

# The study of $N^*(1535)$ and $N^*(1650)$ from the lattice data

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Phys.Rev.D 108 (2023) 9, 094519



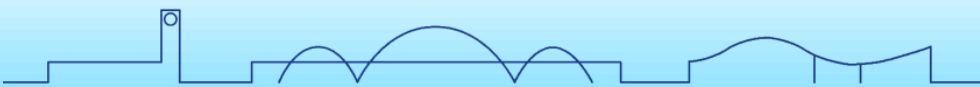
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# Outline

- Introduction of HEFT : Derek's talk
- Analysis of  $N^*(1535)$  and  $N^*(1650)$
- Summary and Outlook



# Introduction of HEFT

J. M. M. Hall *etc.* PRD 87(2013), 094510  
 J.-j. Wu *etc.* PRC90 (2014), 055206  
 Y. Li *etc.* PRD 101(2020), 114501  
 PRD 103(2021), 094518

$$H = H_0 + H_I$$

$$H_0 = \sum_{i=1,n} |B_i\rangle m_i \langle B_i| + \sum_{\alpha} |\alpha(k_{\alpha})\rangle \left[ \sqrt{m_{\alpha 1}^2 + k_{\alpha}^2} + \sqrt{m_{\alpha 2}^2 + k_{\alpha}^2} \right] \langle \alpha(k_{\alpha})|$$

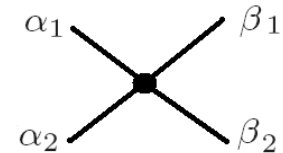
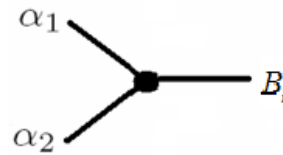
$|B_i\rangle$  bare state, bare mass  $m_i$

$|\alpha(k_{\alpha})\rangle$  non-interaction channels

$$H_I = \hat{g} + \hat{v}$$

$$\hat{g} = \sum_{\alpha} \sum_{i=1,n} \left[ |\alpha(k_{\alpha})\rangle g_{i,\alpha}^+ \langle B_i| + |B_i\rangle g_{i,\alpha} \langle \alpha(k_{\alpha})| \right]$$

$$\hat{v} = \sum_{\alpha,\beta} |\alpha(k_{\alpha})\rangle v_{\alpha,\beta} \langle \beta(k_{\beta})|$$



**T matrix**  
 (Phase Shifts,  
 inelasticity)

**Lattice  
 Spectrum**  
 3

**Resonance**  
 (Mass, Width, Pole position, Coupling)

**HEFT**



# What is the Bare State ?

$$H = H_0 + H_I$$

$$H_0 = \sum_{i=1,n} |B_i\rangle m_i \langle B_i| + \sum_{\alpha} |\alpha(k_{\alpha})\rangle \left[ \sqrt{m_{\alpha 1}^2 + k_{\alpha}^2} + \sqrt{m_{\alpha 2}^2 + k_{\alpha}^2} \right] \langle \alpha(k_{\alpha})|$$

$|B_i\rangle$  bare state, bare mass  $m_i$

$|\alpha(k_{\alpha})\rangle$  non-interaction channels

ANALOGY ( maybe inappropriate )

Quark  $\longrightarrow$  Person

Hadron  $\searrow$   
 $\longrightarrow$  Country



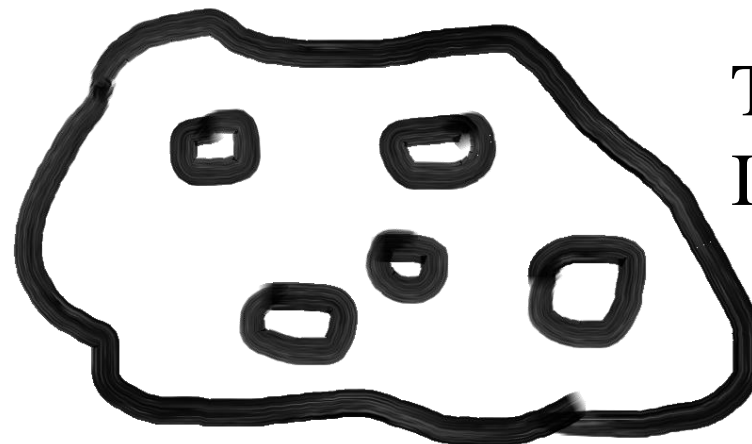
# Why Bare State ?

$$H = H_0 + H_I$$

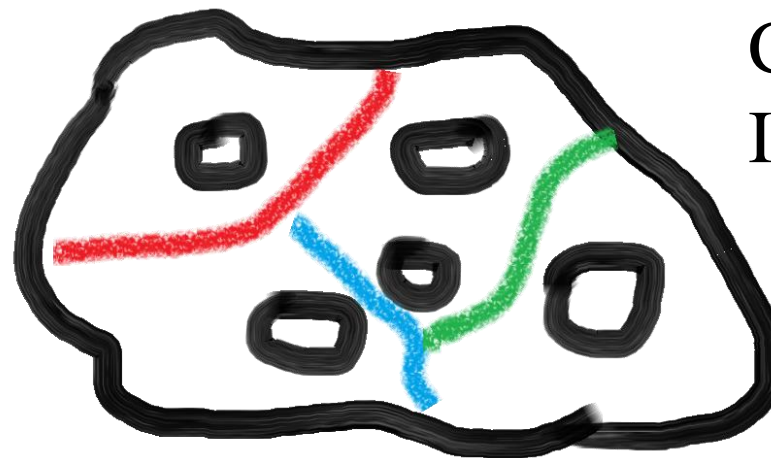
$$H_0 = \sum_{i=1,n} |B_i\rangle m_i \langle B_i| + \sum_{\alpha} |\alpha(k_{\alpha})\rangle \left[ \sqrt{m_{\alpha 1}^2 + k_{\alpha}^2} + \sqrt{m_{\alpha 2}^2 + k_{\alpha}^2} \right] \langle \alpha(k_{\alpha})|$$

$|B_i\rangle$  bare state, bare mass  $m_i$

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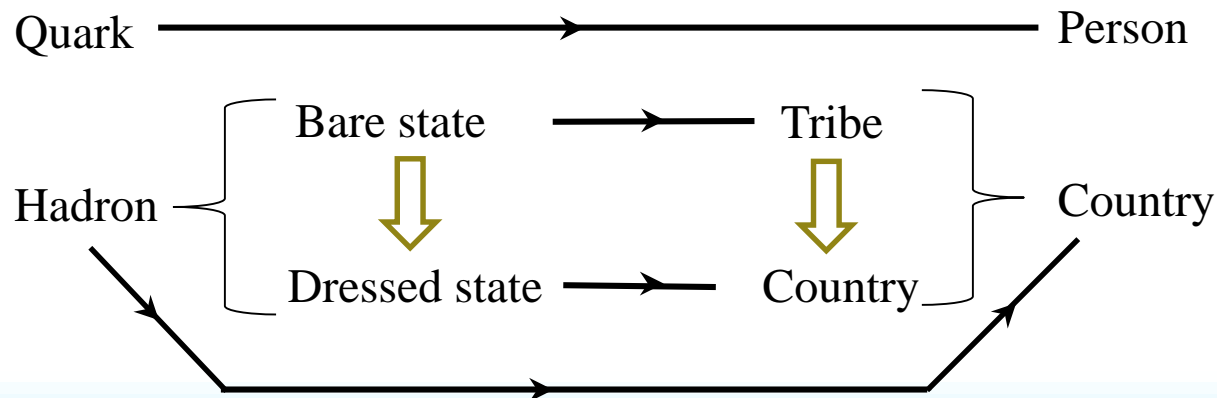


Tribes:  
Independent



Countries:  
Interacted

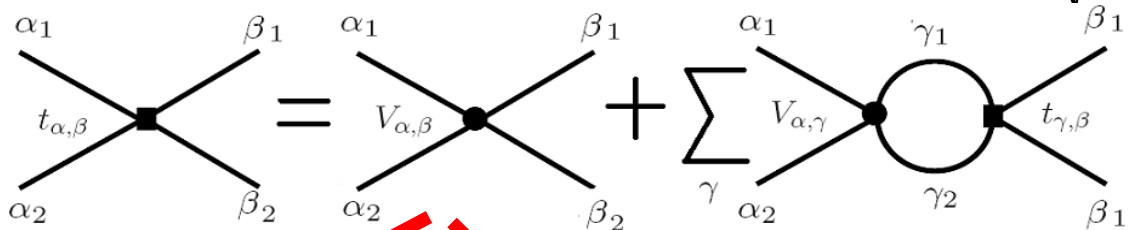
ANALOGY ( maybe inappropriate )



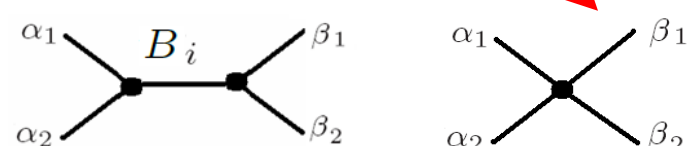
# Introduction of HEFT

- T Matrix:**

$$t_{\alpha,\beta}(k_\alpha, k_\beta, E) = V_{\alpha,\beta}(k_\alpha, k_\beta) + \sum_\gamma \int k_\gamma^2 dk_\gamma \frac{V_{\alpha,\gamma}(k_\alpha, k_\gamma) t_{\gamma,\beta}(k_\gamma, k_\beta, E)}{E - \sqrt{m_{\gamma 1}^2 + k_\gamma^2} - \sqrt{m_{\gamma 2}^2 + k_\gamma^2} + i\varepsilon}$$



**Resonance**  
(Mass, Width, Pole position, Coupling)

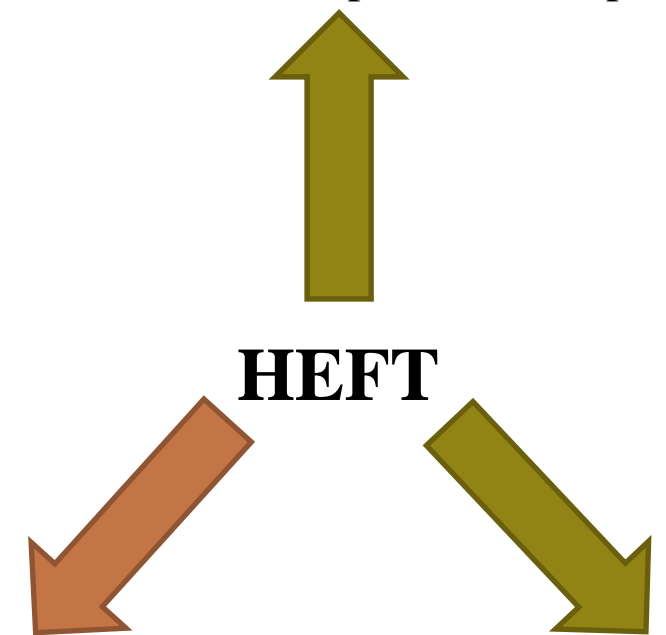


$$S_{\alpha,\beta} = 1 - i2\sqrt{\rho_\alpha} t_{\alpha,\beta}(k_{0\alpha}, k_{0\beta}, E) \sqrt{\rho_\beta}$$

$$\rho_\alpha = \frac{\pi k_{0\alpha} \sqrt{m_{\alpha 1}^2 + k_{0\alpha}^2} \sqrt{m_{\alpha 2}^2 + k_{0\alpha}^2}}{E}$$

$$\eta e^{2i\delta_\alpha} = S_{\alpha,\alpha}$$

$$g_{i,\alpha}^* \frac{1}{E - m_i} g_{i,\beta} \quad V_{\alpha,\beta}$$



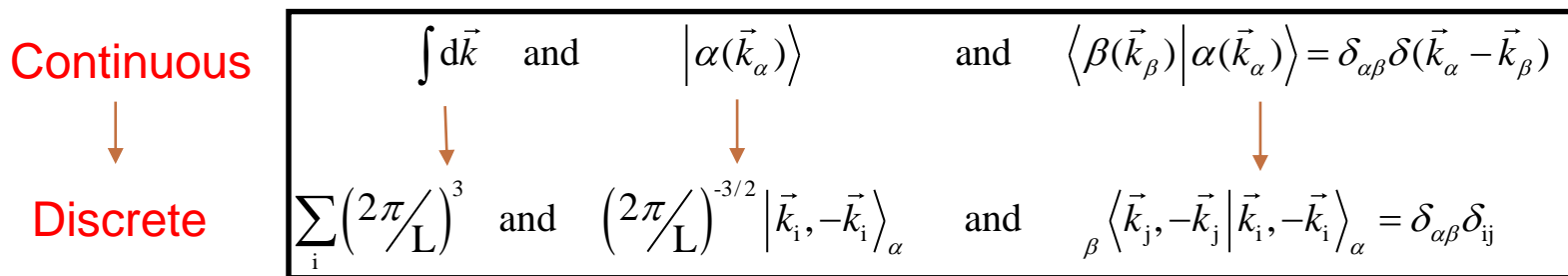
**T matrix**  
(Phase Shifts, inelasticity)

**Lattice Spectrum**

# Introduction of HEFT

J. M. M. Hall etc. PRD 87(2013), 094510  
 J.-j. Wu etc. PRC90 (2014), 055206

- Hamiltonian with discrete momentum



$$H_0 = \sum_{i=1,n} |B_i\rangle m_i \langle B_i| + \sum_{\alpha,i} |\vec{k}_i, -\vec{k}_i\rangle_\alpha \left[ \sqrt{m_{\alpha_B}^2 + k_\alpha^2} + \sqrt{m_{\alpha_M}^2 + k_\alpha^2} \right] \langle\vec{k}_i, -\vec{k}_i|$$

$$H_I = \sum_j (2\pi/L)^{3/2} \sum_{\alpha} \sum_{i=1,n} \left[ |\vec{k}_j, -\vec{k}_j\rangle_\alpha g_{i,\alpha}^+ \langle B_i| + |B_i\rangle g_{i,\alpha} \langle\vec{k}_j, -\vec{k}_j| \right] + \sum_{i,j} (2\pi/L)^3 \sum_{\alpha} |\vec{k}_i, -\vec{k}_i\rangle_\alpha v_{\alpha,\beta} \langle\vec{k}_j, -\vec{k}_j|$$

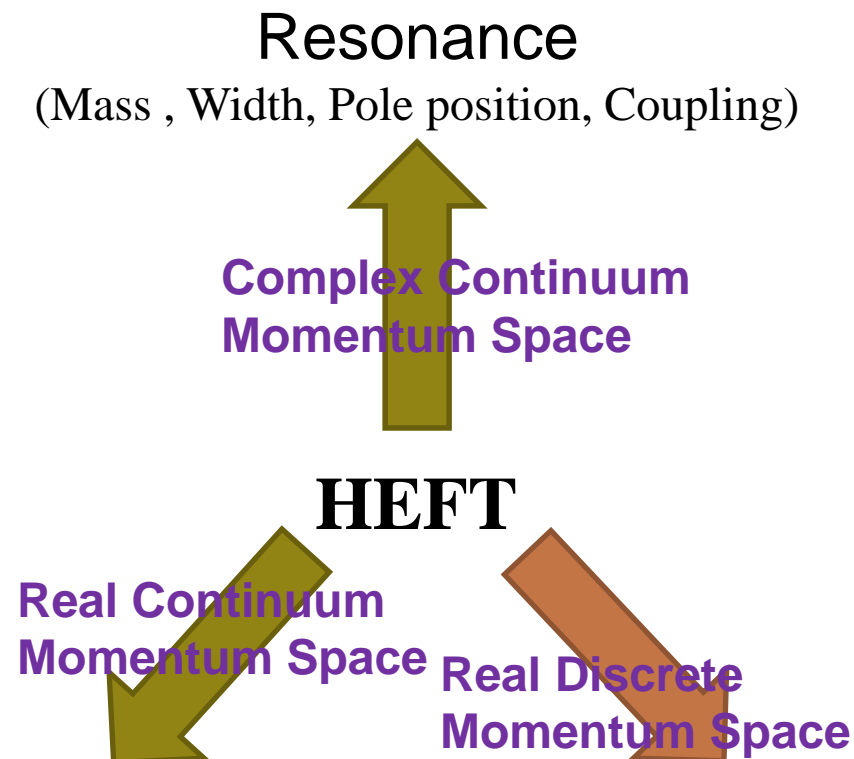
$$[H_0]_{N_c+1} = \begin{pmatrix} m_0 & 0 & \dots & 0 & 0 & \dots \\ 0 & \epsilon_1(k_0) & 0 & \dots & 0 & \dots \\ 0 & 0 & \epsilon_2(k_0) & \dots & 0 & \dots \\ 0 & 0 & 0 & \dots & 0 & \dots \\ 0 & 0 & 0 & \dots & \epsilon_{n_c}(k_0) & \dots \\ 0 & 0 & 0 & \dots & 0 & \epsilon_1(k_1) & \dots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix} \quad [H_I]_{N_c+1} = \begin{pmatrix} 0 & g_1^V(k_0) & g_2^V(k_0) & \dots & g_{n_c}^V(k_0) & g_1^V(k_1) & \dots \\ g_1^V(k_0) & v_{1,1}^V(k_0, k_0) & v_{1,2}^V(k_0, k_0) & \dots & v_{1,n_c}^V(k_0, k_0) & v_{1,1}^V(k_0, k_1) & \dots \\ g_2^V(k_0) & v_{2,1}^V(k_0, k_0) & v_{2,2}^V(k_0, k_0) & \dots & v_{2,n_c}^V(k_0, k_0) & v_{2,1}^V(k_0, k_1) & \dots \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \dots \\ g_{n_c}^V(k_0) & v_{n_c,1}^V(k_0, k_0) & v_{n_c,2}^V(k_0, k_0) & \dots & v_{n_c,n_c}^V(k_0, k_0) & v_{n_c,1}^V(k_0, k_1) & \dots \\ g_1^V(k_1) & v_{1,1}^V(k_1, k_0) & v_{1,2}^V(k_1, k_0) & \dots & v_{1,n_c}^V(k_1, k_0) & v_{1,1}^V(k_1, k_1) & \dots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

$$(H_0 + H_I) |\Psi\rangle = E |\Psi\rangle \quad \text{Eigen-Value} \longleftrightarrow \text{Lattice Spectrum}$$

Eigen-Vector

T matrix  
(Phase Shifts, inelasticity)

Lattice Spectrum



# Introduction of HEFT

**T matrix**  
(Phase Shifts,  
inelasticity)

**Lattice  
Spectrum**

**HEFT**

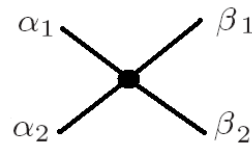
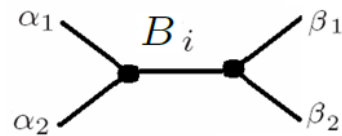
**Resonance**  
(Mass , Width, Pole  
position, Coupling)

**1. Build Model**

**2. Fix Para.**

**3. Extract Phys.**

**Unphysical  $\pi$   
mass ??**



HEFT:

1. Build a Hamiltonian model;

2. If Experimental data available, we fit  
Experimental data to fix the parameters in  
the model;

If Lattice data available (close to physical  
pion mass), we fit these data;

If both, we can use both of them constraint  
the model parameters.

If we only have Lattice data with  
unphysical pion mass, we need another  
parameter for the mass dependence, such  
as mass slope.

3. From the fixed Hamiltonian, we can study  
the properties of Resonance. Especially,  
from the eigenvector in the finite volume, we  
can estimate the internal structure of the  
hadron.

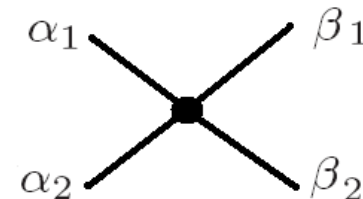
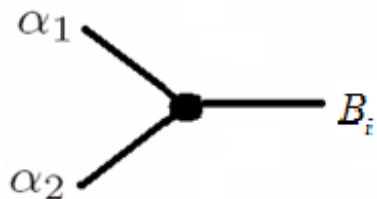




# N\*(1535)

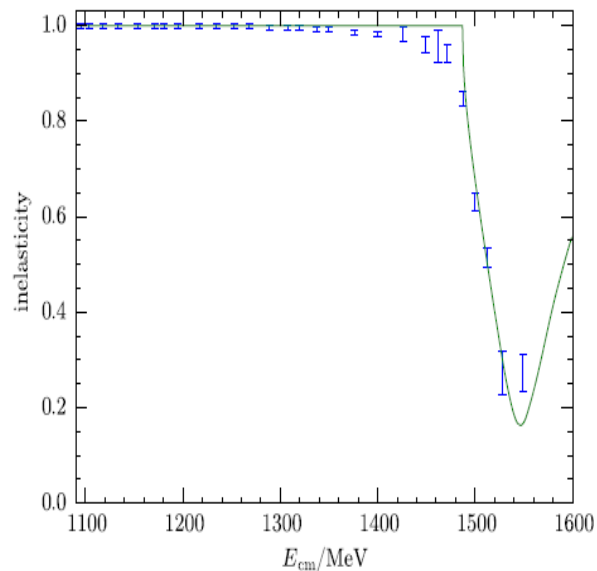
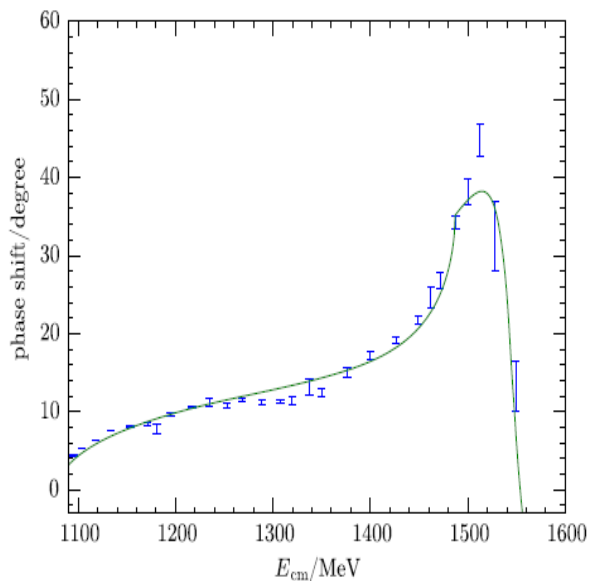
Zhan-wei Liu etc. Phys.Rev.Lett. 116 (2016) no.8, 082004

2 Channels:  $\pi N$  and  $\eta N$



$$G_{iN}^2(k) = \left( 3g_{N_0^*iN}^2 / 4\pi^2 f^2 \right) \omega_i(k) u^2(k)$$

$$\frac{3g_{\pi N}^S \tilde{u}(k) \tilde{u}(k')}{4\pi^2 f^2}$$



$$g_{\pi N}^S = -0.0608 \pm 0.0004$$

$$m_0 = 1601 \pm 14 \text{ MeV}$$

$$g_{N_0^* \pi N} = 0.186 \pm 0.006$$

$$g_{N_0^* \eta N} = 0.185 \pm 0.017$$

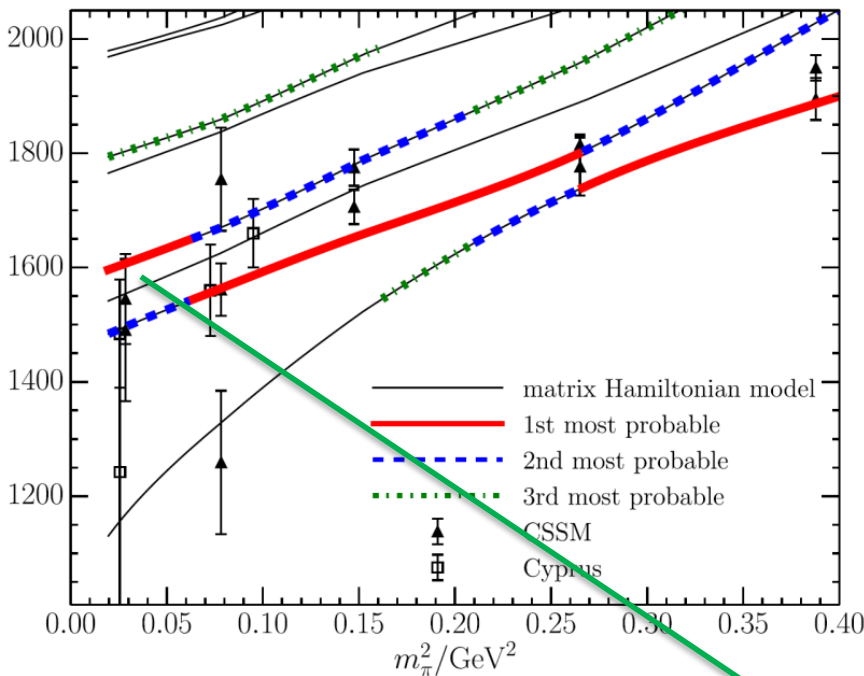
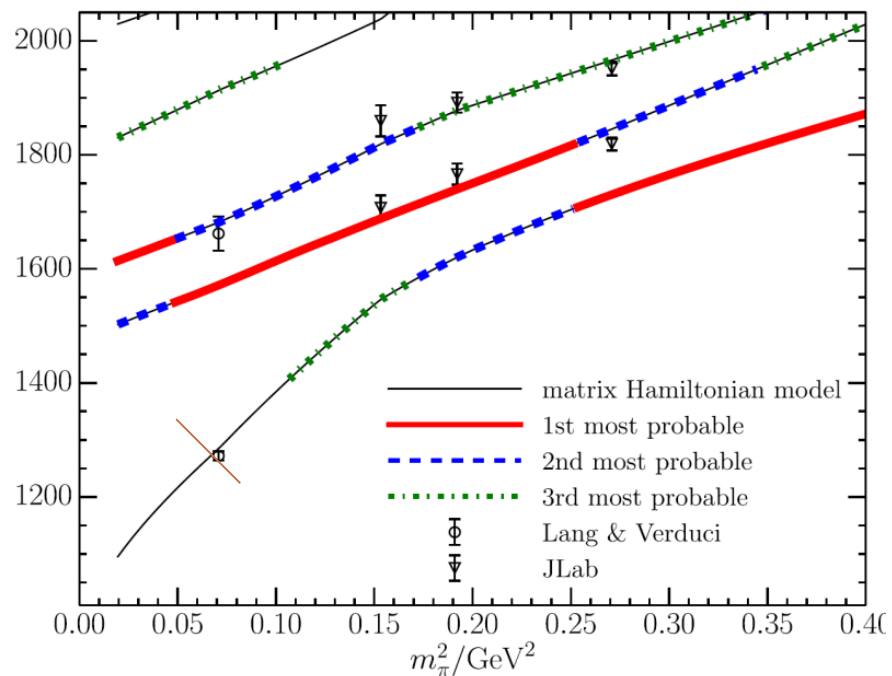
$$\chi_{\text{DOF}}^2 = 6.8$$

$$1531 \pm 29 - i88 \pm 2 \text{ MeV}$$

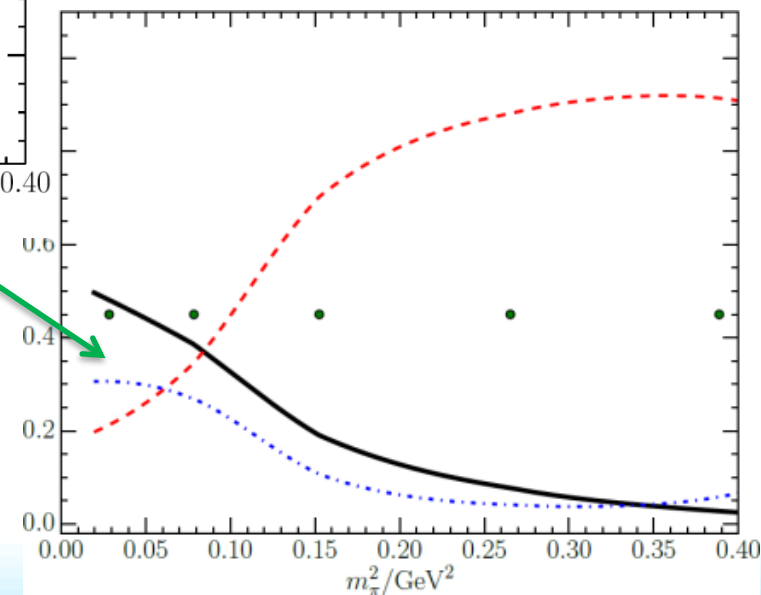


# $N^*(1535)$

Zhan-wei Liu etc. Phys.Rev.Lett. 116 (2016) no.8, 082004

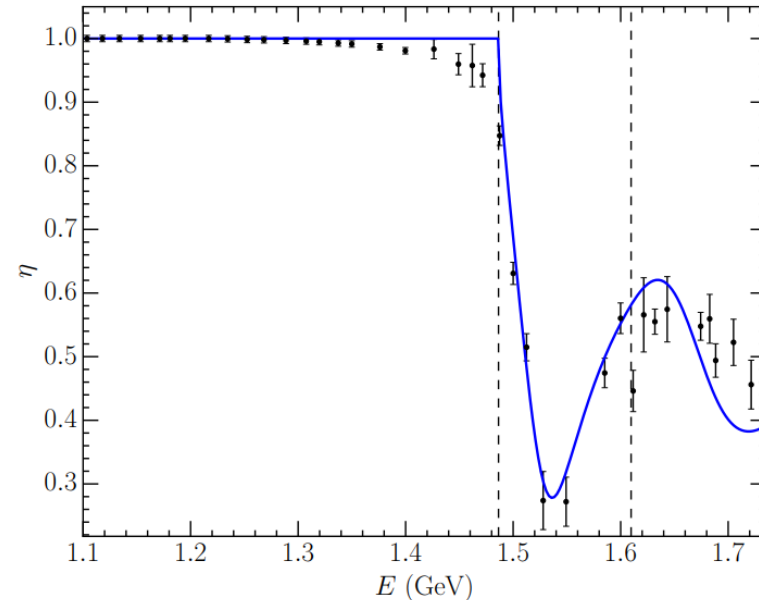
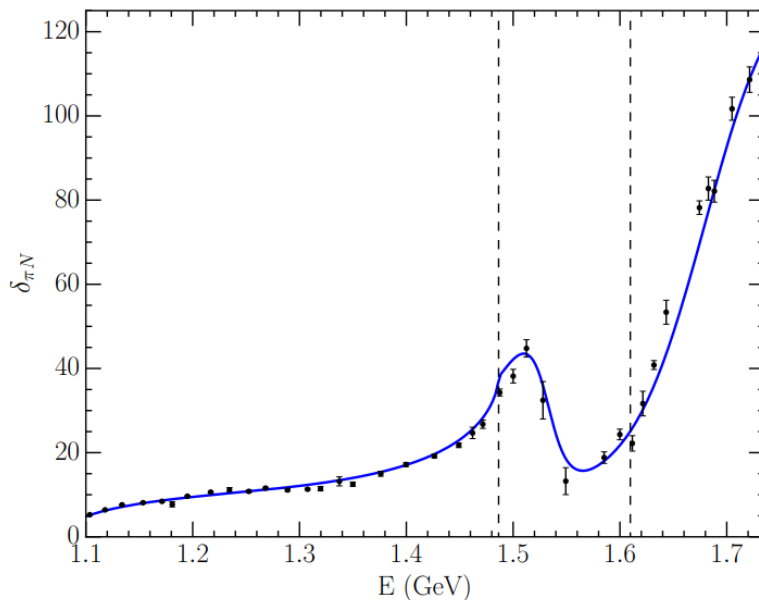


The main components (at least 50% ) of  $N^*(1535)$  is from the 3 quark core.



# For $N^*(1535)/ N^*(1650)$

Parameter	Value	Parameter	Value
$m_{N_1}^{(0)} / \text{GeV}$	1.6301	$m_{N_2}^{(0)} / \text{GeV}$	1.8612
$g_{\pi N}^{N_1}$	0.0898	$g_{\pi N}^{N_2}$	0.2181
$g_{\eta N}^{N_1}$	0.1525	$g_{\eta N}^{N_2}$	0.0009
$g_{K\Lambda}^{N_1}$	0.0000	$g_{K\Lambda}^{N_2}$	-0.2367
$\Lambda_{\pi N}^{N_1} / \text{GeV}$	1.2335	$\Lambda_{\pi N}^{N_2} / \text{GeV}$	1.4000
$\Lambda_{\eta N}^{N_1} / \text{GeV}$	1.2642	$\Lambda_{\eta N}^{N_2} / \text{GeV}$	0.9521
$\Lambda_{K\Lambda}^{N_1} / \text{GeV}$	...	$\Lambda_{K\Lambda}^{N_2} / \text{GeV}$	0.7283
$v_{\pi N, \pi N}$	-0.0655	$v_{\eta N, \eta N}$	-0.0245
$v_{\pi N, \eta N}$	0.0388	$v_{\eta N, K\Lambda}$	0.0320
$v_{\pi N, K\Lambda}$	-0.0757	$v_{K\Lambda, K\Lambda}$	0.1371
$\Lambda_{v, \pi N} / \text{GeV}$	0.6000	$\Lambda_{v, \eta N} / \text{GeV}$	0.9036
$\Lambda_{v, K\Lambda} / \text{GeV}$	0.6060		



$$G_{\alpha}^{N_i}(k) = \frac{\sqrt{3} g_{\alpha}^{N_i}}{2\pi f_{\pi}} \sqrt{\omega_{M_{\alpha}}(k)} u(k), \quad V_{\alpha\beta}(k, k') = \frac{3 v_{\alpha\beta}}{4\pi^2 f_{\pi}^2} \tilde{u}(k) \tilde{u}(k'), \quad \tilde{u}(k) = \frac{\omega_{\pi}(k) + m_{\pi}^{\text{phys}}}{\omega_{\pi}(k)} u(k).$$

We consider three channels:  
 $\pi N, \eta N, K\Lambda$

$$E_{N^*(1535)} = 1510 \pm 10 - (65 \pm 10)i \text{ MeV},$$

$$E_1 = 1500 - 50i \text{ MeV},$$

$$E_{N^*(1650)} = 1655 \pm 15 - (67 \pm 18)i \text{ MeV}.$$

$$E_2 = 1658 - 56i \text{ MeV},$$

EXP

HEFT

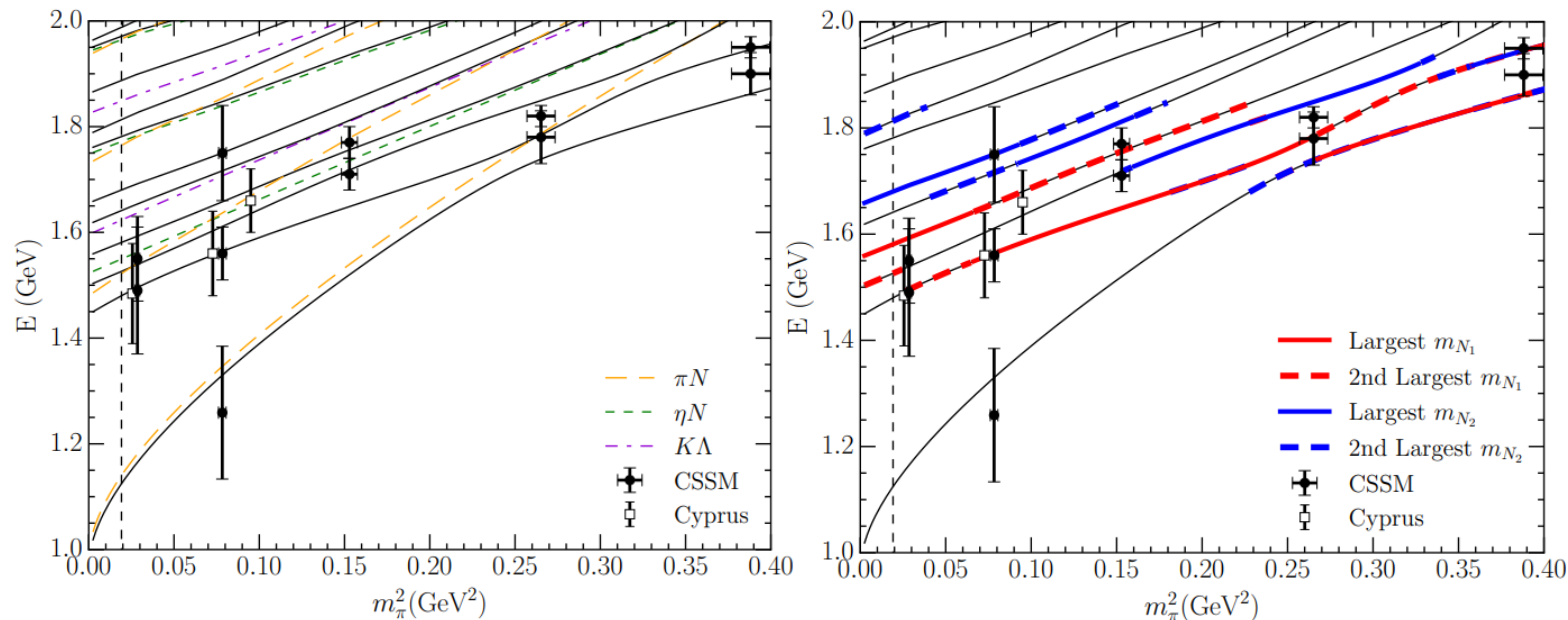
11



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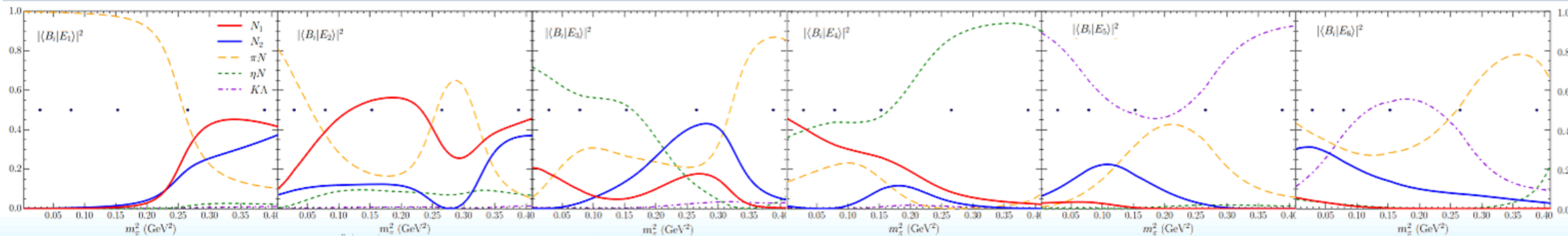
L ~ 3 fm **Fitting**



$$m_{N_i}(m_\pi^2) = m_{N_i}^{(0)} + \alpha_{N_i} (m_\pi^2 - m_\pi^2|_{\text{phys}}),$$

$$\alpha_{N_1} = 0.944 \text{ GeV}^{-1}, \quad \alpha_{N_2} = 0.611 \text{ GeV}^{-1}.$$

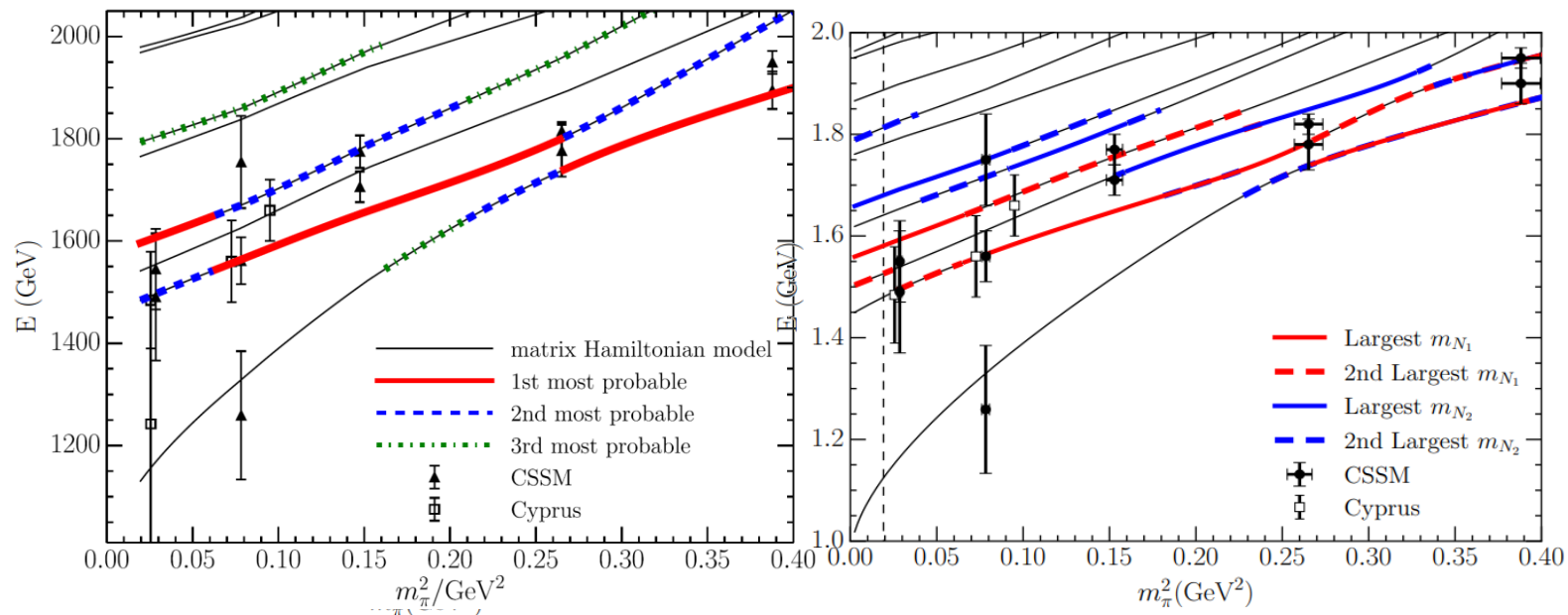
**Fitting**



Parameter	Value	Parameter	Value
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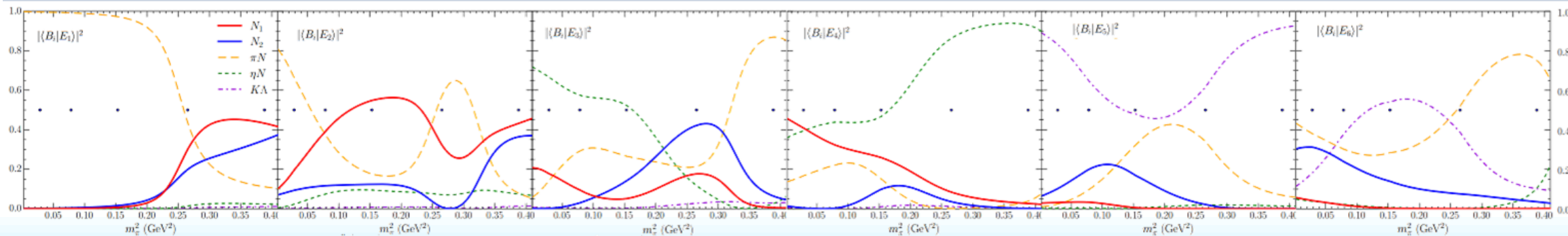
L ~ 3 fm **Fitting**



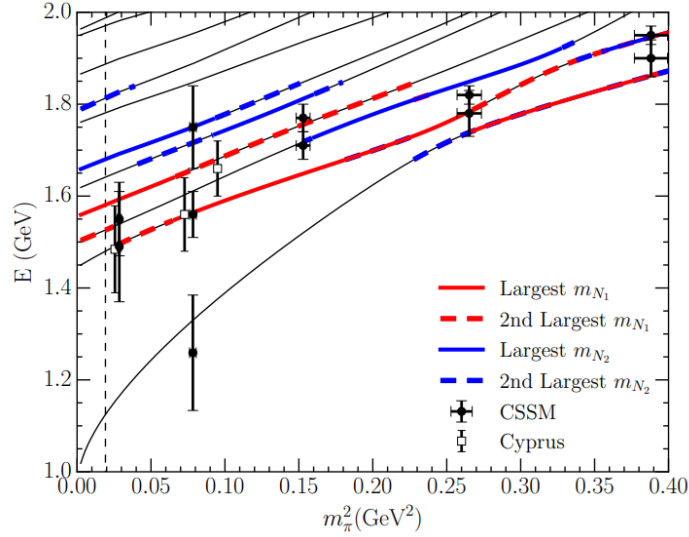
$$m_{N_i}(m_\pi^2) = m_{N_i}^{(0)} + \alpha_{N_i} (m_\pi^2 - m_\pi^2|_{\text{phys}}),$$

$$\alpha_{N_1} = 0.944 \text{ GeV}^{-1}, \quad \alpha_{N_2} = 0.611 \text{ GeV}^{-1}.$$

**Fitting**



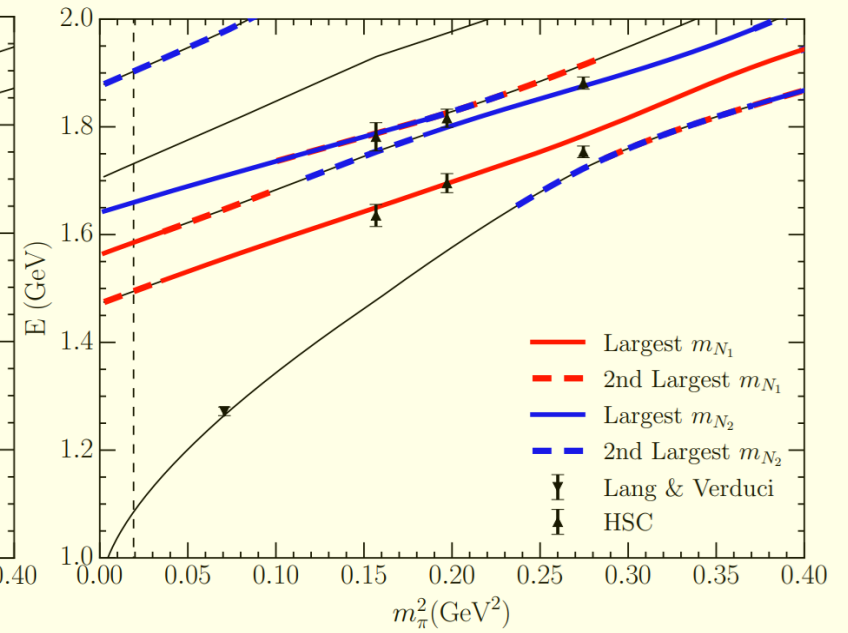
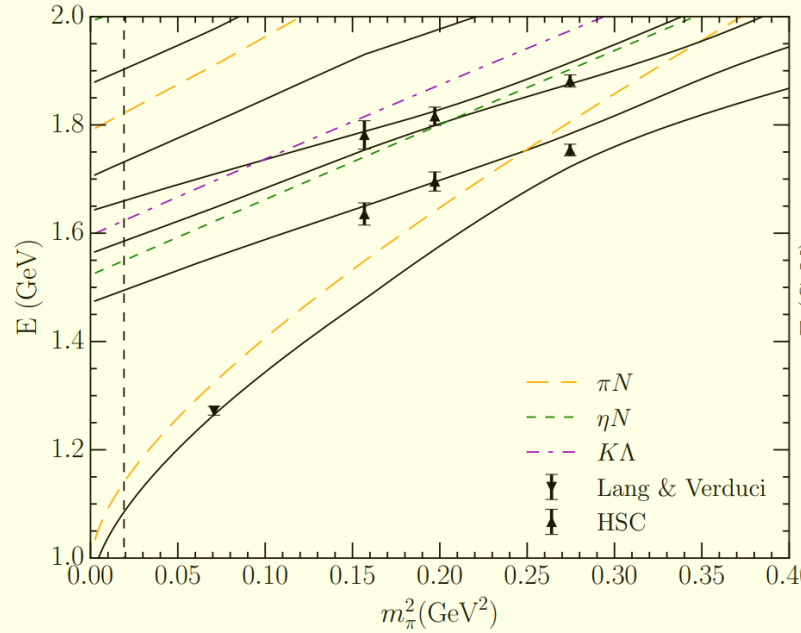
# L ~ 3 fm **Fitting**



# For N\*(1535)/ N\*(1650)

1. Scattering state  
3q state
2. Two clear states

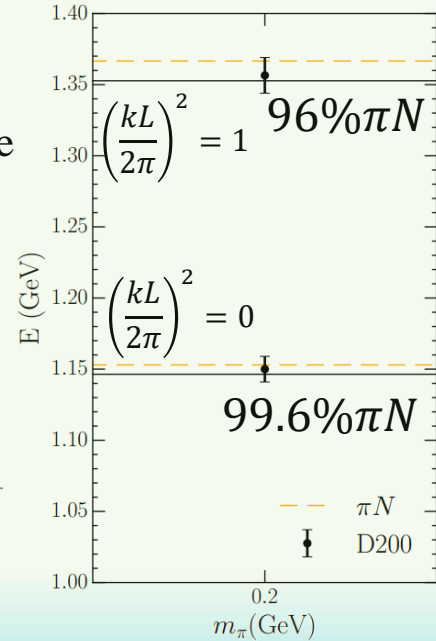
# L ~ 2 fm **Not Fit**



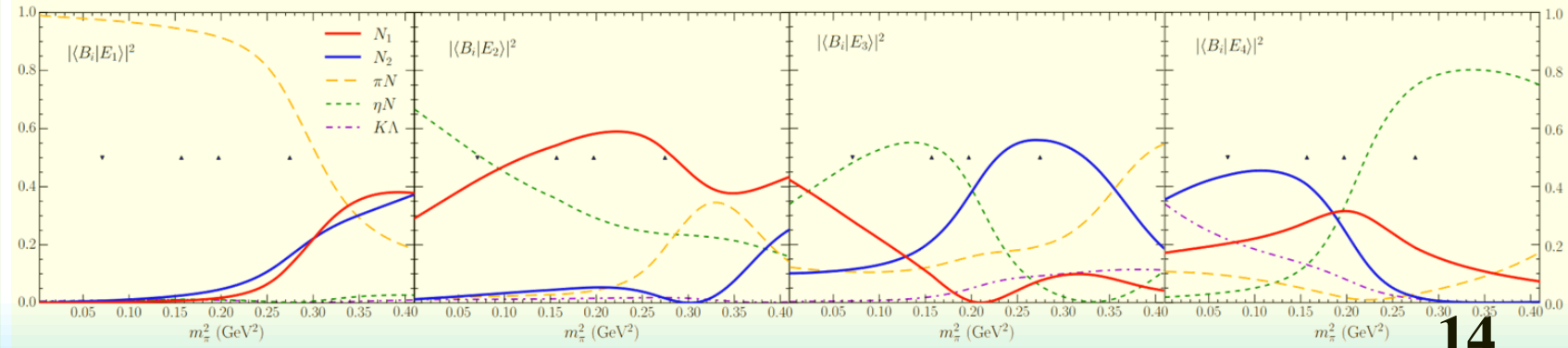
# L ~ 4.05 fm **Not Fit**

Coordinated Lattice Simulations (CLS) consortium

D200 ensemble



$a[\text{fm}]$	$(L/a)^3 \times T/a$	
0.0633(4)(6)	$64^3 \times 128$	
$am_\pi$	$am_K$	$am_N$
0.06617(33)	0.15644(16)	0.3148(23)
$af_\pi$	$af_K$	
0.04233(16)	0.04928(21)	



# For $N^*(1535)$

Define “Contamination Function”  
to compare HEFT VS LQCD

For One bare states:  $\bar{\chi}(0) |\Omega\rangle = |B_0\rangle$

Correlation Function:  $G_\chi(t, \mathbf{p}) = \sum_{\mathbf{x}} e^{-i\mathbf{p}\cdot\mathbf{x}} \langle \bar{\Omega} | \chi(\mathbf{x}, t) \bar{\chi}(0, 0) | \Omega \rangle$ ,

$$G_\chi(t) = \sum_i |\langle \Omega | \chi | E_i \rangle|^2 e^{-E_i t}, \quad G_{B_0}(t) = \sum_i |\langle B_0 | E_i \rangle|^2 e^{-E_i t}$$

Contamination function:  $C_{B_0}(t) = \frac{1}{G_{B_0}(t)} \sum_{i \neq B_0} |\langle B_0 | E_i \rangle|^2 e^{-E_i t}$

If  $|E_{B_0}\rangle$  is a ground state,  
 $C_{B_0} \sim 0$ .

If  $|E_{B_0}\rangle$  is an excited state,  
 $C_{B_0}(t)$  will have a minimal  
value as a function of  $t$ .



# For $N^*(1535)$

Define “Contamination Function”  
to compare HEFT VS LQCD

For One bare states:  $\bar{\chi}(0) |\Omega\rangle = |B_0\rangle$

Correlation Function:  $G_\chi(t, \mathbf{p}) = \sum_{\mathbf{x}} e^{-i\mathbf{p}\cdot\mathbf{x}} \langle \bar{\Omega} | \chi(\mathbf{x}, t) \bar{\chi}(0, 0) | \Omega \rangle$ ,

$$G_\chi(t) = \sum_i |\langle \Omega | \chi | E_i \rangle|^2 e^{-E_i t}, \quad G_{B_0}(t) = \sum_i |\langle B_0 | E_i \rangle|^2 e^{-E_i t}$$

Contamination function:  $C_{B_0}(t) = \frac{1}{G_{B_0}(t)} \sum_{i \neq B_0} |\langle B_0 | E_i \rangle|^2 e^{-E_i t}$

If  $|E_{B_0}\rangle$  is a ground state,  
 $C_{B_0} \sim 0$ .

If  $|E_{B_0}\rangle$  is a excited state,  
 $C_{B_0}(t)$  will have a minimal  
value as function of t.

For Two bare states:  $(\alpha^* \bar{\chi}_1 + \beta^* \bar{\chi}_2) |\Omega\rangle = \alpha^* |N_1\rangle + \beta^* |N_2\rangle$ ,

Correlation Function:  $G_j(\mathbf{p}, t) = \sum_{\mathbf{x}} e^{-i\mathbf{p}\cdot\mathbf{x}} \langle \bar{\Omega} | (\alpha_j \chi_1(\mathbf{x}, t) + \beta_j \chi_2(\mathbf{x}, t)) (\alpha_j^* \bar{\chi}_1(0) + \beta_j^* \bar{\chi}_2(0)) | \Omega \rangle$ .

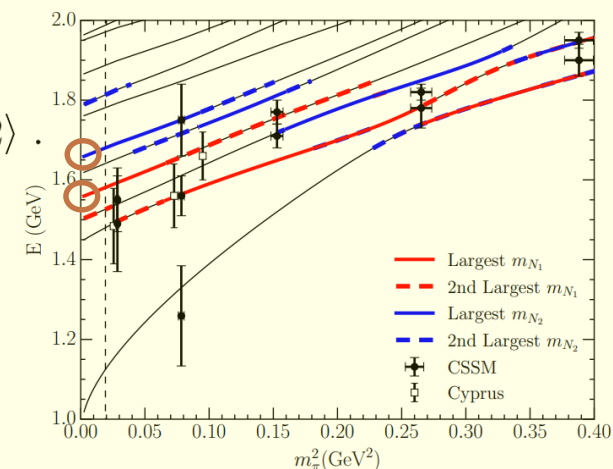
$$G_j(t) = \sum_i (\alpha_j \langle N_1 | + \beta_j \langle N_2 |) | E_i \rangle \langle E_i | (\alpha_j^* | N_1 \rangle + \beta_j^* | N_2 \rangle) e^{-E_i t},$$

$$= \sum_i |\alpha_j \langle N_1 | E_i \rangle + \beta_j \langle N_2 | E_i \rangle|^2 e^{-E_i t}.$$

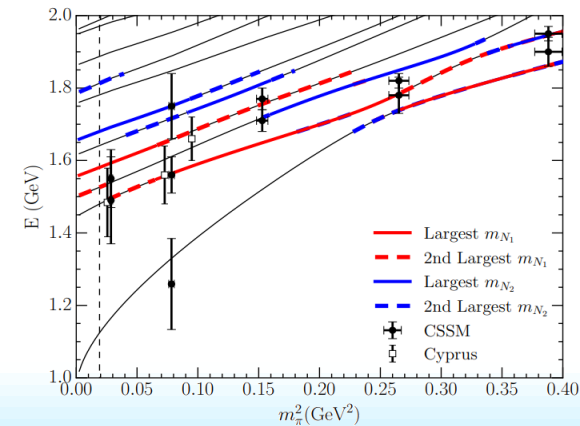
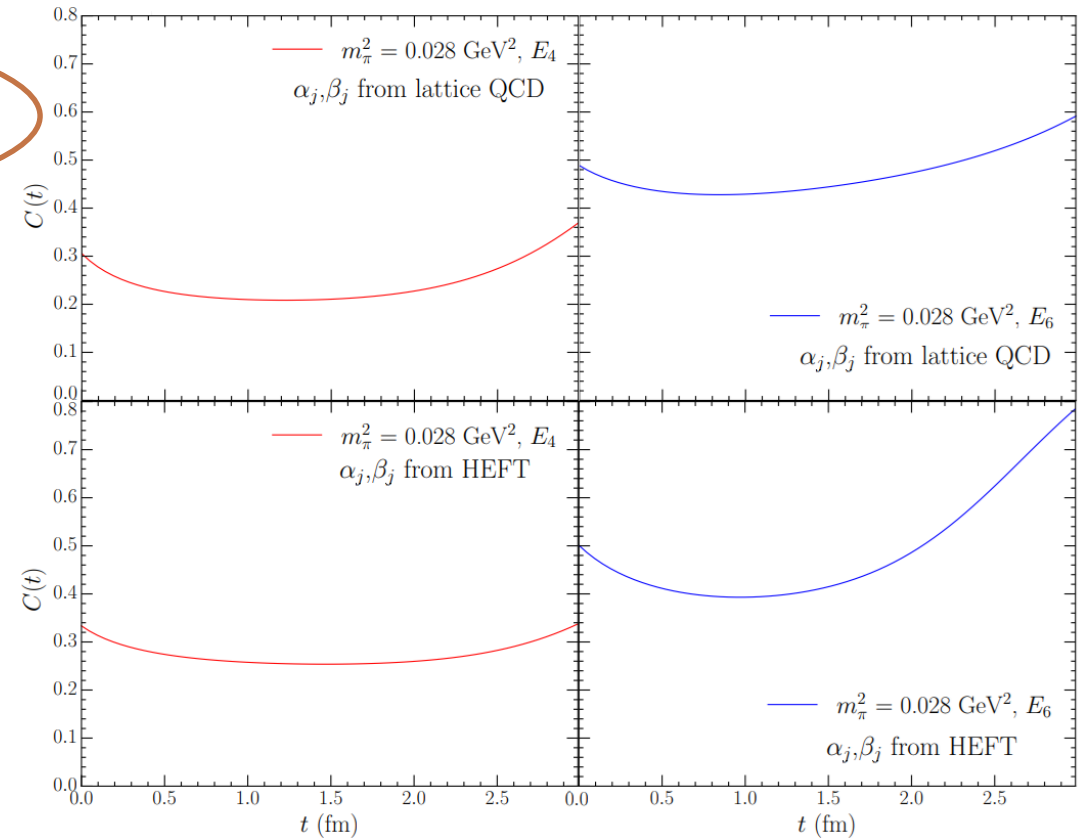
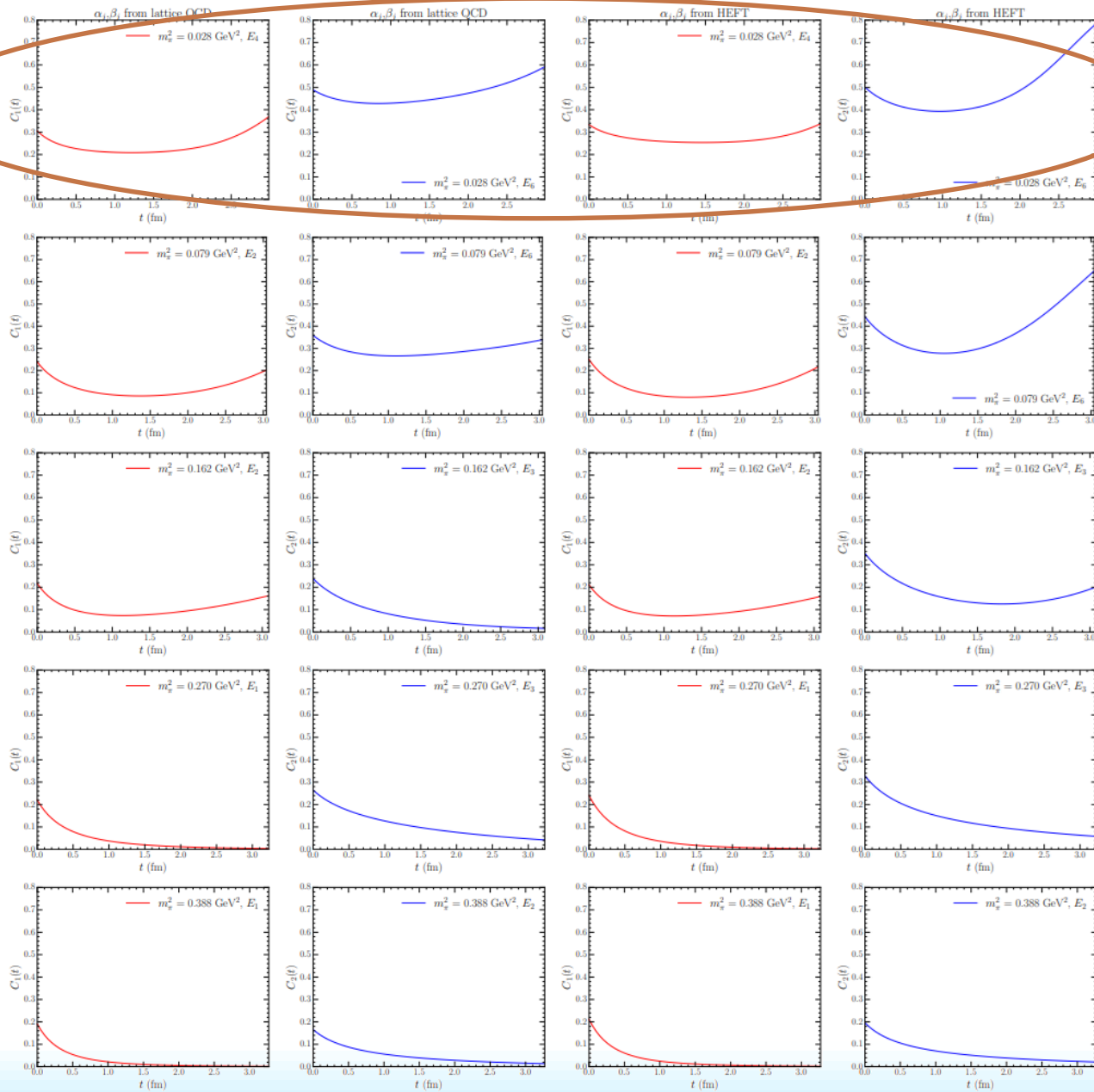
Contamination function:  $C_j(t) = \frac{1}{G_j(t)} \sum_{i \neq N_1, N_2} (\alpha_j \langle N_1 | E_i \rangle + \beta_j \langle N_2 | E_i \rangle)^2 e^{-E_i t}$

$$\alpha_1 = \langle N_1 | E_{N_1} \rangle, \quad \beta_1 = \langle N_2 | E_{N_1} \rangle,$$

$$\alpha_2 = \langle N_1 | E_{N_2} \rangle, \quad \beta_2 = \langle N_2 | E_{N_2} \rangle.$$





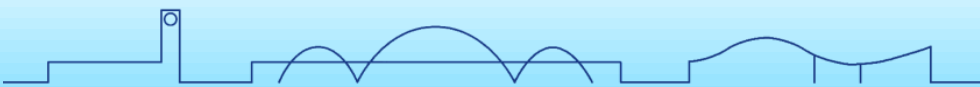


This result shows that HEFT is quiet consistent with Lattice data.



# Summary

- **Here we find that the interpretation of the two resonances as three-quark cores dressed by scattering-state dynamics is consistent with the  $L \sim 3$  fm lattice calculations.**
- **To extend to the  $L \sim 2$  fm and 4 fm are both quiet good.**
- **We define a “contamination function” related to the overlap of bare states with eigenstates, then we compare this function by Lattice input and HEFT results.**
- **All of these consistent comparisons show that the results of HEFT correctly reflect the structure of hadron from experimental and lattice data.**



# Outlook

- The most big problem is too many parameters.
- We have 21 free parameters! Data driven motivation.
- Now we want to improve our model by including  $K\Sigma$  channel, i.e., it will be four coupled channels.
- By using SU(3) symmetry, we set the fixed relationship of coupling constants of bare state with channels, and channels.
- Correspondingly, the fitting will be much hard.

Parameter	Value	Parameter	Value
$m_{N_1}^{(0)} / \text{GeV}$	1.6301	$m_{N_2}^{(0)} / \text{GeV}$	1.8612
$g_{\pi N}^{N_1}$	0.0898	$g_{\pi N}^{N_2}$	0.2181
$g_{\eta N}^{N_1}$	0.1525	$g_{\eta N}^{N_2}$	0.0009
<del><math>g_{K\Lambda}^{N_1}</math></del>	<del>0.0000</del>	<del><math>g_{K\Lambda}^{N_2}</math></del>	<del>-0.2367</del>
$\Lambda_{\pi N}^{N_1} / \text{GeV}$	1.2335	$\Lambda_{\pi N}^{N_2} / \text{GeV}$	1.4000
$\Lambda_{\eta N}^{N_1} / \text{GeV}$	1.2642	$\Lambda_{\eta N}^{N_2} / \text{GeV}$	0.9521
<del><math>\Lambda_{K\Lambda}^{N_1} / \text{GeV}</math></del>	<del>...</del>	<del><math>\Lambda_{K\Lambda}^{N_2} / \text{GeV}</math></del>	<del>0.7283</del>
$v_{\pi N, \pi N}$	-0.0655	$v_{\eta N, \eta N}$	-0.0245
$v_{\pi N, \eta N}$	0.0388	$v_{\eta N, K\Lambda}$	0.0320
$v_{\pi N, K\Lambda}$	-0.0757	$v_{K\Lambda, K\Lambda}$	0.1371
$\Lambda_{v, \pi N} / \text{GeV}$	0.6000	$\Lambda_{v, \eta N} / \text{GeV}$	0.9036
$\Lambda_{v, K\Lambda} / \text{GeV}$	0.6060		

Bare masses (2):

$$m_{N^*1} \sim 1535, \quad m_{N^*2} \sim 1650$$

Bare-Channel couplings (2+2):

$$g_{01}, \quad \Lambda_1, \quad g_{02}, \quad \Lambda_2$$

Channel-Channel coupling:

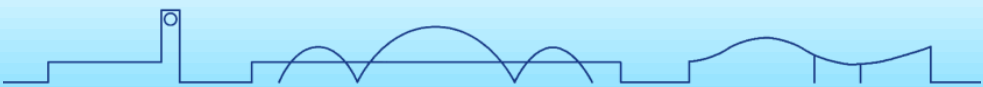
1 couplings:  $\pi N, \eta N, K\Lambda, K\Sigma$

4 (or 1) cuts for each channel

Total: 8 or 11 parameters !



# Thanks for attention!



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