

Analysis of the Complex NMR Lineshape of the Deuteron

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PSTP

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Outline

- Background
- Functional form for deuteron NMR lineshape
- Complex NMR signals
- Scaling complications
- Enhancement techniques
- Future plans

Background

- Solid Polarized Target Group at UNH
	- developing dynamically polarized target
	- measure spin-structure & tensor spin observables
		- A_{77} , T_{20} and b_{1}

E. Long et al., C-12-15-005 PAC 44 (2016)

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Proton Signals

- Spin polarized protons: low T, high $B \rightarrow$ energy gap between spin states \rightarrow uneven distribution
	- Thermal equilibrium (TE) signal can be enhanced
	- Detectable through nuclear magnetic resonance (NMR)
	- Area under curve \rightarrow % vector polarization (P)

Deuteron Signals

- Deuteron signal (with quadrupole splitting) more complex:
	- One curve for $-1 \rightarrow 0$, another for $0 \rightarrow +1$ sum gives P, difference gives Q
	- Curves span quadrupole angle distrubtion, 0-90° (up to 180° doubles back)
- Ideally, NMR signal would be compared to TE signal and scaled
	- TE integral is physical quantity dictated by B, T
	- At 5T & 1K: P for deuteron $\sim 0.08\%$ on level of noise for our lab

Simulating Spin Flips

- Macro that directly simulates spin flips for a given polarization
- Can use to explore new avenues for increasing Q

Plots from code by Elena Long, drawing on: M. H. Cohen et al., Solid State Physics 5, 321 (1957)

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NMR Analysis

$$
R, A, \eta, \phi
$$
 \longrightarrow compeating variables

$$
Y = \sqrt{3 - \eta \cos(2\phi)}
$$

\n
$$
\rho^2 = \sqrt{A^2 + [1 - \epsilon R - \eta \cos(2\phi)]^2}
$$

\n
$$
\rho = \frac{\omega - \omega_d}{3\omega_q}
$$

\n
$$
\rho = \frac{\omega - \omega_d}{3\omega_q}
$$

\n
$$
-3 \le R \le 3
$$

$$
\hbox{functional form of signal} \quad ,
$$

$$
f_{\epsilon}(R, A, \eta, \phi) = \frac{1}{2\pi\rho} \left\{ 2\cos\left(\frac{\alpha}{2}\right) \left[\arctan\left(\frac{Y^2 - \rho^2}{2Y\rho\sin\left(\frac{\alpha}{2}\right)}\right) + \frac{\pi}{2} \right] \right\}
$$

$$
\epsilon = \pm 1 \qquad \qquad + \sin\left(\frac{\alpha}{2}\right) \ln\left(\frac{Y^2 + \rho^2 + 2Y\rho\cos\left(\frac{\alpha}{2}\right)}{Y^2 + \rho^2 - 2Y\rho\cos\left(\frac{\alpha}{2}\right)} \right) \right\}
$$

$$
\text{phi average } \cdot
$$

$$
F_{\epsilon} \approx \frac{1}{J+1} \sum_{j=0}^{J} \frac{\sqrt{3} f_{\epsilon}(R, A, \eta, \phi_j)}{\sqrt{3 - \eta \cos(2\phi_j)}}
$$

positive $\&$ negative spin flips \bigcup

$$
\chi''(r,R) \propto \frac{1}{\omega_q} \left\{ \left[\frac{r^2 - r^{1-3\theta R}}{r^{1-\theta R}} \right] F_+(R) + \left[\frac{r^{1+3\theta R} - 1}{r^{1+\theta R}} \right] F_-(R) \right\}
$$

$$
\theta = \omega_r/\omega_A
$$

- Without reliable TE, curve fit NMR lineshape to known formula from Dulya, et al.
- At high P, physical parameters can be extracted – B, T not required!
- With parameters, curve fit for *r* across timespan for data taking
	- Naively, ratio of peak heights of signal
	- Instead of <u>area method</u>, *r* is used for ratio method to find both P and Q

$$
\mathcal{P} = \frac{r^2 - 1}{r^2 + r + 1} \quad \text{or} \quad \mathcal{Q} = \frac{r^2 - 2r + 1}{r^2 + r + 1}
$$

 In materials with different bonds to deuterium (e.g. butanol – C and O bonds), set of curves for each bond type

C. Dulya et al., Nuclear Instruments and Methods in Physics Research A 398 (1997)

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NMR Analysis R, A, η, ϕ compacting
variables $Y = \sqrt{3 - \eta \cos(2\phi)}$ $R = \frac{\omega - \omega_d}{3\omega_a}$ $\rho^2 = \sqrt{A^2 + [1 - \epsilon R - n \cos(2\phi)]^2}$ $-3 \leq R \leq 3$ $\cos(\alpha) = \frac{1 - \epsilon R - \eta \cos(2\phi)}{\sigma^2}$ Complete Fit **Negative Spin Flips** 80 Positive Spin Flips functional form of signal χ'' $f_{\epsilon}(R,A,\eta,\phi) = \frac{1}{2\pi\rho} \left\{ 2cos(\frac{\alpha}{2}) \left[arctan\left(\frac{Y^2-\rho^2}{2Y\rho sin(\frac{\alpha}{2})}\right) + \frac{\pi}{2} \right] \right\}$ 60 Signal (a.u.) $+ \sin(\frac{\alpha}{2}) \ln\left(\frac{Y^2+\rho^2+2Y\rho\cos(\frac{\alpha}{2})}{Y^2+\rho^2-2Y\rho\cos(\frac{\alpha}{2})}\right)\}$ 40 F_{+1} phi average **U** $\mathsf{F}_{\mathcal{A}}$ 20 $F_{\epsilon} \approx \frac{1}{J+1} \sum_{i=0}^{J} \frac{\sqrt{3} f_{\epsilon}(R, A, \eta, \phi_j)}{\sqrt{3 - \eta cos(2\phi_i)}}$ \circ positive $\&$ negative spin flips \bigcup -1 -3 -2 0 $\mathbf{1}$ 2 3 R $\chi''(r,R) \propto \frac{1}{\omega_c} \left\{ \left[\frac{r^2 - r^{1-3\theta R}}{r^{1-\theta R}} \right] F_+(R) + \left[\frac{r^{1+3\theta R} - 1}{r^{1+\theta R}} \right] F_-(R) \right\}$ $\theta = \omega_q/\omega_d$

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Data Analysis

- Passing deuteron signal data through curve fitting routine based on Dulya yields good fit
	- Component spin-flip curves can be reconstructed from curve fitting parameters
	- Can perform both area and ratio methods!
- Passing simulated data through routine extracts P that agrees with input!

Real & Imaginary Analysis

- Issue with Dulya "false asymmetry" seems artificial why not just fit both real and imaginary parts of signal?
- Two different parts to signal absorptive and dispersive rotate between through phase angle
	- $-$ Real = χ'' cosθ χ' sinθ
	- $-$ Imag. = χ'' sin θ + χ' cos θ
- Changes for dispersive: functional form of signal \bigcup

$$
f_{\epsilon}(R, A, \eta, \phi) = \frac{1}{2\pi\rho} \left\{ 2\sin(\frac{\alpha}{2}) \left[arctan\left(\frac{Y^2 - \rho^2}{2Y\rho sin(\frac{\alpha}{2})}\right) + \frac{\pi}{2} \right] \right\}
$$

$$
\epsilon = \pm 1 \qquad \qquad + \cos(\frac{\alpha}{2}) \ln\left(\frac{Y^2 + \rho^2 + 2Y\rho cos(\frac{\alpha}{2})}{Y^2 + \rho^2 - 2Y\rho cos(\frac{\alpha}{2})}\right) \}
$$

phi average
$$
\sqrt{}
$$

$$
F_{\epsilon} \approx \frac{\mathbf{g}}{J+1} \sum_{j=0}^{J} \frac{\sqrt{3} f_{\epsilon}(R,A,\eta,\phi_j)}{\sqrt{3-\eta cos(2\phi_j)}}
$$

Real & Imaginary Analysis

 Preliminary agreement between P and Q taken from real and imaginary signals!

Fully Imaginary Analysis

Ran full spin-up curve entirely in the imaginary (88.5° phase!)

Peaked at P = 37% , Q = 10.5%

Why Imaginary?

Complex is more accurate – tuning is difficult!

Real area ∝ P, simulations suggest Imaginary area ∝ Q

Very first purely imaginary spin-up!

Keeping Accuracy at High P

Dulya function is missing a final outer parameter: Ξ, which varies with P, more noticeably at high P

Seems to be quartic-ly dependent on r – more analysis needed

Ξ assumed constant: Ξ allowed to float:

Quartic Relationship: Ξ v P

 $E \propto -0.38$ P^4 – 0.61 P^2 + 1

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Increasing Q

- Increasing P via mm-wave enhancement naturally gives Q
- To increase further, can use hole-burning: separate solenoid fires at specific frequency to drive spin flips – increase area difference
	- Lose some P but for A_{77} , Q is more important
- System built & signal found by Nathaly Santiesteban/David Ruth!

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Increasing Q

- Solid ND_3 crystal instead of sine distribution of quadrupole angles (powder pattern), only 3 possible angles
- SUPER preliminary simulations only, no tests so far

Future Goals

- Reach Q of ~30% in deuterated ammonia
	- Continue with hole-burning
	- Use EPR to focus mm-wave frequencies
- Add hole-burning to curve-fitting method in progress!
- Confirm relationship between Q and Imaginary area
- Refine functional form of *Ξ*, find physical cause for it

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Deuteron-Electron Energy Levels

