# The Spin Asymmetries of the Nucleon Experiment - SANE 

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## Probing the Nucleon with Polarized Electromagnetic Scattering

## Inelastic e-nucleon Scattering

- Inclusive EM scattering is described by hadronic and leptonic tensors
- Symmetries reduce hadronic tensor to four structure functions (SF's):
- Symmetric part: unpolarized $\boldsymbol{W}_{\mathbf{1}}, \boldsymbol{W}_{2}$
- Anti-symmetric part:
spin-dependent $\boldsymbol{G}_{\mathbf{1}}, \boldsymbol{G}_{\mathbf{2}}$
- Anti-symmetric part:


Target fragments

Current jet (http://www.desy.de/~gbrandt/feyn/)

$$
\begin{aligned}
W_{\mu v}^{A}= & 2 \epsilon_{\mu v \lambda \sigma} q^{\lambda} \\
& \left\{M^{2} S^{\sigma} \boldsymbol{G}_{1}\left(v, Q^{2}\right)+\left[M v S^{\sigma}-p^{\sigma} S \cdot q\right] \boldsymbol{G}_{\mathbf{2}}\left(v, Q^{2}\right)\right\}
\end{aligned}
$$

Inclusive scattering: undetected final state

- $\boldsymbol{G}_{\mathbf{1}}, \boldsymbol{G}_{\mathbf{2}}, \boldsymbol{W}_{\mathbf{1}}$ and $\boldsymbol{W}_{2}$, contain all the information on nucleon structure that can be extracted from inclusive data


## Structure Functions in Inclusive DIS

- In high energy DIS, $\boldsymbol{G}_{\mathbf{1}}, \boldsymbol{G}_{\mathbf{2}}, \boldsymbol{W}_{\mathbf{1}}$ and $\boldsymbol{W}_{\mathbf{2}}$, become scaling functions of only one variable, up to log violations.

$$
\begin{array}{cl}
\lim _{Q^{2}, v \rightarrow \infty} M W_{1}\left(v, Q^{2}\right)=F_{1}(x) & \lim _{Q^{2}, v \rightarrow \infty} M^{2} v G_{1}\left(v, Q^{2}\right)=g_{1}(x) \\
\lim _{Q^{2}, v \rightarrow \infty} v W_{2}\left(v, Q^{2}\right)=F_{2}(x) & \lim _{Q^{2}, v \rightarrow \infty} M v^{2} G_{2}\left(v, Q^{2}\right)=g_{2}(x) \\
\text { Bjorken } x=Q^{2} /(2 M v) &
\end{array}
$$

- In the quark parton model, $\boldsymbol{g}_{1}$ and $\boldsymbol{F}_{1}$ can be related to parton distribution functions - PDF's:

$$
F_{1}(x)=\frac{1}{2} \sum e_{f}^{2}\left(q_{f}^{\uparrow}(x)+q_{f}^{\downarrow}(x)\right) \quad g_{1}(x)=\frac{1}{2} \sum e_{f}^{2}\left(q_{f}^{\uparrow}(x)-q_{f}^{\downarrow}(x)\right)
$$

(Index $f$ runs over active flavors)

## Virtual Compton Asymmetries

- The spin SF's are also related to virtual photon absorption crosssections and spin asymmetries (SA)
- the helicity of the virtual photon-nucleon system is $3 / 2$ or $1 / 2$ for transverse photons, $1 / 2$ for longitudinal ones
- SA $\boldsymbol{A}_{1}$ is defined in terms of the difference for $3 / 2$ and $1 / 2$ helicity cross sections

$$
\begin{aligned}
& A_{1}=\frac{\sigma_{T}^{(3 / 2)}-\sigma_{T}^{(1 / 2)}}{\sigma_{T}^{(3 / 2)}+\sigma_{T}^{1 / 2)}} \\
& A_{1}=\frac{1}{F_{1}}\left(g_{1}-\gamma^{2} g_{2}\right) ; \quad \gamma=\frac{2 x M}{\sqrt{Q^{2}}}
\end{aligned}
$$

- $\boldsymbol{A}_{2}$ represents the interference between initial transverse and final longitudinal amplitudes

$$
\begin{aligned}
& \boldsymbol{A}_{2}=\frac{\sigma_{T L}^{(1 / 2)}}{\sigma_{T}^{(3 / 2)}+\sigma_{T}^{(1 / 2)}} \leq \sqrt{\frac{A_{1}+1}{2} R} \leq \boldsymbol{R}=\frac{\sigma_{L}}{\sigma_{T}} \\
& A_{2}=\frac{\gamma}{F_{1}}\left(g_{1}+g_{2}\right)=\frac{\gamma}{F_{1}} \boldsymbol{g}_{T}
\end{aligned}
$$

## Model Independent Extraction of Spin Structure Functions

- $\boldsymbol{G}_{\mathbf{1}}$ and $\boldsymbol{G}_{\mathbf{2}}$ can be separated by measuring cross section differences for opposite beam helicities with target spins parallel and transverse to the beam

$$
\begin{array}{r}
\Delta \sigma\left(\theta, \theta_{N}, \phi\right)=\frac{4 \alpha^{2} E^{\prime}}{Q^{2} E}\left[\left(E \cos \theta_{N}+E^{\prime} \cos \alpha\right) M \boldsymbol{G}_{1}+2 E E^{\prime}\left(\cos \alpha-\cos \theta_{N}\right) \boldsymbol{G}_{2}\right] \\
\\
\cos \alpha=\sin \theta_{N} \sin \theta \cos \phi+\cos \theta_{N} \cos \theta, \quad(\theta, \phi: \text { final lepton angles })
\end{array}
$$

- transverse target spin $\theta_{N}$ : comparable $\boldsymbol{G}_{\mathbf{1}}, \boldsymbol{G}_{\mathbf{2}}$ terms

$$
\frac{d^{2} \sigma^{(\uparrow \rightarrow)}}{d \Omega d E^{\prime}}-\frac{d^{2} \sigma^{(\downarrow \rightarrow)}}{d \Omega d E^{\prime}}=\frac{4 \alpha^{2} E^{\prime}}{Q^{2} E} E^{\prime} \sin \theta \cos \phi\left[M \boldsymbol{G}_{1}\left(\nu, \boldsymbol{Q}^{2}\right)+2 E \boldsymbol{G}_{2}\left(\nu, \boldsymbol{Q}^{2}\right)\right]
$$

- $G_{1}$ is twist-2 (plus corrections)

4/10/15 $-\boldsymbol{G}_{2}$ has both twist-2 and twist-3 contributions

## Transverse Polarized Scattering: Unlocking Twist-3

- Twist-2 and twist-3 operators contribute at same order in transverse polarized scattering
- twist-2: handbag diagram
- twist-3: $q g q$ correlations
- direct access to twist-3 via $\boldsymbol{g}_{2}$ :
- interacting $q g q$ is first step to understanding confinement
- "Unique feature of spin-dependent scattering" (R. Jaffe)


## Why is $\boldsymbol{g}_{2}$ interesting?

- tests twist-3 effects = quark-gluon correlations
- higher twist corrections to $\boldsymbol{g}_{1}$ with $3^{\text {rd }}$ moment's $\boldsymbol{d}_{2}$ matrix element
- also test lattice $\mathrm{QCD}, \mathrm{QCD}$ sum rules, quark models
- $\boldsymbol{d}_{2}$ related to color Lorentz force on transverse polarized quark (M. Burkardt, AIP Conf. Proc. 1155 (2009) 26)
- sign of $\boldsymbol{d}_{2}$ related to sign of transverse deformation (anomalous $\kappa^{q}$ )
- polarizabilities of color fields (with twist-4 matrix element $f_{2}$ )
- magnetic $\chi_{\mathrm{B}}=\left(4 \boldsymbol{d}_{2}+f_{2}\right) / 3$ and electric $\chi_{\mathrm{E}}=\left(4 \boldsymbol{d}_{2}-2 f_{2}\right) / 3$.
- contains chiral odd twist-2 $=$ quark transverse spin (mass term)

4/10/15 - test quark masses (covariant parton models)

## $\boldsymbol{g}_{\mathbf{2}}$ and $\boldsymbol{g}_{\mathrm{T}}$ Spin Structure Functions

Experimentally measured quantities

$$
g_{T}(x)=g_{1}(x)+g_{2}(x)=\frac{1}{2} \sum e_{q}^{2} g_{T}^{q}(x)
$$

$g_{T}^{q}$ in terms of Transverse Momentum Dependent distributions ${ }_{[1]}$

$$
\begin{array}{r}
g_{T}(x)=\int d^{2} \vec{k}_{t} \frac{\vec{k}_{t}^{2}}{2 M^{2}} \frac{g_{1 \mathrm{~T}}^{q}\left(x, \vec{k}_{t}^{2}\right)}{x}+\frac{m}{M} \frac{h_{1}(x)}{x}+\tilde{g}_{T}(x) \\
\text { twist-3 TMD quark mass term qgq interaction }
\end{array}
$$

Applying twist-2 Wandzura-Wilczek approximation of $g_{2}$

$$
g_{2}^{w W}(x)=-g_{1}(x)+\int_{x}^{1} \frac{d y}{y} g_{1}(y)
$$

Twist-3 for the nucleon (neglecting quark mass)

$$
\bar{g}_{2}=\frac{1}{2} \sum e_{q}^{2}\left[\tilde{g}_{T}^{q}-\int_{x}^{1} \frac{d y}{y}\left(\tilde{g}_{T}^{q}(y)-\hat{g}_{T}^{q}(y)\right)\right] ; \quad \tilde{g}_{T}=q g \text { term, } \hat{\mathrm{g}}_{T}=\text { Lorentz invariance }[2]
$$

## Proton world $\mathrm{A}_{\| \mid}, \mathrm{A}_{\perp}$ data before SANE



- Two beam energies: 5.9 GeV , 4.7 GeV
- Very good high $\boldsymbol{x}$ coverage with detector at $40^{\circ}$


## Experiment

## Spin Asymmetries of the Nucleon Experiment (TJNAF E07-003)

SANE Collaboration
Argonne National Lab., Christopher Newport U., Florida International U., Hampton U., Jefferson Lab., U. of New Hampshire, Norfolk S. U., North Carolina A\&T S. U., Mississippi S. U., Ohio U., IHEP - Protvino, U. of Regina, Rensselaer Polytechnic I., Rutgers U., Seoul National U., Southern U. New Orleans,

Temple U., Tohoku U., U. of Virginia, Yerevan Physics I., Xavier U.
Spokespersons:
S. Choi (Seoul), M. Jones (JLab), Z-E. Meziani (Temple), O. A. Rondon (U. of Virginia)

Goal: Measure the proton spin structure function $\boldsymbol{g}_{\mathbf{2}}\left(x, Q^{2}\right)$ and spin asymmetry $\mathbf{A}_{1}\left(x, Q^{2}\right)$ for $2.5 \leq \boldsymbol{Q}^{2} \leq 6.5 \mathrm{GeV}^{2}$ and $0.3 \leq \boldsymbol{x} \leq 0.8$

Method: Measure parallel and near-transverse inclusive double spin asymmetries, detecting the electrons with novel non-magnetic large solid angle telescope BETA

## SANE Layout in JLab's Hall C


[1] Big Electron Telescope Array $\Delta \Omega \sim 190 \mathrm{msr} ; \Delta \theta= \pm 10^{\circ}$

## BETA with DIS electron simulation



[^0][4] Norfolk State U. and U. of Regina
[5] UVA- JLab

## Polarized Target



Data

## DATA

| Detector | Detected particle | Scatifering Type | Beam <br> Energy [GeV] | Field Direction | Target |
| :---: | :---: | :---: | :---: | :---: | :---: |
| BETA | $e, \pi^{0}$ | Inclusive inelastic | 5.9, 4.7 | $180^{\circ}, 80^{\circ}$ | NH3 |
| HMS | $e$ | Inclusive inelastic | 5.9, 4.7 | $180^{\circ}, 80^{\circ}$ | NH3 <br> C, $\mathrm{LHe}^{[1]}$ |
|  |  | Inclusive elastic | 5.9 | $80^{\circ}$ | NH3 |
| BETA - HMS | $e-p$ | Coincidence elastic | 5.9 | $80^{\circ}$ | NH3 |

- Data taken in January - March 2009


## BETA and HMS data



- $Q^{2}-x$ phase space of BETA's $80^{\circ}$ data

- Central kinematics of HMS inclusive asymmetry data
- cut on $\mathrm{E}^{\prime} \geq 1.3 \mathrm{GeV}$


## Measured Asymmetries $\mathrm{A}\left(80^{\circ}\right), \mathrm{A}\left(180^{\circ}\right)$

$$
\begin{gathered}
A_{m}=\frac{\epsilon}{f P_{b} P_{t} C_{N}} ; \epsilon=\frac{N^{-}-N^{+}}{N^{-}+N^{+}} \\
A_{p h y s}=\frac{1}{f_{r c}}\left(\frac{A_{m}-f_{b} A_{b}}{1-f_{b}}\right)+A_{r c}
\end{gathered}
$$

- $\boldsymbol{N}^{+,-}=$charge normalized, dead time corrected yields

- $\boldsymbol{P}_{\mathbf{b}^{\prime}} \boldsymbol{P}_{\mathrm{t}}=$ beam, target polarizations
- $\boldsymbol{f}=$ polarized dilution factor
- $\boldsymbol{C}_{\mathrm{N}}={ }^{14} \mathrm{~N}$ polarization correction
- $\boldsymbol{A}_{b}, f_{b}=$ background corrections

${ }_{4 / 10 / 15} \boldsymbol{A}_{r c} \boldsymbol{f}_{r c}=$ radiative corrections


## Preliminary Results

## Spin Asymmetries $\boldsymbol{A}_{\mathbf{1}}$ and $\boldsymbol{A}_{\mathbf{2}}$

- HMS single arm data in the resonances, $\left\langle Q^{2}\right\rangle \sim 1.8 \mathrm{GeV}^{2}$

$$
A_{1}=\frac{1}{D^{\prime}}\left(\frac{E-E^{\prime} \cos \theta}{E+E^{\prime}} A_{180}+\frac{E^{\prime} \sin \theta}{\left(E+E^{\prime}\right) \cos \phi} \frac{A_{180} \cos 80^{\circ}+A_{80}}{\sin 80^{\circ}}\right)
$$

- Model independent separation from

$$
A_{2}=\frac{1}{D^{\prime}} \frac{1}{2 E}\left(\sqrt{Q^{2}} A_{180}-\sqrt{Q^{2}} \frac{E-E^{\prime} \cos \theta}{E^{\prime} \sin \theta \cos \phi} \frac{A_{180} \cos 80^{\circ}+A_{80}}{\sin 80^{\circ}}\right)
$$ measured asymmetries




## Spin Asymmetries $\boldsymbol{A}_{\mathbf{1}}$ and $\boldsymbol{A}_{\mathbf{2}}$

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$$
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$$

- Model independent separation from

$$
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$$ measured asymmetries



(H-y. Kang)

## DIS Spin Asymmetry $\boldsymbol{A}$

- $\boldsymbol{A}_{1}(W)$ shows clear decreasing trend
- SANE BETA data
- statistical errors only
- SLAC data plotted for individual spectrometers
- very broad $Q^{2}$ range
- CLAS data at same $W$ but different $Q^{2}$ merged for
 clarity


## $\boldsymbol{g}_{1}$ and $\boldsymbol{g}_{\mathbf{2}}$ in DIS and Resonances

- BETA proton data
- DIS and Resonances
- $\boldsymbol{g}_{1}, \boldsymbol{g}_{2}{ }^{\mathrm{wW}}$ curves from PDF's at $4 \mathrm{GeV}^{2}$
- $E^{\prime} \geq 0.6 \mathrm{GeV}$
- more data at $1.6 \mathrm{GeV}^{2}$ coming
- SLAC E143, E155, E155x, SMC and HERMES DIS



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## Operator Product Expansion for Spin SF's

- OPE relates Cornwall-Norton moments to matrix elements of twist-2 $\boldsymbol{a}_{\mathrm{N}}$ and twist- $3 \boldsymbol{d}_{\mathrm{N}}$ $\int_{0}^{1} x^{N} g_{1}\left(x, Q^{2}\right) d x=\frac{\boldsymbol{a}_{N}}{2}+t m c, \quad N=0,2,4, .$.
$\int_{0}^{1} x^{N} g_{2}\left(x, Q^{2}\right) d x=\frac{N\left(\boldsymbol{d}_{N}-\boldsymbol{a}_{N}\right)}{2(N+1)}+t m c, N=2,4, .$.
(tme :target mass corrections)
- $\boldsymbol{d}_{2}$ is mean color-magnetic
 field along spin
- Nachtmann moments needed to get twist-3 free of tmc

$$
\boldsymbol{d}_{2}\left(\boldsymbol{Q}^{2}\right)=\int_{0}^{1} d x \xi^{2}\left(2 \frac{\xi}{x} g_{1}+3\left(1-\frac{\xi^{2} M^{2}}{2 Q^{2}}\right) g_{2}\right) \Rightarrow_{Q^{2} \rightarrow \infty} \int_{0}^{1} d x x^{2}\left(2 g_{1}+3 g_{2}\right)
$$

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- $\boldsymbol{d}_{2}$ is mean color-magnetic field along spin
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- SANE analysis final version
- Publications in preparation

$$
\boldsymbol{d}_{2}\left(\boldsymbol{Q}^{\mathbf{2}}\right)=\int_{0}^{1} d x \xi^{2}\left(2 \frac{\xi}{x} g_{1}+3\left(1-\frac{\xi^{2} M^{2}}{2 Q^{2}}\right) g_{2}\right) \Rightarrow_{Q^{2} \rightarrow \infty} \int_{0}^{1} d x x^{2}\left(2 g_{1}+3 g_{2}\right)
$$

## $G_{\mathrm{E}}{ }^{\mathrm{p}} / G_{\mathrm{M}}{ }^{\mathrm{p}}$ from inclusive and coincidence

## data

Ratio from:

- SANE inclusive HMS data at $Q^{2}=2.06 \mathrm{GeV}^{2}$
- $A_{\text {el }}^{\text {p }}=-0.20 \pm 0.02$
- $G_{\mathrm{E}}^{\mathrm{p}} / G_{\mathrm{M}}^{\mathrm{p}}=0.60 \pm 0.18 \pm 0.06$
(statistical + systematic error)
- BETA-HMS $e-p$ coincidences at $Q^{2}=5.66 \mathrm{GeV}^{2}$
- $G_{\mathrm{E}}^{\mathrm{p}} / G_{\mathrm{M}}^{\mathrm{p}}=0.67 \pm 0.36$
(statistical error only)

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## Extras

## DIS Transverse $\operatorname{Spin} \operatorname{SF} \boldsymbol{g}_{\mathrm{T}}{ }^{\mathrm{p}}$



- $\boldsymbol{g}_{\mathrm{T}}{ }^{\mathrm{p}}=\boldsymbol{F}_{1} \boldsymbol{A}_{2} / \gamma$ measures spin distribution normal to $\gamma^{*}$
- $\operatorname{SANE}\left\langle\boldsymbol{g}_{\mathrm{T}}{ }^{\mathrm{p}}(x>.3)>=0.023 \pm 0.006\right.$

- Bag Model (1990's)
- Data scaled $\times 2.5$
- Model updates needed


## Sample of Normalizations and corrections






## Pair-symmetric background - I

- BigCal detects both charge signs
- Significant background from $e^{+} e^{-}$from $\pi^{0}$ decays
- Partial control with cut on $E^{\prime} \geq 1.3 \mathrm{GeV}$; worst dilution $<\sim 0.2$
- Estimate with GEANT simulation of $\pi^{0}$ production
- Need inclusive pion photo- and electro-production cross sections
- Existing D. Wiser parametrization only for H, D targets
- Parameterized Yerevan pion photoproduction data on C at 4.5 GeV
- Cross section scales with pion $P_{\mathrm{T}}$ : use simple exponential scaling fit
- Included fit in J. O'Connell EPC code for single arm hadron photo and electroproduction
- Compared with DESY electroproduction on C at 5 GeV


## Pair-symmetric background - II



K. Alanakian et al., JETP Lett. 32 (1980) 652

- Fitted $\pi^{+}, \pi^{-}$data at $20^{\circ}, 40^{\circ}$, $60^{\circ}$ to $\boldsymbol{\sigma}\left(\boldsymbol{P}_{\mathrm{T}}\right)=\boldsymbol{a} \boldsymbol{e}^{-b \boldsymbol{T} \mathrm{~T}}$

- Wiser $\pi^{-}$data on $H$, scaled $12 \times$, along with $\pi^{-}$data on C and scaling fit to C data.
- $\pi^{0}$ fit from average of $\pi^{+}$and $\pi^{-}$


## Pair-symmetric background - III




- Test of scaling fit with DESY $\mathrm{C}\left(\mathrm{e}, \pi^{-}\right)$data at $5 \mathrm{GeV}, 13^{\circ}$
[1] J. O'Connell CEBAF Summer Workshop, F. Gross and J. Lightbody, ed., Newport News, 1988, p. 345


## Jefferson Angular Momentum - JAM Collaboration



- Joint theorists and experimentalists effort to "study the quark and gluon spin structure of the nucleon by performing global fits of PDFs".
- JAM's spin PDFs are tailored for studies at large Bjorken $\underline{x}$, as well as the resonance-DIS transition region at low and intermediate $\boldsymbol{W}$ and $\boldsymbol{Q}^{2}$. http://wwwold.jlab.org/theory/jam/


## Big Electron Telescope Array - BETA

- BETA specs
- Effective solid angle 0.194 sr
- Energy resolution

$$
10 \% / \sqrt{ } E(\mathrm{GeV})
$$

- 1000:1 pion rejection
- angular resolution $\sim 1 \mathrm{mr}$
- Non-magnetic detector
- detects DIS $e$ and $e^{+} e^{-}$pairs: need to cut on minimum $E^{\prime}$
- Target field helps sweep lowest $E$ background ( $180 \mathrm{MeV} / \mathrm{c}$ cutoff)



## Nucleon Spin beyond $G_{1}$ and $G_{2}$

- Need to go beyond $a_{0}$ to understand nucleon spin
- Orbital angular momentum (OAM) $\boldsymbol{L}$ is needed.
- Partons have transverse momentum, implies OAM
- Mulders et al., Transverse Momentum dependent Distributions - TMDs
- functions of $x$ and $k_{t}$
- Semi-inclusive scattering (detect final $e$, one hadron)

| Transverse Momentum Distributions by Polarization |  |  |  |
| :---: | :---: | :---: | :---: |
| Target $\downarrow$ \quark $\rightarrow$ | U | L | $T$ |
| U | $f_{1}\left(x, k_{t}\right)$ |  | $h_{1}{ }^{\perp}$ |
| L |  | $g_{1}$ | $h_{11}{ }^{\text {d }}$ |
| $T$ | $f_{1 T}{ }^{\text {d }}$ | $g_{1 T}{ }^{\text {- }}$ | $h_{1} h_{1 T}{ }^{\text {d }}$ |

Longitudinal SSF (leading twist)
$g_{1}(x)=\sum g_{1}^{q}(x)=\sum \int d^{2} \vec{k}_{t} g_{1 L}\left(x, \vec{k}_{t}^{2}\right)$
Transverse SSF (twist-3)
$g_{1 \mathrm{~T}}^{(1)}(x)=\sum g_{1 \mathrm{~T}}^{q(1)}(x)=\sum \int d^{2} \vec{k}_{t} \frac{\vec{k}_{t}^{2}}{2 M^{2}} g_{1 \mathrm{~T}}^{q}\left(x, \vec{k}_{t}^{2}\right)$
$g_{T}(x)=g_{1}(x)+\frac{d}{d x} g_{1 \mathrm{~T}}^{(1)}=g_{1}(x)+g_{2}(x)$


[^0]:    [1] BigCal Collaboration
    [2] North Carolina A\&T U.
    [3] Temple U

