

HIGHER TWIST PROTON DISTRIBUTIONS IN THE LIGHT-FRONT QUARK-DIQUARK MODEL

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in collaboration with
Dr. Harleen Dahiya

Plan

Introduction

Light-Front Quark-Diquark Model

TMD Correlator and Parameterization

Result Discussion



Plan

Introduction

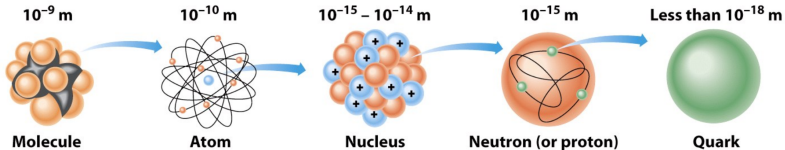
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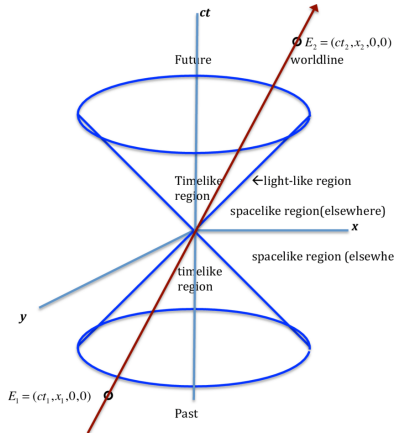
Introduction

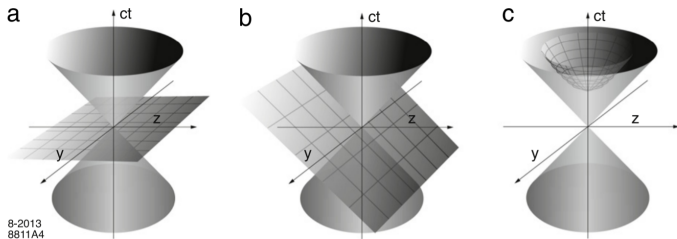


- The theory of the strong interaction which provides the fundamental description of hadronic structure and dynamics in terms of their elementary quarks and gluons degrees of freedom is **Quantum Chromodynamics (QCD)**.
- The foremost problem of hadron physics is to unravel the internal structure of hadron.



- From Special Theory of Relativity:
 - Space and time independently are not invariant quantities.
 - Rather space-time is an invariant object.





(a) the instant form, (b) the front form, (c) the point form.

Their initial surfaces are

a) $x^0 = 0$

b) $x^0 + x^3 = 0$

c) $x^2 = a^2 > 0, x^0 > 0$

- S. J. Brodsky, G. F. de Teramond, H. G. Dosch and J. Erlich, *Phys. Rept.* 584, (2015)



Why Light Front?

- It is an Ideal Framework to describe theoretically the hadronic structure in terms of quarks and gluons. It can overcome many obstacles and has many advantages:
 - Simple vacuum structure $\sim \langle 0|0\rangle = 0$
 - A dynamical system is characterized by ten fundamental quantities: energy, momentum, angular momentum and boost.
 \sim **seven out of which are kinematical.** It allows unambiguous definition of the partonic content of a hadron, exact formulae for form factors, physics of angular momentum of constituents.
 - Dispersion Relation (for ON shell particles)

$$k^- = \frac{(k_\perp)^2 + m^2}{k^+}$$

\sim **no square root factor.**

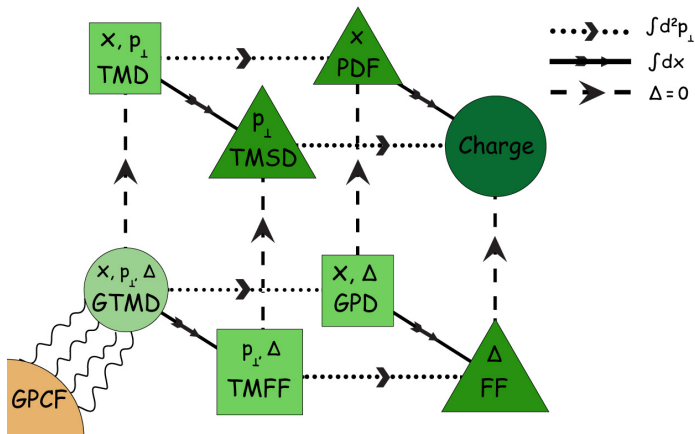


Light-Front Coordinates

- A generic four Vector x^μ in light-cone coordinates is describe as $x^\mu = (x^-, x^+, x_\perp)$.
- $x^+ = x^0 + x^3$ is called as **light-front time**.
- $x^- = x^0 - x^3$ is called as **light-front longitudinal space variable**.
- $x^\perp = (x^1, x^2)$ is the **transverse variable**.
- Similarly, we can define the **longitudinal momentum** $k^+ = k^0 + k^3$ and **light-front energy** $k^- = k^0 - k^3$.



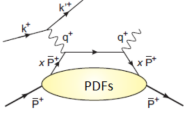
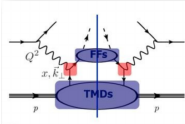
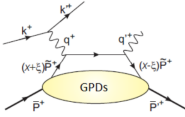
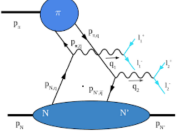
Distribution Functions



- S. Sharma and H. Dahiya, Eur. Phys. J. A 59, 235 (2023)

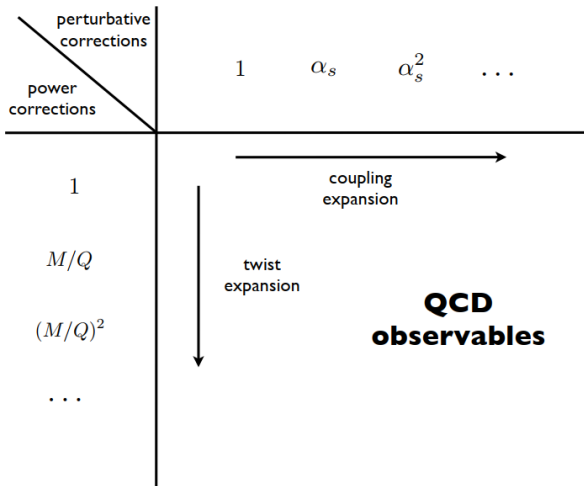


Experiments related with Distribution Functions

PDFs	TMDs	GPDs	GTMDs
1-D information	3-D information	3-D information	6-D information
 <p>DIS</p>	 <p>SIDIS</p>	 <p>DVCS</p>	 <p>DDY</p>



Twist?



- <https://inspirehep.net/literature/1493030>



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Light-Front Quark-Diquark Model

- In this model the proton is described as an aggregate of an **active quark** and a **diquark spectator** of definite mass.
- The proton has spin-flavor $SU(4)$ structure and it has been expressed as a made up of **isoscalar-scalar diquark singlet** $|u S^0\rangle$, **isoscalar-vector diquark** $|u A^0\rangle$ and **isovector-vector diquark** $|d A^1\rangle$ states as

$$|P; \pm\rangle = C_S |u S^0\rangle^\pm + C_V |u A^0\rangle^\pm + C_{VV} |d A^1\rangle^\pm.$$

Here, the scalar and vector diquark has been denoted by S and A respectively. Their isospin has been represented by the superscripts on them.

- *T. Maji and D. Chakrabarti, Phys. Rev. D 94, 094020 (2016)*



Light-Front Quark-Diquark Model

- For the scalar $|\nu S\rangle^\pm$ and vector diquark $|\nu A\rangle^\pm$ case, the expansion of the Fock-state in two particles for $J^z = \pm 1/2$ can be specified as

$$|\nu S\rangle^\pm = \sum_{\lambda^q} \int \frac{dx d^2\mathbf{p}_\perp}{2(2\pi)^3 \sqrt{x(1-x)}} \psi_{\lambda^q}^{\pm(\nu)}(x, \mathbf{p}_\perp) \left| \lambda^q, \lambda^s; xP^+, \mathbf{p}_\perp \right\rangle,$$

$\lambda^s = 0$ (singlet)

$$|\nu A\rangle^\pm = \sum_{\lambda^q} \sum_{\lambda^A} \int \frac{dx d^2\mathbf{p}_\perp}{2(2\pi)^3 \sqrt{x(1-x)}} \psi_{\lambda^q \lambda^A}^{\pm(\nu)}(x, \mathbf{p}_\perp) \left| \lambda^q, \lambda^A; xP^+, \mathbf{p}_\perp \right\rangle.$$

$\lambda^A = \pm 1, 0$ (triplet)



Light-Front Quark-Diquark Model

	λ^q	λ^{Sp}	LFWFs for $J^z = +1/2$	LFWFs for $J^z = -1/2$
S	+1/2	0	$\psi_+^{+(\nu)} = N_S \varphi_1^{(\nu)}$	$\psi_+^{-(\nu)} = N_S \left(\frac{p^1 - ip^2}{xM} \right) \varphi_2^{(\nu)}$
	-1/2	0	$\psi_-^{+(\nu)} = -N_S \left(\frac{p^1 + ip^2}{xM} \right) \varphi_2^{(\nu)}$	$\psi_-^{-(\nu)} = N_S \varphi_1^{(\nu)}$
A	+1/2	+1	$\psi_{++}^{+(\nu)} = N_1^{(\nu)} \sqrt{\frac{2}{3}} \left(\frac{p^1 - ip^2}{xM} \right) \varphi_2^{(\nu)}$	$\psi_{++}^{-(\nu)} = 0$
	-1/2	+1	$\psi_{-+}^{+(\nu)} = N_1^{(\nu)} \sqrt{\frac{2}{3}} \varphi_1^{(\nu)}$	$\psi_{-+}^{-(\nu)} = 0$
	+1/2	0	$\psi_{+0}^{+(\nu)} = -N_0^{(\nu)} \sqrt{\frac{1}{3}} \varphi_1^{(\nu)}$	$\psi_{+0}^{-(\nu)} = N_0^{(\nu)} \sqrt{\frac{1}{3}} \left(\frac{p^1 - ip^2}{xM} \right) \varphi_2^{(\nu)}$
	-1/2	0	$\psi_{-0}^{+(\nu)} = N_0^{(\nu)} \sqrt{\frac{1}{3}} \left(\frac{p^1 + ip^2}{xM} \right) \varphi_2^{(\nu)}$	$\psi_{-0}^{-(\nu)} = N_0^{(\nu)} \sqrt{\frac{1}{3}} \varphi_1^{(\nu)}$
	+1/2	-1	$\psi_{+-}^{+(\nu)} = 0$	$\psi_{+-}^{-(\nu)} = -N_1^{(\nu)} \sqrt{\frac{2}{3}} \varphi_1^{(\nu)}$
	-1/2	-1	$\psi_{--}^{+(\nu)} = 0$	$\psi_{--}^{-(\nu)} = N_1^{(\nu)} \sqrt{\frac{2}{3}} \left(\frac{p^1 + ip^2}{xM} \right) \varphi_2^{(\nu)}$



Light-Front Quark-Diquark Model

- Generic ansatz of LFWFs $\varphi_i^{(\nu)}(x, \mathbf{p}_\perp)$ is being adopted from the **soft-wall AdS/QCD prediction** and the parameters a_i^ν , b_i^ν and δ^ν are established as

$$\varphi_i^{(\nu)}(x, \mathbf{p}_\perp) = \frac{4\pi}{\kappa} \sqrt{\frac{\log(1/x)}{1-x}} x^{a_i^\nu} (1-x)^{b_i^\nu} \exp \left[-\delta^\nu \frac{\mathbf{p}_\perp^2}{2\kappa^2} \frac{\log(1/x)}{(1-x)^2} \right].$$



Input Parameters

Parameter	u	d
C_S^2	1.3872	0
C_V^2	0.6128	0
C_{VV}^2	0	1
N_S	2.0191	0
N_0^ν	3.2050	5.9423
N_1^ν	0.9895	1.1616
a_1^ν	0.280 ± 0.001	0.5850 ± 0.0003
b_1^ν	0.1716 ± 0.0051	0.7000 ± 0.0002
a_2^ν	0.84 ± 0.02	$0.9434^{+0.0017}_{-0.0013}$
b_2^ν	0.2284 ± 0.0035	$0.64^{+0.0082}_{-0.0022}$
δ^ν	1	1
κ	0.4	0.4



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TMD Correlator

- The **unintegrated quark-quark correlator** in the light-front formalism for **SIDIS** is defined as

$$\Phi_{[\Lambda^{N_i} \Lambda^{N_f}]}^{\nu[\Gamma]} = \frac{1}{2} \int \frac{dz^- d^2 z_T}{2(2\pi)^3} e^{ip \cdot z} \langle P; \Lambda^{N_f} | \bar{\psi}^{\nu}(0) \Gamma \mathcal{W}_{[0,z]} \psi^{\nu}(z) | P; \Lambda^{N_i} \rangle \Big|_{z^+=0}.$$

- $|P; \Lambda^{N_i}\rangle$ and $|P; \Lambda^{N_f}\rangle$ are the **initial** and **final** states of the proton having momentum P with helicities Λ^{N_i} and Λ^{N_f} , respectively.
- The **momentum** of the proton (P), struck quark (p) and diquark (P_X) are

$$P \equiv \left(P^+, \frac{M^2}{P^+}, \mathbf{0} \right),$$

$$p \equiv \left(xP^+, \frac{p^2 + |\mathbf{p}_{\perp}|^2}{xP^+}, \mathbf{p}_{\perp} \right),$$

$$P_X \equiv \left((1-x)P^+, P_X^-, -\mathbf{p}_{\perp} \right).$$

- *S. Meissner, A. Metz and M. Schlegel, JHEP 08, 056 (2009)*



TMD Parameterization for proton at twist-3

$$\Phi_{[\Lambda^{N_i} \Lambda^{N_f}]}^{[1]} = \frac{1}{2P^+} \bar{u}(P, \Lambda^{N_f}) \left[e^\nu(x, \mathbf{p}_\perp^2) - \frac{i\sigma^{i+} p_T^j}{P^+} e_T^{\perp\nu}(x, \mathbf{p}_\perp^2) \right] u(P, \Lambda^{N_i}),$$

$$\Phi_{[\Lambda^{N_i} \Lambda^{N_f}]}^{[\gamma_5]} = \frac{1}{2P^+} \bar{u}(P, \Lambda^{N_f}) \left[-\frac{i\sigma^{i+} \gamma_5 p_T^j}{P^+} e_T^{\nu}(x, \mathbf{p}_\perp^2) - i\sigma^{+-} \gamma_5 e_L^{\nu}(x, \mathbf{p}_\perp^2) \right] u(P, \Lambda^{N_i}),$$

$$\Phi_{[\Lambda^{N_i} \Lambda^{N_f}]}^{[\gamma^j]} = \frac{1}{2P^+} \bar{u}(P, \Lambda^{N_f}) \left[\frac{p_T^j}{M} f^{\perp\nu}(x, \mathbf{p}_\perp^2) + \frac{M i\sigma^{j+}}{P^+} f_T^{\prime\nu}(x, \mathbf{p}_\perp^2) \right. \\ \left. + \frac{p_T^j i\sigma^{k+} p_T^k}{M P^+} f_T^{\perp\nu}(x, \mathbf{p}_\perp^2) + \frac{i\sigma^{ij} p_T^i}{M} f_L^{\perp\nu}(x, \mathbf{p}_\perp^2) \right] u(P, \Lambda^{N_i}),$$

$$\Phi_{[\Lambda^{N_i} \Lambda^{N_f}]}^{[\gamma^j \gamma_5]} = \frac{1}{2P^+} \bar{u}(P, \Lambda^{N_f}) \left[\frac{i\varepsilon_T^{ij} p_T^i}{M} g^{\perp\nu}(x, \mathbf{p}_\perp^2) + \frac{M i\sigma^{j+} \gamma_5}{P^+} g_T^{\prime\nu}(x, \mathbf{p}_\perp^2) \right. \\ \left. + \frac{p_T^j i\sigma^{k+} \gamma_5 p_T^k}{M P^+} g_T^{\perp\nu}(x, \mathbf{p}_\perp^2) + \frac{p_T^j i\sigma^{+-} \gamma_5}{M} g_L^{\perp\nu}(x, \mathbf{p}_\perp^2) \right] u(P, \Lambda^{N_i}),$$

T-even TMDs

T-odd TMDs



TMD Parameterization for proton at twist-3

$$\Phi_{[\Lambda^{N_i} \Lambda^{N_f}]}^{[i\sigma^{ij} \gamma_5]} = -\frac{i\epsilon_T^{ij}}{2P^+} \bar{u}(P, \Lambda^{N_F}) \left[-h^\nu(x, \mathbf{p}_\perp^2) + \frac{i\sigma^{k+} p_T^k}{P^+} h_T^{\perp\nu}(x, \mathbf{p}_\perp^2) \right] u(P, \Lambda^{N_i}),$$
$$\Phi_{[\Lambda^{N_i} \Lambda^{N_f}]}^{[i\sigma^{+-} \gamma_5]} = \frac{1}{2P^+} \bar{u}(P, \Lambda^{N_F}) \left[\frac{i\sigma^{i+} \gamma_5 p_T^i}{P^+} h_T^\nu(x, \mathbf{p}_\perp^2) + i\sigma^{+-} \gamma_5 h_L^\nu(x, \mathbf{p}_\perp^2) \right] u(P, \Lambda^{N_i}).$$

T-even TMDs

T-odd TMDs



TMD Parameterization for proton at twist-4

$$\Phi_{[\Lambda^{N_i} \Lambda^{N_f}]}^{[\gamma^-]} = \frac{M}{2(P^+)^2} \bar{u}(P, \Lambda^{N_f}) \left[f_3^\nu(x, \mathbf{p}_\perp^2) - \frac{i\sigma^{i+} p_T^i}{P^+} f_{3T}^{\perp\nu}(x, \mathbf{p}_\perp^2) \right] u(P, \Lambda^{N_i}),$$

$$\Phi_{[\Lambda^{N_i} \Lambda^{N_f}]}^{[\gamma^- \gamma_5]} = \frac{M}{2(P^+)^2} \bar{u}(P, \Lambda^{N_f}) \left[\frac{i\sigma^{i+} \gamma_5 p_T^i}{P^+} g_{3T}^\nu(x, \mathbf{p}_\perp^2) + i\sigma^{+-} \gamma_5 g_{3L}^\nu(x, \mathbf{p}_\perp^2) \right] u(P, \Lambda^{N_i}),$$

$$\Phi_{[\Lambda^{N_i} \Lambda^{N_f}]}^{[i\sigma^{j-} \gamma_5]} = \frac{M}{2(P^+)^2} \bar{u}(P, \Lambda^{N_f}) \left[\frac{i\varepsilon_T^{ij} p_T^i}{M} h_3^{\perp\nu}(x, \mathbf{p}_\perp^2) + \frac{M i\sigma^{j+} \gamma_5}{P^+} h_{3T}^\nu(x, \mathbf{p}_\perp^2) \right. \\ \left. + \frac{p_T^j i\sigma^{k+} \gamma_5 p_T^k}{M P^+} h_{3T}^{\perp\nu}(x, \mathbf{p}_\perp^2) + \frac{p_T^j i\sigma^{+-} \gamma_5}{M} h_{3L}^{\perp\nu}(x, \mathbf{p}_\perp^2) \right] u(P, \Lambda^{N_i}).$$

T-even TMDs

T-odd TMDs



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TMD parameterization at higher twist

Γ	$\Phi_{++}^{[\Gamma]} + \Phi_{--}^{[\Gamma]}$	$\Phi_{++}^{[\Gamma]} - \Phi_{--}^{[\Gamma]}$	$\Phi_{-+}^{[\Gamma]} + \Phi_{+-}^{[\Gamma]}$	$\Phi_{-+}^{[\Gamma]} - \Phi_{+-}^{[\Gamma]}$
1	$\frac{2M}{P^+} e$	0	$\frac{-2ip_y}{P^+} e_T^\perp$	$\frac{2p_x}{P^+} e_T^\perp$
γ_5	0	$\frac{-4M}{P^+} e_L$	$\frac{-2p_x}{P^+} e_T$	$\frac{2ip_y}{P^+} e_T$
γ^1	$\frac{2p_x}{P^+} f^\perp$	$\frac{2ip_y}{P^+} f_L^\perp$	$\frac{2ip_x p_y}{MP^+} f_T^\perp$	$\frac{-2M}{P^+} [f_T^\perp + \frac{p_x^2}{M^2} f_T^\perp]$
γ^2	$\frac{2p_y}{P^+} f^\perp$	$\frac{2ip_x}{P^+} f_L^\perp$	$\frac{2iM}{P^+} [f_T^\perp + \frac{p_y^2}{M^2} f_T^\perp]$	$\frac{-2ip_x p_y}{MP^+} f_T^\perp$
$\gamma^1 \gamma_5$	$\frac{-2ip_y}{P^+} g^\perp$	$\frac{4p_x}{P^+} g_L^\perp$	$\frac{2M}{P^+} [g_T^\perp + \frac{p_x^2}{M^2} g_T^\perp]$	$\frac{-2ip_x p_y}{MP^+} g_T^\perp$
$\gamma^2 \gamma_5$	$\frac{2ip_x}{P^+} g^\perp$	$\frac{4p_y}{P^+} g_L^\perp$	$\frac{2p_x p_y}{MP^+} g_T^\perp$	$\frac{-2iM}{P^+} [g_T^\perp + \frac{p_y^2}{M^2} g_T^\perp]$
$i\sigma^{12} \gamma_5$	$\frac{2iM}{P^+} h$	0	$\frac{2p_y}{P^+} h_T^\perp$	$\frac{2ip_x}{P^+} h_T^\perp$
$i\sigma^{21} \gamma_5$	$\frac{-2iM}{P^+} h$	0	$\frac{-2p_x}{P^+} h_T^\perp$	$\frac{-2ip_y}{P^+} h_T^\perp$
$i\sigma^{+-} \gamma_5$	0	$\frac{4M}{(P^+)^2} h_L$	$\frac{2p_x}{P^+} h_T$	$\frac{-2ip_y}{P^+} h_T$
γ^-	$\frac{2M^2}{(P^+)^2} f_3$	0	$\frac{-2ip_y M}{(P^+)^2} f_{3T}^\perp$	$\frac{2p_x M}{(P^+)^2} f_{3T}^\perp$
$\gamma^- \gamma_5$	0	$\frac{4M^2}{(P^+)^2} g_{3L}$	$\frac{2p_x M}{(P^+)^2} g_{3T}$	$\frac{-2ip_y M}{(P^+)^2} g_{3T}$
$i\sigma^{1-} \gamma_5$	$\frac{-2ip_y M}{(P^+)^2} h_{3^\perp}$	$\frac{4p_x M}{(P^+)^2} h_{3L}^\perp$	$\frac{2M^2}{(P^+)^2} [h_{3T}^\perp + \frac{p_x^2}{M^2} h_{3T}^\perp]$	$\frac{-2ip_x p_y}{(P^+)^2} h_{3T}^\perp$
$i\sigma^{2-} \gamma_5$	$\frac{2ip_x M}{(P^+)^2} h_{3^\perp}$	$\frac{4p_y M}{(P^+)^2} h_{3L}^\perp$	$\frac{2p_x p_y}{(P^+)^2} h_{3T}^\perp$	$\frac{-2iM^2}{(P^+)^2} [h_{3T}^\perp + \frac{p_y^2}{M^2} h_{3T}^\perp]$

T-even TMDs
T-odd TMDs



Explicit Expressions of TMDs

$$xe^\nu(x, \mathbf{p}_\perp^2) = \frac{1}{16\pi^3} \left(C_S^2 N_s^2 + C_A^2 \left(\frac{2}{3} |N_1^\nu|^2 + \frac{1}{3} |N_0^\nu|^2 \right) \right) \frac{m}{M} \left[|\varphi_1^\nu|^2 + \frac{p_\perp^2}{x^2 M^2} |\varphi_2^\nu|^2 \right],$$

$$xf^{\perp\nu}(x, \mathbf{p}_\perp^2) = \frac{1}{16\pi^3} \left(C_S^2 N_s^2 + C_A^2 \left(\frac{2}{3} |N_1^\nu|^2 + \frac{1}{3} |N_0^\nu|^2 \right) \right) \left[|\varphi_1^\nu|^2 + \frac{p_\perp^2}{x^2 M^2} |\varphi_2^\nu|^2 \right],$$

$$xg_L^{\perp\nu}(x, \mathbf{p}_\perp^2) = \frac{1}{32\pi^3} \left(C_S^2 N_s^2 + C_A^2 \left(-\frac{2}{3} |N_1^\nu|^2 + \frac{1}{3} |N_0^\nu|^2 \right) \right) \left[|\varphi_1^\nu|^2 - \frac{p_\perp^2}{x^2 M^2} |\varphi_2^\nu|^2 - \frac{2m}{xM} |\varphi_1^\nu| |\varphi_2^\nu| \right],$$

$$xg_T^{\prime\nu}(x, \mathbf{p}_\perp^2) = \frac{1}{16\pi^3} \left(C_S^2 N_s^2 - \frac{1}{3} C_A^2 |N_0^\nu|^2 \right) \frac{m}{M} \left[|\varphi_1^\nu|^2 + \frac{p_\perp^2}{x^2 M^2} |\varphi_2^\nu|^2 \right],$$

$$xg_T^{\perp\nu}(x, \mathbf{p}_\perp^2) = \frac{1}{8\pi^3} \left(C_S^2 N_s^2 - \frac{1}{3} C_A^2 |N_0^\nu|^2 \right) \left[\frac{1}{x} |\varphi_1^\nu| |\varphi_2^\nu| - \frac{m}{x^2 M} |\varphi_2^\nu|^2 \right],$$



Explicit Expressions of TMDs

$$xh_L^\nu(x, \mathbf{p}_\perp^2) = \frac{1}{16\pi^3} \left(C_S^2 N_s^2 + C_A^2 \left(-\frac{2}{3} |N_1^\nu|^2 + \frac{1}{3} |N_0^\nu|^2 \right) \right) \frac{1}{M} \left[m \left(|\varphi_1^\nu|^2 - \frac{p_\perp^2}{x^2 M^2} |\varphi_2^\nu|^2 \right) + \frac{2p_\perp^2}{xM} |\varphi_1^\nu| |\varphi_2^\nu| \right],$$

$$xh_T^\nu(x, \mathbf{p}_\perp^2) = \frac{1}{8\pi^3} \left(-C_S^2 N_s^2 + \frac{1}{3} C_A^2 |N_0^\nu|^2 \right) \left[|\varphi_1^\nu|^2 - \frac{p_\perp^2}{x^2 M^2} |\varphi_2^\nu|^2 - \frac{2m}{xM} |\varphi_1^\nu| |\varphi_2^\nu| \right],$$

$$xh_T^{\perp\nu}(x, \mathbf{p}_\perp^2) = \frac{1}{16\pi^3} \left(C_S^2 N_s^2 - \frac{1}{3} C_A^2 |N_0^\nu|^2 \right) \left[|\varphi_1^\nu|^2 + \frac{p_\perp^2}{x^2 M^2} |\varphi_2^\nu|^2 \right].$$

- S. Sharma, N. Kumar and H. Dahiya, *Nucl. Phys. B* (2023)



Explicit Expressions of TMDs

$$x^2 f_3^\nu(x, \mathbf{p}_\perp^2) = \frac{1}{16\pi^3} \left(C_S^2 N_s^2 + C_A^2 \left(\frac{2}{3} |N_1^\nu|^2 + \frac{1}{3} |N_0^\nu|^2 \right) \right) \left(\frac{p_\perp^2 + m^2}{M^2} \right) \left[|\varphi_1^\nu|^2 + \frac{p_\perp^2}{x^2 M^2} |\varphi_2^\nu|^2 \right],$$

$$x^2 g_{3L}^\nu(x, \mathbf{p}_\perp^2) = \frac{1}{32\pi^3} \left(C_S^2 N_s^2 + C_A^2 \left(-\frac{2}{3} |N_1^\nu|^2 + \frac{1}{3} |N_0^\nu|^2 \right) \right) \left[(p_\perp^2 - m^2) \left[|\varphi_1^\nu|^2 - \frac{p_\perp^2}{x^2 M^2} |\varphi_2^\nu|^2 \right] - \frac{4mp_\perp^2}{xM} |\varphi_1^\nu| |\varphi_2^\nu| \right],$$

$$x^2 h_{3L}^{\perp\nu}(x, \mathbf{p}_\perp^2) = \frac{1}{16\pi^3 M} \left(C_S^2 N_s^2 + C_A^2 \left(-\frac{2}{3} |N_1^\nu|^2 + \frac{1}{3} |N_0^\nu|^2 \right) \right) \left[m \left[|\varphi_1^\nu|^2 - \frac{p_\perp^2}{x^2 M^2} |\varphi_2^\nu|^2 \right] + \frac{(p_\perp^2 - m^2)}{xM} |\varphi_1^\nu| |\varphi_2^\nu| \right],$$



Explicit Expressions of TMDs

$$x^2 g_{3T}^\nu(x, \mathbf{p}_\perp^2) = \frac{1}{8\pi^3 M} \left(C_S^2 N_s^2 - \frac{1}{3} C_A^2 |N_0^\nu|^2 \right) \left[m \left[|\varphi_1^\nu|^2 - \frac{p_\perp^2}{x^2 M^2} |\varphi_2^\nu|^2 \right] + \frac{(p_\perp^2 - m^2)}{xM} |\varphi_1^\nu| |\varphi_2^\nu| \right],$$

$$x^2 h_{3T}^\nu(x, \mathbf{p}_\perp^2) = \frac{1}{16\pi^3} \left(C_S^2 N_s^2 - \frac{1}{3} C_A^2 |N_0^\nu|^2 \right) \left(\frac{p_\perp^2 + m^2}{M^2} \right) \left[|\varphi_1^\nu|^2 + \frac{p_\perp^2}{x^2 M^2} |\varphi_2^\nu|^2 \right],$$

$$x^2 h_{3T}^{\perp\nu}(x, \mathbf{p}_\perp^2) = -\frac{1}{8\pi^3} \left(C_S^2 N_s^2 - \frac{1}{3} C_A^2 |N_0^\nu|^2 \right) \left[|\varphi_1^\nu|^2 + \frac{m^2}{x^2 M^2} |\varphi_2^\nu|^2 - \frac{2m}{xM} |\varphi_1^\nu| |\varphi_2^\nu| \right].$$

- **S. Sharma and H. Dahiya, *Int. J. Mod. Phys. A* 37, 2250205 (2022)**



Relations between twist-2 and twist-3

$$xe^q(x, \mathbf{p}_\perp^2) = x\tilde{e}^q(x, \mathbf{p}_\perp^2) + \frac{m}{M} f_1^q(x, \mathbf{p}_\perp^2),$$

$$xf^{\perp q}(x, \mathbf{p}_\perp^2) = x\tilde{f}^{\perp q}(x, \mathbf{p}_\perp^2) + f_1^q(x, \mathbf{p}_\perp^2),$$

$$xg_L^{\perp q}(x, \mathbf{p}_\perp^2) = x\tilde{g}_L^{\perp q}(x, \mathbf{p}_\perp^2) + g_1^q(x, \mathbf{p}_\perp^2) + \frac{m}{M} h_{1L}^{\perp q}(x, \mathbf{p}_\perp^2),$$

$$xg_T^{\perp q}(x, \mathbf{p}_\perp^2) = x\tilde{g}_T^{\perp q}(x, \mathbf{p}_\perp^2) + \frac{m}{M} h_1^q(x, \mathbf{p}_\perp^2) - \frac{m}{M} \frac{p_\perp^2}{M^2} h_{1T}^{\perp q}(x, \mathbf{p}_\perp^2),$$

$$xg_T^{\perp q}(x, \mathbf{p}_\perp^2) = x\tilde{g}_T^{\perp q}(x, \mathbf{p}_\perp^2) + g_{1T}^{\perp q}(x, \mathbf{p}_\perp^2) + \frac{m}{M} h_{1T}^{\perp q}(x, \mathbf{p}_\perp^2),$$

$$\frac{xh_L^q(x, \mathbf{p}_\perp^2)}{2} = x\tilde{h}_L^q(x, \mathbf{p}_\perp^2) - \frac{p_\perp^2}{M^2} h_{1L}^{\perp q}(x, \mathbf{p}_\perp^2) + \frac{m}{M} g_1^q(x, \mathbf{p}_\perp^2),$$

$$\frac{xh_T^q(x, \mathbf{p}_\perp^2)}{2} = x\tilde{h}_T^q(x, \mathbf{p}_\perp^2) - h_1^q(x, \mathbf{p}_\perp^2) - \frac{p_\perp^2}{2M^2} h_{1T}^{\perp q}(x, \mathbf{p}_\perp^2) + \frac{m}{2M} g_{1T}^{\perp q}(x, \mathbf{p}_\perp^2),$$

$$xh_T^{\perp q}(x, \mathbf{p}_\perp^2) = x\tilde{h}_T^{\perp q}(x, \mathbf{p}_\perp^2) + h_1^q(x, \mathbf{p}_\perp^2) - \frac{p_\perp^2}{2M^2} h_{1T}^{\perp q}(x, \mathbf{p}_\perp^2).$$



Relations between twist-2 and twist-4

$$x^2 f_3(x, \mathbf{p}_\perp^2) = \left(\frac{p_\perp^2 + m^2}{M^2} \right) f_1^\nu(x, \mathbf{p}_\perp^2),$$

$$x^2 h_{3T}^\nu(x, \mathbf{p}_\perp^2) = \left(\frac{p_\perp^2 + m^2}{M^2} \right) h_{1T}^\nu(x, \mathbf{p}_\perp^2),$$

$$x^2 h_{3L}^{\nu\perp}(x, \mathbf{p}_\perp^2) = - \left(\frac{p_\perp^2 - m^2}{M^2} \right) h_{1L}^{\nu\perp}(x, \mathbf{p}_\perp^2) + \frac{2m}{M} g_{1L}^\nu(x, \mathbf{p}_\perp^2),$$

$$x^2 g_{3L}^\nu(x, \mathbf{p}_\perp^2) = \left(\frac{p_\perp^2 - m^2}{M^2} \right) g_{1L}^\nu(x, \mathbf{p}_\perp^2) + \frac{2mp_\perp^2}{M^3} h_{1L}^{\nu\perp}(x, \mathbf{p}_\perp^2),$$

$$x^2 h_{3T}^{\nu\perp}(x, \mathbf{p}_\perp^2) = \frac{m^2}{M^2} h_{1T}^{\nu\perp}(x, \mathbf{p}_\perp^2) - 2 h_1^\nu(x, \mathbf{p}_\perp^2) + \frac{2m}{M} g_{1T}^\nu(x, \mathbf{p}_\perp^2),$$

$$x^2 g_{3T}^\nu(x, \mathbf{p}_\perp^2) = \left(\frac{p_\perp^2 - m^2}{M^2} \right) g_{1T}^\nu(x, \mathbf{p}_\perp^2) + \frac{mp_\perp^2}{M^3} h_{1T}^{\nu\perp}(x, \mathbf{p}_\perp^2) + \frac{2m}{M} h_1^\nu(x, \mathbf{p}_\perp^2).$$



TMD Amplitude matrix for scalar diquark

- S. Sharma *et al.* , *arXiv:2405.13727 [hep-ph] (2024)*

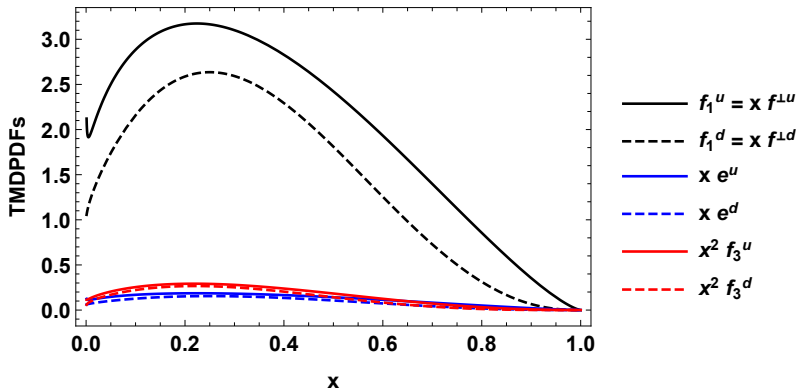
$$A^{\lambda_{Sp}} = \begin{bmatrix} A_{++,++}^{\lambda_{Sp}} & A_{++,+-}^{\lambda_{Sp}} & A_{++,-+}^{\lambda_{Sp}} & A_{++,--}^{\lambda_{Sp}} \\ A_{+-,++}^{\lambda_{Sp}} & A_{+-,+-}^{\lambda_{Sp}} & A_{+-,-+}^{\lambda_{Sp}} & A_{+-,--}^{\lambda_{Sp}} \\ A_{-+,++}^{\lambda_{Sp}} & A_{-+,+-}^{\lambda_{Sp}} & A_{-+,-+}^{\lambda_{Sp}} & A_{-+,-}^{\lambda_{Sp}} \\ A_{--,++}^{\lambda_{Sp}} & A_{--,+-}^{\lambda_{Sp}} & A_{--, -+}^{\lambda_{Sp}} & A_{--,--}^{\lambda_{Sp}} \end{bmatrix}$$

$$A^{\lambda_{Sp}} = \frac{1}{\frac{C_S^2}{16\pi^3}} \begin{bmatrix} h_1 & -\frac{(\mathbf{p}_x + \iota \mathbf{p}_y)}{2M} \varrho & \frac{(\mathbf{p}_x - \iota \mathbf{p}_y)}{2M} \varrho & h_1 \\ -\frac{(\mathbf{p}_x - \iota \mathbf{p}_y)}{2M} \varrho & -\frac{p_\perp^2}{2M^2} h_{1T}^\perp & \frac{(\mathbf{p}_x - \iota \mathbf{p}_y)^2}{2M^2} h_{1T}^\perp & -\frac{(\mathbf{p}_x - \iota \mathbf{p}_y)}{2M} \varrho \\ \frac{(\mathbf{p}_x + \iota \mathbf{p}_y)}{2M} \varrho & \frac{(\mathbf{p}_x + \iota \mathbf{p}_y)^2}{2M^2} h_{1T}^\perp & -\frac{p_\perp^2}{2M^2} h_{1T}^\perp & \frac{(\mathbf{p}_x + \iota \mathbf{p}_y)}{2M} \varrho \\ h_1 & -\frac{(\mathbf{p}_x + \iota \mathbf{p}_y)}{2M} \varrho & \frac{(\mathbf{p}_x - \iota \mathbf{p}_y)}{2M} \varrho & h_1 \end{bmatrix}$$

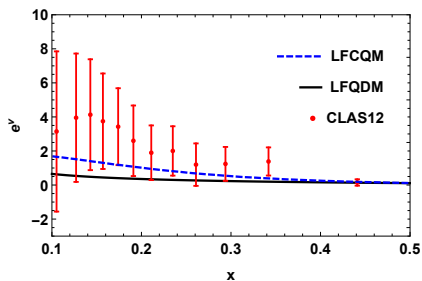
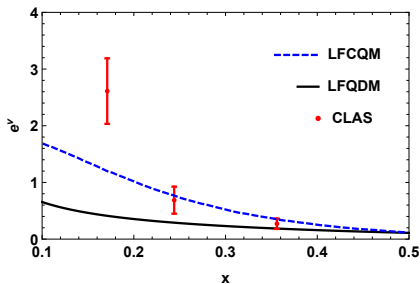
$$A_{\Lambda^{N_f} \lambda^{q_f}, \Lambda^{N_i} \lambda^{q_i}}^{\lambda_{Sp}} = \psi_{\lambda^{q_f} \lambda_{Sp}}^{\Lambda^{N_f \dagger}}(x, \mathbf{p}_\perp) \psi_{\lambda^{q_i} \lambda_{Sp}}^{\Lambda^{N_i}}(x, \mathbf{p}_\perp) \\ \varrho(\text{say}) = g_{1T} = -2h_{1L}^\perp$$



Unpolarized TMDPDFs



Comparison with Phenomenology



LFQDM - 0.09 GeV^2 **S. Sharma, N. Kumar and H. Dahiya, Nucl. Phys. B**
(2023)

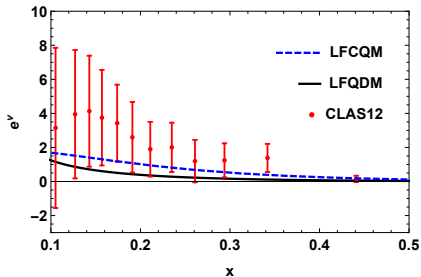
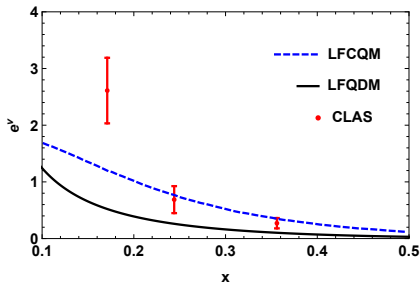
LFCQM - 1 GeV^2 **S. Rodini and B. Pasquini, Nuovo Cim. C 42, 112 (2019)**

CLAS - 1 GeV^2 **A. Courtoy, arXiv:1405.7659 [hep-ph] (2014)**

CLAS12 - 1 GeV^2 **A. Courtoy et al. , Phys. Rev. D 106, 014027 (2022)**



Comparison with Phenomenology



LFQDM - 1 GeV^2 **S. Sharma, N. Kumar and H. Dahiya, Nucl. Phys. B (2023)**

LFCQM - 1 GeV^2 **S. Rodini and B. Pasquini, Nuovo Cim. C 42, 112 (2019)**

CLAS - 1 GeV^2 **A. Courtoy, arXiv:1405.7659 [hep-ph] (2014)**

CLAS12 - 1 GeV^2 **A. Courtoy et al. , Phys. Rev. D 106, 014027 (2022)**





thank you!



Average Transverse Momenta

- The average transverse momenta ($r = 1$) and the average square transverse momenta ($r = 2$) for TMD $\Upsilon^\nu(x, \mathbf{p}_\perp^2)$ in LFQDM is defined as

$$\langle \mathbf{p}_\perp^r(\Upsilon) \rangle^\nu = \frac{\int dx \int d^2 p_\perp p_\perp^r \Upsilon^\nu(x, \mathbf{p}_\perp^2)}{\int dx \int d^2 p_\perp \Upsilon^\nu(x, \mathbf{p}_\perp^2)},$$

TMD w	ν	e^ν	$f^{\perp\nu}$	$g_L^{\perp\nu}$	$g_T^{\prime\nu}$	$g_T^{\perp\nu}$	h_L^ν	h_T^ν	$h_T^{\perp\nu}$	g_T^ν
LFQDM	$\langle p_\perp \rangle^u$	0.99	0.99	1.25	0.99	0.88	1.23	1.25	0.99	1.17
LFQDM	$\langle p_\perp \rangle^d$	1.00	1.00	1.13	1.00	1.07	1.25	1.13	1.00	1.18
LFCQM	$\langle p_\perp \rangle^y$	0.92	0.92

LFQDM - S. Sharma, N. Kumar and H. Dahiya, Nucl. Phys. B (2023)

LFCQM - C. Lorcé, B. Pasquini and P. Schweitzer, JHEP 01, 103 (2015)



Average Transverse Momenta

TMD w	v	e^v	$f^{\perp v}$	$g_L^{\perp v}$	$g_T^{\prime v}$	$g_T^{\perp v}$	h_L^v	h_T^v	$h_T^{\perp v}$	g_T^v
LFQDM	$\langle p_{\perp}^2 \rangle^u$	0.94	0.94	1.30	0.94	0.77	1.39	1.30	0.94	1.28
LFQDM	$\langle p_{\perp}^2 \rangle^d$	0.96	0.96	1.15	0.96	1.06	1.43	1.15	0.96	1.30
LFCQM	$\langle p_{\perp}^2 \rangle^v$	0.86	0.86
Bag Model	$\langle p_{\perp}^2(x_v) \rangle_{Gauss}$	0.68	0.94	1.11	1.11	1.01	1.11	0.94	0.84

LFQDM - S. Sharma, N. Kumar and H. Dahiya, Nucl. Phys. B (2023)

LFCQM - C. Lorcé, B. Pasquini and P. Schweitzer, JHEP 01, 103 (2015)

Bag Model - H. Avakian, A. Efremov, P. Schweitzer and F. Yuan, Phys. Rev. D 81, 074035 (2010)



Average Transverse Momenta

MODEL	f_3^ν (LFQDM)	f_3^ν (LFCQM)
$\langle p_\perp \rangle^u$	0.26	0.28
$\langle p_\perp \rangle^d$	0.27	0.28
$\langle p_\perp^2 \rangle^u$	0.073	0.11
$\langle p_\perp^2 \rangle^d$	0.078	0.11

LFQDM - S. Sharma and H. Dahiya, Int. J. Mod. Phys. A 37, 2250205 (2022)

LFCQM - C. Lorcé, B. Pasquini and P. Schweitzer, JHEP 01, 103 (2015)

