# Strange-Meson Spectroscopy with COMPASS

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# The Strange-Meson Spectrum





### PDG lists 25 strange mesons

- 16 established states, 9 need further confirmation
- Missing states with respect to quark-model predictions
- Many measurements performed more than 30 years ago

### Strange-Meson Spectroscopy with COMPASS COMPASS Setup for Hadron Beams

TAR Agatt [COMPASS, Nucl. Instrum, Methods 779 (2015) 69]



# Strange-Meson Spectroscopy with COMPASS

Production of Strange Mesons





- Diffractive scattering of high-energy kaon beam
- Strange mesons appear as intermediate resonances X<sup>-</sup>
- Decay to multi-body hadronic final states
- $\blacktriangleright K^-\pi^-\pi^+$  final state
  - Study in principle all strange mesons
  - Study a wide mass range
  - COMPASS measured world'

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  - Study in principle all strange mesons
  - Study a wide mass range
  - COMPASS measured world's largest data set of about 720 k events



### Partial wave: $J^P M^{\varepsilon} \xi b^- L$

- ► *J<sup>P</sup>* spin and parity
- ▶ *M<sup>ε</sup>* spin projection
- ξ isobar resonance
- ▶ b<sup>−</sup> bachelor particle
- L orbital angular momentum





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**Data**: 720 k diffractively produced  $K^-\pi^-\pi^+$  candidates







S. Wallner







### Partial-Wave Decomposition

$$\mathcal{I}(\tau, m_{K\pi\pi}, t') = \sum_{a,b \in \mathbb{W}_z(m_{K\pi\pi}, t')} \Psi_a(\tau) \rho_{ab}(m_{K\pi\pi}, t') \left[\Psi_b(\tau)\right]^*$$

- Measure spin-density matrix  $\rho_{ab}(m_{K\pi\pi}, t')$  in independently  $(m_{K\pi\pi}, t')$  cells
  - ▶ No assumption about  $K^-\pi^-\pi^+$  resonances
- Wave set W<sub>z</sub>(m<sub>Kππ</sub>, t') inferred from data using regularization-based model-selection techniques
- Bootstrap resampling to improve uncertainty estimates
  - Performed about 20 M fits
- Detailed Monte Carlo input-output studies





#### Resonance-Model Fit

$$\hat{\rho}_{ab}^{K\pi\pi}(m_{K\pi\pi},t') = \hat{\mathcal{T}}_{a}(m_{K\pi\pi},t') \left[\hat{\mathcal{T}}_{b}(m_{K\pi\pi},t')\right]^{*}$$
$$\hat{\mathcal{T}}_{a}(m_{K\pi\pi},t') = \sum_{k \in \mathbb{S}_{a}} \mathcal{K}(m_{K\pi\pi},t')^{k} \mathcal{C}_{a}(t') \mathcal{D}_{k}(m_{K\pi\pi};\zeta_{k})$$

- Model  $m_{K\pi\pi}$  dependence of partial-wave amplitudes
- Breit-Wigner amplitudes for K<sup>-</sup>π<sup>-</sup>π<sup>+</sup> resonance components
- Coherent non-resonant component parameterizing other *K*<sup>-</sup>π<sup>-</sup>π<sup>+</sup> production mechanisms





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#### Incoherent Backgrounds

- ▶ Incoherent background from  $\pi^-$  diffraction to  $\pi^-\pi^-\pi^+$ and other reactions (in total about 10%)
  - Very good model for dominant  $\pi^-\pi^-\pi^+$  background from COMPASS  $\pi^-\pi^-\pi^+$  analysis
    - Study background in partial waves by
      - Generate pseudodata from  $\pi^-\pi^-\pi^+$  model
      - Apply  $K^-\pi^-\pi^+$  reconstruction event selection
      - Project into  $K^-\pi^-\pi^+$  partial wave
    - Large in some waves, e.g. with ho(770) isobar
    - Small in other waves, e.g. with K\*(892) isobar



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### Handling of Incoherent Backgrounds

- Challanging to explicitly treat in partial-wave decomposition
  - ➡ Effectively taken into account

$$\rho_{ab} = \sum_{z} \mathcal{T}_{a}^{z} \left[ \mathcal{T}_{b}^{z} \right]^{*}$$

- → Measured  $\rho_{ab}$  include background Explicitly model them in resonance-model fi  $\hat{\rho}_{ab}(m_{\kappa\pi\pi}, t') = \hat{\rho}_{ab}^{\kappa\pi\pi}(m_{\kappa\pi\pi}, t')$ 
  - $\blacksquare \pi^-\pi^-\pi^+$  background modeled by partial-wave projection of  $\pi^-\pi^-\pi^+$  pseudodata
    - Yield is only free parameter
  - Incoherent effective background component for other background processes



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Explicitly model them in resonance-model fit

$$\hat{\rho}_{ab}(m_{K\pi\pi}, t') = \hat{\rho}_{ab}^{K\pi\pi}(m_{K\pi\pi}, t') + \hat{\rho}_{ab}^{3\pi}(m_{K\pi\pi}, t')$$

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- Simultaneously included 14 partial waves in resonance-model fit
- Modeled by 13 strange-meson resonance components
- Using measured intensities and interference terms (relative phases)





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### Partial-Wave Analysis of the $K^-\pi^-\pi^+$ Final State







### PDG

►  $K_2^*(1430)$  well known resonance





- ρ(770) K D
- K\*(892) π D
- In agreement with previous measurements
- Cleaner signal in COMPASS data
- Fitted yield of  $\pi^-\pi^-\pi^+$  background consistent with expectation



total resonance model, resonances, non-resonant,  $\pi\pi\pi$  background, effective background





- K<sub>2</sub><sup>\*</sup>(1430) signal
  - $m_0 = (1430.9 \pm 1.4^{+3.1}_{-1.5}) \text{ MeV}/c^2$ •  $\Gamma_0 = (111 \pm 3^{+4}_{-16}) \text{ MeV}/c^2$
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- ► K(1630)
  - Unexpectedly small width of only  $16 \text{ MeV}/c^2$
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### COMPASS $K^-\pi^-\pi^+$ data

- Peak at about 1.4 GeV/ $c^2$ 
  - Established K(1460)
  - But,  $m_{K\pi\pi} \lesssim 1.5 \,\text{GeV}/c^2$  region weakly affected by known analysis artifacts

### Second peak at about 1.7 GeV/c<sup>2</sup>

K(1630) signal with 8.3 σ statistical significance
 Accompanied by rising phase

Weak signal at about 2.0 GeV/c

K(1830) signal with 5.4  $\sigma$  statistical significance



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- ► K(1830) parameters in good agreement with LCHb measurement [PRL 118 (2017) 022003]
- Expected K(1630) width of about 140 MeV/ $c^2$
## Searching for Exotic Strange Mesons with $J^P = 0^-$





- Indications for 3 excited K from a single analysis
- Quark-model predicts only two excited states: potentially K(1460) and K(1830)
- ➡ K(1630) supernumerary signal
- Solution Candidate for exotic non- $q\bar{q}$  state; other explanations possible ( $K^*(892) \omega$  threshold nearby)

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The Strange-Meson Spectrum

- Many strange mesons require further confirmation
- Search for strange partners of exotic non-strange light mesons





### COMPASS

- World's largest data sample on  $K^-\pi^-\pi^+ \Rightarrow$  Most detailed and comprehensive analysis
- Candidate for exotic strange-meson signal with  $J^P = 0^-$





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AMBER: Proposal for High-Precision Strange-Meson Spectroscopy

▶ Goal: Collect  $10 - 20 \times 10^6 \ K^- \pi^- \pi^+$  events using high-energy kaon beam

AMBER is open for interested collaborators to join





#### COMPASS

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## Backup



- Treating the  $\pi^-\pi^-\pi^+$  and Other Backgrounds
- 8 Resonance-Model Fit
  - Modeling the  $K^-\pi^-\pi^+$  Signal
  - Modeling the  $\pi^-\pi^-\pi^+$  Background
  - Modeling the Effective Background
  - $\chi^2$  Fit Procedure
- 9 Wave-Set Selection
  - Regularization: LASSO
  - Regularization: Generalized Pareto
  - Regularization: Cauchy
  - For the  $K^-\pi^-\pi^+$  Final State
- 14-Wave Resonance-Model Fit

- Searching for Exotic Strange Mesons with  $J^P = 0^-$
- Partial Waves with  $J^P = 2^+$ Partial Waves with  $J^P = 2^-$ Partial Waves with  $J^P = 4^+$

- 11 Kinematic Distribution of  $K^-\pi^-\pi^+$  Events
  - Subsystem
  - $m_{\kappa^-\pi^-}$
  - t' Spectrum
  - Exclusivity
- 12 Systematic Studies of the Partial-Wave Decomposition
  - 14 Waves
  - Leakage Waves
- 13 Leakage Effect
- 14 Incoherent  $\pi^-\pi^-\pi^+$  Background





- ►  $J^P M^{\varepsilon}$ : Spin, parity, and spin projection of  $X^-$
- ► ξ: Isobar
- ▶ b: Bachelor particle. Here: Spectator K<sup>-</sup>
- L: Angular momentum between bachelor and isobar





#### Model intensity

$$\mathcal{I}(\tau, m_{K\pi\pi}, t') = \sum_{z} \left| \sum_{a \in \mathbb{W}_{z}(m_{K\pi\pi}, t')} \mathcal{T}_{a}^{z}(\tau; m_{K\pi\pi}) \right|$$

#### Model intensity distribution

- ▶ in 5D  $K^-\pi^-\pi^+$  phase-space
- for a given  $(m_{\kappa\pi\pi}, t')$  cell
- as incoherent sum over coherent sectors z
  - "Rank" of the partial-wave model = number of coherent sectors
- $\Psi_a^z$  known, assuming the isobar model
- Wave set  $\mathbb{W}_{z}(m_{K\pi\pi}, t')$  inferred from data using regularization-based model-selection techniques
- $T_a^z$  extracted in maximum-likelihood fit, independently for each  $(m_{K\pi\pi}, t')$  cell



Spin-Density Matrix

$$\rho_{\textit{ab}} = \sum_{\textit{z}} \mathcal{T}_{\textit{a}}^{\textit{z}} \, \left[ \mathcal{T}_{\textit{b}}^{\textit{z}} \right]^{*}$$

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#### Approach

Effectively take into account in partial-wave decomposition by incoherently adding additional coherent sectors z

(Model background by  $K^-\pi^-\pi^+$  partial waves)

- $\blacktriangleright$  Increasing the rank of the spin-density matrix  $\rho_{ab}$
- Signal not separated from background in partial-wave decomposition
- ➡ Partial-wave amplitudes include background
- Model signal and background contributions in resonance-model fit using more constrained signal model
  - Separate signal from background

$$\mathcal{I}(\tau, m_{K\pi\pi}, t') = \sum_{z} \left| \sum_{a \in \mathbb{W}_{z}(m_{K\pi\pi}, t')} \mathcal{T}_{a}^{z}(\tau; m_{K\pi\pi}) \right|^{2}$$

$$\rho_{ab} = \sum_{z} \mathcal{T}_{a}^{z} \left[ \mathcal{T}_{b}^{z} \right]^{*}$$



#### True physics intensity distribution

$$\mathcal{I}$$
  $(\tau$  ) =  $\left|\sum_{a}^{\text{waves}} \mathcal{T}_{a} \Psi_{a}(\tau)\right|^{2}$ 

#### Experimentally measured intensity distribution

$$\mathcal{I}_{ ext{measured}}( au \quad ) = \quad \eta \ ( au \ ) \mathcal{I} \ ( au \ )$$

- Take into account different processes p
  - Different model intensities *I*<sup>p</sup>
  - **b** Different experimental acceptance  $\eta^{\mathfrak{p}}$
  - \* Formulated in terms of different phase-space variables  $au^{
    m P}$ 
    - $\blacktriangleright$  Jacobian terms  $J(\tau^{K\pi\pi} \rightarrow \tau^{\mathfrak{p}})$  from variable transformation



# True physics intensity distribution for process $\mathfrak{p}$ Experimentally measured intensity distribution $\mathcal{I}^{\mathfrak{p}}(\tau) = \left| \sum_{a}^{\mathsf{waves}} \mathcal{T}^{\mathfrak{p}}_{a} \Psi^{\mathfrak{p}}_{a}(\tau) \right|^{2}$ $\mathcal{I}_{\mathrm{measured}}(\tau) = \sum_{\mathfrak{p}} \eta^{\mathfrak{p}}(\tau) \mathcal{I}^{\mathfrak{p}}(\tau)$

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- $\mathcal{I}^{\pi\pi\pi}$  known by COMPASS analysis
- $\eta^{\pi\pi\pi}$  from detector simulation

- ▶  $\eta^{\pi\pi\pi}$  computationally expensive
- ▶ Different  $m_{3\pi}$  bins enter one  $m_{K\pi\pi}$  bin
- Other background channels:  $K^-K^-K^+$ , ...
  - I<sup>p</sup> unknown
  - Unknown background channels



#### True physics intensity distribution for process p

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True physics intensity distribution for process p

$$\mathcal{I}^{\mathfrak{p}}(\tau^{\mathfrak{p}}) = \left|\sum_{a}^{\mathsf{waves}} \mathcal{T}^{\mathfrak{p}}_{a} \Psi^{\mathfrak{p}}_{a}(\tau^{\mathfrak{p}})\right|^{2}$$

#### Experimentally measured intensity distribution

$$\mathcal{I}_{\text{measured}}(\tau^{\kappa_{\pi\pi}}) = \sum_{\mathfrak{p}} \eta^{\mathfrak{p}}(\tau^{\mathfrak{p}}) \, \mathcal{I}^{\mathfrak{p}}(\tau^{\mathfrak{p}}) \, J(\tau^{\kappa_{\pi\pi}} \to \tau^{\mathfrak{p}})$$

•  $\mathcal{I}^{\pi\pi\pi}$  known by COMPASS analysis

•  $\eta^{\pi\pi\pi}$  from detector simulation

- $\eta^{\pi\pi\pi}$  computationally expensive
- Different  $m_{3\pi}$  bins enter one  $m_{K\pi\pi}$  bin
- Other background channels:  $K^-K^-K^+$ , ...
  - $\blacktriangleright \mathcal{I}^{\mathfrak{p}}$  unknown
  - Unknown background channels

Treating the  $\pi^-\pi^-\pi^+$  and Other Backgrounds



#### Approximate model for process $\mathfrak{p}$ by $K^-\pi^-\pi^+$ partial waves

$$\eta^{\mathfrak{p}}(\tau^{\mathfrak{p}}) \left| \sum_{a}^{\mathsf{waves}} \mathcal{T}^{\mathfrak{p}}_{a} \Psi^{\mathfrak{p}}_{a}(\tau^{\mathfrak{p}}) \right|^{2} \approx \eta^{K\pi\pi}(\tau^{K\pi\pi}) \left| \sum_{a}^{\mathsf{waves}} \tilde{\mathcal{T}}^{\mathfrak{p}}_{a} \Psi^{K\pi\pi}_{a}(\tau^{K\pi\pi}) \right|^{2}$$

Experimentally measured intensity distribution

$$\mathcal{I}(\tau^{K\pi\pi}) = \sum_{\mathbf{p}} \left| \sum_{a}^{\mathsf{waves}} \mathcal{T}_{a}^{\mathbf{p}} \Psi_{a}^{K\pi\pi}(\tau^{K\pi\pi}) \right|^{2}$$

$$\mathcal{I}_{ ext{measured}}( au^{K\pi\pi}) = \eta^{K\pi\pi}( au^{K\pi\pi})\mathcal{I}( au^{K\pi\pi})$$

- ls the set of  $K^-\pi^-\pi^+$  partial waves sufficient?
  - ► Automatic wave-set selection using model-selection techniques

Treating the  $\pi^-\pi^-\pi^+$  and Other Backgrounds



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Experimentally measurable quantities are spin-density matrix elements

- $\blacktriangleright$  Transition amplitudes  $\mathcal{T}_a^p$  are only effective parameters
- $\blacktriangleright$  Cannot determine  $\mathcal{T}_a^p$  of individual processes
- ➡ Cannot separate different processes

Treating the  $\pi^-\pi^-\pi^+$  and Other Backgrounds



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- Large number of fit parameters:  $N_{\text{para}} = N_{\text{r}}(2N_{\text{waves}} N_{\text{r}})$
- Sufficient rank of spin-density matrix must be determined
  - ▶ Rank two needed to describe pure  $\pi^-\pi^-\pi^+$  Monte Carlo sample using  $K^-\pi^-\pi^+$  partial waves
  - Used rank three to model  $K^-\pi^-\pi^+$  sample







- Spin-density matrix  $\rho_{ab}(m_{K\pi\pi}, t')$  measured in partial-wave decomposition
- Model spin-density matrix in resonance-model fit

$$\hat{
ho}_{ab}(m_{K\pi\pi},t')=\hat{
ho}_{ab}^{K\pi\pi}(m_{K\pi\pi},t')+\hat{
ho}_{ab}^{3\pi}(m_{K\pi\pi},t')+\hat{
ho}_{ab}^{
m Bkg}(m_{K\pi\pi},t')$$



$$\hat{\mathcal{T}}_a^z(m_{K\pi\pi},t') = \sum_{k\in\mathbb{S}_a} \mathcal{K}(m_{K\pi\pi},t')^k \mathcal{C}_a^{K\pi\pi}(t') \mathcal{D}_k(m_{K\pi\pi};\zeta_k)$$

- Dynamic functions  $\mathcal{D}_k(m_{K\pi\pi}; \zeta_k)$ 
  - For resonances: rel. Breit-Wigner
  - For non-resonant terms:  $\mathcal{D}_k^{\mathrm{NR}}(m_{K\pi\pi}; a_k, c_k) = (m_{K\pi\pi} m_{\mathrm{thr}})^{a_k} e^{-b(c_k)\tilde{q}_k^2(m_{K\pi\pi})}$
- "Coupling amplitudes":  ${}^{k}C_{a}^{z}(t')$ 
  - Independent coupling amplitude for each t' bin
- Kinematic factor  $K(m_{K\pi\pi}, t')$
- Coherently summed over all assumed model components



$$\hat{\mathcal{T}}_{a}^{z}(m_{K\pi\pi},t') = \sum_{k\in\mathbb{S}_{a}} \mathcal{K}(m_{K\pi\pi},t')^{k} \mathcal{C}_{a}^{K\pi\pi}(t') \mathcal{D}_{k}(m_{K\pi\pi};\zeta_{k})$$

- Dynamic functions  $\mathcal{D}_k(m_{\kappa\pi\pi}; \zeta_k)$ 
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#### $3\pi$ spin-density matrix

$$\hat{\rho}_{ab}^{\pi\pi\pi}(m_{K\pi\pi},t') = \left| \mathcal{C}^{\pi\pi\pi} \right|^2 \rho_{ab}^{\pi\pi\pi}(m_{K\pi\pi},t')$$

•  $\rho_{ab}^{\pi\pi\pi}(m_{K\pi\pi}, t')$  obtained from PWD of  $\pi^-\pi^-\pi^+$  pseudodata sample

- $\blacktriangleright$   $m_{K\pi\pi}$  dependence fixed
- t' dependence fixed
- Rel. strength between partial waves fixed (freed in a study)
- One global real-valued yield parameter  $|C^{\pi\pi\pi}|^2$



#### Background spin-density matrix

- ▶ Additional incoherent contribution form other processes:  $K^-K^-K^+$ , ...
- ► Transition amplitudes modeled by non-resonant parameterizations for each partial wave  $\hat{\mathcal{T}}_{a}^{\text{eBKG}}(m_{K\pi\pi}, t') = K(m_{K\pi\pi}, t') \ \mathcal{C}_{a}^{\text{eBKG}}(t') \mathcal{D}_{k_{a}}^{\text{eBKG}}(m_{K\pi\pi}; a_{k_{a}}, c_{k_{a}})$



- $\blacktriangleright$   $\chi^2$  fit of the real and imaginary parts of the spin-density matrix
  - Taking into account correlations between spin-density matrix elements
  - Shape parameters ( $m_0$ ,  $\Gamma_0$ , ...) and coupling amplitudes are free parameters
- For the main fit, we performed 2000 fit attempts with random start-parameter values for the shape parameters, e.g. mass and width parameters, and the coupling and branching amplitudes.
- Start-parameter ranges for the shape parameters are chosen according to previous measurements (see note)
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$$\mathcal{I}(\tau, m_{K\pi\pi}, t') = \left| \sum_{a \in \mathbb{W}(m_{K\pi\pi}, t')} \mathcal{T}_{a}(m_{K\pi\pi}, t') \Psi_{a}(\tau; m_{K\pi\pi}) \right|^{2}$$

### Challenge: Find the "best" set of waves that describes the data

- If the wave set is too large
  - ➡ Starting to describe statistical fluctuations
- If waves that contribute to the data are missing
  - ➡ Intensity can be wrongly attributed to other waves
  - ➡ Model leakage



#### Infer wave set from data

- Systematically construct large set of allowed partial waves
  - ➡ "Wave pool"
- Fit wave pool to data
  - Impose penalty on  $|\mathcal{T}_a|^2 \Rightarrow$  regularization
  - Suppress insignificant waves
- Select waves that significantly contribute to data
  - "Best" subset of waves that describe the data





▶  $\pi^{-}\pi^{-}\pi^{+}$  Monte Carlo mock data set with 126 partial waves

- Massive overfitting
- Almost all waves pick up intensity

#### Courtesy F. Kaspar, TUM

Fitting wave pool of 753 waves

# Wave-Set Selection





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$$\ln \mathcal{L}_{\rm fit} = \ln \mathcal{L}_{\rm extended} + \sum_{a}^{\rm waves} \ln \mathcal{L}_{\rm reg}(|\mathcal{T}_a|; \{c_{\rm para}\})$$

#### LASSO/L1 regularization<sup>1</sup>

$$\ln \mathcal{L}_{\mathrm{reg}}(|\mathcal{T}_a|;\lambda) = -\lambda |\mathcal{T}_a|$$

- Maximum at  $|\mathcal{T}_a| = 0$
- Well established<sup>2</sup>
- "Smoothing" at  $|\mathcal{T}_a| = 0$

 $<sup>\</sup>frac{1}{2}$  Robert Tibshirani. "Regression Shrinkage and Selection via the Lasso". In: Journal of the Royal Statistical Society. Series B 58.1 (1996) Baptiste Guegan et al. "Model selection for amplitude analysis". In: JINST 10.09 (2015), P09002



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$$\mathcal{T}_{\rm a}| \rightarrow \sqrt{|\mathcal{T}_{\rm a}|^2 + \varepsilon}$$



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 $\lambda = 0.3$  $arepsilon = 10^{-5}$ 

- Bias also on large transition amplitudes
- Some additional waves
- Some waves missing

#### Courtesy F. Kaspar, TUM

S. Wallner



### Generalized Pareto<sup>1</sup>

$$\ln \mathcal{L}_{\mathrm{reg}}(|\mathcal{T}_{\boldsymbol{a}}|; \Gamma, \zeta) = -\frac{1}{\zeta} \ln \left[ 1 + \zeta \frac{|\mathcal{T}_{\boldsymbol{a}}|}{\Gamma} \right]$$

- Wave intensities spread over orders of magnitudes
- Use logarithmic prior
  - ➡ Heavy-tailed
  - ➡ Less bias on large waves
- ▶ LASSO-like for  $|\mathcal{T}_a| \rightarrow 0$
- ► "Smoothing" at  $|\mathcal{T}_a| = 0$  $|\mathcal{T}_a| \rightarrow \sqrt{|\mathcal{T}_a|^2 + \varepsilon}$



<sup>1</sup> Artin Armagan, David B. Dunson, and Jaeyong Lee. "Generalized double Pareto shrinkage". In: Statistica Sinica (2013). doi: 10.5705/ss.2011.048





- Less bias on large transition amplitudes
- $\blacktriangleright$  Clear kink in intensity distribution to smoothing scale  $\ \Rightarrow$  Selection
- Less additional waves
- Some small waves missing

#### Courtesy F. Kaspar, TUM



## "Cauchy"

$$\ln \mathcal{L}_{\mathrm{reg}}(|\mathcal{T}_{a}|;\Gamma) = -\ln \left[1 + \frac{|\mathcal{T}_{a}|^{2}}{\Gamma_{a}^{2}}\right]$$

- Logarithmic prior
- ▶ L2-like for  $|\mathcal{T}_a| \to 0$









- Less bias on large transition amplitudes
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- Few additional waves
- Few small waves missing

#### Courtesy F. Kaspar, TUM

S. Wallner



#### Wave pool

- ▶ Spin  $J \le 7$
- Angular momentum  $L \leq 7$
- Positive naturality of exchange particle
- 12 isobars
  - $[K\pi]_{S}^{K\pi}$ ,  $[K\pi]_{S}^{K\eta}$ ,  $K^{*}(892)$ ,  $K^{*}(1680)$ ,  $K_{2}^{*}(1430)$ ,  $K_{3}^{*}(1780)$
  - $[\pi\pi]_{s}$ ,  $f_0(980)$ ,  $f_0(1500)$ ,  $\rho(770)$ ,  $f_2(1270)$ ,  $\rho_3(1690)$

 $\Rightarrow$  "Wave pool" of 596 waves

"only" 720 k events



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## Regularization

$$\ln \mathcal{L}_{ ext{reg}}(|\mathcal{T}_{a}|; arGamma) = - \ln \left[ 1 + rac{|\mathcal{T}_{a}|^{2}}{arGamma_{a}^{2}} 
ight]$$

- Use Cauchy regularization
- Scale of |T<sub>a</sub>| depends on experimental acceptance
  - Apply penalty on expected number N<sub>a</sub> of observed events

 $\Gamma_{s} = \frac{\Gamma}{\sqrt{\eta_{s}}} \Rightarrow \frac{|\mathcal{T}_{s}|^{2}}{\Gamma_{s}^{2}} = \frac{\bar{N}_{s}}{\Gamma^{2}}$ 





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## COMPASS

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## COMPASS

#### Regularization

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#### Imposing continuity of the wave set

- Wave-set inferred independently for each  $(m_{K\pi\pi}, t')$  cell
- Impose continuity of the wave set in  $m_{K\pi\pi}$  by adding additional regularization term

$$\ln \mathcal{L}_{\text{cont}}(\{\mathcal{T}_{a}(m_{K\pi\pi},t')\};\lambda) = \sum_{j=i-3}^{j=i+3} \lambda \left|\mathcal{T}_{a}(m_{K\pi\pi},t')(m_{K\pi\pi}^{j+1}) - \mathcal{T}_{a}(m_{K\pi\pi},t')(m_{K\pi\pi}^{j})\right|^{2},$$

which suppresses fluctuations among neighboring  $m_{K\pi\pi}$  bins





#### Wave-set size

- ▶ 5 to 90 waves per  $(m_{K\pi\pi}, t')$  cell
- Larger wave set for larger binning in  $m_{K\pi\pi}$
- Larger wave set in t' bins with more events





- Selection of large signals
- as well as of signals at per-mil level





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- as well as of signals at per-mil level

Searching for Exotic Strange Mesons with  $J^P = 0^-$ 





### PDG

- ▶ *K*(1460) and *K*(1830)
- ► K(1630)
  - Unexpectedly small width of only  $16 \text{ MeV}/c^2$
  - $\blacktriangleright$   $J^P$  of K(1630) unclear

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  - But,  $m_{K\pi\pi} \lesssim 1.5 \, {
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- Second peak at about 1.7 GeV/c<sup>2</sup>
  - K(1630) signal with 8.3 σ statistical significance
     Accompanied by rising phase
- Weak signal at about 2.0 GeV/c<sup>2</sup>

K(1830) signal with 5.4  $\sigma$  statistical significance





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total resonance model, resonances, non-resonant,  $\pi\pi\pi$  background, effective background



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- ► K(1830) parameters in good agreement with LCHb measurement [PRL 118 (2017) 022003]
- Realistic K(1630) width of about 140 MeV/ $c^2$





- Indications for 3 excited K from a single analysis
- Quark-model predicts only two excited states: potentially K(1460) and K(1830)
- ➡ K(1630) supernumerary signal
- Solution Candidate for exotic non- $q\bar{q}$  state; other explanations possible ( $K^*(892) \omega$  threshold nearby)

Searching for Exotic Strange Mesons with  $J^P = 0^-$ 





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Searching for Exotic Strange Mesons with  $J^P = 0^-$ 





Searching for Exotic Strange Mesons with  $J^P = 0^-$ 







#### 500 0**-→**ρK ACCMOR 400 $K^{-}\pi^{-}\pi^{+}$ from ACCMOR ≥ 300 ¥ ▶ Potential K(1630) signal already in ACCMOR analysis / 20 Events 200 100 1 00 1.20 1 40 1.60 1.80 MKTT

Limited by kinematic range

WA03 (CERN), 200 000 events, ACCMOR, Nucl. Phys. B 187 (1981)

Searching for Exotic Strange Mesons with  $J^P = 0^-$ 



### $K^-\pi^-\pi^+$ from ACCMOR

• Potential K(1630) signal already in ACCMOR analysis

#### $K^{-}\pi^{-}\pi^{+}$ from LHCb

- ▶ Measurement of  $D^0 \to K^{\mp} \pi^{\pm} \pi^{\pm} \pi^{\mp}$  at LHCb
  - Study strange mesons in  $K\pi\pi$  subsystem
  - MIPWA of  $J^P = 0^-$  amplitude
  - Potential signal above  $1.6 \,\mathrm{GeV}/c^2$
  - Limited by kinematic range



Searching for Exotic Strange Mesons with  $J^P = 0^-$ 



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#### PDG

►  $K_2^*(1430)$  well known resonance





- ▶ Signal in K<sup>\*</sup><sub>2</sub>(1430) mass region
- In different decays
  - ρ(770) K D
  - K\*(892) π D
- In agreement with previous measurements
- Cleaner signal in COMPASS data







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Cleaner signal in COMPASS data





WA03 (CERN), 200 000 events, ACCMOR, Nucl. Phys. B 187 (1981)

Partial Waves with  $J^P = 2^+$ 





- $K_2^*(1430)$  parameters consistent with previous observations
- ▶ Better agreement with PDG average values for neutral  $K_2^*(1430)$

Partial Waves with  $J^P = 2^-$ 





#### PDG

- Established  $K_2(1770)$  and  $K_2(1820)$
- $\blacktriangleright$  K<sub>2</sub>(2250) need further confirmation

A+ Ag>it



- ▶ Simultaneously fit 4 waves with  $J^P = 2^-$
- 1.8 GeV/c<sup>2</sup> peak modeled by K<sub>2</sub>(1770), K<sub>2</sub>(1820)
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Partial Waves with  $J^P = 2^-$ 





### $K_2(1770)$ and $K_2(1820)$

- ▶ Two states were considered by only three measurements ACCMOR, LASS, LHCb
- Only LHCb measurement could confirm two states (3  $\sigma$  statistical significance)
- We observe two sates with  $11 \sigma$  statistical significance

Partial Waves with  $J^P = 2^-$ 





#### *K*<sub>2</sub>(2250)

- Studied so far mainly in  $\overline{\overline{A}}^{\circ}\overline{p}$  final states
- First simultaneous measurement of  $K_2(1770)$ ,  $K_2(1820)$ , and  $K_2(2250)$
- Resonance parameters consistent with previous observations

Partial Waves with  $J^P = 2^-$ 





Partial Waves with  $J^P = 2^-$ 









### PDG





Signal K<sub>4</sub><sup>\*</sup>(2045) signal in K<sup>\*</sup>(892) π and ρ(770) K decays





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A+ Ag >it

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- Also, real and imaginary parts of interference terms described well, including their magnitude
- ▶ Intensities and real and imaginary parts of interference terms not directly related as  $\operatorname{Rank}[\rho_{ab}] > 1$  $|\rho_{ab}| \neq \sqrt{|\rho_{aa}| |\rho_{bb}|}$ 
  - Analysis artifacts in intensities of small waves, which are the least constrained by data
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▶ Also structure in  $\pi^-\pi^+$  and  $K^-\pi^+$  subsystems

- Successive 2-body decay via  $\pi^-\pi^+$  /  $K^-\pi^+$  resonance called isobar
- Also structure in angular distributions





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 $m_{K^-\pi^-} \, [{\rm GeV}/c^2]$ 




### Kinematic Distribution of $K^-\pi^-\pi^+$ Events































































### Systematic Studies of the Partial-Wave Decomposition $_{\mbox{\tiny Leakage Waves}}$





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- Unexpected low-mass enhancement in 3<sup>+</sup> 1<sup>+</sup> K\*(892) π D wave
- Similar to dominant 1<sup>+</sup> wave
- Sensitive to systematic effects
- Decay amplitudes of different J<sup>P</sup> are orthogonal
- Event selection requires to identify one of the two negative particles
  - Limited acceptance due to limited kinematic range of final-state PID
- Loss of orthogonality taking acceptance into account
  - Reduced differentiability of certain partial waves
- Only a sub-set of partial waves affected





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#### • $K^-\pi^-\pi^+$ and $\pi^-\pi^-\pi^+$ similar experimental footprint

#### Distinguishable only by

- Beam particle identification
- Final-state particle identification
- Excellent beam PID:
  - Expect small contamination from beam  $\pi^-$
- Final-state PID does not suppress π<sup>-</sup>π<sup>-</sup>π<sup>+</sup> background
  - ► Non-negligible  $\pi^{-}\pi^{-}\pi^{+}$  background in  $K^{-}\pi^{-}\pi^{+}$  sample of about 7 %
  - ⇒ Dominant background in  $K^-\pi^-\pi^+$  sample





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#### • Well established model for $\pi^- + p \rightarrow \pi^- \pi^- \pi^+ + p$

- From very same data set
- Measured with high precision
- Acceptance corrected
- Generate  $\pi^-\pi^-\pi^+$  Monte Carlo sample
- Mis-interpret  $\pi^-\pi^-\pi^+$  Monte Carlo events as  $\mathcal{K}^-\pi^-\pi^-$ 
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  - Study  $\pi^{-}\pi^{-}\pi^{+}$  background in individual  $K^{-}\pi^{-}\pi^{+}$  partial waves





- Significant contribution to waves with  $\rho(770)$  isobar
- ▶  $\pi^{-}\pi^{-}\pi^{+}$  produces peaking structures
- Largest relative contribution to  $2^+ 1^+ \rho(770) \ K D$  wave
- Small contribution to waves with  $K^*(892)$  isobar
- Also significant contribution to waves with f<sub>2</sub>(1270) and K<sup>\*</sup><sub>2</sub>(1430) isobars
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 $K^{-}\pi^{-}\pi^{+}$  data.  $\pi^{-}\pi^{-}\pi^{+}$  pseudo data



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