

Deeply Virtual Compton Scattering at JLab 12GeV

First Studies of the Fall 2014 Run at Hall A

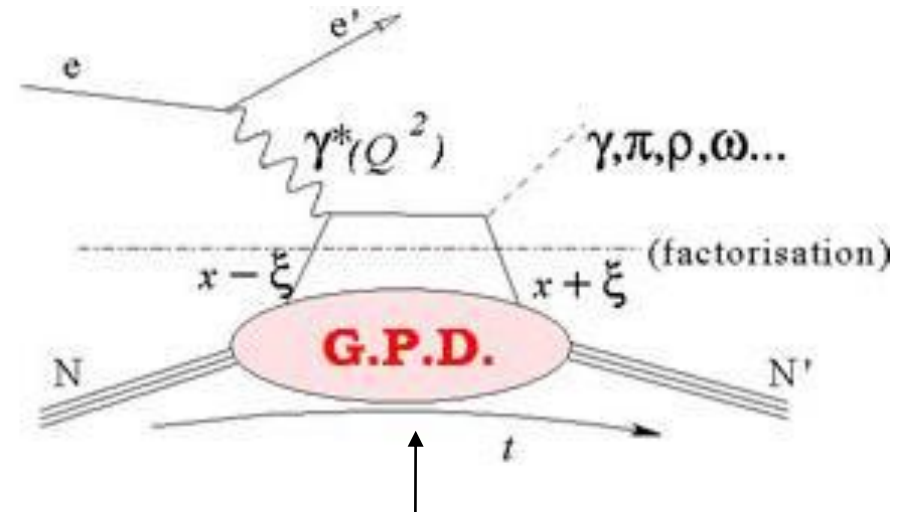
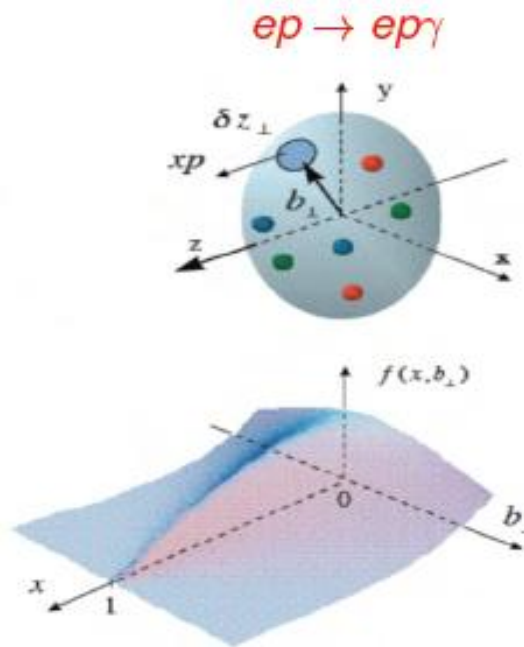
Marco Carmignotto

JLab Hall A DVCS Collaboration

Topics

- Generalized Parton Distributions and Deeply Virtual Compton Scattering
- DVCS at Jefferson Laboratory, Hall A
- Instrumentation:
 - Hall A spectrometers
 - DVCS: dedicated calorimeter FPGA based data acquisition
- Preliminary studies:
 - Data acquisition deadtimes monitoring
 - Detector efficiencies
 - Preliminary missing mass checks: DVCS and π^0

Generalized Parton Distributions



GPDs describe the soft part, i.e., with non-perturbative functions.

Generalized Parton Distributions (GPD) encode information on the distribution of partons both in the transverse plane and in the longitudinal direction

Nucleon structure described by 4 GPDs:
 H , E (unpolarized), \tilde{H} , \tilde{E} , polarized)

Quark GPDs E and H are connected to the elastic form factors:

$$\int_0^1 dx H^q(x, \xi, t, \mu^2) = F_1^q(t)$$

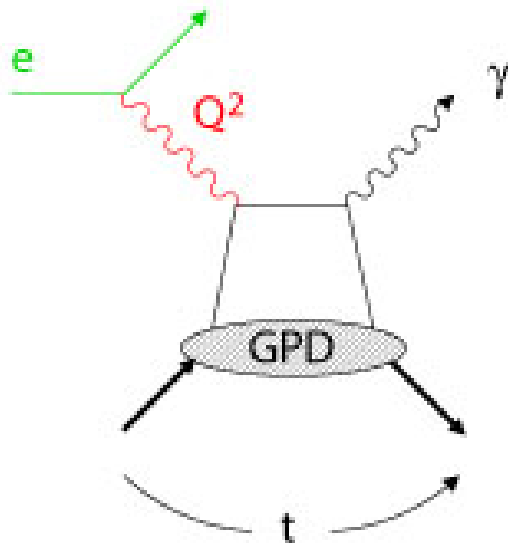
$$\int_0^1 dx E^q(x, \xi, t, \mu^2) = F_2^q(t)$$

And also to usual PDFs.

Accessing GPDs experimentally

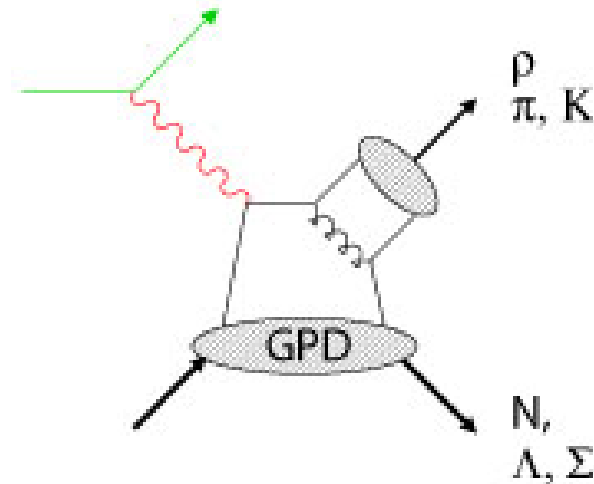
We can access GPDs through hard exclusive reactions

Deeply Virtual Compton Scattering (DVCS)



- Simplest and cleanest way to access GPDs experimentally. The spin asymmetry studies in DVCS give access to imaginary part of the scattering amplitude.
- DVCS (flavor blind) probes GPD H and provides additional information on singlet quarks

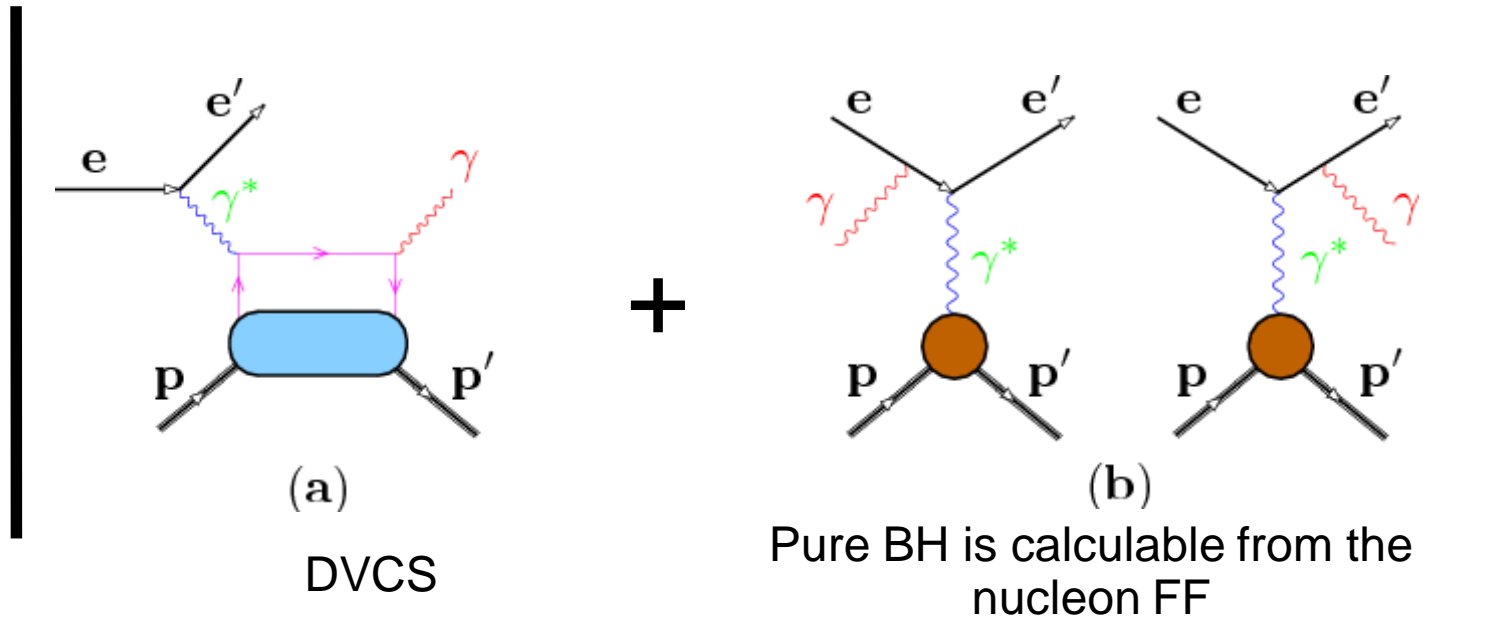
Deeply virtual meson production (DVMP)



- Mesons select definite charge, spin, flavor component of GPD
- Quantum numbers in DVMP probe individual GPD components selectively
- Need good understanding of reaction mechanism
 - QCD factorization for mesons is complex (additional interaction of the produced meson)

Accessing GPDs with DVCS

DVCS amplitudes interfere with E&M Bethe-Heitler:



2

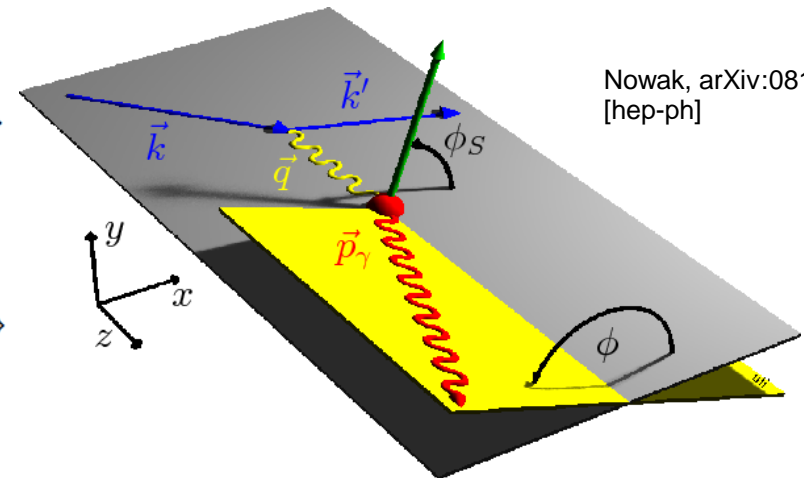
Interference term:

$$\mathcal{I} = \frac{e^6}{x_{BY}^3 \mathcal{P}_1(\phi_{\gamma\gamma}) \mathcal{P}_2(\phi_{\gamma\gamma}) t} \left\{ c_0^{\mathcal{I}} + \sum_{n=1}^3 (-1)^n \left[c_n^{\mathcal{I}}(\lambda) \cos(n\phi_{\gamma\gamma}) - \lambda s_n^{\mathcal{I}} \sin(n\phi_{\gamma\gamma}) \right] \right\}$$

DVCS² term:

$$|\mathcal{T}^{DVCS}(\lambda)|^2 = \frac{e^6}{y^2 Q^2} \left\{ c_0^{DVCS} + \sum_{n=1}^2 (-1)^n \left[c_n^{DVCS} \cos(n\phi_{\gamma\gamma}) + \lambda s_1^{DVCS} \sin(n\phi_{\gamma\gamma}) \right] \right\}$$

Encodes GPDs



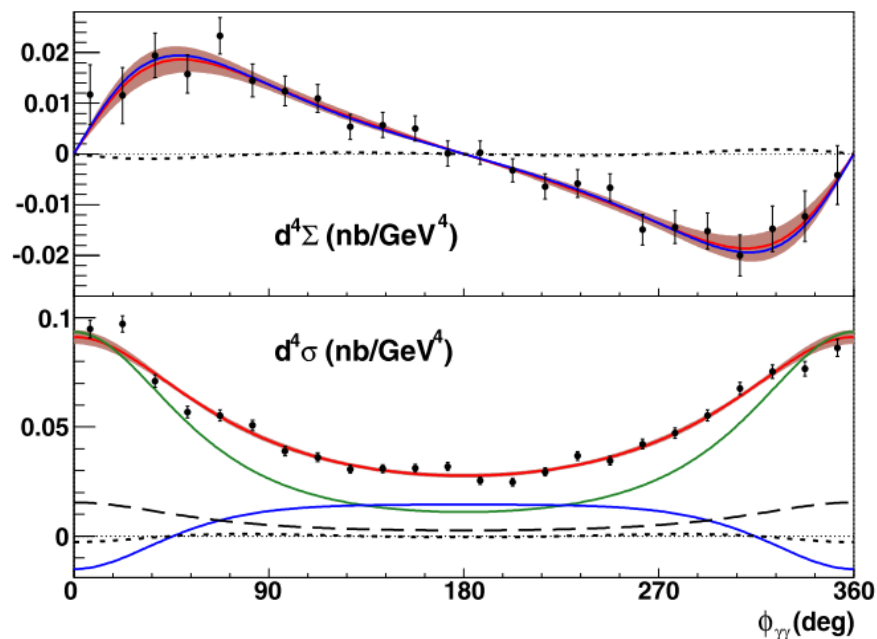
Nowak, arXiv:0812.2679 [hep-ph]

5

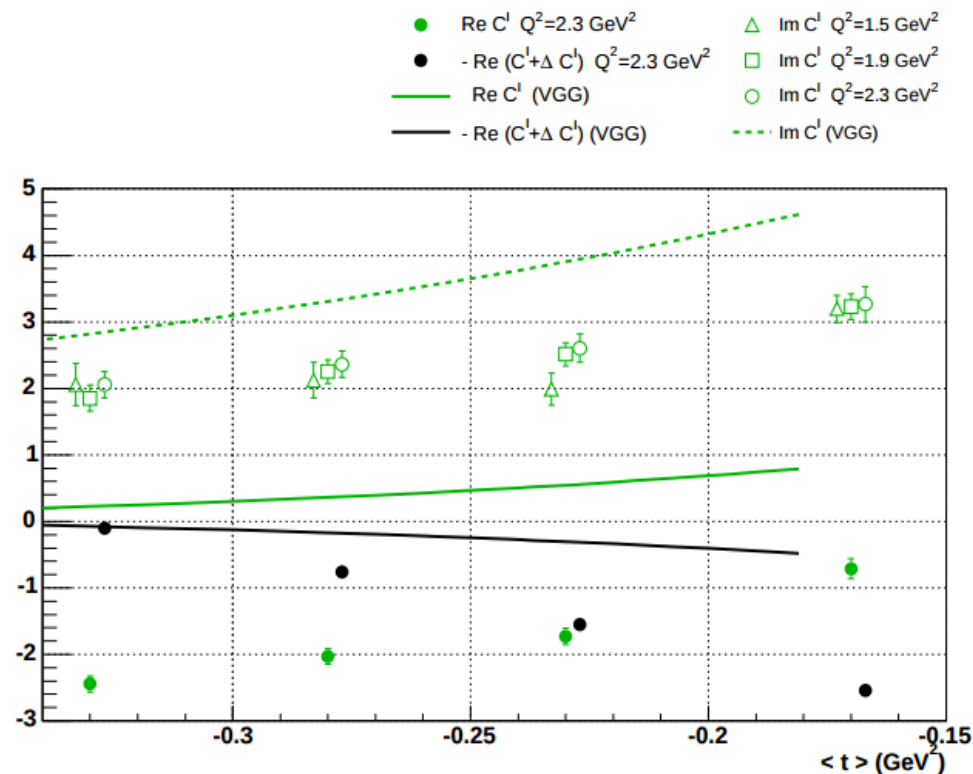
DVCS at JLab Hall A - 6 GeV

6 GeV

DVCS helicity-dependent ($d^4\Sigma$) and helicity-independent ($d^4\sigma$) cross sections measured in E00-110 for $Q^2 = 2.3 \text{ GeV}^2$ and $t = -0.28 \text{ GeV}^2$.

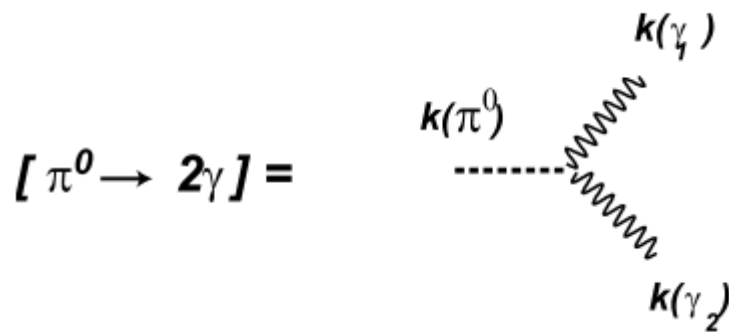
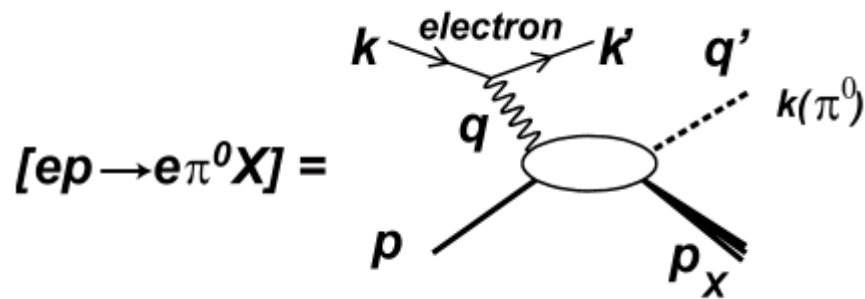


Some Fourier coefficients measured



π^0 electroproduction

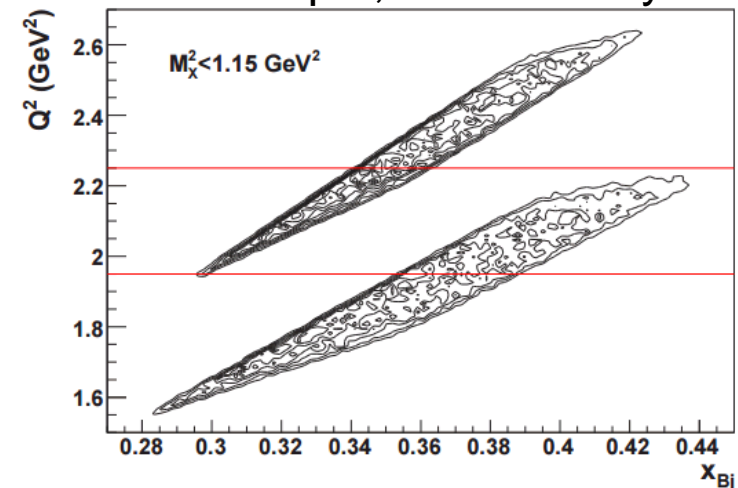
Measurements of the same final state particles: e^- and photon



Cross section measured at same Q^2 , but different epsilon, allow for L/T separation.

$$\left. \frac{d\sigma_v}{d\Omega_\pi} \right|_{unpol} = \frac{d\sigma_T}{d\Omega_\pi} + \epsilon_L \frac{d\sigma_L}{d\Omega_\pi} + \sqrt{2\epsilon_L(1+\epsilon)} \frac{d\sigma_{TL}}{d\Omega_\pi} \cos \phi_\pi + \epsilon \frac{d\sigma_{TT}}{d\Omega_\pi} \cos 2\phi_\pi$$

For example, 6 GeV analysis:

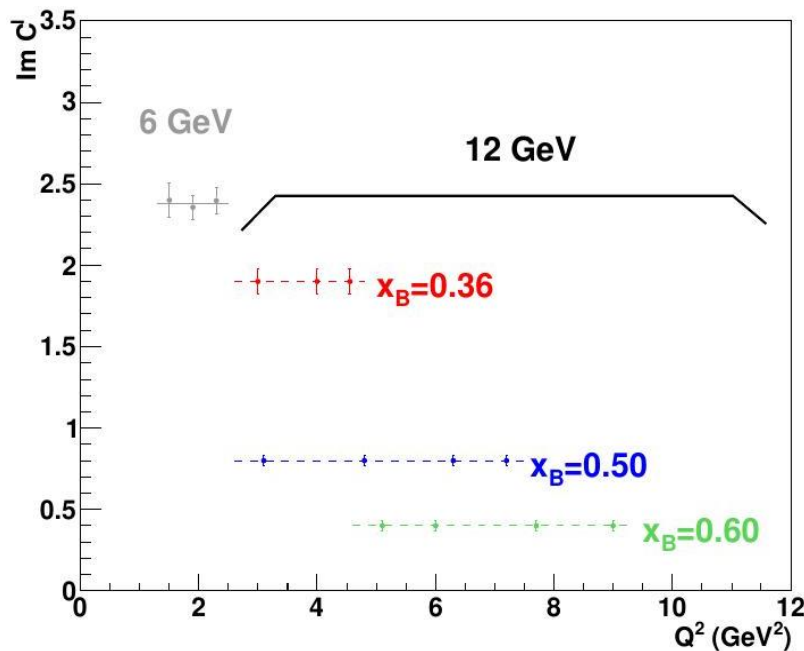


π^0 events are in the acceptance:
 We will measure the π^0 electroproduction cross section
 in the same kinematics, that may allow for L/T separation

DVCS kinematics: JLab @ 12 GeV

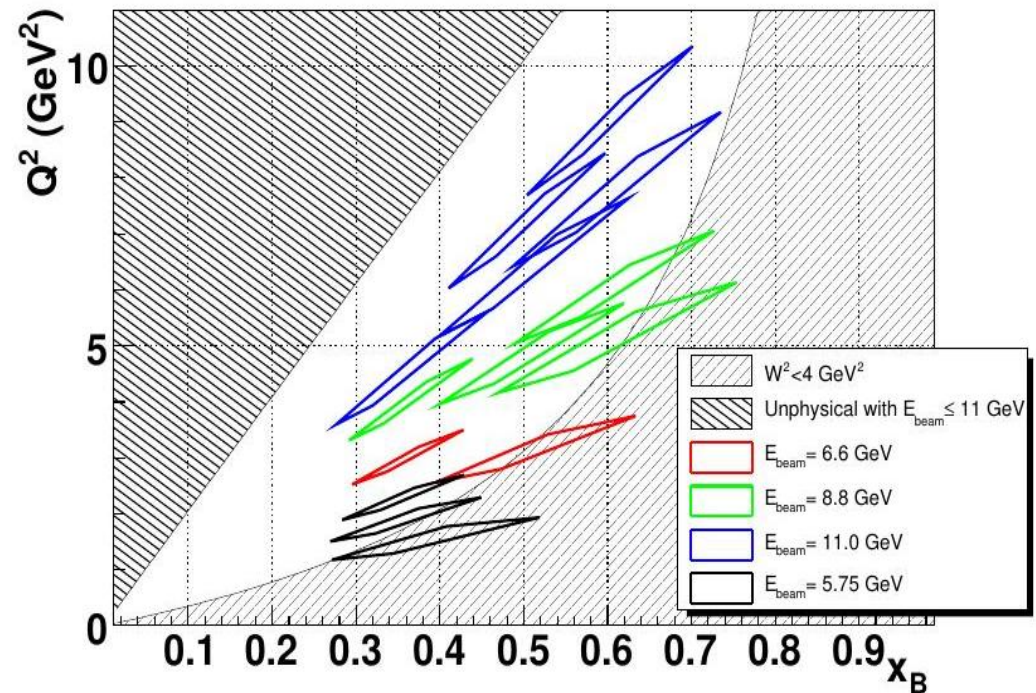
- Scaling test of DVCS cross-sections to 5% precision over large arm in Q^2
- Separation of Re and Im part of DVCS amplitude (polarized and total cross section)

3rd generation of DVCS experiments at JLab Hall A



We have a hint for leading order domination. We will expand the Q^2 and x_{Bj} scan to nail that down.

Planned scans in Q^2 and x_{Bj}



JLab at 12 GeV and DVCS instrumentation

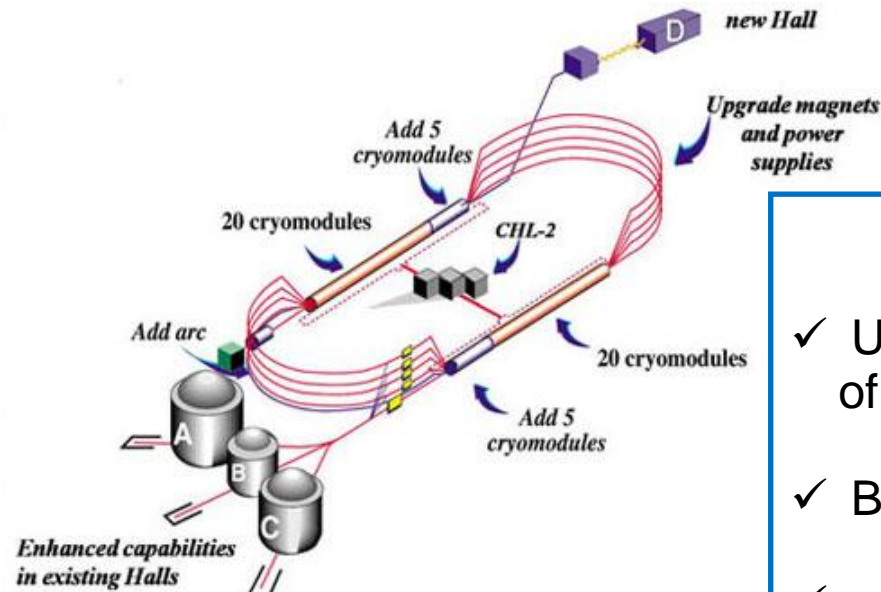
Jefferson Laboratory at 12 GeV



JLab during the upgrade
Commissioning phase now!



Added new
cryomodules to
double the electron
beam energy

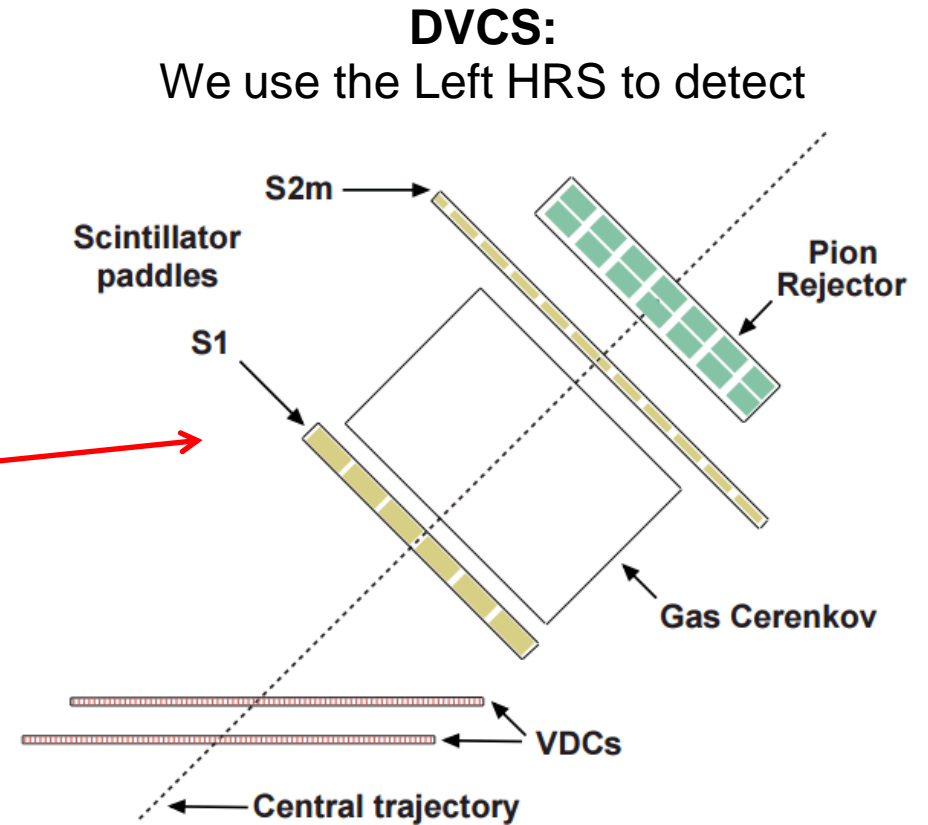
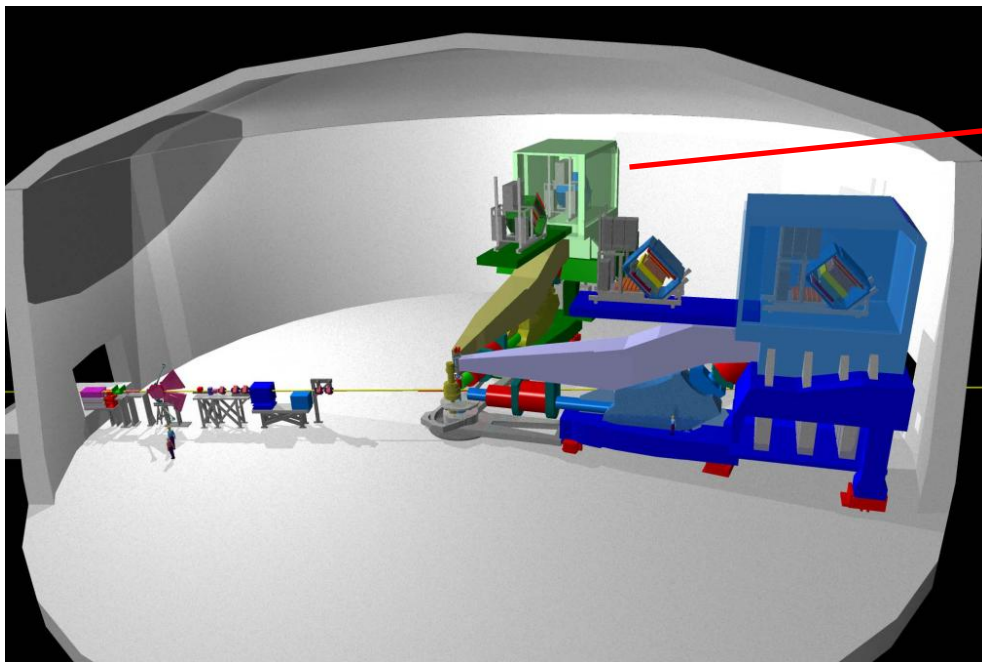


Upgrades:

- ✓ Upgraded equipment of existing Halls
- ✓ Built Hall D (GlueX)
- ✓ ...
- Commissioning phase

JLab Hall A

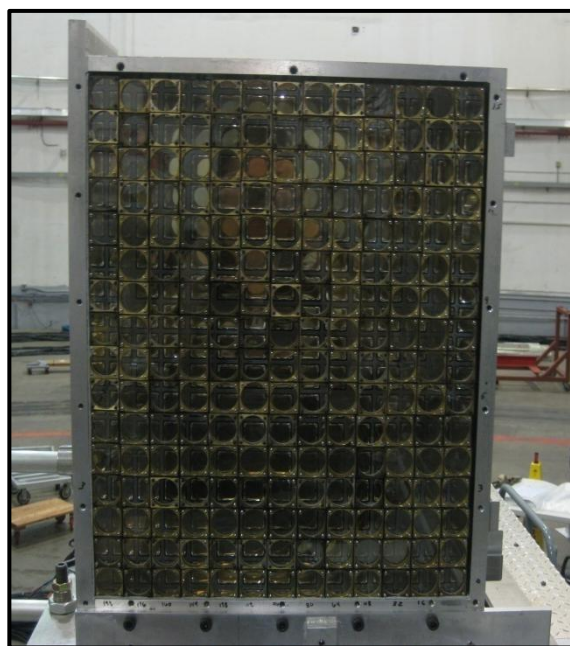
Existing two High Resolution Spectrometers (HRS)



- s2m and s0: scintillators pannels for triggering
- PID: Gas Cherenkov, Pion Rejector calorimeter
- Drift chamber, FPP chamber

Dedicated detector for DVCS

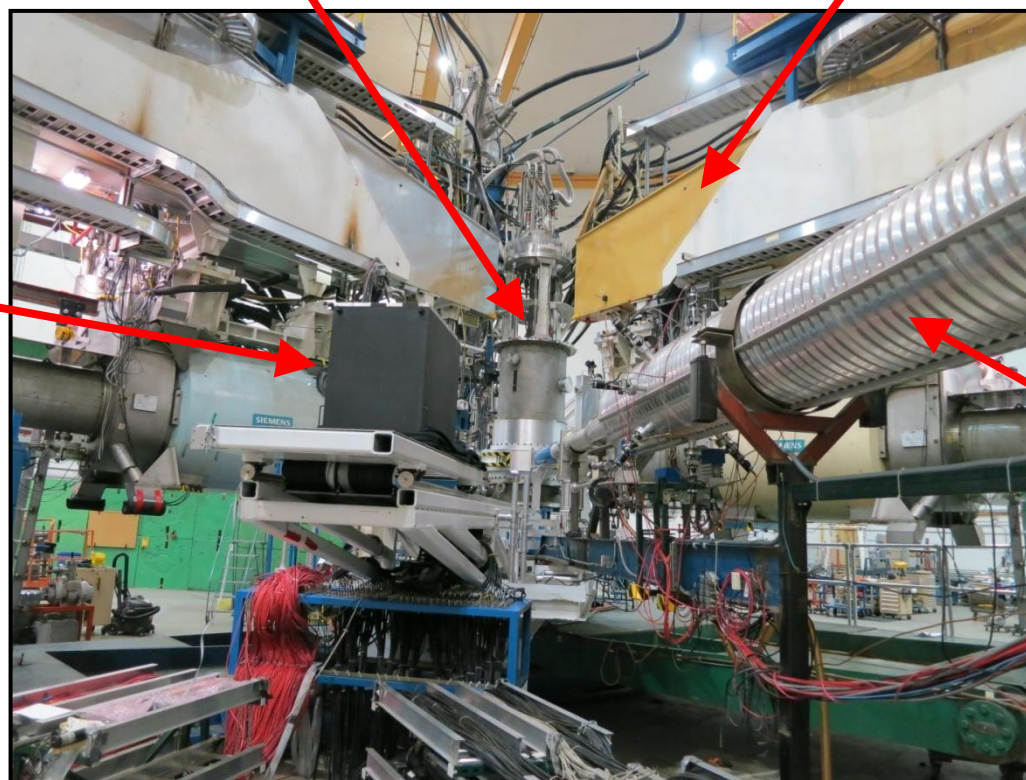
DVCS calorimeter



208 PbF_2 blocks calorimeter
Resolution: 3% ~ modest

LH2 Target

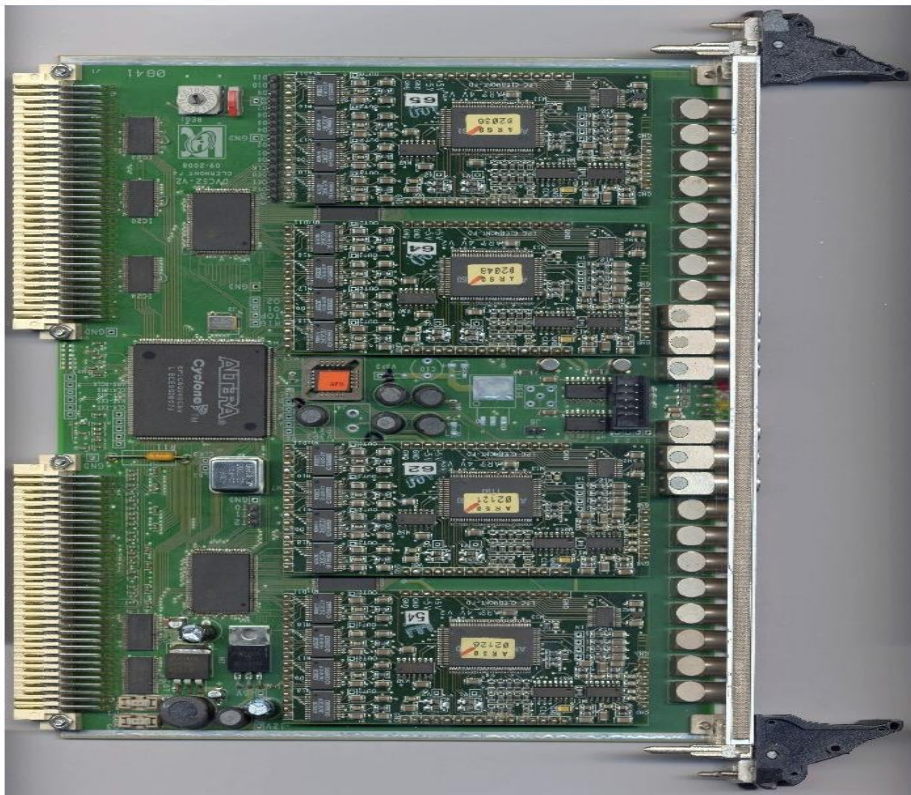
Hall A Left HRS



Beam dump

Calorimeter photon energy resolution is our limiting factor in the missing mass reconstruction

Dedicated FPGA electronics for trigger and signal sampling

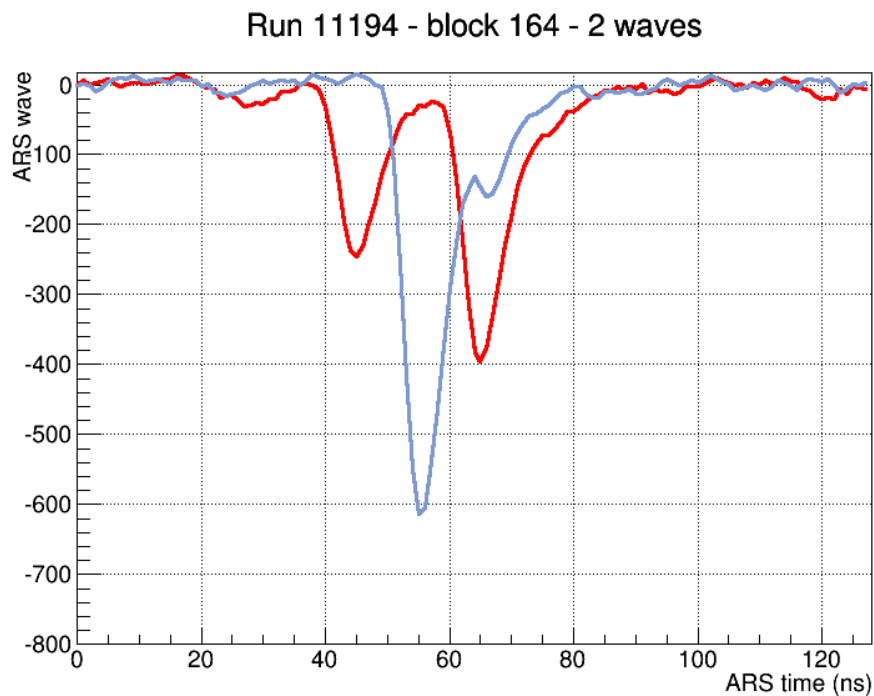


Analog Ring Sampler (ARS)

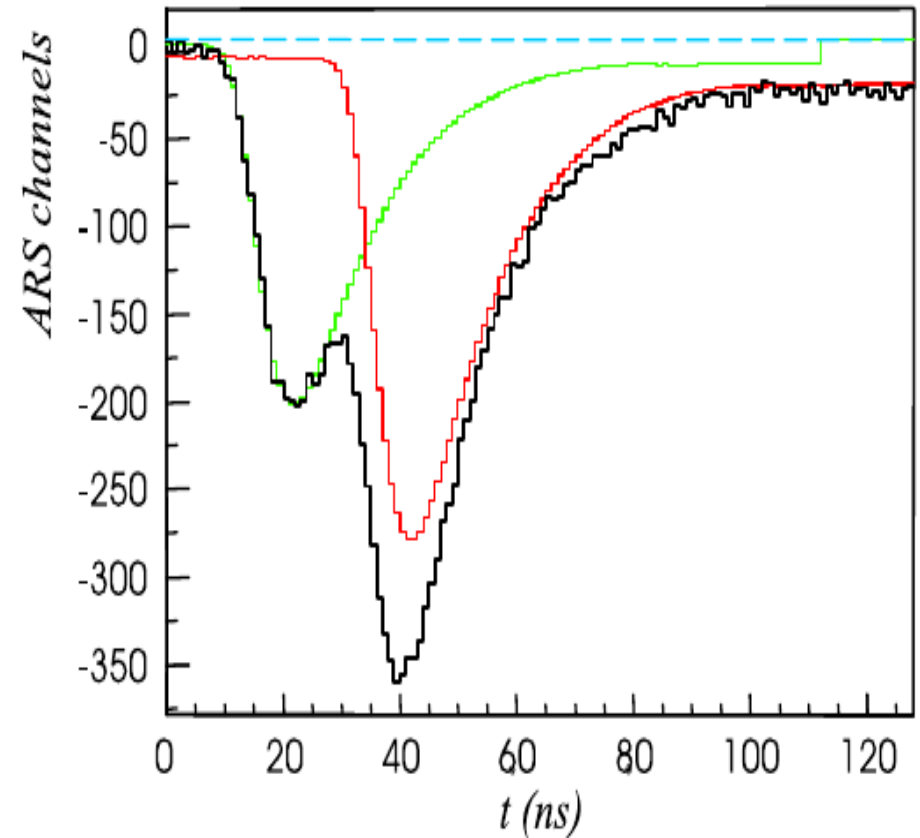
- 208 channels
- 1 GHz sampling rate
- Digitalization of 128 ns window

**Allows for waveform analyses
and better separate pile-up
events off-line**

Sampling and disentangling pile-up events



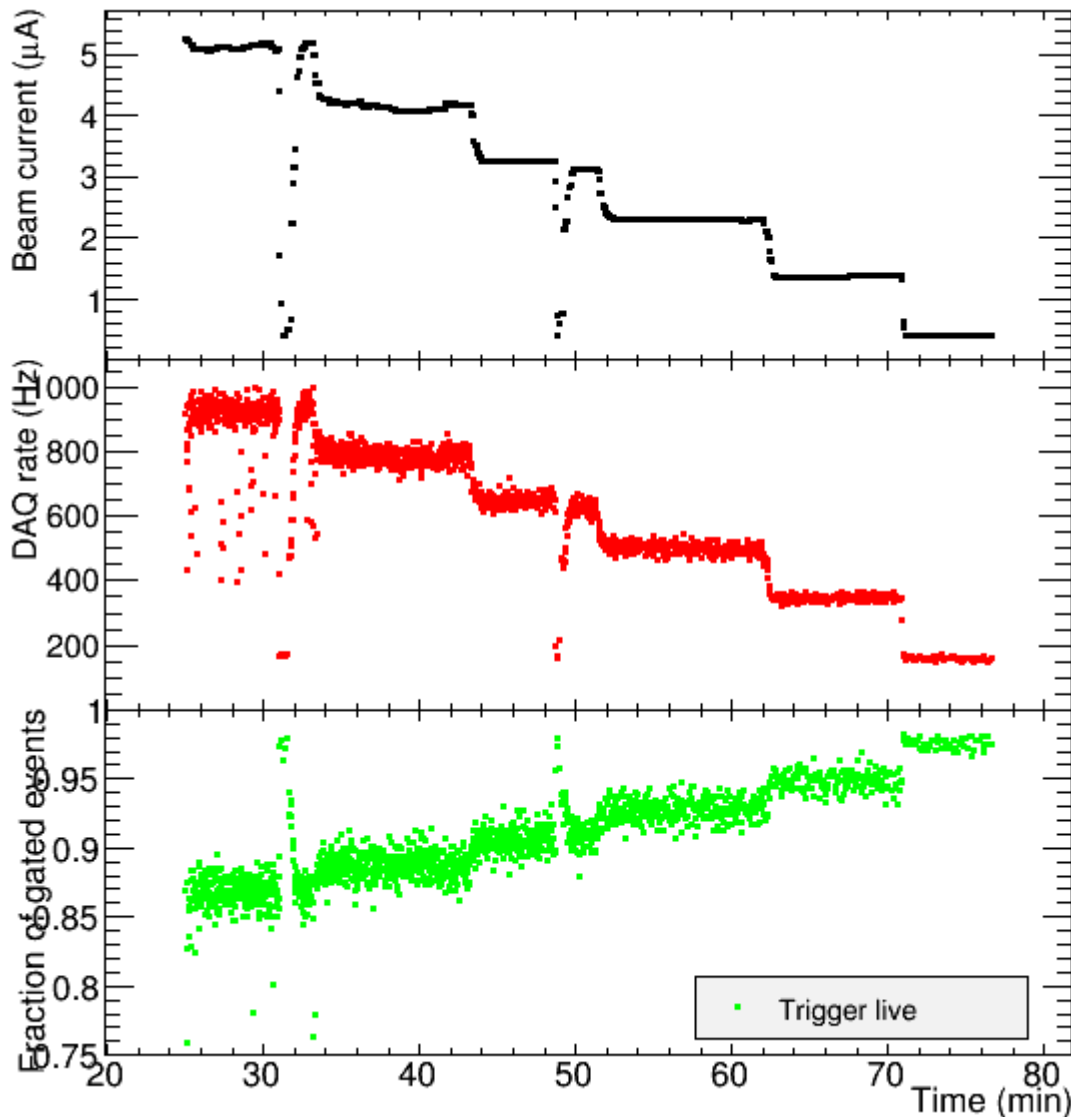
Two actual examples of pile up events
Taken few days ago, during a
production run



Algorithm built to fit and disentangle
these events

Preliminary studies

Production runs: Monitoring deadtimes closely

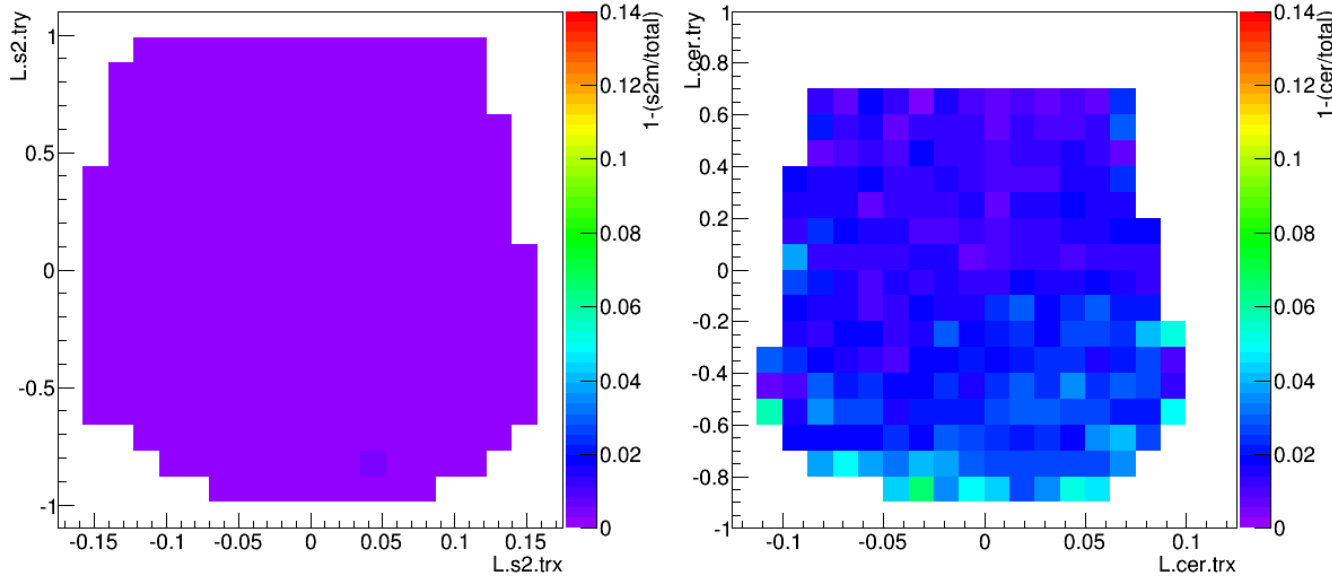


$$Yield = \frac{N_{events}}{Q \cdot DAQ_{LIVE} \cdot eff}$$

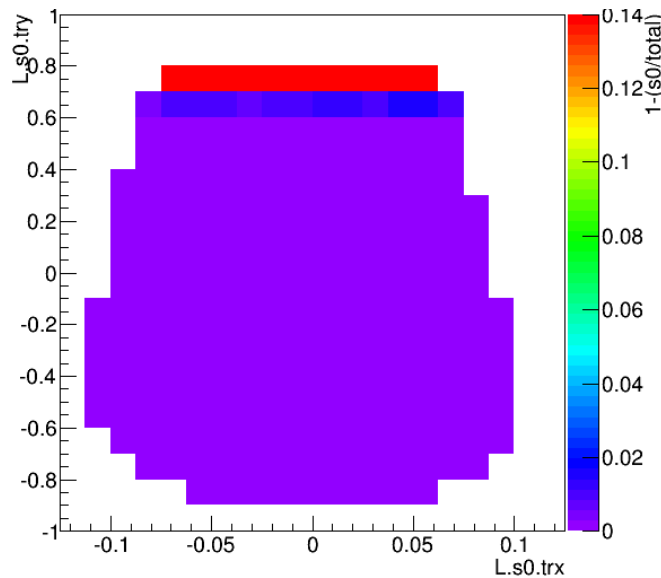
- DAQ deadtime is a direct normalization factor to the cross sections extractions
- Monitoring of DAQ deadtimes during the runs, for the calculation of deadtime weighted beam current
- In the example, a run taken at with several beam currents for deadtime studies.
- Few beam trips observed (will be removed later for the analysis)
- Mean deadtime of ~5% in a production run

Monitoring details of the DAQ: Trigger efficiencies

run 10374 - cer inefficiency



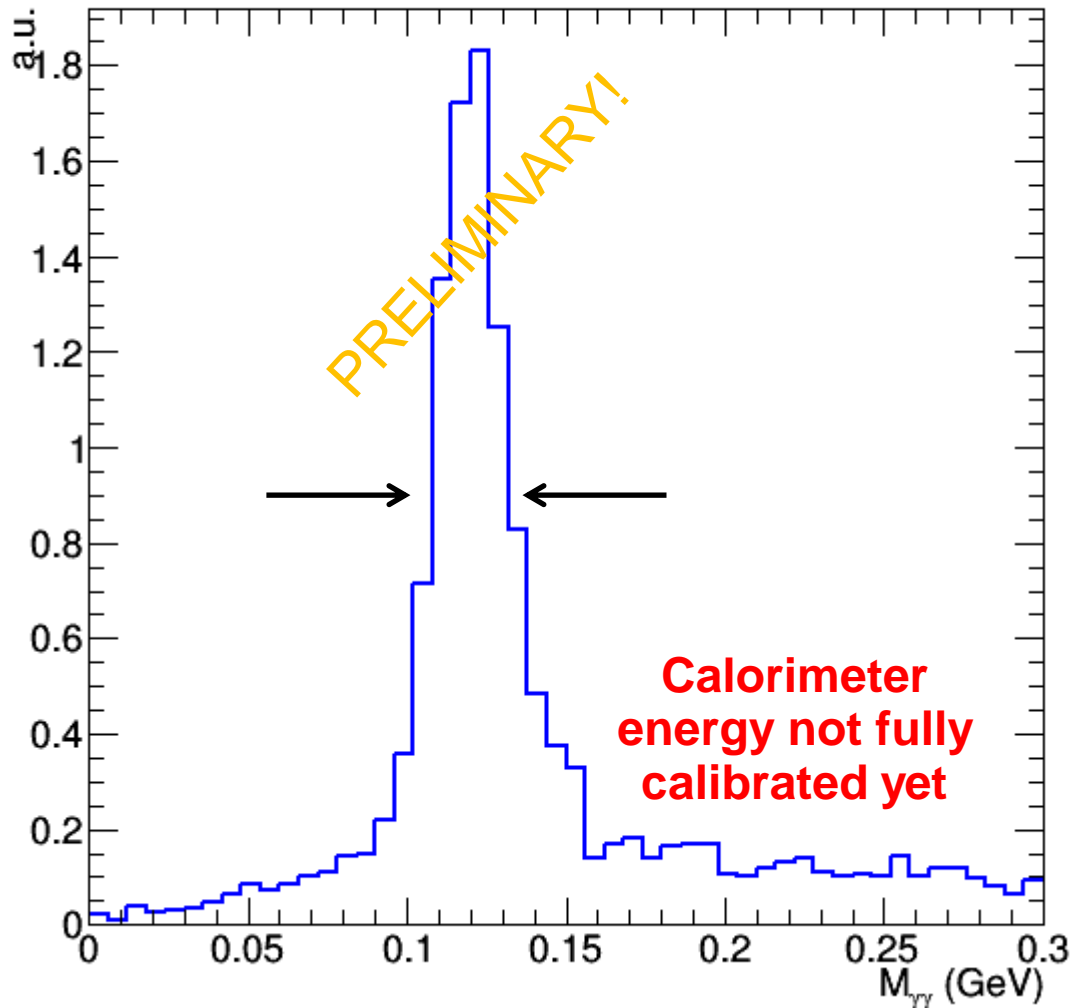
$$Yield = \frac{N_{events}}{Q \cdot DAQ_{LIVE}} \cdot \boxed{eff}$$



- Detector efficiencies are also a direct normalization factor to the cross sections extractions
- Monitoring trigger efficiencies for the different kinematics of the experiment

Very preliminary first data analysis

2 clusters events

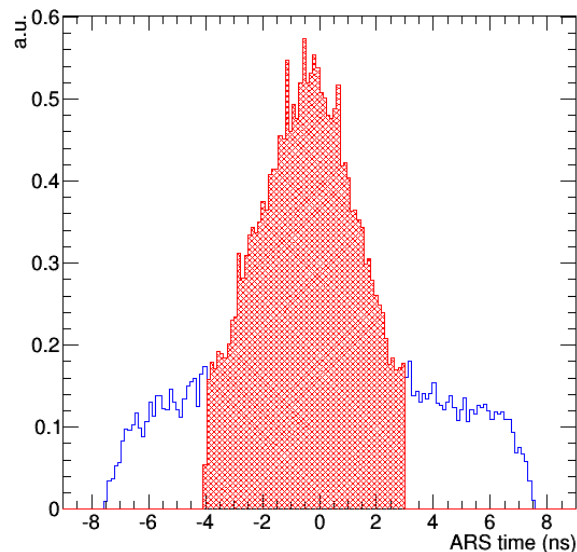


- Invariant mass reconstruction of 2 clusters events
- π^0 peaks at slightly shifted mass, since we still didn't tune all the calibration coefficients

Very preliminary analysis: One cluster events

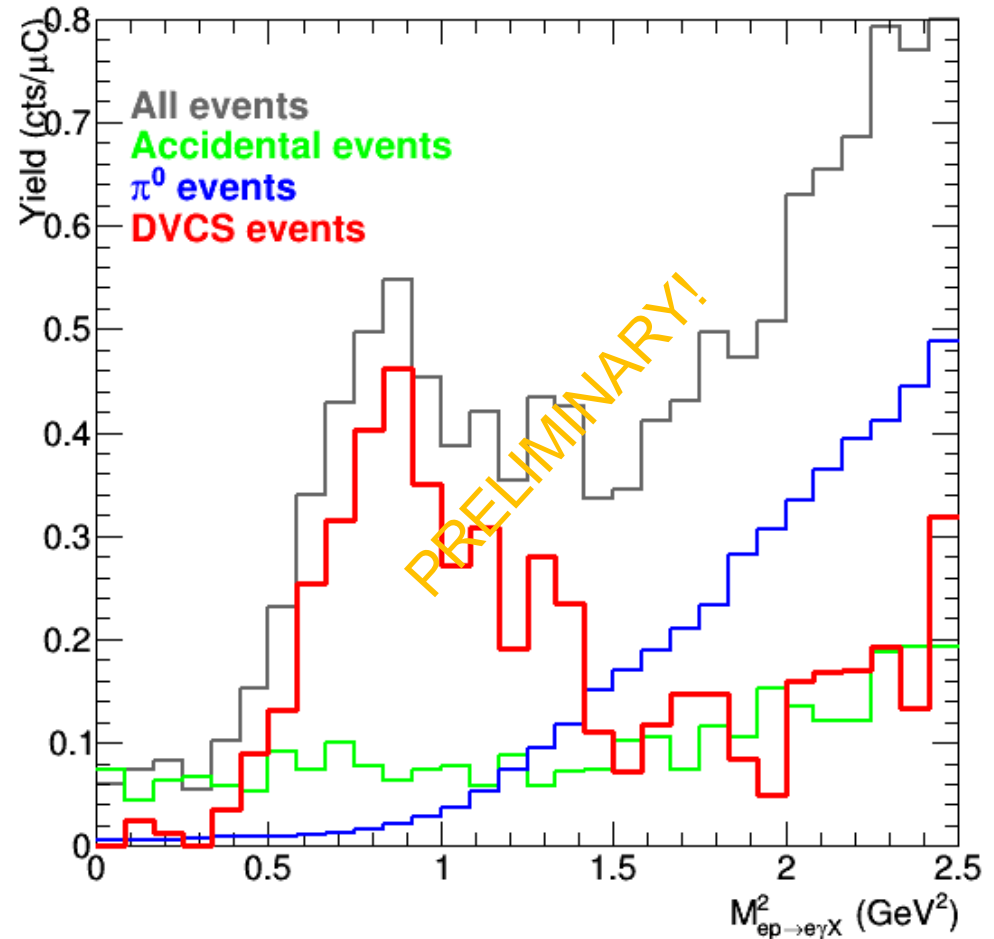
Some preliminary cuts applied:

- Vertex reconstructed in the target
- Photon energy above 0.6 GeV (expected DVCS range)
- PID: pion rejection
- Time of event at the calorimeter



Cut example: time of event in the calorimeter

Exclusive peak in missing mass DVCS events



2014/2015 data taking continues...

- Running the experiment together with the JLab 12 GeV commissioning
- 100 days approved to run the DVCS experiment at JLab Hall A
- Run time already scheduled/planned for 2016
- Scaling test of DVCS cross section for leading order factorization confirmation
- Extraction of t -dependent polarized and unpolarized DVCS cross section (and π^0 electroproduction) over a wide kinematic range
 - Q^2 from 2 to 9 GeV^2
 - $x_B = 0.36, 0.5, \text{ and } 0.6$

Thank you!

More about DVCS at Hall A at the APS meeting:
Lee Allison, "The Spring 2015 JLab Hall A Deeply Virtual Compton Scattering Run"
Sunday, April 12, 2015 - 1:30PM