




17.–21. June 2024, York

NSTAR2024



Exploring the 3D Structure of Baryon Resonances with Transition Generalized Parton Distributions

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GIESSEN



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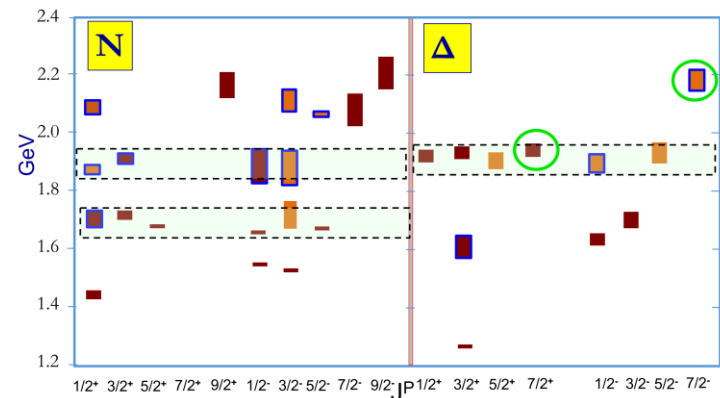
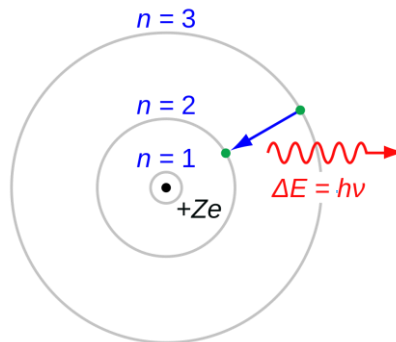
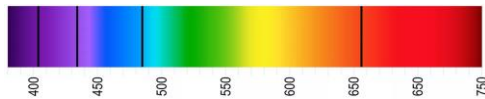
University of Connecticut

June 18th 2024

Introduction

- Baryons and their resonances are emergent phenomena of QCD
- Elementary QCD dynamics expressed in spectrum and structure of the hadronic states
 - ➔ Both aspects are essential for a complete understanding of strong interactions

- Connection between the internal motion and the excitation spectrum:



➔ A rich spectrum of baryon resonances emerges from QCD

- Baryon resonances play an important role:
 - Behavior of matter at high densities and temperatures (early universe, stellar structure)
 - Existence of hypothetical strange matter (neutron stars)
 - Theory of nuclear forces (e.g. the Δ isobar)
 - Description of neutrino interactions with nuclei at a few GeV

➔ **Understanding the internal structure of the baryon resonances is of fundamental interest and practical importance**

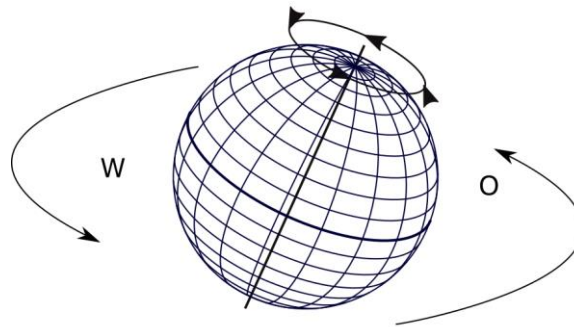
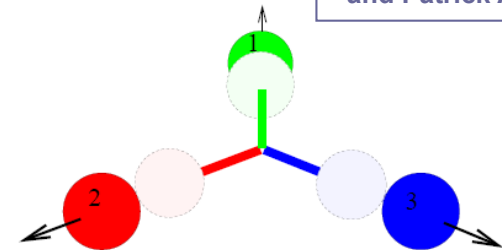
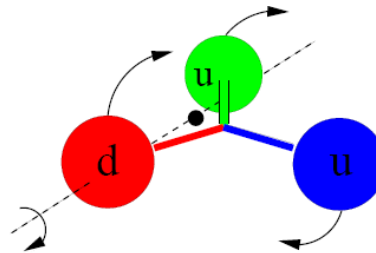
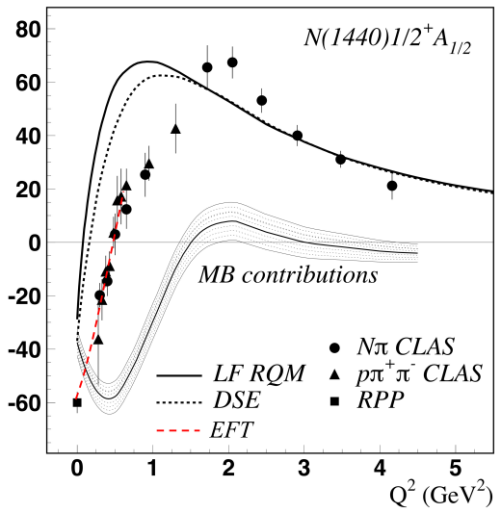


Introduction

Electroproduction experiments: Information on the structure of baryon resonances from **electromagnetic transition form factors**

→ Spatial (2D) distribution of charge and current in dynamical systems

talks by Victor Mokeev and Patrick Achenbach



• Not all phenomena can be explained in a 2D picture

$$L \sim r \times p$$

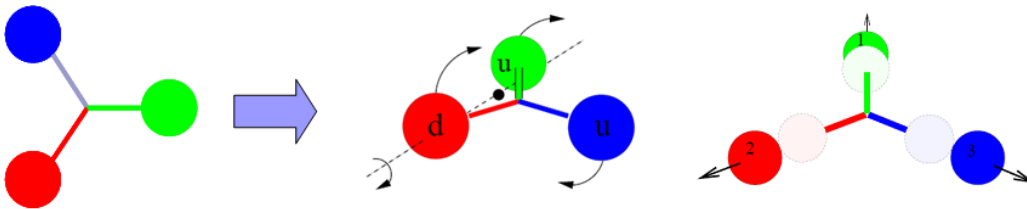


source: www.prosieben.de/serien/galileo

source: www.weltkugel-globus.de/die-erde

Generalized Parton Distributions (GPDs)

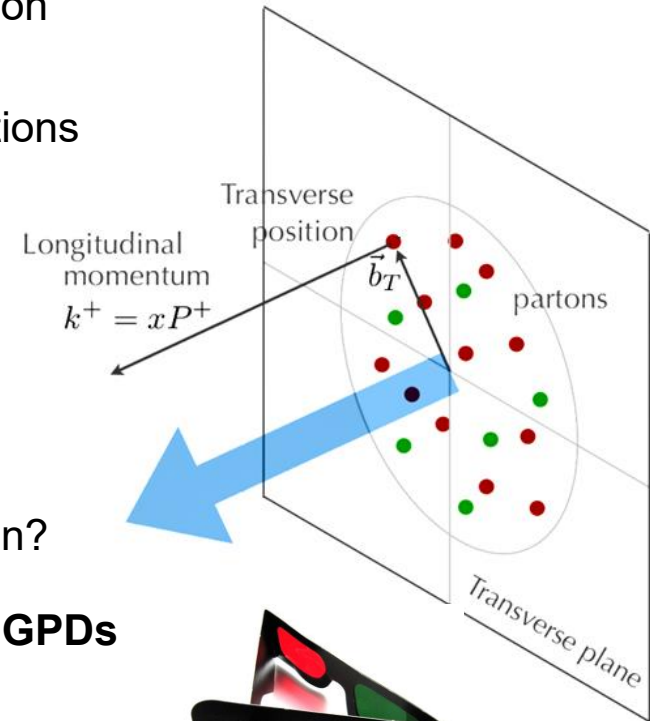
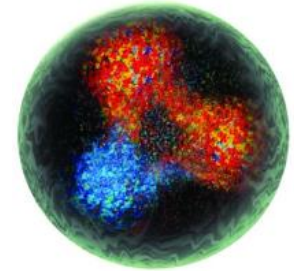
- Significant advances in the study of the ground-state nucleon structure over the last two decades
- **Generalized Parton Distributions (GPDs)** connect the transverse position and the longitudinal momentum space
 - ➔ Unification of elastic nucleon form factors and quark/gluon particle densities (PDFs)
 - ➔ 3D tomographic images of the quark and gluon distributions + Access to the mechanical properties of the nucleon



How does the excitation affect the 3D structure of the nucleon?

3D picture of the excitation process: Encoded in **transition GPDs**

- ➔ Characterisation of baryon resonance structure based on quark/gluon tomography



www.wikipedia.de

Transition GPDs for the $N \rightarrow \Delta$ Transition

$N \rightarrow \Delta$ transition: 16 transition GPDs

- 8 twist-2 helicity non-flip transition GPDs

unpolarized:

$$\left. \begin{aligned} \int_{-1}^1 dx G_1(x; \xi; t) &\propto G_M^*(t), \\ \int_{-1}^1 dx G_3(x; \xi; t) &\propto G_C^*(t), \\ \int_{-1}^1 dx G_2(x; \xi; t) &\propto G_E^*(t), \\ \int_{-1}^1 dx G_4(x; \xi; t) &= 0, \end{aligned} \right\} \begin{array}{l} \text{for the } N \rightarrow \Delta \text{ transition} \\ \text{Jones-Scardon EM FF} \end{array}$$

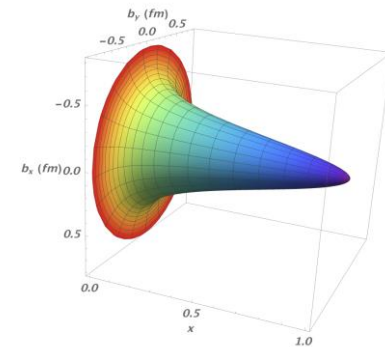
polarized:

$$\left. \begin{aligned} \int_{-1}^1 dx \tilde{G}_1(x; \xi; t) &\propto C_5^A(t) \\ \int_{-1}^1 dx \tilde{G}_3(x; \xi; t) &\propto C_3^A(t) \\ \int_{-1}^1 dx \tilde{G}_2(x; \xi; t) &\propto C_6^A(t) \\ \int_{-1}^1 dx \tilde{G}_4(x; \xi; t) &\propto C_4^A(t) \end{aligned} \right\} \begin{array}{l} \text{Adler form factors} \end{array}$$

+ 8 helicity flip transition GPDs (twist-3, transversity)

Physics Content of Transition GPDs

- 3D (x-dependent) imaging of the excitation process
 - $N \rightarrow N^*$ transition charge / magnetization densities
- Transition GPDs connect the spin and angular momentum of resonances to the motion and distribution of the partons within the excited baryon



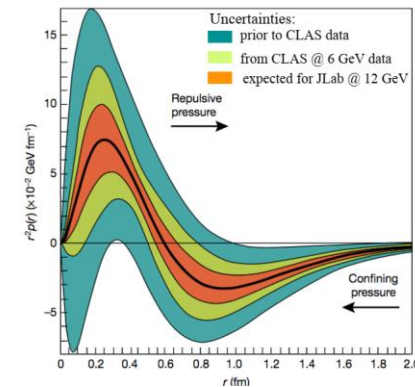
$$\int_{-1}^1 dx x h_M(x, \xi, 0) = 2J_{p \rightarrow \Delta^+}^{u-d}$$

C. C. Granados, C. Weiss, Phys. Lett. B **797**, 134847 (2019)

J.Y. Kim, H.Y. Won, J. Goity, C. Weiss, Phys. Lett. B **844**, 138083 (2023)

- Access to shear forces and pressure distributions within nucleon resonances via gravitational form factors

$$\int dx x H(x, \xi, t) = M_2(t) + \frac{4}{5} \xi^2 d_1(t)$$



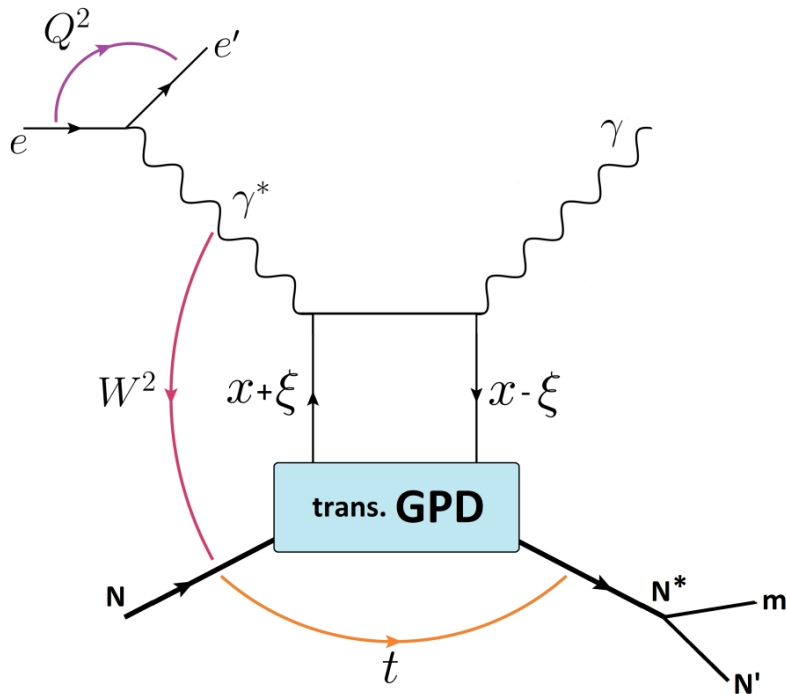
V.D. Burkert, L. Elouadrhiri,
F.X. Girod, Nature **557**, 396 (2018)

- Access to the anomalous magnetic moment and to the tensor charge of resonance

$$k_T^{u,d} = \int dx \bar{E}_T^{u,d}(x, \xi = 0, t = 0) \quad \delta_T^{u,d} = \int dx H_T^{u,d}(x, \xi = 0, t = 0)$$

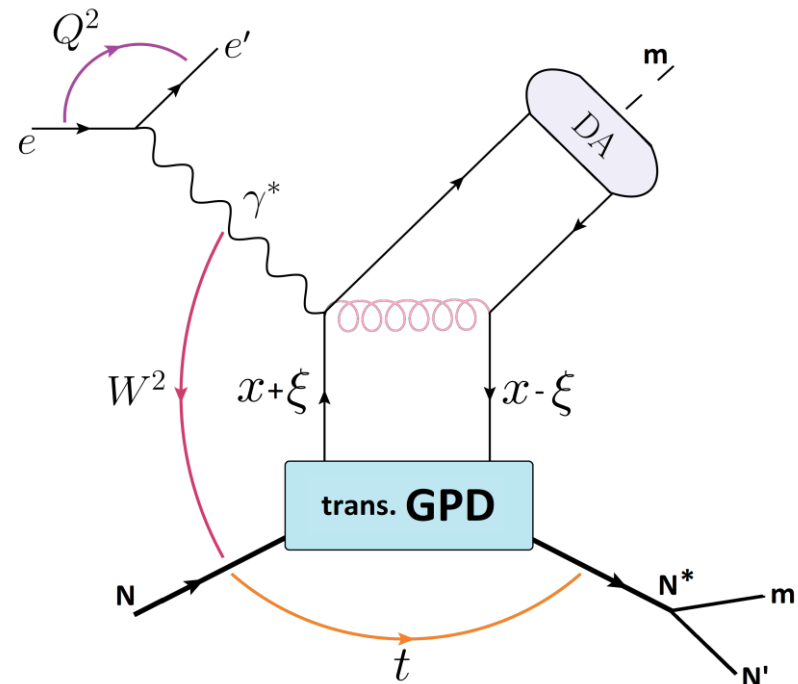
Experimental Access to Transition GPDs

$N \rightarrow N^*$ DVCS



Access to the helicity
non-flip transition GPDs

$N \rightarrow N^*$ DVMP



+ Access to the helicity
flip transition GPDs

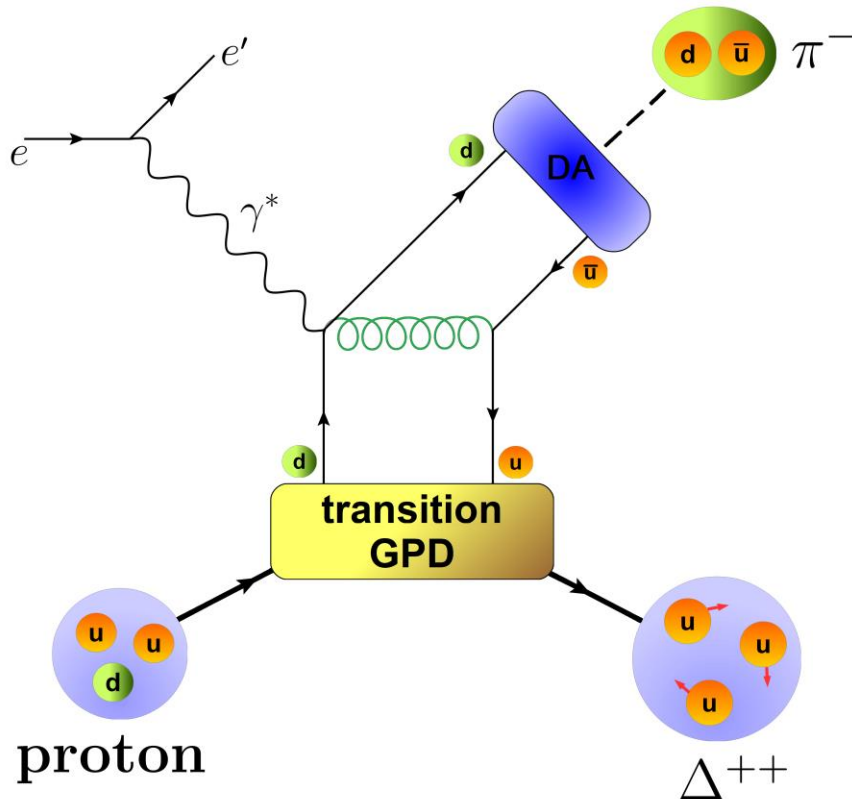
$W > 2 \text{ GeV}$

Factorisation expected for: $-t / Q^2 \ll 1$, x_B fixed and $Q^2 > M_{N^*}^2$

The $N \rightarrow N^*$ DVMP Processes

$$ep \rightarrow eN^{*0}\pi^+ \rightarrow e(p\pi^-)\pi^+ \\ \rightarrow e(n\pi^0)\pi^+$$

$$ep \rightarrow e\Delta^+\pi^0 \rightarrow e(n\pi^+)\pi^0 \\ \rightarrow e(p\pi^0)\pi^0$$



$$ep \rightarrow e\Delta^{++}\pi^- \rightarrow ep\pi^+\pi^-$$

- ➔ Provides access to the **d-quark** content of the nucleon
- ➔ Provides access to **p- Δ transition GPDs**

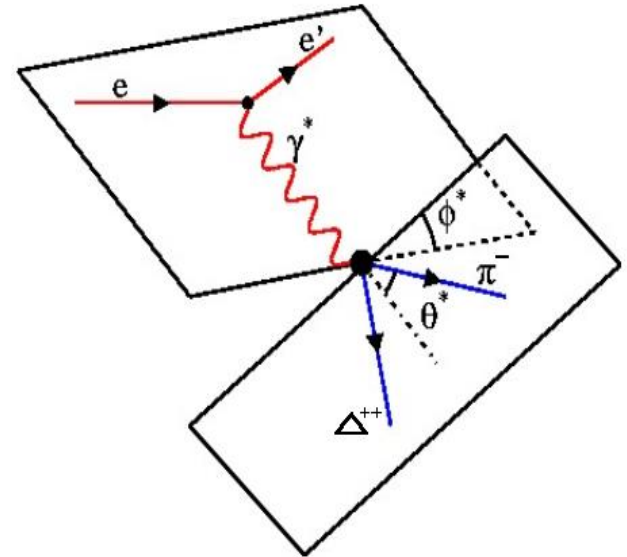
Peter Kroll, Kornelija Passek-Kumerički, Phys. Rev. D 107, 054009 (2023)
<https://doi.org/10.1103/PhysRevD.107.054009>

S. Diehl et al. (CLAS Collaboration), Phys. Rev. Lett. 131, 021901 (2023)
<https://doi.org/10.1103/PhysRevLett.131.021901>

DVMP Electroproduction Cross-Section and BSA

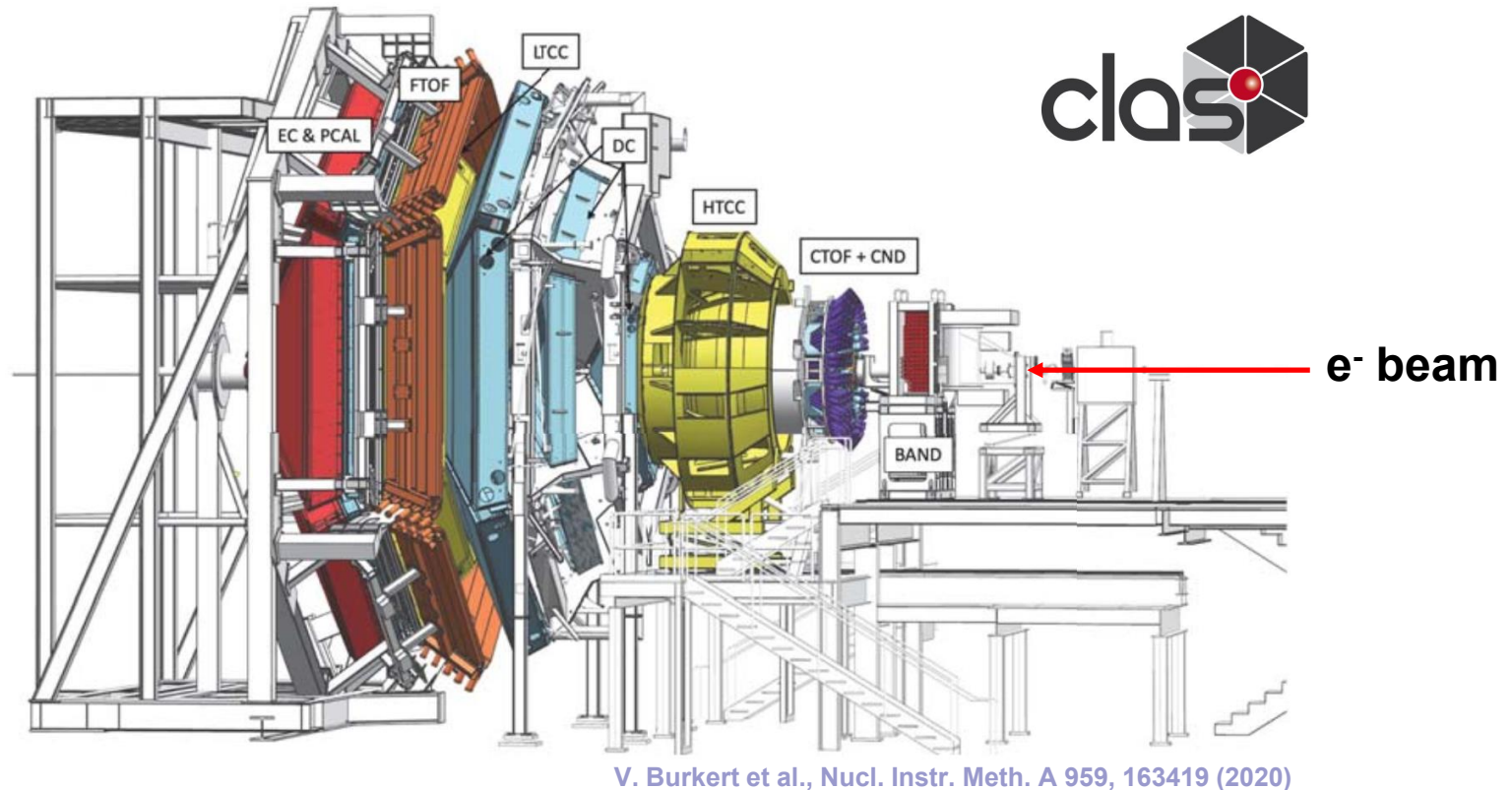
Cross section (longitudinally pol. beam and unpol. target):

$$\begin{aligned}
 2\pi \frac{d^2\sigma}{dt d\phi} = & \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \epsilon \cdot \cos(2\phi) \frac{d\sigma_{TT}}{dt} \\
 & + \sqrt{2\epsilon(1+\epsilon)} \cdot \cos(\phi) \frac{d\sigma_{LT}}{dt} \\
 & + h \cdot \sqrt{2\epsilon(1-\epsilon)} \cdot \sin(\phi) \frac{d\sigma_{LT'}}{dt}
 \end{aligned}$$



$$BSA(t, \phi, x_B, Q^2) = \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-} = \frac{\sqrt{2\epsilon(1-\epsilon)} \frac{\sigma_{LT'}}{\sigma_0} \sin \phi}{1 + \sqrt{2\epsilon(1+\epsilon)} \frac{\sigma_{LT}}{\sigma_0} \cos \phi + \epsilon \frac{\sigma_{TT}}{\sigma_0} \cos 2\phi}$$

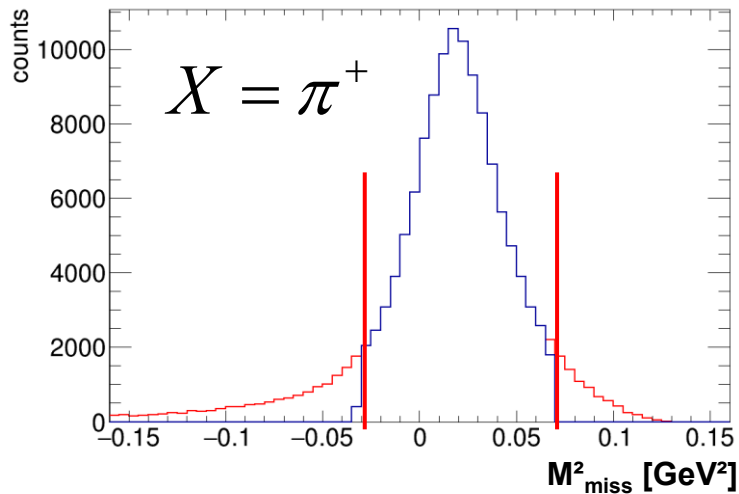
Experiment: CLAS12 at JLAB



- ➔ Data recorded with CLAS12 during fall 2018 and spring 2019 (RG-A)
 - ➔ 10.6 GeV / 10.2 GeV electron beam ~ 86 % average polarization
 - ➔ liquid H₂ target

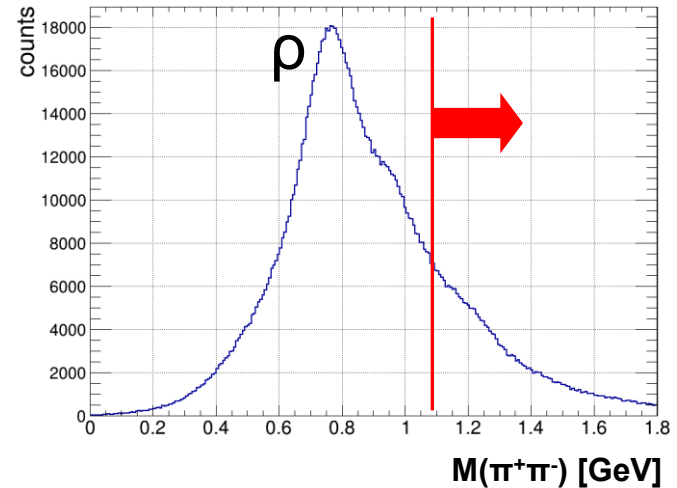
$\pi^-\Delta^{++}$ DVMP: Event Selection and Kinematic Cuts

Event selection: $ep \rightarrow ep\pi^- X$



→ 2 sigma cut around the missing π^+

BG: $ep \rightarrow ep\rho \rightarrow ep\pi^+\pi^-$



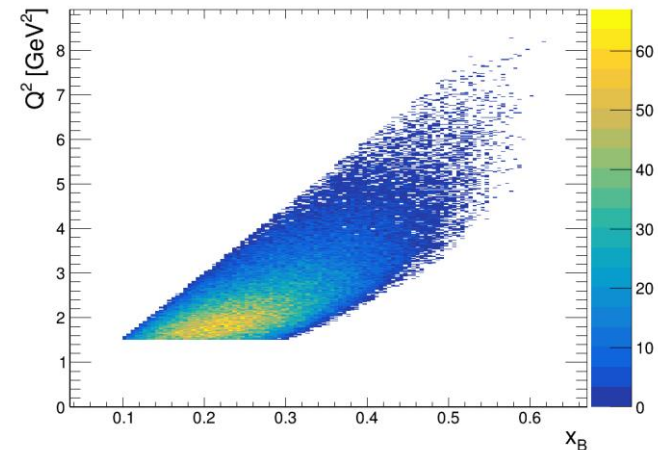
$M(\pi^+\pi^-) > 1.1 \text{ GeV}$ (ρ cont. < 0.8%)

Kinematic cuts:

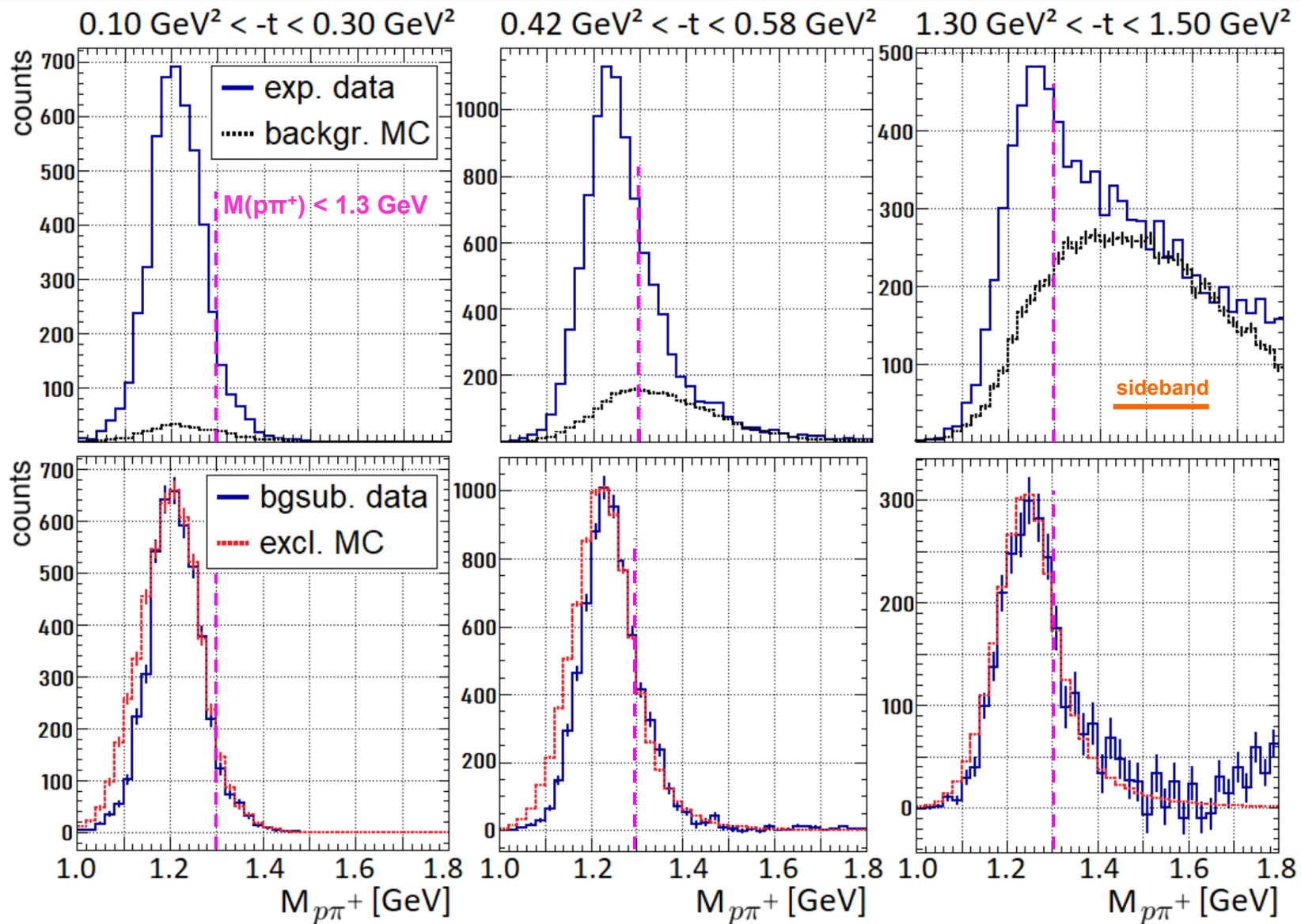
$$Q^2 > 1.5 \text{ GeV}^2$$

$$W > 2 \text{ GeV}$$

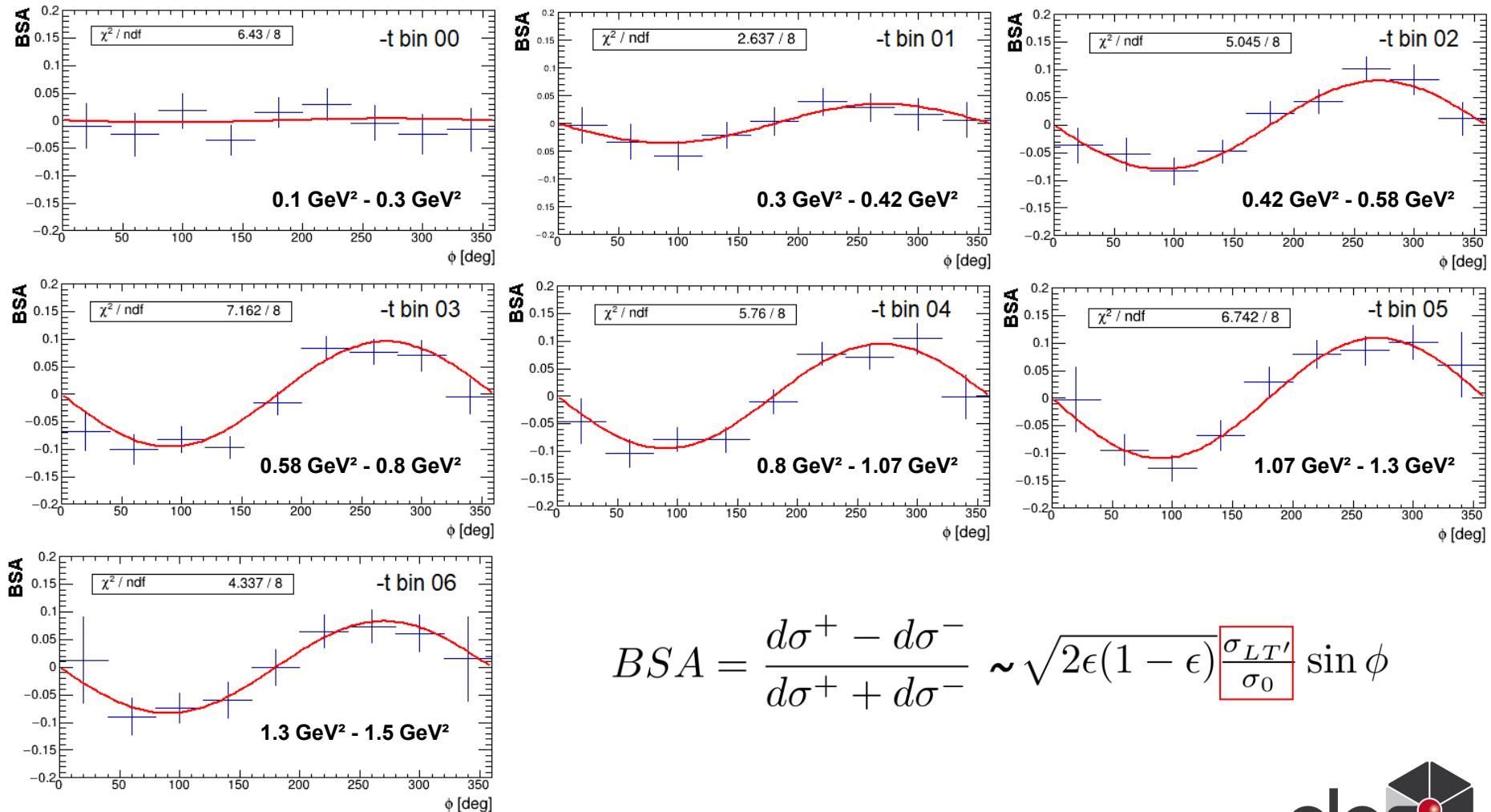
$$-t < 1.5 \text{ GeV}^2$$



$\pi^-\Delta^{++}$ DVMP: Signal and Background Separation



$\pi^-\Delta^{++}$ DVMP: Beam Spin Asymmetries (Q^2 - x_B integrated)



$$BSA = \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-} \sim \sqrt{2\epsilon(1-\epsilon)} \frac{\sigma_{LT'}}{\sigma_0} \sin \phi$$

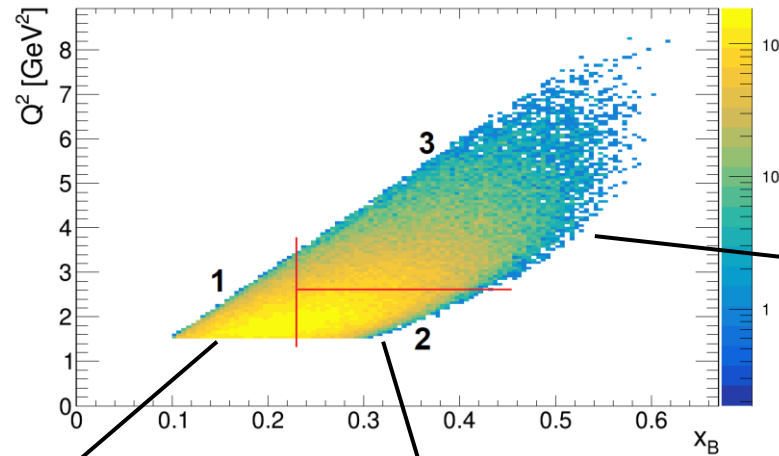


$\pi^- \Delta^{++}$ DVMP: Results



S. Diehl et al. (CLAS Collaboration),
 Phys. Rev. Lett. 131, 021901 (2023).
<https://doi.org/10.1103/PhysRevLett.131.021901>

■ sys. uncertainty

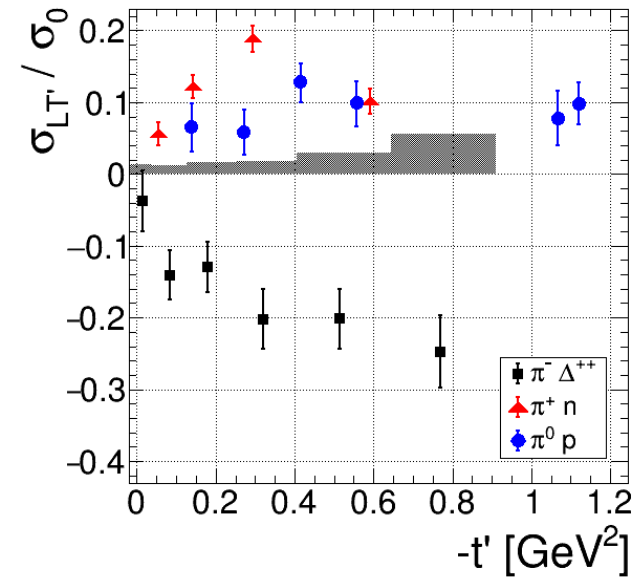
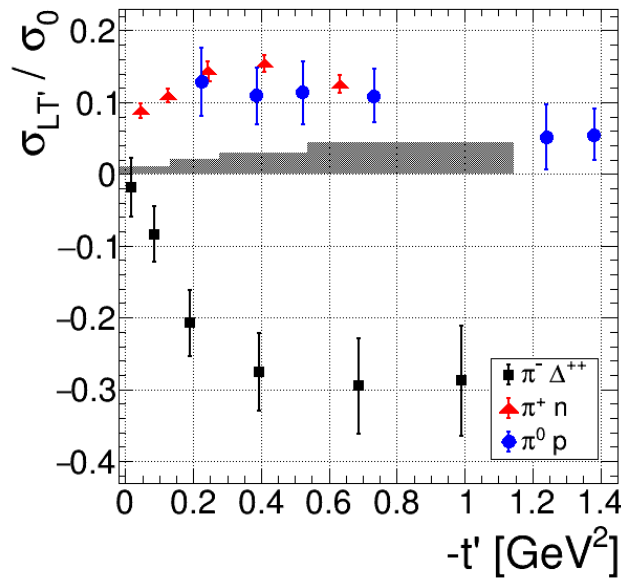
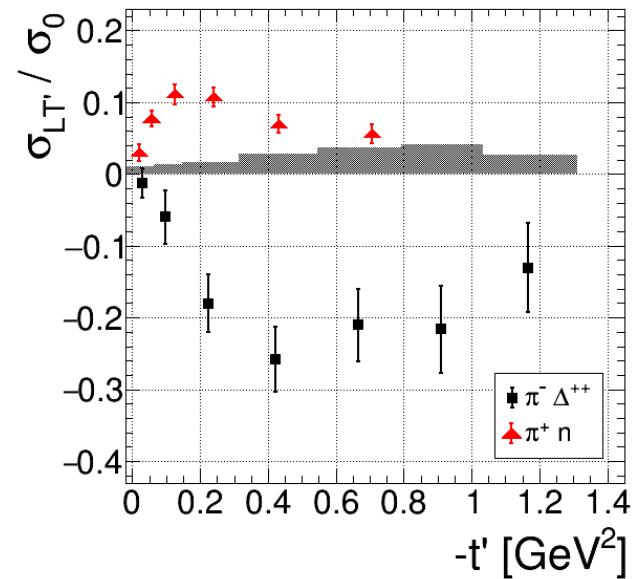


proton (uud)
 \rightarrow neutron (udd)
 π^+ ($|u\bar{d}\rangle$)
 $\rightarrow \Delta^{++}$ (uuu)
 π^- ($|d\bar{u}\rangle$)

bin 1 ($\langle Q^2 \rangle = 1.95 \text{ GeV}^2$, $\langle x_B \rangle = 0.19$)

bin 2 ($\langle Q^2 \rangle = 2.11 \text{ GeV}^2$, $\langle x_B \rangle = 0.28$)

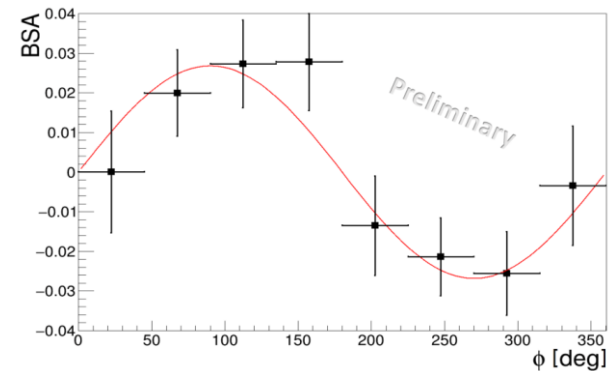
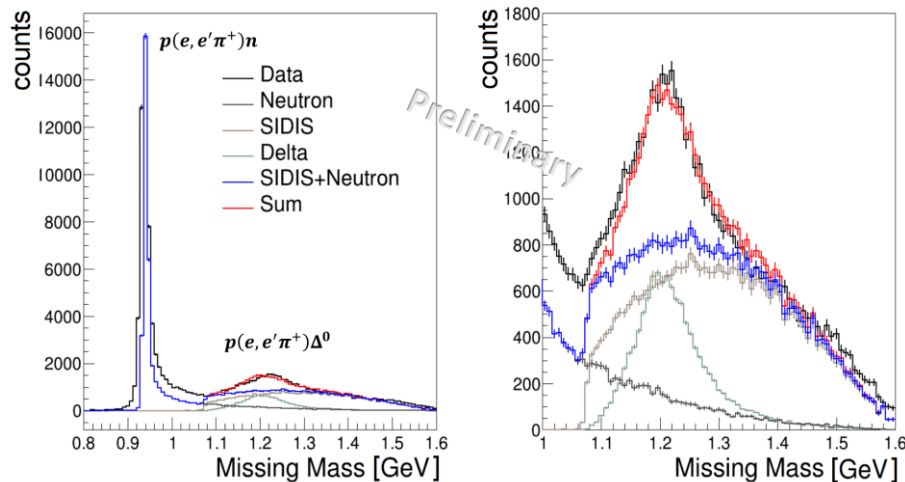
bin 3 ($\langle Q^2 \rangle = 3.38 \text{ GeV}^2$, $\langle x_B \rangle = 0.34$)



First Results on Other $N \rightarrow N^*$ DVMP Channels

JLAB hall C (high resolution two arm spectrometer): $e' \pi^+ \Delta^0$

- Electron and π^+ detected, final state reconstructed through missing mass
- Capability to perform high Q^2 longitudinal / transverse separation of the cross-section

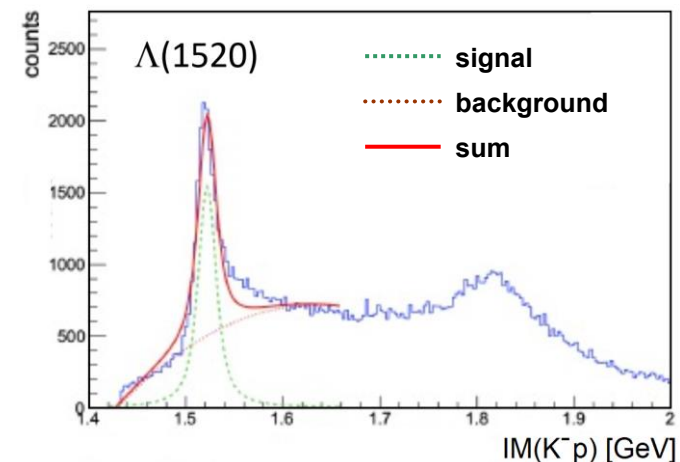


A. Usman (U. Regina)

CLAS12: Hyperon transitions

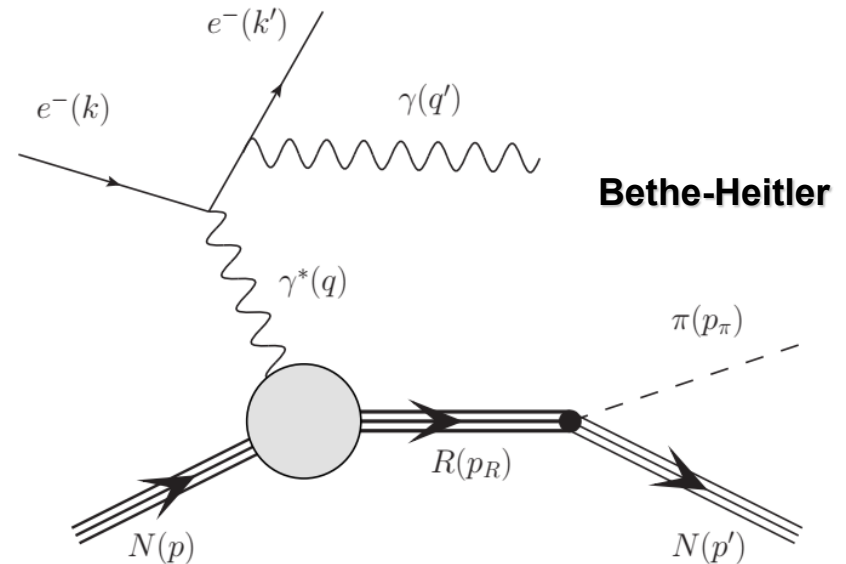
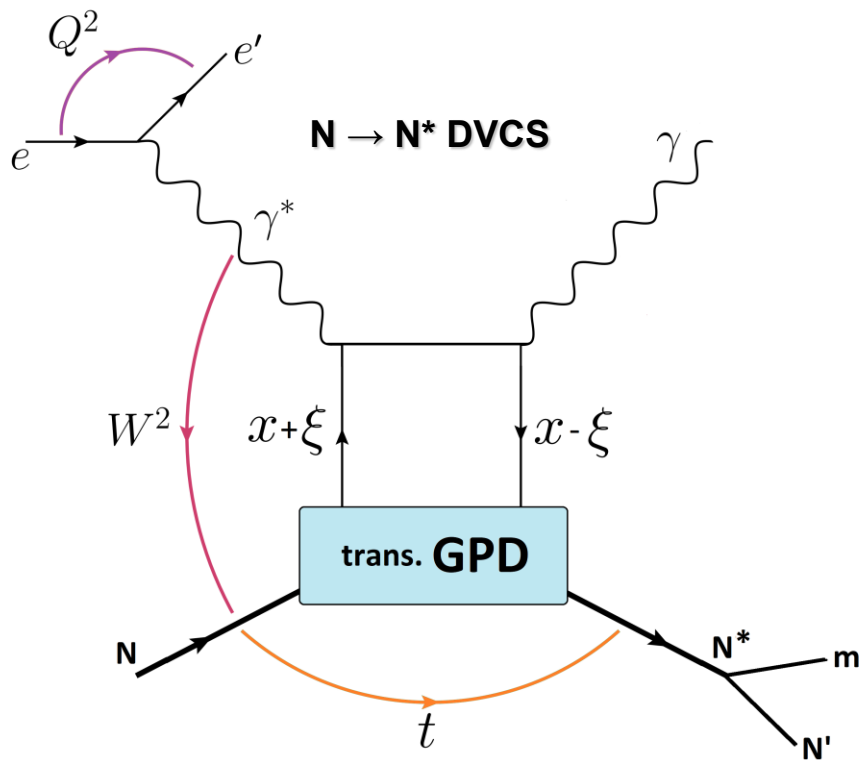
- Access to strange quark chiral-odd GPDs
- Exclusive $e p \rightarrow e' p K^+ K^-$ events

U. Shrestha (UConn)



The $N \rightarrow N^*$ DVCS Process

$$\gamma^* p \rightarrow N^* \gamma \rightarrow N \text{ meson } \gamma$$



➔ Sensitive to twist-2 transition GPDs

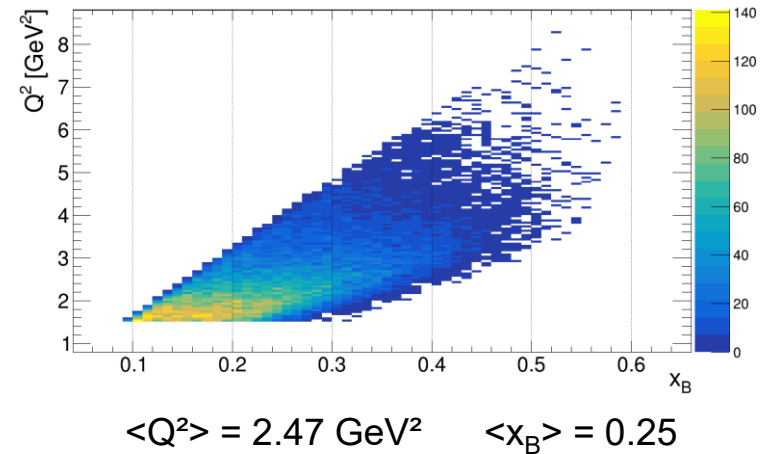
K. M. Semenov-Tian-Shansky, M. Vanderhaeghen,
Phys. Rev. D **108**, 034021 (2023)

N → N* DVCS: Event Selection

Event selection: $e p \rightarrow e' N^{*+} \gamma \rightarrow e' n \pi^+ \gamma$

- Exclusivity cuts on the missing masses, missing energy, missing transverse momentum and missing cone angle

Kinematic cuts: $W > 2 \text{ GeV}$ $Q^2 > 1.5 \text{ GeV}^2$
 $y < 0.8$ $-t < 2 \text{ GeV}^2$

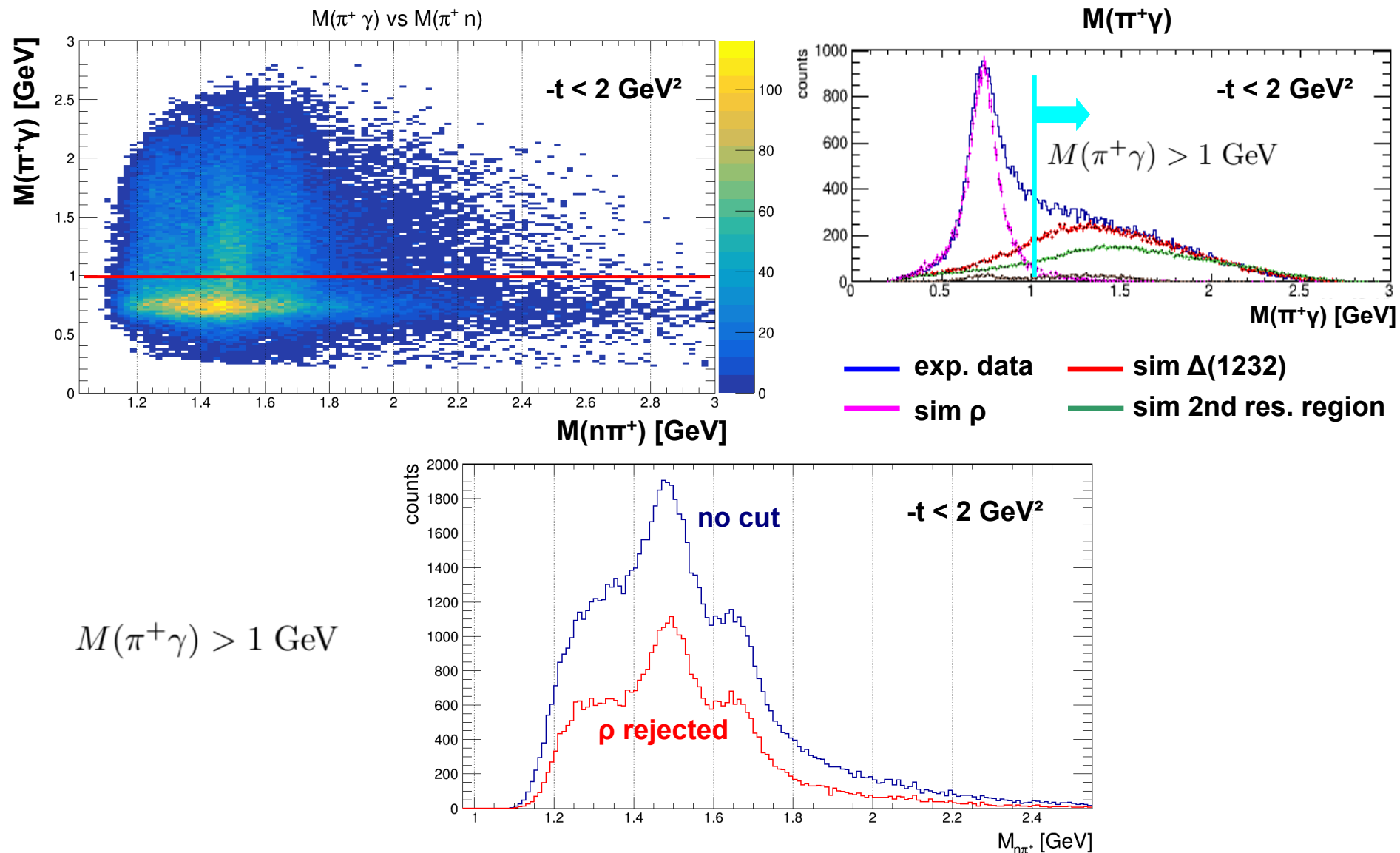


Physics background: $e p \rightarrow e' n \rho^+ \rightarrow e' n \pi^+ \gamma$ (very rare)

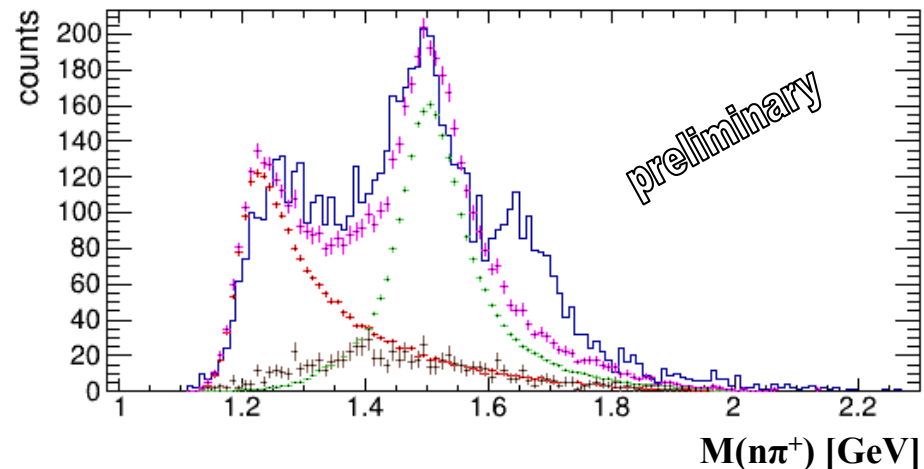
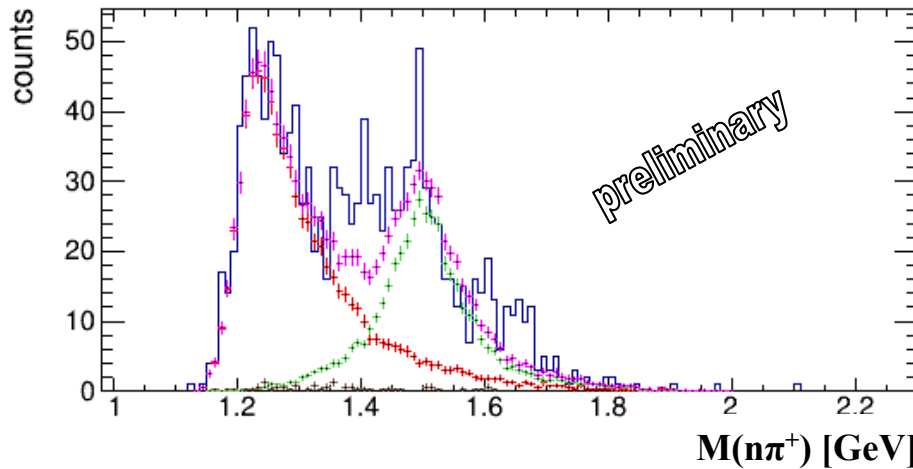
Event selection background: $e p \rightarrow e' n \rho^+ \rightarrow e' n \pi^+ \pi^0 \rightarrow e' n \pi^+ \gamma (\gamma)$
 → Can be suppressed (next slide)

$e p \rightarrow e' N^{*+} \pi^0 \rightarrow e' n \pi^+ \pi^0 \rightarrow e' n \pi^+ \gamma (\gamma)$
 → Needs to be subtracted bin by bin

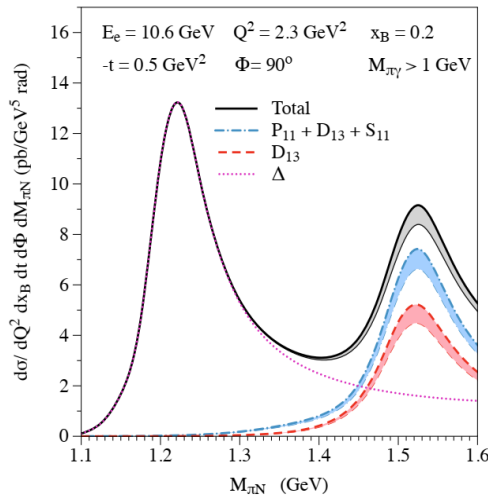
$N \rightarrow N^*$ DVCS: ρ^+ Background Rejection



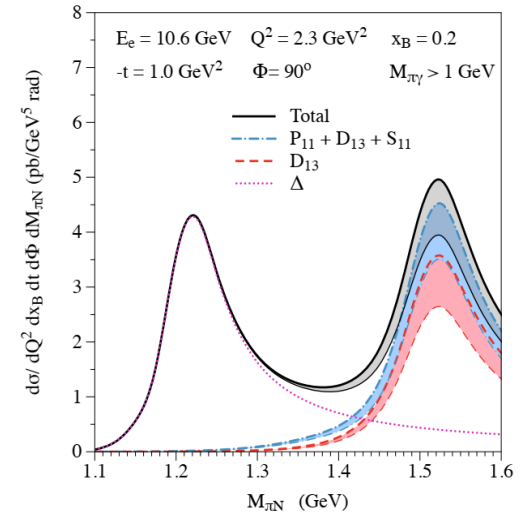
$N \rightarrow N^*$ DVCS: Resonance Mass Spectra

 $-t < 0.5 \text{ GeV}^2$
 $W > 2 \text{ GeV}, Q^2 > 1.5 \text{ GeV}^2$
 $0.5 \text{ GeV}^2 < -t < 1.0 \text{ GeV}^2$


— exp. data
 — sim. $\Delta(1232)$ / 2nd res. region
 — sim. full DIS
 — sim. sum



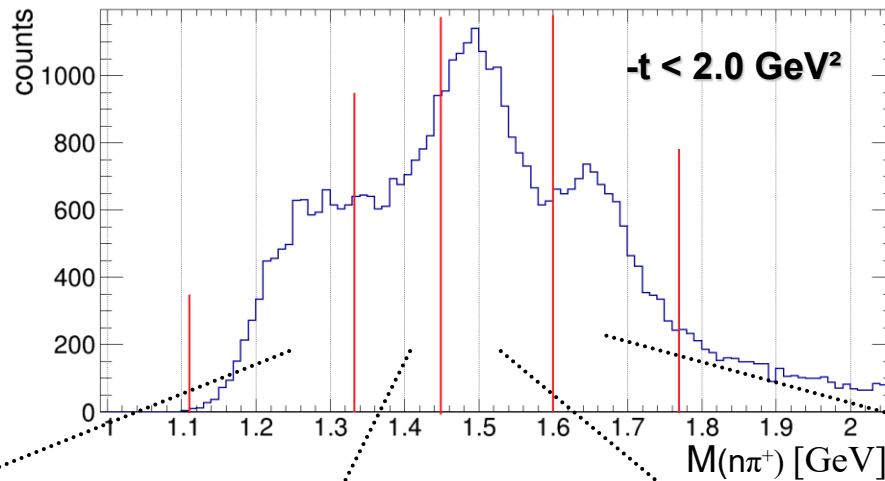
K. Semenov-Tian-Shansky,
M. Vanderhaeghen,
Phys. Rev. D 108, 034021 (2023)



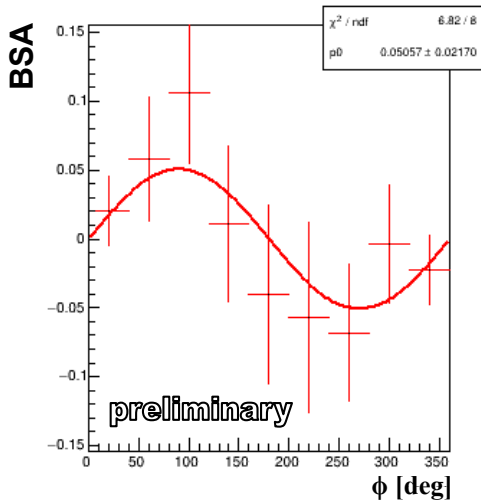
N → N* DVCS: Beam Spin Asymmetries



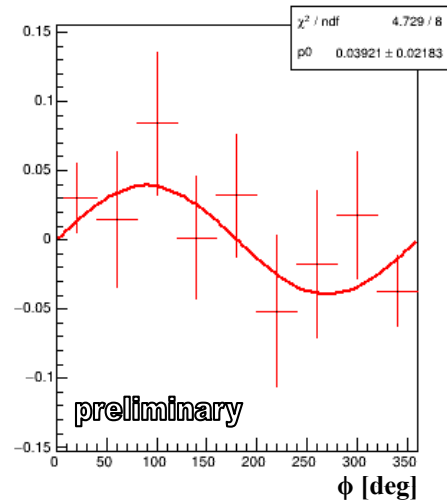
$$BSA = \frac{d\sigma^+ - d\sigma^-}{d\sigma^+ + d\sigma^-}$$



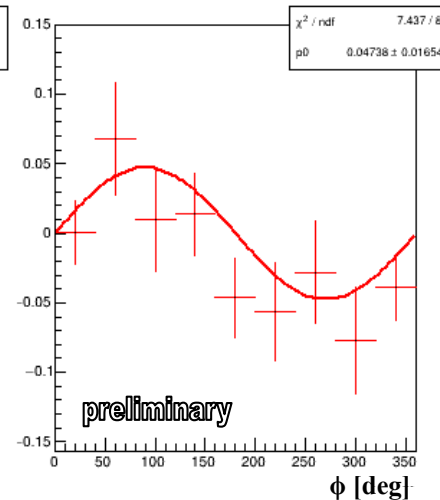
1.13 GeV < $M(n\pi^+) < 1.33$ GeV



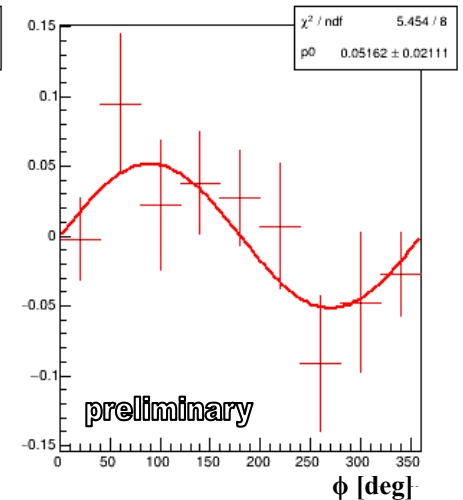
1.33 GeV < $M(n\pi^+) < 1.45$ GeV



1.45 GeV < $M(n\pi^+) < 1.60$ GeV

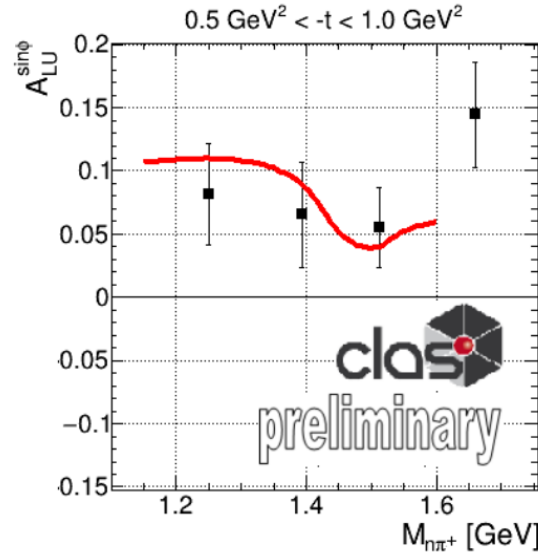
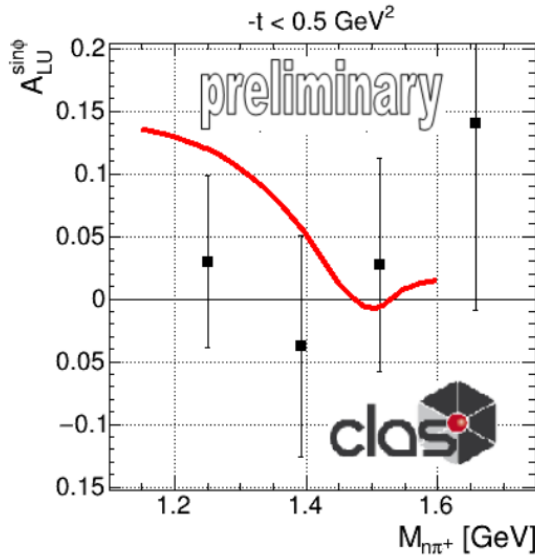


1.60 GeV < $M(n\pi^+) < 1.77$ GeV



$$BSA \sim A_{LU}^{\sin(\phi)} \cdot \sin(\phi)$$

Results for $N \rightarrow N^* \text{ DVCS}$ ($\langle Q^2 \rangle = 2.47 \text{ GeV}^2$, $\langle x_B \rangle = 0.25$)



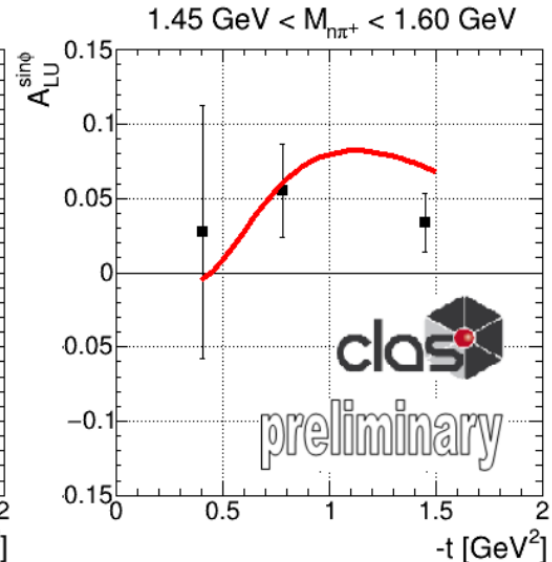
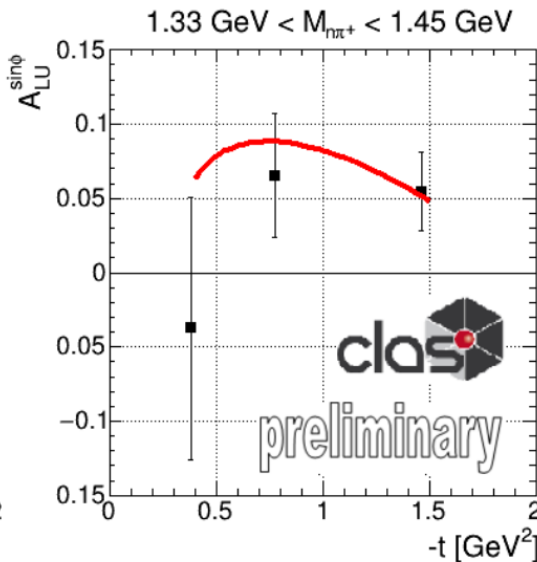
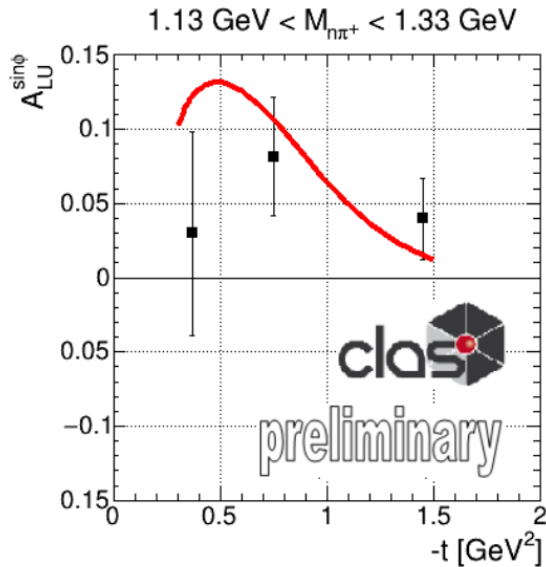
Theory curves in CLAS12

kinematics from:

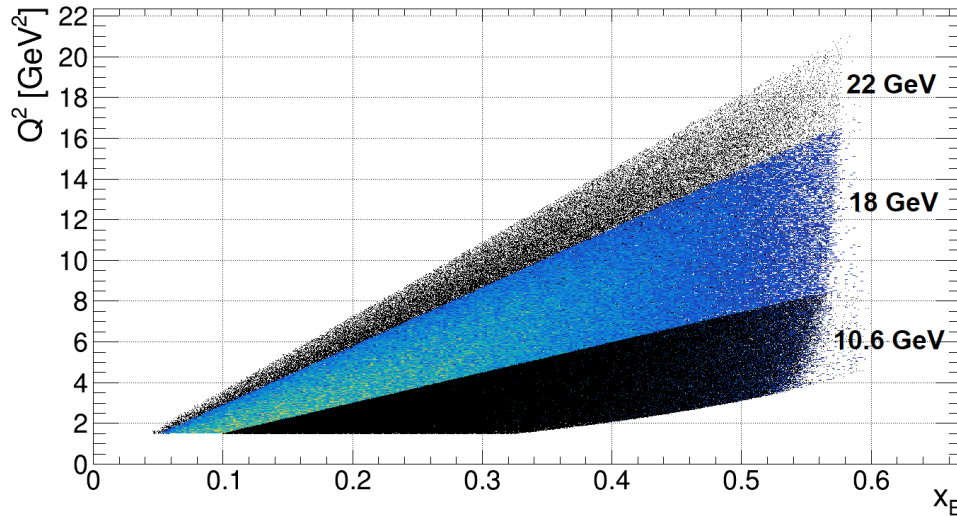
K. M. Semenov-Tian-Shansky, M. Vanderhaeghen, Phys. Rev. D **108**, 034021 (2023)

Data:

No π^0 background subtraction so far!



Future: Perspectives for a 22 GeV JLAB Upgrade



$$ep \rightarrow e\Delta^{++}\pi^- \rightarrow ep\pi^+\pi^-$$

$$ep \rightarrow e\Delta^+\gamma \rightarrow en\pi^+\gamma$$

Extended Q^2 range

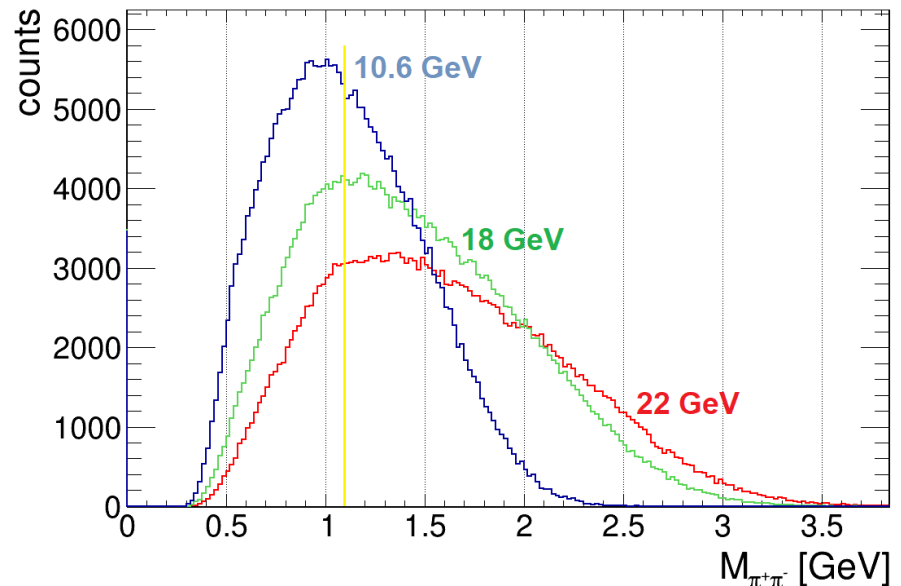
➔ **Advantage for factorisation**

Better signal / background separation

➔ **Higher efficiency**

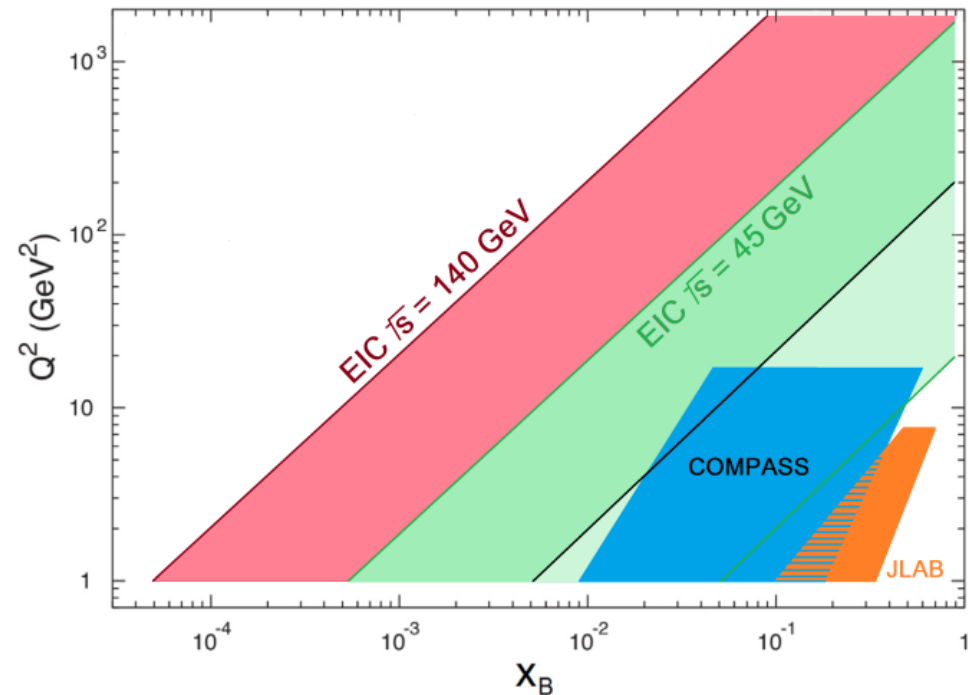
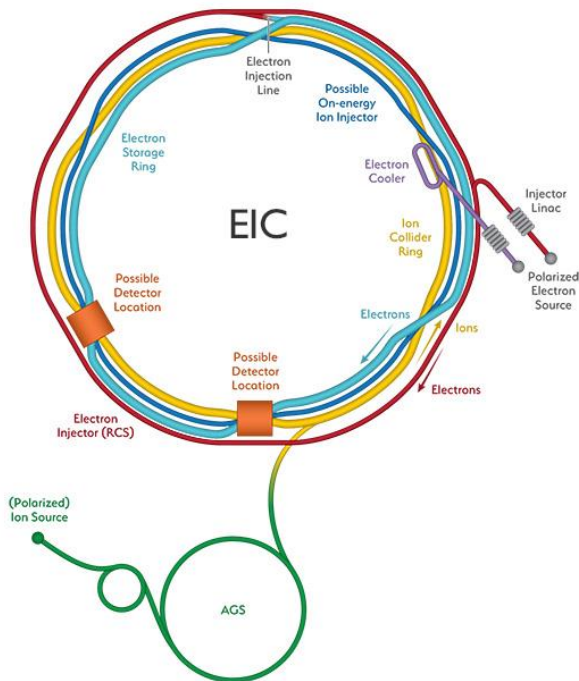
Transition GPDs are a potential part of the science program for a 22 GeV JLAB upgrade:

A. Accardi, P. Achenbach, D. Adhikari et al., Strong Interaction Physics at the Luminosity Frontier with 22 GeV Electrons at Jefferson Lab (2023). <https://doi.org/10.48550/arXiv.2306.09360>



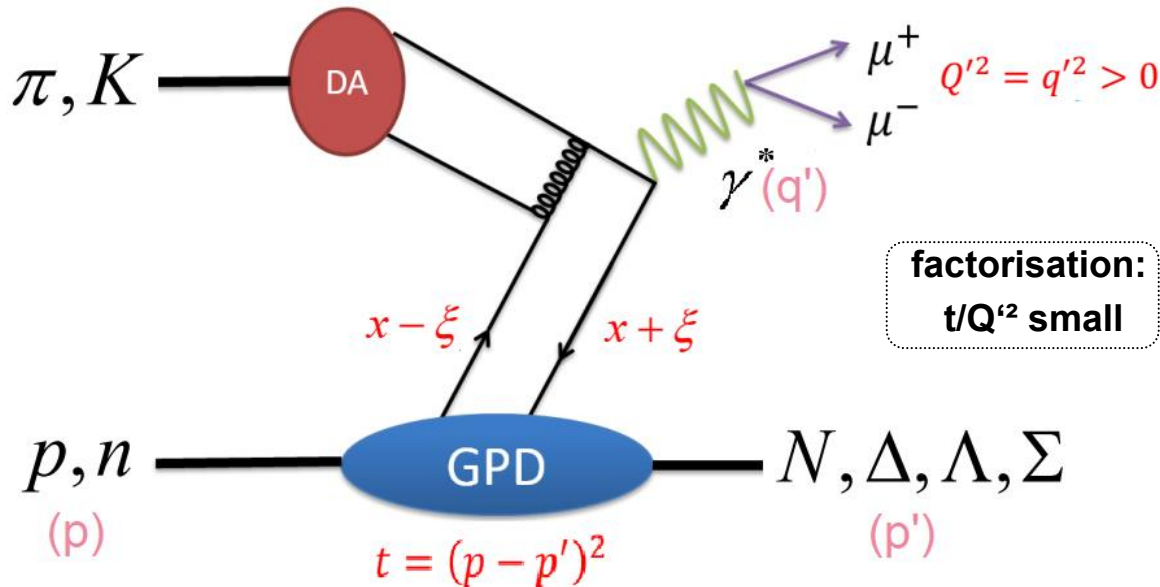
Future: Transition GPDs at EIC and EicC

- Extension of the kinematic regime to the sea-quark and gluonic sector
 → Low x_B and higher Q^2 values



- Potential for unique insights into the contributions of sea quarks and gluons to the excitation process and to the characteristics of baryon resonances

Future: Transition Processes in Hadron Scattering



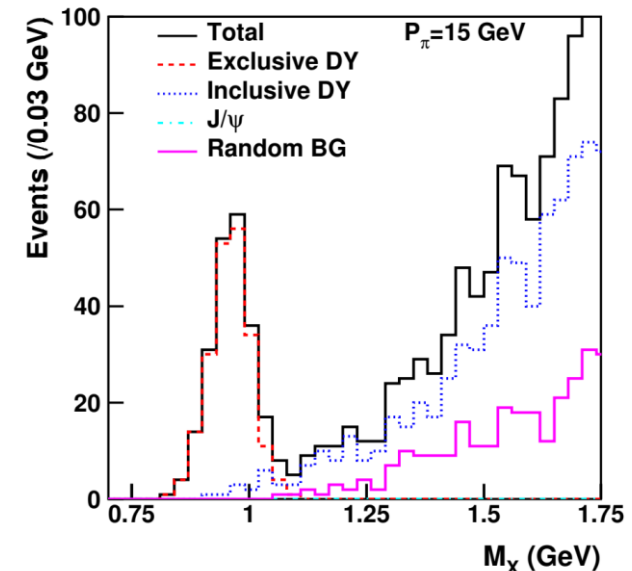
- $\pi^- p \rightarrow \gamma^* \Delta^0$
- $\pi^- n \rightarrow \gamma^* \Delta^-$
- $\pi^+ p \rightarrow \gamma^* \Delta^{++}$
- $\pi^+ n \rightarrow \gamma^* \Delta^+$
- $K^- p \rightarrow \gamma^* \Lambda$
- $K^- p \rightarrow \gamma^* \Lambda(1405)$
- $K^- p \rightarrow \gamma^* \Lambda(1520)$
- $K^- n \rightarrow \gamma^* \Sigma^-$
- $K^+ n \rightarrow \gamma^* \Theta^+$

→ Probes universality of GPDs and trans. GPDs
+ Different kinematic regime in x and ξ than e-scatt.

J-PARC (up to 20 GeV/c π, K beams):

- Large acceptance and good momentum resolution spectrometer system

$e\mu^+\mu^-X$ ($X = p$) missing mass:



S. Kumano + W.C. Chang

- Potentially also with **AMBER @ CERN** (up to 190 GeV/c π, K beams)

Summary and Outlook

- Transition GPDs can help us to better understand baryon resonances and the excitation process itself, by relating their properties to the 3D motion and distribution of the partons
- Hard exclusive $\pi\Delta^{++}$ production and $N\rightarrow N^*$ DVCS can be well measured with CLAS12 (first published observable sensitive to transition GPDs)
- First results on $N\rightarrow N^*$ DVCS BSAs show a very promising agreement with theory predictions
- A JLAB energy upgrade will help to significantly improve these measurements and the extraction of transition GPDs
- Further opportunities: EIC, COMPASS / AMBER, J-PARC

Summary and Outlook



The first workshop on transition GPDs took place in August 2023 at ECT* Trento



<https://indico.ectstar.eu/event/176/>

- A whitepaper has been submitted to EPJ-A (under review) and to arXiv: <http://arxiv.org/abs/2405.15386>

Exploring Baryon Resonances with Transition Generalized Parton Distributions: Status and Perspectives

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