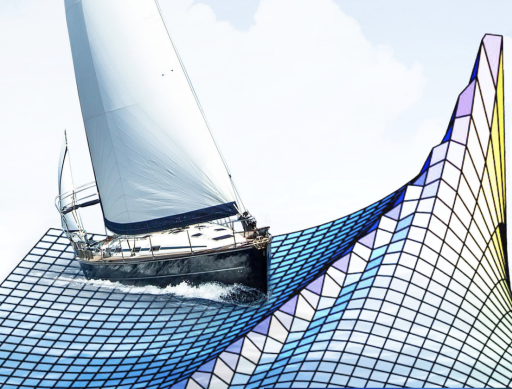


FIRST INTERNATIONAL SCHOOL OF HADRON FEMTOGRAPHY

Jefferson Lab | September 16 - 25, 2024

Experimental Methods

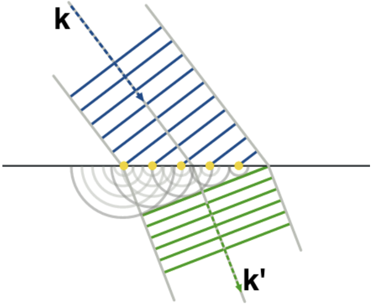
Sep 21st 2024

F.-X. Girod

Lecture 3
with Charles Hyde

Diffraction and Imaging

Huygens-Kirchhoff-Fresnel principle

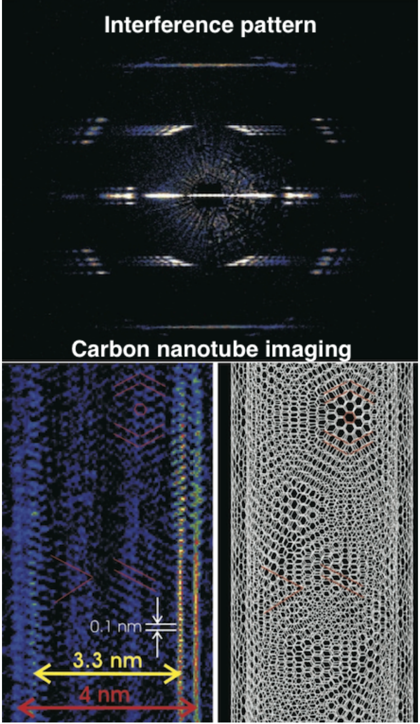


$\vec{q} = \vec{k} - \vec{k'}$

The interference pattern is given by the superposition of spherical wavelets

$$f(\Omega_{\vec{q}}) = \int \frac{d^3\vec{r}}{(2\pi)^3} F(\vec{r}) e^{i\vec{q}\cdot\vec{r}}$$

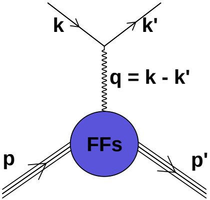
Fourier imaging



Elastic scattering

Form Factors

Probing deeper using virtual photons



$$J^{\mu}_{\mathrm{EM}} = F_1 \gamma^{\mu} + \frac{\kappa}{2M} F_2 i \sigma^{\mu\nu} q_{\nu}$$

$$\frac{d\sigma}{d\Omega} = \frac{\sigma_{\mathrm{Mott}}}{\epsilon(1+\tau)} \left[\tau G_{\mathrm{M}}^2 + \epsilon G_{\mathrm{E}}^2\right]$$

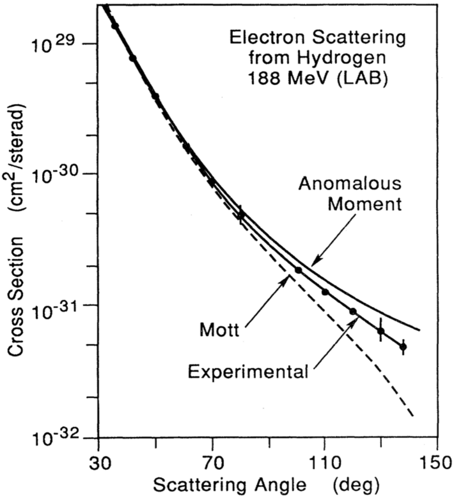
$$\tau = \frac{Q^2}{4M^2}$$

$$Q^2 = -\left(k - k'\right)^2 = -m^2_{\gamma^*}$$

$$\frac{1}{\epsilon} = 1 + 2(1+\tau)\tan^2\frac{\theta_e}{2}$$

$$G_{\mathrm{E}} = F_1 - \tau F_2$$

$$G_{\mathrm{M}} = F_1 + F_2$$



Hofstadter Nobel prize 1961

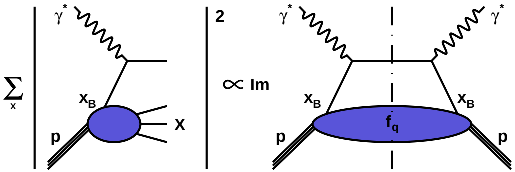
"The best fit in this figure indicates an rms radius close to $0.74 \pm 0.24 \times 10^{-13}$ cm."

Imaging in transverse impact parameter space

Deeply Inelastic Scattering

Parton Distributions

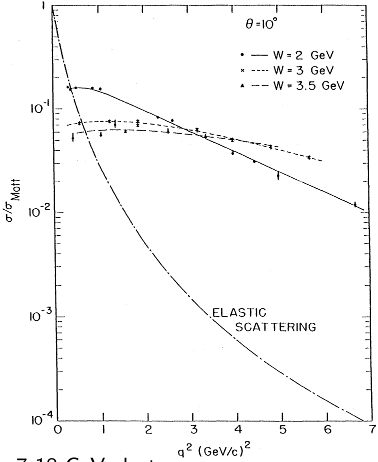
Optical theorem



The total cross section is given by
the imaginary part of the forward amplitude

$$\nu = E_{\gamma^*} \quad , \quad x_B = \frac{Q^2}{2M\nu}$$

$$\sigma_{\rm DIS}(x_B, \not{Q^2}) \rightarrow \text{ scaling, point-like constituents}$$



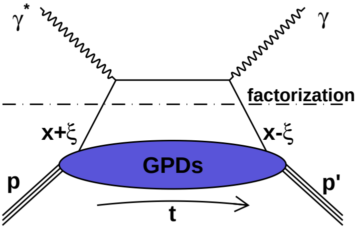
Discovery of quarks, SLAC-MIT group, 7-18 GeV electron
Friedman, Kendall, Taylor, Nobel prize 1990

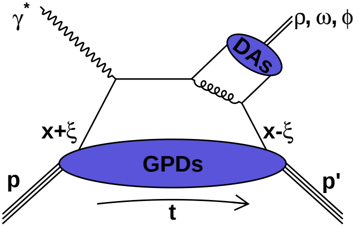
$$\lim_{Q^2\to\infty} \sigma_{\rm DIS}(x_B) = \int_{x_B}^{1} \frac{d\xi}{\xi} \sum_a f_a(\xi,\mu) \hat{\sigma}^a \left(\frac{x_B}{\xi}, \frac{Q}{\mu}\right)$$

1-D distribution in longitudinal momentum space

Deep Exclusive Scattering

Generalized Parton Distributions





$$\gamma^* p \to \gamma p', \; \gamma^* p \to \begin{cases} \rho p' \\ \omega p' \\ \phi p' \end{cases}$$

Bjorken regime :

$Q^2 \to \infty$, x_B fixed

t fixed $\ll Q^2$, $\xi \to \frac{x_B}{2-x_B}$

$$\frac{P^+}{2\pi} \int dy^- \, e^{ixP^+y^-} \langle p' | \bar{\psi}_q(0) \gamma^+ (1+\gamma^5) \psi(y) | p \rangle$$

$$= \bar{N}(p') \left[H^q(x,\xi,t)\gamma^+ + E^q(x,\xi,t) i\sigma^{+\nu} \frac{\Delta_\nu}{2M} \right.$$

$$\left. + \tilde{H}^q(x,\xi,t)\gamma^+\gamma^5 + \tilde{E}^q(x,\xi,t)\gamma^5 \frac{\Delta^+}{2M} \right] N(p)$$

spin	N no flip	N flip
q no flip	H	E
q flip	$\tilde{H}$	$\tilde{E}$

3-D Imaging conjointly in transverse impact parameter and longitudinal momentum

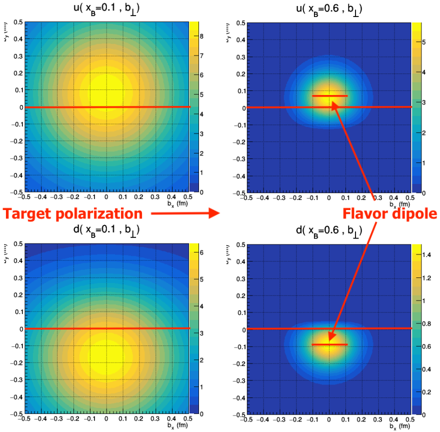
GPDs and Transverse Imaging

(x_B, t) correlations

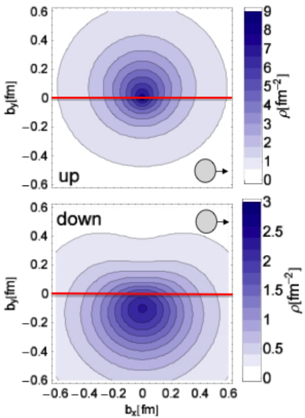
$$q_X(x, \vec{b}_\perp) = \int \frac{d^2\vec{\Delta}_\perp}{(2\pi)^2} \left[H(x,0,t) - \frac{E(x,0,t)}{2M} \frac{\partial}{\partial b_y} \right] e^{-i\vec{\Delta}_\perp \cdot \vec{b}_\perp}$$

M. Burkardt, Int. J. Mod. Phys. **A 18** 173 (2003)

QCDSF coll. PRL**98** 222001 (2007)



Lattice calculation

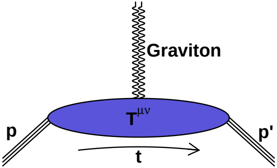


Energy Momentum Tensor

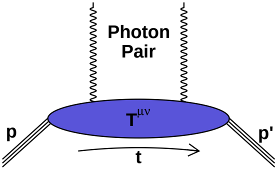
Gravitational Form Factors definition :

$$\langle p'|\hat{T}^q_{\mu\nu}|p\rangle = \bar{N}(p')\left[\ M^q_2(t)\ \frac{P_\mu P_\nu}{M} + J^q(t)\ \frac{\imath(P_\mu\sigma_{\nu\rho}+P_\nu\sigma_{\mu\rho})\Delta^\rho}{2M} + d^q_1(t)\ \frac{\Delta_\mu\Delta_\nu - g_{\mu\nu}\Delta^2}{5M}\ \right]N(p)$$

Confinement forces from space-space components of EMT
The graviton with spin 2 couples directly to EMT
But gravity is too weak to produce count rates in the detector



We can construct a spin 2 operator using two spin 1 operators
$\rightarrow$ use a process with two photons to measure the EMT?
X. Ji *PRL* **78** 610 (1997) ; M. Polyakov *PLB* **555** 57 (2003)



GPDs and Energy Momentum Tensor

(x, ξ) correlations

Form Factors accessed *via* second x-moments :

$$\langle p'|\hat{T}^q_{\mu\nu}|p\rangle = \bar{N}(p')\left[M_2^q(t)\,\frac{P_\mu P_\nu}{M} + J^q(t)\,\frac{\imath(P_\mu\sigma_{\nu\rho}+P_\nu\sigma_{\mu\rho})\Delta^\rho}{2M} + d_1^q(t)\,\frac{\Delta_\mu\Delta_\nu - g_{\mu\nu}\Delta^2}{5M}\right] N(p)$$

Angular momentum distribution

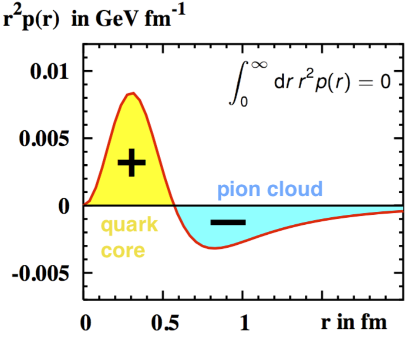
$$J^q(t) = \frac{1}{2}\int_{-1}^{1} \mathrm{d}x\, x\,[H^q(x,\xi,t) + E^q(x,\xi,t)]$$

Mass and force/pressure distributions

$$M_2^q(t) + \frac{4}{5} d_1(t)\xi^2 = \frac{1}{2}\int_{-1}^{1} \mathrm{d}x\, x H^q(x,\xi,t)$$

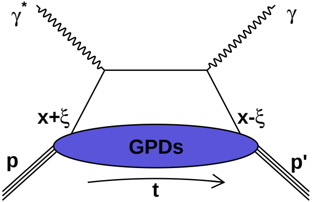
$$d_1(t) = 15M \int \mathrm{d}^3\vec{r}\,\frac{j_0(r\sqrt{-t})}{2t}\, p(r)$$

Distribution of pressure

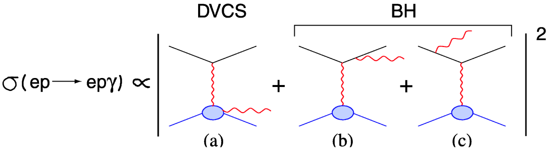


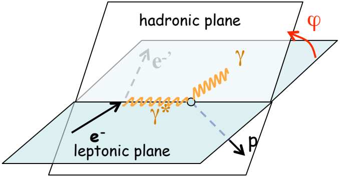
Deeply Virtual Compton Scattering

The cleanest GPD probe at low and medium energies



$$ep \to ep\gamma$$





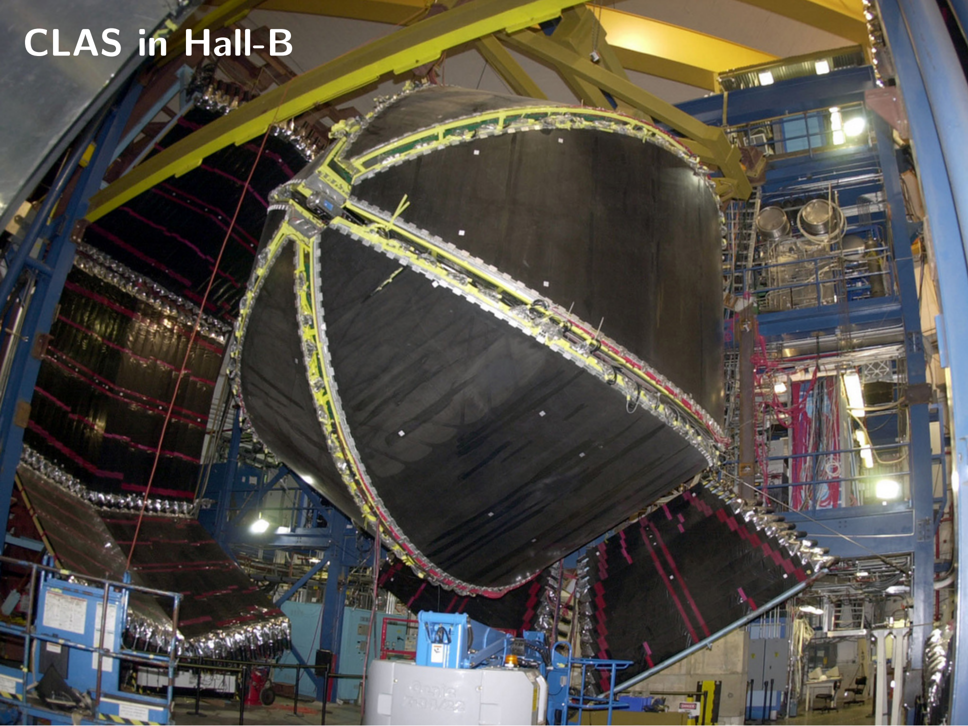
$$A_{\mathrm{LU}} = \frac{d^4\sigma^{\rightarrow} - d^4\sigma^{\leftarrow}}{d^4\sigma^{\rightarrow} + d^4\sigma^{\leftarrow}} \overset{\text{twist-2}}{\approx} \frac{\alpha \sin\phi}{1 + \beta\cos\phi}$$

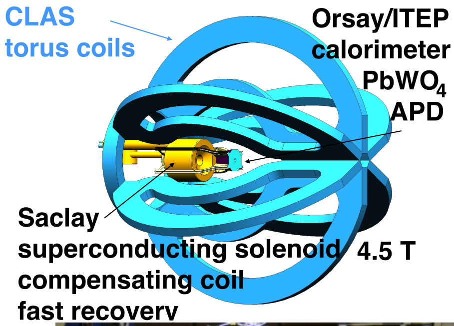
$$\alpha \propto \mathcal{I}\mathrm{m}\left(F_1\mathcal{H} + \xi G_M\tilde{\mathcal{H}} - \frac{t}{4M^2}F_2\mathcal{E}\right)$$

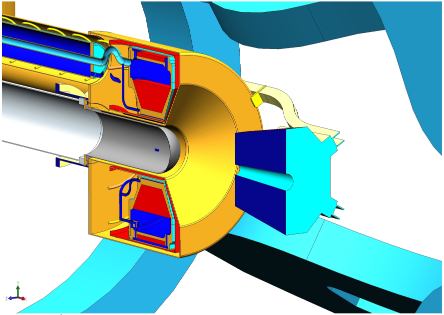
$$\mathcal{H}(\xi, t) = i\pi H(\xi, \xi, t) + \mathcal{P}\int_{-1}^{1} \mathrm{d}x \, \frac{H(x, \xi, t)}{x - \xi}$$

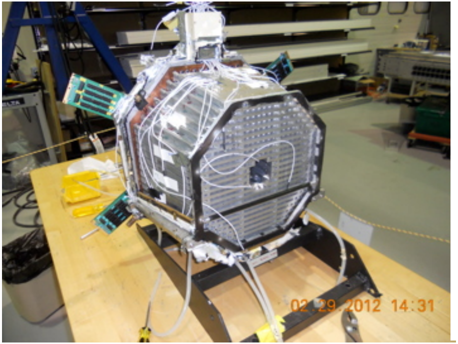
$$A_{\mathrm{UL}} \propto \mathcal{I}\mathrm{m}\left(F_1\tilde{\mathcal{H}} + \xi G_M\mathcal{H} + G_M\frac{\xi}{1+\xi}\mathcal{E} + \cdots\right)\sin\phi$$

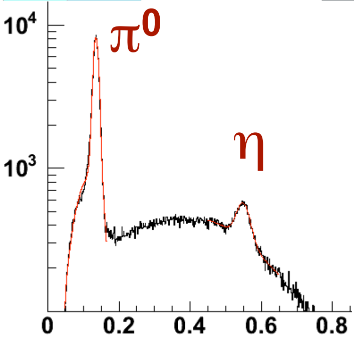
CLAS in Hall-B





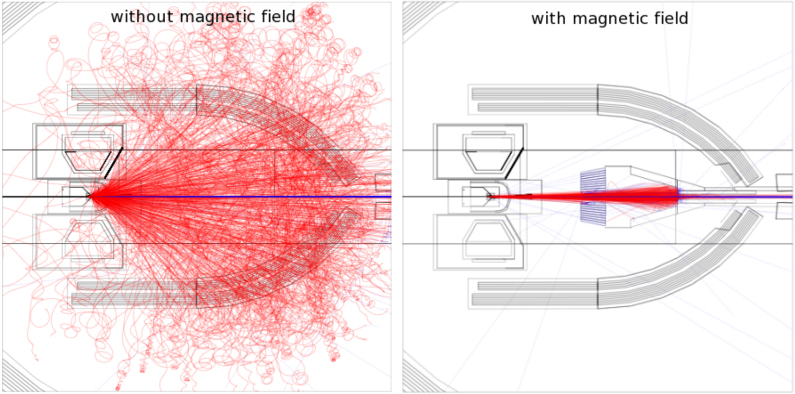




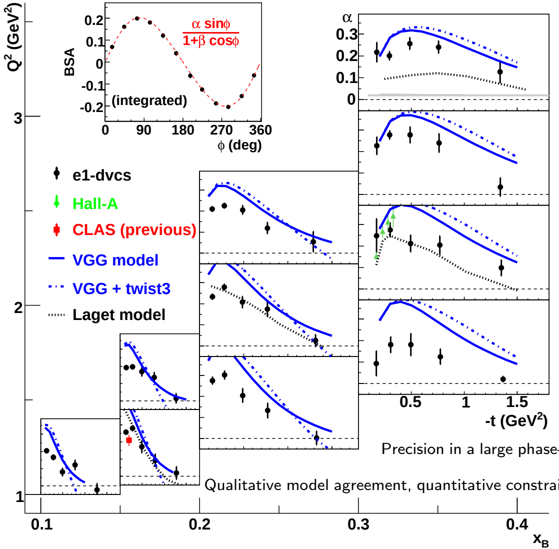


$$M_{\gamma\gamma} = \sqrt{2E_{\gamma_1}E_{\gamma_2}\left(1-\cos\theta_{\gamma\gamma}\right)}$$

Solenoid and Inner Calorimeter



CLAS proton Beam Spin Asymmetry

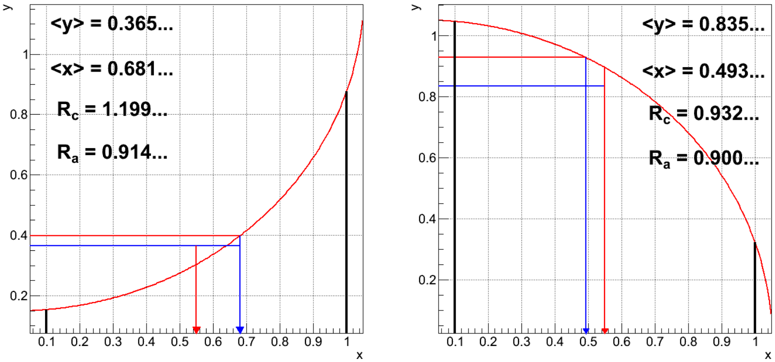


Precision in a large phase-space (x_B, Q^2, t)

Qualitative model agreement, quantitative constraints on parameters

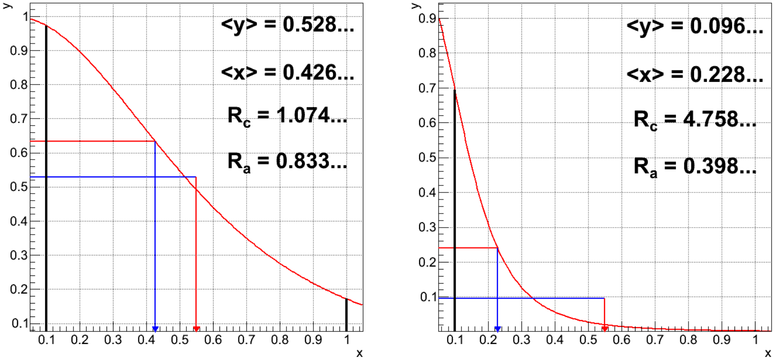
F.-X. G. *et al.*, PRL **100** (2008) 162002

Finite Bin Size Corrections



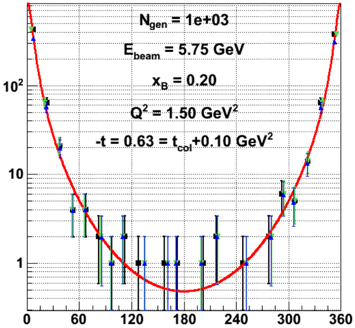
Bin volume	$V = \int_{\text{bin}} \mathrm{d}x$
Average value (observable) in bin	$< y >= y_a = \frac{1}{V} \int_{\text{bin}} y(x)\mathrm{d}x$
Bin center	$x_c = \frac{1}{V} \int_{\text{bin}} x\mathrm{d}x$
Bin average	$< x >= x_a = \frac{1}{V y_a} \int_{\text{bin}} x y(x)\mathrm{d}x$

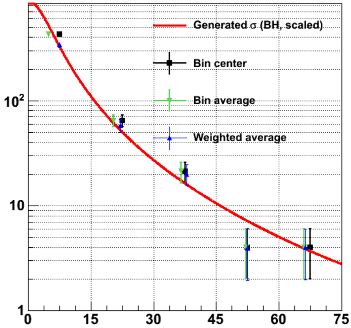
Finite Bin Size Corrections



Bin volume	$V = \int_{\text{bin}} \mathrm{d}x$
Average value (observable) in bin	$< y >= y_a = \frac{1}{V} \int_{\text{bin}} y(x)\mathrm{d}x$
Bin center	$x_c = \frac{1}{V} \int_{\text{bin}} x\mathrm{d}x$
Bin average	$< x >= x_a = \frac{1}{V y_a} \int_{\text{bin}} x y(x)\mathrm{d}x$

Finite Bin Size Corrections





Bin volume	$V = \int_{\text{bin}} \mathrm{d}x$
Average value (observable) in bin	$< y >= y_a = \frac{1}{V} \int_{\text{bin}} y(x)\mathrm{d}x$
Bin center	$x_c = \frac{1}{V} \int_{\text{bin}} x\mathrm{d}x$
Bin average	$< x >= x_a = \frac{1}{Vy_a} \int_{\text{bin}} xy(x)\mathrm{d}x$

CLAS proton cross-section

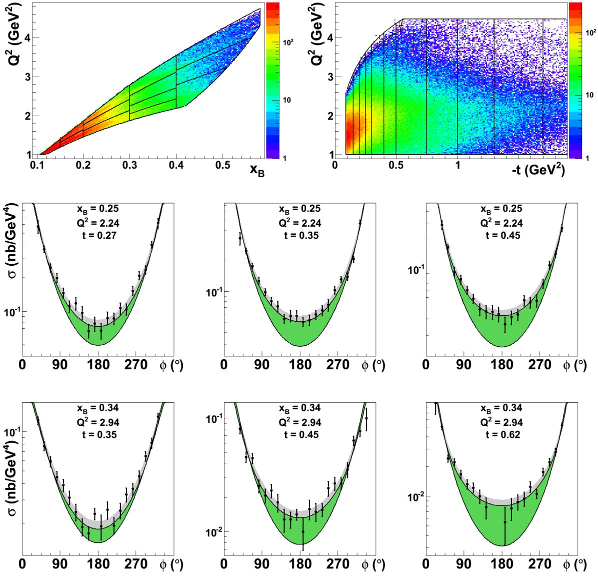
More than 3k bins

Dispersion relation :

$$\mathcal{R}e\ \mathcal{H} = \left[\int \mathcal{I}m\ \mathcal{H} \right] + \Delta$$

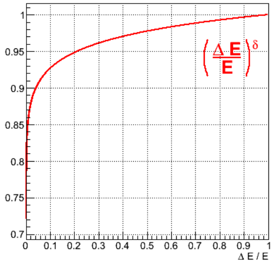
green band shows difference with BH

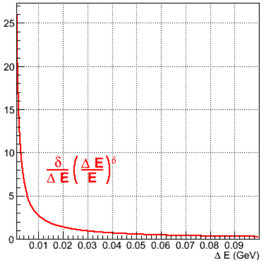
$\rightarrow$ sensitivity to d_1

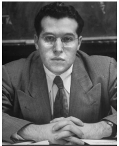


H.-S. Jo *et al.* PRL **115** (2015) 212003

Radiative corrections (semiclassical)





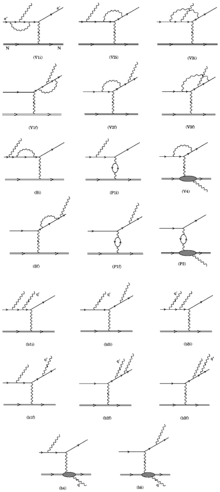


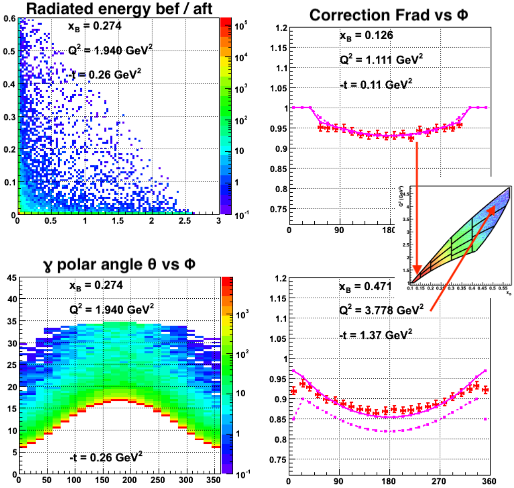
$$\left.\frac{d\sigma}{d\Omega}\right|_{\text{meas}} = (1+\delta_{\text{virt}}+\delta_{\text{real}}) \left.\frac{d\sigma}{d\Omega}\right|_{\text{Born}} \quad , \quad \delta_{\text{real}} = \ln\left(\frac{\Delta E}{E}\right)^{\delta} \quad , \quad \delta = \frac{2\alpha}{\pi}\left[\ln\frac{Q^2}{m^2} - 1\right]$$

Schwinger resummation : $1+\delta_{\text{real}} \to e^{\delta_{\text{real}}}$, $\left.\frac{d\sigma}{d\Omega}\right|_{\text{meas}} = \left.\frac{d\sigma}{d\Omega}\right|_{\text{Born}} (1+\delta_{\text{virt}}) \left(\frac{\Delta E}{E}\right)^{\delta}$

Peskin & Schroeder chapter 6 (Introduction)

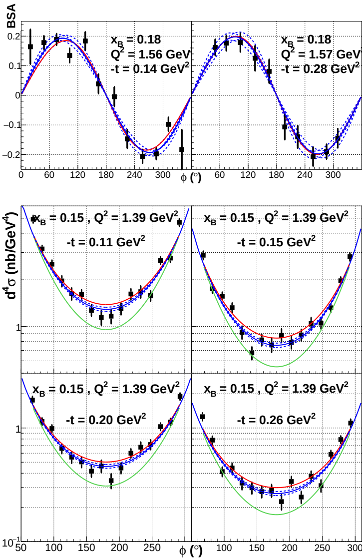
DVCS Radiative Corrections Fast MC





M. Vanderhaeghen *et al.*, Phys. Rev **C62** (2000) 025510
I. Akushevich & A. Ilyichev Phys. Rev **D98** (2018) 013005

Global Fits to extract the D-term



Beam Spin Asymmetries

$$\mathrm{Im}\mathcal{H}(\xi,t) = \frac{N}{1+x}\left(\frac{2\xi}{1+\xi}\right)^{-\alpha_R(t)}\left(\frac{1-\xi}{1+\xi}\right)^{b}\left(\frac{1-\xi}{1+\xi}\,\frac{t}{M^2}\right)^{-1}$$

Unpolarized cross-sections

Use dispersion relation:

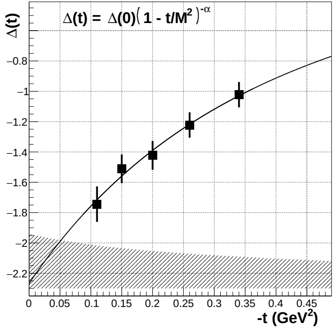
$$\mathrm{Re}\mathcal{H}(\xi,t) = \Delta + \mathcal{P}\int dx\left(\frac{1}{\xi-x}-\frac{1}{\xi+x}\right)\mathrm{Im}\mathcal{H}(\xi,t)$$

pure Bethe-Heitler

local fit + uncertainty range

resulting global fit

DVCS Dispersion: subtraction constant results



$$\mathrm{Im}\mathcal{H}(\xi,t) = \frac{N}{1+x}\left(\frac{2\xi}{1+\xi}\right)^{-\alpha_R(t)}\left(\frac{1-\xi}{1+\xi}\right)^{b}\left(\frac{1-\xi}{1+\xi}\,\frac{t}{M^2}\right)^{-1}$$

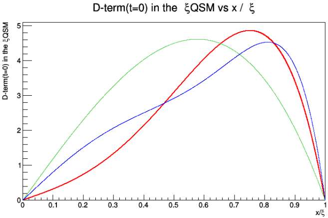
α_R is fixed from small-x Regge phenomenology
b is a free parameter for the large x behavior
p is fixed to 1 for the valence
M is a free parameter controlling the t dependence

$$\Delta(t) = \Delta(0)\left(1 - \frac{t}{M^2}\right)^{-\alpha} = 2\int_{-1}^{1} dz\,\frac{D(z,t)}{1-z}$$

$$D(z,t) = (1-z^2)\sum_{k=1}^{\infty}\left[e_u^2\ d_{2k-1}^u(t) + e_d^2\ d_{2k-1}^d(t)\right]\ C_{2k-1}^{3/2}(z)$$

Hereafter assume $d_{2k-1}^u \approx d_{2k-1}^d$

Separation of the GFF d_1



In the χQSM:
$d_1(0) \approx -4.0\,;\ d_3(0) \approx -1.2\,;\ d_5(0) \approx -0.4$

$$H(x,\xi,t) = \cdots + \theta\left[1-\frac{x^2}{\xi^2}\right] D(\frac{x}{\xi},t)$$

$$D(z,t) = \left(1-z^2\right)\left[d_1(t)C_1^{3/2}(z) + d_3(t)C_3^{3/2}(z) + \cdots\right]$$

$$C_1^{3/2}(z) = 3z$$

$$C_3^{3/2}(z) = \frac{5}{2}\left(7z^3 - 3z\right)$$

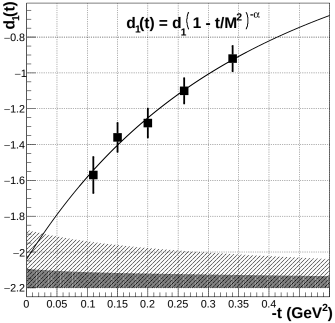
$$C_5^{3/2}(z) = \frac{21}{8}\left(33z^5 - 30z^3 + 5z\right)$$

To separate orthogonal Gegenbauer polynomials:
requires measurement of (x,ξ) dependence (or at least $z = x/\xi$ dependence)
different reaction such as DDVCS, will require higher luminosity

For now, to make progress it is necessary to make assumptions
Use guidance from models such as χQSM or lattice results
Also implement constraints from theory into phenomenological fits

Separation of the GFF d_1

$$D^q(\tfrac{x}{\xi}, t) = \left(1 - \tfrac{x^2}{\xi^2}\right)\left[d_1^q(t)C_1^{3/2}(\tfrac{x}{\xi}) + d_3^q(t)C_3^{3/2}(\tfrac{x}{\xi}) + \cdots\right]$$



t-dependence of the D-term :

Dipole gives singular pressure at $r = 0$

Power law implied by counting rules?

Exponential?

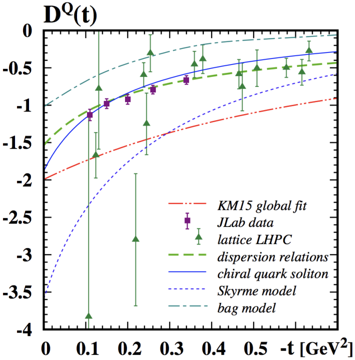
...

$d_1(0) < 0$ dynamical stability of bound state

$d_1(0) = -2.04 \pm 0.14 \pm 0.33$

First Measurement of new fundamental quantity

D-term comparison with theory



Dispersion Relation Analysis

Chiral quark soliton model

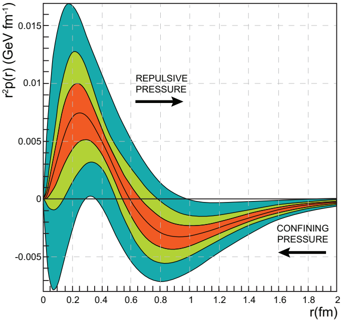
Lattice results LHPC

Global fit

M. V. Polyakov, P. Schweitzer Int.J.Mod.Phys. **A33** (2018)

em: $\partial_\mu J^\mu_{\rm em} = 0$	$\langle N'\|J^\mu_{\rm em}\|N\rangle$	$\longrightarrow$	$Q = 1.602176487(40) \times 10^{-19}$C	
			$\mu = 2.792847356(23)\mu_N$	
weak: PCAC	$\langle N'\|J^\mu_{\rm weak}\|N\rangle$	$\longrightarrow$	$g_A = 1.2694(28)$	
			$g_p = 8.06(55)$	
gravity: $\partial_\mu T^{\mu\nu}_{\rm grav} = 0$	$\langle N'\|T^{\mu\nu}_{\rm grav}\|N\rangle$	$\longrightarrow$	$m = 938.272013(23)\,{\rm MeV}/c^2$	
			$J = \frac{1}{2}$	
			$\boldsymbol{D} = \boldsymbol{?}$	

Proton Pressure distribution results



The pressure at the core of the proton is $\sim 10^{35}$ Pa
About 10 times the pressure at the core of a neutron star

Positive pressure in the core (repulsive force)
Negative pressure at the periphery: pion cloud
Pressure node around $r \approx 0.6$ fm

Stability condition : $\int_0^\infty \mathrm{d}t\, r^2 p(r) = 0$

Rooted into Chiral Symmetry Breaking

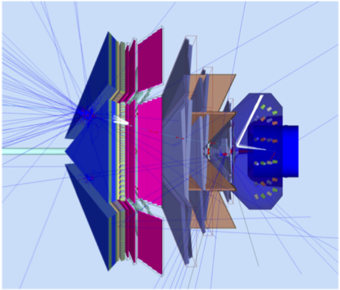
World data fit

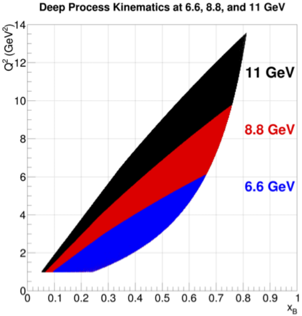
CLAS 6 GeV data

Projected CLAS12 data E12-16-010B

DVCS with CLAS12 E12-16-010B

Experimental Configuration

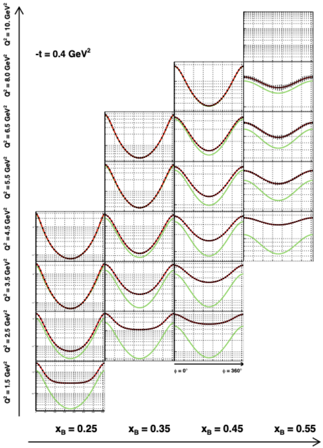


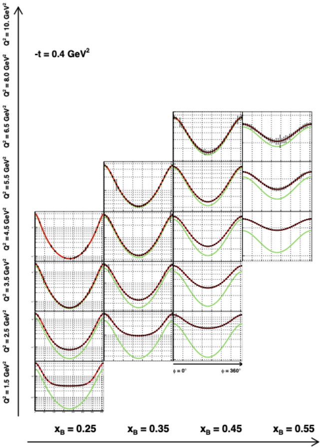


- $\mathcal{L} = 1 \times 10^{35}$ cm^{-2}s^{-1}
- Inclusive electron trigger (all calibration reactions will be analyzed in parallel)
- Electrons in the forward detector
- Protons in the central detector and forward detector
- Photons in the forward detector and forward tagger

DVCS with CLAS12 E12-16-010B

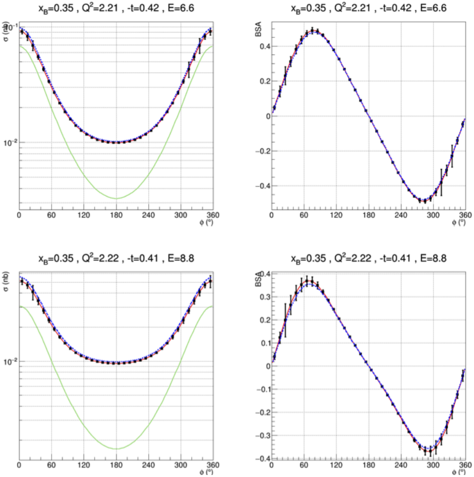
DVCS Projected Cross Section at 11GeV & 8.8GeV





DVCS with CLAS12 E12-16-010B

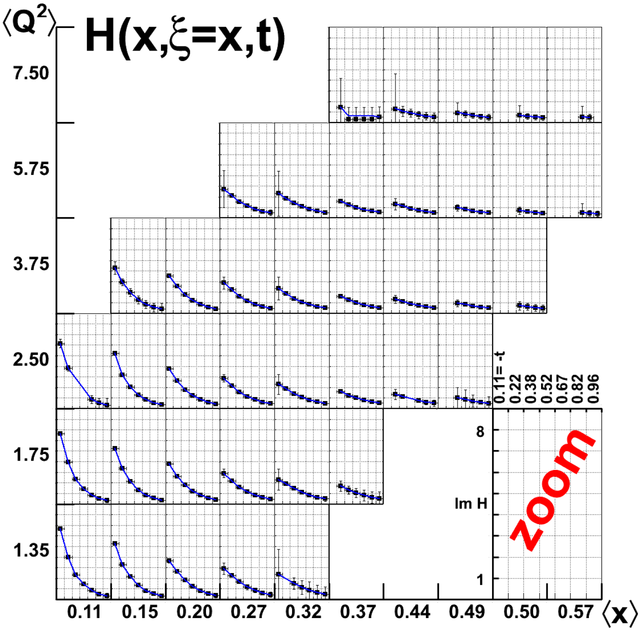
DVCS Projected Results



- Left : cross-sections at fixed kinematics for beam energies 6.6 GeV and 8.8 GeV
- Right : corresponding Beam Spin Asymmetries
- Green : pure Bethe-Heitler
- Red : model fit on 6 GeV data
- Blue : simultaneous fit of the projected data
- Separation of : $\mathcal{I} \sim 1/y^3$ and $\left|\mathcal{T}^{\mathrm{DVCS}}\right|^2 \sim 1/y^2$

Projected impact on GPD extraction methods

Using simulated data based on VGG model. Input GPD H extracted with good accuracy



Projected impact on GPD extraction methods

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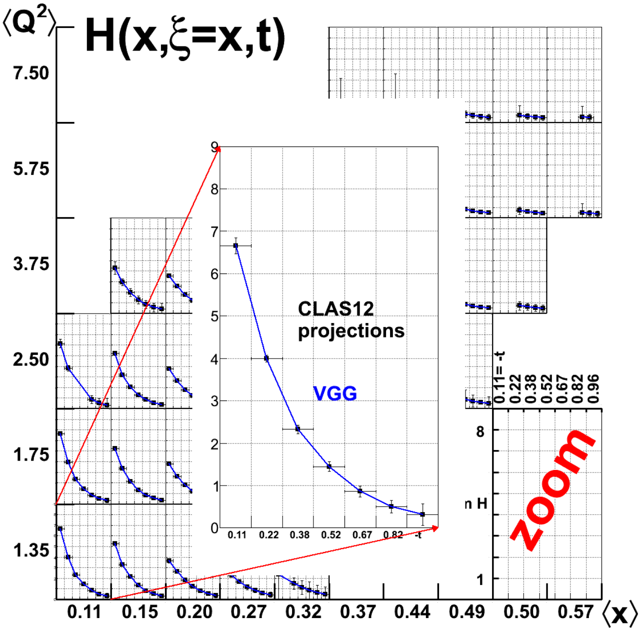

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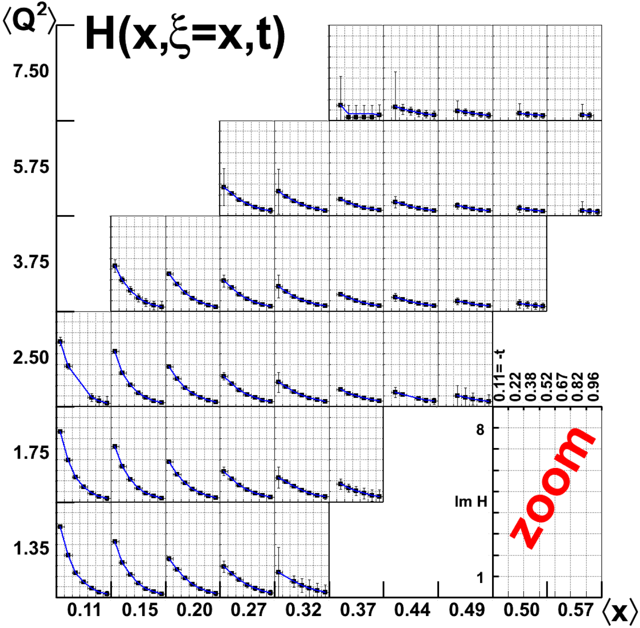

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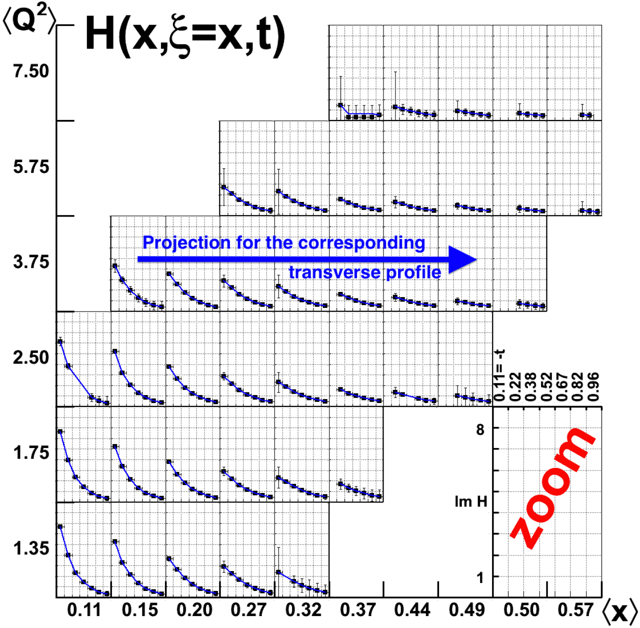
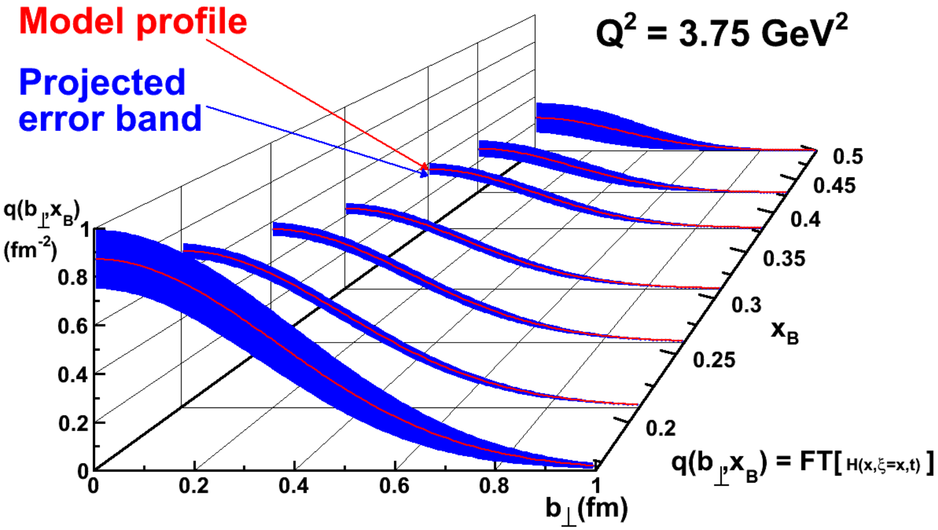
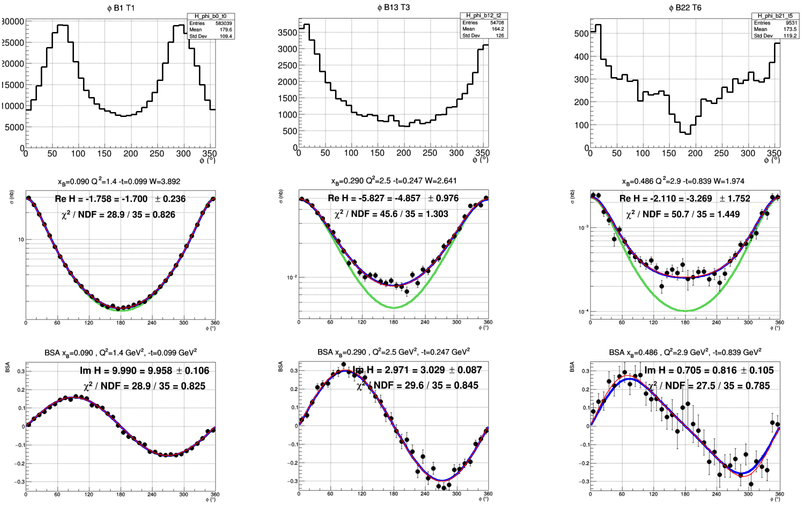


Projection for the Nucleon transverse profile

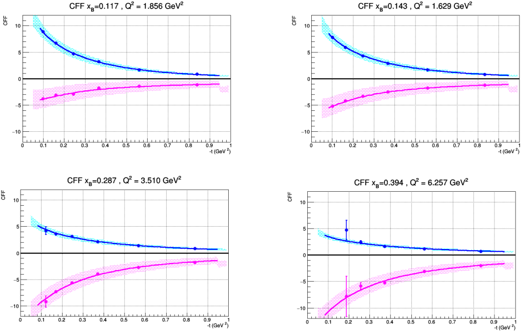


Precision tomography in the valence region

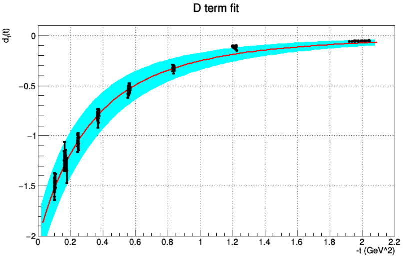
Femtography simulations E12-16-010

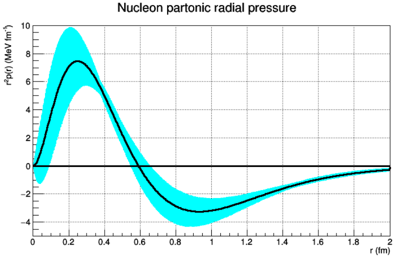


Femtography simulations E12-16-010



High statistics CLAS12 simulations:
sample of extracted Im and Re parts of CFF vs -t





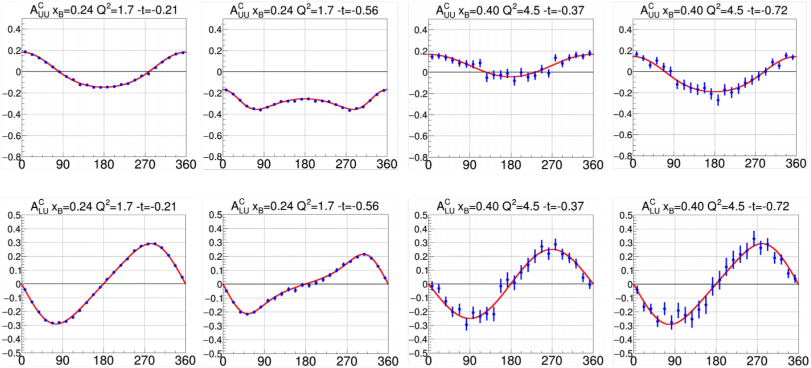
Extraction of the D-term, and partonic pressure
from amplitude dispersion global analysis

DVCS with a Polarized Positron beam

PEPPo production injecting 60 MeV **100 nA positron polarized at 60%**

(PEPPo Collaboration) D. Abbott *et al.* , **PRL**116 (2016) 214801 ; L. Cardman *et al.* AIP CP 1970 (2018) 050001

Proposal 100 days (80+20) at $\mathcal{L} = 0.6 \times 10^{35}$ cm^{-2}s^{-1}



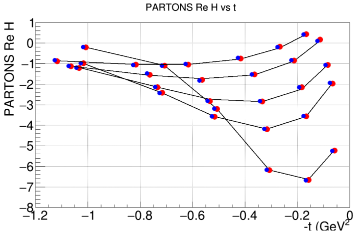
Impact of the CLAS12 Positron data

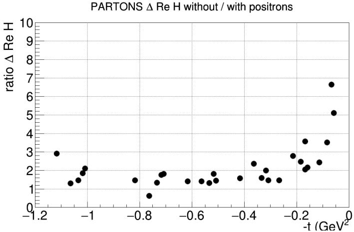
Global analysis of CLAS12 program observables $\{\sigma_{UU}, A_{LU}, A_{UL}, A_{LL}, A^C_{UU}, A^C_{LU}\}$

unpolarized beam charge asymmetry $\mathbf{A^C_{UU}}$ sensitive to the amplitude **real part**

polarized beam charge asymmetry $\mathbf{A^C_{UU}}$ sensitive to the amplitude **imaginary part**

Fitting $\{\mathcal{H}, \tilde{\mathcal{H}}\}$ assuming model values for $\{\mathcal{E}, \tilde{\mathcal{E}}\}$



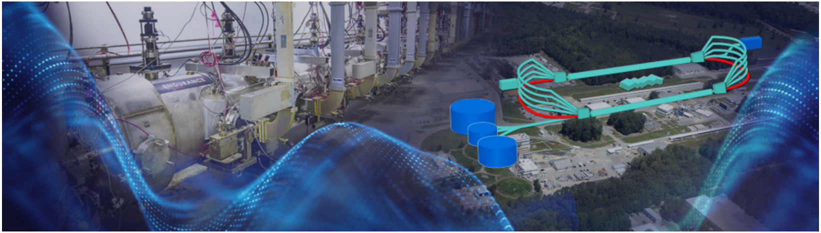


Improvement of the **statistical** and **systematical** uncertainties

Model independent separation of the Interference with BH and DVCS^2

Science at the Luminosity Frontier: JLab at 22 GeV

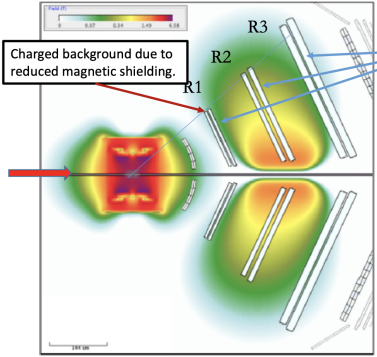
First workshop January 23, 2023 to January 25, 2023



Second workshop

CLAS12 at higher luminosities and energies

- CLAS12 luminosity limited by accidental occupancy of DC R1.



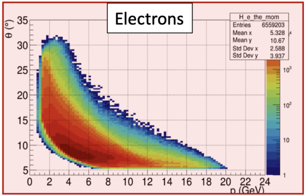
High occupancy in part of R1 limits acceptable operating luminosity.

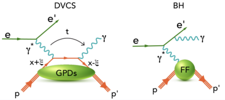
	R1	R2	R3
CLAS12 @ 11 GeV	**2.6%**	**0.76%**	**1.18%**
CLAS12 @ 24 GeV	**2.8%**	**0.77%**	**1.23%**

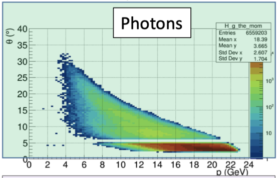
Accidental occupancies increase by less than 10% at 24 GeV compared to 10.6 GeV.
➔ **higher resolution tracking layers to double luminosity.**

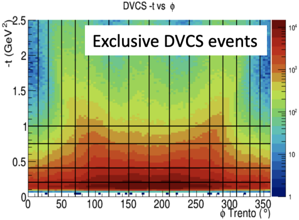
Z. Meador, L. Elouadrhiri

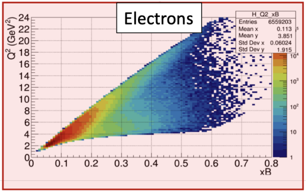
DVCS with CLAS at 24 GeV

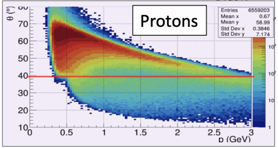


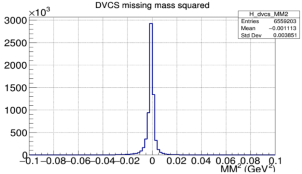




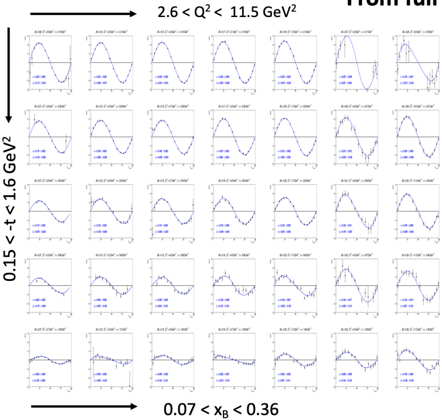






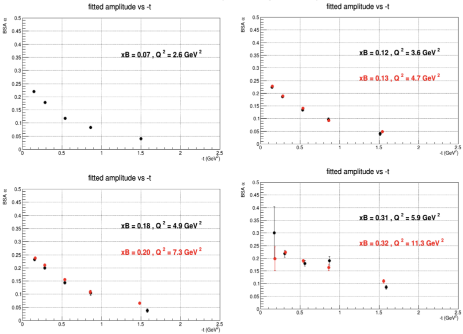


DVCS BSA with CLAS at 24 GeV



From full simulations & reconstructions

$$A_{LU} = \alpha \sin\phi / (1 + \beta \cos 2\phi)$$



Summary - Outlook

- Dispersive analysis of the DVCS amplitude
- Access to the Gravitational Form Factors
- Global analysis with as few parameters as possible
- Questions about model dependance and systematic uncertainties
- Rosenbluth separation with several beam energies
- Direct access to the real part of the amplitude with a positron beam
- Future prospects: precision at the luminosity and energy frontiers

Next lecture: interplay between JLab and the future Electron Ion Collider