# Quantum Chromodynamics Predictions: Structure Formation in the Proton Sea





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# Outline for talk on Fundamental QCD predictions for the internal structures of nucleons (& all baryons)

- Finding structure via quark behavior (Deep Inelastic Scattering overview)
- Quantum Chromodynamics definitions and predictions, including short-range QCD potentials
- Fock states in the proton: 3-quark, 5-quark, 7-quark, ...
- "Diquark capture" model on higher Fock states in the nucleon work in progress with Stan Brodsky
- Diquark capture model vs. data



3-quark Fock state of proton



5-quark Fock state of proton

### <u>Finding Structure via Quark Behavior:</u> <u>Deep Inelastic Scattering variables</u>

- Deep inelastic scattering (DIS) experiments
- Lepton scatters from target, exchanging virtual photon with 4-momentum  $q^2$  given by:  $Q^2 \equiv -q^2 = 2EE'(1 \cos \theta)$
- Fraction of nucleon momentum carried by struck quark is found via **Bjorken scaling variable**  $x_B = \frac{Q^2}{2M_p v}$ , where  $\nu = E - E'$ ,  $M_p$ =mass of nucleon (lepton mass neglected)



Adapted from Kerry Whisnant, 2002 book chapter

Differential cross section for DIS:

$$\frac{d\sigma}{dxdy}\left(e^{-}p \to e^{-}X\right) = \sum_{f} x \ e_{f}^{2}\left[q_{f}(x) + \overline{q}_{\overline{f}}(x)\right] \cdot \frac{2\pi\alpha^{2}s}{Q^{4}}\left(1 + (1-y)^{2}\right)$$

where  $y = \frac{\nu}{E}$  is the fraction of lepton energy transferred to target.  $F_2(x)$  is the **nucleon structure function**, defined as:

$$F_2(x_B) \equiv \sum_f x_B \ e_f^2 \left( q_f(x_B) + \overline{q}_{\overline{f}}(x_B) \right)$$

in terms of quark distribution functions  $q_f(x)$ : probability to find a quark with momentum  $x_i \in [x, x + dx]$ .

### Standard Model of Particle Physics is Group theory!

 Mathematical definition of SM is group theory - these are symmetries of the Lagrangian & path integral:

### $SU(3)_{Color} \otimes SU(2)_{Weak} \otimes U(1)_{Y}$

- But ↑ is in the early Universe, prior to electroweak symmetry breaking by the Higgs boson
- Today, the Standard Model of the cold Universe in a blackbody soup of photons peaked at 2.7 K with the following symmetry structure:

$$SU(3)_{color} \otimes U(1)_{EM}$$

$$\mathscr{L}_{\mathsf{QCD}} = \sum_{q} \left( \overline{\psi}_{qi} \gamma^{\mu} \left[ \delta_{ij} \partial_{\mu} + ig \left( t_a G^a_{\mu} \right)_{ij} \right] \psi_{qj} - m_q \overline{\psi}_{qi} \psi_{qi} \right) - \frac{1}{4} G^a_{\mu\nu} G^{\mu\nu}_a$$

QCD confines color charged objects to distance scales of ~ 1 femtometer. We never detect color charge.

Only red+green+blue = colorless, the 3 quark baryons

and, e.g., red+antired=colorless, the quark-antiquark **mesons** are observed (*plus tetraquarks and pentaquarks now!*)

### Diquarks & Hidden-Color: Background QCD

Building blocks: Quantum chromodynamics, Spin-statistics constrained

- Begin with group theory mathematics of the strong interaction:  $SU(3)_C$
- Next, degrees of freedom (particles carrying strong force charge) indices run over the 3 colors:

$$q_a$$
 (triplet,  $3_c$ ),  $q^a$  (antitriplet,  $\overline{3}_c$ ),  $g_c^b$  (octet,  $8_c$ )

• All combinations of D.o.F. predicted, color charged & color singlet. Combine via  $\delta^a_b$ ,  $d_{abc}$ ,  $\epsilon_{abc}$ :

$$(\bar{q}^a q_a)_{1_{\mathrm{C}}} \quad (\epsilon^{abc} \ q_a q_b q_c)_{1_{\mathrm{C}}} \quad (\epsilon_{abc} \ \bar{q}^a \bar{q}^b \bar{q}^c)_{1_{\mathrm{C}}}, \quad (q_a q_b \epsilon^{abc})_{\bar{3}_{C}} \quad \leftarrow (qq)^c$$

 Higher Fock states (baryon with 3 valence quarks is lowest order Fock states), e.g., the 5quark Fock state for baryons - gluon splitting to meson cloud:

$$|N\rangle \subset (\epsilon^{abc}q_a q_b q_c \ \bar{q}^e q_e)_{1_{\rm C}}$$

• 5-quark Fock using color octets, gluon splitting:

$$|N\rangle \subset \left( (\epsilon^{abf} q_a q_b q_c)_{8_{\rm C}} (q^c q_f)_{8_{\rm C}} \right)_{1_{\rm C}} \longrightarrow$$



### Quantum Chromodynamics prediction: Diquarks

- 3 color charges in  $SU(3)_C$ , the local gauge symmetry  $\equiv QCD$
- QCD ⇒ Diquark creation: Quark-quark bond with single gluon exchange & group theory transformation into a fundamentally different object:

$$3_C \otimes 3_C \to \overline{3}_C$$

#### Up quark in the fundamental rep of $SU(3)_C$ :





Like quarks and gluons, diquarks carry color charge. They cannot be seen directly due to color confinement. Only  $1_C$  (red+green+blue or red-antired etc.) observed.

Therefore there is no direct evidence for diquarks. Work in progress for diquark detection experimental proposals (*e.g.*, diquark jets from DIS increase  $\Lambda$  production)

Strong indirect evidence exists (baryon mass splittings, Regge slopes).



Diquark with anti-color charge: green x blue -> anti-red

### Quark-quark potential in QCD: V(r) calculation

•  $SU(3)_C$  invariant QCD Lagrangian:

$$\mathscr{L}_{\text{QCD}} = -\frac{1}{4}G^{\mu\nu}_{a}G^{a}_{\mu\nu} + \bar{\Psi}_{f}\left(i\gamma^{\mu}D_{\mu} - m\right)\Psi_{f}$$

where covariant derivative  $D_{\mu} = \partial_{\mu} - ig_s A^a_{\mu} t^a$  acts on quark fields,  $t^a$  are the 3x3 traceless Hermitian matrices (e.g., the 8 Gell-Mann matrices),  $g_s$  the strong interaction coupling,  $\alpha_s \equiv \frac{g_s^2}{4\pi}$ 

• QCD potential for states in representations R and R' is given by:

$$V(r) = \frac{g_s^2}{4\pi r} t_R^a \otimes t_R^a$$

• To compute V(r) for a  $3_c \otimes 3_c \to \overline{3}_c$ , we use the definition of the scalar  $C_2(R)$ ,  $t_R^a t_R^a \equiv C_2(R)$  **1**, the *quadratic Casimir operator* (NB:  $R_f$  is the final state representation):

$$V(r) = \frac{g_s^2}{4\pi r} \cdot \frac{1}{2} \cdot \left( C_2 \left( R_f \right) - C_2(R) - C_2(R') \right)$$

• Diquarks combine 2 fundamental representation quarks into an anti-fundamental,  $3_C \otimes 3_c \rightarrow \overline{3}_C$ :

$$V(r) = -\frac{2}{3} \frac{g_s^2}{4\pi r} \Longrightarrow$$
 Diquark is bound!





Compare to color singlet attractive potential:

## Higher QCD Fock states in the Proton



Sea quark distributions contribute most at low-*x*<sub>B</sub> Sea quarks created by gluon splitting, creating higher Fock states

$$|p\rangle_{3q} = |uud\rangle = C_1 |uud\rangle + C_2 |u[ud]\rangle$$
$$|p\rangle_{5q} = |uud\rangle + |uud \,\overline{d}d\rangle$$
$$|p\rangle_{7q} = |uud\rangle + |uud \,\overline{d}d\rangle + |uud \,\overline{d}d \,\overline{u}u\rangle$$
$$|p\rangle = |p\rangle_{3q} + |p\rangle_{5}q + |p\rangle_{7q} + |p\rangle_{9q} + \dots$$

"Physical implications of the extrapolation and statistical bootstrap of nucleon structure function ratios  $\frac{F_2^n}{F_2^p}$  for mirror nuclei <sup>3</sup>He and <sup>3</sup>H" H.Valenty, JRW, F.Benmokhtar, D.Higinbotham, A.Parker, E.Seroka, Phys.Rev.C 2023

### Work in progress: Small- $x_B$ physics in proton via Diquark Capture $\overline{d}/\overline{u}$ effects

- Gluon in nucleon splits into  $q \bar{q}$ , adding to the 3-valence quarks
- Creates 5-quark Fock state in the nucleon wavefunction
- $g 
  ightarrow u ar{u}, \, g 
  ightarrow d ar{d}$  produced equally
- Diquark formation [ud] has ~150 MeV binding energy valence u may capture the d quark from  $d\bar{d}$  pair affects  $d/\bar{u}$  ratio
- Creates residual  $\overline{d}$  in the quark sea in the quark-diquark configuration of the proton,  $\overline{d}$  must carry more of the proton's spin than  $\overline{u}$  STAR experiment seems to contradict this



### Strange quarks: Small- $x_B$ physics in proton via Diquark Capture of s

- Gluon in nucleon splits into q ar q, adding to the 3-valence quarks
- Creates 5-quark Fock state in nucleon wavefunction
- $g \rightarrow s\bar{s}$
- Diquark capture of s from  $s\bar{s}$  creates [us] work in progress to check low mass, high binding energy
- Predicts strange seaquark contributions to mass and spin of proton but not to EM properties!



### Matching diquark capture in the proton sea to data:

- Quark spectator counting rules do **not** give a flat slope from  $0.15 \leq x_B \leq 0.4$
- Diquark capture (at this point) predicts a downturn in the slope towards  $\frac{\bar{d}}{\bar{u}} = 1$
- Disagrees with SeaQuest and agrees with earlier experiments...

• Spin asymmetries match for negative  $\overline{d}$  spin contribution, as mandated by spin-statistics constraints on 5-quark Fock state in proton wavefunction - work in progress



- SeaQuest result agrees well with both meson cloud model (light blue band) and statistical model (green band)
  - Which model is more likely?



FIG. 6. The difference of the light sea-quark polarizations as a function of x at a scale of  $Q^2 = 10 \, (\text{GeV}/c)^2$ . The green band shows the NNPDFpol1.1 results [1] and the blue hatched band shows the corresponding distribution after the STAR 2013  $W^{\pm}$  data are included by reweighting.

# Fin

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### Structures in the proton at both high and low $x_B$

Nucleon structure function :  $F_2(x_B) = \sum Q_i^2 x(q_i(x_B) + \bar{q}_i(x_B))$ 

 Structure functions in <sup>3</sup>He & <sup>3</sup>H from JLab's MARATHON experiment Ratio of deep inelastic (DIS) structure functions provides fundamental information about the quark distributions of nucleons - including in our extrapolated regions  $F_{2}^{n} = \begin{bmatrix} P_{2}^{n} \\ 0.7 \\ 0.6 \\ 0.7 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.6 \\ 0.8 \\ X_B \end{bmatrix}$ 

Sea quark distributions contribute most at low-x<sub>B</sub>
 Sea quarks created by gluon splitting, creating higher Fock states



$$|p\rangle_{5q} = |uud\rangle + |uud dd\rangle$$
$$|p\rangle_{7q} = |uud\rangle + |uud \bar{d}d\rangle + |uud \bar{d}d \bar{u}u\rangle$$

"Physical implications of the extrapolation and statistical bootstrap of nucleon structure function ratios  $\frac{F_2^2}{F_2^p}$  for mirror nuclei <sup>3</sup>He and <sup>3</sup>H" H.Valenty, JRW, F.Benmokhtar, D.Higinbotham, A.Parker, E.Seroka, Phys.Rev.C 2023

### Diquark binding energy from Color hyperfine structure

Use  $\Lambda^0$  baryon to find binding energy of [ud] :

$$\mathsf{B}.\mathsf{E}_{\cdot[ud]} = m_u^b + m_d^b + m_s^b - M_{\Lambda^0}$$

Spin-spin interaction contributes to hadron mass; QCD hyperfine interactions:

1. 
$$M_{\text{(baryon)}} = \sum_{i=1}^{3} m_i + a' \sum_{i < j} \left( \sigma_i \cdot \sigma_j \right) / m_i m_j$$

2. 
$$M_{(\text{meson})} = m_1 + m_2 + a (\sigma_1 \cdot \sigma_2) / m_1 m_2$$

(de Rujula, Georgi & Glashow 1975, Gasiorowicz & Rosner 1981, Karliner & Rosner 2014)

Effective masses of light quarks are found using Eq.1 and fitting to measured baryon masses:

$$m_u^b = m_d^b \equiv m_q^b = 363 \text{ MeV}, \ m_s^b = 538 \text{ MeV}$$

$$\mathsf{B.E.}_{[ud]} = m_u^b + m_d^b + m_s^b - M_\Lambda = 148 \pm 9 \text{ MeV}$$

[NB: Diquark-carrying baryons  $\Lambda_c$ ,  $\Sigma_c^+$ ,  $\Sigma_c^0$ ,  $\Sigma_c^- \Longrightarrow \sim 159 \pm 10 \text{ MeV}$ ]

Relevant diquark-carrying baryons:  $\Lambda$ ,  $\Sigma^+$ ,  $\Sigma^0$ ,  $\Sigma^-$ 



TABLE I: Diquark	properties
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Diquark   Bi	nding Energy (Me	eV)   Mass (MeV)   I	sospir	n $I   \text{Spin } S$
[ud]	$148\pm9$	$\mid 578 \pm 11 \mid$	0	0
(ud)	0	$ 776 \pm 11 $	1	1
(uu)	0	$776 \pm 11$	1	1
(dd)	0	$776 \pm 11$	1	1

Uncertainties calculated using average light quark mass errors  $\Delta m_q = 5 \ MeV \ [37]$ 

TABLE II: Relevant  $SU(3)_C$  hyperfine structure baryons [28]

Baryon	Diquark-Quark content	$   Mass (MeV)   I (J^P) $
Λ	[ud]s	$\left  1115.683 \pm 0.006 \left  0 \left( \frac{1}{2}^+ \right) \right. \right.$
$\Sigma^+$	(uu)s	$1189.37 \pm 0.07 \left  1 \left( \frac{1}{2}^+ \right) \right $
$\Sigma^0$	(ud)s	$1192.642 \pm 0.024 1 \left(\frac{1}{2}^{+}\right)$
$\Sigma^{-}$	(dd)s	$\left  1197.449 \pm 0.030 \right  1 \left( \frac{1}{2}^+ \right)$

 $I(J^P)$  denotes the usual isospin I, total spin J and parity P quantum numbers, all have L=0 therefore J=S

"Diquark Induced Short-Range Correlations & the EMC Effect," JRW, Nucl.Phys.A 2023