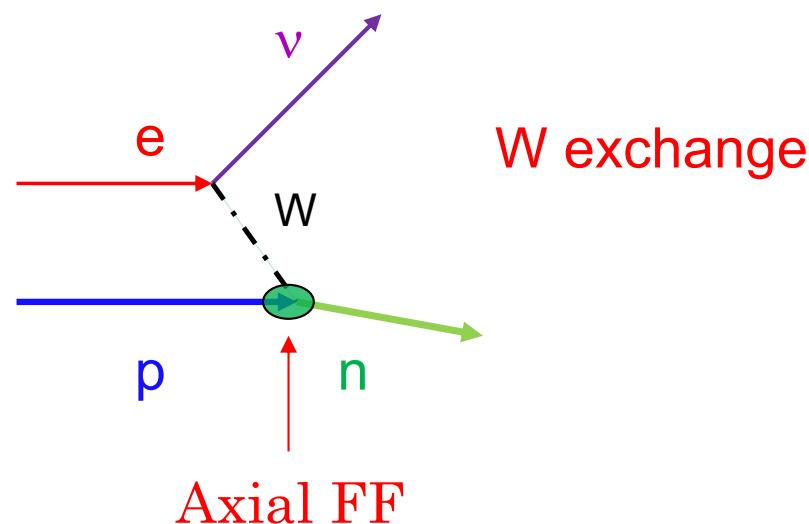


Weak Axial-vector Form Factor

B. Wojtsekhowski, JLab

in collaboration with

P. Deltiarenko, A. Deur, J. Golak,
D. Jones, C. Keppel, E. King, J. Napolitano



Charge current experiments for DIS

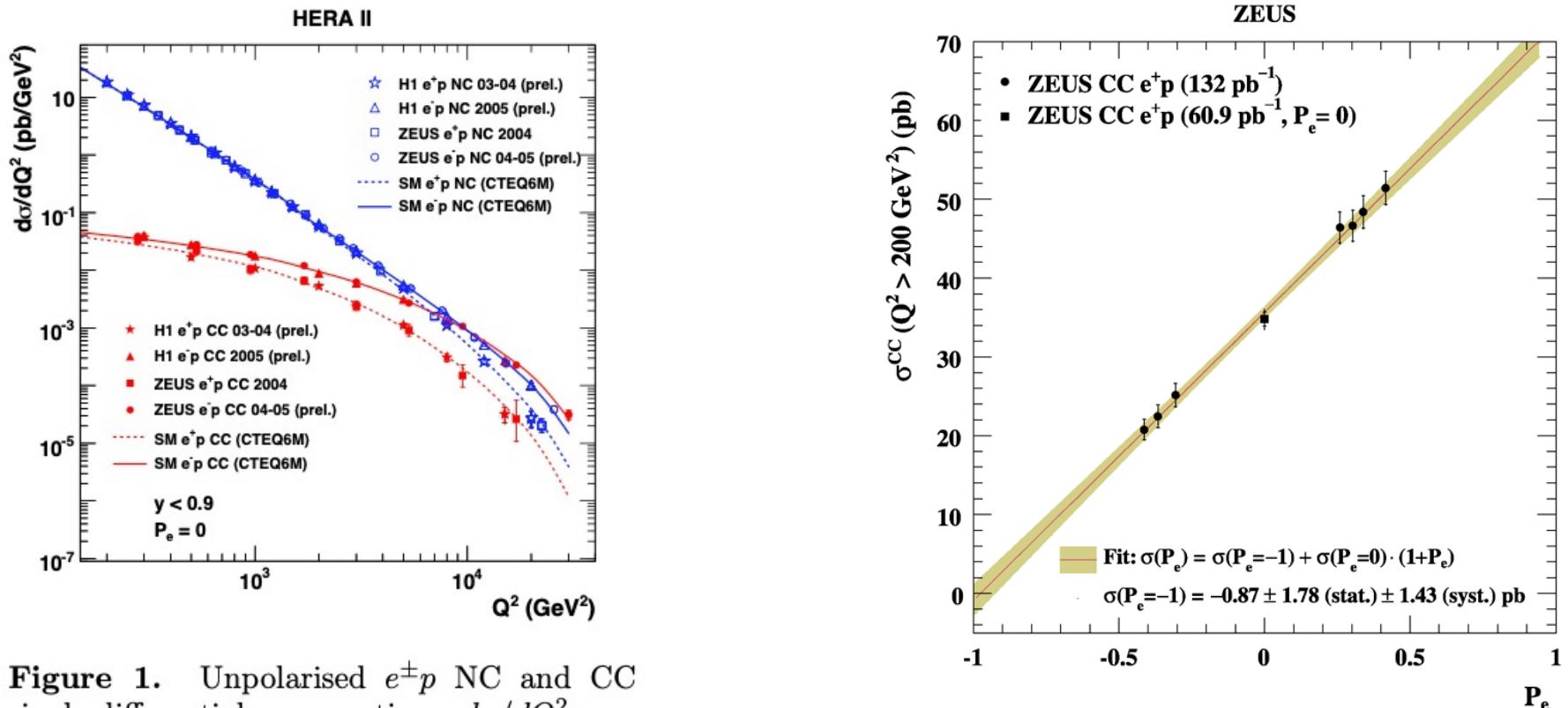


Figure 1. Unpolarised $e^{\pm}p$ NC and CC single differential cross sections, $d\sigma/dQ^2$.

3. Polarised CC cross sections

The longitudinal polarisation has a particularly strong effect on the CC cross sections, as they are predicted to be linearly dependent on the polarisation, independently of kinematic variables:

$$\sigma_{CC}^{e^{\pm}p}(P_e) = (1 \pm P_e)\sigma_{CC}^{e^{\pm}p}(P_e = 0). \quad (4)$$

High-energy quasielastic $\nu_\mu n \rightarrow \mu^- p$ scattering in deuterium

T. Kitagaki, S. Tanaka, H. Yuta, K. Abe, K. Hasegawa,
A. Yamaguchi, K. Tamai, T. Hayashino, Y. Otani, H. Hayano, and H. Sagawa

vector-current (CVC) hypothesis are also assumed to simplify the formulation. Reported values of the axial-vector mass M_A range between 0.65 and 1.07 GeV; the weighted average is somewhat smaller than, but consistent with, the mass value $M_A \sim 1.15$ GeV obtained from electroproduction experiments.⁵ These results, as well as the absolute cross section for the quasielastic reaction, are consistently described by the formulation of the $V-A$ theory in the low-energy $\nu_\mu n \rightarrow \mu^- p$, there has been no experi-

from this experiment have been published elsewhere.⁶⁻¹⁰

II. EXPERIMENTAL DETAILS

A. Neutrino beam and bubble chamber

The wide-band neutrino beam was produced by 350-GeV/c protons striking a 33-cm-long beryllium oxide target. Figure 1 shows the schematic layout of the neutrino beam line. Secondary particles with positive charge were focused by a horn magnet pulsed to a maximum current of 80 kA. The neutrinos were produced from π^+ and K^+ decays in flight in a 400-m-long decay pipe. With the exception of neutrinos, almost all particles which pass through this decay pipe are absorbed in the 900-m-long earth berm and iron shield. Thus at the end of the berm, a beam consisting primarily of ν_μ emerged. The contamination of the neutrino flux by antineutrinos is estimated by a Monte Carlo simulation¹¹ to be about 14%. The neutrino flux has a maximum at 20 GeV and extends above 200 GeV with an average energy of 27 GeV. A total of 328 000 pictures was taken with 4.9×10^{18} extracted protons, averaging about 1.5×10^{13} protons per pulse. A detailed study of this flux is given in Ref. 10.

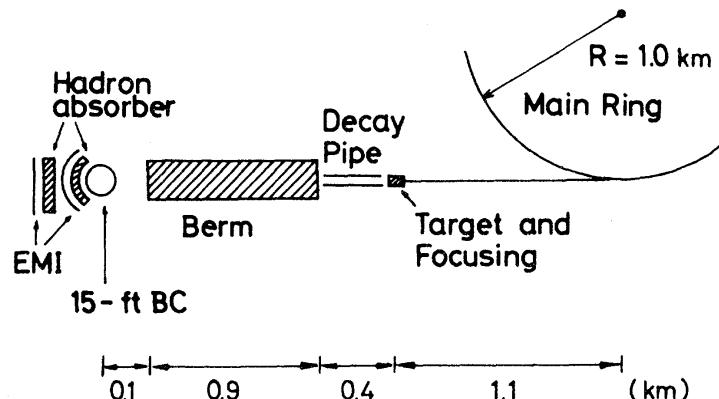


FIG. 1. Schematic layout of the neutrino beam line and the bubble chamber with two-plane external muon identifiers (EMI).

Article

Measurement of the axial vector form factor from antineutrino–proton scattering

48 | Nature | Vol 614 | 2 February 2023

<https://doi.org/10.1038/s41586-022-05478-3>

Received: 19 April 2022

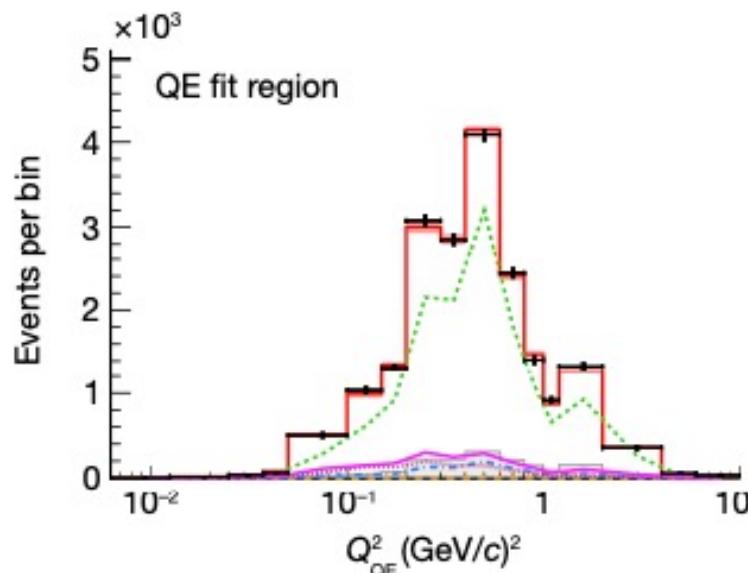
Accepted: 25 October 2022

Published online: 1 February 2023

Open access

 Check for updates

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The CC proposal history at JLab

1988 LOI to PAC3 by J. Napolitano

2003 LOI to PAC25 by A. Deur

2023 LOI to PAC51 by D. Dutta

The LOI 1988 is missing in PAC3 report, discovered just recently

The LOI 2003 had the focus on $Q^2 \sim 1\text{-}3 \text{ GeV}^2$. Interest is very large, some questions about how to proceed, was not updated to proposal

The LOI 2023 proposed a different idea (for low Q^2) based on TDIS proton detector and the reaction with a positron beam:
 $e^+ + d \rightarrow p + p + \nu$

Collaboration buildup

The group started from a concept + LOIs + a few experts

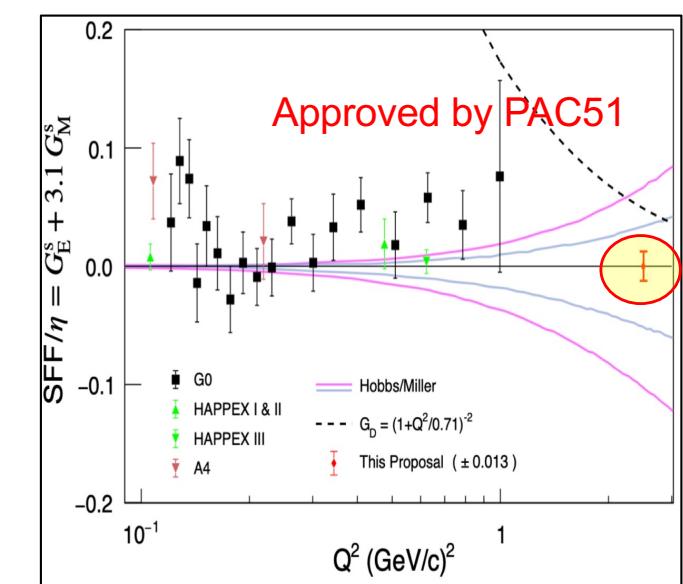
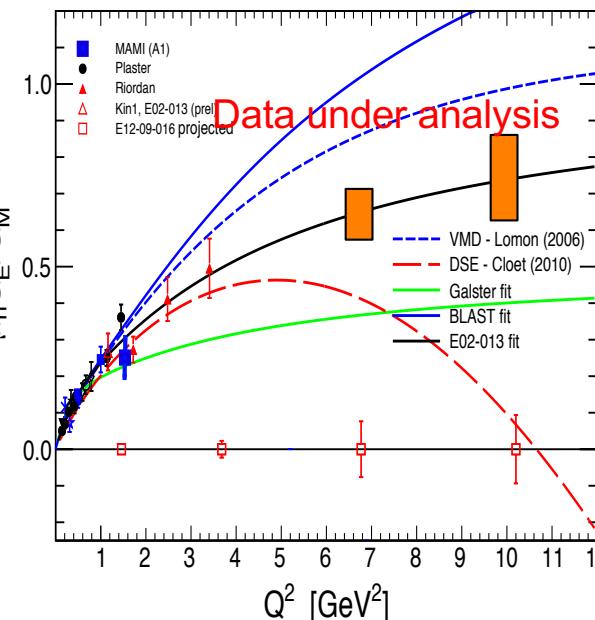
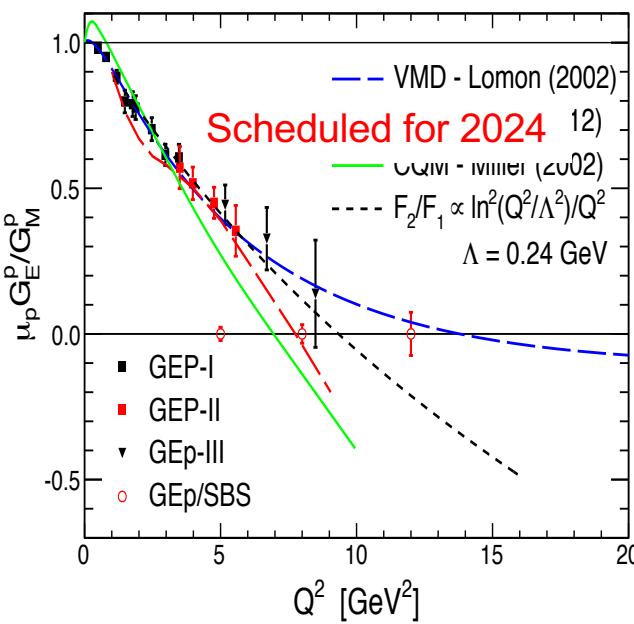
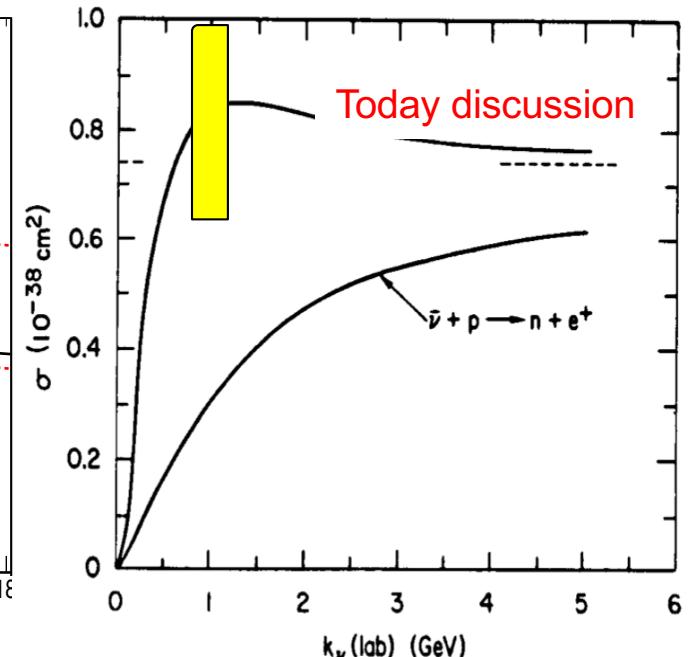
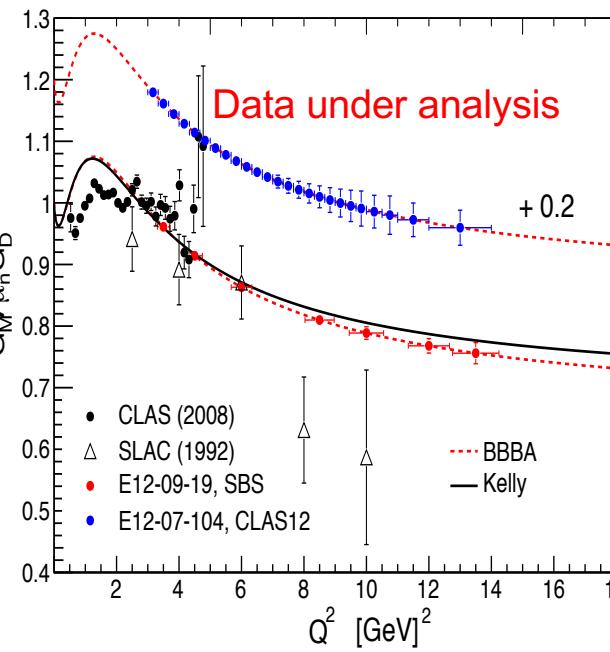
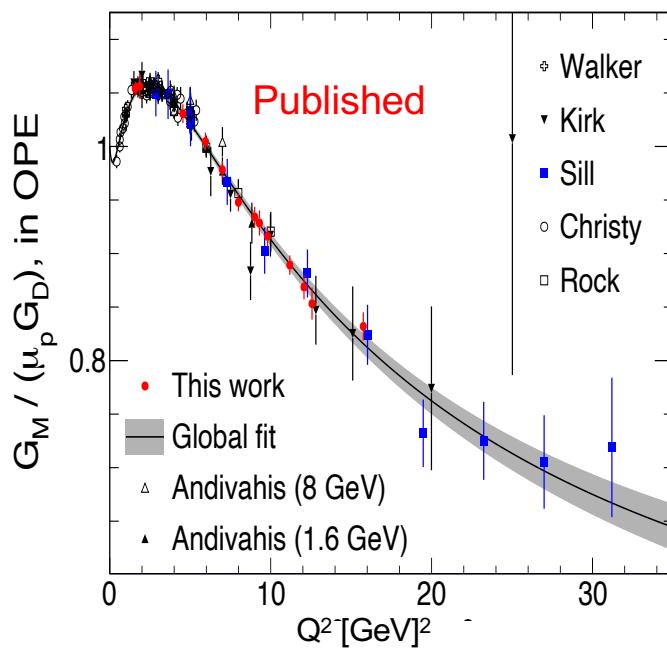
Added INFN: E.Cisbani, R.Perrino, O.Benhar +

Added SBS collaborators: T.Averett

CLAS12 detector expert: D.Carman

CH4 detector under study by M.Bukhari

The nucleon elastic FFs



Cross section calculation

NEUTRINO REACTIONS AT ACCELERATOR ENERGIES *

C.H. LLEWELLYN SMITH

Stanford Linear Accelerator Center, Stanford University, Stanford, California 94305, USA

$$F_A(Q^2) = F_A(0)/(1+Q^2/M_A^2)^2 , \quad (6)$$

where the value of $F_A(0) = -1.23 \pm 0.01$ is taken from β -decay experiments.¹⁶

From these assumptions, the differential cross section for the quasielastic reaction can be expressed in terms of only one parameter, M_A , as

$$\frac{d\sigma}{dQ^2} = \frac{G^2 M^2 \cos^2 \theta_C}{8\pi E_\nu^2} \left[A(Q^2) + B(Q^2) \frac{(s-u)}{M^2} + C(Q^2) \frac{(s-u)^2}{M^4} \right] , \quad (7)$$

where $s-u=4ME_\nu-Q^2-m_\mu^2$, and $M=(M_n+M_p)/2$. The values of the Fermi constant and of the Cabibbo angle are taken to be $G=1.16632 \times 10^{-5}$ GeV $^{-2}$ and $\cos\theta_C=0.9737$, respectively (see Ref. 16). The structure

$$M_A = 1.03 \pm 0.04 \text{ GeV} ,$$

Received 30 August 1971

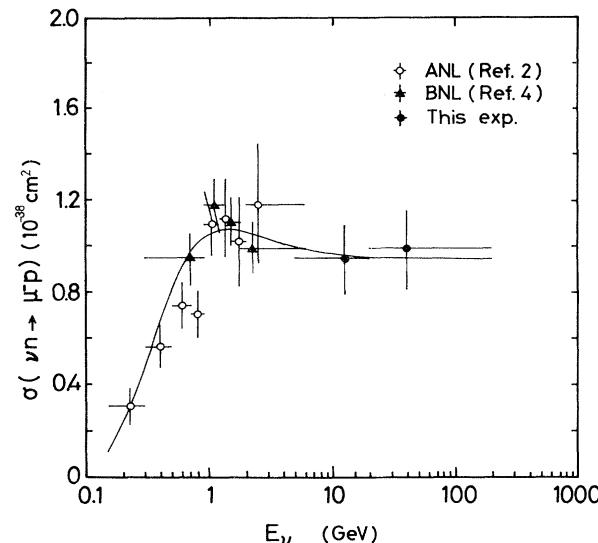


FIG. 10. Quasielastic cross section $\sigma(\nu_\mu n \rightarrow \mu^- p)$ as a function of E_ν . The data points from this experiment and Ref. 4 are calculated from Eq. (7) using the M_A values in Table I. The curve is derived from Eq. (7) with $M_A = 1.05$ GeV.

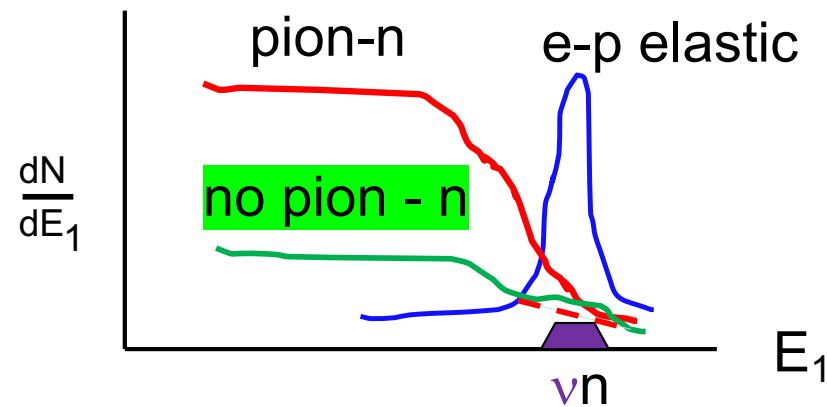
Challenges in the study of $e + p \rightarrow \nu + n$ process

- Cross section for the weak process is $\sim 10^{-39} \text{ cm}^2/\text{sr}$
- Pion photo-production cross section $\sim 10^8$ of the weak one
- Proton rate from electron elastic e-p $\sim 10^7$ of the weak one

The key is a lepton initial energy

$$E_1 = (E_n - m) / \left[1 + \frac{P_n \cos \theta_n - E_n}{m} \right]$$

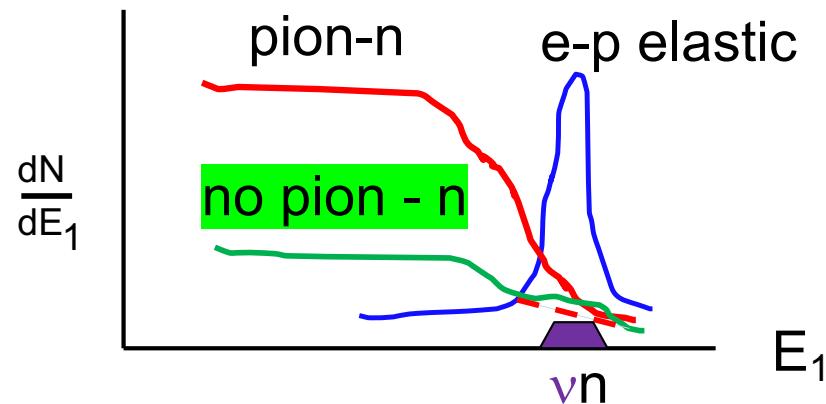
Reconstructed using the momentum and angle of the neutron



The key is a lepton initial energy

$$E_1 = (E_n - m) / \left[1 + \frac{P_n \cos \theta_n - E_n}{m} \right]$$

Reconstructed using the momentum and angle of the neutron

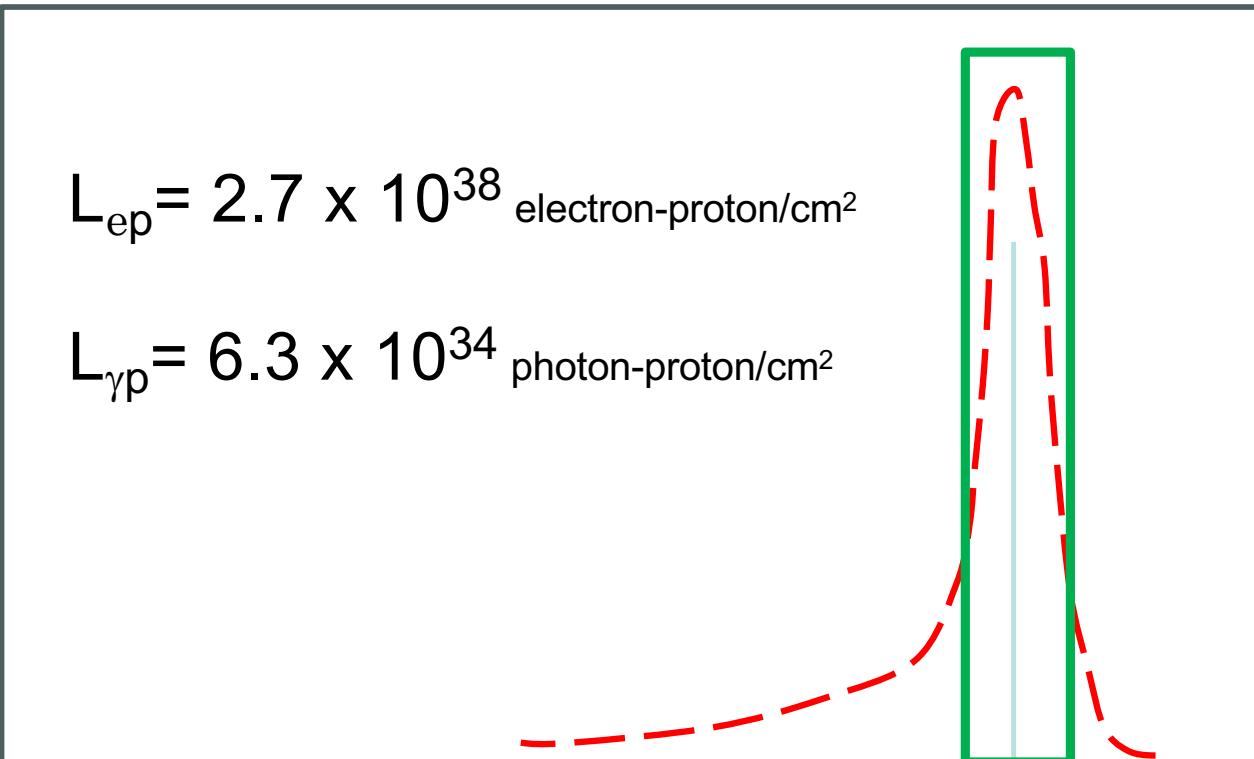


$$E_1 = (E_n - m_n + m_\pi^2/2m_n) / \left[1 + \frac{P_n \cos(\theta_n) - E_n}{m_n} \right]$$

Proposed solution for the $e + p \rightarrow \nu + n$ experiment

1. High **momentum resolution** neutron detector
2. High **angular resolution** neutron detector
3. **Reconstruction of the incident energy to at least 1%**
4. High efficiency of the charge particle spectrometer as a veto
5. Analysis of the distribution (3.) shape
6. Determination of the extra rate at the elastic “peak”
7. **Beam helicity effect is 100% for $e + p \rightarrow \nu + n$**

$H(e,e'p)$



$\sigma_p = 0.5\%$

event rate at 30 deg,

50 msr ~ 190 kHz

$E = 2.2$ GeV

$Q^2 = 1$ GeV²

Electron initial energy from a spectrometer

event rate at (degrees) 30-el / 48-p

Electrons ~ 190 kHz

E-shower trigger with p-veto ~ 2 kHz

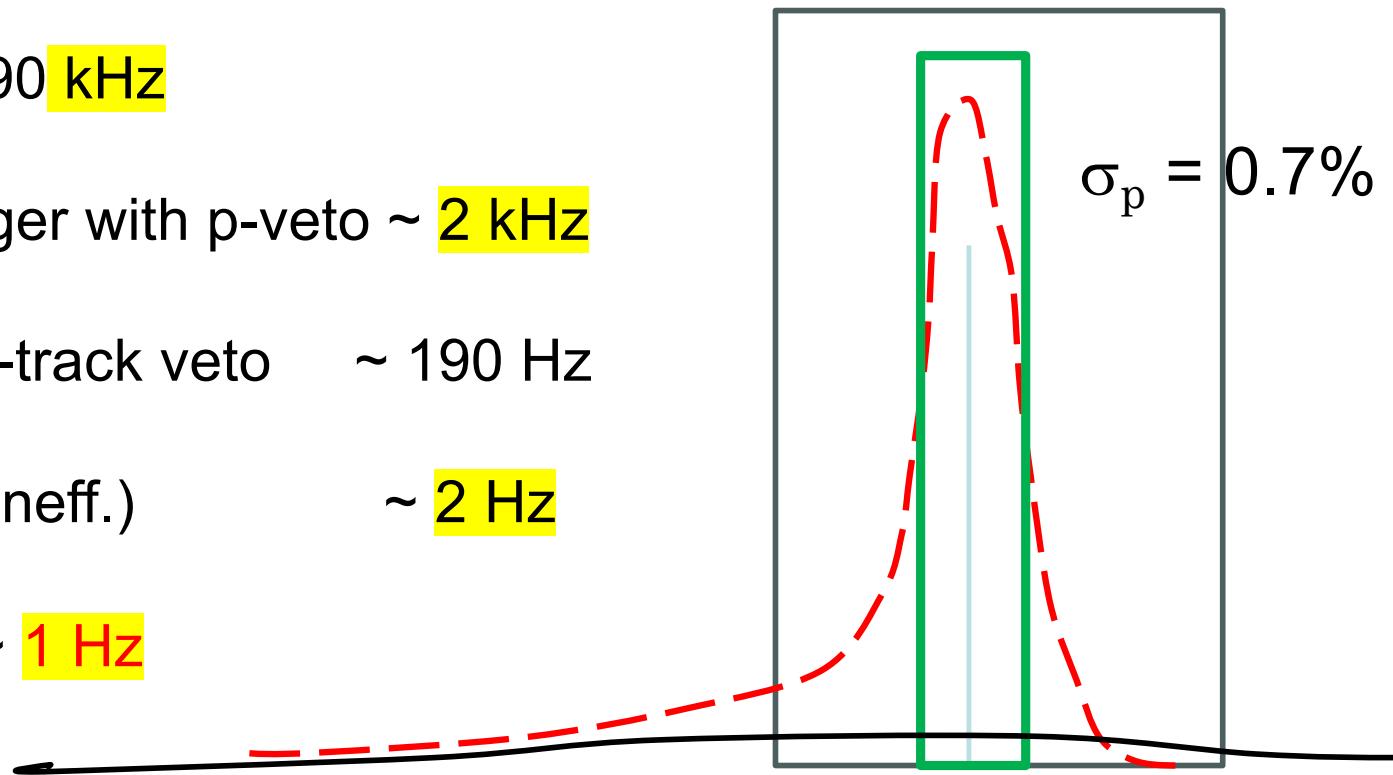
Off-line with e-track veto ~ 190 Hz

+ p-veto (1% ineff.) ~ 2 Hz

+ 50% n-eff. ~ 1 Hz

$H(e, n \backslash no\ p \backslash no\ e')$

$\sigma_p = 0.7\%$



Electron initial energy from T_p and θ_p

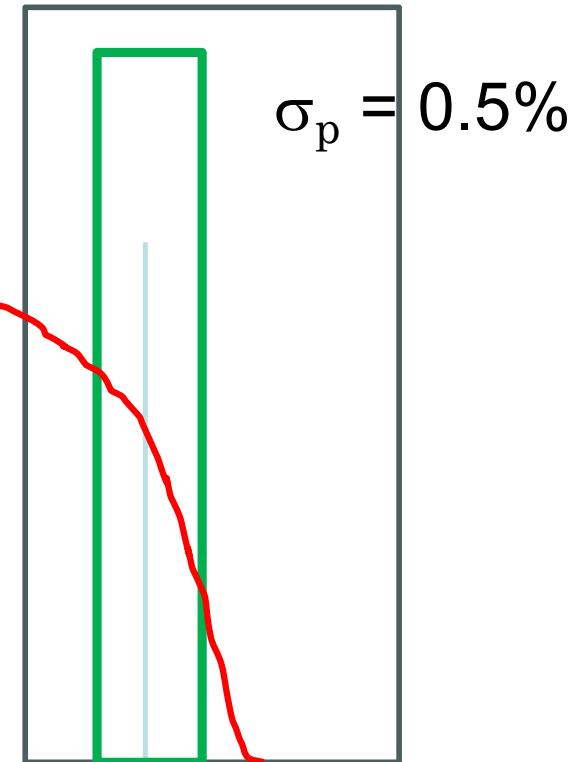
Pion rate calculation:

$H(e, \pi^+ n)$

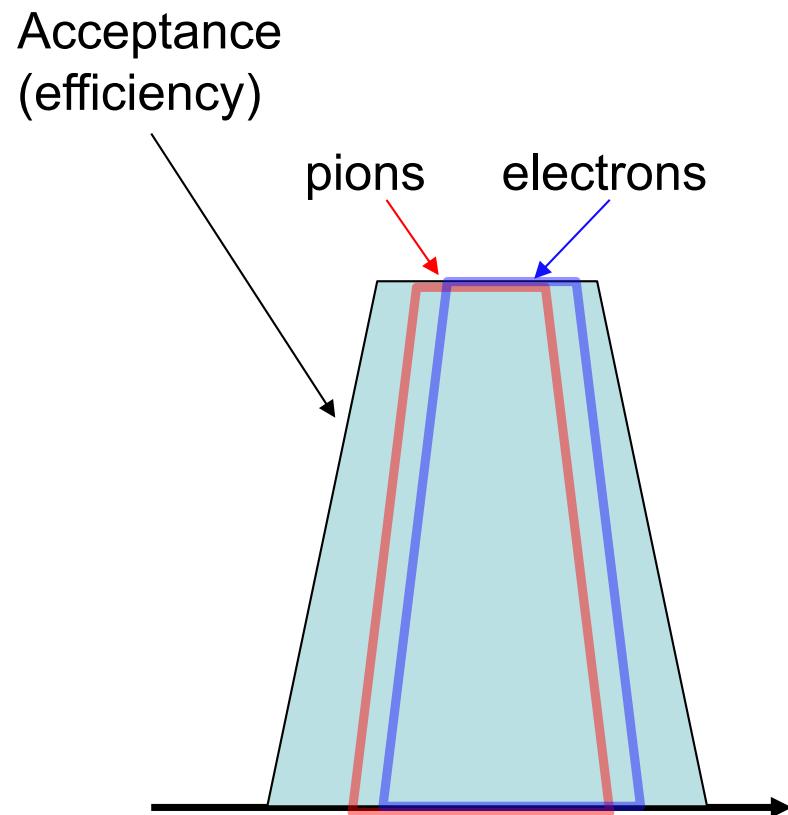
$L_{\gamma p} = 6.3 \times 10^{34}$ (both real and q-real photons in 1.4% interval at end-point)

$d\sigma/d\Omega_{\pi}(\text{cms}) = 0.5 \times 10^{-30} \text{ cm}^2/\text{sr}$
(CLAS arXiv:0903.1110) =>

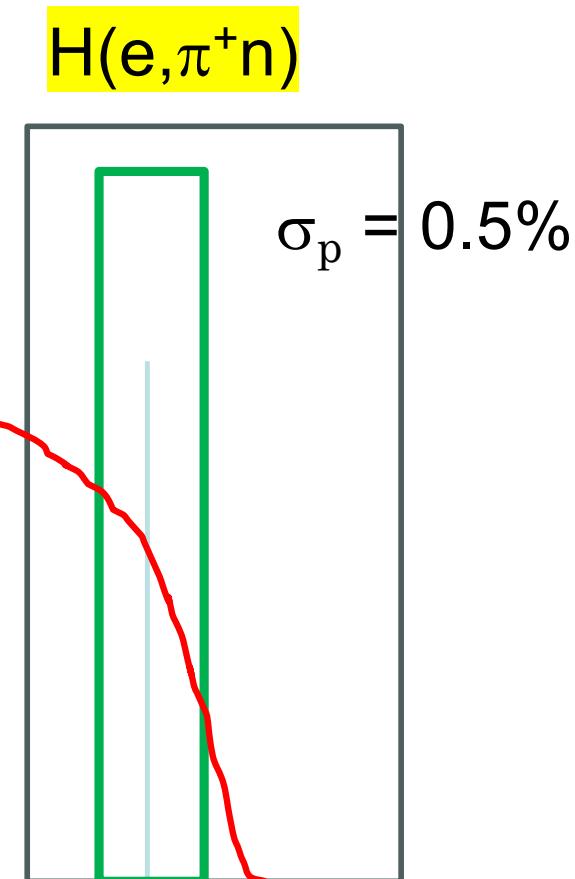
Rate in 50 msr (lab solid angle)
is 400 Hz



Photon energy from a pion E, θ in the spectrometer

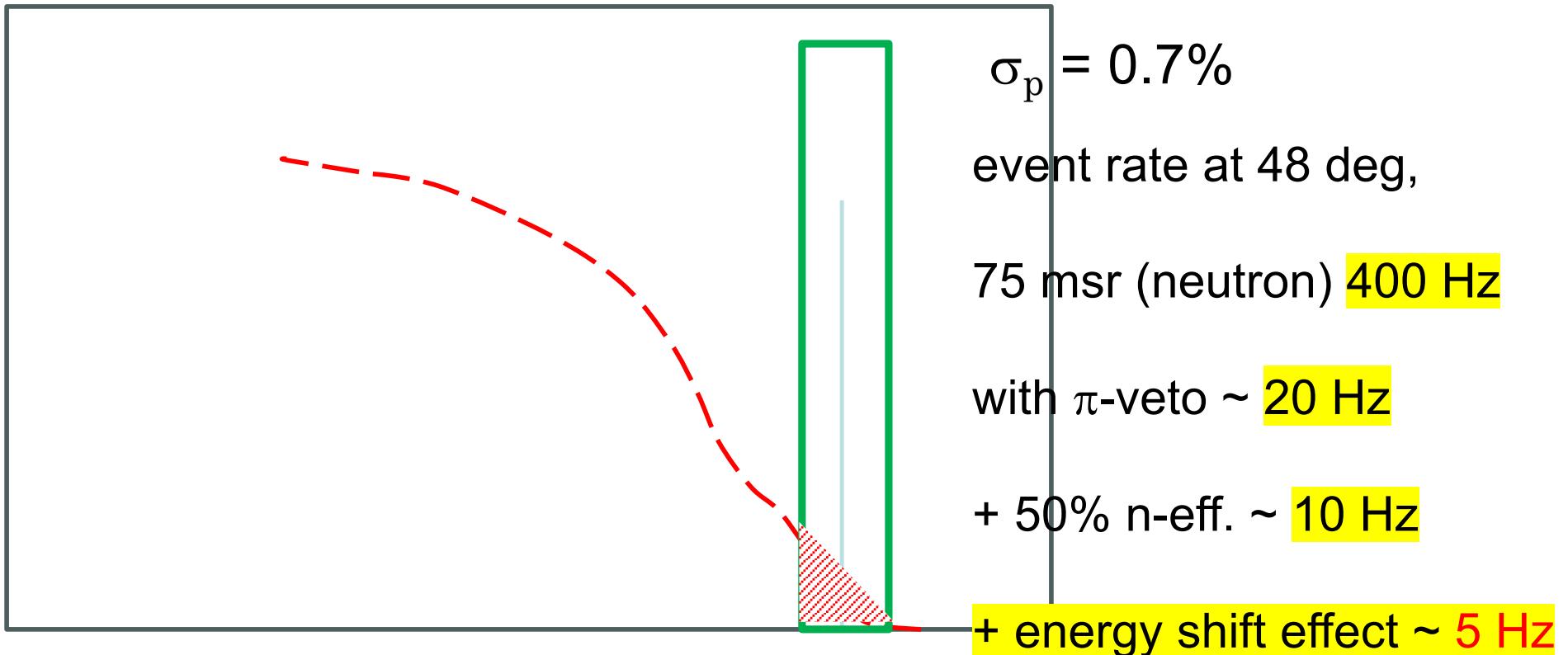


SBS has 70 msr: a large acceptance in angle θ
 Shift ($e-\pi$) = 0.35 deg



Photon energy from a pion E, θ in the spectrometer

$H(e,n \setminus no \pi^+)$



Photon energy from T_n and θ_n (assuming CC process)
It has **42 MeV ($\sim 2\%$) shift** due to the pion mass effect

$$E_1 = (E_n - m_n + m_\pi^2/2m_n) / \left[1 + \frac{P_n \cos(\theta_n) - E_n}{m_n} \right]$$

Factor ~ 2 in the rate reduction

$$L_{ep} = 2.7 \times 10^{38}$$

$d\sigma/d\Omega_\nu = 2.5 \times 10^{-39} \text{ cm}^2/\text{sr}$
the sr in lab for neutrino
from A.Deur's code

We used $1.1 \times 10^{-39} \text{ cm}^2/\text{sr}$
from J.Golak's calculation

Rate in 75 msr (neutron)
is 54 per hour

$H(e,n)\nu$

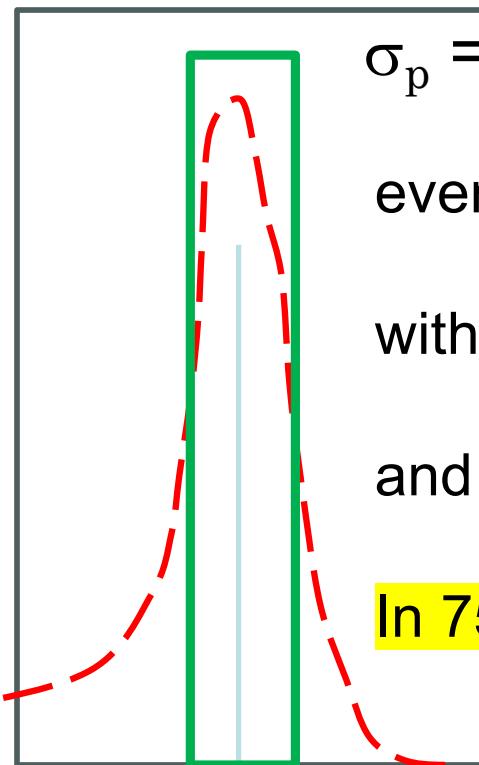
$$\sigma_p = 0.7\%$$

event rate at 48 deg,

with 50% n-efficiency

and 85% P_e

In 75 msr ~ 23 per hour



Electron energy from T_n and θ_n (assuming CC process)

$$H(e,n \setminus no e') + H(e,n \setminus no \pi^+)$$

At $Q^2 = 1 \text{ GeV}^2$

$S/B \sim 23 * / 21600$ in one hour
($A_s = 1$, $P_e = 0.85$)

Beam helicity asymmetry
 $A = 1.1 \times 10^{-3}$

500-hour data taking run

$$A = 11 \times 10^{-4} \pm 3 \times 10^{-4}$$

$$\sigma_p = 0.7\%$$

event rate at 48 deg,

+ 50% n-eff. $\sim 6 \text{ Hz}$
(combined ep, $\gamma\pi$)

or ~ 21600 per hour

Electron energy from T_n and θ_n (for CC process)

Weak Proton Form Factor at 1 GeV²

Estimation of the experiment parameters:

- Beam energy **2.2 GeV** with a 10-cm long LH2 target
- Electron/pion/neutrino angle **30** degrees, $p_e = 1.7 \text{ GeV}/c$
- Recoil proton/neutron **48** degrees, $p_n = 1.1 \text{ GeV}/c$
- π^+ in SBS; efficiency $\sim 90\%+5\%$ (μ are forward) – **need MC**
- Electron detection efficiency **99.9%**; solid angle 50 msr – **need MC**
- Neutron in LND: solid angle 75 msr; at **15 m distance**: **2 m x 8 m**
- For **1%** dE_1/E_1 using $E_1 = (E_n - m) / \left[1 + \frac{P_n \cos \theta_n - E_n}{m} \right]$
angular resolution $\delta\theta \sim 6 \text{ mrad} \Rightarrow 8 \text{ cm}$ coordinate
time resolution **dt => 0.11 ns**

Weak Proton Form Factor at 1 GeV²

- Beam energy **2.2 GeV** with a 10-cm long LH2 target
- Electron/pion/neutrino angle **30** degrees, $p_e = 1.7 \text{ GeV}/c$
- Recoil proton/neutron **48** degrees, $p_n = 1.1 \text{ GeV}/c$
- π^+ in SBS; efficiency $\sim 90\% + 5\%$? (μ are forward) – **need MC**
- Electron detection efficiency **99.9%**; solid angle 50 msr – **need MC**
- Neutron in LND: solid angle 75 msr; at **15m distance**: **2m x 8m**
- For **1%** dE_1/E_1 using $E_1 = (E_n - m) / \left[1 + \frac{P_n \cos \theta_n - E_n}{m} \right]$

Projected result $A = 11 \times 10^{-4} \pm 3 \times 10^{-4}$

Weak process cross section $1.1 \pm 0.3 \times 10^{-39} \text{ cm}^2/\text{sr}$

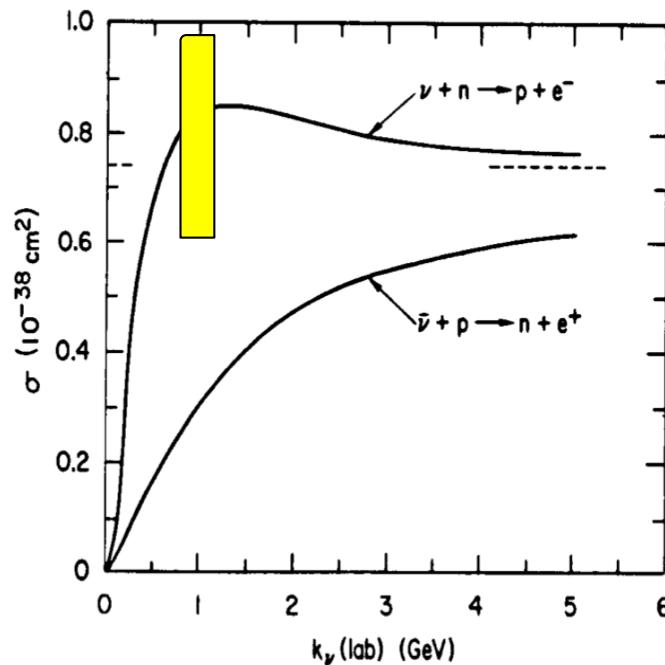
False asymmetry is below 10^{-5} , as it was in MC for sFF: A in e,e'p due to the recoil proton side polarization and the detector as an analyzer

Proton Axial-vector Form Factor at 1 GeV²

Projected result $A = 11 \times 10^{-4} \pm 3 \times 10^{-4}$

Weak process cross section $1.1 \pm 0.3 \times 10^{-39} \text{ cm}^2/\text{sr}$

After 500 hours data taking



Backup slides

Time-of-Flight resolution

$$\frac{\sigma_p}{p} = \gamma^2 \times \frac{\sigma_\beta}{\beta}$$

$$\frac{\sigma_\beta}{\beta} = \frac{\sigma_{ToF}}{ToF}$$

for 10 m path and 0.12 ns time resolution

using $Q^2=1$ GeV 2 ($\gamma=1.5$)

$$\frac{\sigma_p}{p} = \gamma^2 \times 1/275 \sim 0.4\%$$

0.12 ns time resolution is hard

Modern ToF system CLAS12

D.S. Carman, L. Clark, R. De Vita et al.

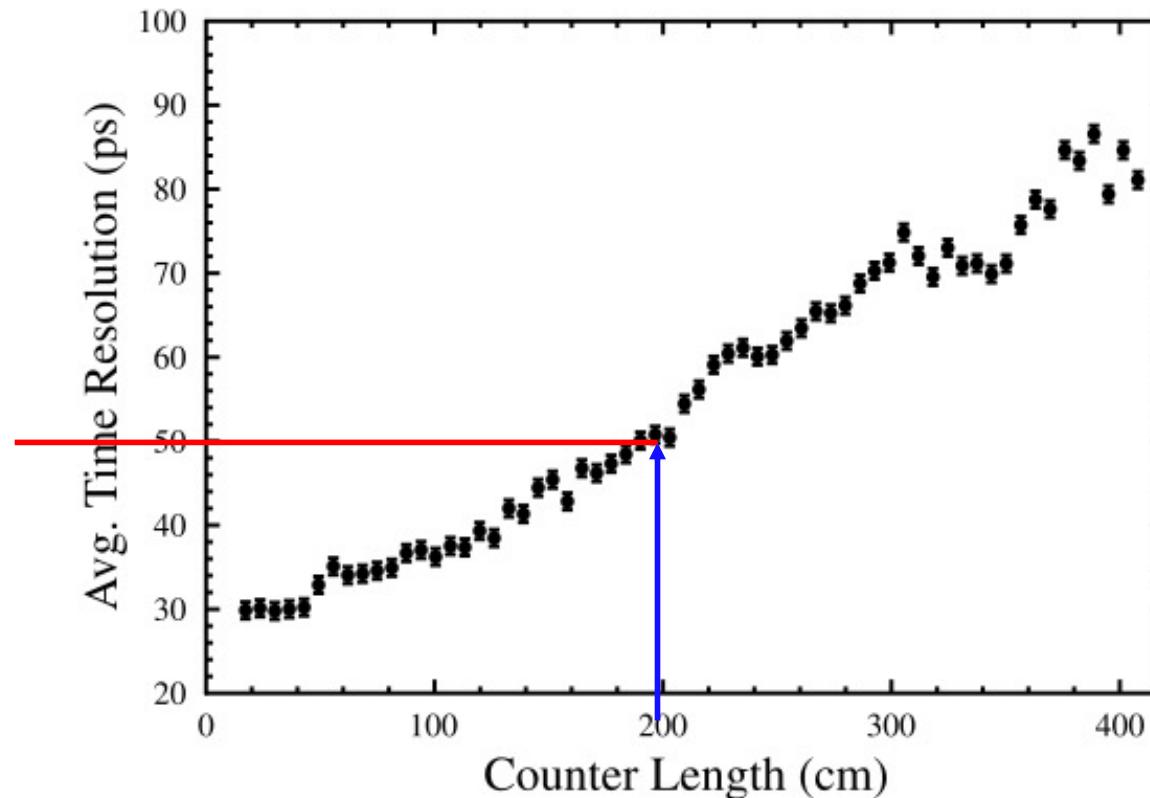


Fig. 14. Measurements of the time resolution (ps) vs. counter length (cm) achieved for the FTOF panel-1b system averaged over the six counters of a given length belonging to each CLAS12 Forward Carriage sector. These data were acquired on the bench using cosmic rays. Full details are included in Ref. [11].

Modern ToF system CLAS12

Parameter	Design Value
Panel-1a	
Angular Coverage	$\theta = 5^\circ \rightarrow 35^\circ, \phi : 50\% \text{ at } 5^\circ \rightarrow 85\% \text{ at } 35^\circ$
Counter Dimensions	$L = 32.3 \text{ cm} \rightarrow 376.1 \text{ cm}, w \times h = 15 \text{ cm} \times 5 \text{ cm}$
Scintillator Material	BC-408
PMTs	EMI 9954A, Philips XP2262
Design Resolution	90 ps \rightarrow 160 ps
Panel-1b	
Angular Coverage	$\theta = 5^\circ \rightarrow 35^\circ, \phi : 50\% \text{ at } 5^\circ \rightarrow 85\% \text{ at } 35^\circ$
Counter Dimensions	$L = 17.3 \text{ cm} \rightarrow 407.9 \text{ cm}, w \times h = 6 \text{ cm} \times 6 \text{ cm}$
Scintillator Material	BC-404 (#1 \rightarrow #31), BC-408 (#32 \rightarrow #62)
PMTs	Hamamatsu R9779
Design Resolution	60 ps \rightarrow 110 ps
Panel-2	
Angular Coverage	$\theta = 35^\circ \rightarrow 45^\circ, \phi : 85\% \text{ at } 35^\circ \rightarrow 95\% \text{ at } 45^\circ$
Counter Dimensions	$L = 371.3 \text{ cm} \rightarrow 426.1 \text{ cm}, w \times h = 22 \text{ cm} \times 5 \text{ cm}$
Scintillator Material	BC-408
PMTs	Photonis XP4312B, EMI 4312KB
Design Resolution	145 ps \rightarrow 160 ps

e 1: Table of parameters for the scintillators, PMTs, and counters for the FTOF panel-1b, and panel-2 arrays.

Weak Proton Form Factor at 2 GeV²

Estimation of the experiment parameters:

- Beam energy **4.4 GeV** with a 10-cm long LH2 target
- Electron/pion/neutrino angle **21** degrees, $p_e = 3.4 \text{ GeV}/c$
- Recoil proton/neutron **43** degrees, $p_n = 1.7 \text{ GeV}/c$
- π^+ in SBS; efficiency $\sim 96\%+3\%$ (μ are forward) – need MC
- Electron detection efficiency **99.9%**; solid angle 50 msr – need MC
- Neutron in LND: solid angle 135 msr; at **15 m distance**: **4m x 8m**

Luminosities

$$L_{ep} = 2.7 \times 10^{38} \text{ for } 10 \text{ cm LH2} \times 100 \mu\text{A}$$

Photon flux estimate:

Quasi real: $I_e \times 0.013 \times dE_\gamma/E_e \times 0.75$ (Budnev-1975)

Real: $I_e \times 0.007 \times dE_\gamma/E_e \quad (10\text{cm LH2} / 2 / 735\text{cm})$

$$I_e \times (0.013 \times 0.014 \times 0.75 + 0.007 \times 0.014)$$

$$L_{\gamma p} = 6.3 \times 10^{34}$$

Parameters of SBS

Solid angle

$\theta_{central}$, degree	Ω , msr	D, meter	Hor. range, degree	Vert. range, degree
3.5	5	9.5	± 1.3	± 3.3
5.0	12	5.8	± 1.9	± 4.9
7.5	30	3.2	± 3	± 8
15	72	1.6	± 4.8	± 12.2
30	76	1.5	± 4.9	± 12.5

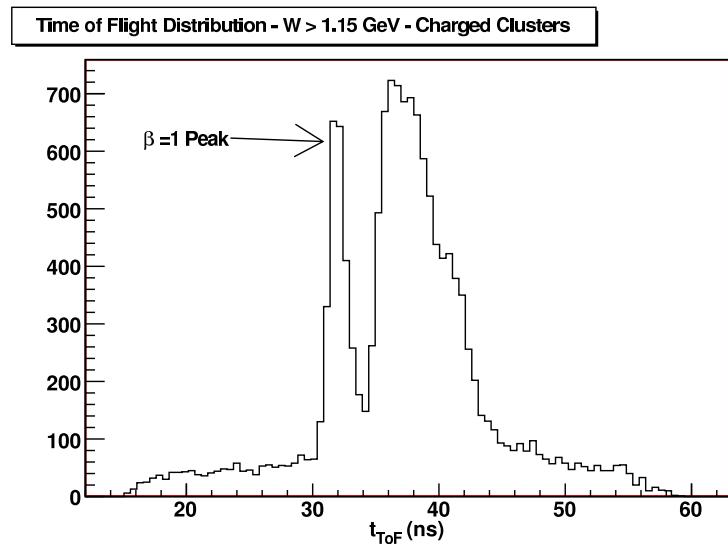
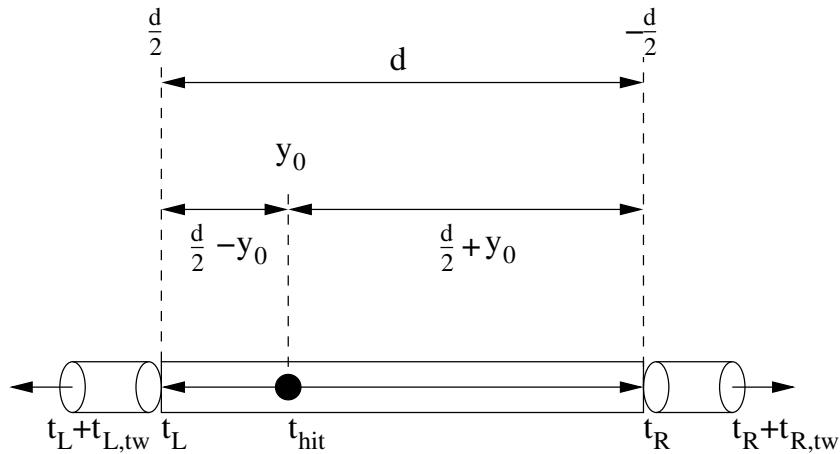
Resolution:

Momentum $\Rightarrow \frac{\sigma_p}{P} = 0.0029 + 0.0003 \times p [\text{GeV}]$

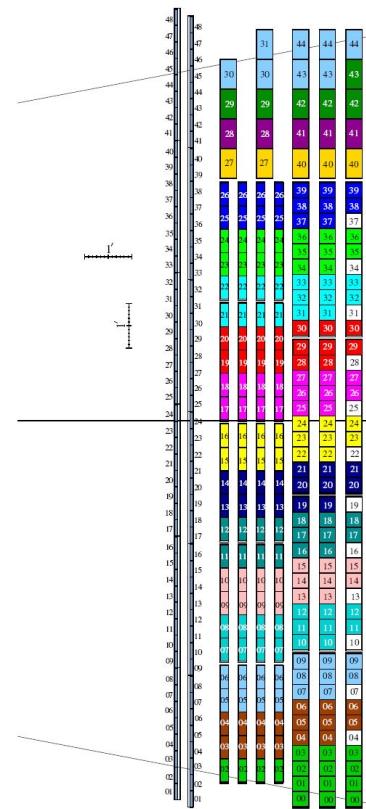
Angular $\Rightarrow \sigma_\theta = 0.14 + 1.3/p [\text{GeV}], \text{ mrad}$

Momentum acceptance $\Rightarrow P$ range from $2 - 10, \text{ GeV}/c$

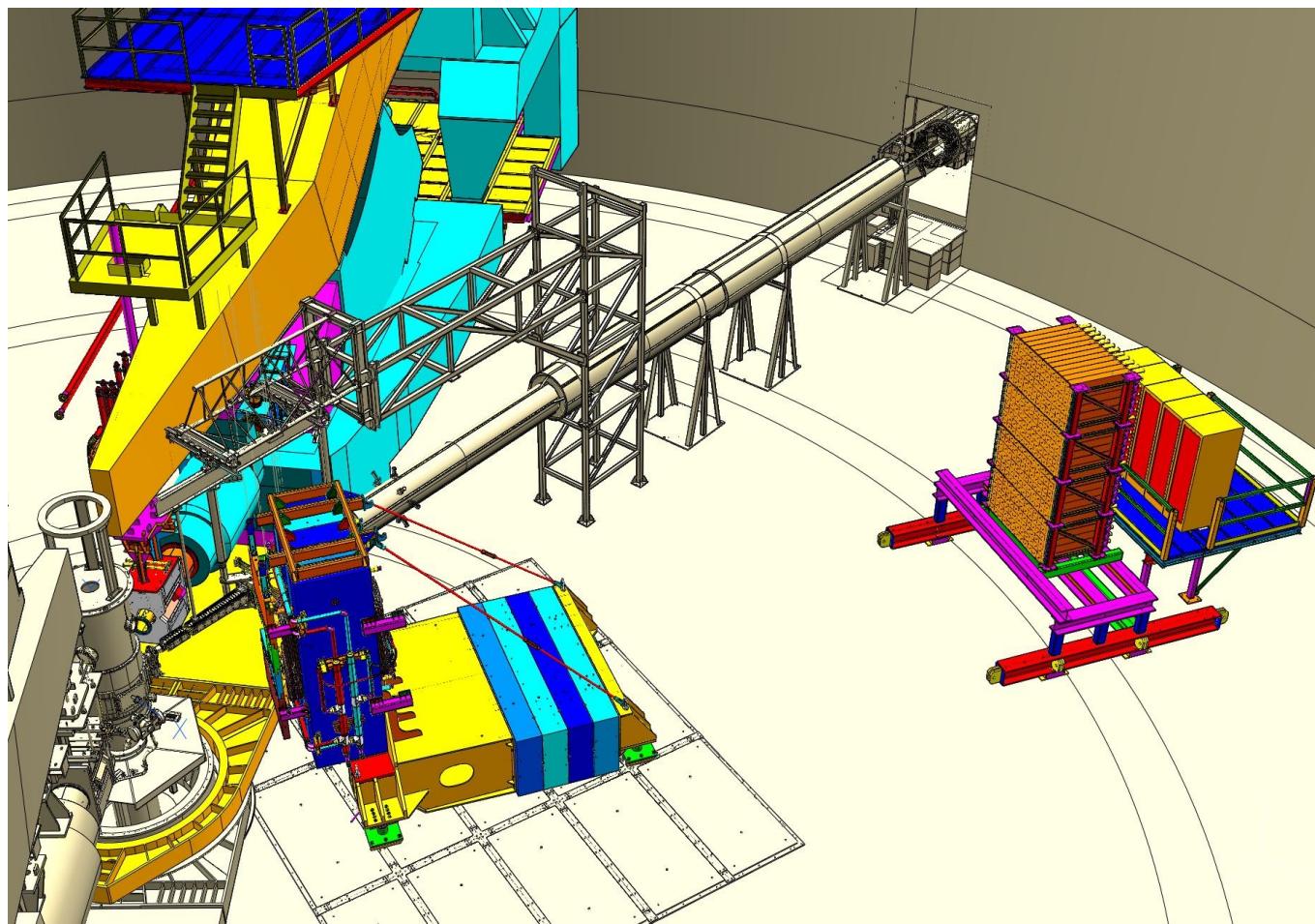
Time-of-flight



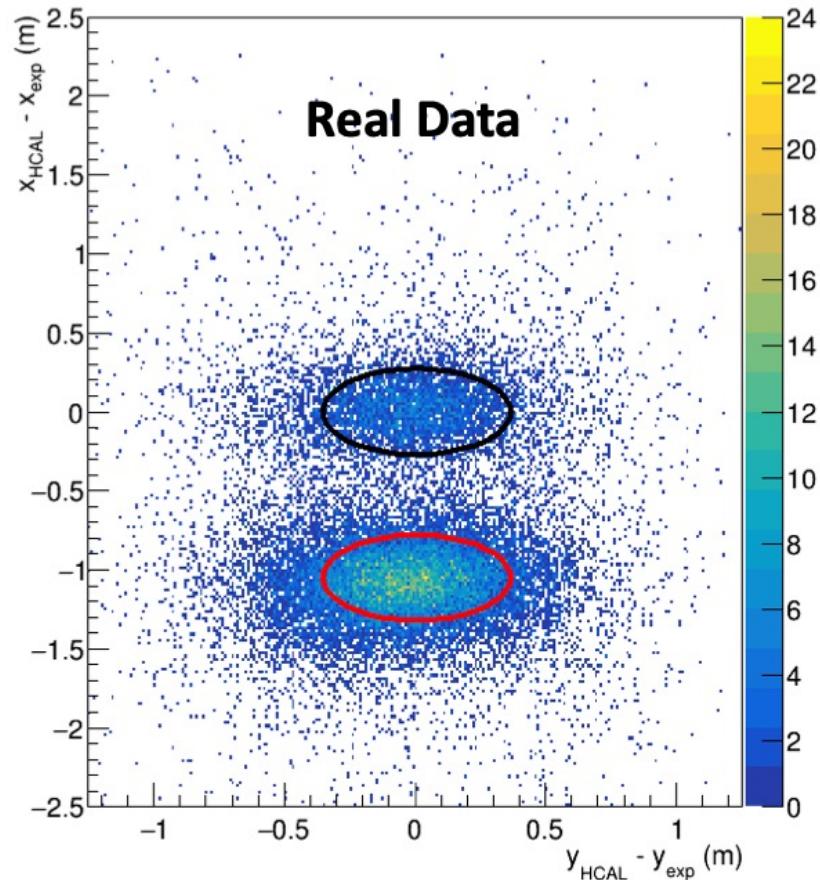
GEn-I “Big Hand”



SBS neutron arm – Hadron Shower Calorimeter



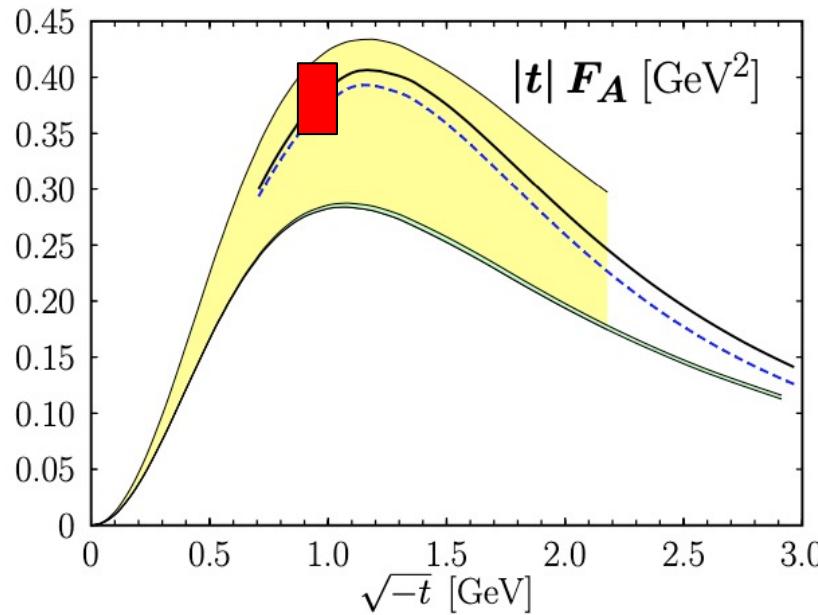
SBS neutron arm



Isovector axial FF and GPD \tilde{H}

M. Diehl, P. Kroll arXiv:1302.4604; [hep-ph](#)>arXiv:1703.05000

JLab - Wide Angle Compton Scattering:



$$\frac{d\sigma}{dt} = \frac{2\pi\alpha_{em}^2}{s^2} \left[-\frac{u}{s} - \frac{s}{u} \right] \left\{ \frac{1}{2} \left(R_V^2(t) + R_A^2(t) \right) - \frac{us}{s^2 + u^2} \left(R_V^2(t) - R_A^2(t) \right) \right\}$$

$$K_{LL} = 2 \frac{-t}{s-u} \frac{R_A}{R_V} \frac{1+\eta\kappa_T}{1+\kappa_T^2} \left[1 + \frac{t^2}{(s-u)^2} \frac{R_A^2}{R_V^2} \frac{1}{1+\kappa_T^2} \right]^{-1}$$

