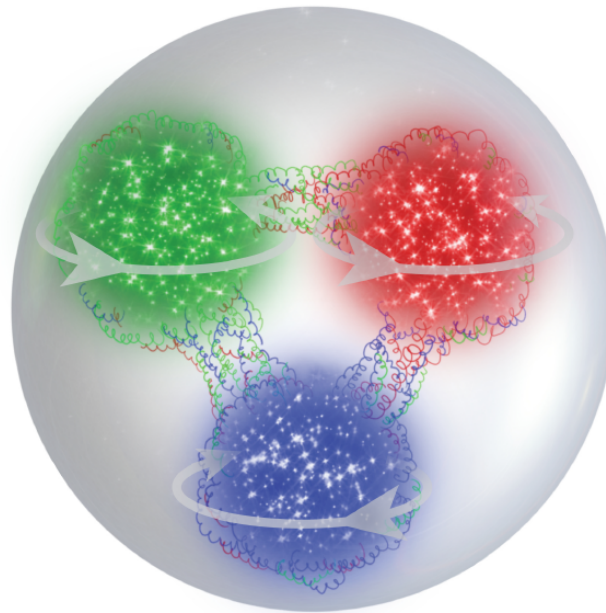
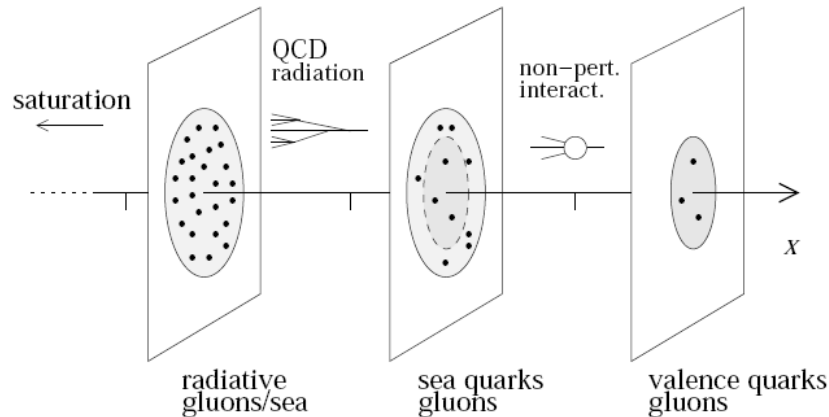
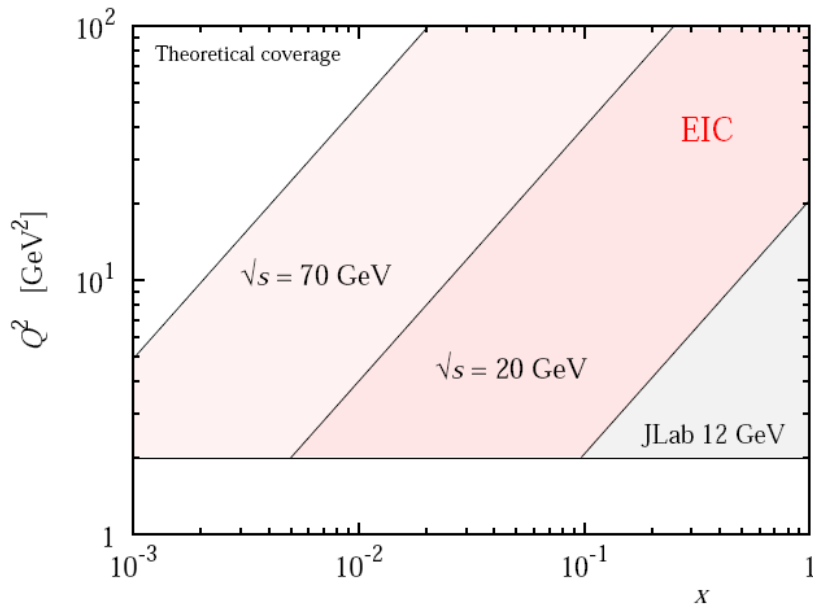


Present status of Phenomenology of the Transverse Spin

Alexei Prokudin



Nucleon landscape



Nucleon is a many body dynamical system of quarks and gluons

Changing x we probe different aspects of nucleon wave function

How **partons move** and how they are distributed in **space** is one of the future directions of development of nuclear physics

Technically such information is encoded into Generalised Parton Distributions and Transverse Momentum Dependent distributions

These distributions are also referred to as **3D (three-dimensional) distributions**

Unified View of Nucleon Structure

Wigner function

5D

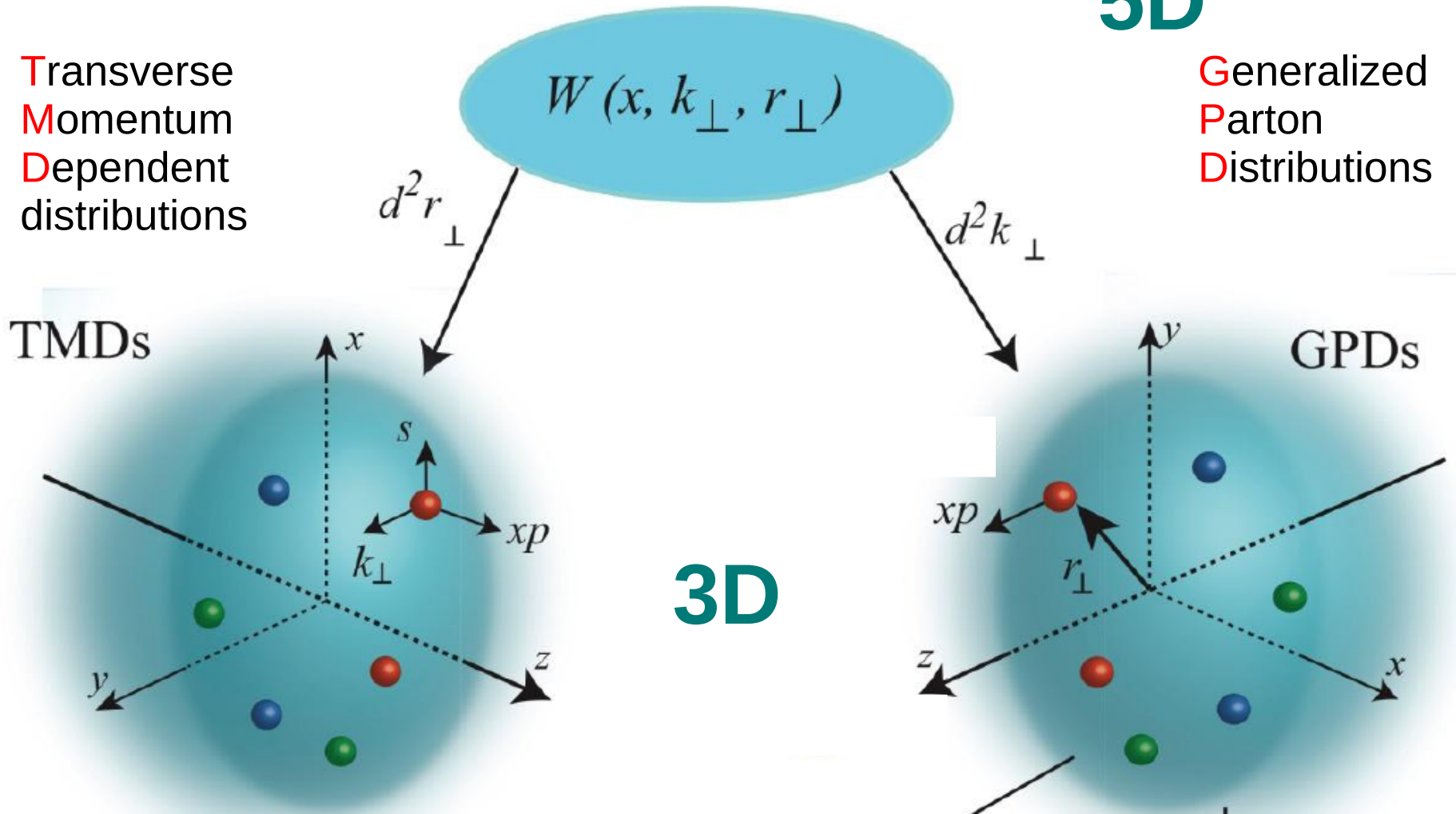
Transverse
Momentum
Dependent
distributions

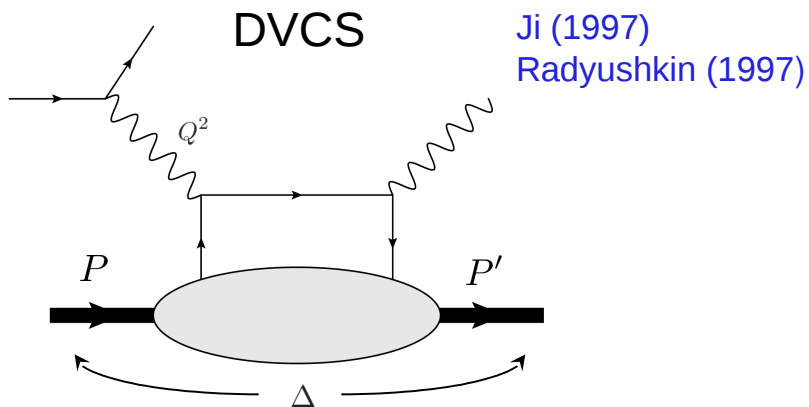
Generalized
Parton
Distributions

TMDs

GPDs

3D





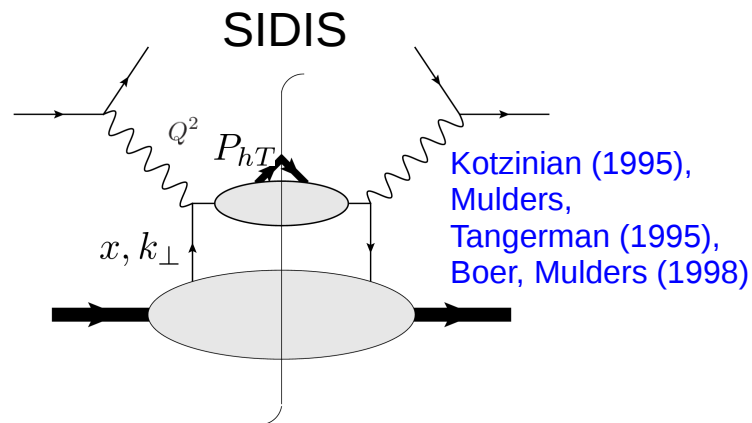
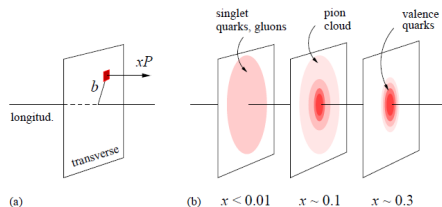
Q^2 ensures hard scale, pointlike interaction

$\Delta = P' - P$ momentum transfer can be varied independently

Connection to 3D structure Burkardt (2000)
Burkardt (2003)

$$\rho(x, \vec{r}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{-i\vec{\Delta}_\perp \cdot \vec{r}_\perp} H_q(x, \xi = 0, t = -\vec{\Delta}_\perp^2)$$

Drell-Yan frame $\Delta^+ = 0$ Weiss (2009)

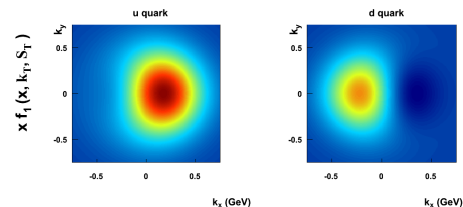


Q^2 ensures hard scale, pointlike interaction

P_{hT} final hadron transverse momentum can be varied independently

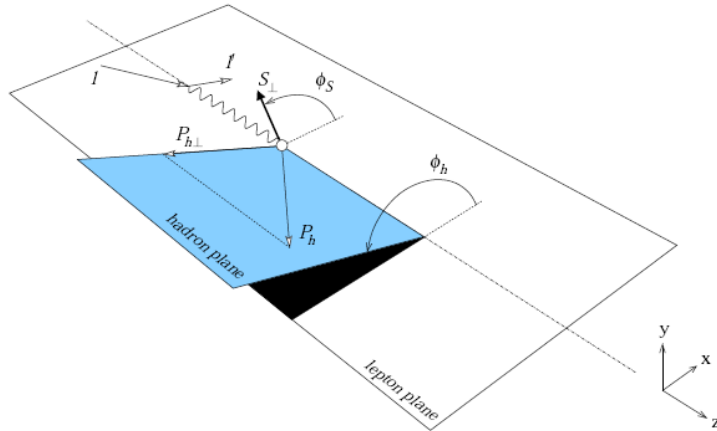
Connection to 3D structure Ji, Ma, Yuan (2004)
Collins (2011)

$$\tilde{f}(x, \vec{b}_T) = \int d^2 k_\perp e^{-i\vec{b}_T \cdot \vec{k}_\perp} f(x, \vec{k}_\perp)$$



AP (2012)

Semi Inclusive Deep Inelastic Scattering



One can rewrite the cross-section in terms of **18** structure functions

Each structure function encodes parton dynamics via convolutions of TMDs when factorization is applicable

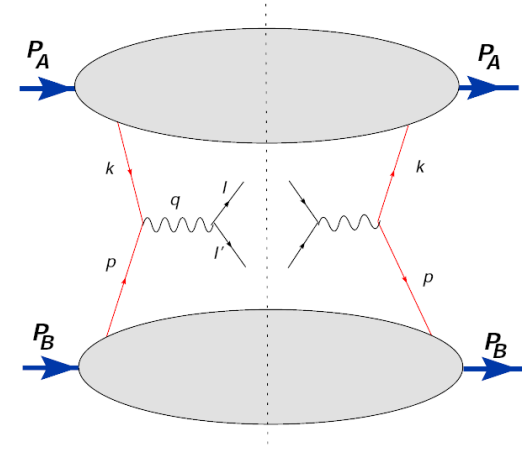
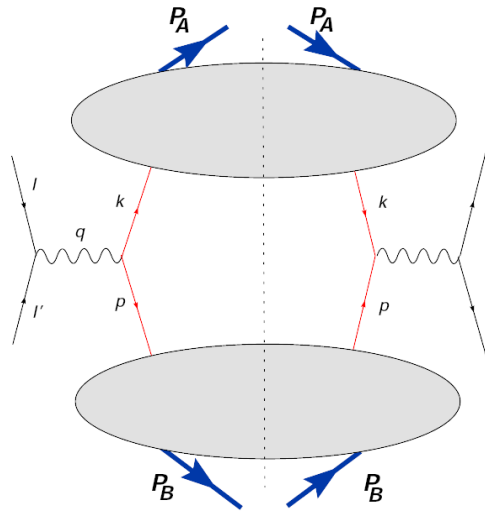
Mulders, Tangerman (1995),
Boer, Mulders (1998)
Bacchetta et al (2007)

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} =$$

$$\frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right.$$

$$\left. + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} + \dots \right.$$

e+e- and Drell-Yan



One can rewrite the cross-section of e+e- in terms of **72** structure functions

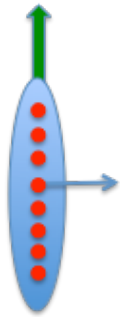
Boer, Jacob, Mulders (1997),
Pitonyak, Schlegel, Metz (2013)

One can rewrite the cross-section of Drell-Yan in terms of **48** structure functions

Tangerman, Mulders (1995),
Boer (1999),
Arnold, Metz, Schlegel (2009)

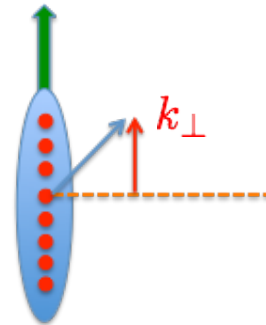
From 1-D to 3-D

Hadron structure: one-dimensional picture to three-dimensional tomography



$$f(x)$$

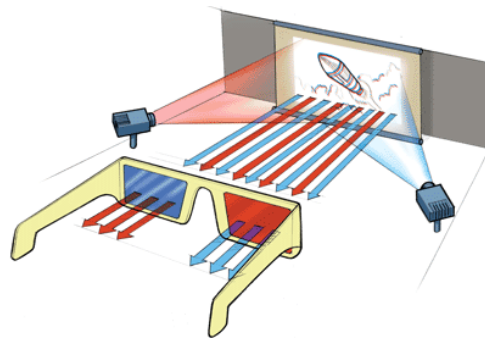
Collinear PDFs



$$f(x, k_{\perp})$$

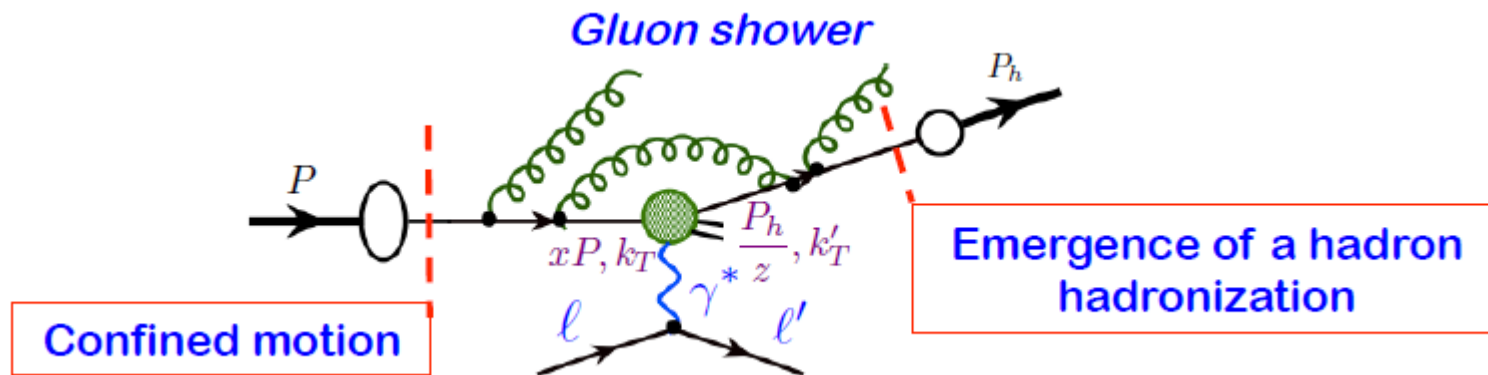
Transverse momentum distributions (TMDs)

Very interesting and non-trivial consequences:
rich QCD dynamics and new insight on hadron structure



Why QCD evolution is interesting?

Study of evolution gives us insight on different aspects and origin of confined motion of partons, gluon radiation, parton fragmentation



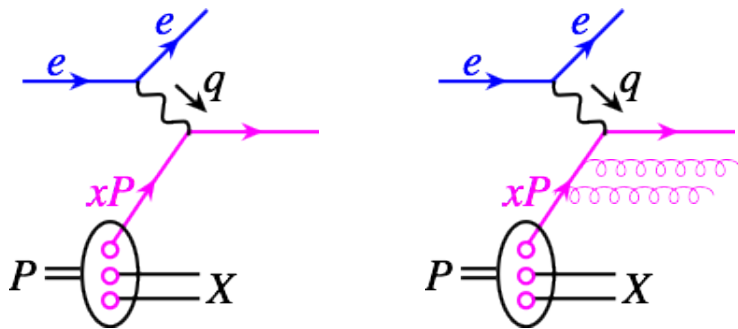
Evolution allows to connect measurements at very different scales

What do we mean by QCD evolution?

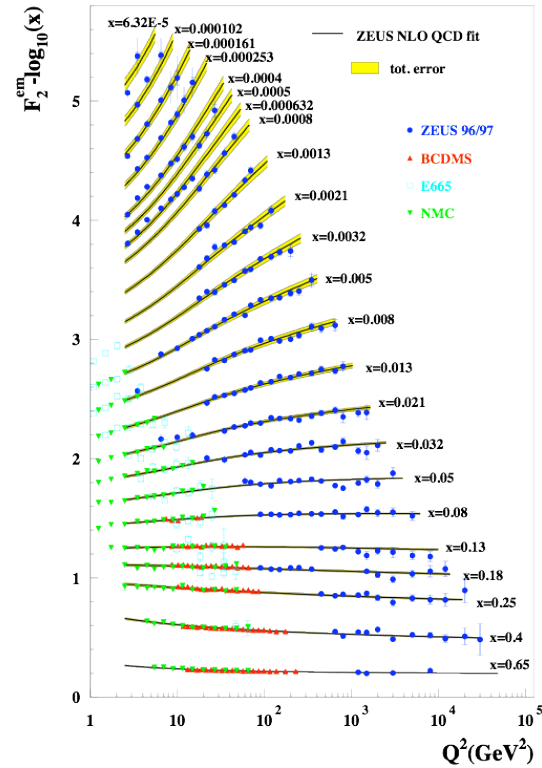
Very well known example:
DGLAP evolution of collinear
parton distributions

Take into account perturbative
corrections

Single logarithms are resummed
order by order in perturbative
calculations



$$\left(\alpha_s \ln \frac{Q^2}{\mu^2} \right)^n$$



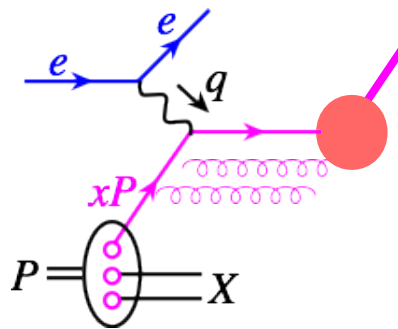
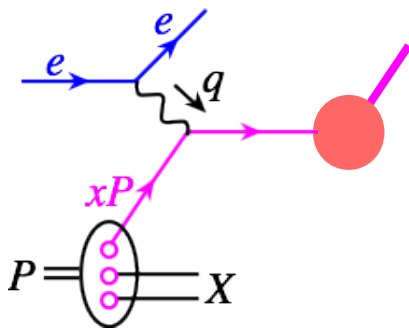
What do we mean by QCD evolution?

TMD factorization is applicable in case two different scales are observed in processes such as SIDIS, Drell-Yan, W/Z production in hadron-hadron collisions. Kinematical regime: $Q_T \ll Q$

For SIDIS Q_T is transverse momentum of final hadron

Again we need to take into account perturbative corrections

Double logarithms are resummed order by order in perturbative calculations



$$\left(\alpha_s \ln^2 \frac{Q^2}{Q_T^2} \right)^n$$

Approaches to TMD evolution

Collins-Soper-Sterman (CSS) resummation framework

Collins-Soper-Sterman 1985
ResBos: C.P. Yuan, P. Nadolsky
Qiu-Zhang 1999, Vogelsang etc...
Kang-Xiao-Yuan 2011, Sun-Yuan 2013
Prokudin-Kang-Sun-Yuan 2014

“New” Collins approach

Collins 2011
Aybat-Rogers 2011,
Aybat-Collins-Rogers-Qiu, 2012
Aybat-Prokudin-Rogers 2012
Anselmino-Boglione-Melis 2012
Prokudin-Bacchetta 2013
Echevarria-Idilbi-Kang-Vitev 2014

Soft Collinear Effective Theory (SCET)

Echevarria-Idilbi-Schafer-Scimemi 2012
D'Alesio-Echevarria-Melis-Scimemi 2014

Approaches to TMD evolution

Different approaches are essentially identical

Phenomenological results vary however due to different treatment of initial conditions

TMD evolution in a nut shell

TMD functions are measured at scale Q $f(x, k_{\perp}; Q)$

Evolution is performed in Fourier space

$$\tilde{f}(x, b; Q) = \int d^2 k_{\perp} e^{-ik_{\perp} b} f(x, k_{\perp}; Q)$$

Standard CSS formalism, evolution starts from $\mu_b = c/b$, $c = 2e^{-\gamma_E}$

$$\tilde{f}(x, b; Q) = \tilde{f}(x, b; \mu_b) e^{-S_{pert}(b)}$$

$$S^{PT}(b) = \int_{\mu_b}^Q \frac{d\mu}{\mu} \left(A \ln \frac{Q^2}{\mu^2} + B \right) \quad \text{Perturbative Sudakov factor}$$

$$A = \sum_{n=1} \left(\frac{\alpha_s}{\pi} \right)^n A^{(n)} \quad B = \sum_{n=1} \left(\frac{\alpha_s}{\pi} \right)^n B^{(n)}$$

TMD evolution in a nut shell

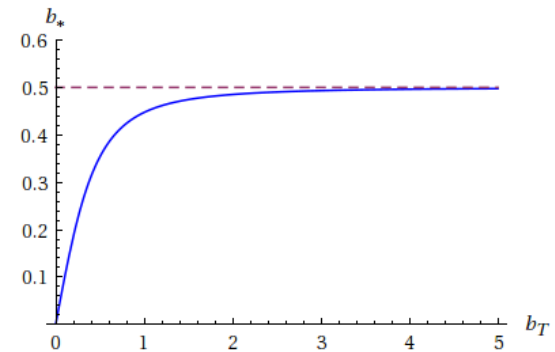
Calculation is perturbative, valid only in region $b \ll 1/\Lambda_{QCD}$

Fourier transform in momentum space involves non-perturbative region

$$f(x, k_{\perp}; Q) = \int_0^{\infty} \frac{bdb}{2\pi} J_0(k_{\perp} b) \tilde{f}(x, b; Q)$$

Non perturbative region needs to be treated.
Common method b^* prescription

$$b_* = \frac{b}{\sqrt{1 + b^2/b_{max}^2}}$$



$$\tilde{f}(x, b; Q) = \tilde{f}(x, b_*; c/b_*) e^{-S_{pert}(b_*)} e^{-S_{NP}(b)}$$

Non perturbative Sudakov factor

TMD evolution in a nut shell

Relation to collinear functions at small values of b:

$$\tilde{f}^j(x, b_*; c/b_*) = \sum_{j'=q,g} \int_x^1 \frac{d\hat{x}}{\hat{x}} C_{j/j'}\left(\frac{x}{\hat{x}}, b_*; c/b_*\right) f^{j'}(x; c/b_*)$$

$$C = \sum_{n=1} \left(\frac{\alpha_s}{\pi}\right)^n C^{(n)} \quad \text{Wilson coefficient}$$

Collinear PDF

For transversity and helicity TMDs:

[Bacchetta-Prokudin 2013](#)

For Collins function (relation to twist-3 function): [Yuan-Zhou 2009, Kang 2011](#)

In future also gluon functions will be important

For gluon twist-3 function: [Dai-Kang-Prokudin-Vitev 2014](#)

*Taking into account Wilson coefficients is very important!
Large K factors of collinear computations between LO and NLO!*

TMD evolution in a nut shell


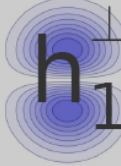




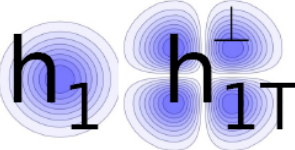
Precision of extraction depends on precision of calculations

$$\begin{array}{llll} \text{Leading Log (LL):} & A^{(1)} & & \\ \text{Next-to Leading Log (NLL):} & A^{(1,2)} & B^{(1)} & C^{(1)} \\ \text{Next-to-Next-to Leading Log (NNLL):} & A^{(1,2,3)} & B^{(1,2)} & C^{(1)} \end{array}$$

Precision is important!

$C^{(1)}$ means that one should use NLO collinear distributions

TMD distributions

$N \backslash q$	U	L	T
U			
L			
T			








8 functions in total (at leading twist)

Each represents different aspects of partonic structure

Each function is to be studied

Kotzinian (1995), Mulders, Tangerman (1995), Boer, Mulders (1998)

TMD Fragmentation Functions

$N \backslash q$	U	L	T
U			
L			
T			

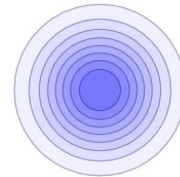
8 functions describing fragmentation of a quark into spin $\frac{1}{2}$ hadron

Mulders, Tangerman (1995), Meissner, Metz, Pitonyak (2010)

TMD distributions

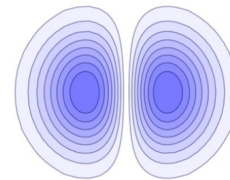
Three types of modulations

$$f(x, \mathbf{k}_{\perp}^2)$$



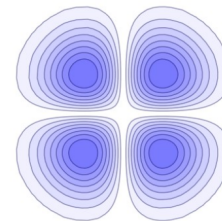
Monopole

$$\frac{\mathbf{k}_{\perp i} S_{T i}}{M} f(x, \mathbf{k}_{\perp}^2)$$



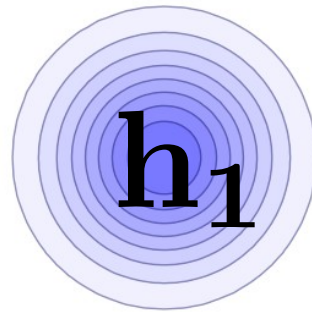
Dipole

$$\frac{\mathbf{k}_{\perp}^i \mathbf{k}_{\perp}^j - \frac{1}{2} \mathbf{k}_{\perp}^2 g_T^{ij}}{M^2} f(x, \mathbf{k}_{\perp}^2)$$

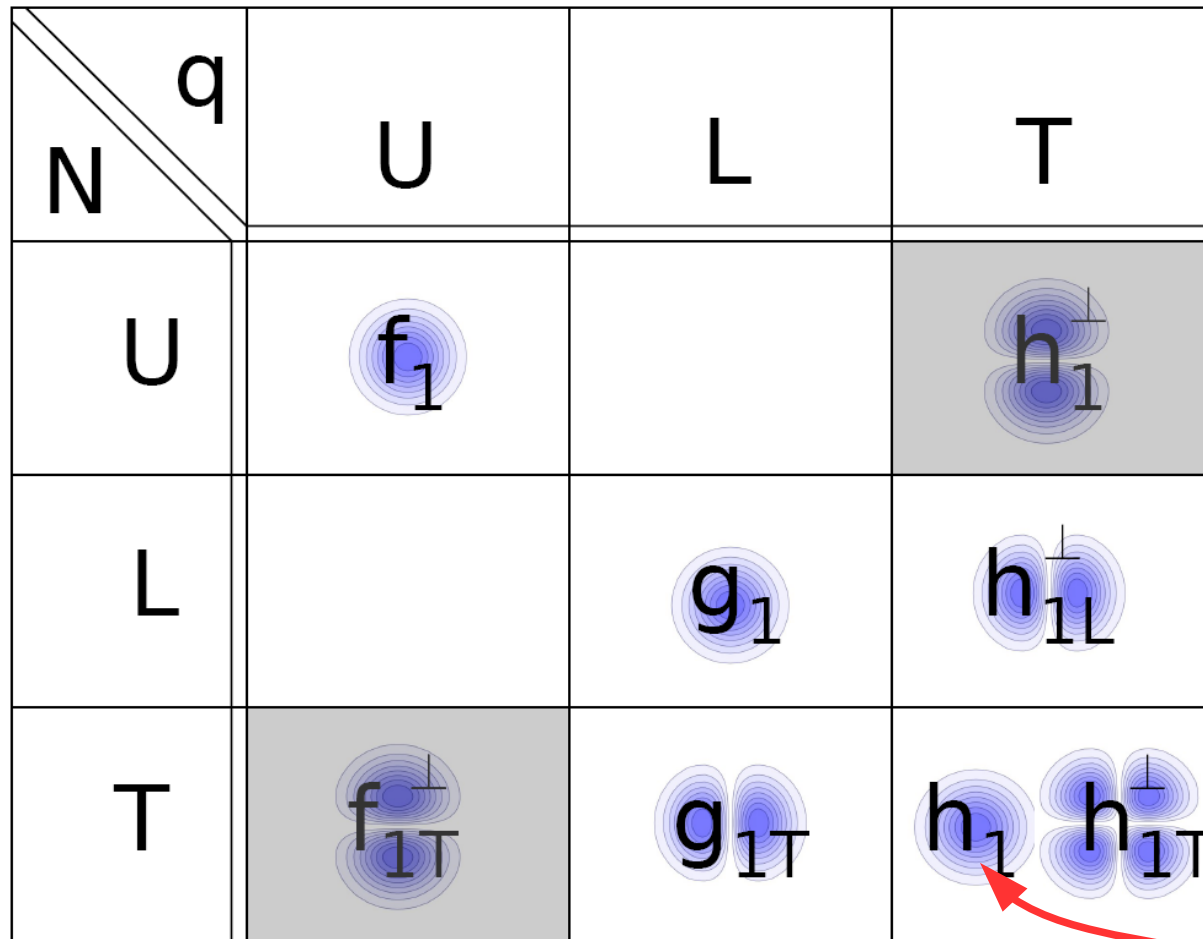


Quadrupole

Kotzinian (1995), Mulders, Tangerman (1995), Boer, Mulders (1998)



TMD distributions



Transversity

Kotzinian (1995), Mulders, Tangerman (1995), Boer, Mulders (1998)

Twist-2 collinear PDFs

Quark-quark correlator can be decomposed by means of
3 Parton Distributions Functions (PDF) in collinear (kt integrated) case

$$\Phi(x; P, S) = \frac{1}{2} \left\{ f_1(x) \not{P} + S_L g_1(x) \gamma_5 \not{P} + \frac{1}{2} h_1(x) \gamma_5 [\not{S}_T, \not{P}] \right\}$$

Twist-2 collinear PDFs

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Unpolarised PDF



Twist-2 collinear PDFs

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Unpolarised PDF

Helicity distribution

Twist-2 collinear PDFs

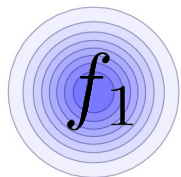
Quark-quark correlator can be decomposed by means of
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Unpolarised PDF

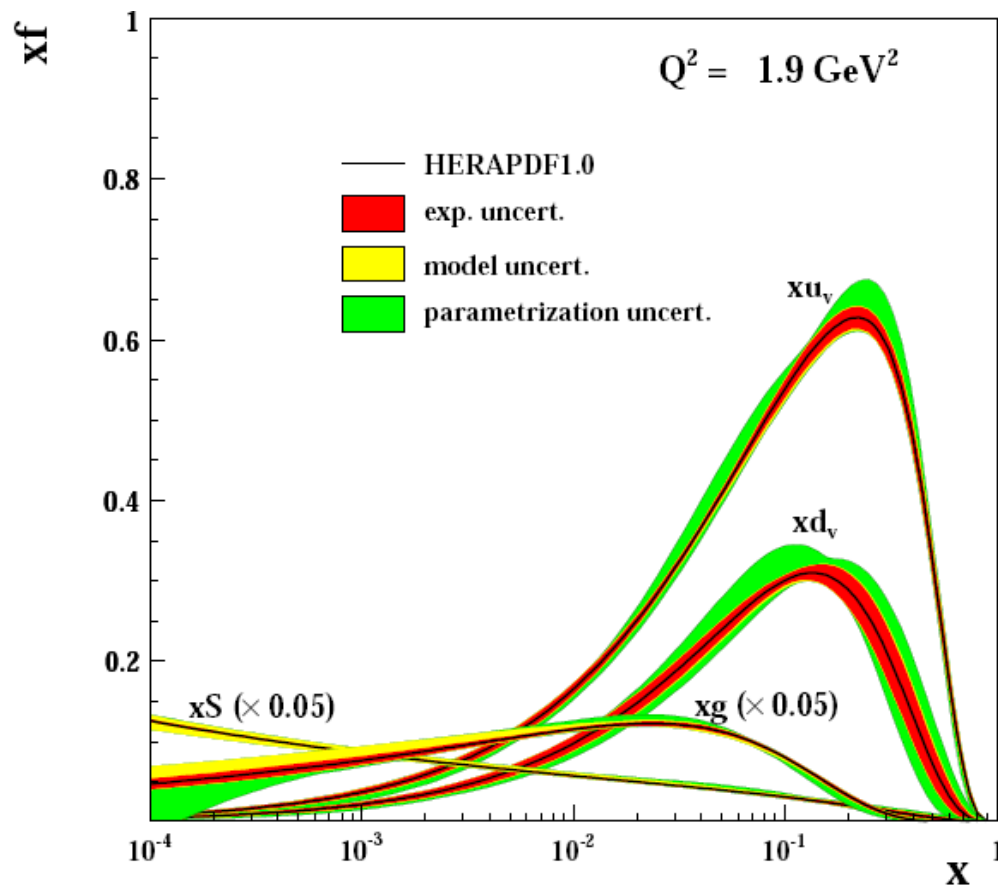
Helicity distribution

Transversity distribution



Unpolarised PDFs

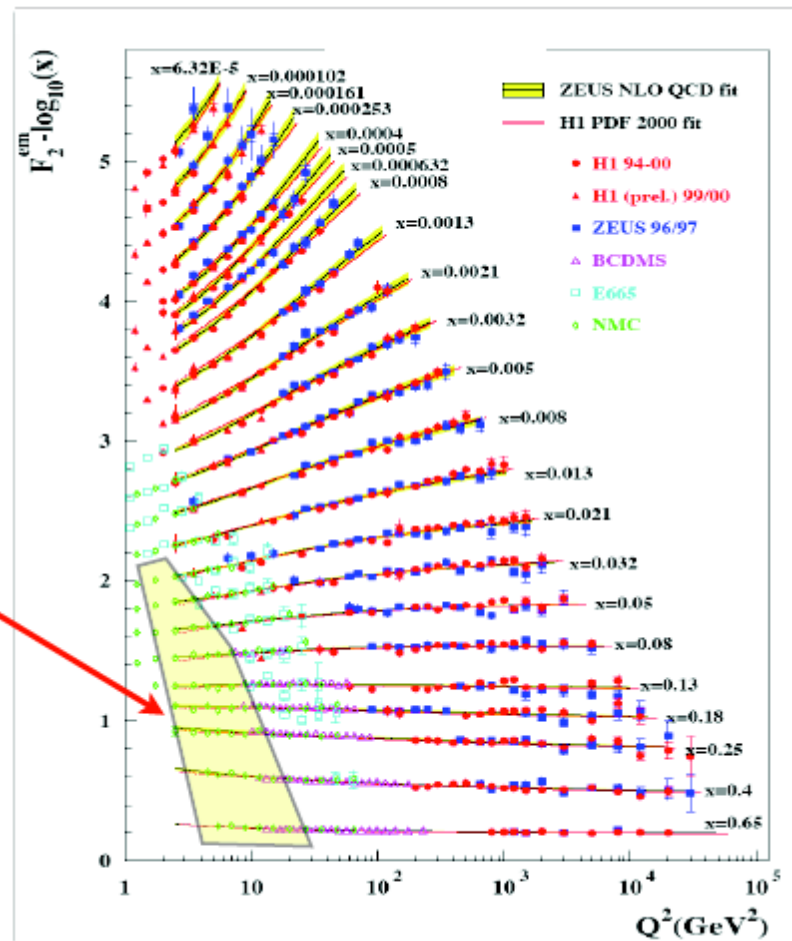
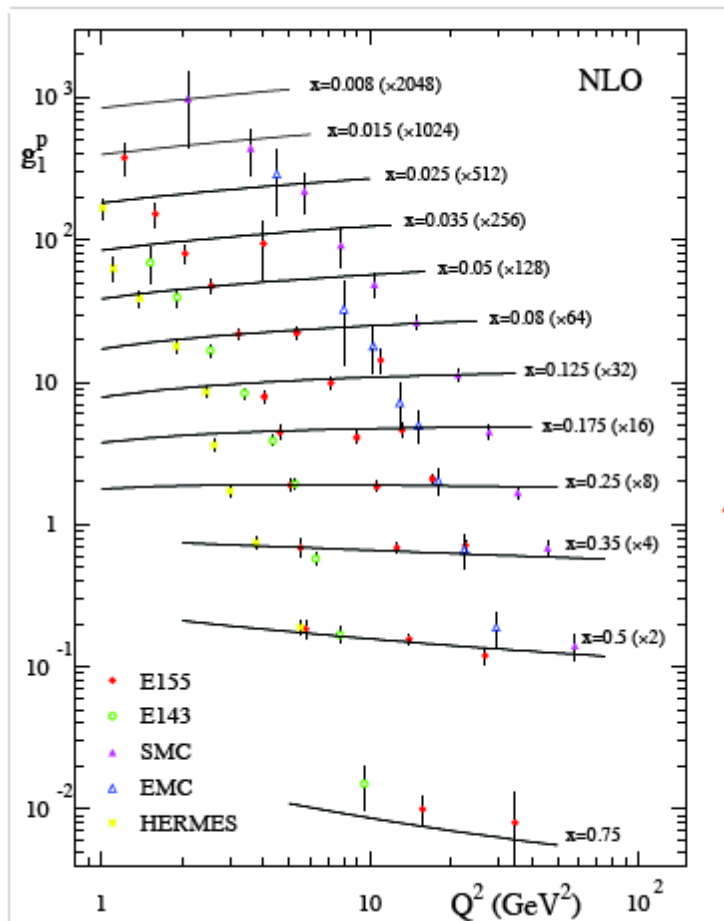
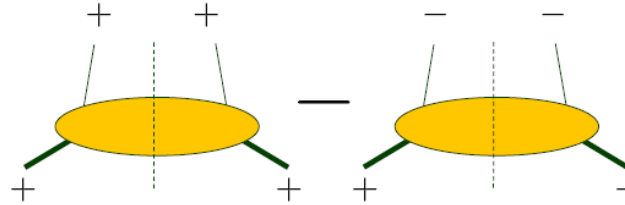
Good knowledge of unpolarised Parton Distribution Functions is acquired

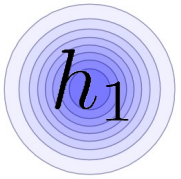


Helicity distributions

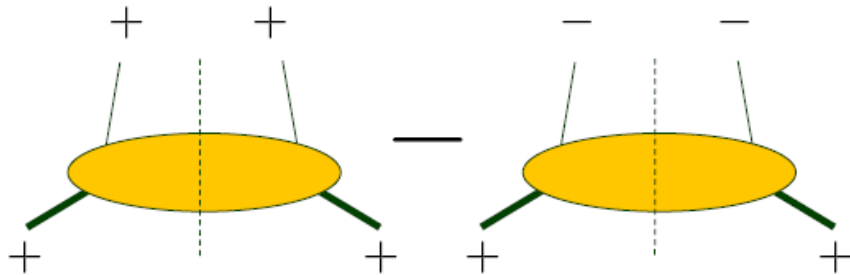


Helicity distributions are relatively well known

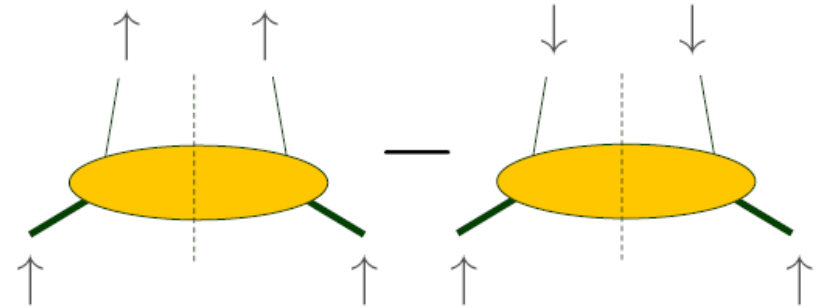




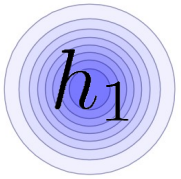
Helicity distribution



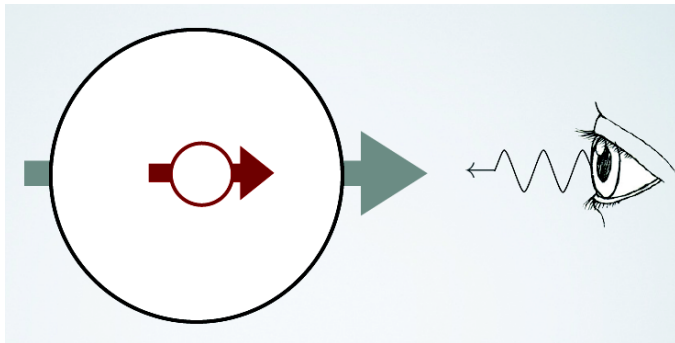
Transversity distribution



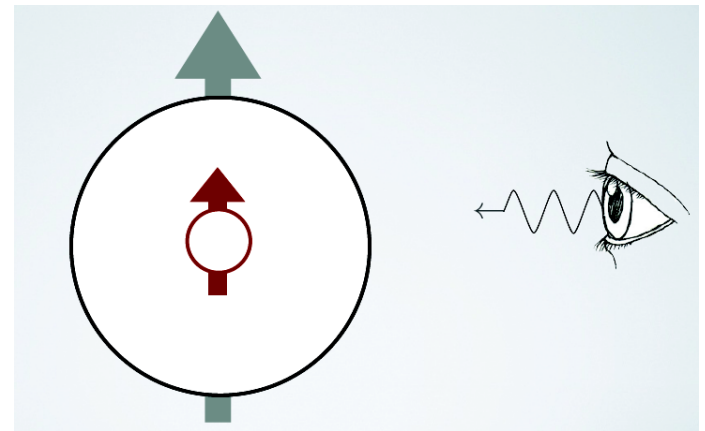
Distribution of transversely polarised quarks inside transversely polarised nucleon



Helicity distribution

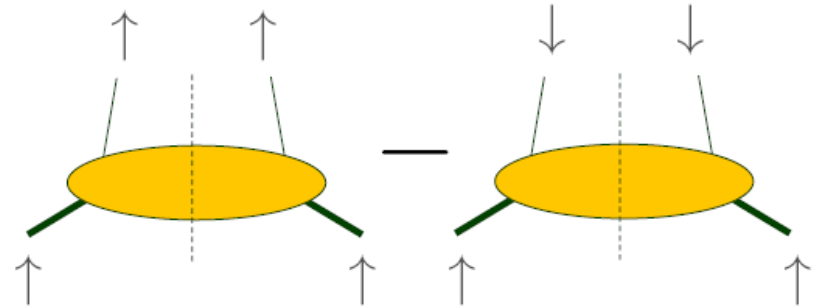


Transversity distribution



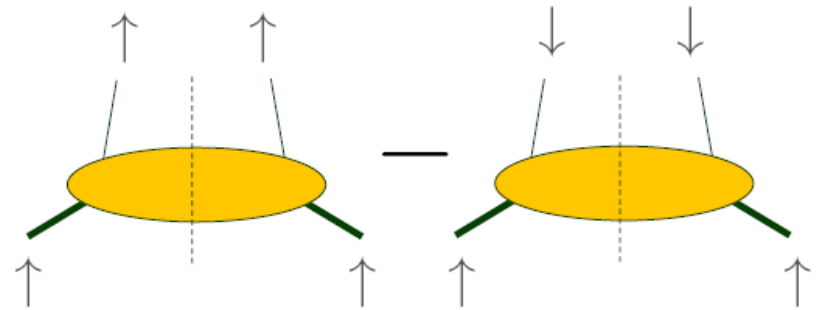
Boost and rotation do not commute \rightarrow helicity and transversity are different!

Transversity distribution



2005: first data on transversity

Transversity distribution

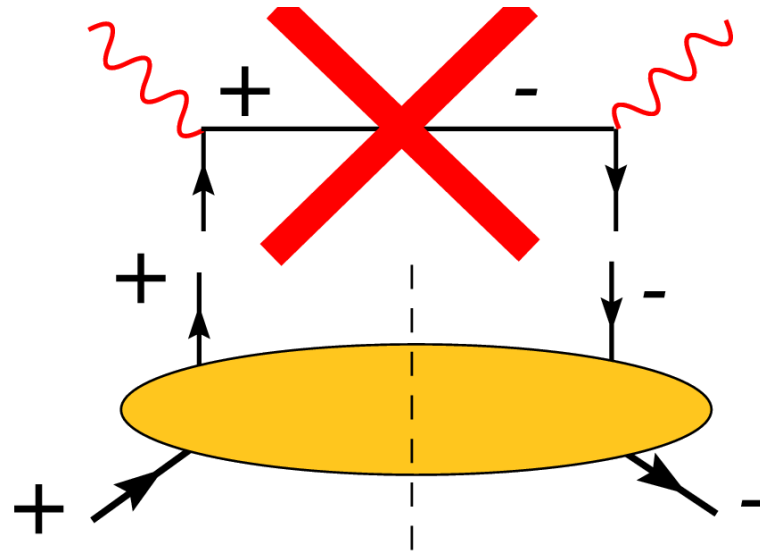


2005: first data on transversity

2012: hundred points from HERMES , COMPASS and JLAB

Transversity distribution

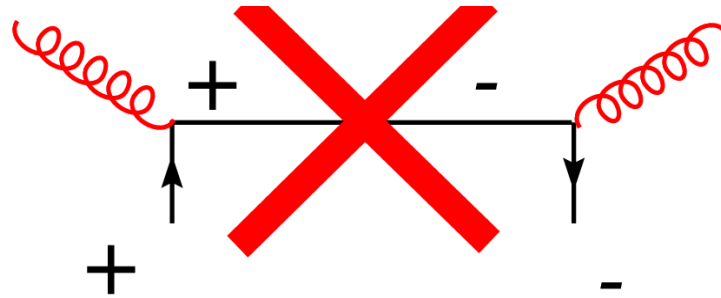
Chiral Odd: it cannot be measured in Deep Inelastic Scattering process



Needs another chiral odd function to be measured

Transversity distribution

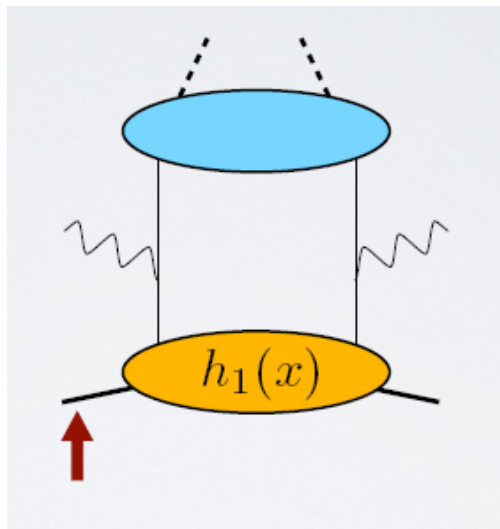
QCD Evolution: no gluon contribution in the evolution



$h_1(x, Q^2)$ is suppressed at low x

JLab 12 is an ideal place to measure transversity → as JLab explores high x region

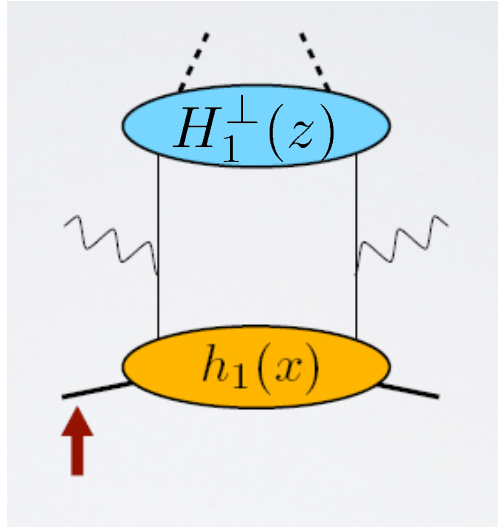
Transversity from SIDIS



First extraction in 2007, [Anselmino et al 07](#)

$$A_{UT}^{\sin(\Phi_h + \Phi_S)} \propto \frac{\sum e_q^2 h_1^q \otimes H_1^{\perp q}}{\sum e_q^2 f_1^q \otimes D_1^q}$$

Transversity from SIDIS

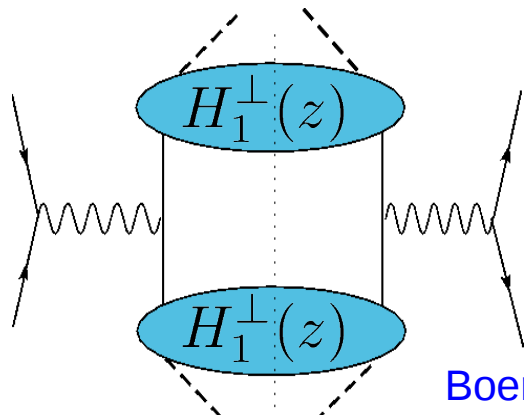


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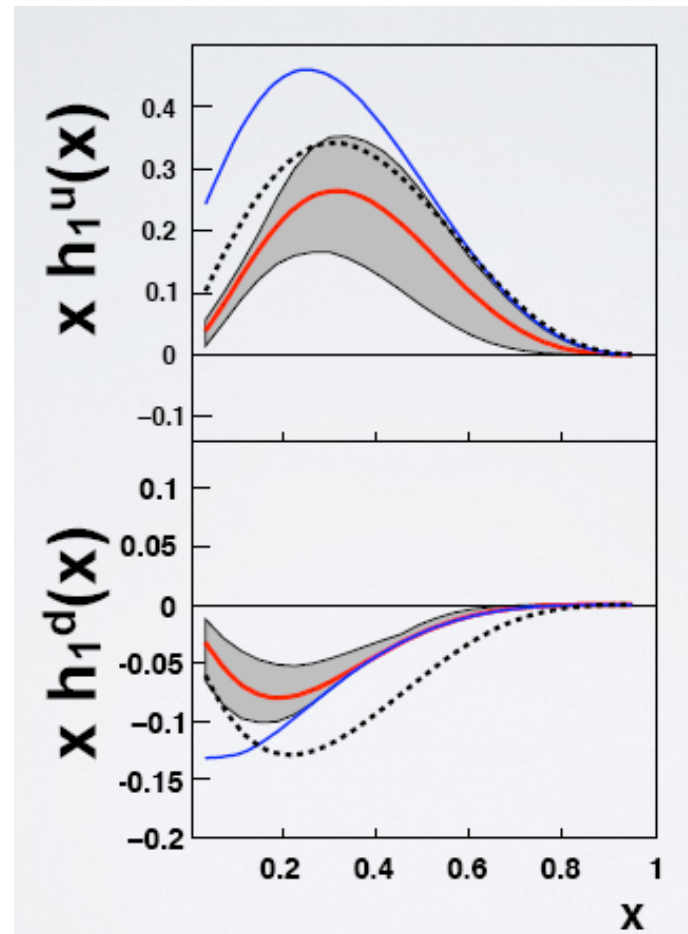
Two unknowns, transversity $h_1(x)$
 Collins Fragmentation Function $H_1^{\perp}(z)$

Information on $H_1^{\perp}(z)$ is available from e^+e^-



Boer (1999)

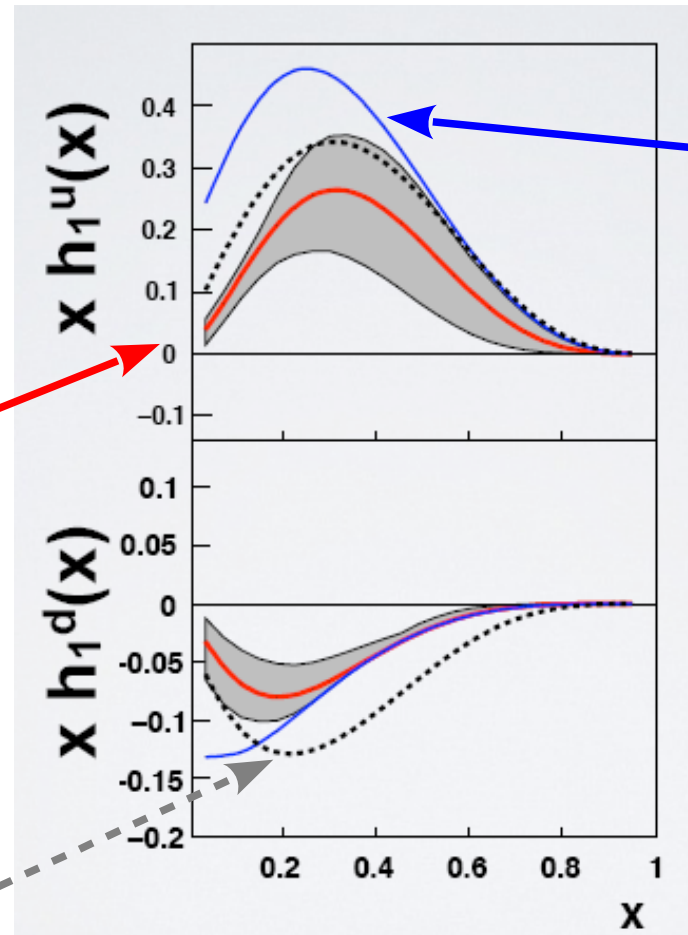
$$A_{e^+e^-} \propto \frac{\sum e_q^2 H_1^{\perp q} \otimes H_1^{\perp \bar{q}}}{\sum e_q^2 D_1^q \otimes D_1^{\bar{q}}}$$



Anselmino et al 09

Transversity

Transversity



Soffer Bound

$$|h_1(x)| \leq \frac{1}{2} (f_1(x) + g_1(x))$$

Valid at LO QCD,
Barone 97, Bourelly et al 98

Valid at NLO QCD, Vogelsang 98

Helicity

Anselmino et al 09

Transversity is the source of information on tensor charge

$$\delta q = \int_0^1 dx (h_1^q(x) - h_1^{\bar{q}}(x))$$

Fundamental quantity

Caveat: no sum rules

First NLL' extraction from the data

$$A^{(1,2)} \quad B^{(1)} \quad C^{(1)}$$

Collins function is related to twist-3 function

$$\tilde{H}_1^{\perp, \alpha}(z_h, b; \mu_b) \sim \left(\frac{-ib^\alpha}{2z_h} \right) H^{(3)}(z_h; \mu_b)$$

We solve also DGLAP equations for transversity and (diagonal) Collins FF

Diagonal part for twist-3 Collins function is:

Yuan-Zhou 2009, Kang 2011

$$\frac{d}{d \ln \mu^2} H_q^{(3)}(z_h, \mu) = \frac{\alpha_s}{2\pi} P_{i \rightarrow q}^H \otimes H_i^{(3)}$$

$$P_{i \rightarrow q}^{H_1}(\hat{x}) = \delta_{iq} C_F \left(\frac{2\hat{z}}{(1-\hat{z})_+} + \frac{3}{2} \delta(1-\hat{z}) \right)$$

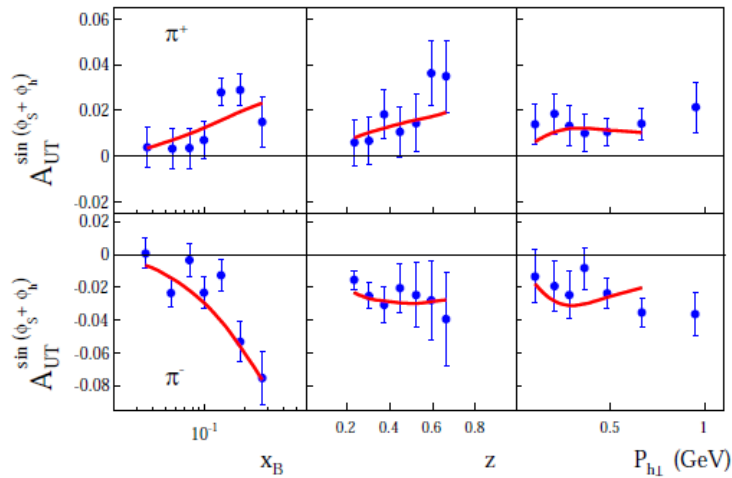
SIDIS data used: HERMES, COMPASS, JLAB – 140 points

e+e- data used: BELLE, BABAR including PT dependence – 122 points

$$\chi^2 / \text{d.o.f.} \simeq 0.88$$

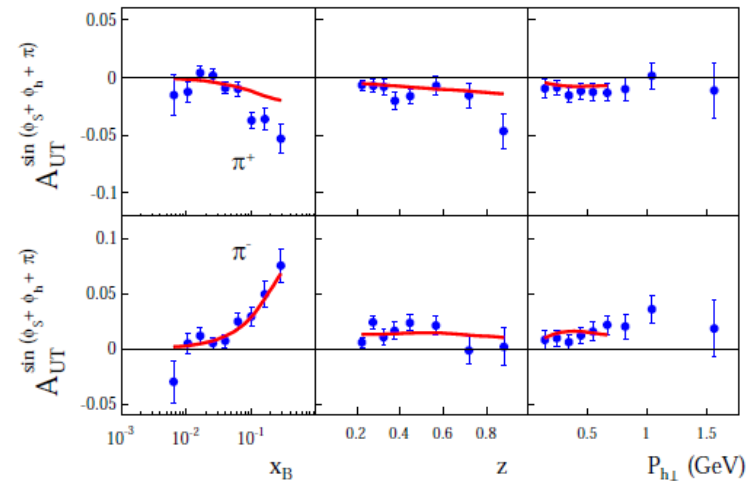
$$\ell P \rightarrow \pi^\pm X$$

HERMES



$$1 \lesssim \langle Q^2 \rangle \lesssim 6 \text{ GeV}^2$$

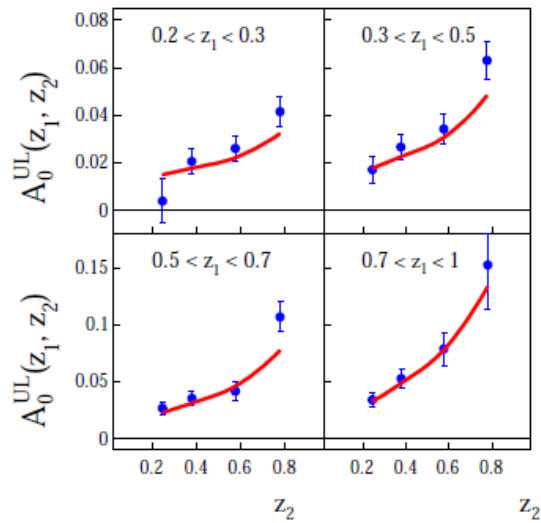
COMPASS



$$1 \lesssim \langle Q^2 \rangle \lesssim 21 \text{ GeV}^2$$

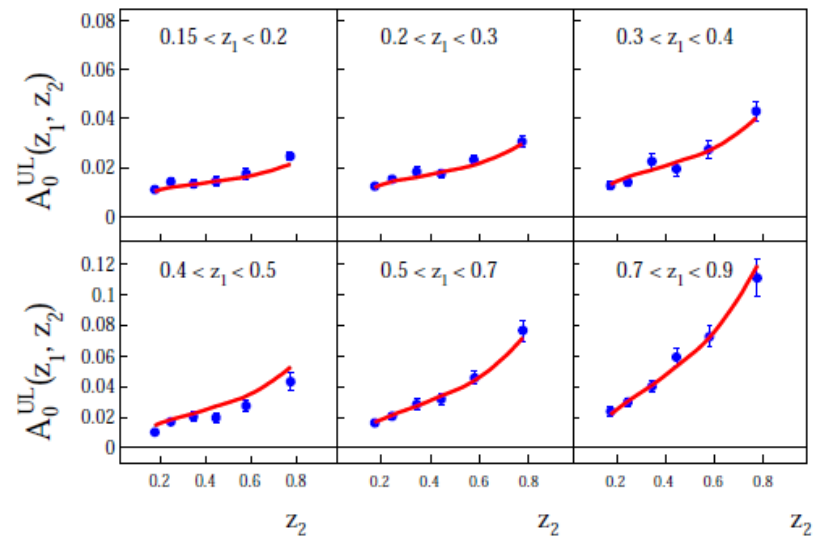
$$e^+e^- \rightarrow \pi\pi X$$

BELLE



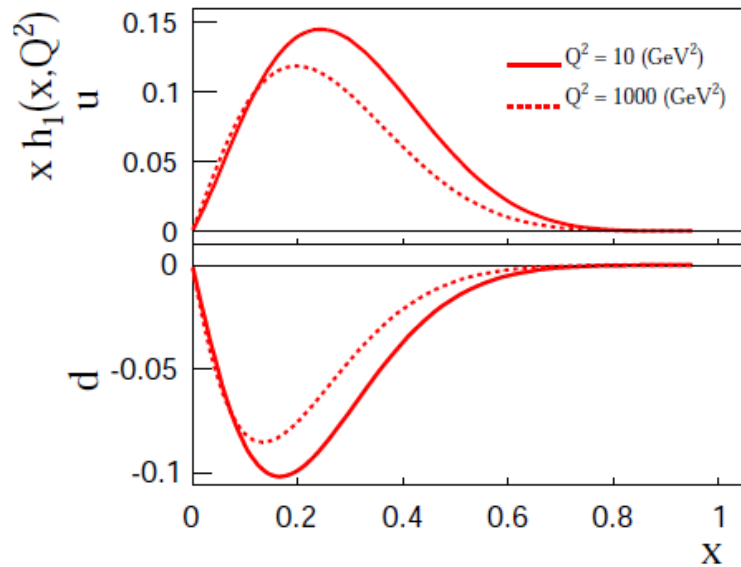
$$Q^2 = 110 \text{ GeV}^2$$

BABAR



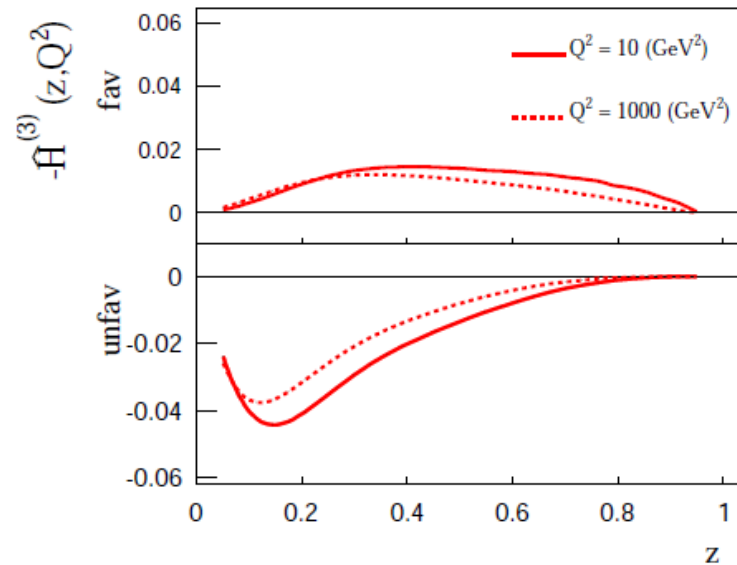
$$Q^2 = 110 \text{ GeV}^2$$

Transversity



Positive u and negative d transversity

Collins



Positive favoured and negative unfavoured Collins FF

Compatible with LO extraction [Anselmino et al 2009](#)

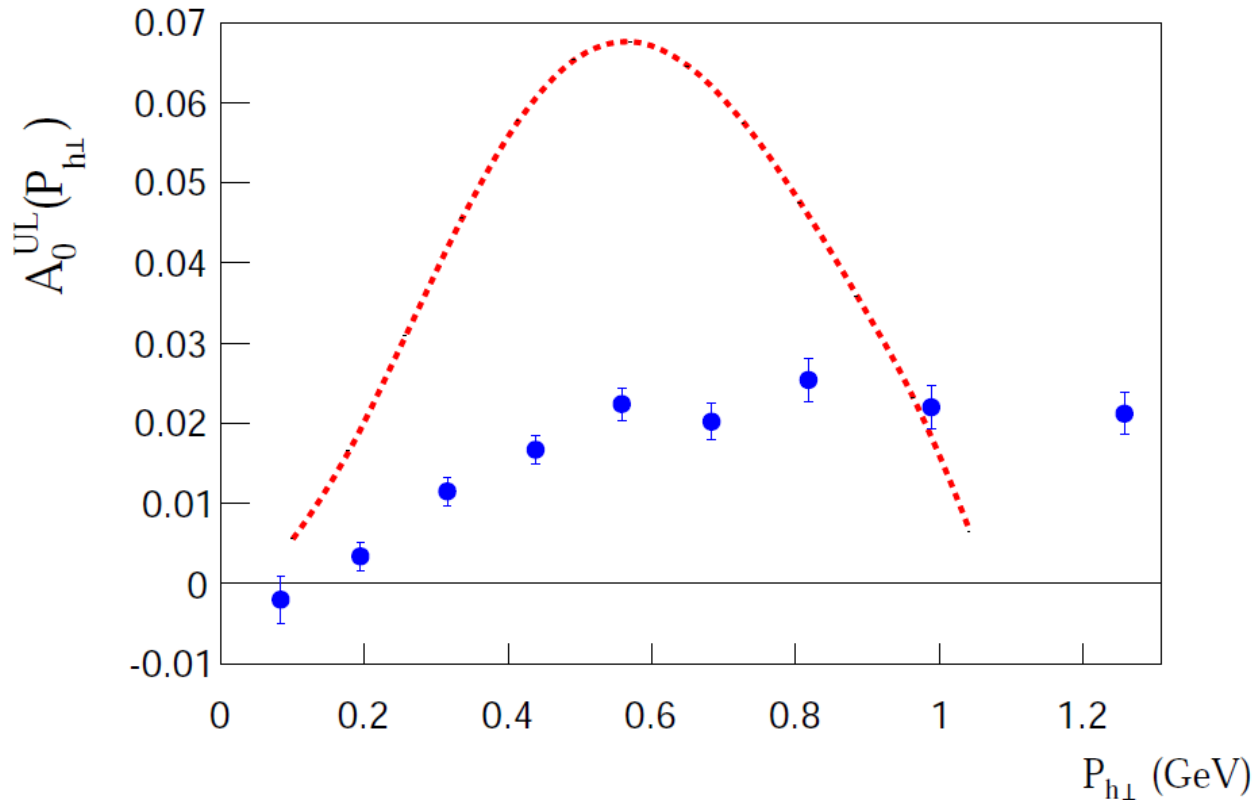
What are evolution effects?

No evolution:

$$e^+e^- \rightarrow \pi\pi X$$



$$Q^2 = 2.4 \text{ GeV}^2$$



What are evolution effects?

$$e^+e^- \rightarrow \pi\pi X$$

No evolution:

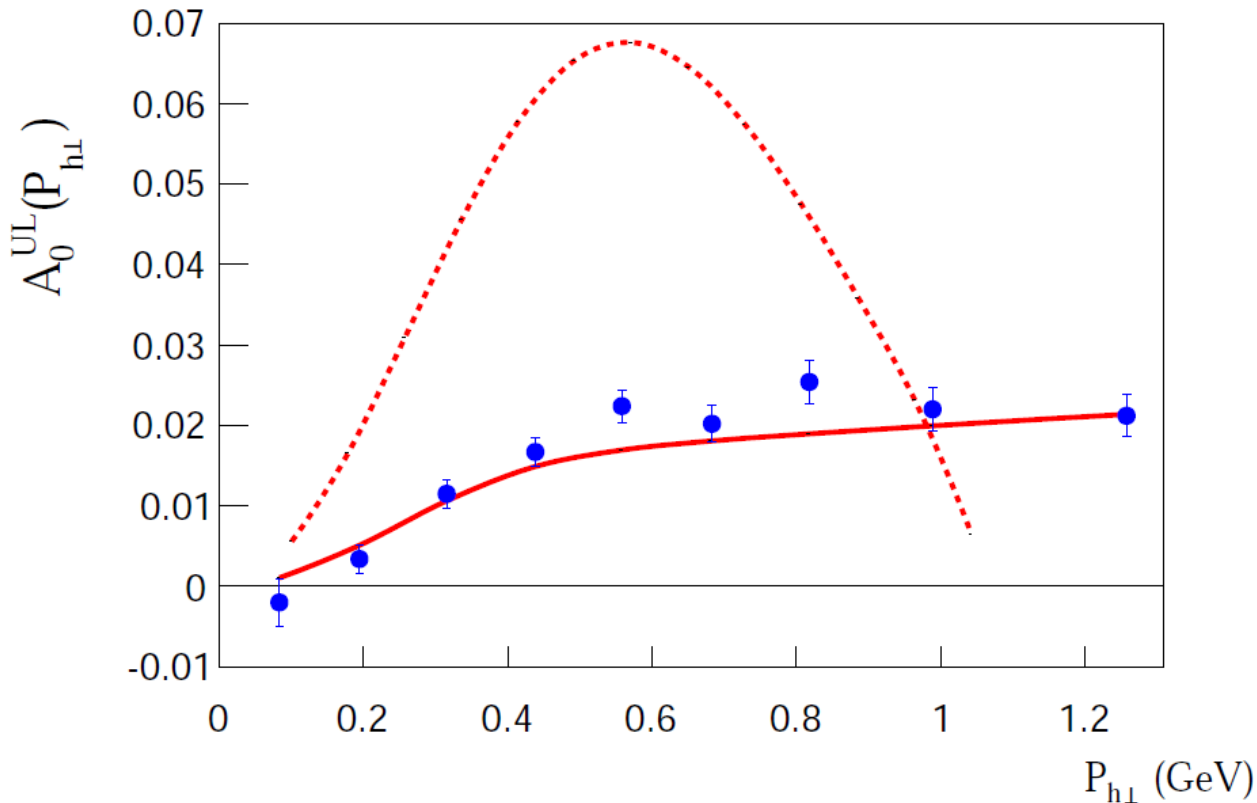


$$Q^2 = 2.4 \text{ GeV}^2$$

NLL' evolution:



$$Q^2 = 110 \text{ GeV}^2$$



$$Q_1^2/Q_2^2 \simeq 50$$

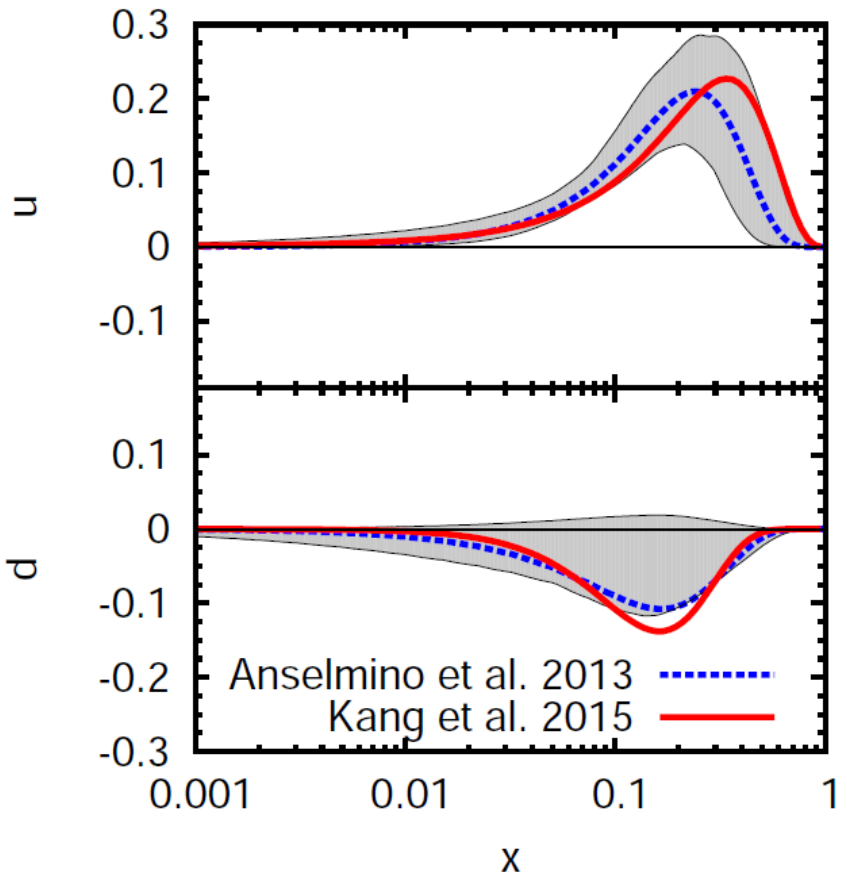
Asymmetry ratio ~ 3.5

Comparison to extractions

Torino 2013

Anselmino et al 2013

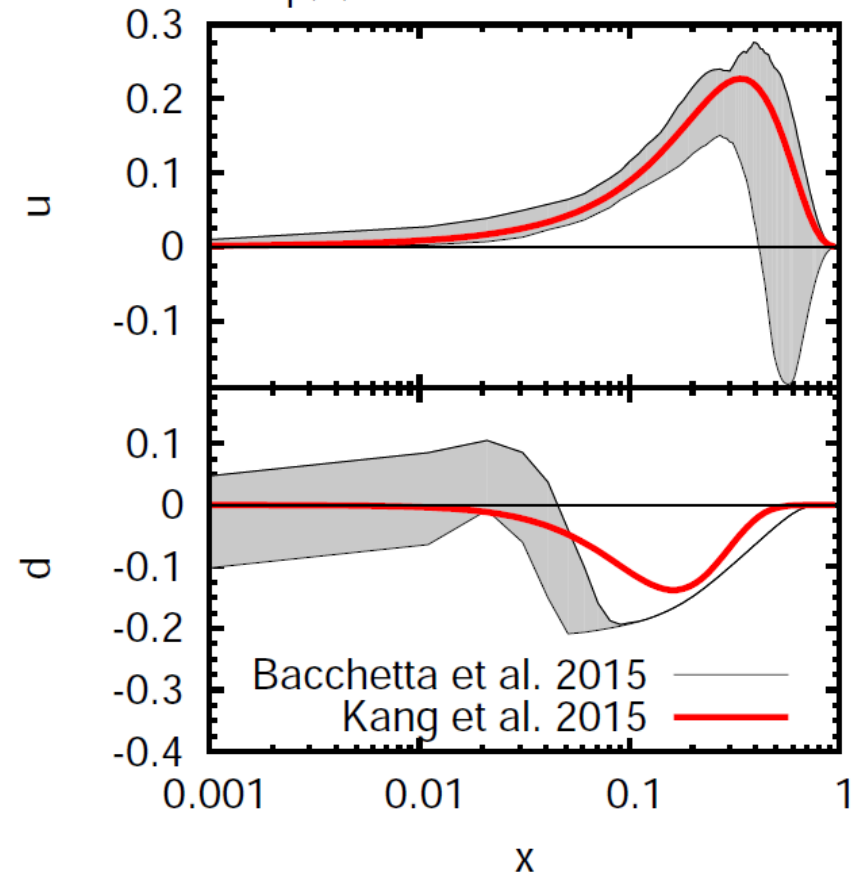
$x h_1(x) \quad Q^2=2.4 \text{ GeV}^2$

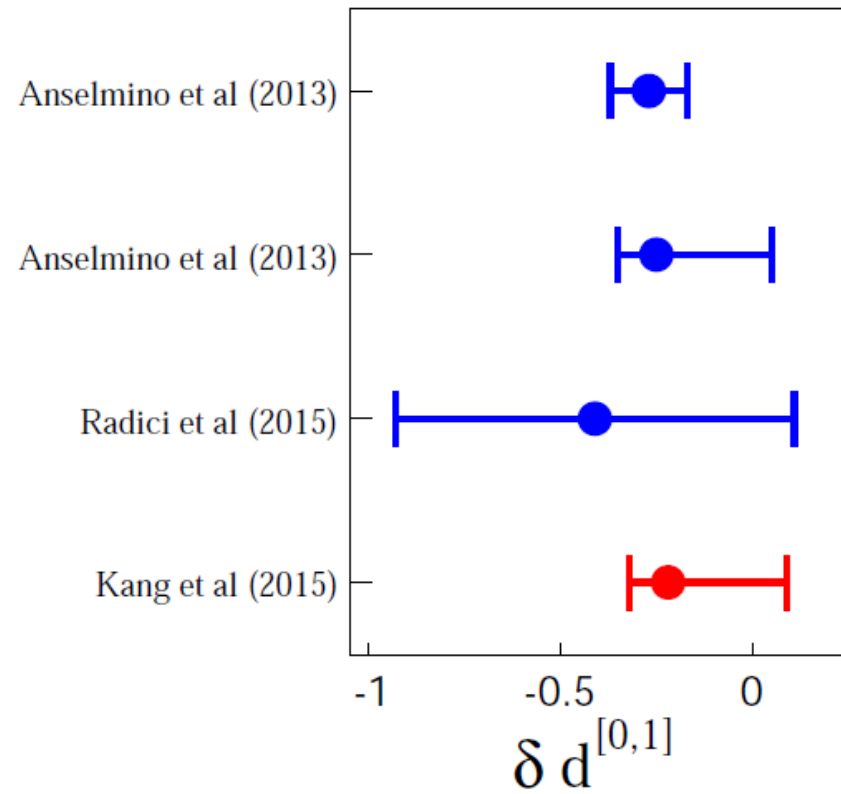
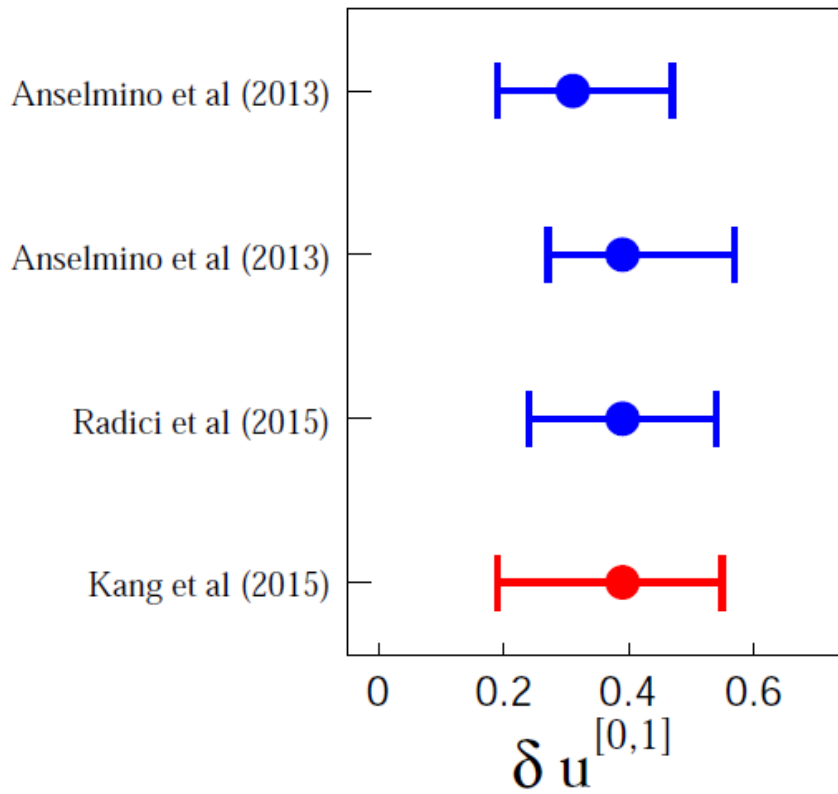


Pavia 2015

Radici et al 2015

$x h_1(x) \quad Q^2=2.4 \text{ GeV}^2$



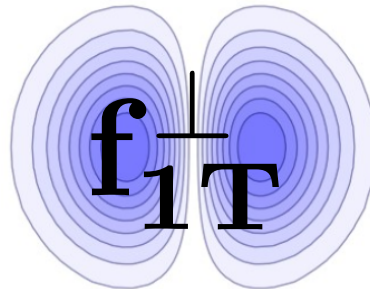


$$\delta u^{[0,1]} = +0.39_{-0.07}^{+0.11}$$

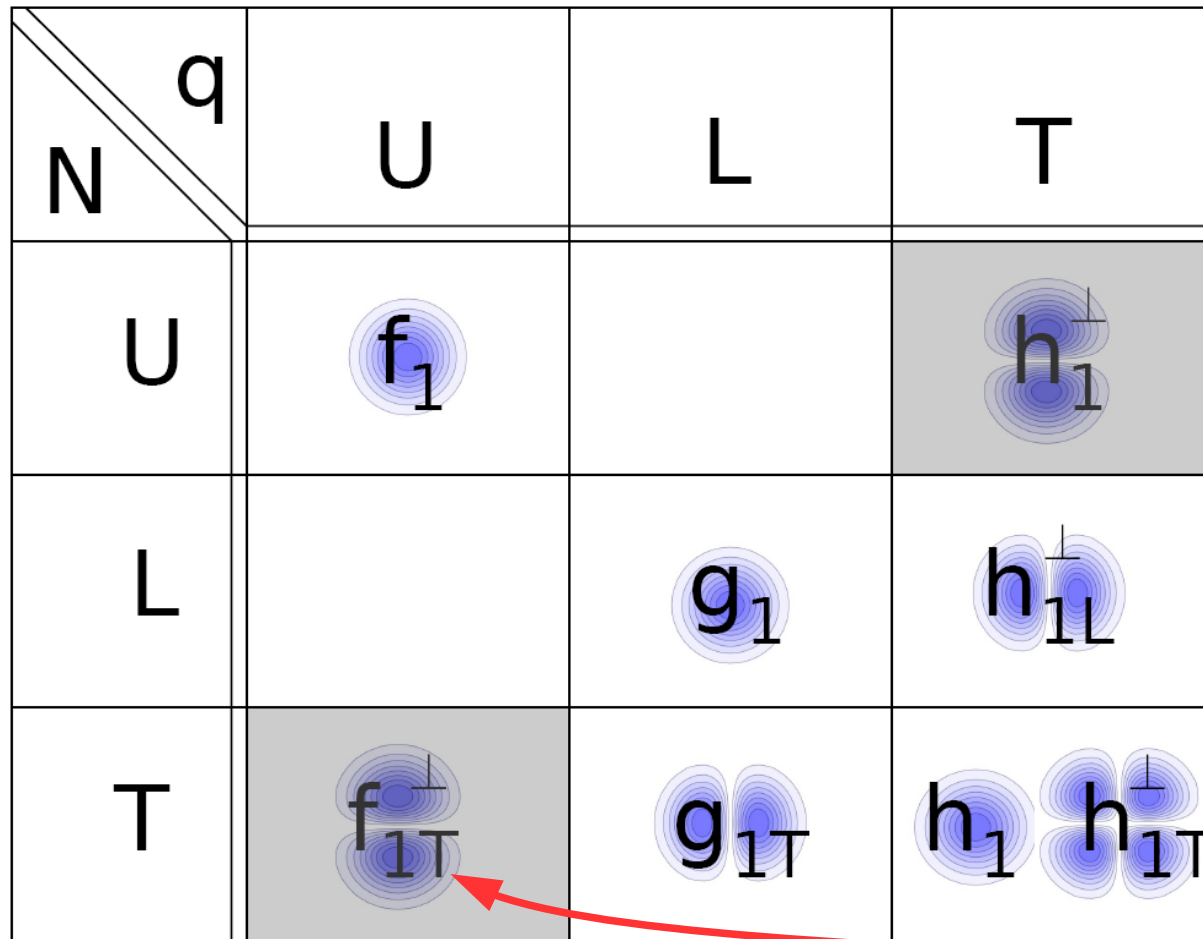
$$\delta d^{[0,1]} = -0.22_{-0.08}^{+0.14}$$

The most accurate modern extraction!

Sivers function



TMD distributions



Sivers function

Kotzinian (1995), Mulders, Tangerman (1995), Boer, Mulders (1998)

$$f(x, \mathbf{k}_T, \mathbf{S}_T) = f_1(x, \mathbf{k}_T^2) - f_{1T}^\perp(x, \mathbf{k}_T^2) \frac{\epsilon_T^{ij} \mathbf{k}_{Ti} \mathbf{S}_{Tj}}{M}$$

Correlation of the **transverse** spin and **motion** of the quarks

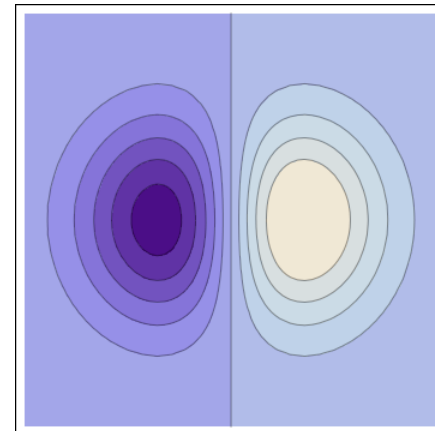
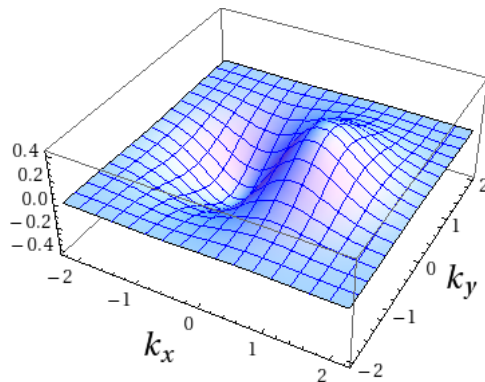
Suppose the spin is along Y direction:

$$\mathbf{S}_T = (0, 1)$$

Deformation in momentum space is:

$$k_x \cdot f(k_x^2 + k_y^2)$$

This is called “dipole” deformation.



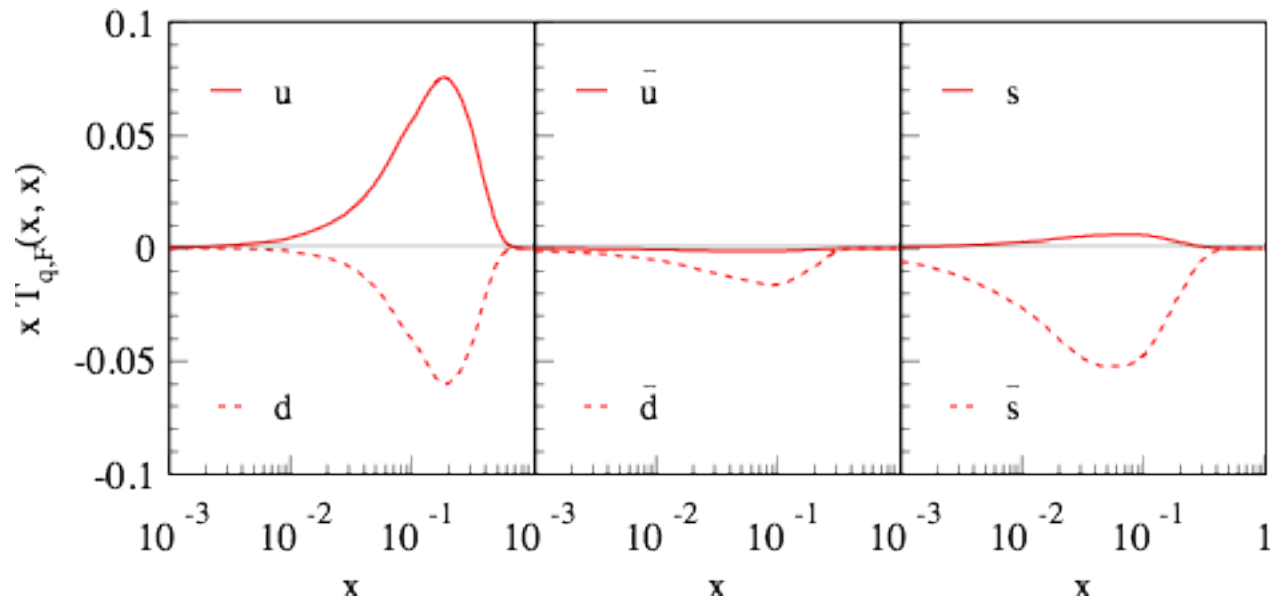
First NLL extraction from the data $A^{(1,2)}$ $B^{(1)}$

Sivers function is related to Qiu-Sterman function [Kang-Xiao-Yuan, 2011](#)
[Aybat-Collins-Rogers-Qiu, 2012](#)

$$\tilde{f}_{1T}^{\perp,\alpha}(x, b; \mu_b) = -\frac{1}{M} \int d^2 k_{\perp} e^{-ik_{\perp} b} k_{\perp}^{\alpha} f_{1T}^{\perp}(x, k_{\perp}; \mu_b)$$

$$\tilde{f}_{1T}^{\perp,\alpha}(x, b; \mu_b) = \left(\frac{-ib^{\alpha}}{2} \right) T(x, x; \mu_b)$$

First NLL extraction from the data



Compatible with previous extractions [Anselmino et al 2007](#)

Only u and d Sivers functions are constrained by existing data

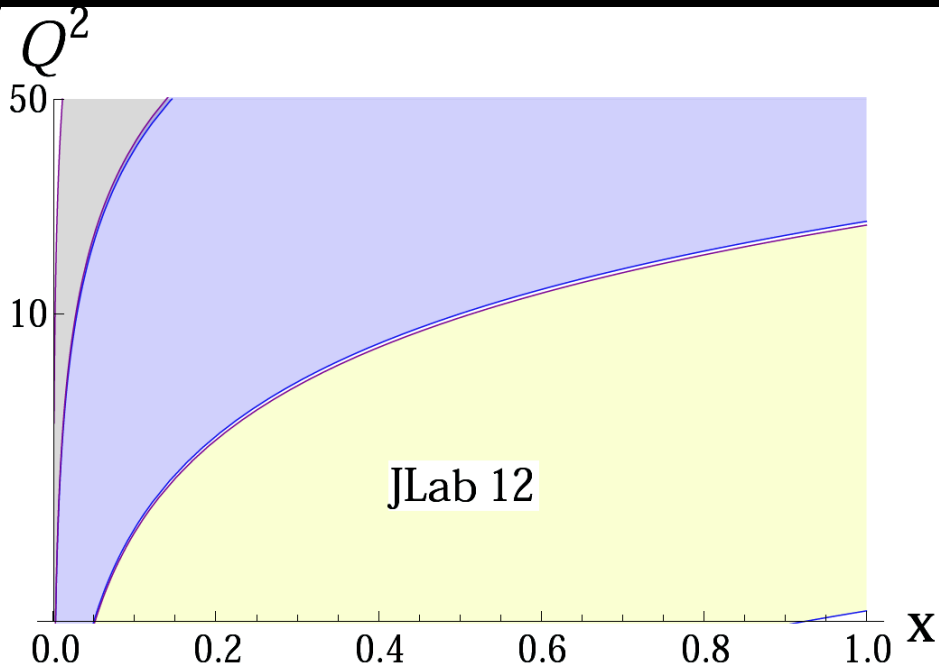
Sea quarks functions are largely unknown – RHIC data and EIC data are needed

Conclusions

~~Conclusions~~

FUTURE

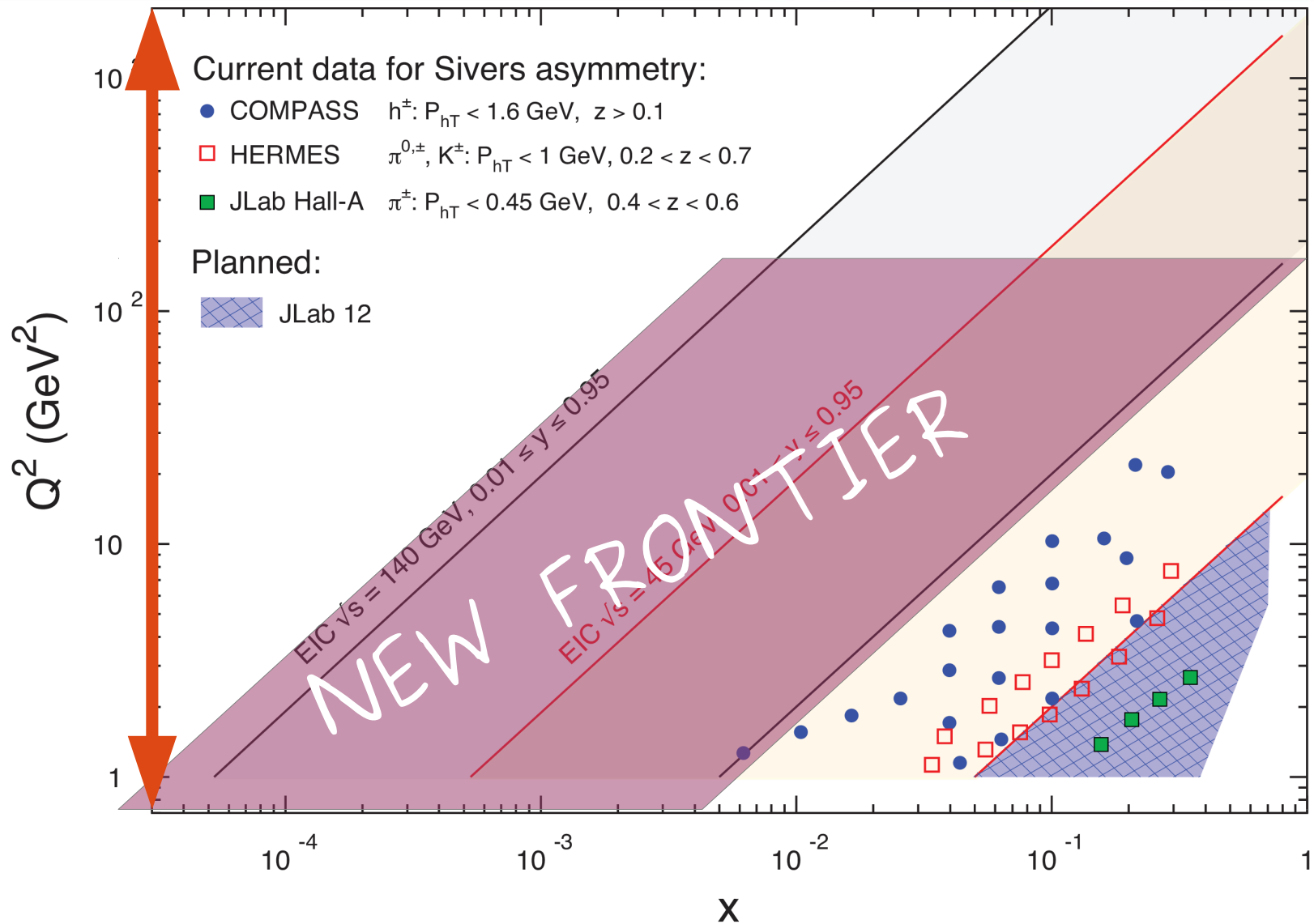
Jefferson Lab 12



Data in large- x region will become available from Jefferson Lab 12

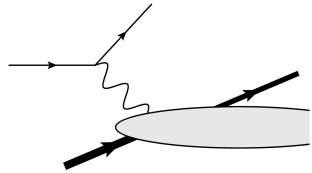
Precision is important – NLL and higher precision in perturbative calculations

Electron Ion Collider



Complementarity of SIDIS, e+e- and Drell-Yan, and hadron-hadron

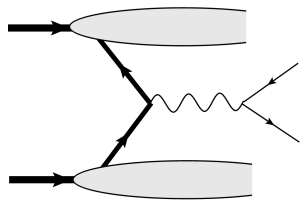
Various processes allow study and test of evolution, universality and extractions of distribution and fragmentation functions. We need information from all of them



$$f(x) \otimes D(z)$$

Semi Inclusive DIS – convolution of distribution functions and fragmentation functions

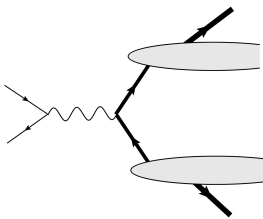
$$l + P \rightarrow l' + h + X$$



$$f(x_1) \otimes f(x_2)$$

Drell-Yan – convolution of distribution functions

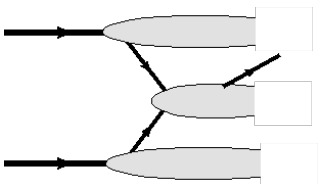
$$P_1 + P_2 \rightarrow \bar{l}l + X$$



$$D(z_1) \otimes D(z_2)$$

e+ e- annihilation – convolution of fragmentation functions

$$\bar{l} + l \rightarrow h_1 + h_2 + X$$



$$f(x_1) \otimes f(x_2) \otimes D(z)$$

Hadron-hadron – convolutions of PDF and fragmentation functions

$$h_1 + h_2 \rightarrow h_3(\gamma, jet, W, \dots) + X$$

Combining measurements from all above is important

THANK YOU!