

# Structure of the $\Lambda(1405)$ From Photoproduction at Give

Reinhard Schumacher Carnegie Mellon University

With Nilanga Wickramaarachchi (Catholic Univ.) & Peter Hurck (Univ. Glasgow) & Other GlueX Collaborators



6-18-24 NSTAR2024, York, UK



- $\blacksquare$  Place of the  $\Lambda(1405)$  in the world
- GlueX measurement for two final states
- $\blacksquare$  K-matrix fits with one or two  $\Lambda(1405)$  resonances & two scattering states
- Pole positions and branchings of the two resonances

### Recall the Motivation

- What is the place of the  $\Lambda(1405)$  in baryonic physics?
  - It's too light, compared to  $\Lambda(1520),$  in the quark model.
  - Close to the  $N\overline{K}$  mass threshold molecular/penta aspect.
  - Decays to  $\Sigma \pi$ , but MUST also decay to  $N\overline{K}$ .
- Chiral unitary models, CPT, LQCD (& others) predict <u>two</u> I=0 states in  $\Lambda(1405)$  mass range.
- GlueX has the best data set making it cleanly in photoproduction:  $\gamma p \rightarrow K^+ \Lambda(1405) \rightarrow K^+ \{\Sigma^0 \pi^0\} \rightarrow K^+ \{p \ K^-\} (> N\overline{K} \text{ threshold})$

### Chiral Unitary Models



Fig. 1. Trajectories of the poles in the scattering amplitudes obtained by changing the SU(3) breaking parameter x gradually. At the SU(3) symmetric limit (x = 0), only two poles appear, one is for the singlet and the other for the octets. The symbols correspond to the step size  $\delta x = 0.1$ .

- SU(3) baryons irreps 1+8<sub>s</sub>+8<sub>a</sub> combine with 0 - Goldstone bosons to generate:
  - Two octets and a singlet of  $\frac{1}{2^{-}}$  baryons dynamically generated in the SU(3) limit
  - SU(3) breaking leads to two S = -1, I = 0 poles near 1405 MeV
    - ~1420 mostly  $N\overline{K}$
    - ~1390 mostly  $\Sigma~\pi$
  - Possible weak I=1 pole also predicted

# Y Rele positions from the literature



- Higher pole ~1430 MeV couples more strongly to  $N\overline{K}$ , lower pole ~1390 MeV couples more to  $\Sigma\pi$
- Many theorists believe:  $N\overline{K}$  quasi-bound state submerged in  $\Sigma\pi$  continuum: coupled-channel dynamics
- Most data from low-energy NK scattering, kaonic atoms not very sensitive to  $\Sigma\pi$  pole position

GlueX approach is new and different

## Pole positions from the literature

#### B. Cid-Mora, HIM Mainz, MENU 2023 Lattice QCD Theory





J. Bulava et al., Phys Rev Lett 132, 051901 (2024)J. Bulava et al., Phys Rev D 109, 014511 (2024)



GLUE System at JLab

The GlueX Beamline and Detector NIM A 987, 164807 (2021)

- ~ 12 GeV e<sup>-</sup> beam converted to
  6.5 11.6 GeV photon beam
- 30 cm LH2 target
- ~ 1.5 T Solenoidal magnet
- Drift chambers
- Start counter/Time of flight
- Electromagnetic Calorimeters
- This analysis: Data from "Phase I" runs 2017, 2018



### GlueX Competitive Advantages:

• GlueX has world's best data set making  $\Lambda(1405)$  cleanly in photoproduction:  $\gamma p \rightarrow K^+ \Lambda(1405)$ 

 $\rightarrow K^+ \{\Sigma^0 \ \pi^0\}$  (pure I = 0, no I = 1 contamination)

$$\rightarrow \mathbf{K}^{+} \{ \{ \gamma \Lambda \} \pi^{0} \} \rightarrow \mathbf{K}^{+} \gamma \mathbf{p} \pi^{-} \gamma \gamma$$

• GlueX also has:  $\gamma p \rightarrow K^+ \Lambda(1405)$ 

 $ightarrow K^+ \ \{p \ K^-\}$  (when above  $N\overline{K}$  threshold)

- Do K-matrix fit to both final states together
  - Never done before...

## Experimental Method I



#### • $\Sigma^0 \pi^0$ channel

- Exclusive kinematic fit to beam photon & final state {K<sup>+</sup> γ p π<sup>-</sup> γ γ} particles
- Constrain  $\Lambda$  and  $\pi^0$ masses, but not  $\Sigma^0$  mass, in each  $\Sigma^0 \pi^0$  mass bin
- Background removal fit under  $\Sigma^0$  in each  $\Sigma^0\pi^0$  mass bin
- Use common GlueX acceptance & photon flux normalizations

### Experimental Method II



#### • $\Sigma^0 \pi^0$ channel

- Clean detection of  $\Lambda(1405)$  &  $\Lambda(1520)$
- Evident pK<sup>-</sup> threshold effect
- Smooth acceptance

- *p*K<sup>-</sup> channel
  - $\Lambda(1520)$  sits on top of  $\Lambda(1405)$  tails
  - Good, smooth acceptance

### Cross Sections Differential in Mass



### K-matrix formalism\* (outline sketch)

- We have two resonances,  $\Lambda(1405)_A$  and  $\Lambda(1405)_B$ , each coupled to  $\Sigma^0 \pi^0$  and p K<sup>-</sup>. The  $\Lambda(1520)$  also decays to the same final states.
- Assume J=  $\frac{1}{2}$  L=0 states do no interfere with J=3/2 L=2 state

$$\widehat{T} = \left(I - i\widehat{K}\rho\right)^{-1}\widehat{K}$$
$$K = \sum_{\alpha} \frac{m_{\alpha}\Gamma_{\alpha}(m)}{m_{\alpha}^2 - m^2}$$

$$\widehat{K_{ij}} = \sum_{\alpha} \frac{\gamma_{\alpha i} \gamma_{\alpha j} m_{\alpha} \Gamma^0_{\alpha}}{m_{\alpha}^2 - m^2} B^l_{\alpha i} B^l_{\alpha j}$$

Lorentz-invariant T-matrix (2 in x 2 out)

Sum over resonances A & B ; real function, preserves unitarity of *T* 

Invariant K-matrix for available decay modes  $i, j = \{\Sigma^0 \pi^0, p K^-\}$ 

\* à la S.U. Chung et al., Ann. Physik 4,404 (1995).

### K-matrix formalism\* (outline sketch)

$$\widehat{P}_{i} = \sum_{\alpha} \frac{\beta_{\alpha} \gamma_{\alpha i} m_{\alpha} \Gamma_{\alpha}^{0}}{m_{\alpha}^{2} - m^{2}}$$
$$\widehat{F}_{i} = \left(I - i\widehat{K}\rho\right)^{-1}\widehat{P}_{i}$$

$$\frac{d\sigma_i(m)}{dm} \sim \rho_i \left| \widehat{F}_i(m) \right|^2$$

Photoproduction vector for decay modes *i* ; same sum over poles as K matrix

Production exp't replacement of *T* matrix "formation exp't" for decay mode *i* 

Fit to experimental data for decay mode *i* 

$$T_{11}(m) = \rho_{\Sigma^0 \pi^0}(m) \widehat{T}_{11}(m)$$

Compute *T*-matrix to be tested for unitarity and to find "*T*-matrix poles"

\* à la S.U. Chung et al., Ann. Physik 4,404 (1995).

13

### K-matrix formalism - issues

- $\blacksquare$  Ignore the possibility of  $\eta\Lambda$  and  $K\Xi$  decays
- Poles "A" & "B" are below threshold for pK<sup>-</sup> channel
- Define "branching ratio" & "branching fractions" in terms of fitted  $\Sigma \pi$  and  $N\overline{K}$  final states
  - Calculate using mass-integrated cross sections to each final state computed for each resonance separately
  - Not computed in terms of pole residues
    - (threshold issues make this difficult)

### $\frac{1}{K^{+}}$ 2-Pole K-matrix Fit to $\Lambda(1405)$ A,B



#### • $\Sigma^0 \pi^0$ channel

- Solid fit to data
- Dashed each A,B resonance separately
- Dotted fit to data:
  - full K-matrix fit with coherent  $\Lambda(1405){\rm A,B}$  states
  - prior to convolving 7.8 MeV GlueX mass resolution

#### ■ pK<sup>-</sup> channel

- Solid fit to data:
  - 2.0 MeV GlueX mass resolution
- Dashed coherent tail of  $\Lambda(1405)A,B$  states
- Dotted incoherent high-mass background
  - 3<sup>rd</sup> order polynomial
- $0.00 \le t \le 1.50 \text{ GeV}^2$  (full range)
- $\Lambda(1520)$  cross section agreement < 5%

### **1-Pole K-matrix Fit to** $\Lambda(1405)B$



- $\Sigma^0 \pi^0$  channel
  - Solid fit to data
  - Dashed single  $\Lambda(1405)$  resonance

#### ■ pK<sup>-</sup> channel

- Solid fit to data
- Dashed  $pK^-$  tail of  $\Lambda(1405)$  state
- Dotted incoherent high-mass background
  - 3<sup>rd</sup> order polynomial
- $0.00 \le t \le 1.50 \text{ GeV}^2$  (full range)
- Poorer fit than 2-pole ansatz: especially in critical threshold region



Λ(1520) PDG

#### arLambda(1520) POLE POSITION

REAL PART 1517 to 1518 ( $\approx 1517.5$ ) MeV

 $-2 \times \text{IMAGINARY PART}$  14 to 18 ( $\approx 16$ ) MeV ( $\rightarrow \sim 2 \times 8 \text{ MeV}$ )

GlueX (preliminary):

 $(1516.5 \pm 0.3) - i (8.3 \pm 0.1) \text{ MeV}$ 

Good agreement with PDG: suggests the GlueX method is sound

# Check Unitarity of the Amplitudes

R. A. Sch./ CMU



- Argand diagram and squared-magnitude for the  $\Sigma^0 \pi^0$  amplitude (red)
  - Two  $\Lambda(1405)$  resonances with  $\Sigma^0 \pi^0$  and  $pK^-$  initial/final states.
  - Each amplitude stays properly bounded.
- Separately,  $\Lambda(1520)$  is a single pK<sup>-</sup> amplitude (blue)

### **Cross Sections & Branching Fractions**



### Systematic Uncertainty in Pole Locations

- Variations on the procedures:
  - Nominal fit defines central values
  - Shift mass bins by 4 MeV and redo the fits: sensitive to  $\ensuremath{pK^{\scriptscriptstyle -}}$  threshold
  - Tighter kinematic fit cut: sensitive to any backgrounds between  $\Lambda 's$
  - Looser kinematic fit cut: " " " " "
  - Slice data by t-bin:  $0.00 < t < 0.35 \text{ GeV}^2$
  - " " :  $0.35 < t < 0.60 \,\mathrm{GeV^2}$
  - $\blacksquare \qquad " \qquad " \qquad : 0.60 < t < 1.50 \, \mathrm{GeV^2}$
  - Resum t-bins to total: sensitive to acceptance modeling
  - Rescale  $\Lambda(1520)$  cross section to exactly match in both channels
- How to combine all this information ?
  - Use "Nominal" fit for central values...
  - Use standard deviations of all fits to define systematic spreads

### Systematic Tests Summary I



- Pole Positions
  - Nominal
  - Shifted bins
  - Tighter KINFIT
  - Looser KINFIT
  - t-bin 1
  - t-bin 2
  - t-bin 3
  - Resum t-bins
  - Rescale A(1520) cross section

### Systematic Tests Summary II



- Branching Fractions
  - Nominal
  - Shifted bins
  - Tight KINFIT
  - Loose KINFIT
  - t-bin 1
  - t-bin 2
  - t-bin 3
  - Resum t-bins
  - Rescale A(1520) cross section

# Y 💦 The Landscape Including GlueX Data/Fit

GlueX (preliminary):

Red dot positions and error bars to be finalized...

Recent (year  $\geq$  2000) predictions: M. Mai – Eur. Phys. J. Spec. Top. 230 6, 1593, (2021)

Thresholds:  $\Sigma^0 \pi^0$  1327.62 MeV *p K*<sup>-</sup> 1431.95 MeV







GLUE

Project each resonance separately onto the real axis • Integrate  $\Sigma^0 \pi^0$  & pK<sup>-</sup> modes across full mass range:  $\sigma_{ii}$ • Define branching ratio as  $\sigma_{\Sigma 0 \pi 0}$  /  $\sigma_{pK-}$  , etc.

# Summary/Conclusions

- First measurement of the  $\Lambda(1405)$ s decaying two separate ways:  $\Sigma^0\pi^0$  &  $pK^-$
- K-matrix fit to two intermediate resonances: A & B
- Two-pole ansatz is superior to single-pole ansatz
- Branching ratio/fractions defined and presented
- To do: systematics to be finalized

GlueX acknowledges the support of several funding agencies and computing facilities (http://gluex.org/thanks)





### **Y** $\mathbf{K}^{+}$ Rescaling of pK<sup>-</sup> and $\Sigma^{0}\pi^{0}$ Data

- Trust that isospin holds exactly
- Trust that PDG branching fractions are all OK
- Part I: Scale (Peter's)  $\Lambda(1520) \rightarrow p K^-$  cross section to match (Nilanga's)  $\Lambda(1520) \rightarrow \Sigma^0 \pi^0$  cross section
  - $p K^-$  branch to  $\Lambda(1520)$  total: x 1/(0.45/2) (scale up)
  - Total  $\Lambda(1520)$  to  $\Sigma^0 \pi^0$ :
- x 0.42 / 3 (scale down)
- Net p K<sup>-</sup> rescaling factor = 0.6222

### **Rescaling of** $pK^-$ and $\Sigma^0\pi^0$ Data

- The p K<sup>-</sup> "background" gets rescaled, too... so...
- Part II: Scale (Reinhard's) computed model  $\Lambda(1405) \rightarrow p K^-$ <u>tail</u> to match rescaled  $\Lambda(1520) \rightarrow \Sigma^0 \pi^0$ 
  - We see only  $\Sigma^0 \pi^0$  but not  $\Sigma^+ \pi^- \& \Sigma^- \pi^+ : \times 3.0$  (scale up)
    - (this is the total strength of  $\Lambda(1405)$  production)
  - Equal  $\Lambda(1405)$  decay to  $nK^0$  and  $pK^-$ : x 0.5 (scale down)
  - Adjust for the  $pK^-$  data rescaling: x 0.622
  - Net pK<sup>-</sup> calculated tail curve rescaling = 0.9333

### **Rescaling of** $pK^-$ and $\Sigma^0\pi^0$ Data

- Our quoted  $\Lambda(1405)$  branching ratio/fractions are for isospin-corrected  $\Sigma \pi$  and  $N\overline{K}$
- Part III: Scale measured cross sections to account for isospin
  - We measure (Nilanga)  $\Lambda(1405) \rightarrow \Sigma^0 \pi^0$ , not  $\Sigma^+ \pi^- \& \Sigma^- \pi^+$ , so correct for isospin: × 3 (scale up)
  - Computed  $N\overline{K}$  tail (Reinhard) from  $\Lambda(1405) \rightarrow \Sigma^0 \pi^0$ , again correct for isospin: × 3 (scale up)
    - (K-matrix fit does not, in itself, distinguish NK modes)