



Toward extracting scattering phase shift from integrated correlation function in lattice QCD

Peng Guo

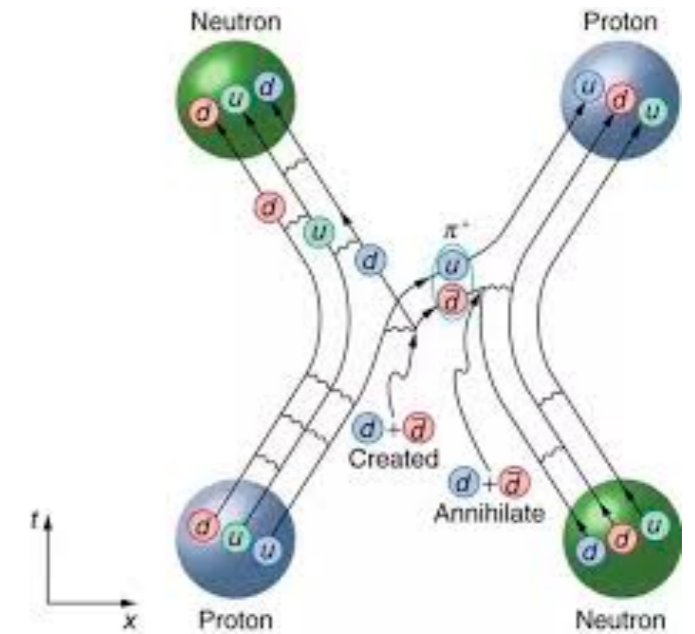
Dakota State University



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Topical Group on Hadronic
Physics, Anaheim, CA
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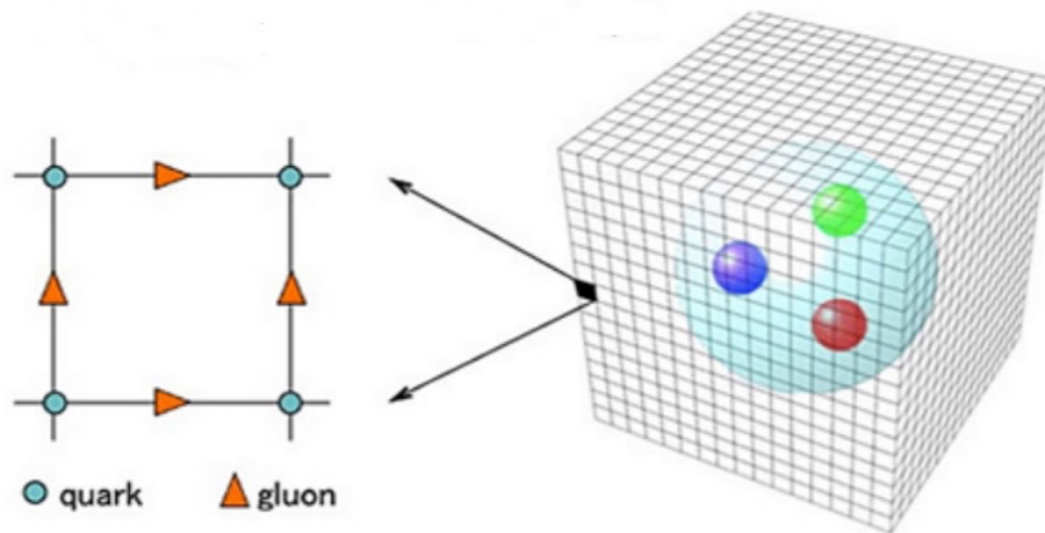
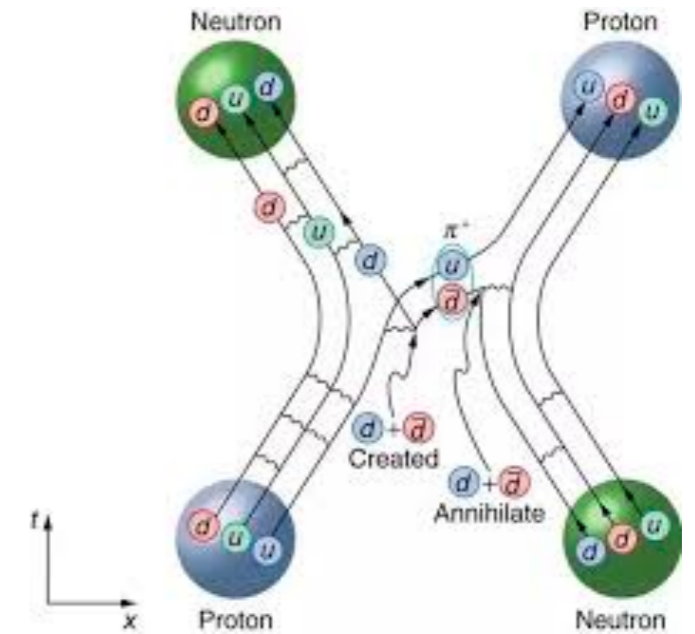
Major tasks in Nuclear/Hadron physics:

- Understand and extract few-body dynamics, such as scattering phase shift
- Understand resonances and bound states



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- Understand and extract few-body dynamics, such as scattering phase shift
- Understand resonances and bound states



☑ Periodic box in LQCD

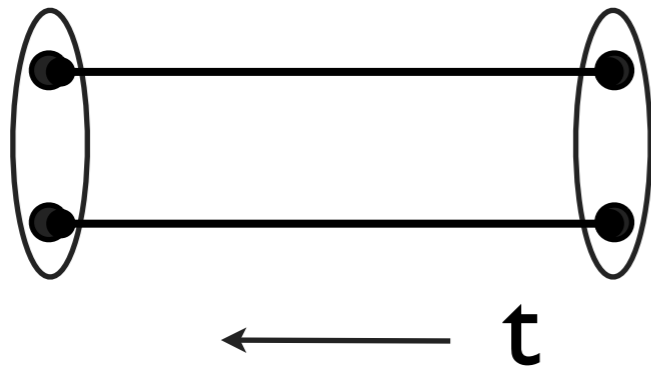


Lattice QCD: one route to non-perturbative dynamics



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$$\langle 0 | \mathcal{O}(t) \mathcal{O}^\dagger(0) | 0 \rangle$$

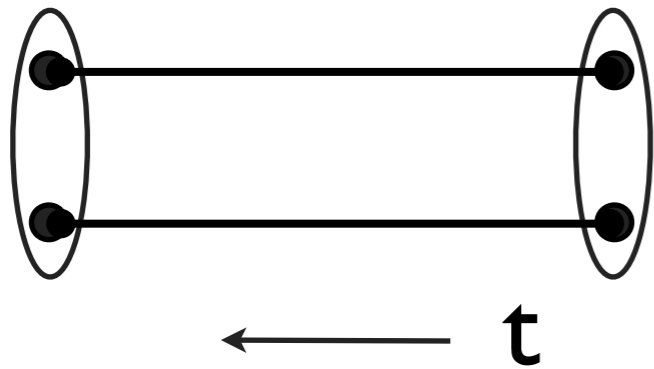




Lattice QCD: one route to non-perturbative dynamics

$$\langle 0 | \mathcal{O}(t) \mathcal{O}^\dagger(0) | 0 \rangle \propto \int \mathcal{D}[U] \mathcal{D}[\psi, \bar{\psi}] e^{-S_G[U] - S_F[\psi, \bar{\psi}, U]} \mathcal{O}(t) \mathcal{O}^\dagger(0)$$

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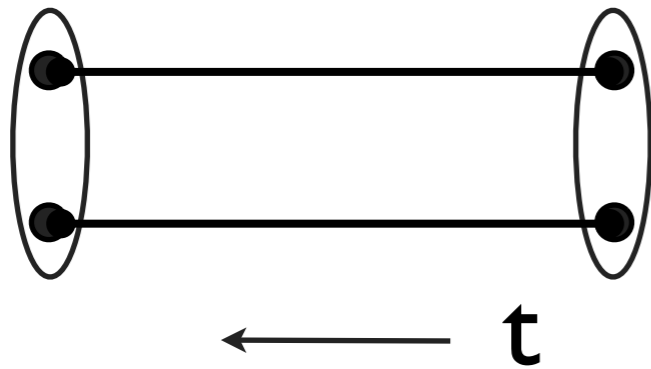


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$$S_G[U] + S_F[\psi, \bar{\psi}, U] \rightarrow \int d^4x \bar{\psi}(x) [\gamma_\mu (\partial_\mu + igA_\mu(x)) + m] \psi(x) + \frac{1}{2} \int d^4x \text{tr}[F_{\mu\nu}(x) F_{\mu\nu}(x)]$$

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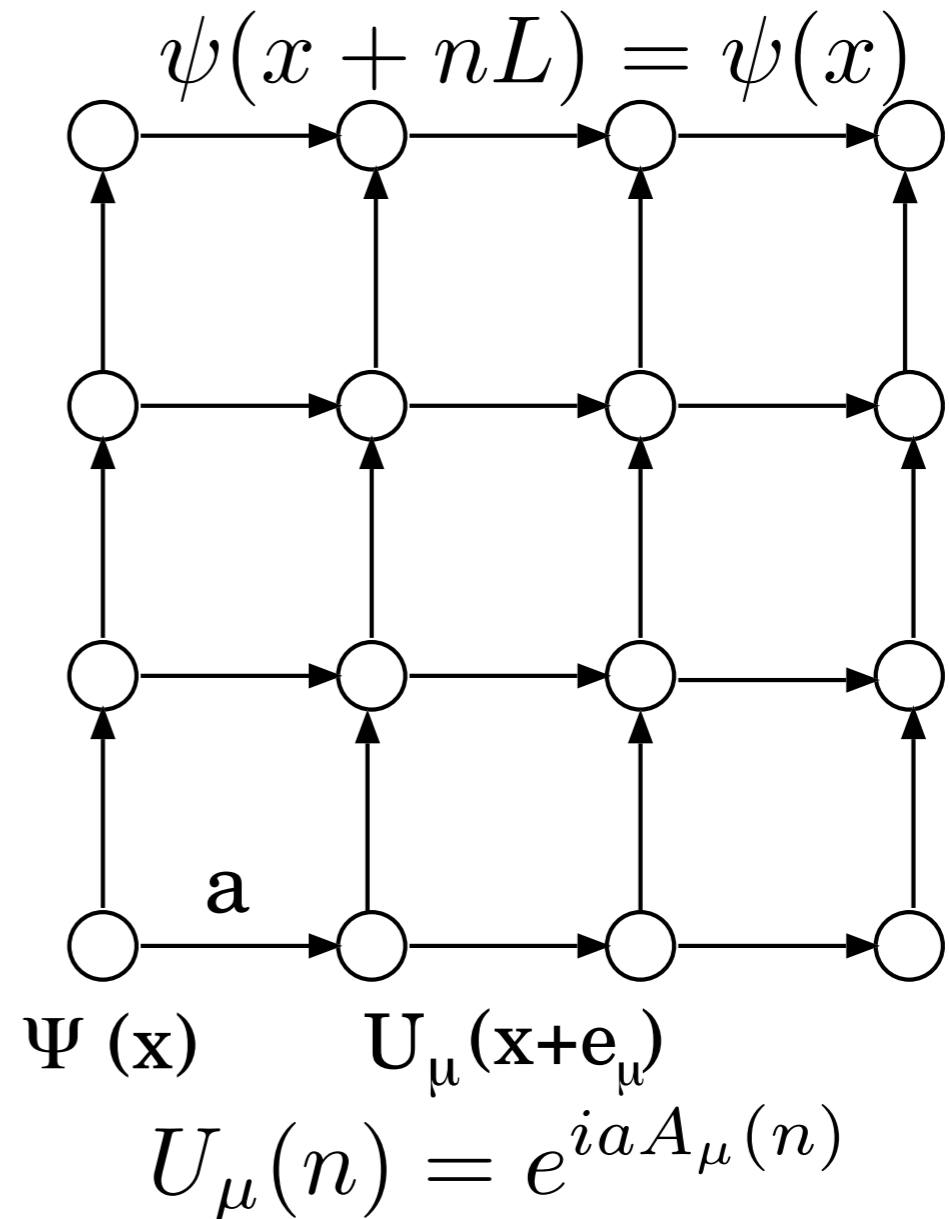
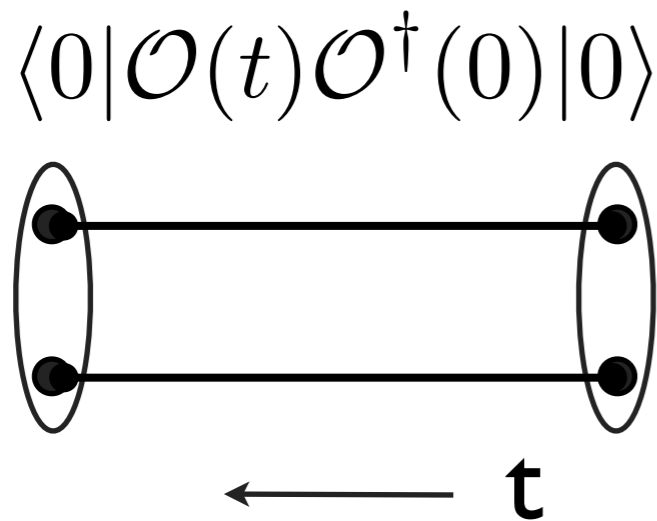




Lattice QCD: one route to non-perturbative dynamics

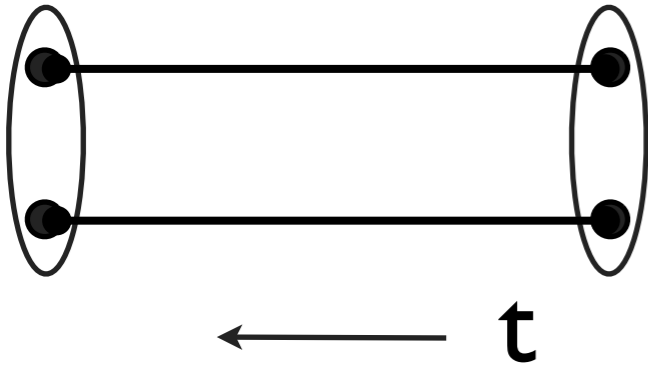
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Spectroscopy from Lattice QCD

$$\langle 0 | \mathcal{O}(t) \mathcal{O}^\dagger(0) | 0 \rangle$$

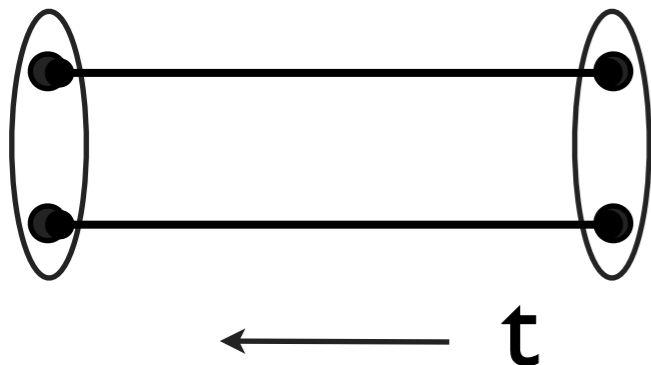


Spectroscopy from Lattice QCD

$$\mathcal{O}(t) = e^{t\hat{H}} \mathcal{O}(0) e^{-t\hat{H}}$$



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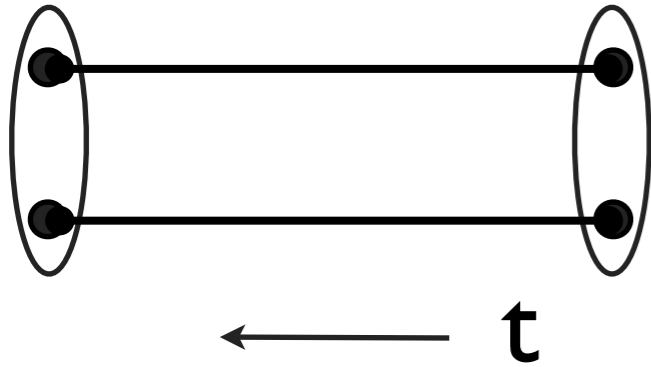


Spectroscopy from Lattice QCD

$$\mathcal{O}(t) = e^{t\hat{H}} \mathcal{O}(0) e^{-t\hat{H}}$$

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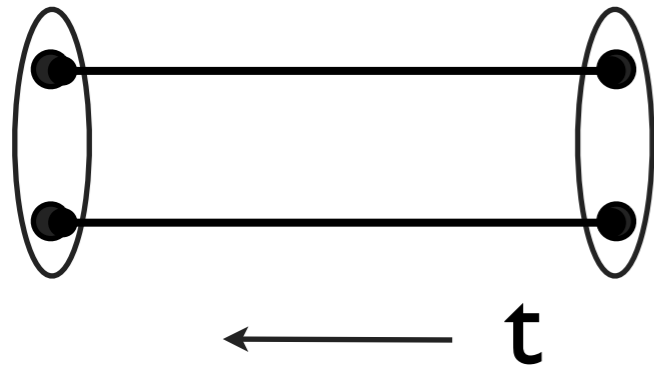


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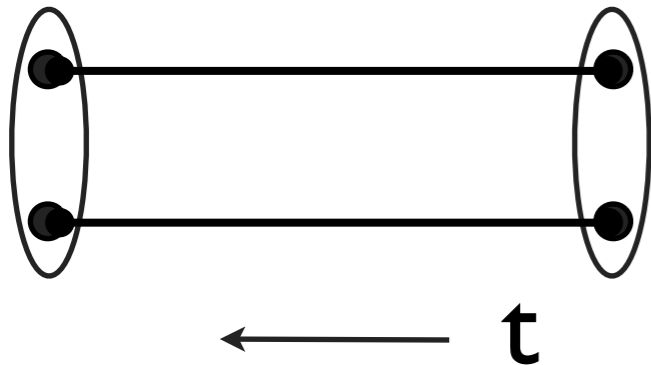
$$\sim \sum_n z_n^* z_n e^{-E_n t}, \quad z_n = \langle n | \mathcal{O}^\dagger(0) | 0 \rangle$$

Spectroscopy from Lattice QCD

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(Note: Red arrows point from the $e^{t\hat{H}}$ and $\sum_n |n\rangle\langle n| = 1$ terms in the equation above to the $\mathcal{O}(t)$ and $\mathcal{O}^\dagger(0)$ terms in this equation, respectively.)



$$\sim \sum_n z_n^* z_n e^{-E_n t}, \quad z_n = \langle n | \mathcal{O}^\dagger(0) | 0 \rangle$$

$$\rightarrow |z_0|^2 e^{-E_0 t}$$

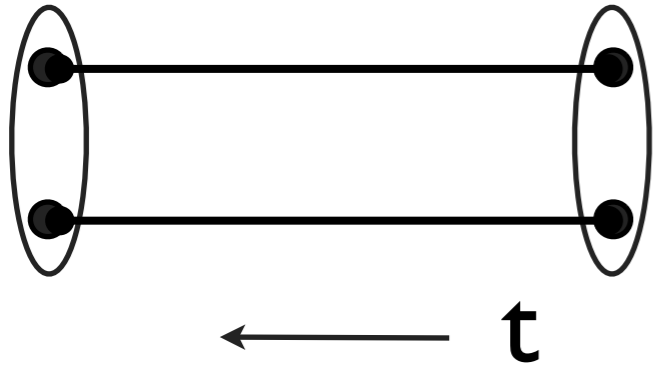
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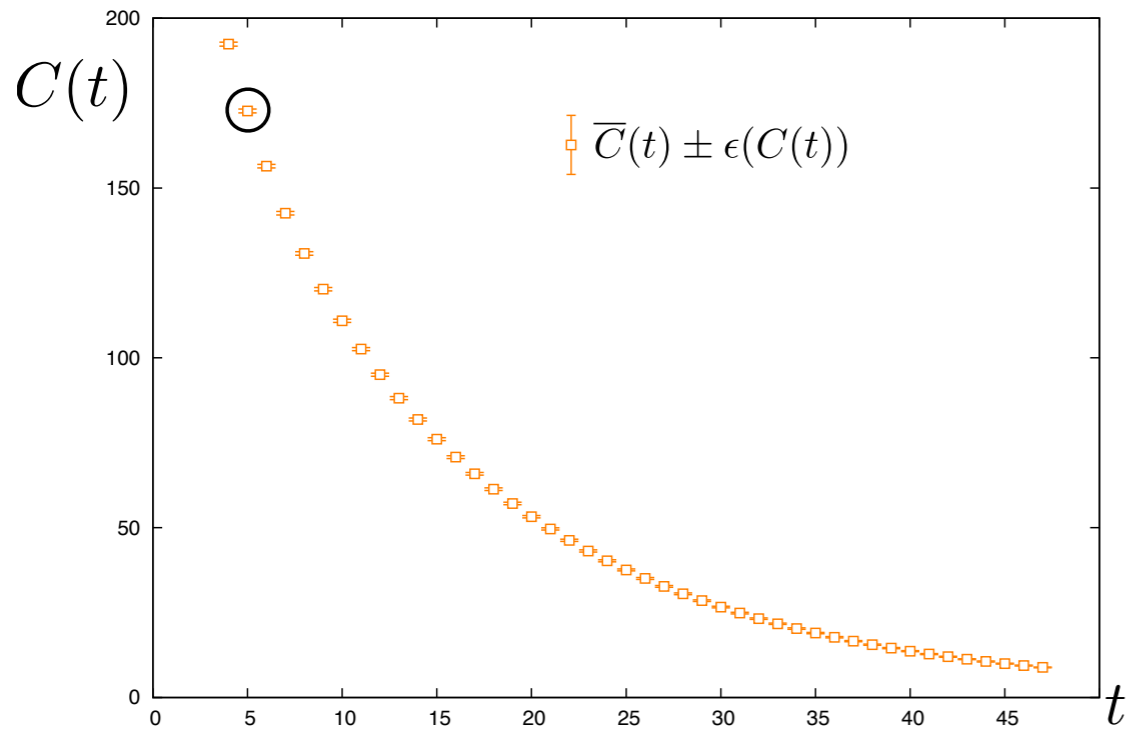
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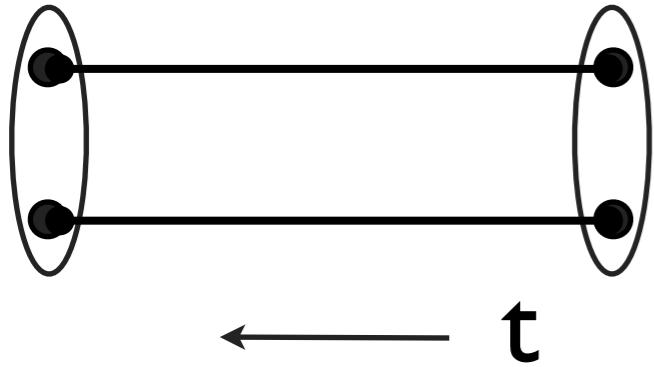


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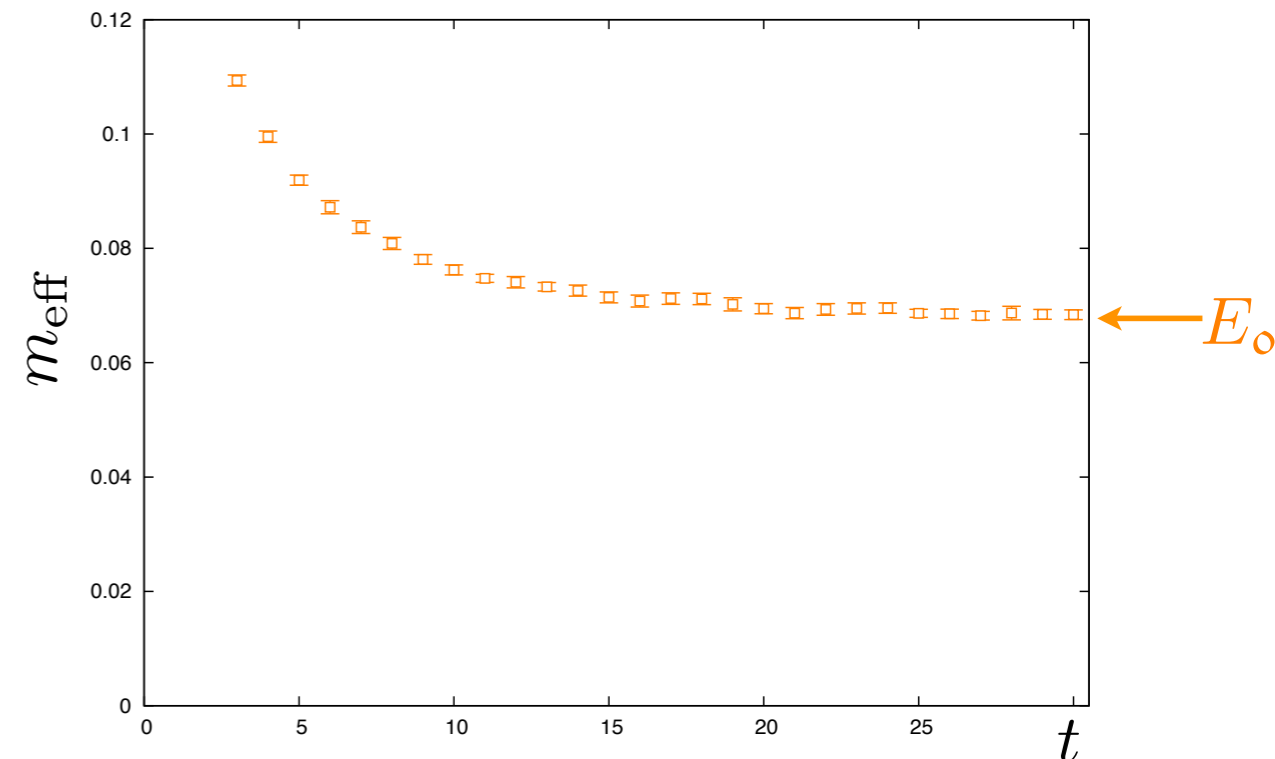
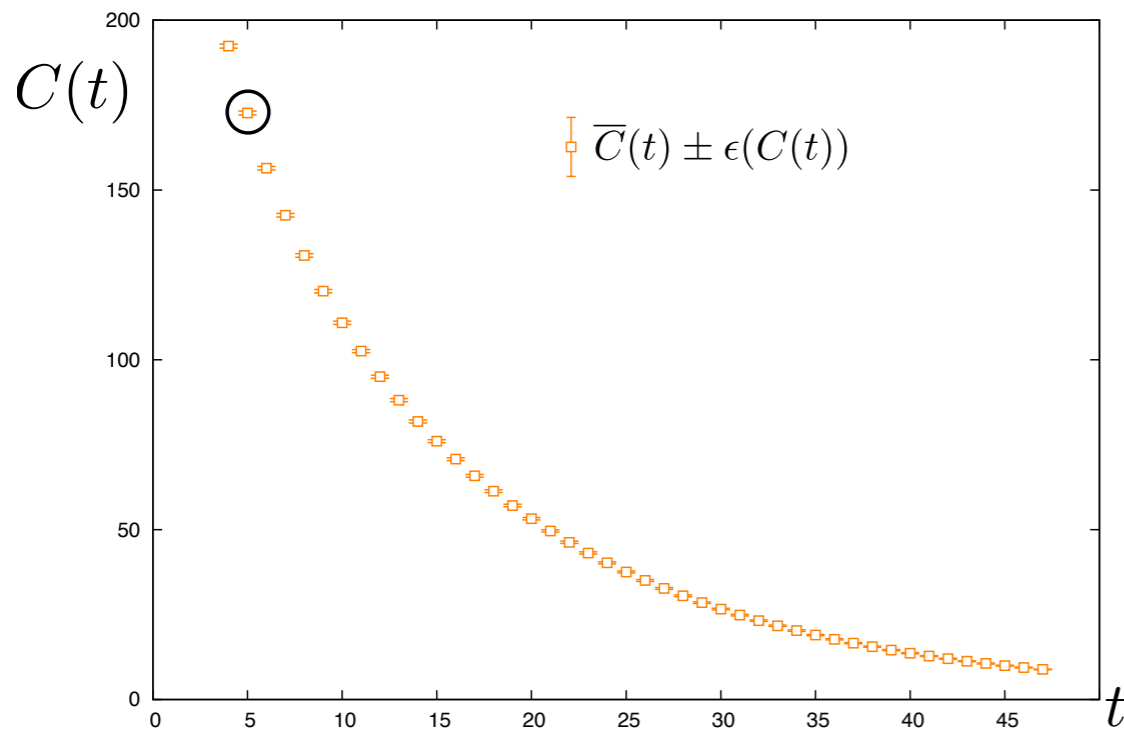
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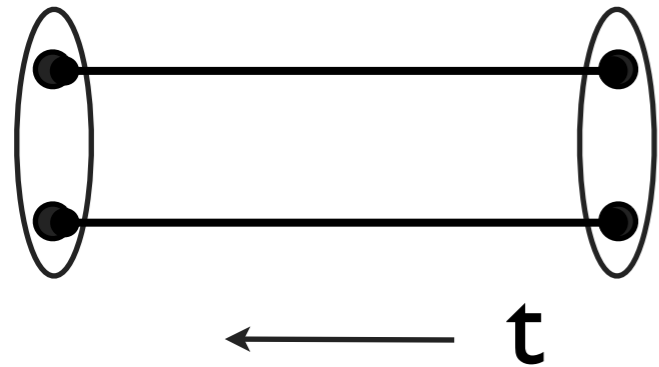
$$m_{\text{eff}}(t) = -\frac{1}{\delta t} \ln \frac{C(t + \delta t)}{C(t)}$$



Spectroscopy from Lattice QCD

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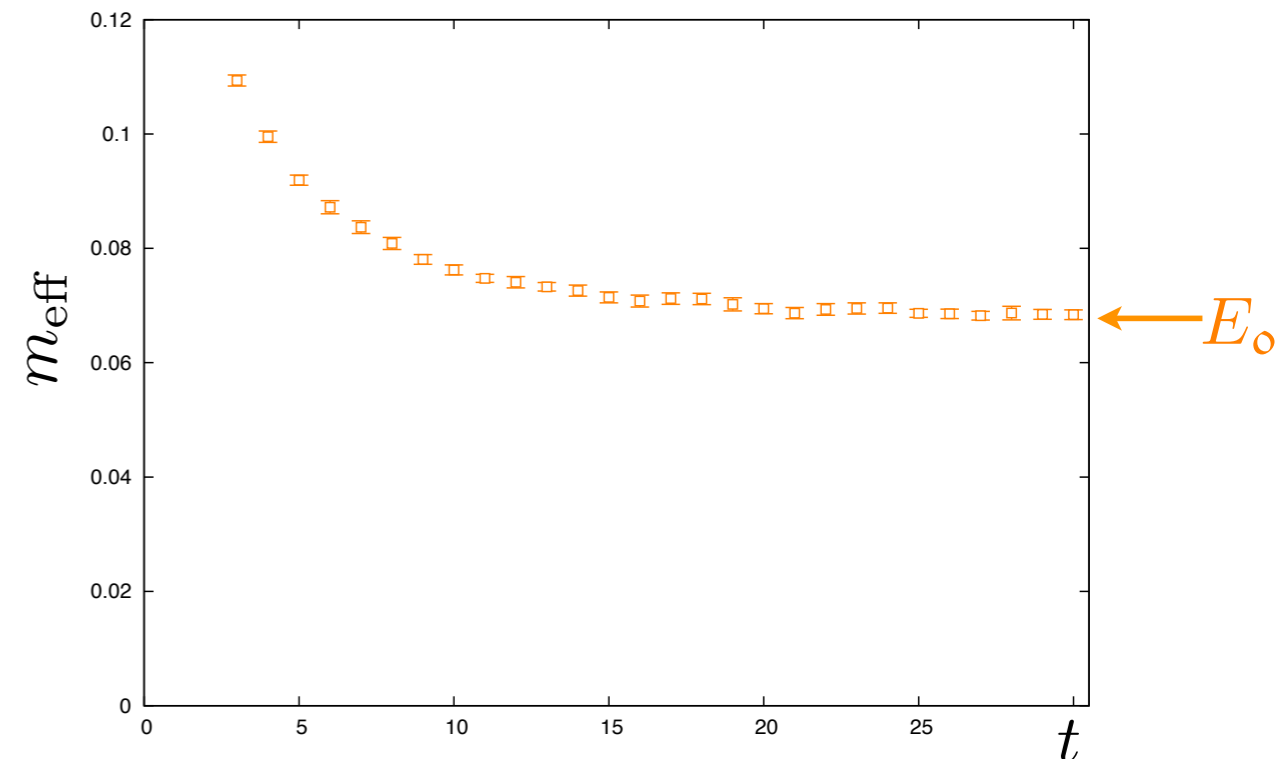
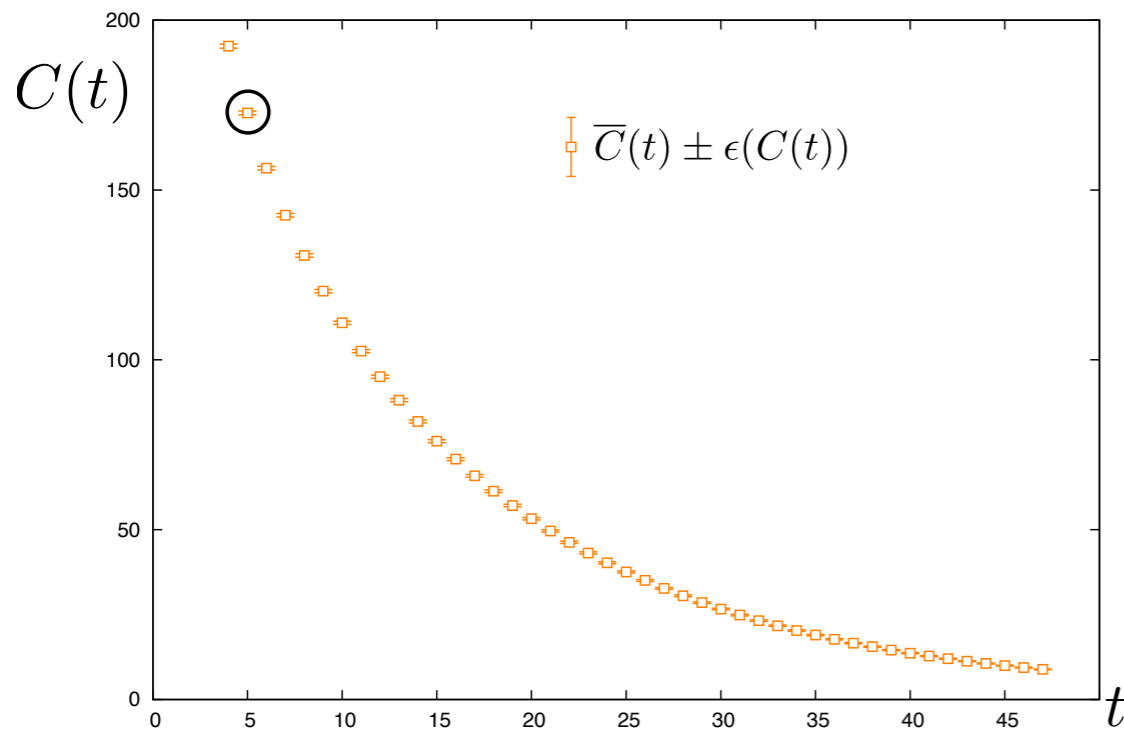
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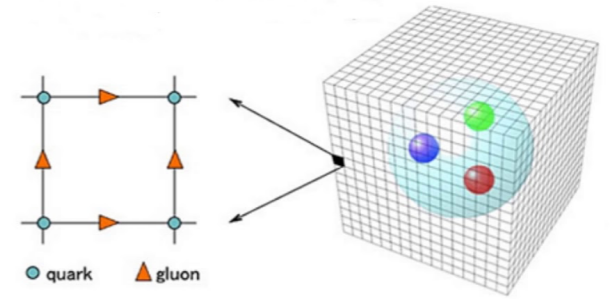
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Few-body interaction in a box:

Challenges:

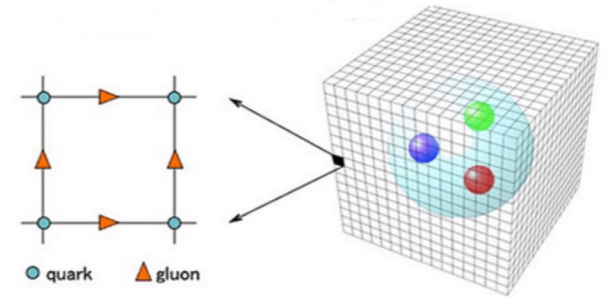
- ☑ No asymptotic states
- ☑ Stationary solutions instead of scattering solutions
- ☑ Discrete energies



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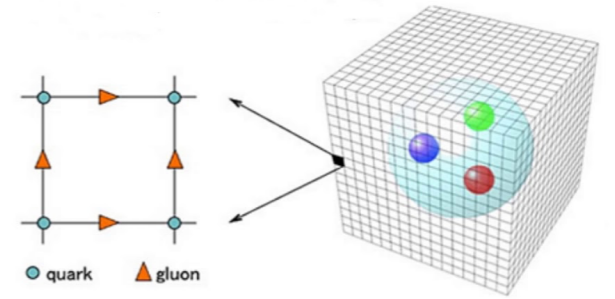


🌐 Extracting Two-body dynamics from discrete energy levels

Few-body interaction in a box:

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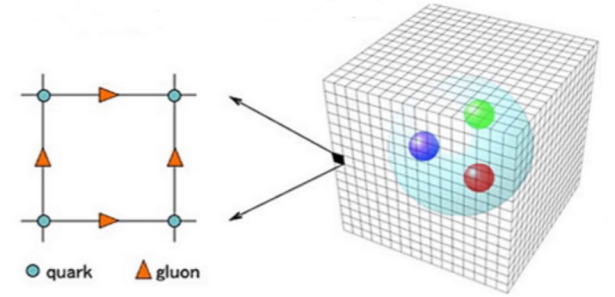


Extracting Two-body dynamics from discrete energy levels

- ✱ Lusecher formula-like QC as result of factorization of long-range effect and short-range dynamics

$$\det [\cot \delta(E) - \mathcal{M}(E)] = 0$$

Few-body interaction in a box:



Challenges:

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- ✓ Stationary solutions instead of scattering solutions
- ✓ Discrete energies



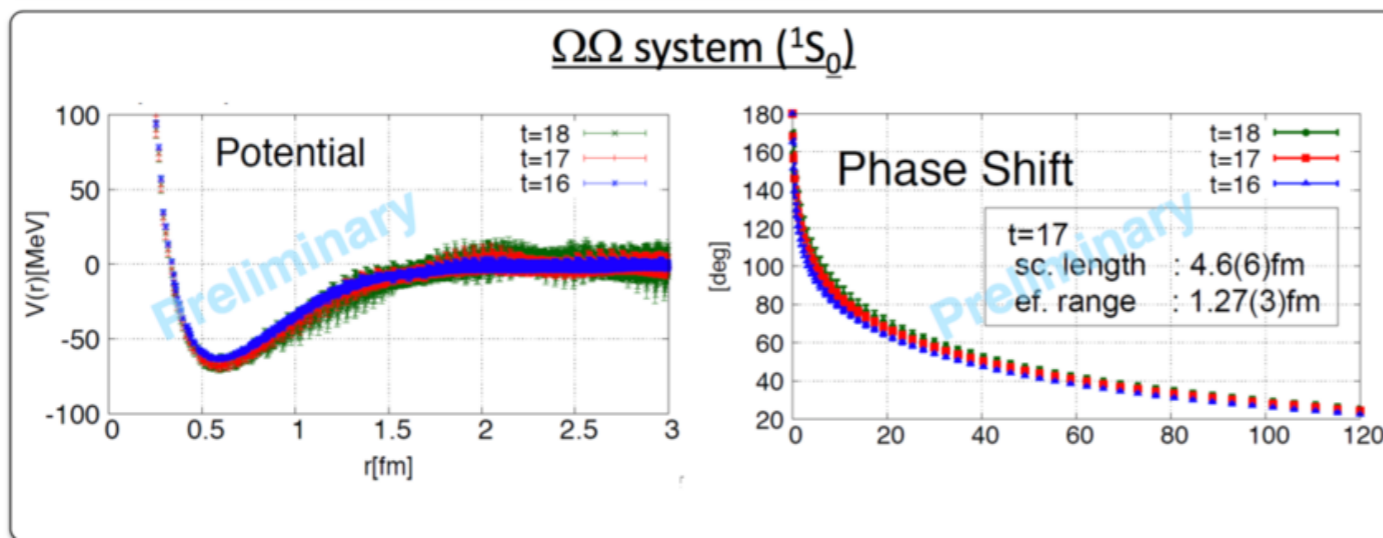
Extracting Two-body dynamics from discrete energy levels

- * Lusecher formula-like QC as result of factorization of long-range effect and short-range dynamics

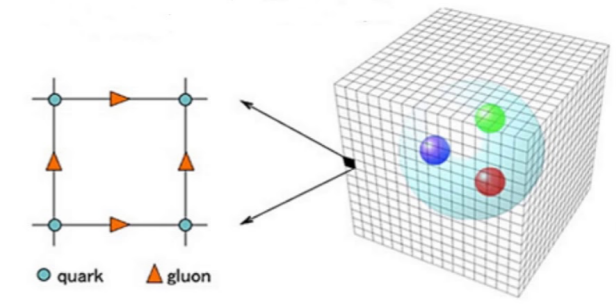
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- * HAL QCD collaboration potential method

$N_f = 2 + 1, m_\pi = 146 \text{ MeV}, a \approx 0.085 \text{ fm}, L \approx 8 \text{ fm}$



NN controversy:



- Two-Nucleon correlation function suffer Signal-to-noise ratio issue:

$$\mathcal{R}(t) \xrightarrow{t \rightarrow \infty} e^{-(m_N - \frac{3}{2}m_\pi)t}$$

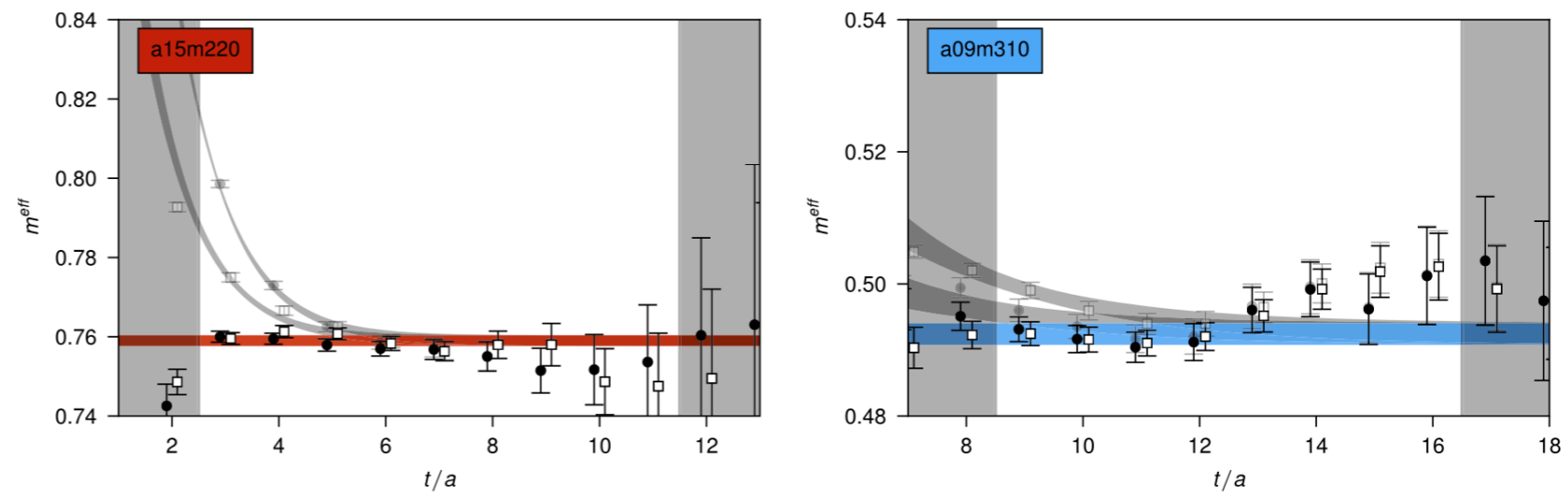


Fig. 1. Effective mass plots of the nucleon from Ref. [64] which suffer from correlated, late time fluctuations, making it more challenging to identify the ground state. The top plot is from a calculation with $a \sim 0.15$ fm and $m_\pi \sim 220$ MeV while the bottom is for $a \sim 0.09$ fm and $m_\pi \sim 310$ MeV.

Progress in Particle and Nuclear Physics 121 (2021) 103888

- Lusecher formula and HAL QCD collaboration potential method disagree on whether or not two-nucleon form a bound state with pion masses as heavy as 800 MeV

Progress in Particle and Nuclear Physics 121 (2021) 103888

- Lusecher formula face difficulties at large volume limit due to increasing density of states

Alternative to Lusecher formula method:

- * The difference of integrated correlation functions between interacting and non-interacting systems approaches rapidly to its infinite volume limit which is related to scattering phase shift

$$C(t) - C_0(t) = \sum_n [e^{-\epsilon_n t} - e^{-\epsilon_n^{(0)} t}] \xrightarrow{L \rightarrow \infty} \frac{t}{\pi} \int_0^\infty d\epsilon \delta(\epsilon) e^{-\epsilon t}$$

where

$$C(t) = \int_0^L dr C(rt; r0)$$

$$C(rt; r'0) = \langle 0 | T \left[\hat{\mathcal{O}}_H(r, t) \hat{\mathcal{O}}_H^\dagger(r', 0) \right] | 0 \rangle$$

Two-particle creation operator

Alternative to Lusecher formula method:

✱ Exactly solvable model with contact interaction:
$$\left[-\frac{1}{2\mu} \frac{d^2}{dr^2} + V_0 \sum_{n \in \mathbb{Z}} \delta(r + nL) \right] \psi_\epsilon^{(\text{rel})}(r) = \epsilon \psi_\epsilon^{(\text{rel})}(r)$$

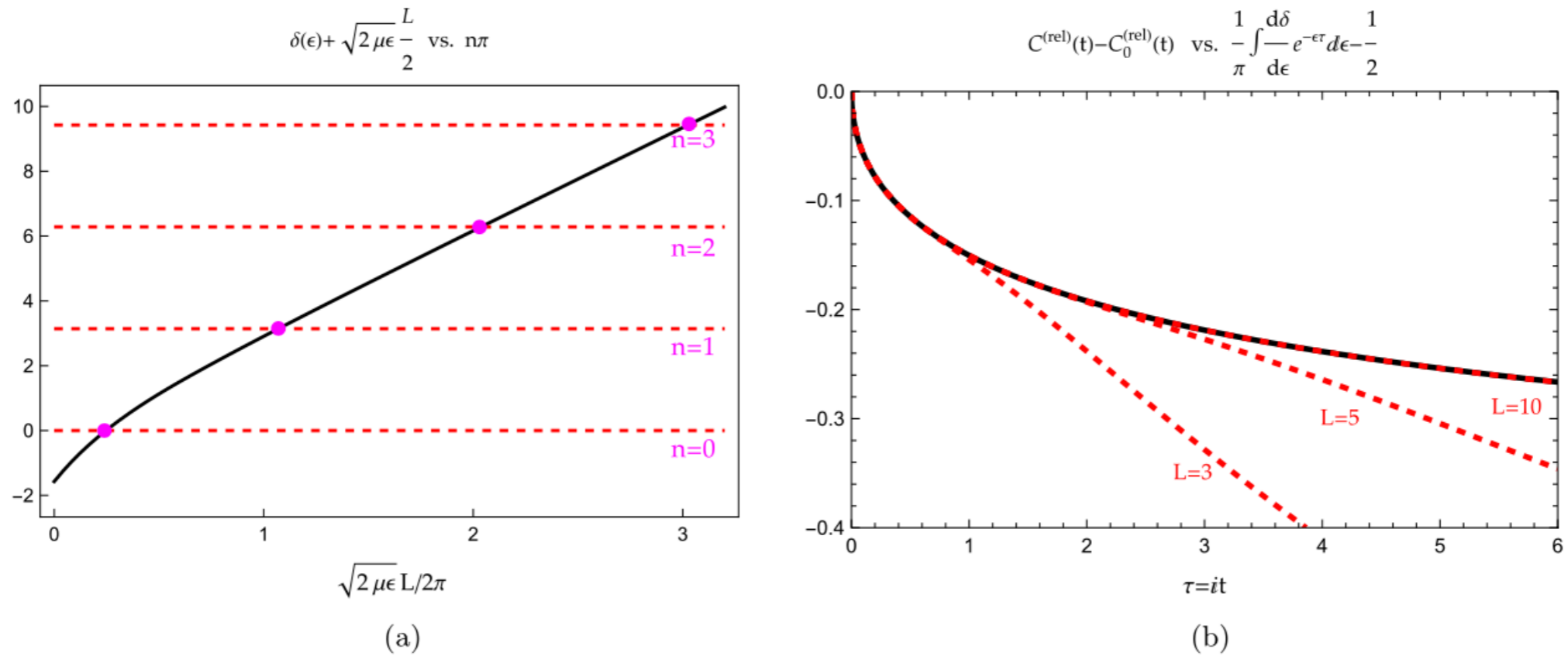
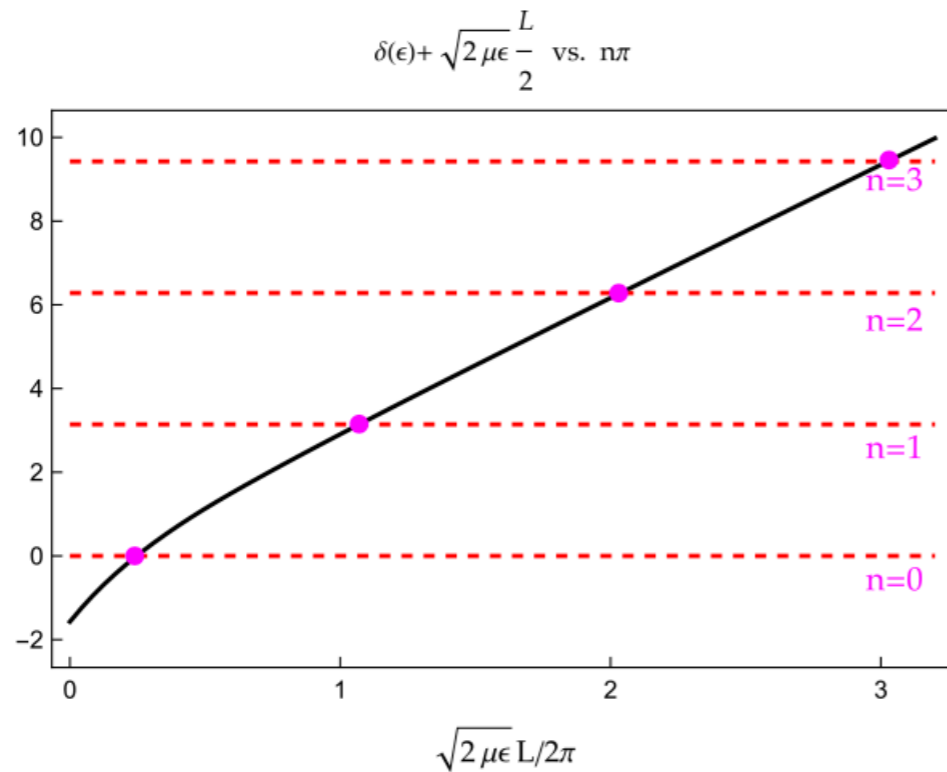


FIG. 1. The energy spectra and difference of integrated correlation function plots for particles interaction in a periodic box: (a) $\delta(\epsilon_n) + \sqrt{2\mu\epsilon_n} \frac{L}{2}$ (solid black) vs $n\pi$ (dashed red) with $L = 3$, energy spectra are located at intersection points of black and red curves; (b) $\frac{1}{\pi} \int_0^\infty d\epsilon \frac{d\delta(\epsilon)}{d\epsilon} e^{-\epsilon\tau} - \frac{1}{2}$ (solid black) vs $C^{(\text{rel})}(t) - C_0^{(\text{rel})}(t)$ (dashed red) with $L = 3, 5, 10$. The rest of parameters are taken as $V_0 = 0.5$ and $\mu = 1$.

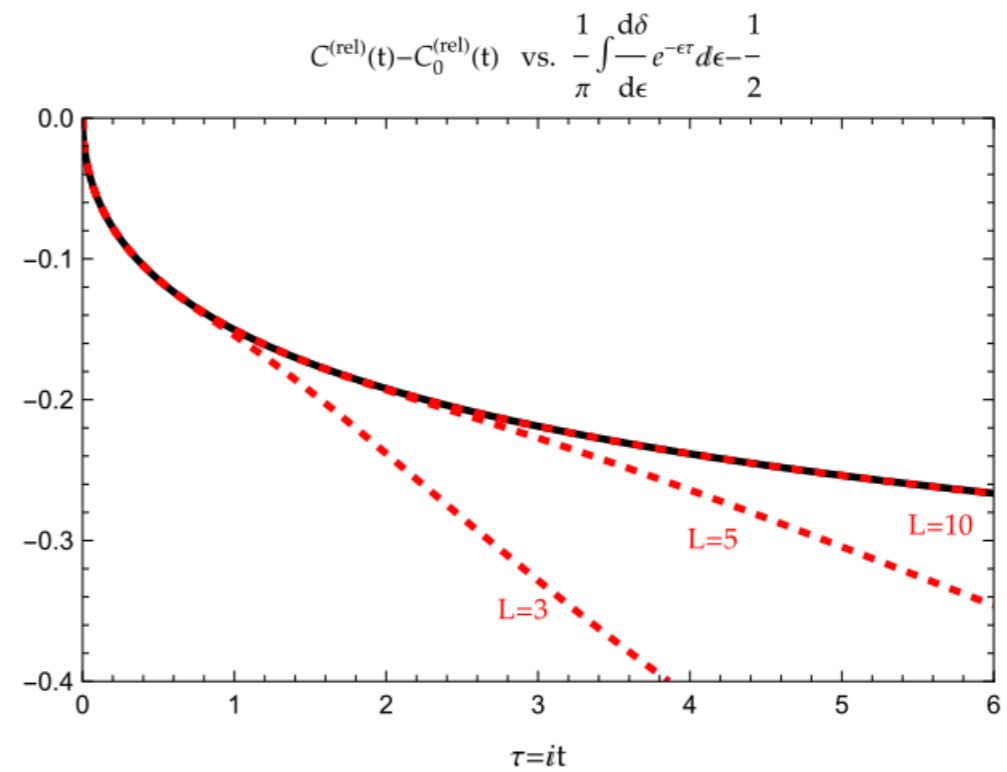
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$$\longrightarrow C(t) - C_0(t) \xrightarrow{L \rightarrow \infty} \frac{1}{2} \text{erfc}\left(\mu V_0 \sqrt{\frac{t}{2\mu}}\right) e^{(\mu V_0)^2 \frac{t}{2\mu}} - \frac{1}{2}$$



(a)



(b)

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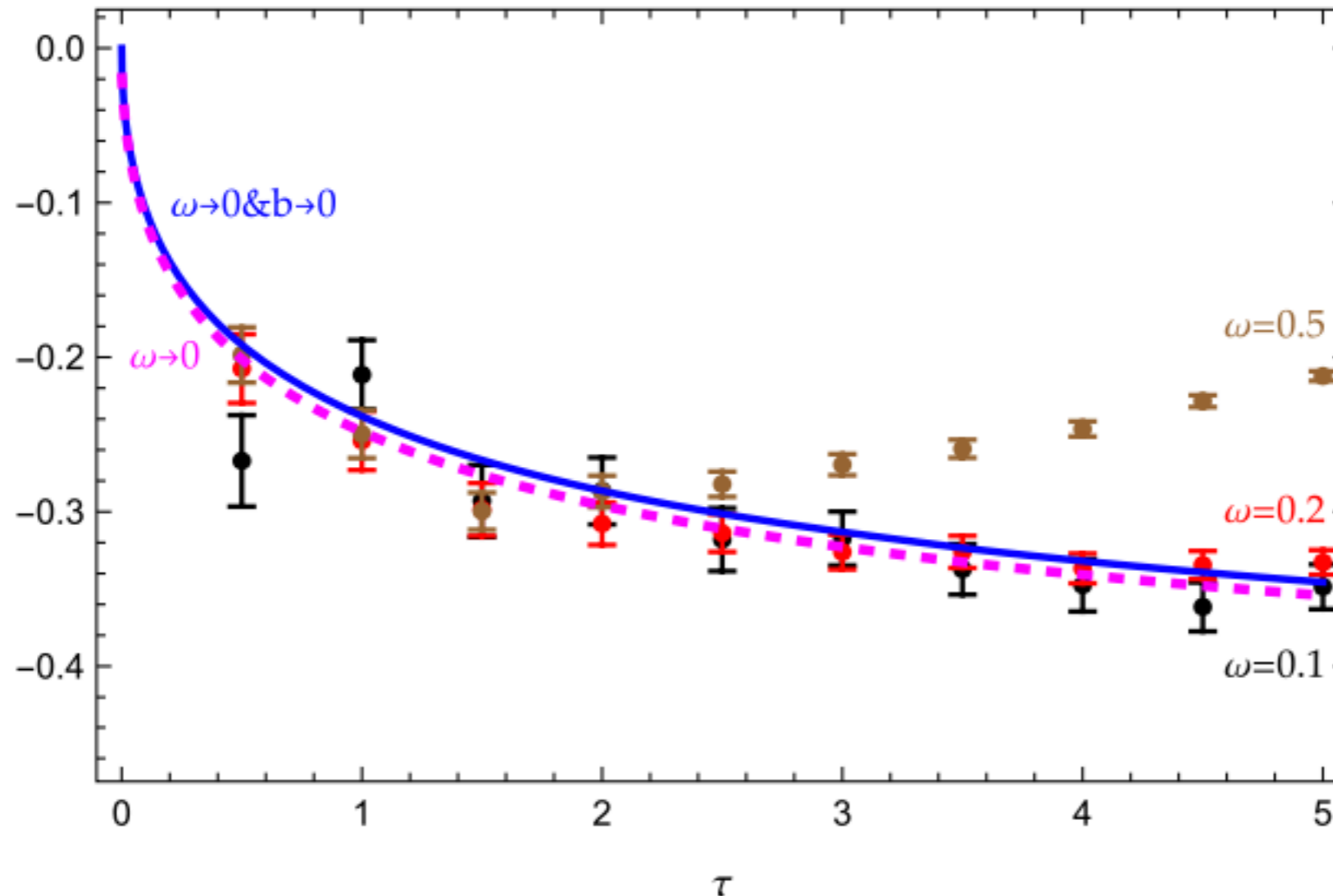
$$\left[-\frac{1}{2\mu} \frac{d^2}{dr^2} + \frac{1}{2} \mu \omega^2 r^2 + V(r) \right] \psi_n(r) = \epsilon_n \psi_n(r),$$

- ✱ Monte Carlo simulation test with square well potential in harmonic trap:

where

$$V(r) = \begin{cases} \frac{V_0}{b}, & r \in [-\frac{b}{2}, \frac{b}{2}] \\ 0, & \text{otherwise} \end{cases}, \quad \xrightarrow{b \rightarrow 0} V_0 \delta(r).$$

$$c(\tau) - c_0(\tau) \text{ vs. } \frac{1}{\pi} \int d\epsilon \frac{d(\delta_+ + \delta_-)}{d\epsilon} e^{-\epsilon\tau} - \frac{1}{2}$$



PHYSICAL REVIEW D **108**, 074504 (2023)

• e-Print: [2402.15628](https://arxiv.org/abs/2402.15628) [hep-lat]

Alternative to Lusecher formula method:

- ✱ Relativistic extension for complex scalar lattice field theory model

$$S_E = -\kappa \sum_{x,t,\hat{n}_x,\hat{n}_t} \hat{\phi}^*(x,t) \hat{\phi}(x + \hat{n}_x, t + \hat{n}_t) + c.c. \\ + (1 - 2\lambda) \sum_{x,t} |\hat{\phi}(x,t)|^2 + \lambda \sum_{x,t} |\hat{\phi}(x,t)|^4$$

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$$\Delta C(t) \stackrel{t=-i\tau}{=} \sum_{n=0}^{\infty} \left[\frac{e^{-E_n \tau}}{E_n} - \frac{e^{-E_n^{(0)} \tau}}{E_n^{(0)}} \right]$$



$$\xrightarrow[t=-i\tau]{L \rightarrow \infty} \frac{1}{\pi} \int_{2m}^{\infty} d\epsilon \delta(\epsilon) \left(\tau + \frac{1}{\epsilon} \right) \frac{e^{-\epsilon \tau}}{\epsilon}$$

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$$C(rt; r'0) = \langle 0 | \mathcal{T} \left[\hat{\mathcal{O}}(r, t) \hat{\mathcal{O}}^\dagger(r', 0) \right] | 0 \rangle.$$

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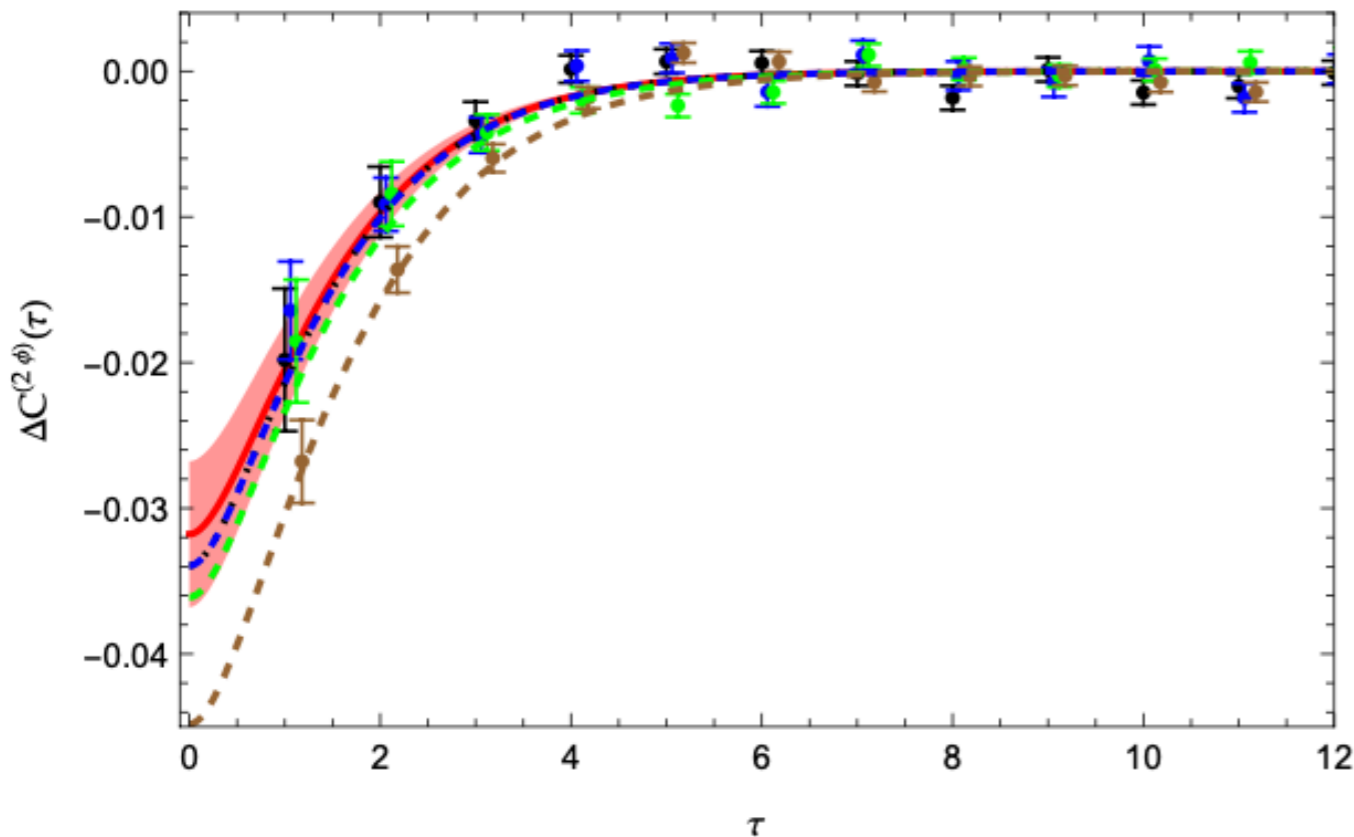
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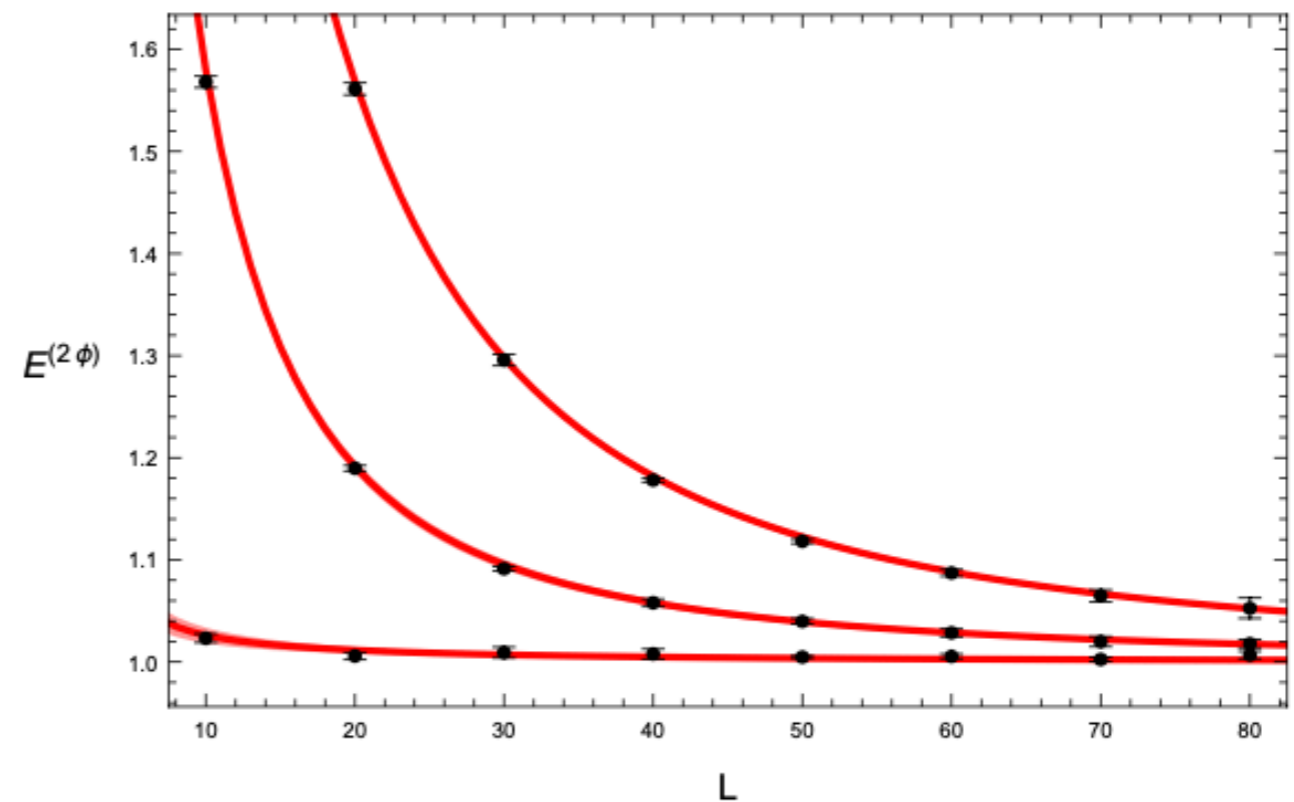
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$$+ (1 - 2\lambda) \sum_{x,t} |\hat{\phi}(x,t)|^2 + \lambda \sum_{x,t} |\hat{\phi}(x,t)|^4$$



(a) Heavy ϕ field with $m \sim 0.500$ vs. lattice data.
 PHYSICAL REVIEW D **110**, 014504 (2024)



(b) $E^{(2\phi)}$ with $m \sim 0.500$ and $V_0 \sim 0.271$
 • e-Print: [2402.15628](https://arxiv.org/abs/2402.15628) [hep-lat]

Summary

Alternative approach to Lusecher formula method

- ✓ **The difference of integrated correlation functions between interacting and non-interacting systems converge rapidly to its infinite volume limit that is related to scattering phase shift through an integral;**



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Summary

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- ☑ **Good candidate to overcome Lusecher formula method at large volume limit;**
- ☑ **Inelastic effect may be important in some cases, and need to be build in;**
- ☑ **May have potential to overcome S/N problem.**



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Outlook

Integrated correlation function approach to scattering phase shifts

☑ Including coupled-channel effect:

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Toward extracting scattering phase shift from integrated correlation functions. III. Coupled channels



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Outlook

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☑ Including coupled-channel effect:

PHYSICAL REVIEW D **111**, 054506 (2025)

Toward extracting scattering phase shift from integrated correlation functions. III. Coupled channels

☑ 3+1 relativistic lattice model test (on-going) and real QCD test;



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Outlook

Integrated correlation function approach to scattering phase shifts

☑ Including coupled-channel effect:

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Toward extracting scattering phase shift from integrated correlation functions. III. Coupled channels

☑ 3+1 relativistic lattice model test (on-going) and real QCD test;

☑ Quantum simulation (coming soon).



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