



Analysis of the Complex NMR Lineshape of the Deuteron

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20TH INTERNATIONAL WORKSHOP ON POLARIZED SOURCES, TARGETS, AND POLARIMETRY

SEPT. 22-27 | JEFFERSON LAB, NEWPORT NEWS, VA

Outline

- Background
- Functional form for deuteron NMR lineshape
- Complex NMR signals
- Scaling complications
- Enhancement techniques
- Future plans



Background

- Solid Polarized Target Group at UNH
 - developing dynamically polarized target
 - measure spin-structure & tensor spin observables
 - A_{ZZ} , T_{20} and b_1



E. Long et al., C-12-15-005 PAC 44 (2016)

Proton Signals

- Spin polarized protons: low T, high B → energy gap between spin states → uneven distribution
 - Thermal equilibrium (TE) signal can be enhanced
 - Detectable through nuclear magnetic resonance (NMR)
 - Area under curve \rightarrow % vector polarization (P)



09/24/24

Deuteron Signals

- Deuteron signal (with quadrupole splitting) more complex:
 - One curve for $-1 \rightarrow 0$, another for $0 \rightarrow +1$ sum gives P, difference gives Q
 - Curves span quadrupole angle distrubtion, 0-90° (up to 180° doubles back)
- Ideally, NMR signal would be compared to TE signal and scaled
 - TE integral is physical quantity dictated by B, T
 - At 5T & 1K: P for deuteron ~ 0.08% on level of noise for our lab



Simulating Spin Flips

- Macro that directly simulates spin flips for a given polarization
- Can use to explore new avenues for increasing Q



Plots from code by Elena Long, drawing on: M. H. Cohen et al., Solid State Physics 5, 321 (1957)

NMR Analysis

$$R, A, \eta, \phi$$
 \swarrow compacting variables

$$Y = \sqrt{3 - \eta \cos(2\phi)} \qquad \qquad R = \frac{\omega - \omega_d}{3\omega_q}$$

$$\rho^2 = \sqrt{A^2 + [1 - \epsilon R - \eta \cos(2\phi)]^2} \qquad \qquad -3 \le R \le 3$$

$$\cos(\alpha) = \frac{1 - \epsilon R - \eta \cos(2\phi)}{\rho^2}$$

$$f_{\epsilon}(R, A, \eta, \phi) = \frac{1}{2\pi\rho} \{ 2\cos(\frac{\alpha}{2}) \left[\arctan\left(\frac{Y^2 - \rho^2}{2Y\rho sin(\frac{\alpha}{2})}\right) + \frac{\pi}{2} \right]$$
$$\epsilon = \pm 1 \qquad + \sin(\frac{\alpha}{2}) ln\left(\frac{Y^2 + \rho^2 + 2Y\rho cos(\frac{\alpha}{2})}{Y^2 + \rho^2 - 2Y\rho cos(\frac{\alpha}{2})}\right) \}$$

$$F_{\epsilon} \approx \frac{1}{J+1} \sum_{j=0}^{J} \frac{\sqrt{3}f_{\epsilon}(R, A, \eta, \phi_j)}{\sqrt{3 - \eta \cos(2\phi_j)}}$$

positive & negative spin flips 🗸

$$\chi''(r,R) \propto \frac{1}{\omega_q} \left\{ \left[\frac{r^2 - r^{1-3\theta R}}{r^{1-\theta R}} \right] F_+(R) + \left[\frac{r^{1+3\theta R} - 1}{r^{1+\theta R}} \right] F_-(R) \right\}$$
$$\theta = \omega_-/\omega_-$$

- Without reliable TE, curve fit NMR lineshape to known formula from Dulya, et al.
- At high P, physical parameters can be extracted B, T not required!
- With parameters, curve fit for *r* across timespan for data taking
 - Naively, ratio of peak heights of signal
 - Instead of <u>area method</u>, *r* is used for <u>ratio</u> <u>method</u> to find both P and Q

$$\mathcal{P} = \frac{r^2 - 1}{r^2 + r + 1}$$
 or $\mathcal{Q} = \frac{r^2 - 2r + 1}{r^2 + r + 1}$

In materials with different bonds to deuterium (e.g. butanol – C and O bonds), set of curves for each bond type

C. Dulya et al., Nuclear Instruments and Methods in Physics Research A 398 (1997)

NMR Analysis R, A, η, ϕ compacting variables $Y = \sqrt{3 - \eta \cos(2\phi)}$ $R = \frac{\omega - \omega_d}{3\omega_a}$ $\rho^2 = \sqrt{A^2 + [1 - \epsilon R - \eta \cos(2\phi)]^2}$ $-3 \le R \le 3$ $\cos(\alpha) = \frac{1 - \epsilon R - \eta \cos(2\phi)}{\rho^2}$ Complete Fit Negative Spin Flips 80 Positive Spin Flips functional form of signal $f_{\epsilon}(R, A, \eta, \phi) = \frac{1}{2\pi\rho} \left\{ 2\cos(\frac{\alpha}{2}) \left[\arctan\left(\frac{Y^2 - \rho^2}{2Y\rho\sin(\frac{\alpha}{2})}\right) + \frac{\pi}{2} \right] \right\}$ 60 Signal (a.u.) $+\sin(\frac{\alpha}{2})ln\left(\frac{Y^2+\rho^2+2Y\rho cos(\frac{\alpha}{2})}{Y^2+\rho^2-2Y\rho cos(\frac{\alpha}{2})}\right)\}$ 40 **F**₊₁ phi average 🞝 20 $F_{\epsilon} \approx \frac{1}{J+1} \sum_{i=0}^{J} \frac{\sqrt{3} f_{\epsilon}(R, A, \eta, \phi_j)}{\sqrt{3 - \eta \cos(2\phi_j)}}$ 0 positive & negative spin flips -3 -2 $^{-1}$ 1 2 0 3 R $\chi''(r,R) \propto \frac{1}{\omega_{\sigma}} \left\{ \left[\frac{r^2 - r^{1-3\theta R}}{r^{1-\theta R}} \right] F_{+}(R) + \left[\frac{r^{1+3\theta R} - 1}{r^{1+\theta R}} \right] F_{-}(R) \right\}$ $\theta = \omega_q / \omega_d$

C. Dulya et al., Nuclear Instruments and Methods in Physics Research A 398 (1997) 09/24/24 Analysis of the Complex NMR Lineshape of the Deuteron

Data Analysis

- Passing deuteron signal data through curve fitting routine based on Dulya yields good fit
 - Component spin-flip curves can be reconstructed from curve fitting parameters
 - Can perform both area and ratio methods!
- Passing simulated data through routine extracts P that agrees with input!



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Real & Imaginary Analysis

- Issue with Dulya "false asymmetry" seems artificial why not just fit both real and imaginary parts of signal?
- Two different parts to signal absorptive and dispersive rotate between through phase angle
 - Real = $\chi'' \cos\theta \chi' \sin\theta$
 - Imag. = $\chi'' \sin\theta + \chi' \cos\theta$
- Changes for dispersive: functional form of signal

$$\begin{split} f_{\epsilon}(R,A,\eta,\phi) &= \frac{1}{2\pi\rho} \{ 2 \sin(\frac{\alpha}{2}) \left[\arctan\left(\frac{Y^2 - \rho^2}{2Y\rho sin(\frac{\alpha}{2})}\right) + \frac{\pi}{2} \right] \\ &\epsilon = \pm 1 \\ &+ \cos(\frac{\alpha}{2}) ln\left(\frac{Y^2 + \rho^2 + 2Y\rho cos(\frac{\alpha}{2})}{Y^2 + \rho^2 - 2Y\rho cos(\frac{\alpha}{2})}\right) \} \end{split}$$

$$F_{\epsilon} \approx \frac{\mathbf{E}}{J+1} \sum_{j=0}^{J} \frac{\sqrt{3}f_{\epsilon}(R, A, \eta, \phi_j)}{\sqrt{3 - \eta cos(2\phi_j)}}$$



Real & Imaginary Analysis

 Preliminary agreement between P and Q taken from real and imaginary signals!



Fully Imaginary Analysis

Ran full spin-up curve entirely in the imaginary (88.5° phase!)

• Peaked at P = 37%, Q = 10.5%



Why Imaginary?

Complex is more accurate – tuning is difficult!

Real area \propto P, simulations suggest Imaginary area \propto Q

Very first purely imaginary spin-up!



Keeping Accuracy at High P

Dulya function is missing a final outer parameter: Ξ , which varies with P, more noticeably at high P

Seems to be quartic-ly dependent on r – more analysis needed

E assumed constant:

Ξ allowed to float:



Quartic Relationship: E v P

 $\Xi \propto -0.38 P^4 - 0.61 P^2 + 1$



Increasing Q

- Increasing P via mm-wave enhancement naturally gives Q
- To increase further, can use hole-burning: separate solenoid fires at specific frequency to drive spin flips increase area difference
 - Lose some P but for A_{zz}, Q is more important
- System built & signal found by Nathaly Santiesteban/David Ruth!



Increasing Q

- Solid ND₃ crystal instead of sine distribution of quadrupole angles (powder pattern), only 3 possible angles
- SUPER preliminary simulations only, no tests so far



Future Goals

- Reach Q of ~30% in deuterated ammonia
 - Continue with hole-burning
 - Use EPR to focus mm-wave frequencies
- Add hole-burning to curve-fitting method in progress!
- Confirm relationship between Q and Imaginary area
- Refine functional form of Ξ , find physical cause for it

Acknowledgments

UNH DNP PIs:

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- Elena Long
- Nathaly Santiesteban
- Karl Slifer
- UNH DNP collaborators:
 - Anchit Arora
 - Hector Chinchay
 - Muhammad Farooq
 - Chhetra Lama
 - Olaiya Olokunboyo
 - Eli Phippard
 - David Ruth
 - Zoe Wolters
 - Allison Zec
- This work is supported by the Dept. of Energy Grant DE-FG02-88ER40410.





















Deuteron-Electron Energy Levels

