

Pythia 8 for Electron-Ion Collisions

Physics Opportunities at an Electron-Ion Collider XI

Ilkka Helenius

February 26, 2025



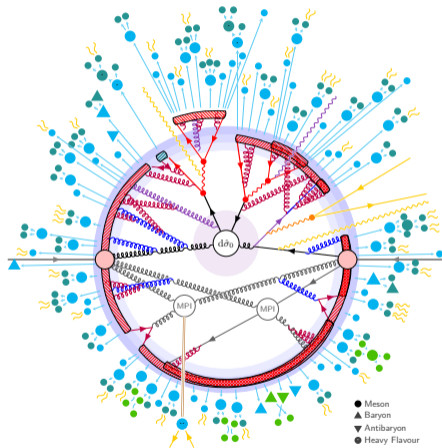
Pythia 8 event generator

Pythia 8: A general purpose event generator

- Latest release 8.313 (Jan 2025)
- A complete physics manual for 8.3
[SciPost Phys. Codebases 8-r8.3 (2022)]

Outline

- Pythia 8 introduction
- Deep inelastic scattering (HERA)
- Photoproduction
 - With proton target (HERA)
 - With nuclear target (UPC@LHC)
- Summary & Outlook



[figure by P. Skands]

Current members (in 8.313 release)

- Javira Altmann (Monash University)
- Christian Bierlich (Lund University)
- Naomi Cooke (University of Glasgow)
- Ilkka Helenius (University of Jyväskylä)
- Philip Ilten (University of Cincinnati)
- Leif Lönnblad (Lund University)
- Stephen Mrenna (Fermilab)
- Christian Preuss (University of Wuppertal)
- Torbjörn Sjöstrand (Lund University)
- Peter Skands (Monash University)



[Pythia Week in Oxford 2024]

- Spokesperson
- Codemaster
- Webmaster

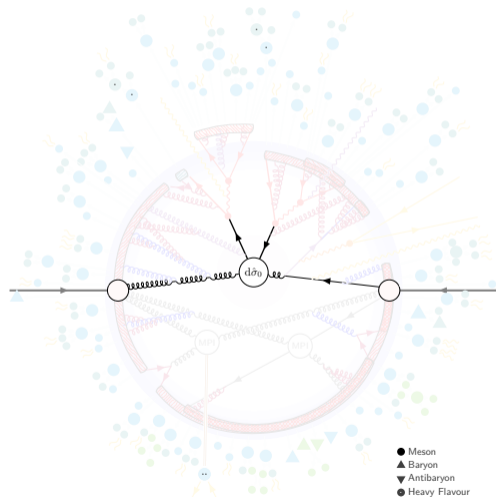
<https://pythia.org>
authors@pythia.org

Physics modelled within Pythia 8

Classify event generation in terms of
“hardness”

1. Hard Process (here $t\bar{t}$)

[figure credit: P. Skands]

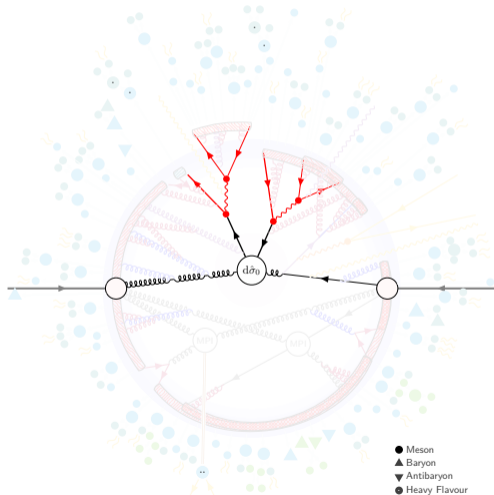


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Classify event generation in terms of “hardness”

1. Hard Process (here $t\bar{t}$)
2. Resonance decays (t, Z, \dots)

[figure credit: P. Skands]

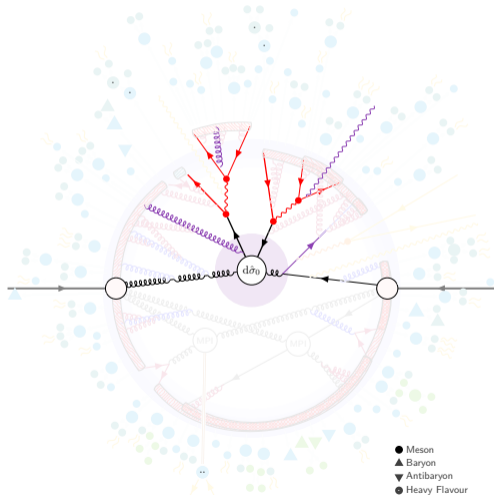


Physics modelled within Pythia 8

Classify event generation in terms of “hardness”

1. Hard Process (here $t\bar{t}$)
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3. Matching, Merging and matrix-element corrections

[figure credit: P. Skands]

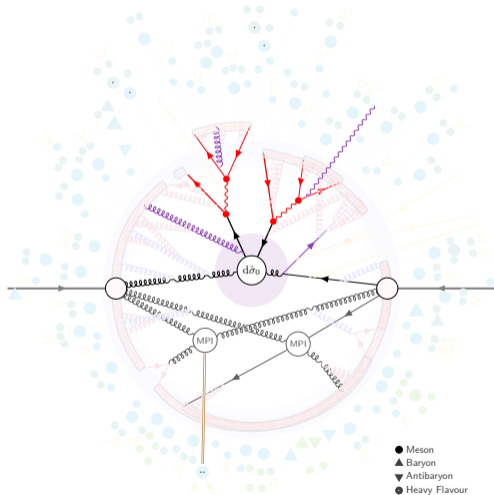


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1. Hard Process (here $t\bar{t}$)
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3. Matching, Merging and matrix-element corrections
4. Multiparton interactions

[figure credit: P. Skands]

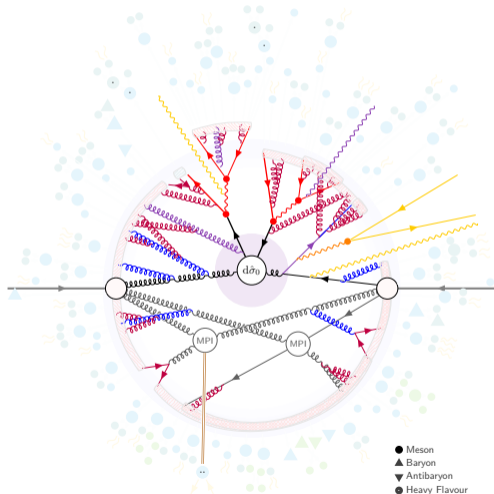


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5. Parton showers:
ISR, FSR, QED, Weak

[figure credit: P. Skands]

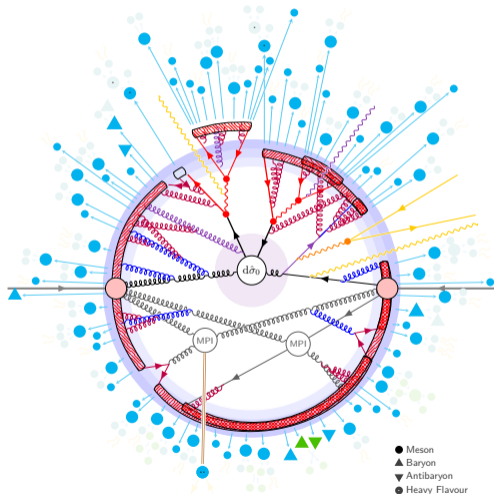


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[figure credit: P. Skands]

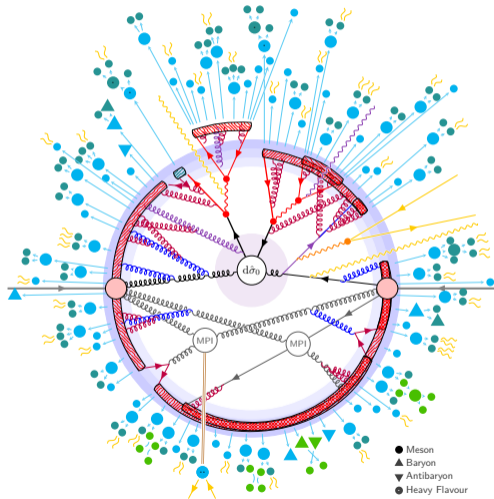


Physics modelled within Pythia 8

Classify event generation in terms of “hardness”

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4. Multiparton interactions
5. Parton showers:
ISR, FSR, QED, Weak
6. Hadronization, Beam remnants
7. Decays, Rescattering

[figure credit: P. Skands]



Available beam configurations in Pythia 8

Hadronic collisions

- p-p: hard, soft and low-energy processes
- h-p, where $h = \pi^{\pm,0}, K^{\pm,0}, \phi^0, \dots$

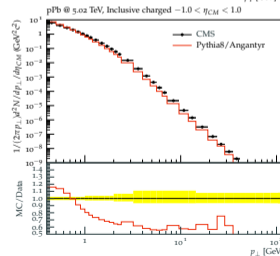
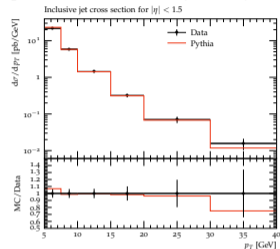
Collisions with leptons

- e^+e^- , including $\gamma\gamma$ (also in p-p)
- e-p: (neutrino) DIS, photoproduction with soft and hard QCD processes

Heavy-ion collisions with Angantyr

- A-A, p-A and h-A
- UPCs with proton target, also VMD-A
- Some cosmic-ray related processes

[OPAL: PLB 658 (2008) 185-192]



Electron-proton collisions

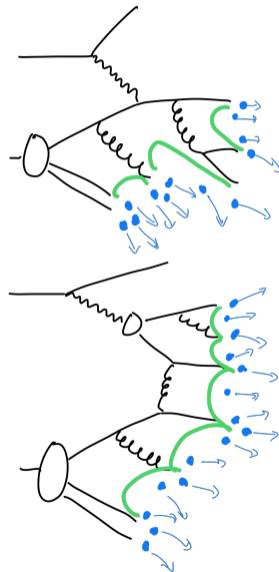
Classified in terms photon virtuality Q^2

Deep inelastic scattering (DIS)

- High virtuality, $Q^2 > \text{a few GeV}^2$
- Lepton scatters off from a parton by exchanging a highly virtual photon

Photoproduction

- Low virtuality, $Q^2 \rightarrow 0 \text{ GeV}^2$
⇒ Direct and resolved contributions
- Factorize γ flux, evolve γp system
- Hard scale provided by the final state
- Also soft QCD processes, diffraction



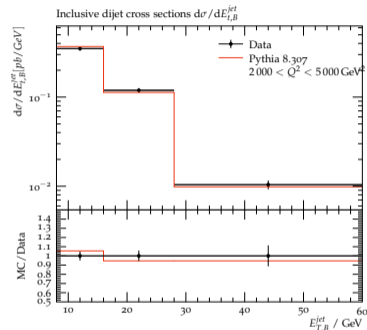
Deep inelastic scattering (DIS)

Jet production in DIS

- PS accurate only for soft and collinear emissions
- Matrix element corrections helps at high- Q^2 but still misses low- Q^2 high- E_T part

Merging in DIS

- Hard events with several final-state partons
- Combine with parton shower emissions using merging algorithms to avoid double counting
- Need a dynamically calculated merging scale to account for multiple scales (Q^2 and p_T^{jet})
- Vincia sector shower with a unique history



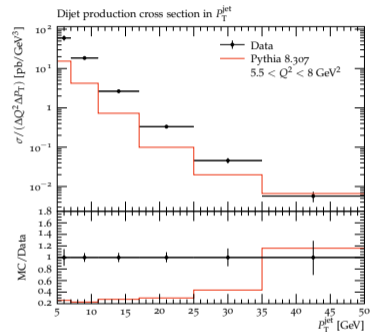
[ZEUS: EPJC 70 (2010) 965-982]

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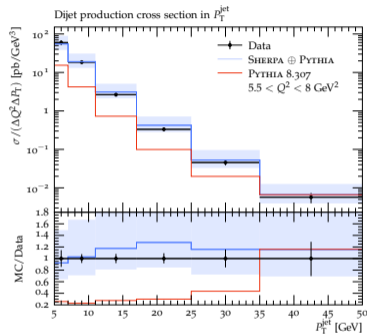
[H1: EPJC 77 (2017) 215]

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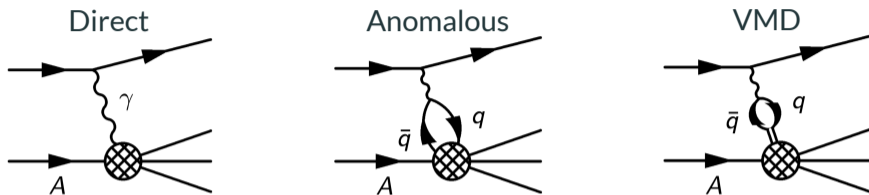


[H1: EPJC 77 (2017) 215]

Part of the 8.313 release,
with an example

Photoproduction

Photon structure at $Q^2 \approx 0 \text{ GeV}^2$



Partonic structure of resolved (anom. + VMD) photon encoded in photon PDFs

$$f_i^\gamma(x_\gamma, \mu^2) = f_i^{\gamma, \text{dir}}(x_\gamma, \mu^2) + f_i^{\gamma, \text{anom}}(x_\gamma, \mu^2) + f_i^{\gamma, \text{VMD}}(x_\gamma, \mu^2)$$

- $f_i^{\gamma, \text{dir}}(x_\gamma, \mu^2) = \delta_{i\gamma} \delta(1 - x_\gamma)$
- $f_i^{\gamma, \text{anom}}(x_\gamma, \mu^2)$: Perturbatively calculable
- $f_i^{\gamma, \text{VMD}}(x_\gamma, \mu^2)$: Non-perturbative, fitted or vector-meson dominance (VMD)

Factorized cross section

$$d\sigma^{\gamma A \rightarrow kl+X} = f_i^\gamma(x_\gamma, \mu^2) \otimes f_j^A(x_p, \mu^2) \otimes d\sigma^{ij \rightarrow kl}$$

Multiparton interactions (MPIs) with resolved photons

- MPIs from $2 \rightarrow 2$ QCD cross sections

$$\frac{d\mathcal{P}_{\text{MPI}}}{dp_T^2} = \frac{1}{\sigma_{\text{nd}}(\sqrt{s})} \frac{d\sigma^{2 \rightarrow 2}}{dp_T^2}$$

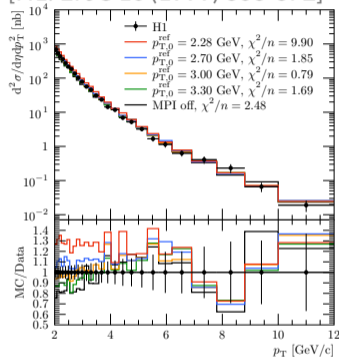
$\sigma_{\text{nd}}(\sqrt{s})$ is the non-diffractive cross section

- Partonic cross section diverges at $p_T \rightarrow 0$
 \Rightarrow Introduce a screening parameter p_{T0}

$$\frac{d\sigma^{2 \rightarrow 2}}{dp_T^2} \propto \frac{\alpha_s(p_T^2)}{p_T^4} \rightarrow \frac{\alpha_s(p_{T0}^2 + p_T^2)}{(p_{T0}^2 + p_T^2)^2}$$

- Energy-dependent parametrization:
 $p_{T0}(\sqrt{s}) = p_{T0}^{\text{ref}}(\sqrt{s}/\sqrt{s_{\text{ref}}})^\alpha$
- Number of interactions: $\langle n \rangle = \sigma_{\text{int}}(p_{T0})/\sigma_{\text{nd}}$

[H1: EPJC 10 (1999) 363-372]



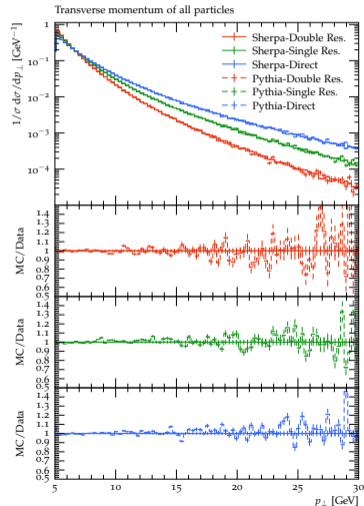
- Use H1 data to (re-)tune parameter(s)
- $\langle W_{\gamma p} \rangle \approx 200 \text{ GeV}$

Comparisons between Pythia, Sherpa and Herwig

[I. Helenius, P. Meinzinger, S. Plätzer, P. Richardson: arXiv:2406.08026 [hep-ph]]

Compare different generators for photoproduction

- Good agreement at ME-level
- Differences build up from inputs and modelling
- Scale variations large at LO



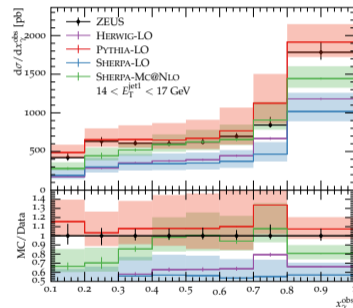
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Dijets in γp (HERA)



[ZEUS: EPJC 23 (2002) 615-631]

Comparisons between Pythia, Sherpa and Herwig

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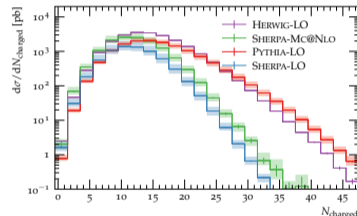
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Solid predictions for EIC require

- Validated inputs: (γ) PDFs, accurate flux
- Improved modelling for PS and remnant handling
- Tuning of models to HERA and LEP data

Predictions for multiplicity distributions in EIC

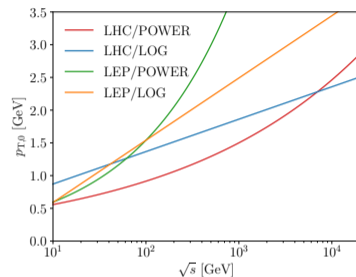


MPI tuning for photoproduction

[J.M. Butterworth, I. Helenius, J.J. Juan Castella, B. Pattengale, S. Sanjrani, M. Wing: SciPost Physics]

Systematic comparisons of MPI tunes

- pp at LHC and Tevatron and for $\gamma\gamma$ from LEP
- Data for jet and charged-particle production for pp, γp and $\gamma\gamma$ (10 data sets in total)



MPI tuning for photoproduction

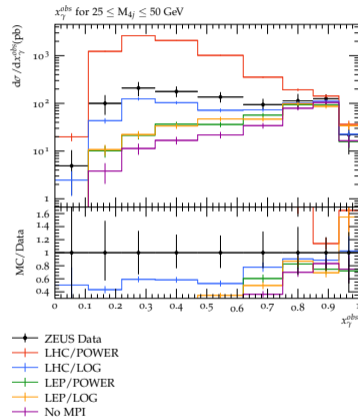
[J.M. Butterworth, I. Helenius, J.J. Juan Castella, B. Pattengale, S. Sanjrani, M. Wing: SciPost Physics]

Systematic comparisons of MPI tunes

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- Data for jet and charged-particle production for pp, γp and $\gamma\gamma$ (10 data sets in total)

Conclusions

- Standard pp tunes generate too many MPIs
- Can find good agreement for $\gamma\gamma$ and γp
- Further constraints from 3- and 4-jet production
- Published new Rivet analyses enabling dedicated tunes for each beam configuration



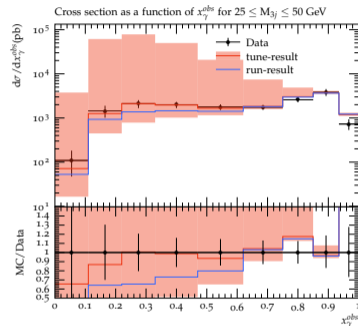
[ZEUS: NPB 792 1 (2008)]

Automized tuning with Professor 2

- Use the 3-/4-jet data from ZEUS
- Vary p_{T0}^{ref} and α , 100 points in parameter space
- Build interpolating function, minimize χ^2

Preliminary findings

- Large variation within reasonable variations
- Small discrepancy between the interpolated and the simulation with the resulting parameters



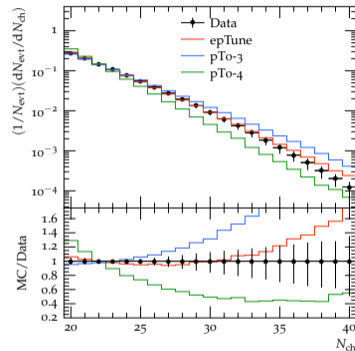
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Preliminary findings

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- Small discrepancy between the interpolated and the simulation with the resulting parameters
- Tune improve agreement with ZEUS multiplicity distribution
- Include more data and test universality



[ZEUS: JHEP 12 (2021) 102]

Nuclear target

Ultraperipheral heavy-ion collisions

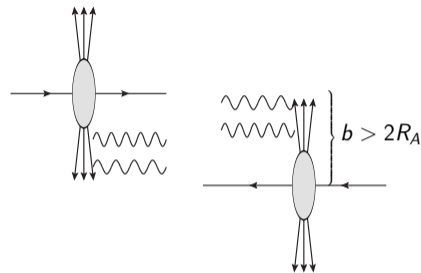
- Large impact parameter ($b \gtrsim 2R_A$)
⇒ No strong interactions
- At LHC relevant for p+p, p+Pb, Pb+Pb
- Large flux due to large EM charge of nuclei
⇒ $\gamma\gamma$ and γA collisions

Photon flux from equivalent photon approximation

- Define flux in impact-parameter space ⇒ Reject hadronic interactions with b_{\min}
- Integrating the point-like approximation we get

$$f_{\gamma}^A(x) = \frac{2\alpha_{\text{EM}}Z^2}{x\pi} \left[\xi K_1(\xi)K_0(\xi) - \frac{\xi^2}{2} (K_1^2(\xi) - K_0^2(\xi)) \right]$$

where $\xi = b_{\min} x m$ where $b_{\min} \approx 2R_A$ and m per nucleon mass



- Nuclear form factor heavily suppresses Q^2 of the photon ⇒ Photoproduction!

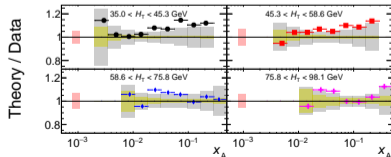
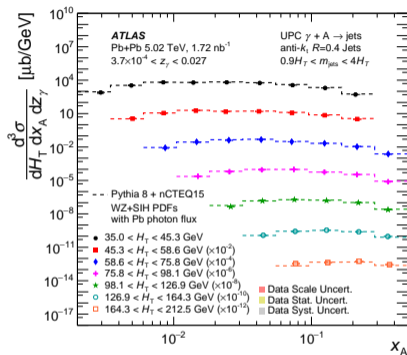
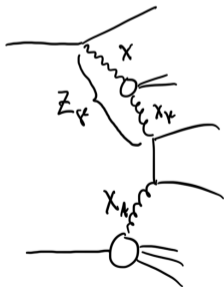
Dijets in ultra-peripheral heavy-ion collisions in XnOn

- Good agreement out of the box when accounting both direct and resolved
- EM nuclear break-up significant
- Pythia setup with nucleon target only
 \Rightarrow Is such a setup enough for $\gamma+A$?

$$H_T = \sum_i p_{T,i}$$

$$Z_\gamma = \frac{M_{\text{jets}}}{\sqrt{s_{NN}}} e^{+y_{\text{jets}}}$$

$$X_A = \frac{M_{\text{jets}}}{\sqrt{s_{NN}}} e^{-y_{\text{jets}}}$$

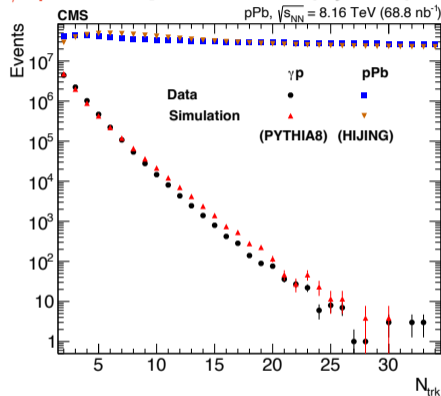


[ATLAS: arXiv:2409.11060]

Multiplicity distributions in UPCs

$\gamma+p$:

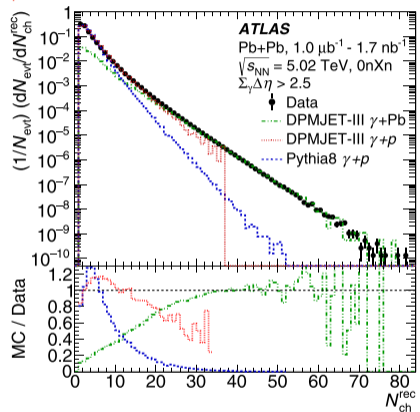
[CMS: Murillo Quijada, QM2022]



- Multiplicity distribution well reproduced in $\gamma+p$ interactions

$\gamma+Pb$:

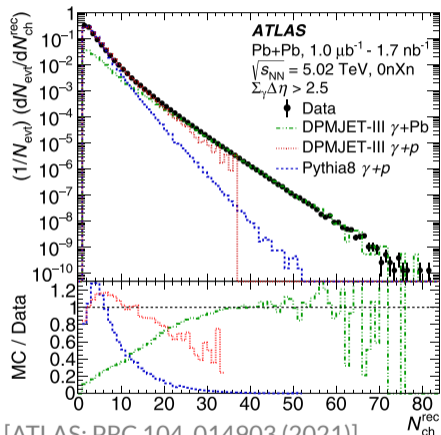
[ATLAS: PRC 104, 014903 (2021)]



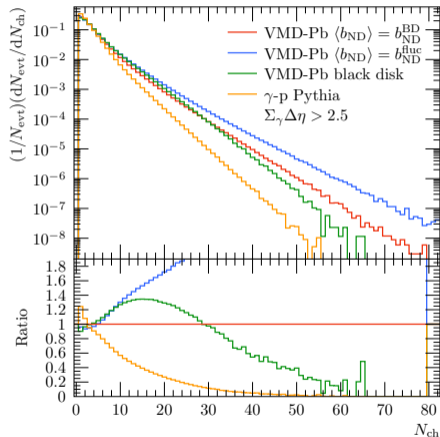
- High multiplicities missed with $\gamma+p$
⇒ Multi-nucleon interactions

Modelling $\gamma+A$ with Pythia

[I. Helenius, M. Uthm: EPJC 84 (2024) 11, 1155]



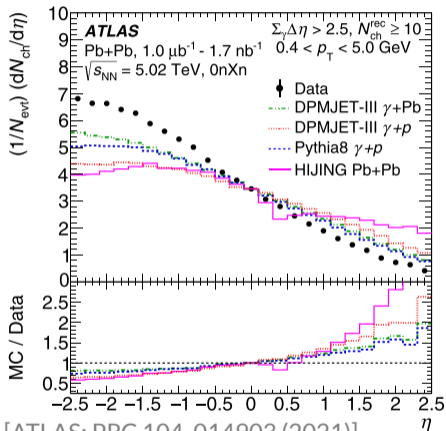
[ATLAS: PRC 104, 014903 (2021)]



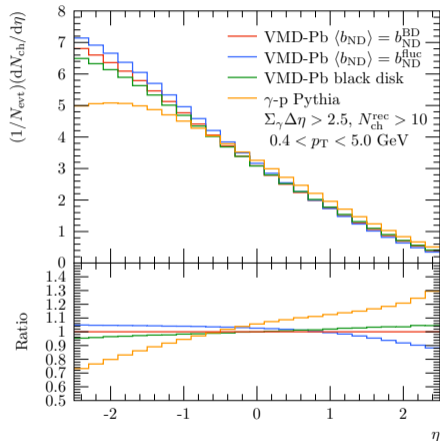
- ATLAS data not corrected for efficiency, estimated with $N_{\text{ch}}^{\text{rec}} \approx 0.8 \cdot N_{\text{ch}}$
- Relative increase in multiplicity well in line with the VMD-Pb setup

Modelling γ +A with Pythia

[I. Helenius, M. Uthmeim: EPJC 84 (2024) 11, 1155]



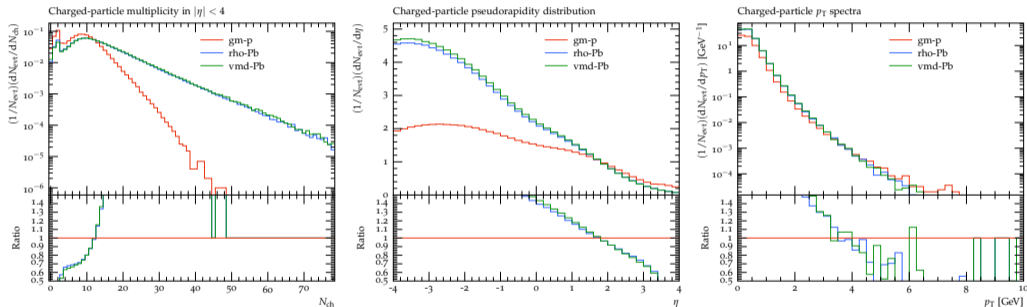
[ATLAS: PRC 104, 014903 (2021)]



- Multiplicity cut adjusted according to the limited efficiency
- Good description of the measured rapidity distribution with the VMD-Pb setup

Photoproduction on nuclear target at the EIC

- Min. bias events with $E_e = 18$ GeV and $E_n = 275$ GeV with $W_{\min} = 50$ GeV
- Compare results with proton and nuclear targets, latter modelled with VMD



- A similar increase of high-multiplicity events as in UPCs at the LHC
- More particles produced in the lead-going direction
- VMD: in 80 % of events the photon fluctuates into a ρ meson

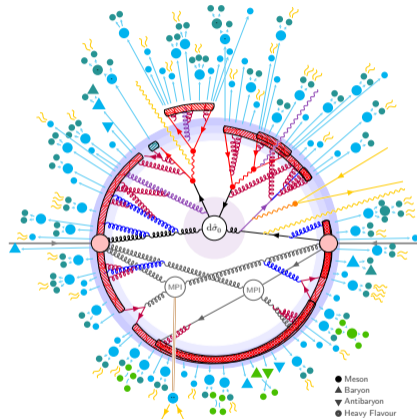
Summary & Outlook

Recent improvements

- Multi-jet merging in DIS
- A comparison study for photoproduction
- Baseline study for tuning
- VMD to model photonuclear collisions, in line with ATLAS UPC data

Ongoing efforts

- Tuning for DIS and photoproduction with HERA data
- Further improvements and validations for photoproduction with nuclear target



[figure by P. Skands]

- A satellite workshop in connection with the EICUG meeting on July 9-11 in JLab

Topics

- General-purpose event generators for electroproduction, photoproduction, and diffractive processes
- Comparisons and tuning to relevant ep and ed data
- Modeling of radiative effects
- Specialized event generators and interfacing

Organizers

- Frank Krauss
- Ilkka Helenius
- Markus Diefenthaler

Backup slides

Recent highlights

New parton shower Apollo

C.T. Preuss: JHEP 07 (2024) 161

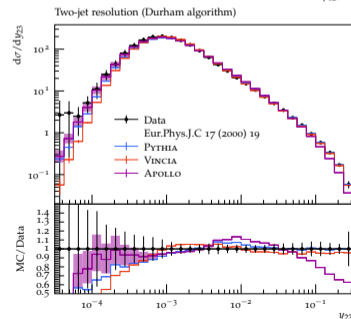
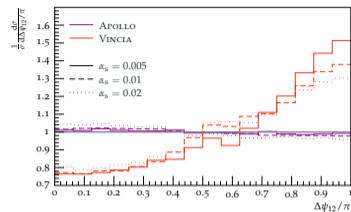
- Improved antenna shower heriting from Vincia
⇒ Easy to combine with fixed-order
- Improved recoil handling similar to Alaric
⇒ First NLL accurate parton shower in Pythia
- Currently only for e^+e^-

Machine-learning based hadronization

C. Bierlich, P. Ilten, S. Mrenna et al. (ML-HAD):

SciPost Phys. 17 (2024) 2, 045, arXiv:2410.06342

- Learn fragmentation functions from data
- Currently tested in a simplified $q\bar{q}$ case



Multiparton interactions (MPIs)

- MPIs from $2 \rightarrow 2$ QCD cross sections

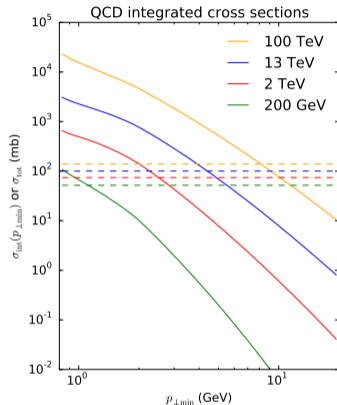
$$\frac{d\mathcal{P}_{\text{MPI}}}{dp_T^2} = \frac{1}{\sigma_{\text{nd}}(\sqrt{s})} \frac{d\sigma^{2 \rightarrow 2}}{dp_T^2}$$

$\sigma_{\text{nd}}(\sqrt{s})$ is the non-diffractive cross section

- Partonic cross section diverges at $p_T \rightarrow 0$
 \Rightarrow Introduce a screening parameter p_{T0}

$$\frac{d\sigma^{2 \rightarrow 2}}{dp_T^2} \propto \frac{\alpha_s(p_T^2)}{p_T^4} \rightarrow \frac{\alpha_s(p_{T0}^2 + p_T^2)}{(p_{T0}^2 + p_T^2)^2}$$

- Energy-dependent parametrization:
 $p_{T0}(\sqrt{s}) = p_{T0}^{\text{ref}} (\sqrt{s}/\sqrt{s_{\text{ref}}})^\alpha$
- Number of interactions: $\langle n \rangle = \sigma_{\text{int}}(p_{T0})/\sigma_{\text{nd}}$



- $\sigma_{\text{int}}(p_{T,\text{min}})$ exceeds σ_{tot}
 \Rightarrow Several interactions

Parton-level evolution

Common evolution scale (p_T) for FSR, ISR and MPIs

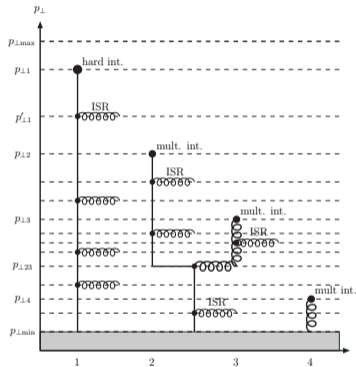
- Probability for something to happen at given p_T

$$\frac{d\mathcal{P}}{dp_T} = \left(\frac{d\mathcal{P}_{\text{MPI}}}{dp_T} + \sum \frac{d\mathcal{P}_{\text{ISR}}}{dp_T} + \sum \frac{d\mathcal{P}_{\text{FSR}}}{dp_T} \right) \times \exp \left[- \int_{p_T}^{p_T^{\text{max}}} dp_T' \left(\frac{d\mathcal{P}_{\text{MPI}}}{dp_T'} + \sum \frac{d\mathcal{P}_{\text{ISR}}}{dp_T'} + \sum \frac{d\mathcal{P}_{\text{FSR}}}{dp_T'} \right) \right]$$

where $\exp[. . .]$ is a Sudakov factor (probability that nothing else has happened before p_T)

Simultaneous partonic evolution

1. Start the evolution from the hard-process scale
2. Sample p_T for each \mathcal{P}_i , pick one with highest p_T
3. Continue until $p_{T\text{min}} \sim \Lambda_{\text{QCD}}$ reached



[T. Sjöstrand, P. Skands:
EPJC 39 (2005) 129-154]

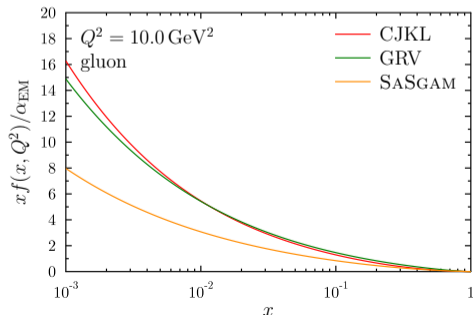
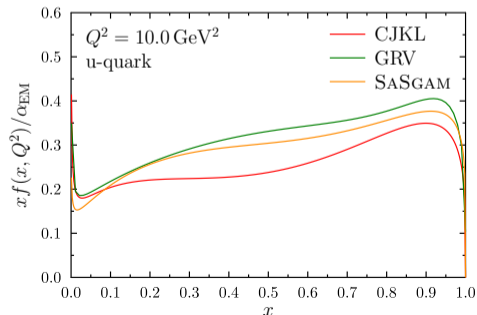
PDFs for resolved photons

DGLAP equation for photons

- Additional term due to $\gamma \rightarrow q\bar{q}$ splittings

$$\frac{\partial f_i^\gamma(x, Q^2)}{\partial \log(Q^2)} = \frac{\alpha_{\text{em}}}{2\pi} e_i^2 P_{i\gamma}(x) + \frac{\alpha_s(Q^2)}{2\pi} \sum_j \int_x^1 \frac{dz}{z} P_{ij}(z) f_j(x/z, Q^2)$$

where $P_{i\gamma}(x) = 3(x^2 + (1-x)^2)$ for quarks, 0 for gluons (LO)



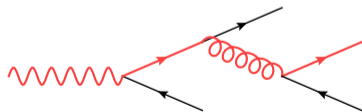
Evolution equation and ISR for resolved photons

ISR probability based on DGLAP evolution

- Add a term corresponding to $\gamma \rightarrow q\bar{q}$ to (conditional) ISR probability

$$d\mathcal{P}_{a \leftarrow b} = \frac{dQ^2}{Q^2} \frac{\alpha_s}{2\pi} \frac{x' f_a^\gamma(x', Q^2)}{x f_b^\gamma(x, Q^2)} P_{a \rightarrow bc}(z) dz + \frac{dQ^2}{Q^2} \frac{\alpha_{em}}{2\pi} \frac{e_b^2 P_{\gamma \rightarrow bc}(x)}{f_b^\gamma(x, Q^2)}$$

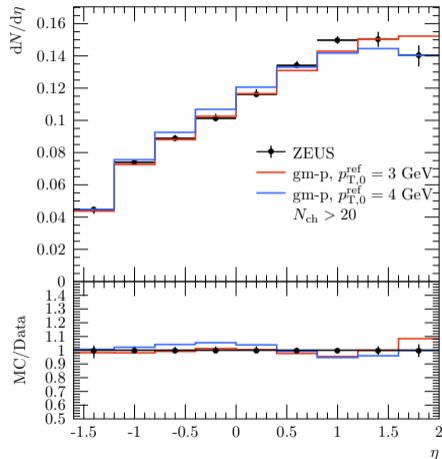
- Corresponds to ending up to the beam photon during evolution
 - ⇒ Parton originated from the point-like (anomalous) part of the PDFs
 - No further ISR or MPIs below the scale of the splitting
 - Implemented for the default Simple Shower in Pythia 8



Comparison to ZEUS data for charged hadrons ($N_{\text{ch}} > 20$)

Pseudorapidity

- Data well reproduced
- Not sensitive to MPI modelling ($p_{\text{T},0}$)



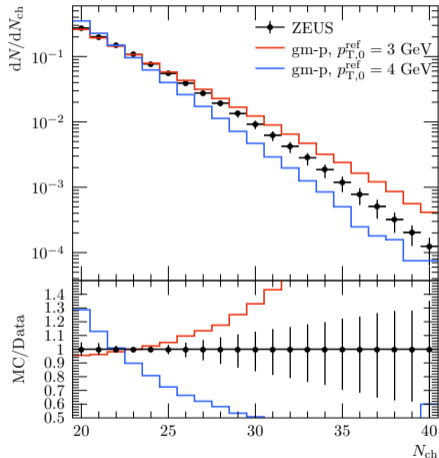
Comparison to ZEUS data for charged hadrons ($N_{ch} > 20$)

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Multiplicity

- Sensitivity to MPI parameters, clear support for MPIs
- Data within $p_{T,0}$ variations



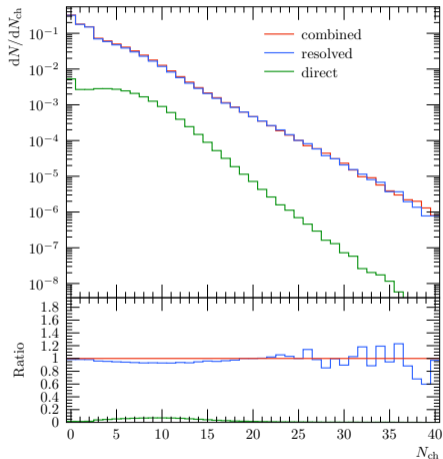
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Pseudorapidity

- Data well reproduced
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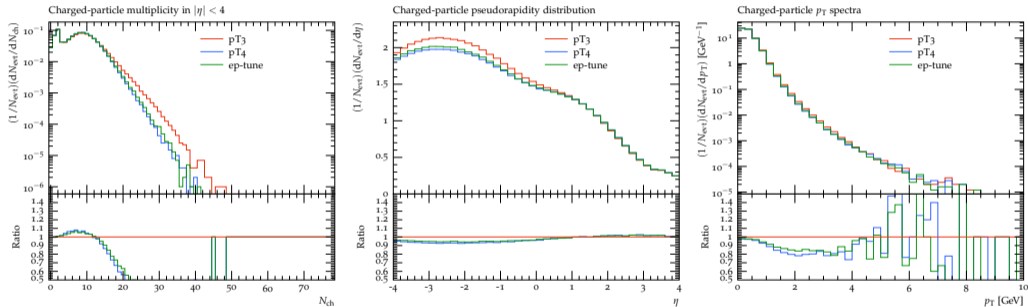
Multiplicity

- Sensitivity to MPI parameters, clear support for MPIs
- Data within $p_{\text{T},0}$ variations
- Direct contribution negligible in high-multiplicity events ($N_{\text{ch}} > 20$)



Photoproduction on proton target at the EIC

- Min. bias events with $E_e = 18$ GeV and $E_n = 275$ GeV with $W_{\min} = 50$ GeV
- Compare results with proton and nuclear targets, latter modelled with VMD



- A similar increase of high-multiplicity events as in UPCs at the LHC
- More particles produced in the lead-going direction
- Only modest effects to the p_T spectra

Comparisons between Pythia, Sherpa and Herwig

[I. Helenius, P. Meizinger, S. Plätzer, P. Richardson: arXiv:2406.08026 [hep-ph]]

- Summary of the modelling differences between the generators

Property	Pythia	Sherpa	Herwig
Flux	LL	NLL	LL
$\alpha_s(M_Z^2)$	0.130, 1-loop running	0.118, 3-loop running	
PDFs	CJKL	SAS2M	SAS2M
Remnants	forced splittings/PS rejection	PS rejection	forced splitting
$\gamma \rightarrow q\bar{q}$ Splitting	yes	no	no
MPI tuning	preliminary γ -p/ γ - γ tune	untuned	untuned

Experimental heavy-ion UPC classification

- Event selection typically relies on Zero-degree calorimeters ($X > 0$)

XnXn: At least one neutron on both sides
⇒ A+A (hadronic interaction)

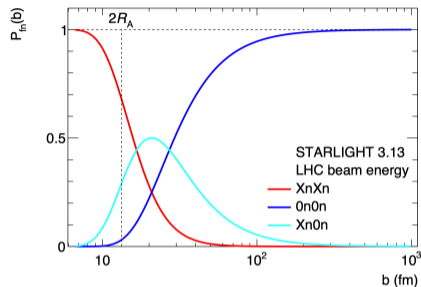
Xn0n: At least one neutron only on one side
⇒ γ +A

0n0n: No neutrons on either side
⇒ γ + γ

Possible caveats

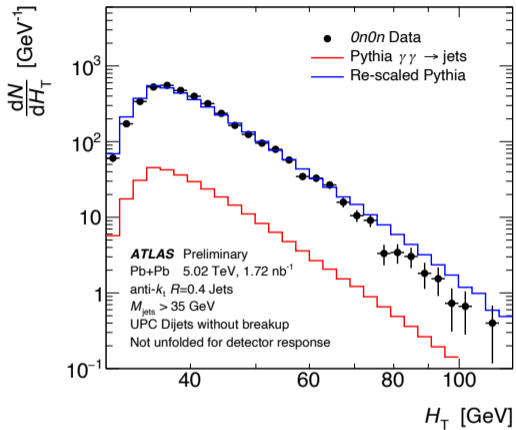
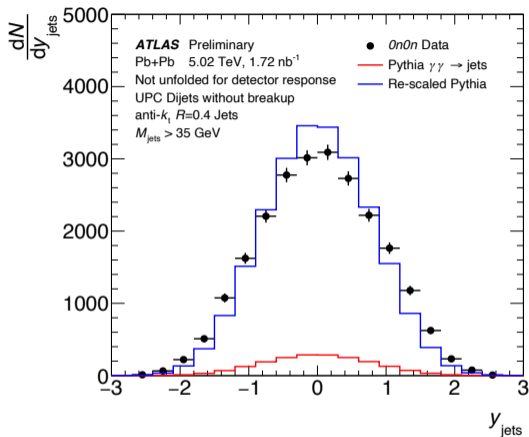
- Additional EM interactions may break up the nuclei in “near-encounter” events
[Eskola, Guzey, Helenius, Paakkinen, Paukkunen; PRC 110 (2024) 054906]
- Also diffractive processes will keep nuclei intact
⇒ Xn0n condition will remove diffractive contribution to γ +A

See e.g. [Guzey, Klasen; PRD 104 (2021) 11 114013]



Ann.Rev.Nucl.Part.Sci. 70 (2020) 323-354

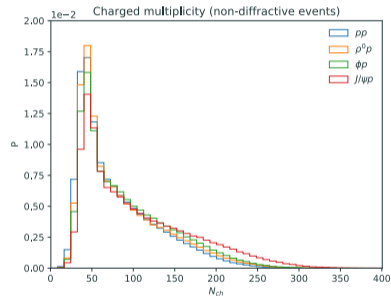
Dijets in ultra-peripheral heavy-ion collisions in 0n0n



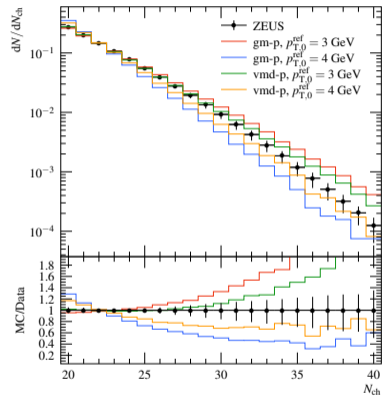
- Per-event yield underestimated by a factor of ten!
- Shape in a reasonable agreement
- $\gamma\gamma \rightarrow \mu^+\mu^-$ ok so likely a QCD effect \Rightarrow Contribution from diffractive events?

[ATLAS-CONF-2022-021]

- Resolved contribution dominates total cross section
- ⇒ Set up an explicit VMD model with linear combination of vector-meson states (ρ , ω , ϕ and J/ψ)
- Use VM PDFs from SU21
[Sjöstrand, Utheim; EPJC 82 (2022) 1, 21]
- Cross sections from SaS
[Schuler, Sjöstrand; PRD 49 (1994) 2257-2267]
- Sample collision energy from flux
- ⇒ Vector meson-proton scatterings

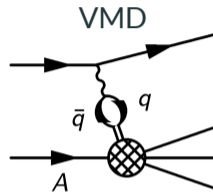
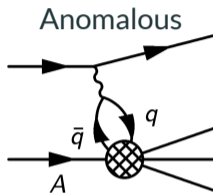
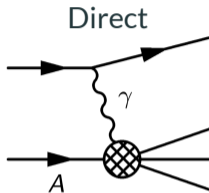


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- Cross sections from SaS
[Schuler, Sjöstrand; PRD 49 (1994) 2257-2267]
- Sample collision energy from flux
- ⇒ Vector meson-proton scatterings
- In line with the full photoproduction



[ZEUS: JHEP 12 (2021) 102]

Vector meson dominance (VMD)



Linear combination of three components

$$|\gamma\rangle = c_{\text{dir}}|\gamma_{\text{dir}}\rangle + \sum_q c_q|q\bar{q}\rangle + \sum_V c_V|V\rangle$$

where the last term includes a linear combination of vector meson states up to J/ψ

$$c_V = \frac{4\pi\alpha_{\text{EM}}}{f_V^2}$$

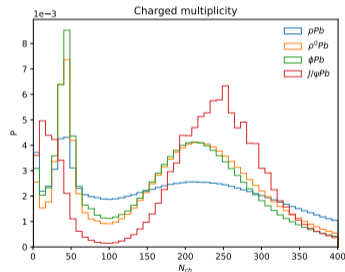
V	$f_V^2/(4\pi)$
ρ^0	2.20
ω	23.6
ϕ	18.4
J/ψ	11.5

Modelling $\gamma+A$ with Pythia

[I. Helenius, M. Uthmeim: arXiv:2406.10403 [hep-ph]; Accepted for publication in EPJC]

Angantyr model for heavy ions in Pythia

- Monte Carlo Glauber to sample nucleon configurations
 - Cross section fluctuations, fitted to partial nucleon-nucleon cross sections
 - Secondary (wounded) collisions as diffractive excitations
 - Can now handle generic hadron-ion and varying energy
- ⇒ VMD-nucleus scatterings

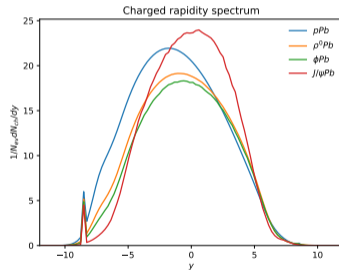


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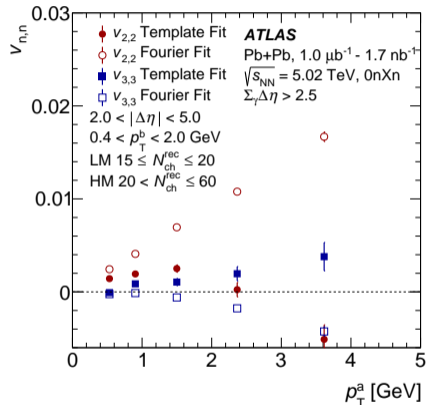
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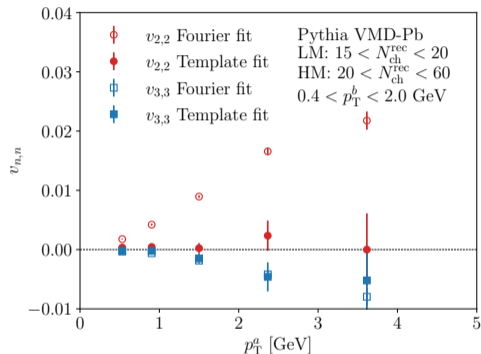


Two-particle correlations in $\gamma+A$ with Pythia

[ATLAS: PRC 104, 014903 (2021)]



[I. Helenius, M. Uthm: EPJC 84 (2024) 11, 1155]

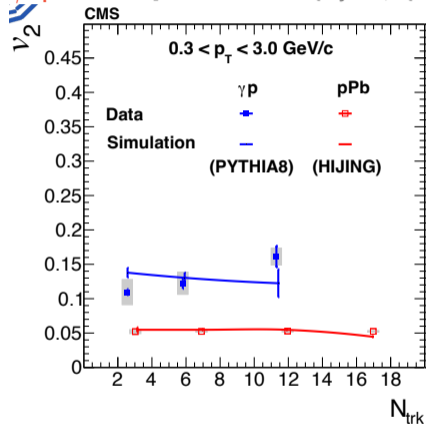


- No finite v_2 left after template fit in the Pythia simulation
⇒ Revisit with final state effects such as rope hadronization and string showing

Collectivity in UPCs at the LHC

$\gamma+p$

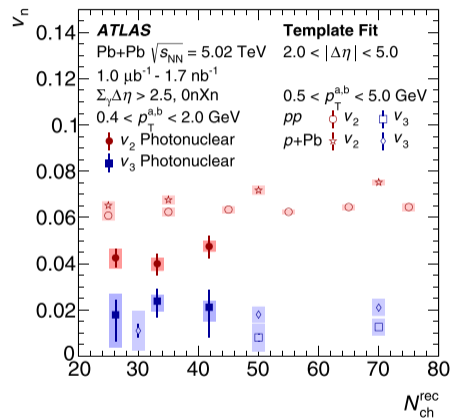
[CMS: Murillo Quijada, QM2022]



- Finite v_2 for $\gamma+p$, in line with Pythia
 \Rightarrow Jet-like correlations?

$\gamma+Pb$

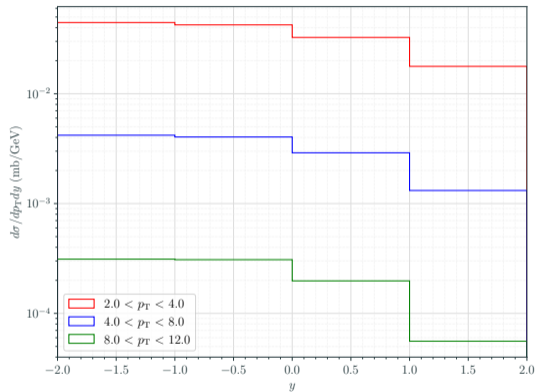
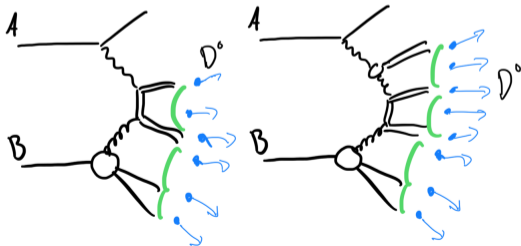
[ATLAS: PRC 104, 014903 (2021)]



- Finite v_n also after Template fit
 subtracting “non-flow”

Inclusive D-meson production in UPCs

- New experimental analysis for open charm production in UPCs ongoing in CMS and ALICE
- Can use Pythia UPC implementation to calculate cross-section predictions



[A.-M. Levälampi: Research training thesis, 2024]