

# Triangle Singularities in a Hilbert's House

NSTAR24, York

Ajay S. Sakthivasan <sup>1</sup>

M. Döring <sup>2,3</sup> M. Mai <sup>1,2</sup> A. Rusetsky <sup>1,4</sup>

<sup>1</sup>HISKP & BCTP, University of Bonn

<sup>2</sup>INS & Department of Physics, GWU

<sup>3</sup>Thomas Jefferson National Accelerator Facility

<sup>4</sup>Tbilisi State University

June 17, 2024

# Presentation Outline

- 1 Introduction
- 2 Analytic Aspects
- 3 Infinite Volume Unitarity Formalism
- 4 Conclusions

# Introduction

- Hadronic resonances correspond to poles of the  $S$ -matrix.

# Introduction

- Hadronic resonances correspond to poles of the  $S$ -matrix.
- Experimentally — peaks in invariant mass distributions.

# Introduction

- Hadronic resonances correspond to poles of the  $S$ -matrix.
- Experimentally — peaks in invariant mass distributions.
- All observed peaks correspond to hadronic states? **No**.

- **Landau Singularities** may not correspond to poles of the  $S$ -matrix, but can *mimic* a resonance.

- **Landau Singularities** may not correspond to poles of the  $S$ -matrix, but can *mimic* a resonance.
- Originally proposed by Landau. Recently used to explain a lot of observed peaks (one-loop level, rescattering effects not taken into account).

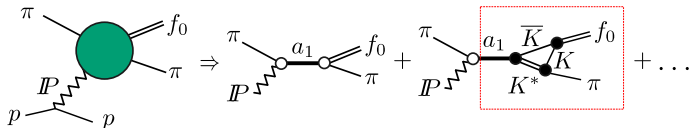
- **Landau Singularities** may not correspond to poles of the  $S$ -matrix, but can *mimic* a resonance.
- Originally proposed by Landau. Recently used to explain a lot of observed peaks (one-loop level, rescattering effects not taken into account).
- We try to achieve the following:



- **Landau Singularities** may not correspond to poles of the  $S$ -matrix, but can *mimic* a resonance.
- Originally proposed by Landau. Recently used to explain a lot of observed peaks (one-loop level, rescattering effects not taken into account).
- We try to achieve the following:
  - ① To systematically study the effect of **final-state rescattering**.

- **Landau Singularities** may not correspond to poles of the  $S$ -matrix, but can *mimic* a resonance.
- Originally proposed by Landau. Recently used to explain a lot of observed peaks (one-loop level, rescattering effects not taken into account).
- We try to achieve the following:
  - ① To systematically study the effect of **final-state rescattering**.
  - ② To **consistently describe the unstable particles**.

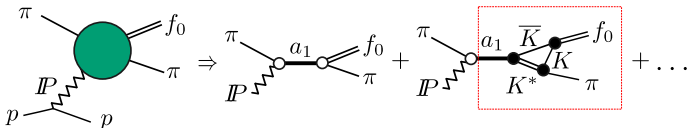
- **Landau Singularities** may not correspond to poles of the  $S$ -matrix, but can *mimic* a resonance.
- Originally proposed by Landau. Recently used to explain a lot of observed peaks (one-loop level, rescattering effects not taken into account).
- We try to achieve the following:
  - ① To systematically study the effect of **final-state rescattering**.
  - ② To **consistently describe the unstable particles**.
  - ③ Apply this to an  $a_1(1420)$ -like system, and eventually the  $a_1(1420)$  system<sup>1</sup>.



<sup>1</sup>Figure from: 10.1103/physrevlett.127.082501

- **Landau Singularities** may not correspond to poles of the  $S$ -matrix, but can *mimic* a resonance.
- Originally proposed by Landau. Recently used to explain a lot of observed peaks (one-loop level, rescattering effects not taken into account).
- We try to achieve the following:
  - ① To systematically study the effect of **final-state rescattering**.
  - ② To **consistently describe the unstable particles**.
  - ③ **Apply this to an  $a_1(1420)$ -like system**, and eventually the  $a_1(1420)$  system<sup>1</sup>.

} **IVU**



<sup>1</sup>Figure from: 10.1103/physrevlett.127.082501

- $a_1(1420)$ : resonance like structure in COMPASS experiment.  
(10.1103/PhysRevLett.115.082001)

- $a_1(1420)$ : resonance like structure in COMPASS experiment.  
(10.1103/PhysRevLett.115.082001)
- Interpretation as a kinematical singularity — Mikhasenko et al.  
(10.1103/PhysRevD.91.094015)

- $a_1(1420)$ : resonance like structure in COMPASS experiment. (10.1103/PhysRevLett.115.082001)
- Interpretation as a kinematical singularity — Mikhasenko et al. (10.1103/PhysRevD.91.094015)
- Triangle singularities in systems with non-zero spin particles — Bayar et al. (10.1103/PhysRevD.94.074039)

- $a_1(1420)$ : resonance like structure in COMPASS experiment. (10.1103/PhysRevLett.115.082001)
- Interpretation as a kinematical singularity — Mikhasenko et al. (10.1103/PhysRevD.91.094015)
- Triangle singularities in systems with non-zero spin particles — Bayar et al. (10.1103/PhysRevD.94.074039)
- Review — Guo et al. (10.1016/j.pnpnp.2020.103757)



- $a_1(1420)$ : resonance like structure in COMPASS experiment.  
(10.1103/PhysRevLett.115.082001)
- Interpretation as a kinematical singularity — Mikhasenko et al.  
(10.1103/PhysRevD.91.094015)
- Triangle singularities in systems with non-zero spin particles — Bayar et al. (10.1103/PhysRevD.94.074039)
- Review — Guo et al. (10.1016/j.pnpnp.2020.103757)
- Relevance to lattice QCD — Korpa et al.  
(10.1103/PhysRevD.107.L031505), Isken et al.  
(10.1103/PhysRevD.109.034032)

# The Landau Equations

- Can one say something about the non-analyticities of an integral, without actually evaluating it? E.g.,

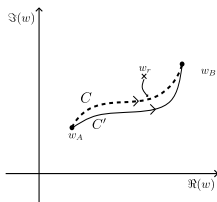
$$f(z) = \int_C dw g(z, w).$$

# The Landau Equations

- Can one say something about the non-analyticities of an integral, without actually evaluating it? E.g.,

$$f(z) = \int_C dw g(z, w).$$

- Singularities in  $g$  can be avoided — *contour deformation* (to a contour  $C'$ ).

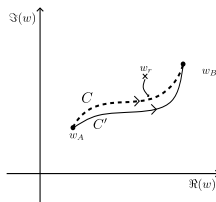


# The Landau Equations

- Can one say something about the non-analyticities of an integral, without actually evaluating it? E.g.,

$$f(z) = \int_C dw g(z, w).$$

- Singularities in  $g$  can be avoided — *contour deformation* (to a contour  $C'$ ).



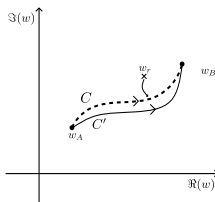
- When is this not possible?

# The Landau Equations

- Can one say something about the non-analyticities of an integral, without actually evaluating it? E.g.,

$$f(z) = \int_C dw g(z, w).$$

- Singularities in  $g$  can be avoided — *contour deformation* (to a contour  $C'$ ).



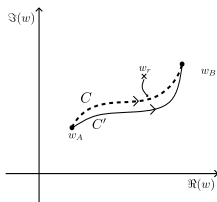
- When is this not possible?
  - 1 Endpoint singularities.

# The Landau Equations

- Can one say something about the non-analyticities of an integral, without actually evaluating it? E.g.,

$$f(z) = \int_C dw g(z, w).$$

- Singularities in  $g$  can be avoided — *contour deformation* (to a contour  $C'$ ).



- When is this not possible?
  - Endpoint singularities.
  - Pinch singularities.

- **The Landau Equations:** singularities corresponding to Feynman integrals (*Landau singularities*).

- **The Landau Equations:** singularities corresponding to Feynman integrals (*Landau singularities*).
- For the Feynman integral with  $N$  internal propagators and  $l$  loop momenta,

$$I_1 = \int \prod_{i=1}^l d^4 k_i \frac{1}{\prod_{r=1}^N (q_r(k_i)^2 - m_r^2)}.$$



- **The Landau Equations:** singularities corresponding to Feynman integrals (*Landau singularities*).
- For the Feynman integral with  $N$  internal propagators and  $l$  loop momenta,

$$I_1 = \int \prod_{i=1}^l d^4 k_i \frac{1}{\prod_{r=1}^N (q_r(k_i)^2 - m_r^2)}.$$

- For some real parameters  $\alpha_i$ , the Landau equations read,

$$\alpha_i (q_i^2 - m_i^2) = 0 \implies \alpha_i = 0 \text{ or } q_i^2 = m_i^2,$$
$$\sum_{i \in \text{loop}} \alpha_i q_i^\mu(k_j) = 0.$$

- $\alpha_i \neq 0 \implies$  all propagators go on-shell for the **leading singularity**.  
One also needs  $\alpha_i > 0$ .

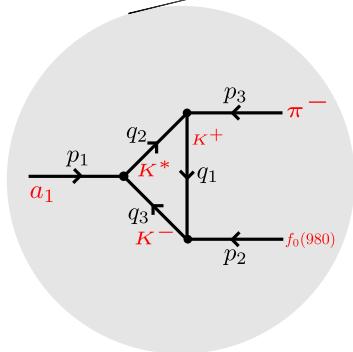
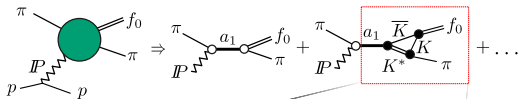
- $\alpha_i \neq 0 \implies$  all propagators go on-shell for the **leading singularity**.  
One also needs  $\alpha_i > 0$ .
- $\alpha_i = 0$  case is called the **subleading singularity** — equivalent to contracting the propagator.

- $\alpha_i \neq 0 \implies$  all propagators go on-shell for the **leading singularity**. One also needs  $\alpha_i > 0$ .
- $\alpha_i = 0$  case is called the **subleading singularity** — equivalent to contracting the propagator.
- Interesting analogy: Feynman diagrams  $\rightarrow$  electrical circuits; propagators  $\rightarrow$  wires, with current  $q_i$  and resistance  $\alpha_i$ . Then the Landau equations are identical to Kirchhoff's law!<sup>2</sup>

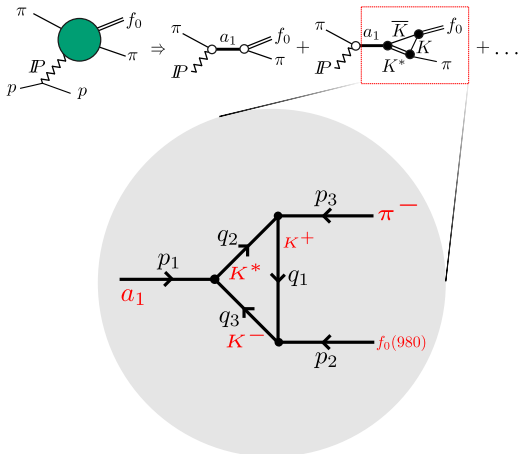
---

<sup>2</sup> $\Delta V = \sum IR$ , for a triangle circuit with  $\Delta V = 0$ .

- The leading Landau singularity associated with the graph:



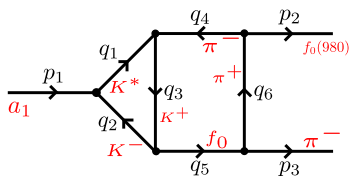
- The leading Landau singularity associated with the graph:



- Landau singularity for  $p_3 = m_\pi$ :  $p_2/\text{GeV} \equiv \sqrt{\sigma}/\text{GeV} \in [0.986, 1.024]$  and  $p_1/\text{GeV} \equiv \sqrt{s}/\text{GeV} \in [1.385, 1.436]$ .

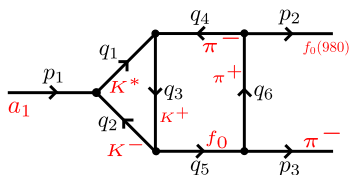
- What happens to singularities after rescattering? **Ladder diagrams.**

- What happens to singularities after rescattering? **Ladder diagrams.**
- 1 + 1-loop diagram:





- What happens to singularities after rescattering? **Ladder diagrams.**
- 1 + 1-loop diagram:

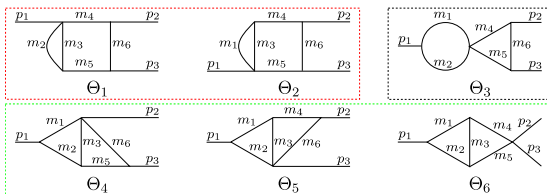


- Check for leading and subleading singularities. . .

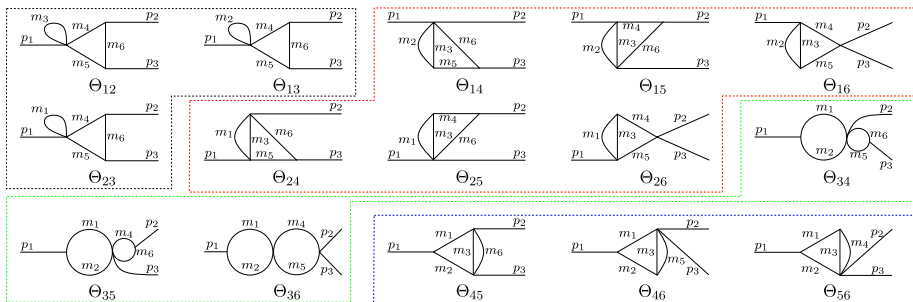
- Leading singularity is not present.

- Leading singularity **is not present**.
- For subleading singularities go through all contractions and check for singular graphs. . .

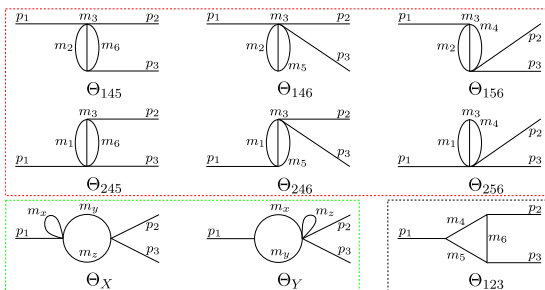
- Leading singularity **is not present**.
- For subleading singularities go through all contractions and check for singular graphs. . .



- Leading singularity **is not present**.
- For subleading singularities go through all contractions and check for singular graphs. . .

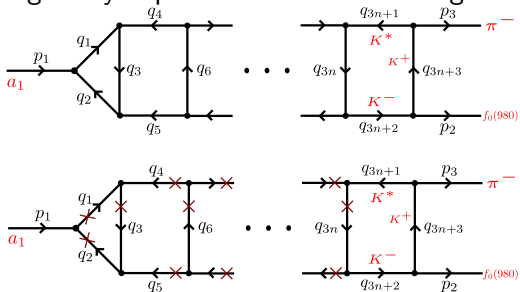


- Leading singularity **is not present**.
- For subleading singularities go through all contractions and check for singular graphs. . .



- Accommodating **infinite rescatterings**.

- Accommodating **infinite rescatterings**.
- The triangle singularity is present and no other singularities.





- Landau equations give the conditions for the singularity to be present.

- Landau equations give the conditions for the singularity to be present.
- One still needs to evaluate the corresponding Feynman integrals.

- Study the system using **Infinite Volume 3-Body Unitary Formalism**, due to Mai et al.<sup>3</sup>

---

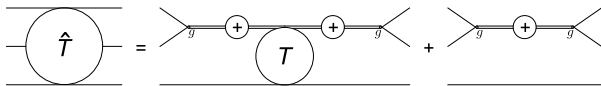
<sup>3</sup>[10.1140/epja/i2017-12368-4](https://arxiv.org/abs/10.1140/epja/i2017-12368-4)

- Study the system using **Infinite Volume 3-Body Unitary Formalism**, due to Mai et al.<sup>3</sup>
- **In brief**: the system is split into a two-body subsystem — the **isobar**, and the **spectator**, which is on-shell.

---

<sup>3</sup>[10.1140/epja/i2017-12368-4](https://doi.org/10.1140/epja/i2017-12368-4)

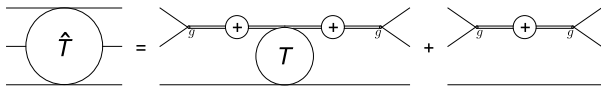
- Study the system using **Infinite Volume 3-Body Unitary Formalism**, due to Mai et al.<sup>3</sup>
- **In brief**: the system is split into a two-body subsystem — the **isobar**, and the **spectator**, which is on-shell.
- The full three-body amplitude is split into a **connected** and a **disconnected** amplitude.<sup>4</sup>



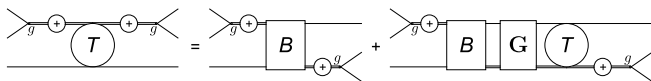
<sup>3</sup>[10.1140/epja/i2017-12368-4](https://doi.org/10.1140/epja/i2017-12368-4)

<sup>4</sup>Figure from: [10.1140/epja/i2017-12368-4](https://doi.org/10.1140/epja/i2017-12368-4)

- Study the system using **Infinite Volume 3-Body Unitary Formalism**, due to Mai et al.<sup>3</sup>
- **In brief**: the system is split into a two-body subsystem — the **isobar**, and the **spectator**, which is on-shell.
- The full three-body amplitude is split into a **connected** and a **disconnected** amplitude.<sup>4</sup>



- One considers the **Bethe-Salpeter ansatz**,  $T = B + BGT$ , to determine  $T$ .<sup>4</sup>



<sup>3</sup>10.1140/epja/i2017-12368-4

<sup>4</sup>Figure from: 10.1140/epja/i2017-12368-4

- The equation to be solved,

$$T(p, q) = B(p, q) + \int \frac{d^3l}{(2\pi)^3} \frac{1}{2E_l} B(p, l) \tau(\sigma(l)) T(l, q).$$

- The equation to be solved,

$$T(p, q) = B(p, q) + \int \frac{d^3\mathbf{l}}{(2\pi)^3} \frac{1}{2E_l} B(p, l) \tau(\sigma(l)) T(l, q).$$

- We have the explicit forms of  $B$  and  $\tau$ .

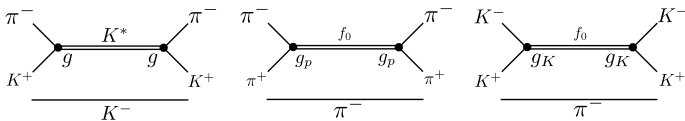


- The equation to be solved,

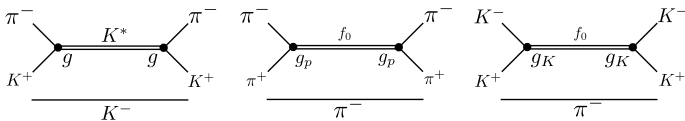
$$T(p, q) = B(p, q) + \int \frac{d^3\mathbf{l}}{(2\pi)^3} \frac{1}{2E_l} B(p, l) \tau(\sigma(l)) T(l, q).$$

- We have the explicit forms of  $B$  and  $\tau$ .
- In this work, we consider scalar propagators, in relative  $s$ -wave.

- Need to carry out a coupled channel analysis.

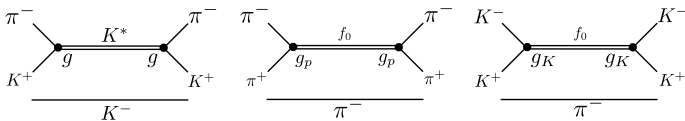


- Need to carry out a coupled channel analysis.



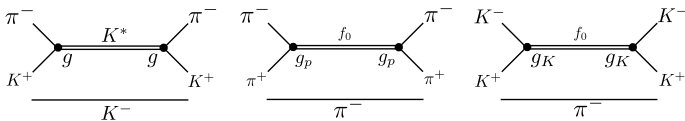
- There are two unknowns corresponding to each isobar — coupling and bare mass.

- Need to carry out a coupled channel analysis.



- There are two unknowns corresponding to each isobar — coupling and bare mass.
- These are fixed through PDG values of physical mass and width.

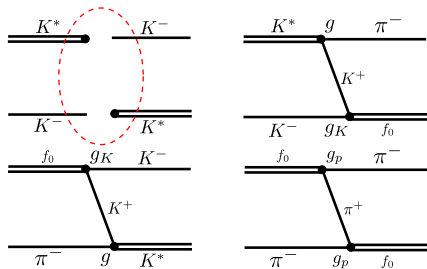
- Need to carry out a coupled channel analysis.



- There are two unknowns corresponding to each isobar — coupling and bare mass.
- These are fixed through PDG values of physical mass and width.
- For  $f_0$  we also make use of  $g_K/g_p$  ratio from BaBar collaboration.

- $B$  is parametrised through the spectator masses.

- $B$  is parametrised through the spectator masses.
- Diagrammatically,

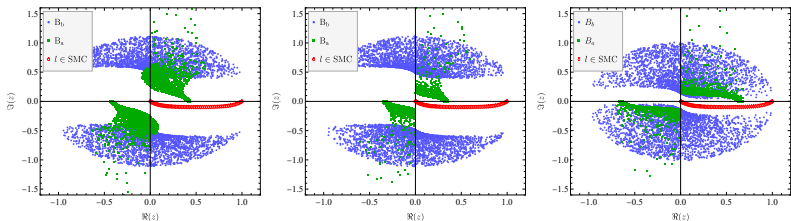


- The integral is solved numerically, using Gaussian quadrature.

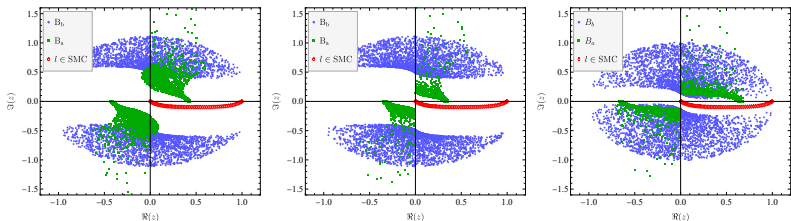


- The integral is solved numerically, using Gaussian quadrature.
- But, singularities on the integration contour may prevent numerical implementation.

- The integral is solved numerically, using Gaussian quadrature.
- But, singularities on the integration contour may prevent numerical implementation.
- **Contour deformation**: Evaluate the integral along a complex contour, that is far from the singularities  $\rightarrow$  obtain  $T(\sqrt{s}, q \in \mathbb{C})$ .



- The integral is solved numerically, using Gaussian quadrature.
- But, singularities on the integration contour may prevent numerical implementation.
- **Contour deformation**: Evaluate the integral along a complex contour, that is far from the singularities  $\rightarrow$  obtain  $T(\sqrt{s}, q \in \mathbb{C})$ .



- But, we need the amplitudes for  $q \in \mathbb{R} \rightarrow$  can be done in different ways!

- The solution of the Bethe-Salpeter equation can be given in form of the **Born series**.

The diagram shows the Born series for the transition operator  $T$ . On the left, a circle labeled  $T$  is connected to two external lines. This is equal to the sum of three terms: 1) a single vertex  $B$  connected to two external lines; 2) a vertex  $B$  connected to two external lines, with a loop consisting of two  $B$  vertices and two internal lines; 3) a vertex  $B$  connected to two external lines, with a loop consisting of two  $B$  vertices, two internal lines, and a small circle with a minus sign. The series continues with an ellipsis.

$$T = B + B \circlearrowleft B + B \circlearrowleft B \circlearrowleft B + \dots$$

- The solution of the Bethe-Salpeter equation can be given in form of the **Born series**.

$$\begin{array}{c} \text{---} \\ \text{---} \end{array} \circlearrowleft T \begin{array}{c} \text{---} \\ \text{---} \end{array} = \begin{array}{c} \text{---} \\ \text{---} \end{array} \text{---} B \begin{array}{c} \text{---} \\ \text{---} \end{array} + \begin{array}{c} \text{---} \\ \text{---} \end{array} \text{---} B \begin{array}{c} \text{---} \\ \text{---} \end{array} \circlearrowleft T \begin{array}{c} \text{---} \\ \text{---} \end{array} \text{---} B \begin{array}{c} \text{---} \\ \text{---} \end{array} + \dots$$

- This defines the **loop expansion** for the invariant mass distributions.

- The solution of the Bethe-Salpeter equation can be given in form of the **Born series**.

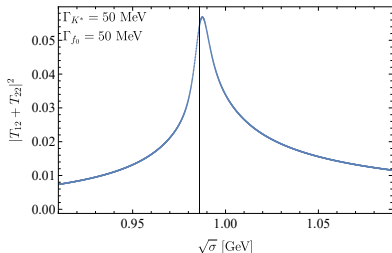
$$\begin{array}{c} \text{---} \\ \text{---} \end{array} \circlearrowleft T \begin{array}{c} \text{---} \\ \text{---} \end{array} = \begin{array}{c} \text{---} \\ \text{---} \end{array} \diagdown B \begin{array}{c} \text{---} \\ \text{---} \end{array} + \begin{array}{c} \text{---} \\ \text{---} \end{array} \diagdown B \begin{array}{c} \text{---} \\ \text{---} \end{array} \circlearrowleft T \begin{array}{c} \text{---} \\ \text{---} \end{array} \diagdown B \begin{array}{c} \text{---} \\ \text{---} \end{array} + \dots$$

- This defines the **loop expansion** for the invariant mass distributions.
- **Complex Contour Analytic Continuation:** We use a **continued fraction** to obtain the isobar-spectator amplitude for  $q \in \mathbb{R}$ .

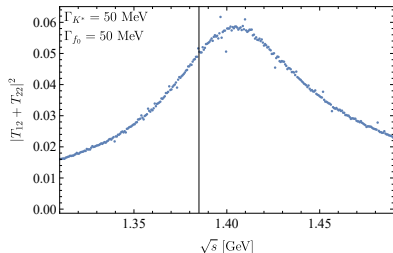
- The solution of the Bethe-Salpeter equation can be given in form of the **Born series**.

$$\text{---} \circlearrowleft T \text{---} = \text{---} B \text{---} + \text{---} B \text{---} B \text{---} + \dots$$

- This defines the **loop expansion** for the invariant mass distributions.
- Complex Contour Analytic Continuation:** We use a **continued fraction** to obtain the isobar-spectator amplitude for  $q \in \mathbb{R}$ .
- At 1-loop level (the triangle),



Amplitude squared vs.  $\sqrt{\sigma}$ , for  $\sqrt{s} = 1.42$  GeV.

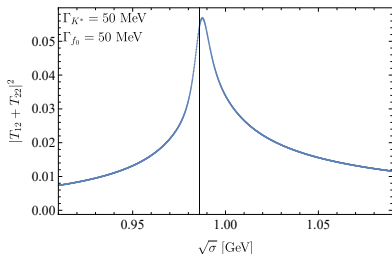


Amplitude squared vs.  $\sqrt{s}$ , for  $\sqrt{\sigma} = 0.99$  GeV.

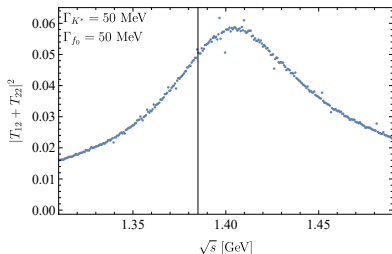
- The solution of the Bethe-Salpeter equation can be given in form of the **Born series**.

$$\text{---} \circlearrowleft T \text{---} = \text{---} B \text{---} + \text{---} B \text{---} B \text{---} + \dots$$

- This defines the **loop expansion** for the invariant mass distributions.
- Complex Contour Analytic Continuation:** We use a **continued fraction** to obtain the isobar-spectator amplitude for  $q \in \mathbb{R}$ .
- At 1-loop level (the triangle),



Amplitude squared vs.  $\sqrt{\sigma}$ , for  $\sqrt{s} = 1.42$  GeV.

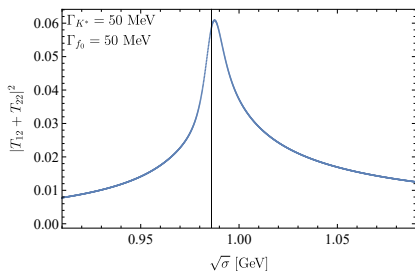


Amplitude squared vs.  $\sqrt{s}$ , for  $\sqrt{\sigma} = 0.99$  GeV.

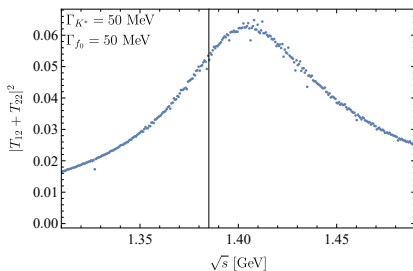
- We see the (leading) Landau singularity in both the plots.



- At  $1 + \infty$ -loop level,

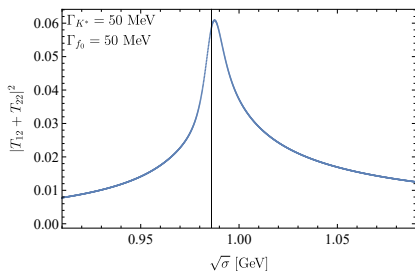


Amplitude squared vs.  $\sqrt{\sigma}$ , for  $\sqrt{s} = 1.42$  GeV.



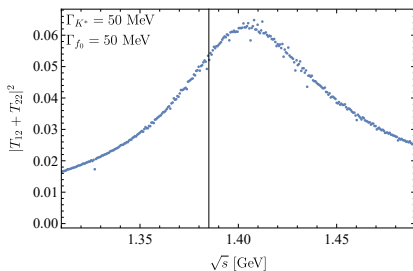
Amplitude squared vs.  $\sqrt{s}$ , for  $\sqrt{\sigma} = 0.99$  GeV.

- At  $1 + \infty$ -loop level,



Amplitude squared vs.  $\sqrt{\sigma}$ , for  $\sqrt{s} = 1.42$  GeV.

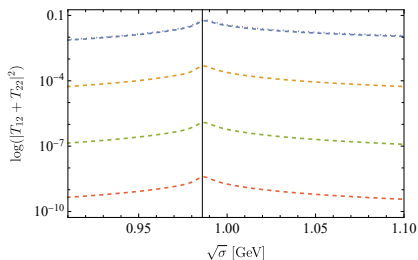
- The magnitudes are quite similar!



Amplitude squared vs.  $\sqrt{s}$ , for  $\sqrt{\sigma} = 0.99$  GeV.

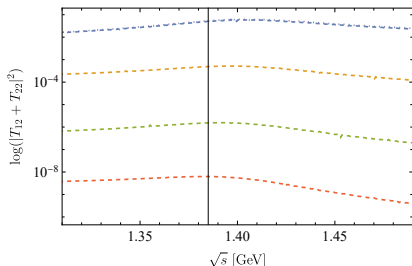
- The order of magnitude of the amplitudes at the triangle level is much larger than the subsequent levels.

- The order of magnitude of the amplitudes at the triangle level is much larger than the subsequent levels.
- Triangle diagram contributes the most to the amplitude:



--- Triangle    --- Triangle+1box    --- Triangle+2box  
 --- Triangle+3box    ..... IVU

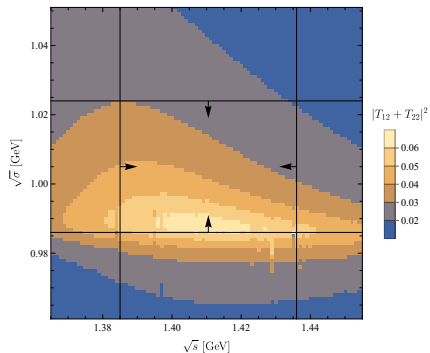
Amplitude squared vs.  $\sqrt{\sigma}$ , for  
 $\sqrt{s} = 1.42$  GeV.



--- Triangle    --- Triangle+1box    --- Triangle+2box  
 --- Triangle+3box    ..... IVU

Amplitude squared vs.  $\sqrt{s}$ , for  
 $\sqrt{\sigma} = 0.99$  GeV.

- Characteristic of triangle singularities: **very sensitive to invariant masses!**



# Conclusions

- Triangle singularity at 1-loop level and at higher loops as **subleading singularity**.

# Conclusions

- Triangle singularity at 1-loop level and at higher loops as **subleading singularity**.
- The isobar-spectator model numerically implemented and the triangle singularity observed at  $1 + \infty$ -loop level.

# Conclusions

- Triangle singularity at 1-loop level and at higher loops as **subleading singularity**.
- The isobar-spectator model numerically implemented and the triangle singularity observed at  $1 + \infty$ -loop level.
- The singularity is of **logarithmic** nature at all levels.



# Conclusions

- Triangle singularity at 1-loop level and at higher loops as **subleading singularity**.
- The isobar-spectator model numerically implemented and the triangle singularity observed at  $1 + \infty$ -loop level.
- The singularity is of **logarithmic** nature at all levels.
- The higher order diagrams did not have significant contributions  $\implies$  explains why no need to take final state interactions into account.

# Conclusions

- Triangle singularity at 1-loop level and at higher loops as **subleading singularity**.
- The isobar-spectator model numerically implemented and the triangle singularity observed at  $1 + \infty$ -loop level.
- The singularity is of **logarithmic** nature at all levels.
- The higher order diagrams did not have significant contributions  $\implies$  explains why no need to take final state interactions into account.
- Outlook:
  - 1 Realistic model — spin with pseudovector source.

# Conclusions

- Triangle singularity at 1-loop level and at higher loops as **subleading singularity**.
- The isobar-spectator model numerically implemented and the triangle singularity observed at  $1 + \infty$ -loop level.
- The singularity is of **logarithmic** nature at all levels.
- The higher order diagrams did not have significant contributions  $\implies$  explains why no need to take final state interactions into account.
- Outlook:
  - 1 Realistic model — spin with pseudovector source.
  - 2  $a_1(1420)$  on the lattice? Implement finite volume unitarity.<sup>5</sup>

---

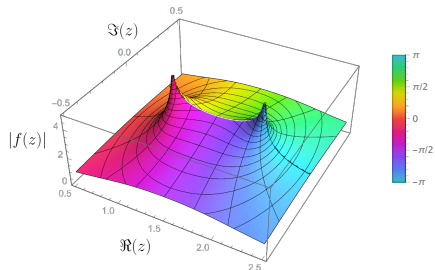
<sup>5</sup>[10.1140/epja/i2017-12440-1](https://arxiv.org/abs/10.1140/epja/i2017-12440-1)

Thanks for listening!

# Landau singularities

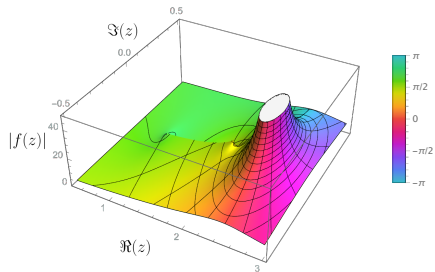
**Endpoint singularities:** singularities in  $g$ , hitting the contour  $C$ . E.g.,

$$f(z) = \int_1^2 dw \frac{1}{w-z}. \quad (1)$$

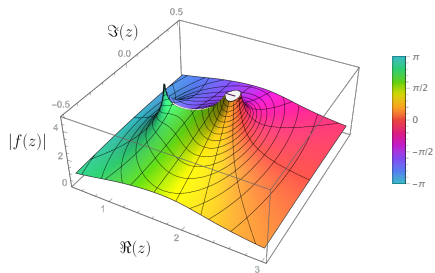


**Pinch singularities:** Contour  $C$ , gets trapped between two singularities in  $g$ . E.g.,

$$f(z) = \int_1^2 dw \frac{1}{(w-z)(w-5/2)}. \quad (2)$$



Pinch singularity avoided when the singularities approach from the same side of the contour.



# The Other Method

- **Complex Contour on Riemann Surfaces (Cahill & Sloan)<sup>6</sup>:**  
Instead of a continued fraction, use the knowledge of the **location of the branch cuts** to analytically continue the amplitude for  $q \in \mathbb{R}$ .<sup>7</sup>

---

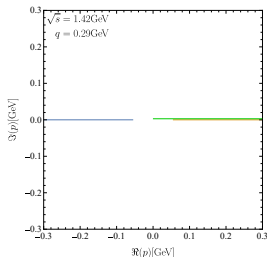
<sup>6</sup>10.1016/0375-9474(71)90156-4

<sup>7</sup>For a comprehensive review: *The Quantum Mechanical Three-Body Problem*, Schmid & Ziegelmann.

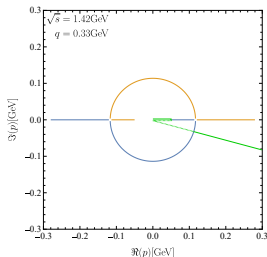


# The Other Method

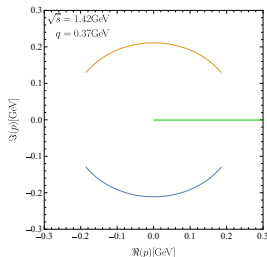
- Complex Contour on Riemann Surfaces (Cahill & Sloan)<sup>6</sup>:**  
 Instead of a continued fraction, use the knowledge of the **location of the branch cuts** to analytically continue the amplitude for  $q \in \mathbb{R}$ .<sup>7</sup>



Branch cuts of  $B$  for small  $q$ .



Branch cuts of  $B$  for intermediate  $q$ .

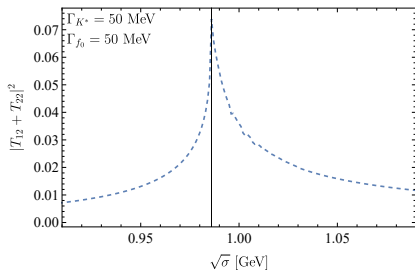


Branch cuts of  $B$  for large  $q$ .

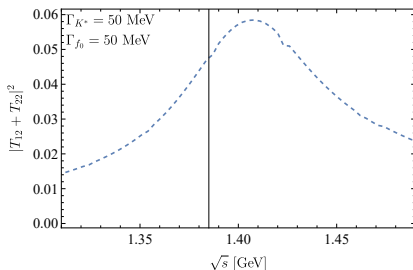
<sup>6</sup>10.1016/0375-9474(71)90156-4

<sup>7</sup>For a comprehensive review: *The Quantum Mechanical Three-Body Problem*, Schmid & Ziegelmann.

- At  $1 + \infty$ -loop level,

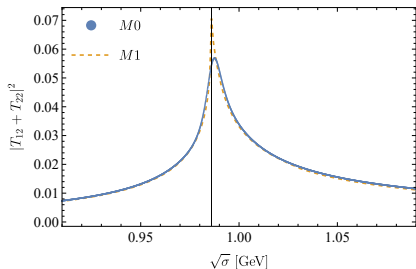


Amplitude squared vs.  $\sqrt{\sigma}$ , for  $\sqrt{s} = 1.42$  GeV.

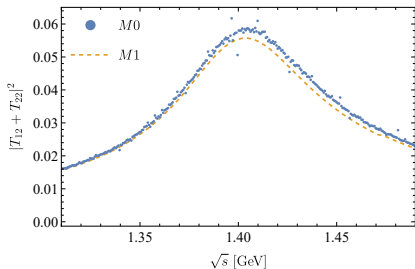


Amplitude squared vs.  $\sqrt{s}$ , for  $\sqrt{\sigma} = 0.99$  GeV.

- Matches with the other method.



Amplitude squared vs.  $\sqrt{\sigma}$ , for  $\sqrt{s} = 1.42$  GeV.



Amplitude squared vs.  $\sqrt{s}$ , for  $\sqrt{\sigma} = 0.99$  GeV.