Energy-energy correlators in jets across collision systems

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Jet substructure observables describe the dynamics of high-energy partons and their confinement into hadrons.

Collision

Hadronization hadrons $\pi^{\pm} K^{\pm} \dots$



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Fragmentation partons **QUd** ...

Detection

Energy correlators are a novel jet substructure observable.

Energy-energy correlators measure statistical correlations of energy flow within jets.

- correlations of energy flow operators
- IRC-safe and well-defined in QFT!
- no need for declustering or grooming

$$\frac{d\sigma_{EEC}}{dR_L} = \sum_{i,j} \int d\sigma(R'_L) \frac{p_{T,i} p_{T,j}}{p_{T,jet}^2} \delta(R'_L - R_{L,ij})$$

EECs are fundamental theoretical objects *and* experimentally measurable!



The two-point EEC is calculated from pairs of tracks.

 $R_L = \sqrt{\Delta \phi_{ij}^2 + \Delta \eta_{ij}^2}$

EEC definition:

$$\frac{d\sigma_{EEC}}{dR_L} = \sum_{i,j} \int d\sigma(R'_L) \frac{p_{T,i} p_{T,j}}{p_{T,jet}^2} \delta(R'_L - R_{L,ij})$$

- For each pair of tracks inside the jet, calculate the energy weight.
- 2. Count the number of weighted track pairs as a function of R_{L} .



EECs show a clear transition region.

- Transition between perturbative (large R_L, partonic) and non-perturbative regimes (hadronic, small R₁)
- Time evolution of jet formation is imprinted onto the EEC angular scaling
- EECs let us probe jet formation and confinement!



In this talk: differential measurements of EECs

- 1. EECs in pp probing hadronization
- 2. EECs in D⁰-tagged jets probing flavor effects
- 3. EECs in p-Pb probing jets in higher-multiplicity environments
- 4. EECs in Pb-Pb

probing jets in the presence of the QGP

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EECs in pp

Clear separation of perturbative and nonperturbative regions.

EEC peak is visibly dependent on jet p_{T} .

Universal transition region after rescaling the x-axis to $\langle p_T^{ch jet} \rangle R_L$ (common peak position and height).

 $\Sigma_{EEC}(R_L)$







We can probe hadronization with pp EECs...



- PYTHIA & Herwig perform well, Herwig captures peak position
- Sherpa Lund does well, AHADIC does not
- both cluster models peak at smaller R,

Lund string models: PYTHIA 8, Sherpa Lund Cluster models: Herwig 7, Sherpa AHADIC





cluster models

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... or with charged energy correlators!



- Are there more unlike-sign or like-sign correlations? $\langle \mathcal{E}_+ \mathcal{E}_- \rangle$ or $\langle \mathcal{E}_+ \mathcal{E}_+ \rangle + \langle \mathcal{E}_- \mathcal{E}_- \rangle$
- Exploring correlations in angle and charge increases sensitivity to different hadronization mechanisms.

STAR measured charged correlators.

- The charged EECs align in the perturbative region with MC predictions.
- MCs underpredict like-sign EEC at small R₁!
- Cannot resolve the predicted shift in EEC peak

- Herwig (cluster) and PYTHIA (string) predict similar behavior in the charge-weighted ratio
- PYTHIA captures large R_L behavior, but neither model describes small R_L de-correlation!
 - resembles STAR r_c measurement



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probing jets in the presence of the QGP

How does the shower depend on flavor?

- At the largest R₁, the scaling behavior in heavy-flavor jets is identical to light quark jets.
- Turnover exhibits a mass dependence!
- Change of shape at small angles is a consequence of the dead cone.

gluon jets

Flavor effects can be probed with ratios to inclusive jets.

quark iets

heavy guark jets



arXiv: 2210.09311 (Craft, Lee, Mecaj, Moult)

Comparing D⁰-tagged jet to inclusive jet EECs

Upper panel:

- p_{T} cut on leading track in incl. jets to study fragmentation bias
- significant suppression at all R_{L} , slopes at large R_{L} seem different
- peak positions are similar due to gluon contribution to inclusive

From the ratios:

- D^0 /inclusive \rightarrow mass + Casimir
- $D^0/LF \rightarrow$ isolated mass effects

Clear mass effect in D⁰ jets!



pQCD calculation from K. Lee

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Are EECs modified in p-Pb?

Differences from pp:

- initial state (nPDF, isospin)
- final-state interactions?
- comovers? collectivity??



EECs in p-Pb are a window into interactions in small systems.



EECs are modified in p-Pb, in the lowest jet p_{τ} range.



- Significant difference between EECs in p-Pb compared to pp!
 - jet structure appears to be altered only in the lowest jet $p_{\rm T}$ range
- Initial state effect?
 - some models lead to a similar effect
- Final state effect?
 - some calculations reproduce this
- Modification is qualitatively consistent with ALICE measurement^{*} of HM/MB z_{ch} in pp

nPDF models do not fully capture the enhancement and suppression seen in data.

- Comparing to PYTHIA with an nPDF turned on, and PYTHIA Angantyr
- PYTHIA results use:
 - nPDF: EPPS21nlo_CT18Anlo
 - PDF: CT14nlo
- nPDFs are within ~1 σ at small R_{L} but these are very large uncertainties
- Neither captures the behavior at large $R_{\rm L}$.



Some theory models can reproduce the enhancement.

- Fu et al. use a higher-twist model and require comovers to capture the enhancement
 - p_{T}^{jet} of 30 GeV/*c*, R_{pPb} chosen to be 0.85
- Barata et al. capture the NP region as well with multiple scatterings and transverse momentum broadening
 - HT correction to jet's perturbative evolution
 - p_{T} broadening in TMD fragmentation function



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Jets lose energy as they traverse the quark-gluon plasma.



EECs in Pb-Pb can probe various medium effects.

Color coherence



- *large angle emission*: medium resolves emitted gluon as a separate object
- *small angle emission*: gluon and emitter resolved as single object
- *critical angle*: minimum separation to resolve separately



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Color coherence

- *large angle emission*: medium resolves emitted gluon as a separate object
- *small angle emission*: gluon and emitter resolved as single object
- critical angle: minimum separation to resolve separately
 Medium response
 - Jets can induce a *medium response* (recoil partons and back-reaction).
 - Energetic partons can pull the medium, leaving a depletion called the *"jet wake"*.





άΣ/dθ



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CMS measured EECs in Pb-Pb...



Pb-Pb loosely has the same features as observed in pp.

CMS measured EECs in Pb-Pb... and found differences.



Pb-Pb loosely has the same features as observed in pp.



- Pb-Pb peaks at smaller R₁ than pp
 - energy loss: Pb-Pb jets have higher initial virtuality
 - flat trend at lowest R_{L} (free hadron scaling)
- Pb-Pb EEC has a different shape at large R₁
 - at $R_{\rm L}$ ~ 0.1, we see another enhancement

How do calculations stack up against this data?

Color coherence

- Holguin et al. calculation* has two free parameters: k and normalization
- Calculation is normalized to data in $0.042 < R_1 < 0.126$
- Turn-on angle is similar in data and calculation! _

* arXiv:2407.07936 [hep-ph] (Andres, Dominguez, Holguin, Marguet, Moult)

Medium response

- Modeling medium response using jet wake in Hybrid
- Hybrid only predicts the large- R_1 enhancement with the wake included
 - does not capture the magnitude or onset of the effect
- JEWEL (not shown) requires recoils for enhancement



Summary and outlook

- Universality of EEC shape and turnover in pp no very obvious conclusions about hadronization
- EECs are altered in HF jets dramatic reduction in amplitude – clear mass/flavor effect
- EECs are modified in p-Pb models with final-state interactions capture large R₁ trends
- Energy loss is visible in Pb-Pb EECs along with some interesting large R₁ effects (medium response?)

Not shown here: N-point energy correlators, recent theory developments, and more!

Many more results are coming at QM25 – it's an exciting time in the world of EECs.

Backup

TRACK FUNCTIONS

Chang, Procura, Thaler, Waalewijn `13 Li, Moult, Waalewijn, Zhu et al `21, 22

Track functions are non-perturbative functions describing the momentum fraction of initial parton converted to hadrons with a particular property R.

$$T_q(x) = \int dy^+ d^2 y_\perp e^{ik^- y^+/2} \frac{1}{2N_c} \sum_X \delta\left(x - \frac{P_R}{k_-}\right) \operatorname{tr}\left[\frac{\gamma^-}{2} \left\langle 0 \left|\psi\left(y^+, 0, y_\perp\right)\right| X \right\rangle \left\langle X |\bar{\psi}(0)|0 \right\rangle\right]$$

Track functions are technically what is needed for calculating the usual jet shape observables on track (jet mass, jet angularities, jet charge, etc...), but is technically challenging

$$\delta(e - \hat{e}(\{p_i^{\mu} \in X_J\})) \to \delta(e - \hat{e}(\{x_i p_i^{\mu} \in X_J\}))$$

requires simultaneous knowledge of all the tracks in jet.

NLO ···· LO NLO ---- LO NLO ---- LC 2.0 T(x)1.5 1.0 0.5 $\mu = 100 \text{ GeV}$ 0.0 0.2 0.4 0.6 0.8 0.0 1.0



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angle \langle X | ar{\psi}(0) | 0
ight
angle
ight]$$



bservables. For all of these observables, the uncertainties for the track-based observables are significantly maller than those for the calorimeter-based observables, particularly for higher values of β , where more oft radiation is included within the jet. However, since no track-based calculations exist at the present me, calorimeter-based measurements are still useful for precision QCD studies. [ATLAS Collaboration, 1912.09837] selection of charged particle jets. Note that track-based observables are IRC-unsafe. In general, perturbative track functions can be used to directly compare track-based measurements to analytical ulations [67–69]; however, such an approach has not yet been developed for jet angularities. Two





ENERGY CORRELATORS ON TRACK

From this "detector" as a fundamental operator perspective, track function formalism provides the essential matching between partonic and hadronic detectors.

 $\mathcal{E}_{R}(\bar{n})$ \mathcal{E}

 $\langle \mathcal{E}_{R}\left(\vec{n}_{1}\right) \mathcal{E}_{R}\left(\vec{n}_{2}\right) \cdots \mathcal{E}_{R}\left(\vec{n}_{k}\right) \rangle \\ = \sum_{i_{1},i_{2},\cdots i_{k}} T_{i_{1}}(1) \cdots T_{i_{k}}(1) \left\langle \mathcal{E}_{i_{1}}\left(\vec{n}_{1}\right) \mathcal{E}_{i_{2}}\left(\vec{n}_{2}\right) \cdots \mathcal{E}_{i_{k}}\left(\vec{n}_{k}\right) \right\rangle$

+ contact terms. Requires up to k-th moment

Aside from the fact that this is technically much simpler, it only involves NP *numbers*, not functions.



D⁰ reconstruction steps

D⁰ candidates were D⁰-tagged charged jets K. were created using the reconstructed from daughter anti- k_{τ} algorithm (R=0.4) tracks using topological and for each candidate. particle identification selections $(D^0 \rightarrow K^- + \pi^+, and charge conjugates).$ Corrected the EECs for the charged jets, anti-k_T, R = 0.4 Invariant mass analysis was performed to with $D^0 \rightarrow K^-\pi^+$ and charge coni efficiency of D⁰-tagged jet remove combinatorial $K^{-}\pi^{+}$ pairs surviving reconstruction and the D⁰ selections. removed the contribution Promot D⁰ from beauty decays. deband (SB) and charge conj. signal region signal + background iets anti-k-, R = 0.4 - background sideband (SB) 30 p_____(GeV/c) nd charge coni $10 \le p_{\pi}^{\text{ch. jet}} \le 20 \text{ GeV}/c, |\eta_{-1}| \le 0.5$ iets, anti- k_{π} , R = 0.4 $B \le \rho_{1}^{D^{0}} < 12 \text{ GeV}/c, |v| \le 0.1$ < 20 GeV/c In 1<0.5 12 GeV/c, ly _l ≤ 0.8 Corrected the EECs for detector effects with a bin-by-bin correction method. ΔR_{STD-D^0}

Probing hadronization with pp EECs



Lund string models: PYTHIA 8, Sherpa Lund Cluster models: Herwig 7, Sherpa AHADIC

Inclusive jets:

- PYTHIA & Herwig perform well, Herwig captures peak position
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Heavy flavor jets:

- data favors PYTHIA's implementation
- Herwig overpredicts inclusive jet EECs; underpredicts in HF
- Sherpa Lund underpredicts the data
- AHADIC fails to describe the peak

Cluster hadronization for higher p_T jets, string-breaking models for D⁰ jets?

EEC measurements in more complex systems require background subtraction techniques.

1. Subtract UE energy density from the jet p_{T} :

$$\rho = \text{median} \left\{ \frac{p_{\text{T,jet}}^{k_{\text{T}}}}{A_{\text{jet}}^{k_{\text{T}}}} \right\} \cdot C \qquad C = \frac{\sum_{j} A_{j}}{A_{\text{acc}}}$$

- 2. Correct the EEC distribution for combinatorial background:
 - signal-signal (what we want!)
 - signal-background
 - background-background

hard scattering

We use the perpendicular cone to estimate the latter two contributions.

These subtraction steps are also performed for the pp baseline.

Perp. cone for combinatorial EEC background

- Particles from jet: sig + bkg
- Particles from perp cone: bkg'
- Pairs in the combined cone:

(sig + bkg + bkg')(sig + bkg + bkg') = sig*sig + 2sig*bkg + bkg*bkg + 2sig*bkg' + 2bkg*bkg' + bkg'*bkg' jet-perp perp-perp



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- Particles from jet: sig + bkg
- Particles from perp cone: bkg'
- Pairs in the combined cone:

(sig + bkg + bkg')(sig + bkg + bkg') = sig*sig + 2sig*bkg + bkg*bkg + 2sig*bkg' + 2bkg*bkg' + bkg'*bkg' jet-perp perp-perp

- Sig-bkg pairs: jet-perp 2 perp-perp
- Bkg-bkg pairs: perp-perp
- Total background: jet-perp perp-perp



p-Pb and pp comparison



The transition region in p-Pb resembles pp.

- Universality of the EEC peak position across jet p_{τ} and collision system.
- EEC peak height for 20-40 GeV/c jets is slightly lower than for other jets.



2

X

ALICE Preliminary

p-Pb $\sqrt{s_{NN}}$ = 5.02 TeV

anti- k_{T} ch jets, R = 0.4

all pairs, $p_{\tau}^{trk} > 1.0 \text{ GeV}/c$

Transition region

peak $\approx 2.43 \text{ GeV}/c$

- (20, 40) GeV/c

← (60, 80) GeV/*c*

38

± 0.07 GeV/c

range

Varying the track cut changes the EEC behavior at large R_{L} .



Strong sensitivity to track p_{T} cut in low p_{T} jets!

non-perturbative effects increase for lower jet p_{T}

Track cut modifies the enhancement in ratio

 but not the small-R_L suppression

Suggests that the origin of the effect lies in softer interactions at small x!

$log(R_1)^2$ shape of the p-Pb/pp ratio — why?

- The ratio appears to follow a $log(R_{L})^{2}$ scaling. Why: change in q/g ratio?



Quark-jet and gluon-jet EECs



Quark-jet and gluon-jet and D⁰-tagged jet EECs

