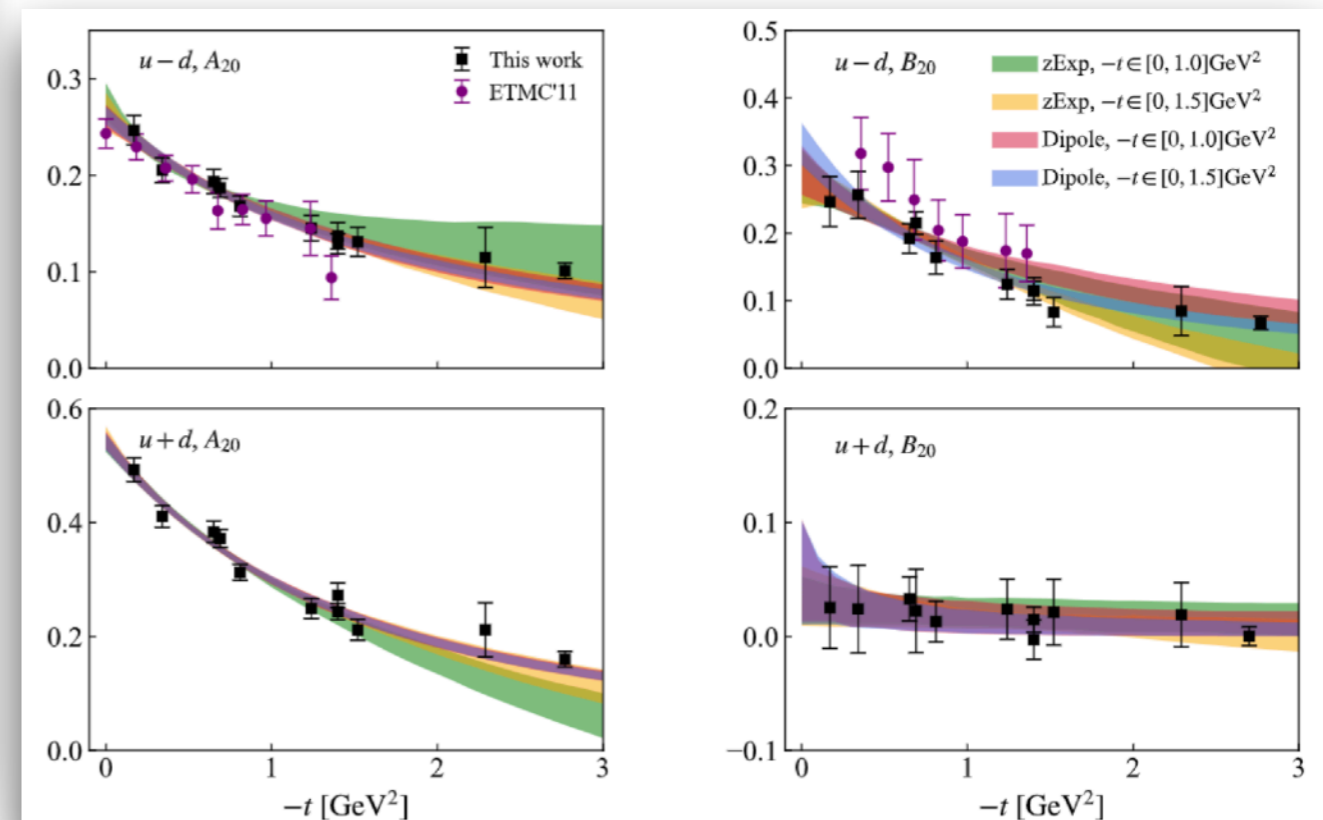
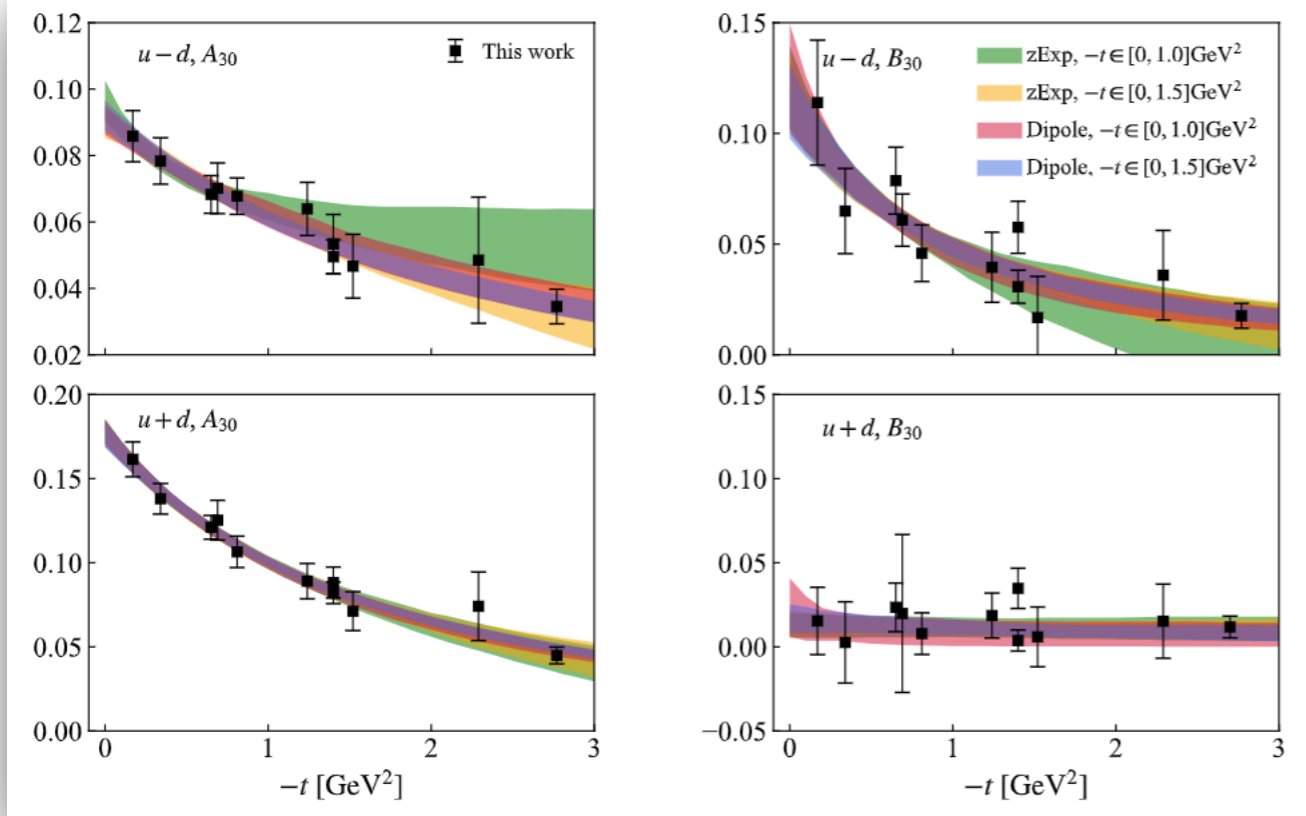
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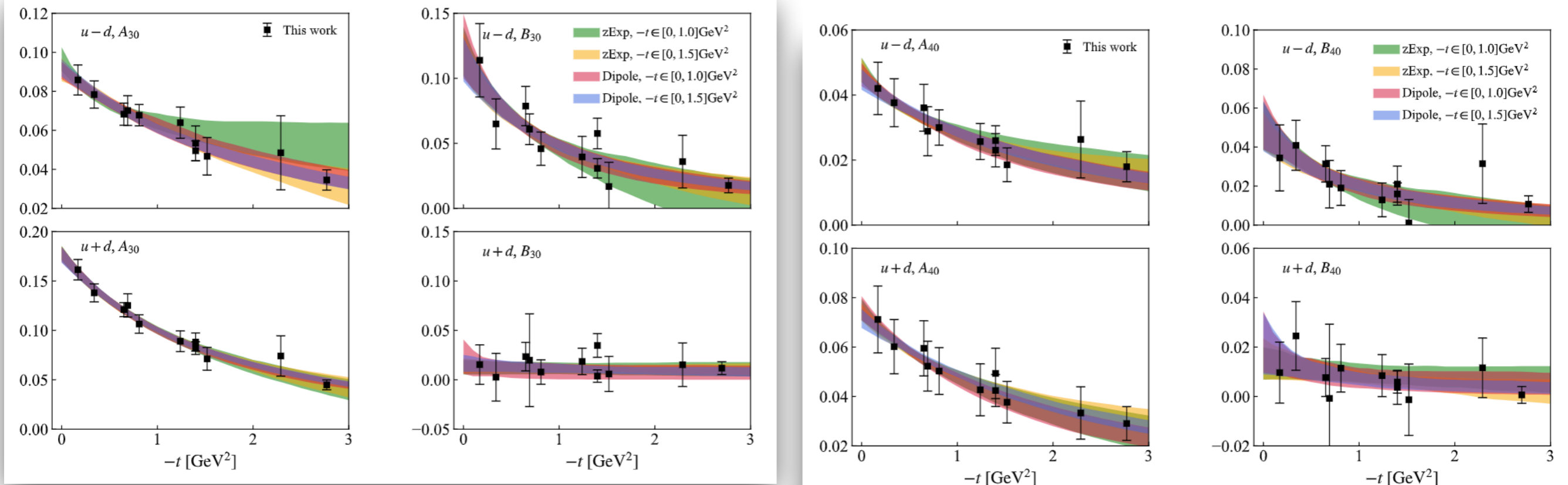
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Beyond leading twist

Extraction twist-3 very challenging
both experimentally and theoretically

Theoretical setup



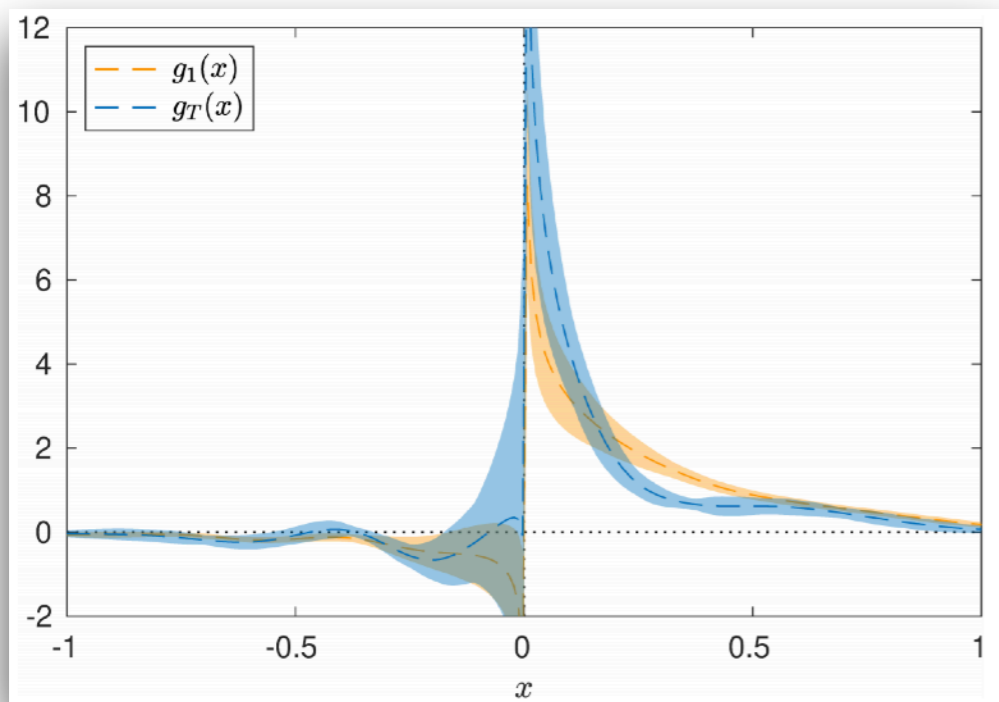
[D. Kiptily and M. Polyakov, Eur. Phys. J. C37 (2004) 105]

[F. Aslan et al., Phys. Rev. D 98, 014038 (2018)]

$$F^{[\gamma^\mu \gamma_5]}(x, \Delta; P^3) = \frac{1}{2P^3} \bar{u}(p_f, \lambda') \left[P^\mu \frac{\gamma^3 \gamma_5}{P^0} F_{\tilde{H}}(x, \xi, t; P^3) + P^\mu \frac{\Delta^3 \gamma_5}{2mP^0} F_{\tilde{E}}(x, \xi, t; P^3) \right. \\ \left. + \Delta_\perp^\mu \frac{\gamma_5}{2m} F_{\tilde{E}+\tilde{G}_1}(x, \xi, t; P^3) + \gamma_\perp^\mu \gamma_5 F_{\tilde{H}+\tilde{G}_2}(x, \xi, t; P^3) \right. \\ \left. + \Delta_\perp^\mu \frac{\gamma^3 \gamma_5}{P^3} F_{\tilde{G}_3}(x, \xi, t; P^3) + i\varepsilon_\perp^{\mu\nu} \Delta_\nu \frac{\gamma^3}{P^3} F_{\tilde{G}_4}(x, \xi, t; P^3) \right] u(p_i, \lambda)$$

★ Forward limit for twist-3: only $\tilde{H} + \tilde{G}_2 \equiv g_T$ survives

[S. Bhattacharya et al., PRD 102 (2020) 11 (Editors Selection)]



Twist-3 very important and have physical interpretation:

- as sizable as twist-2
- contain information about quark-gluon correlations inside hadrons
- appear in QCD factorization theorems for various observables (e.g. g_2)

Talk by S. Li

Theoretical setup



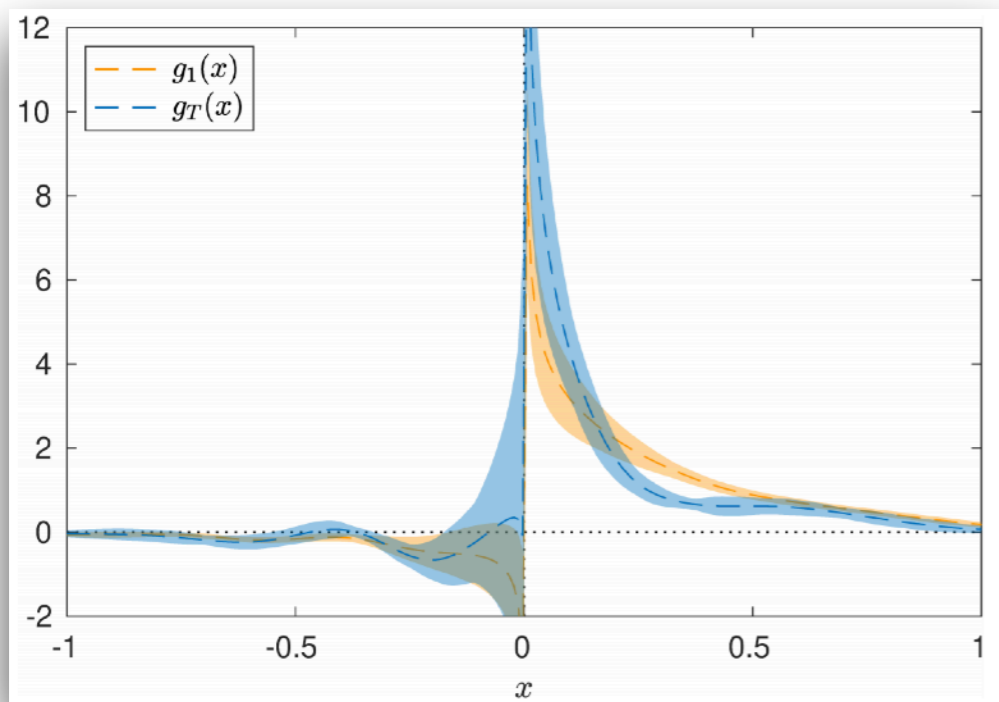
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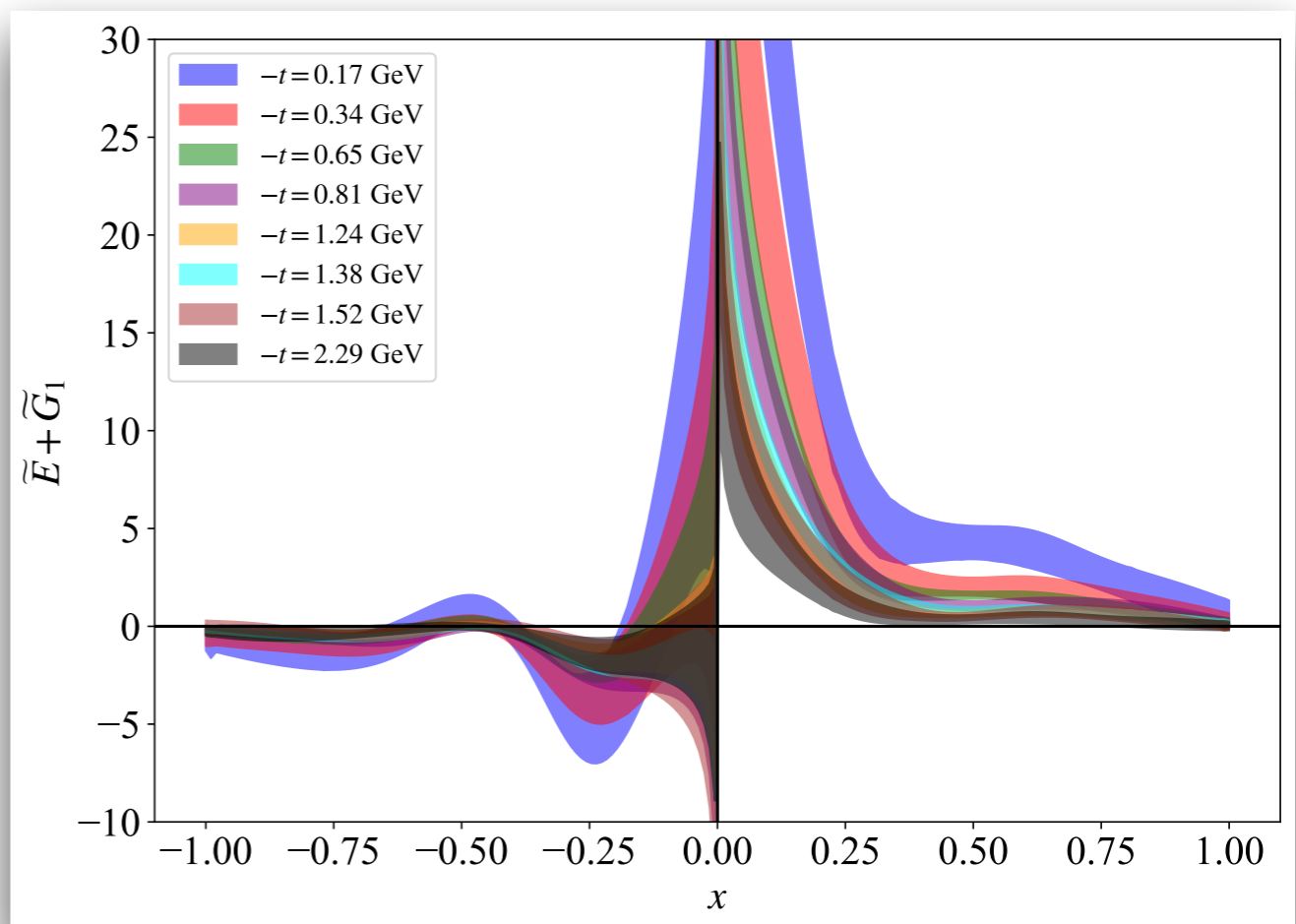
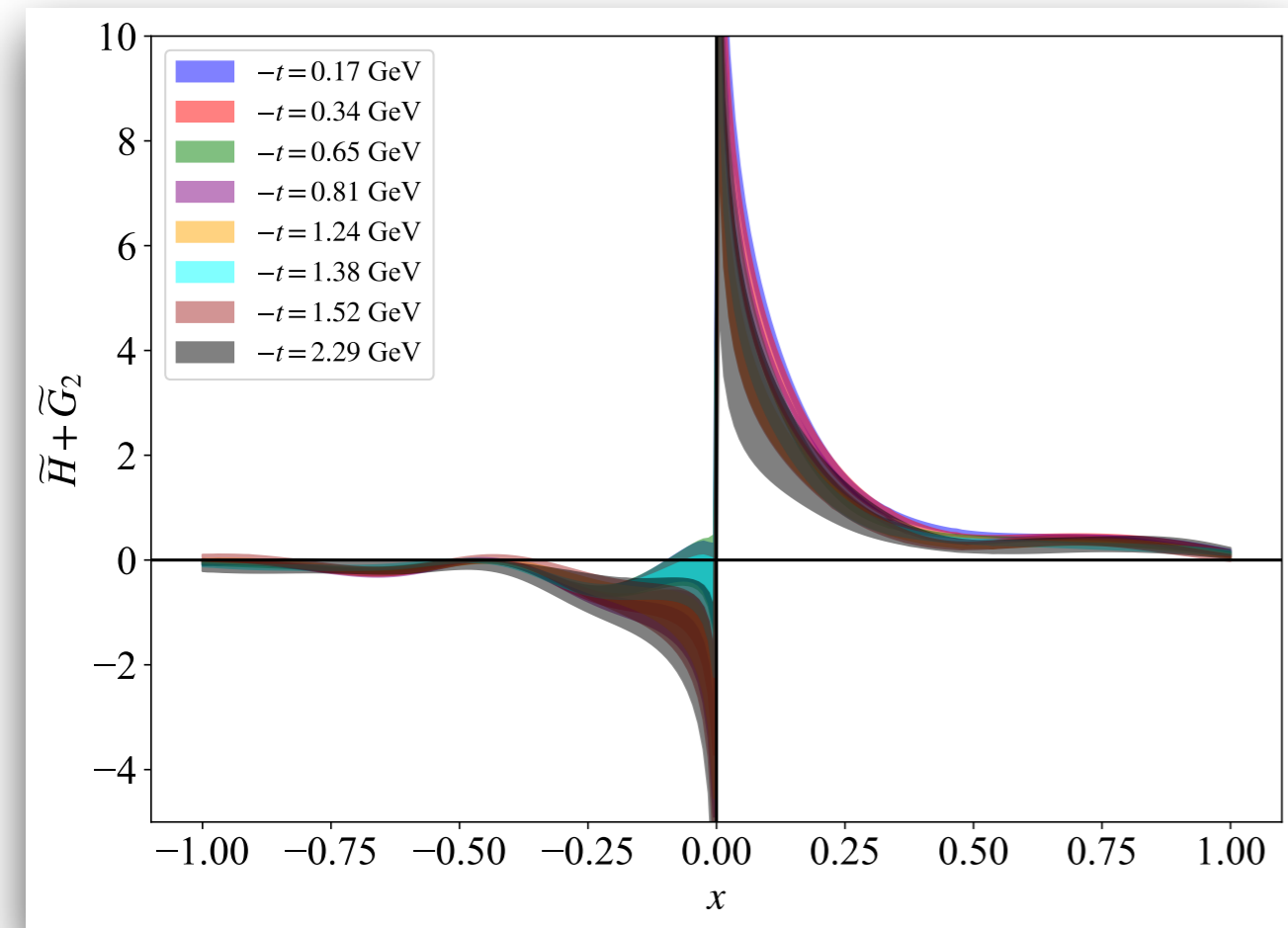
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$$= \bar{u}(p_f, \lambda') \left[\frac{i\varepsilon^{\mu P z \Delta}}{m} \tilde{A}_1 + \gamma^\mu \gamma_5 \tilde{A}_2 + \gamma_5 \left(\frac{P^\mu}{m} \tilde{A}_3 + m z^\mu \tilde{A}_4 + \frac{\Delta^\mu}{m} \tilde{A}_5 \right) \right. \\ \left. + m \not{z} \gamma_5 \left(\frac{P^\mu}{m} \tilde{A}_6 + m z^\mu \tilde{A}_7 + \frac{\Delta^\mu}{m} \tilde{A}_8 \right) \right] u(p_i, \lambda),$$

[D. Kiptily and M. Polyakov, Eur. Phys. J. C37 (2004) 105]

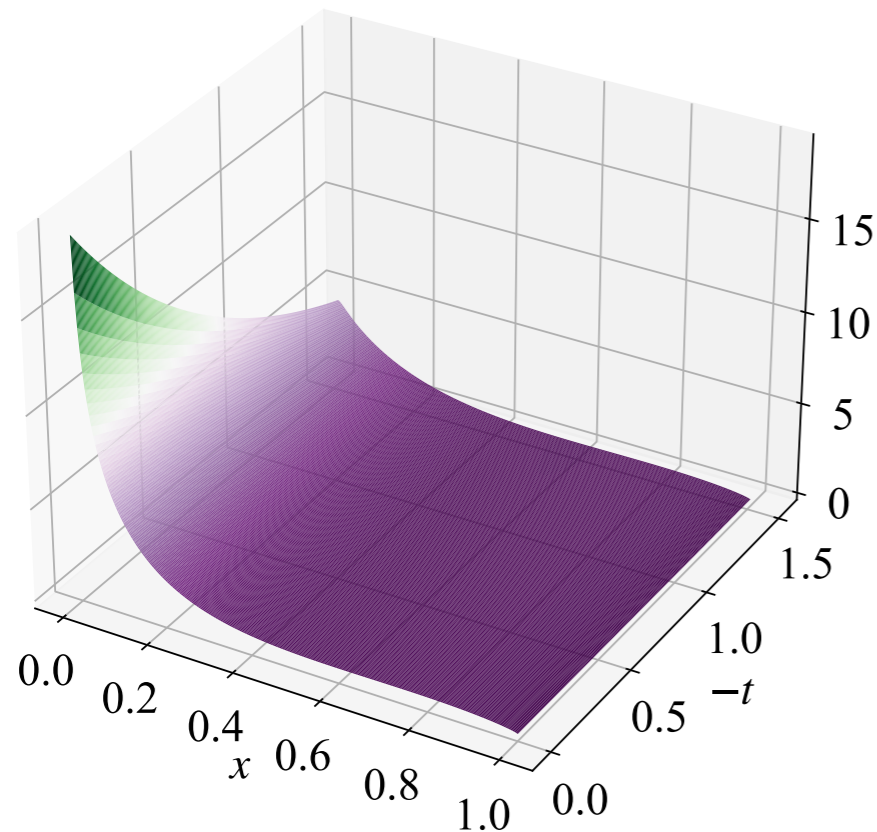
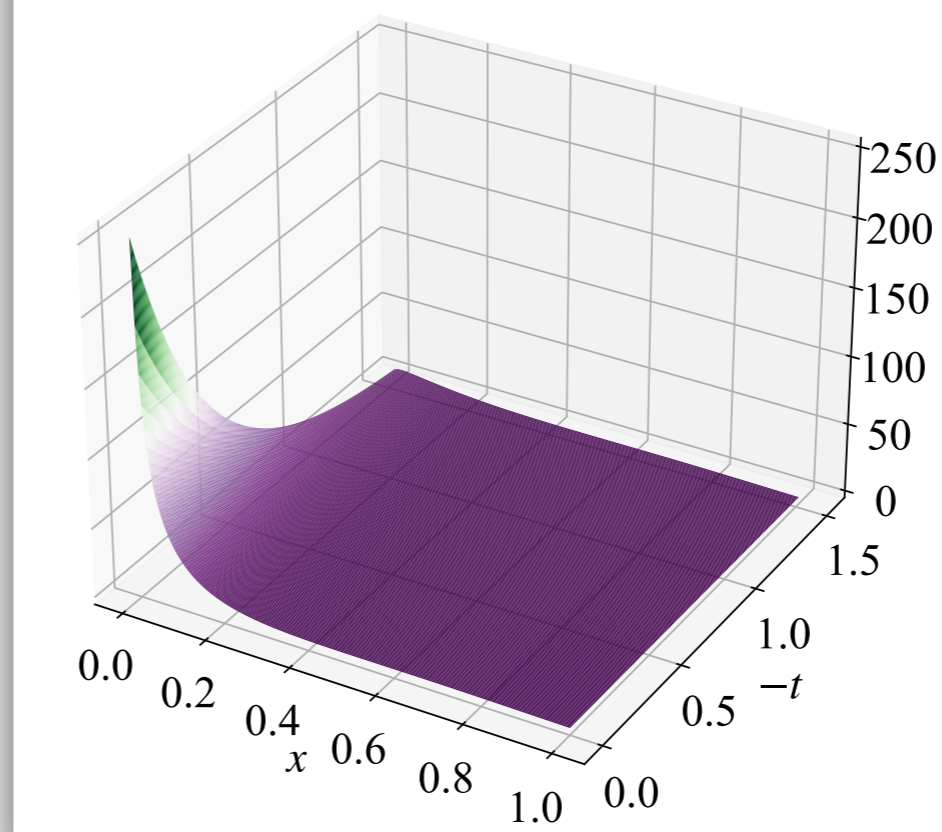
[S. Bhattacharya et al., 109 (2024) 3, 034508]



★ Parametrization of -t dependence (preliminary)

$$\text{GPD}(x, -t, 0) = A x^{\alpha_0 - \alpha_1 t} (1 - x)^\beta$$

Ademollo & Del Giudice Gatto & Preparata

$\tilde{H} + \tilde{G}_2$  $\tilde{E} + \tilde{G}_1$ 

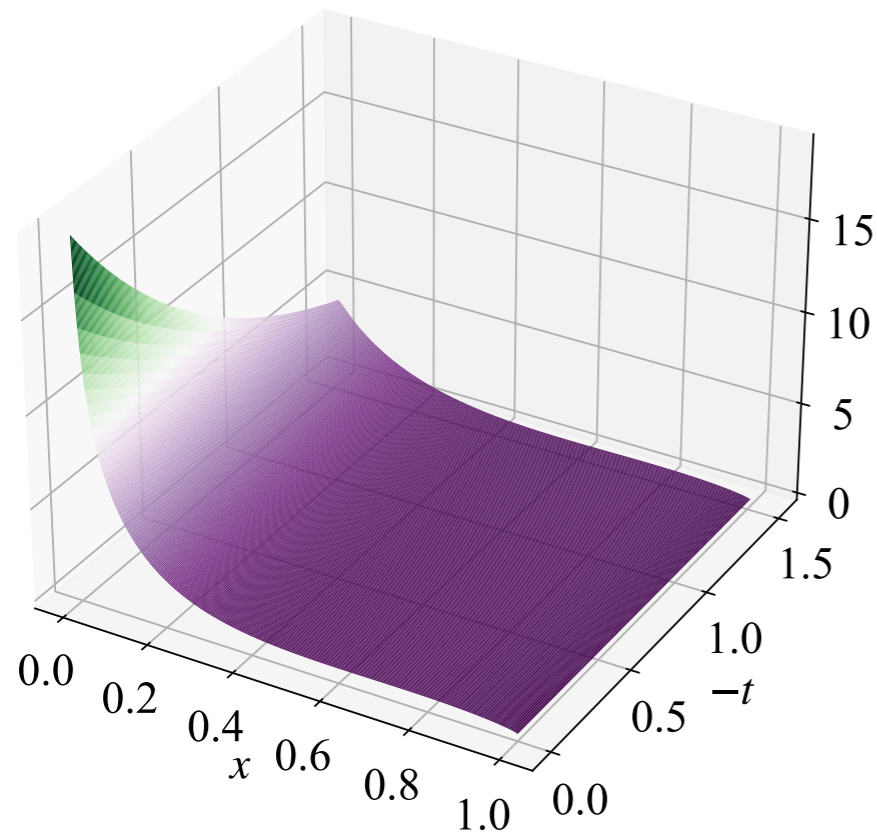
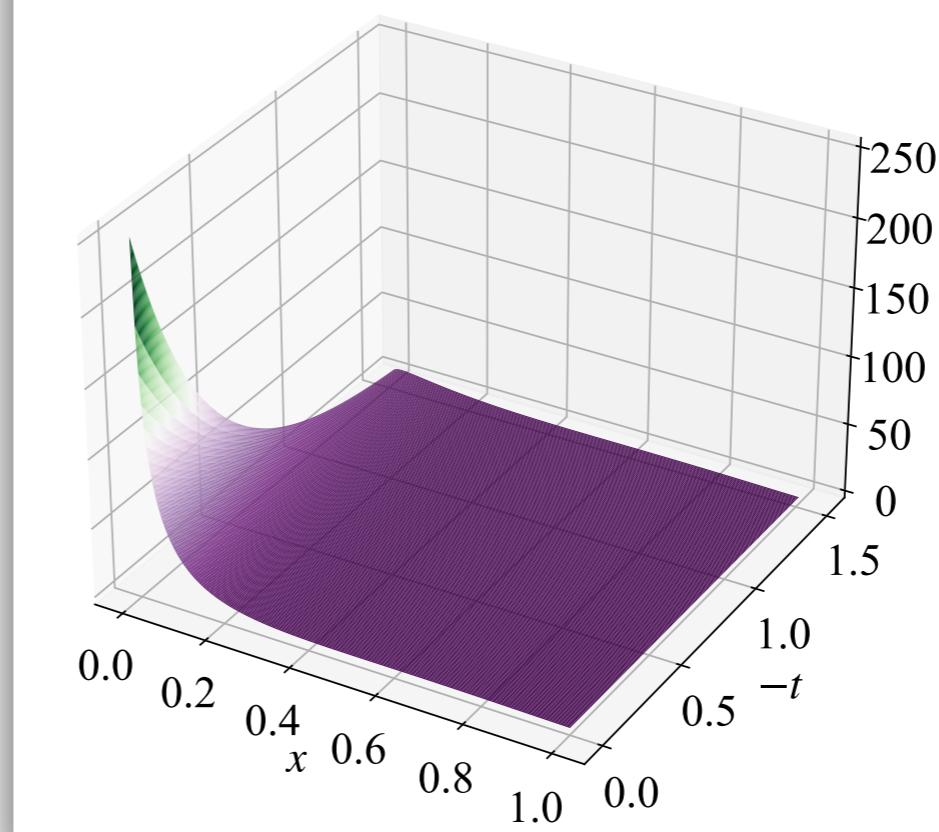
- ★ Direct access to \tilde{E} -GPD not possible for zero skewness

$$P^\mu \frac{\Delta^3 \gamma_5}{2mP^0} \underline{F_{\tilde{E}}}(x, \xi, t; P^3)$$

- ★ Glimpse into \tilde{E} -GPD through twist-3 :

$$\int_{-1}^1 dx \tilde{E}(x, \xi, t) = G_P(t)$$

$$\int_{-1}^1 dx \tilde{G}_i(x, \xi, t) = 0, \quad i = 1, 2, 3, 4$$

$\widetilde{H} + \widetilde{G}_2$  $\widetilde{E} + \widetilde{G}_1$ 

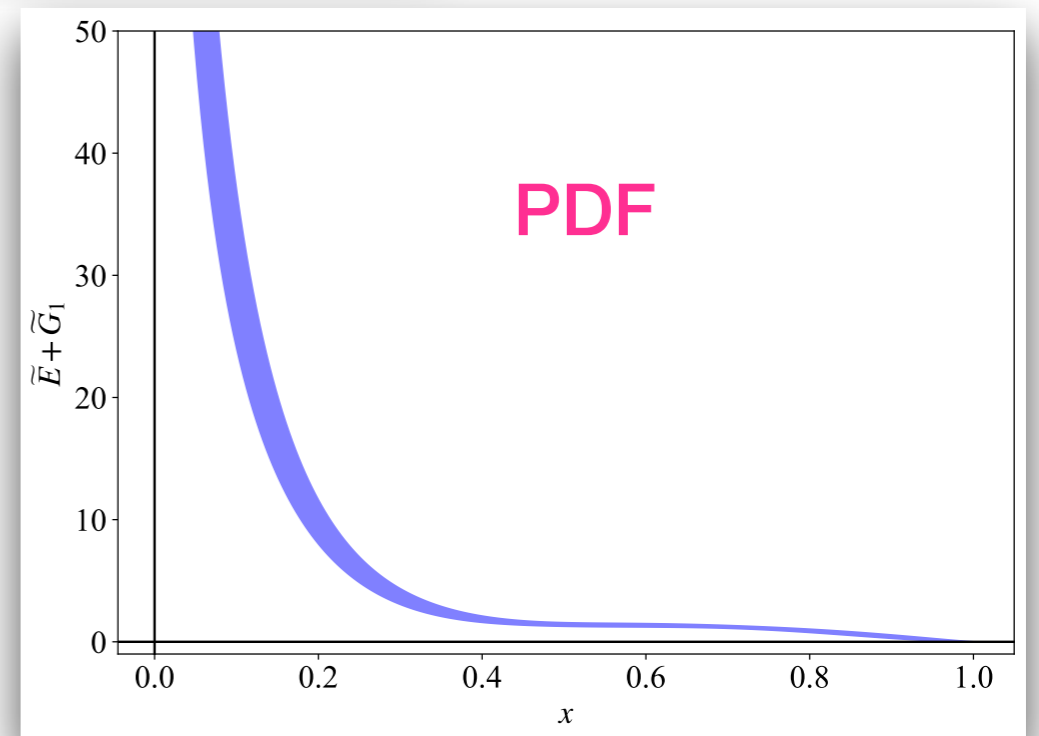
- ★ Direct access to \widetilde{E} -GPD not possible for zero skewness

$$P^\mu \frac{\Delta^3 \gamma_5}{2mP^0} \underline{F_{\widetilde{E}}}(x, \xi, t; P^3)$$

- ★ Glimpse into \widetilde{E} -GPD through twist-3 :

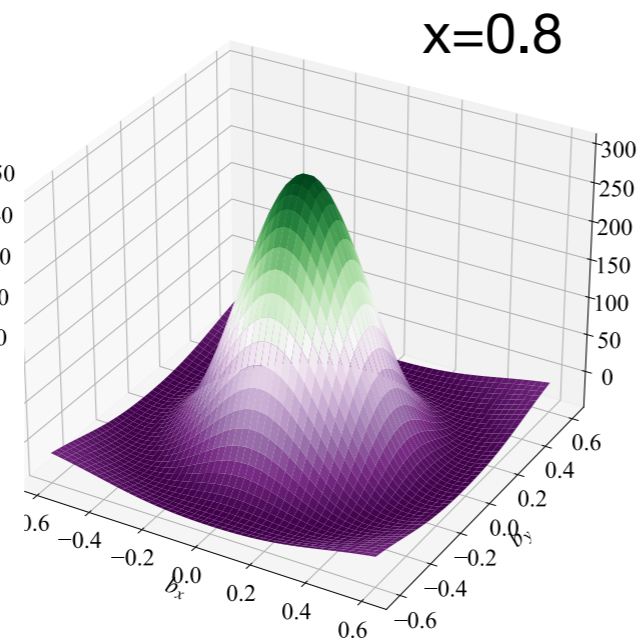
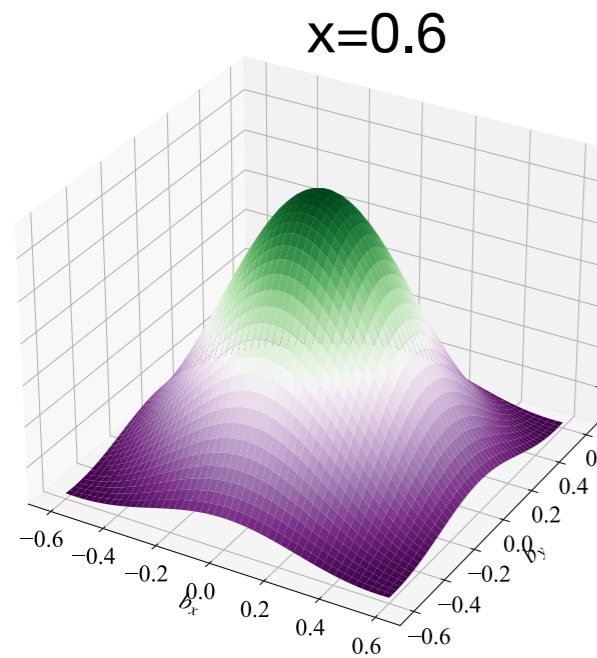
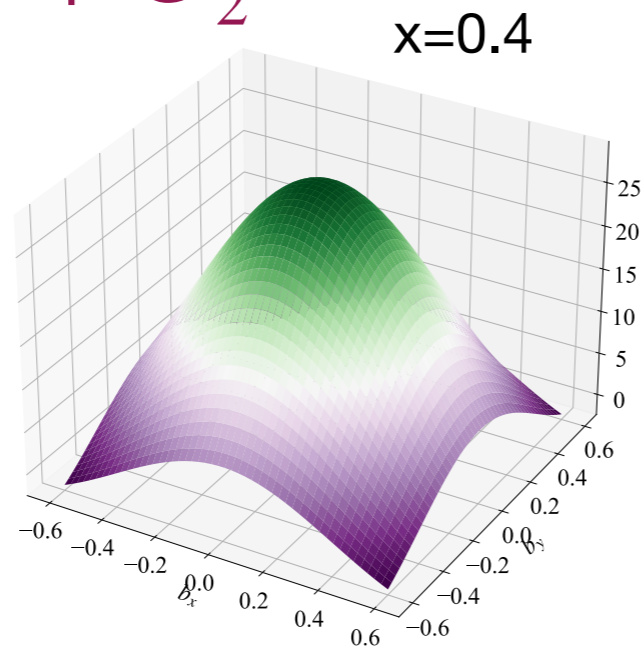
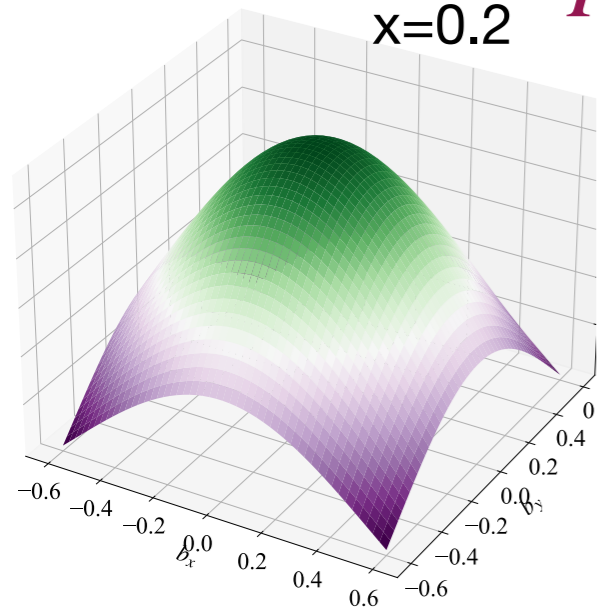
$$\int_{-1}^1 dx \widetilde{E}(x, \xi, t) = G_P(t)$$

$$\int_{-1}^1 dx \widetilde{G}_i(x, \xi, t) = 0, \quad i = 1, 2, 3, 4$$

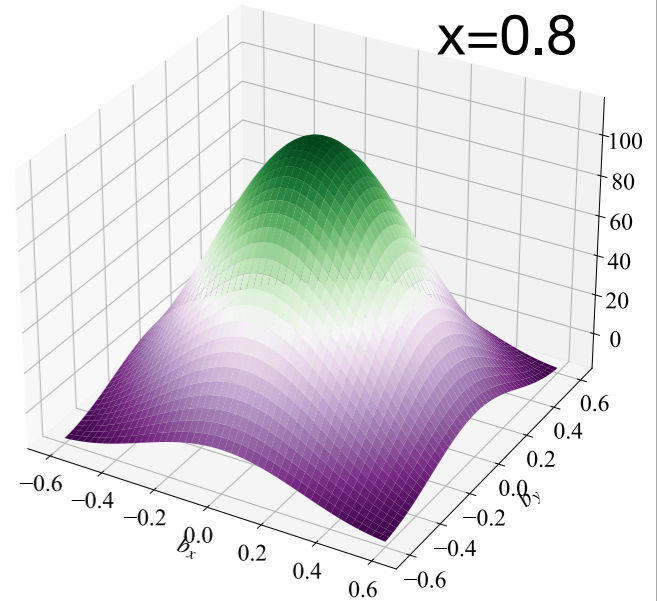
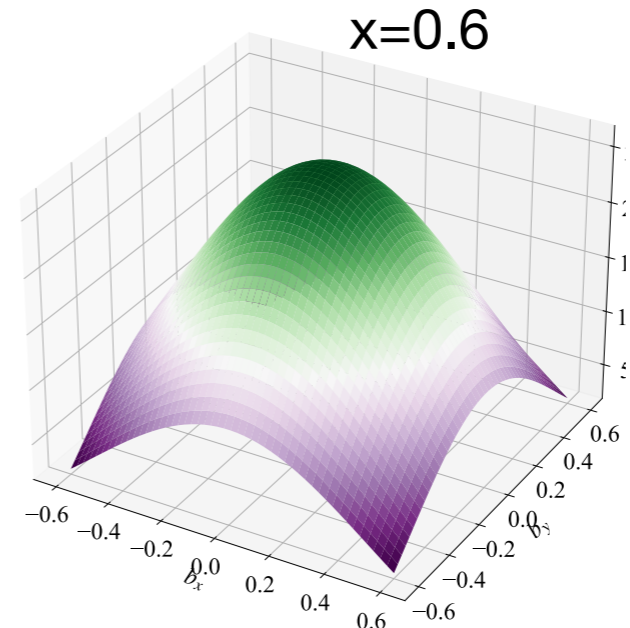
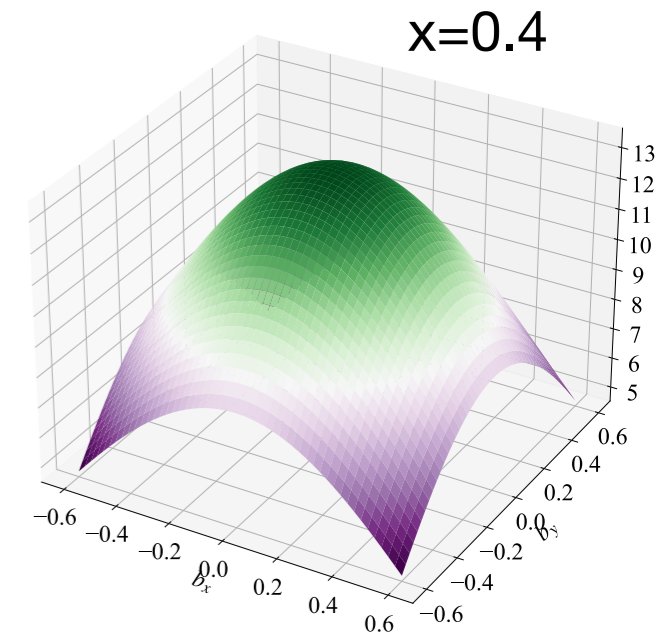
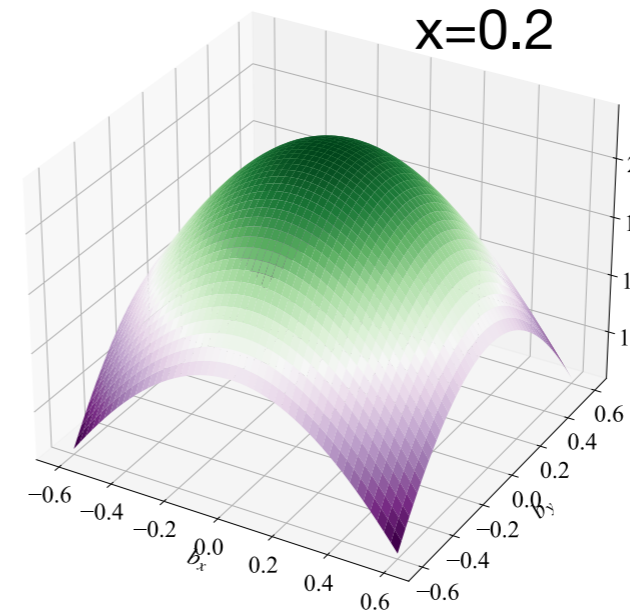


Impact parameter space

$$\widetilde{H} + \widetilde{G}_2$$



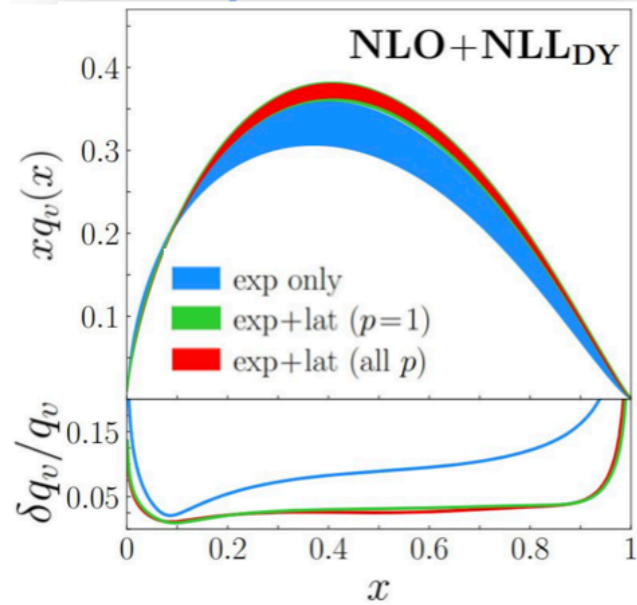
$$\widetilde{E} + \widetilde{G}_1$$



Synergy/Complementarity of lattice and phenomenology

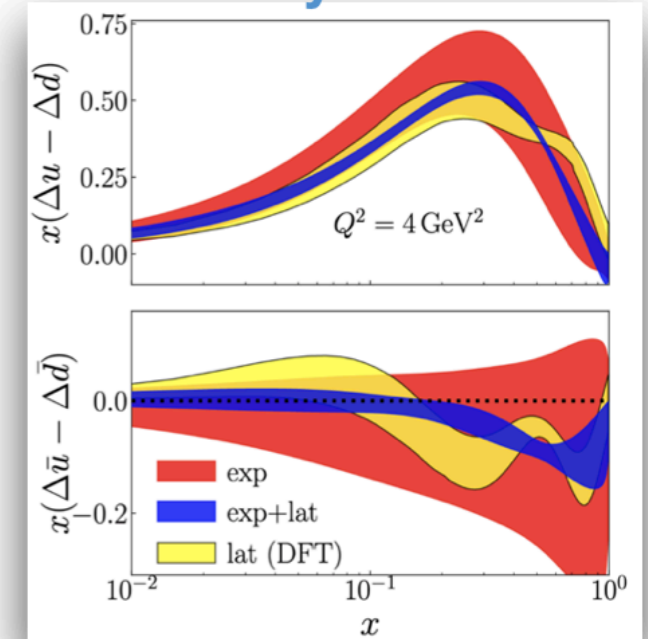
Synergies: constraints & predictive power of lattice QCD

pion PDF

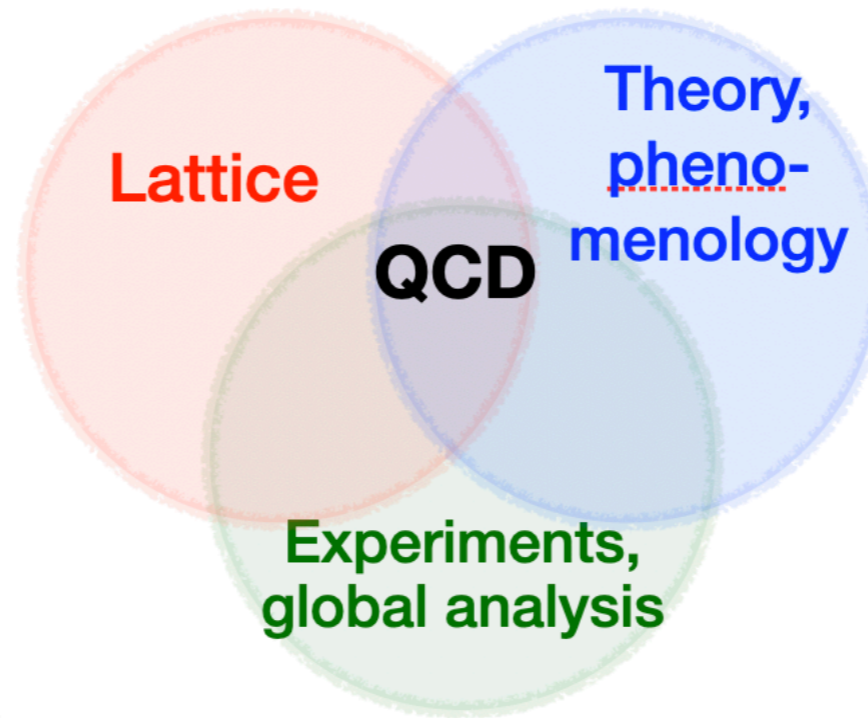


[JAM/HadStruc, PRD105 (2022) 114051]

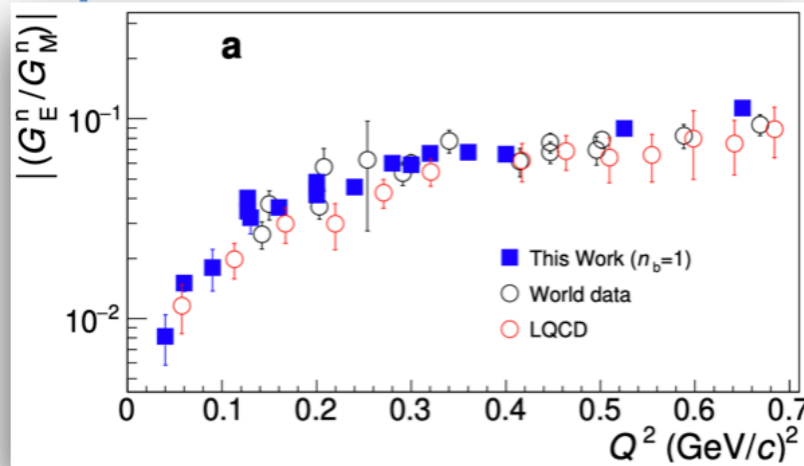
helicity PDF



[JAM & ETMC, PRD 103 (2021) 016003]

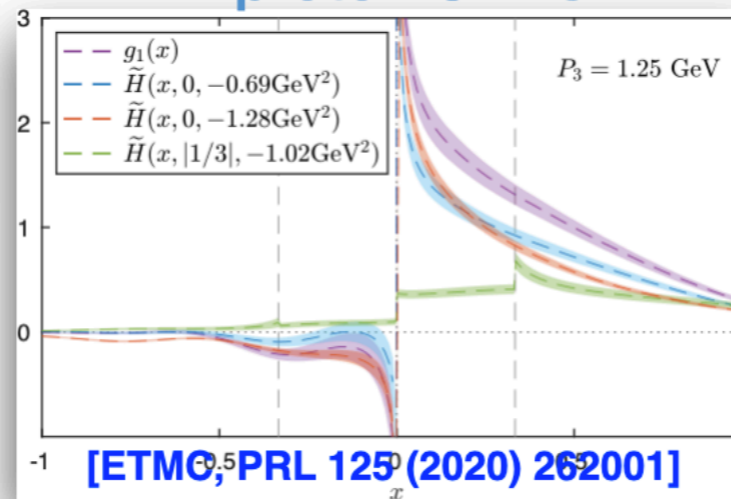


proton & neutron radius



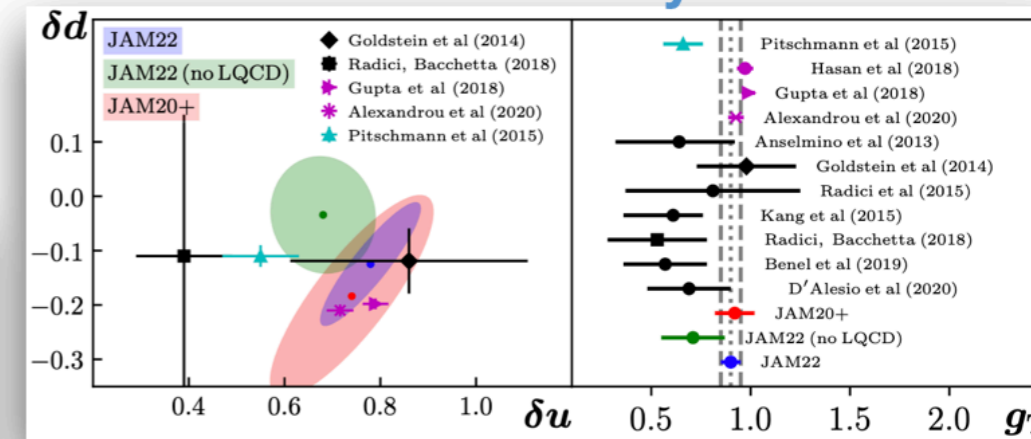
[Atac et al., Nature Comm. 12, 1759 (2021)]

proton GPDs



[ETMC, PRL 125 (2020) 262001]

transversity PDF



[JAM, PRD 106 (2022) 3, 034014]

And many more!



Toward synergy for GPDs

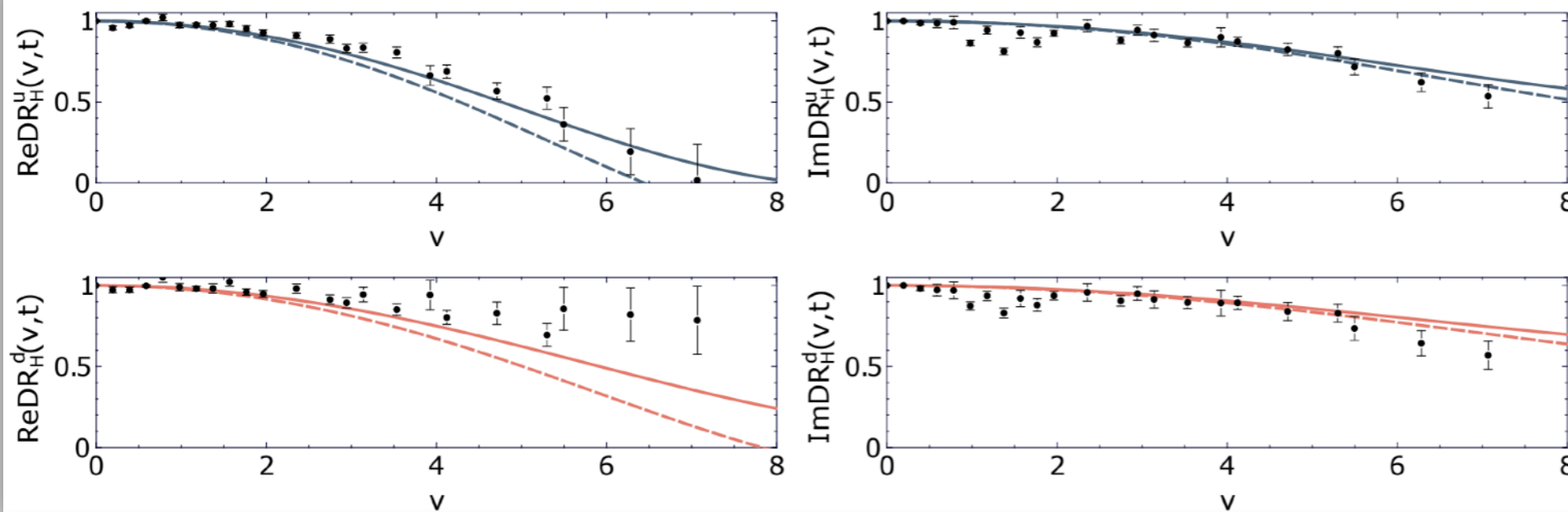
- ★ Forming ratios of GPDs seems to suppress systematic uncertainties

[K. Cichy et al., PRD 110 (2024) 11, 114025]

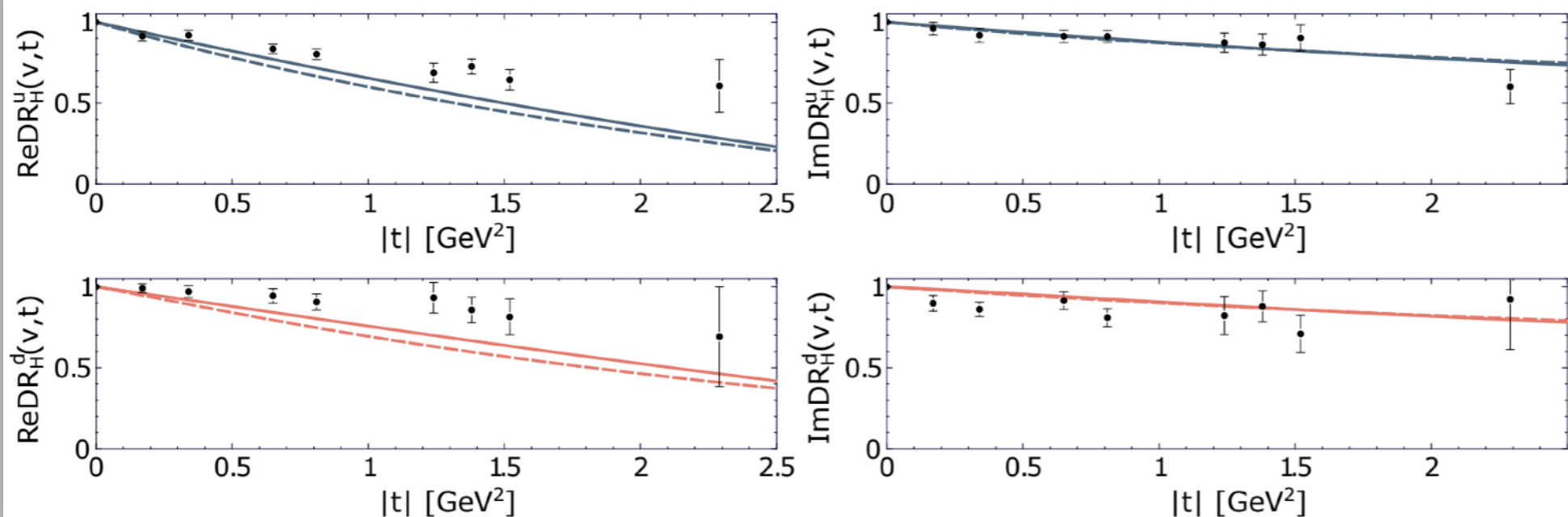
$$\text{DR}_{\text{Re}}^{\hat{H}^q}(\nu, t) = \frac{\text{Re}\hat{H}^q(\nu, t)}{\text{Re}\hat{H}^q(\nu, 0)} \frac{\text{Re}\hat{H}^q(0, 0)}{\text{Re}\hat{H}^q(0, t)},$$

$$\text{DR}_{\text{Im}}^{\hat{H}^q}(\nu, t) = \lim_{\nu' \rightarrow 0} \frac{\text{Im}\hat{H}^q(\nu, t)}{\text{Im}\hat{H}^q(\nu, 0)} \frac{\text{Im}\hat{H}^q(\nu', 0)}{\text{Im}\hat{H}^q(\nu', t)}$$

(a) As a function of ν for $|t| = 0.65 \text{ GeV}^2$.



(b) As a function of $|t|$ for $\nu = 3.14$.

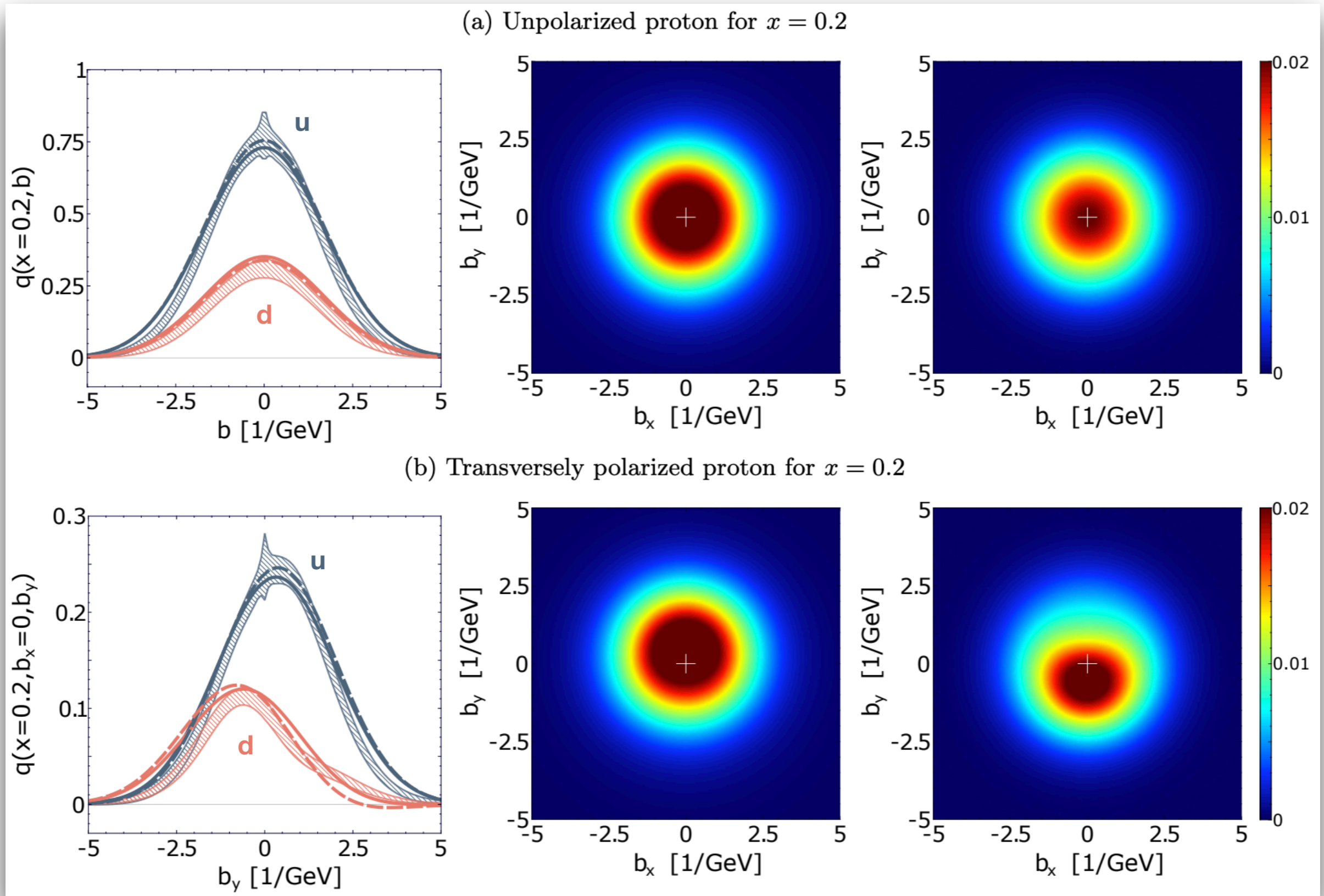


- GK (solid curve)
- VGG (dashed curve)

- Good agreement for up quark
- Reasonable agreement for down quark

- Further study needed on how to combine lattice results with data

Tomographic Images



- GK (solid line), VGG (dashed line)

[K. Cichy et al., PRD 110 (2024) 11, 114025]

How to lattice QCD data fit into the overall effort for hadron tomography

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- ★ Lattice data may be incorporated in global analysis of experimental data and may influence parametrization of t and ξ dependence

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1. **Theoretical studies** of high-momentum transfer processes using perturbative QCD methods and study of GPDs properties
2. **Lattice QCD** calculations of GPDs and related structures
3. **Global analysis** of GPDs based on experimental data using modern data analysis techniques for inference and uncertainty quantification

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- ★ Lattice data may be incorporated in global analysis of experimental data and may influence parametrization of t and ξ dependence



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Other GPD global analysis efforts:

- Gepard [<https://gepard.phy.hr/>]
- PARTONS [<https://partons.cea.fr>]
- EXCLAIM [<https://exclaimcollab.github.io/web.github.io/#/>]

See talk by S. Liuti

Concluding Remarks

- ★ Impressive progress in the extraction of PDFs from Lattice QCD
- ★ Extensive programs in Gluon PDFs
- ★ New Developments in several promising directions:
DA, GPDs, TMDs
- ★ Synergy with phenomenology has the potential to enhance the impact of lattice QCD data and complement data sets

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