Tensor-Polarized DVCS: Prospects & Capabilities

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Univ. of New Hampshire

2024-07-17





What Deuterons Do That Protons Don't

Proton

 $\mathsf{Spin}\text{-}\frac{1}{2} \; \mathsf{System}$





$$m = +\frac{1}{2}$$

$$m=-rac{1}{2}$$

"Typical" Vector Polarization



$$P_z = p_+ - p_-$$

Deuteron

Spin-1 System

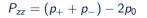






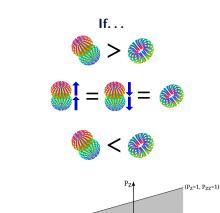
Vector and Tensor Polarization





J Forest, et al, PRC **54** 646 (1996)

Tensor Polarization Properties



 $(P_Z=0, P_{ZZ}=0)$

Then...

$$0 < P_{zz} \le 1$$

$$P_{zz}=0$$

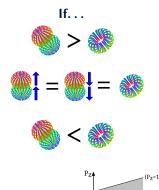
$$-2 \le P_{zz} < 0$$

- P_z ranges from -1 to +1
- \bullet P_{zz} ranges from -2 to +1
- In deuterons both P_z and P_{zz} can be nonzero simultaneously

(Pz=0, Pzz=-2)

(Pz=-1, Pzz=1)

Tensor Polarization Properties

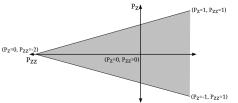


Then...

$$0 < P_{zz} \le 1$$

$$P_{zz}=0$$

$$-2 \le P_{zz} < 0$$



A high-luminosity tensor-polarized target has promise as a novel probe of nuclear physics

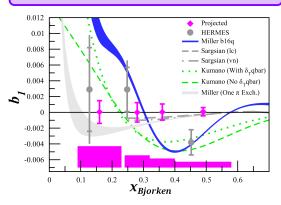
Approved Experiment: b_1

- Intended to improve upon HERMES' 2005 data
- Verifications of zero-crossing
 - Implications for Close-Kumano sum rule
- Tensor physics at quark level
- Better understanding of b₁ allows discrimination of different deuteron components by spin (e.g., quarks vs gluons)

Approved by JLab with A-physics rating!

E12-13-011

The Deuteron Tensor Structure Function b_1



K. Slifer et al, JLab C12-13-011 Spokespersons: K. Slifer, O.R. Aramayo, J.P. Chen, N. Kalantrians, D. Keller, E. Long, P. Solvignon

Approved Experiment: A_{zz}

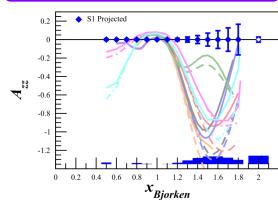
- First-of-its-kind quasielastic
 A_{zz} measurement
- Implications for SRC physics and deuteron wavefunction
- Widest range of x covered by a single measurement
- Measurement of T₂₀ included!

Spokespersons: E. Long, K. Slifer, P. Solvignon, D. Day, D. Keller, D. Higinbotham

Approved by JLab with Apphysics rating!

E12-15-005

Quasi-Elastic and Elastic Deuteron Tensor Asymmetries



E. Long et al, JLab C12-15-005

Propsects: Deuterons, GPDs and DVCS

Spin-1 GPDs

Vector and Axial Elements for Deuteron GPDs

$$\begin{split} V_{\lambda'\lambda} &= -(\epsilon'^* \cdot \epsilon) H_1 + \frac{(\epsilon \cdot n)(\epsilon'^* \cdot P) + (\epsilon'^* \cdot n)(\epsilon \cdot P)}{P \cdot n} H_2 - \frac{(\epsilon \cdot P)(\epsilon'^* \cdot P)}{2M_D^2} H_3 \\ &+ \frac{(\epsilon \cdot n)(\epsilon'^* \cdot P) - (\epsilon'^* \cdot n)(\epsilon \cdot P)}{P \cdot n} H_4 + \left[4M_D^2 \frac{(\epsilon \cdot n)(\epsilon'^* \cdot n)}{(P \cdot n)^2 + \frac{1}{3}(\epsilon'^* \cdot \epsilon)} \right] H_5 \end{split} \tag{1}$$

$$\begin{split} A_{\lambda'\lambda} &= -i \frac{\epsilon_{\mu\alpha\beta\gamma} \mathsf{n}^{\mu} \epsilon'^{**\alpha} \epsilon^{\beta} P^{\gamma}}{P \cdot \mathsf{n}} \tilde{H}_{1} \\ &+ i \frac{\epsilon_{\mu\alpha\beta\gamma} \mathsf{n}^{\mu} \Delta^{\alpha} \epsilon^{\beta}}{P \cdot \mathsf{n}} \frac{\epsilon^{\gamma} (\epsilon'^{*} \cdot P) + \epsilon'^{**\gamma} (\epsilon \cdot P)}{M_{D}^{2}} \tilde{H}_{2} \\ &+ i \frac{\epsilon_{\mu\alpha\beta\gamma} \mathsf{n}^{\mu} \Delta^{\alpha} P^{\beta}}{P \cdot \mathsf{n}} \frac{\epsilon^{\gamma} (\epsilon'^{*} \cdot P) - \epsilon'^{**\gamma} (\epsilon \cdot P)}{M_{D}^{2}} \tilde{H}_{3} \\ &+ i \frac{\epsilon_{\mu\alpha\beta\gamma} \mathsf{n}^{\mu} \Delta^{\alpha} P^{\beta}}{P \cdot \mathsf{n}} \frac{\epsilon^{\gamma} (\epsilon'^{*} \cdot \mathsf{n}) + \epsilon'^{**\gamma} (\epsilon \cdot \mathsf{n})}{P \cdot \mathsf{n}} \tilde{H}_{4} \end{split}$$

where P=(p+p')/2, ϵ is the target deuteron polarization and ϵ' is the recoil deuteron polarization

E. Berger, F. Cano, M. Diehl, B. Pire PRL 87 (2001) 142302

Spin-1 GPDs

Relevant Sum Rules

$$\int_{-1}^{1} dx H_i(x,\xi,t) = G_i(t) \qquad (i=1,2,3) \qquad (3)$$

$$\int_{-1}^{1} dx \tilde{H}_i(x,\xi,t) = \tilde{G}_i(t) \qquad (i=1,2) \qquad (4)$$

$$\int_{-1}^{1} dx H_4(x,\xi,t) = \int_{-1}^{1} \tilde{H}_3(x,\xi,t) = 0 \qquad (5)$$

$$\int_{-1}^{1} dx H_5(x,\xi,t) = \int_{-1}^{1} \tilde{H}_4(x,\xi,t) = 0 \qquad (6) \qquad \begin{aligned} H_4(x,\xi,t) &= -H_4(x,-\xi,t) \quad (9) \\ \tilde{H}_3(x,\xi,t) &= -\tilde{H}_3(x,\xi,t) \quad (10) \end{aligned}$$

Time-Reversal Behavior

$$H_i(x,\xi,t) = H_i(x,-\xi,t)$$

(i = 1,2,3,5) (7)

$$\tilde{H}(x,\xi,t) = \tilde{H}_i(x,-\xi,t)
(i = 1,2,4)$$
(8)

$$H_4(x,\xi,t) = -H_4(x,-\xi,t)$$
 (9)

$$\tilde{H}_3(x,\xi,t) = -\tilde{H}_3(x,\xi,t)$$
 (10)

E. Berger, F. Cano, M. Diehl, B. Pire PRL 87 (2001) 142302 M. Diehl Phys. Rept. 388 (2003) 41-277

Form Factors & the Forward Limit

Form Factors

Forward Limit

$$\int_{-1}^{1} dx H_{1}(x,\xi,t) = G_{C}(t) + \frac{tG_{Q}(t)}{6M_{D}^{2}} \qquad \Rightarrow \int_{-1}^{1} dx H_{1}(x,0,0) = Z \qquad (11)$$

$$\int_{-1}^{1} dx H_{2}(x,\xi,t) = G_{M}(t) \qquad \Rightarrow \int_{-1}^{1} dx H_{2}(x,0,0) = Z\mu_{D} \qquad (12)$$

$$\int_{-1}^{1} dx H_{3}(x,\xi,t) = \frac{1}{1 - \frac{t}{4M_{D}^{2}}} \left[G_{M}(t) - G_{C}(t) + \left(1 - \frac{t}{6M_{d}^{2}} \right) G_{Q}(t) \right]$$

$$\Rightarrow \left[\int_{-1}^{1} dx H_{3}(x,0,0) = Z(\mu_{D} + Q_{D} - 1) \right]$$
(13)

S Liuti, K Kathuria 2014 J. Phys. Conf. Ser. 543 012005; A. Kirchner, D. Müller Eur.Phys.J. C32 (2003) 347-375 $H_1\Rightarrow$ momentum, $H_2\Rightarrow$ ang. mom., $H_3\Rightarrow$ quadrupole. \mathcal{Q}_D is the deuteron quadrupole moment

Forward Limit and b_1

First tensor structure function b_1

$$\int_{0}^{1} dx b_{1} = 2 \sum_{i=u,d,s} e_{i}^{2} \int_{0}^{1} dx \delta \bar{q}_{i}(x) \quad (14)$$

for unpolarized sea quarks:

$$\int_0^1 dx b_1(x) = 0 \qquad (15)$$

F. E. Close, S. Kumano PRD 42 (1990) 2377

A. Kirchner, D. Müller Eur.Phys.J. C32 (2003) 347-375

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Tensor GPD H₅

In forward limit:

$$H_5(x,0,0) \equiv b_1(x)$$
 (16)

with

$$b_1(x) = \frac{1}{2} \sum_{q} e_q^2 \left[q^0(x) - \frac{q^1(x) + q^{-1}(x)}{2} \right]$$
(17)

where $q^{i}(x)$ are quark densities

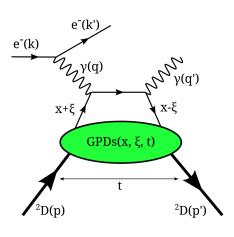
S. Liuti, K. Kathuria 2014 J. Phys. Conf. Ser. 543 012005

Coherent Deuteron DVCS

$$eD
ightarrow e \gamma D$$

- Cross-section dependent on \mathcal{M}_{BH} and \mathcal{M}_{DVCS} and interference terms
- Rates not yet examined
- Unpolarized deuteron DVCS: E08-25
- Two approved inclusive tensor experiments: b₁ and A_{zz}
- New experimental group for exclusive measurement of Azz
- No current proposals for DVCS with a tensor-polarized target!

M. Benali *et. al.* Nat. Phys. 16 (2020) 191-198



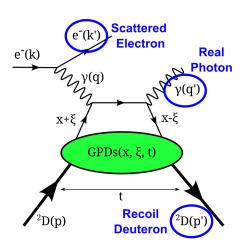
Above: Handbag diagram for coherent deuteron DVCS

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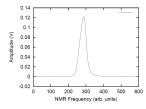


Above: Handbag diagram for coherent deuteron DVCS, with outgoing particles labeled.

UNH Polarized Targets

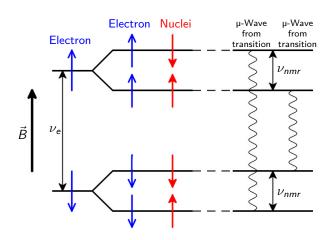
Dynamic Nuclear Polarization (DNP)

- Using µwaves, drive spin transitions of unpaired electrons
- Electrons transfer spin to nuclei
- Nuclear absorption spectrum gives polarimetry info



Above: Characteristic lineshape of the proton

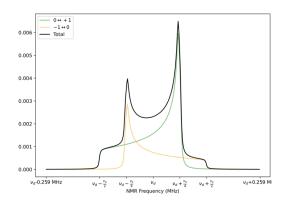
C.D. Keith et al, NIM A 501



Above: Diagram of the energy level transitions in the DNP process. Adapted from Annu. Rev. Nucl. Part. Sci. 1997. 47:67-109

Deuteron Polarization

- NMR at nuclear spin transition frequency drives further spin transitions
- Proton lineshape from $-1/2 \leftrightarrow 1/2$ transition
- $\begin{array}{c} \bullet \ \ \ \ \, \text{Deuteron lineshape has} \\ -1 \leftrightarrow 0 \ \ \text{and} \ \ 0 \leftrightarrow 1 \\ \ \ \ \ \text{components} \end{array}$
 - But NMR only gives the sum of the two
- Signal shape affected by material properties and magnetic field angle
- P_{zz} extracted from area difference of transition curves!



Above: Simulated deuteron lineshape showing the contributions from both the $-1 \to 0$ transition and the $0 \to 1$ transition. Lineshape generated by E. Long.

NMR Curve Fitting

- Fit NMR lineshape with procedure from C. Dulya et al, NIM A 398 (1997) 109-125
- Includes effects from molecular bond quadrupole terms
- Can naively use peak height ratio r to estimate polarization

$$P_z = \frac{r^2 - 1}{r + r^2 + 1}$$

$$P_{zz} = \frac{r^2 - 2r + 1}{r^2 + r + 1}$$
(18)

 Then compare ratio and area methods for P_{zz} measurement consistency

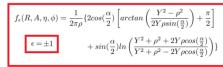
Right: Parts of the curve fitting method suggested by C. Dulya et al.

R, A, η, ϕ compacting variables

$$\rho^2 = \sqrt{A^2 + [1 - \epsilon R - \eta cos(2\phi)]^2} \qquad \qquad R = \frac{\omega - \omega_d}{3\omega_q}$$

$$cos(\alpha) = \frac{1 - \epsilon R - \eta cos(2\phi)}{\rho^2} \qquad \qquad -3 \le R \le 3$$

functional form of signal



phi average 🞝

$$F_{\epsilon} \approx \frac{1}{J+1} \sum_{j=0}^{J} \frac{\sqrt{3} f_{\epsilon}(R,A,\eta,\phi_{j})}{\sqrt{3-\eta cos(2\phi_{j})}}$$

positive & negative spin flips \bigcup



UNH Polarized Target Lab



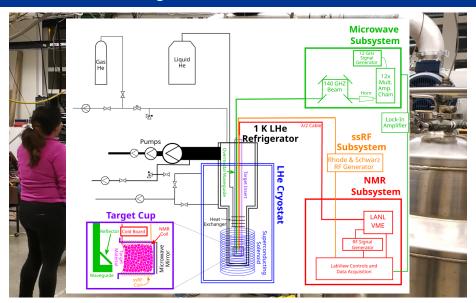
The UNH polarized target group is hard at work!

UNH Polarized Target Lab



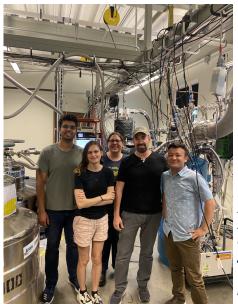
The UNH polarized target group is hard at work!

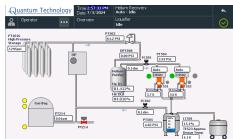
UNH Polarized Target Lab



The UNH polarized target group is hard at work!

NEW Helium Reliquefaction System





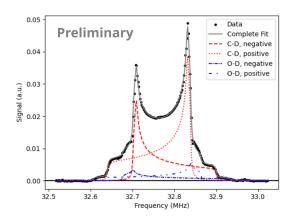
Left: group photo of the reliquefier installation team. Above: Reliquifier system software overview. Right: A grad student adds LN2 to our helium purifier dewar.



Tensor Target & Development

Tensor Polarization (UNH)

- Fit with Dulya procedure closely matches data from recent UNH cooldown
 - C. Dulya et al, NIM A 398 (1997) 109-125
- Reconstruct spin-flip and quadrupole curves from fit parameters
- With reconstruction can do more in-depth polarization analysis
- Fit method works very well for UNH data!



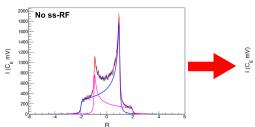
Above: Curve fit of NMR lineshape from recent target cooldown at UNH.

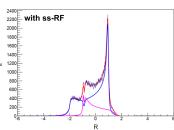
Tensor Polarization Measurement (UVA)

- UVA-pioneered tensor enhancement technique
- Additional RF coils drive spin flips
- Manipulates area of NMR curve

- Small frequency range
- UVA lab achieved $P_{zz}{=}31.1{\pm}8.5\%$ with ssRF technique

Right: ssRF coil schematic. Below: NMR lineshapes without and with ssRF applied. Figures reproduced from D. Keller, et al. NIM A **981** 164504 (2020)



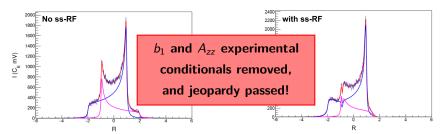


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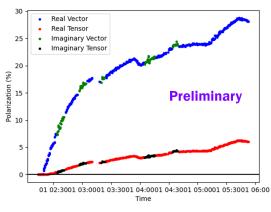
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Real & Imaginary NMR Signals

- Switch from real to imaginary lineshape by tuning phase
- Use fitting for real and imaginary lineshapes differently
- Demonstrated resilience to having phase not tuned perfectly
- Real and imaginary measurements match each other well!



Above: Data from recent UNH cooldown with both real and imaginary line data for both vector and tensor polarization. Figure courtesy of M. McClellan.

Summary

Professors

















Nathaly

Santiestehan

David Ruth

70e Wolters

Thank you to the UNH **PolTarg** Group and our collaborators at UVA!

- Inclusive tensor program approved!
- Tensor-polarized DVCS has implications for H_3 and H_5 GPDs
- Possible DVCS expansion to tensor experiment program, actively seeking collaborators!
- Target development for current tensor experiments
- UNH group polarizing more (publications upcoming. . .)
- Exciting new developments upcoming!

Trans & Nonbinary Physicists



The Trans and Nonbinary Physicists
Discord server is an online
community for transgender and
nonbinary physicists — from
enthusiasts to professors! — to
socialize, network, and support one
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Follow

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Questions, comments, concerns, observations?

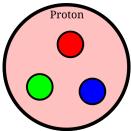
Backup Slides



Protons & Deuterons

Proton





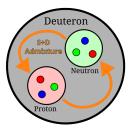
Three valence quarks + gluons and sea quarks

No nucleon-nucleon interactions

$$m=\pmrac{1}{2}$$
 S. Kumano, IOP Proc. Tens. Pol. Targ. (2014)

Deuteron

Spin-1 System



Proton-Neutron bound state

Simplest nuclear system: nucleon interaction effects

$$m=\pm 1,0$$

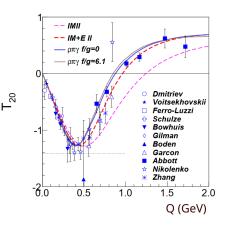
Elastic Tensor Analyzing Power

For
$$1.5 \le x \le 2.0$$

$$T_{20} \approx \frac{A_{zz}}{\sqrt{2}d_{20}} \tag{19}$$

- Third of three elastic scattering functions of deuteron
- Extracted by measuring A_{zz} near elastic peak
- Current data doesn't constrain models well at high x

M. Kohl Nucl Phys A 805 (2008)



Above: T_{20} with current measurements and theoretical models.

R. Holt, R. Gilman Rept.Prog.Phys. 75 (2012)

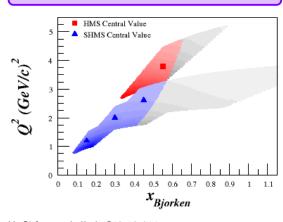
b₁ Kinematics

- Approved for 30 days of physics running + 10.8 days overhead
- 11 GeV beam incident on polarized target
- 9.2% systematic error on A_{zz}
- Foward scattering angles

	ХВj	Q^2 [GeV ²]	E'_0 [GeV]	$\begin{bmatrix} \theta_{e'} \\ [^{\circ}] \end{bmatrix}$
SHMS	0.15	1.21	6.70	7.35
SHMS	0.30	2.00	7.45	8.96
SHMS	0.452	2.58	7.96	9.85
HMS	0.55	3.81	7.31	12.50

E12-13-011

The Deuteron Tensor Structure Function b_1



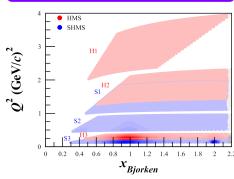
K. Slifer et al, JLab C12-13-011

A_{zz} Kinematics

- Approved for 34 days of physics running + 10.3 days overhead
 - 25 days 8.8 GeV beam
 - 8 days 6.6 GeV beam
 - 1 day of 2.2 GeV beam
- 9.2% systematic error on A_{zz} , 7.4% on T_{20}
- Forward scattering angles

	E_0 [GeV]	Q^2 [GeV ²]	E_0' [GeV]	$\begin{bmatrix} \theta_{e'} \\ [^{\circ}] \end{bmatrix}$
SHMS (S1)	8.8	1.5	8.36	8.2
HMS (H1)	8.8	2.9	7.26	12.2
SHMS (S2)	6.6	0.7	6.35	7.5
HMS (H2)	6.6	1.8	5.96	12.3
SHMS (S3)	2.2	0.2	2.15	10.9
HMS (H3)	2.2	0.3	2.11	14.9

E12-15-005 Quasi-Elastic and Elastic Deuteron Tensor Asymmetries

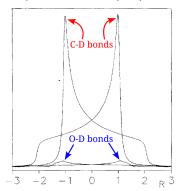


E. Long et al, JLab C12-15-005

BACKUP: Tensor Polarization & DNP

ND₃ and Other Target Materials

C. Dulya, et al, NIM A 398 (1997)



- Both b_1 and A_{zz} experiments call for solid ND₃ targets
- Polarization also done with frozen chemically-doped deuterated alcohols
- Lineshape affected by quadrupole splitting of molecule
 - Different for ND₃ vs butanol

Left: C-D, O-D bond contribution to the deuteron NMR lineshape in d-butanol

Material	Dopant & method	Polarizable nucleons % by weight	
ND ₃	\overline{ND}_2	~30%	
d-ammonia	Irradiation	7~3076	
C ₄ D ₉ OD	TEMPO	23.7%	
d-butanol	Chemical		

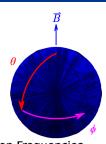
D. Crabb, W. Meyer, Annu. Rev. Nucl. Part. Sci 47 67-109 (1997)

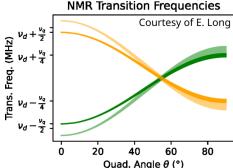
BACKUP: UNH NMR Simulation

From Spin Flips To Lineshape

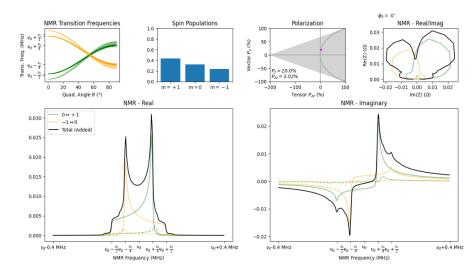
- Simulation derived from Cohen & Reif model (1957)
- Assume random distribution of quadrupole angles θ from 0°-90°
- Then calculates transition frequency based on angle and energy level
- Performs 125,000 spin flips into 300 angle bins, giving NMR signal
- Simulation can reproduce results from UVA polarized target lab, as well as previous UNH cooldowns

Right: Quadrupole angles randomly distributed in B-field. Below: Spin transition frequencies for each angle θ and each transition.





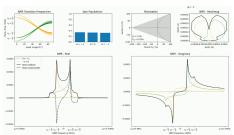
Simulated NMR Spectrum

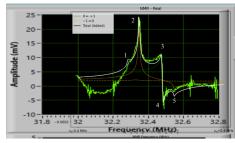


Generates real & imaginary components of NMR signal plus spin transition components! (Figure courtesy M. McClellan)

NMR Lineshape: Real & Imaginary

- Simulation can be used to retroactively understand previous cooldowns
- First UNH deuteron "ugly" NMR signal now understood to be from a mistuned phase





Above: Lineshape of first deuteron NMR signal recorded by UNH group (Fall 2020). *Left:* Simulation showing matching real and imaginary components.

BACKUP: Tensor Polarization Analysis

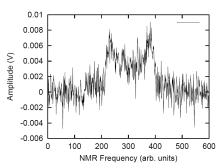
Thermal Equilibrium & Enhancement

Deuteron thermal equilibrium (TE) polarization before microwave irradiation:

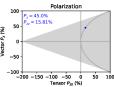
$$P(1) = \frac{4 \tanh\left(\frac{g_i \mu_i B}{2k_B T}\right)}{3 + \tanh^2\left(\frac{g_i \mu_i B}{2k_B T}\right)} \tag{20}$$

Only 0.1% polarization at 5 T and 1 K.

TE signal can be used for calibration if detected. Signal is then enhanced with microwaves.



Above: Deuteron TE signal from CLAS target. From C. Keith et al, NIM A 501 (2003). Right: Polarization curve during enhancement.



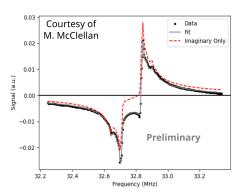
Real & Imaginary Fits

- Can now manually set NMR phase angle ϕ during cooldowns
- Fit using a rotation of the absorptive (χ'') and dispersive (χ') around phase angle:

Real =
$$\chi'' \cos \phi - \chi' \sin \phi$$

Imag = $\chi'' \sin \phi + \chi' \cos \phi$ (21)

- Can fit a simultaneous mixture of real and imaginary
- First fits with the new method match data well, look very promising!



Above: Fit of recent cooldown data using real and imaginary parts. Fit is compared with an "imaginary only" signal and then fitted for a phase mistune.