

## Unraveling the unknown

Nearly all the visible matter in the Universe consists of hadrons, and yet much is still unknown about them.

Their internal structure is described by a **non abelian gauge field theory** of quarks and gluons with **SU(3) color symmetry**, and no direct observation is allowed because of their **confined nature**.

## 3-dimensional maps

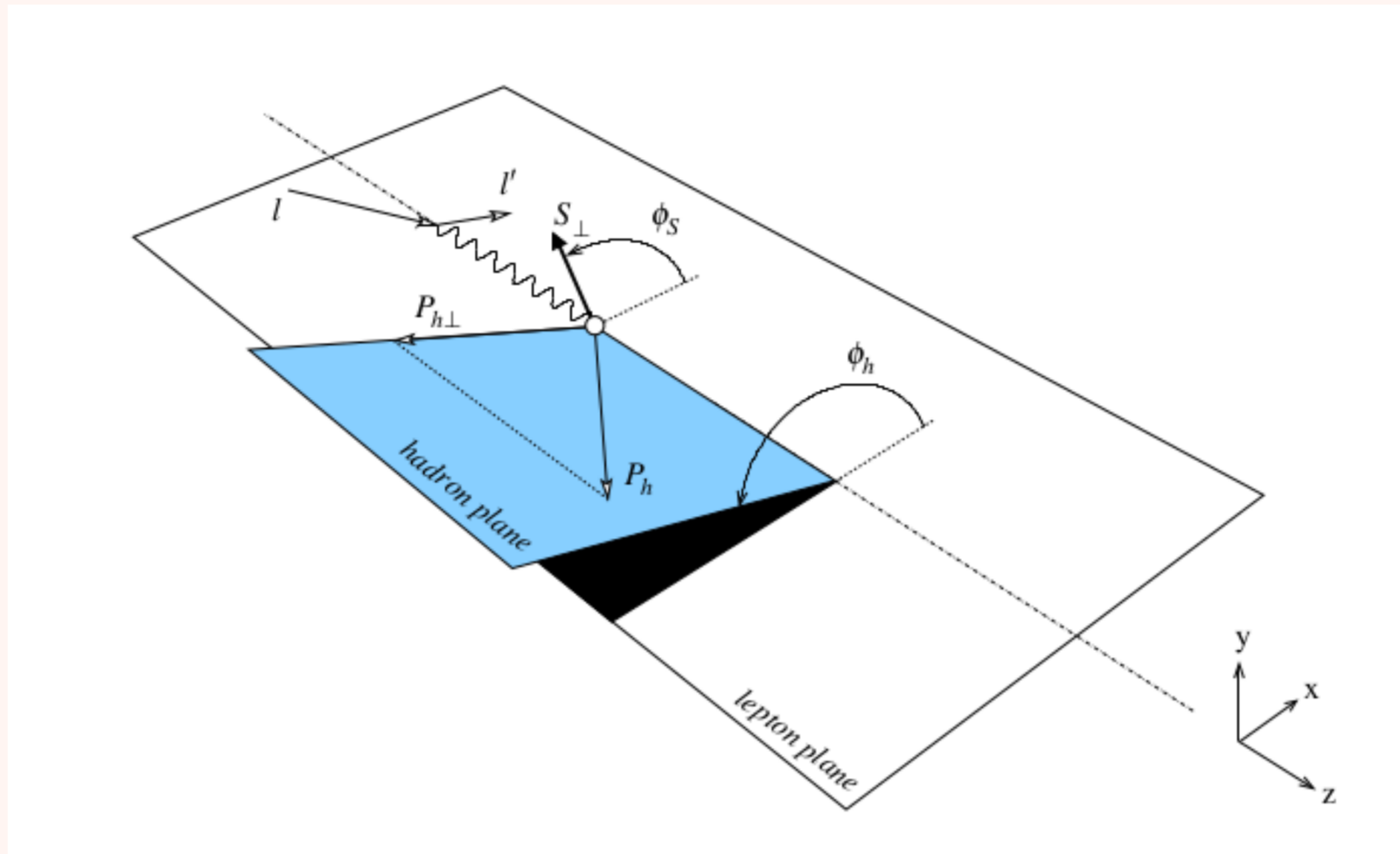
In order to get informations on the internal structure, physicists can

- probe the hadrons in **high-energy scattering processes**
- develop **maps in position and momentum space** to build an accurate description

These maps are called **PDFs** and **FFs**, and they show a dependence on the *longitudinal momentum fraction*  $x$  and the *intrinsic transverse momentum*  $k_T$  of partons

## SIDIS

$$\ell(l) + N(P) \rightarrow \ell(l') + h(P_h) + X$$



## SIDIS variables

$$x = \frac{Q^2}{2P \cdot q} \quad y = \frac{P \cdot q}{P \cdot l} \quad z = \frac{P \cdot P_h}{P \cdot l} \quad \gamma = \frac{2Mx}{Q}$$

- $q = l - l'$ , with  $Q^2 = -q^2 > 0$
- In the region where  $q_T \ll Q$ , factorization theorems are valid

The aim of my thesis is to improve the knowledge of the 3D structure of hadrons exploiting the phenomenology of **longitudinally polarized SIDIS** through the implementation of **Nanga Parbat**, a fitting framework developed in Pavia by A. Bacchetta et al.

## Defining the observables

Assuming one photon exchange, the **cross section** can be expressed in terms of a set of structure functions.

$$F_{UT,L} \quad \begin{cases} U \rightarrow \text{beam polarization} \\ T \rightarrow \text{target polarization} \\ L \rightarrow \text{photon polarization} \end{cases}$$

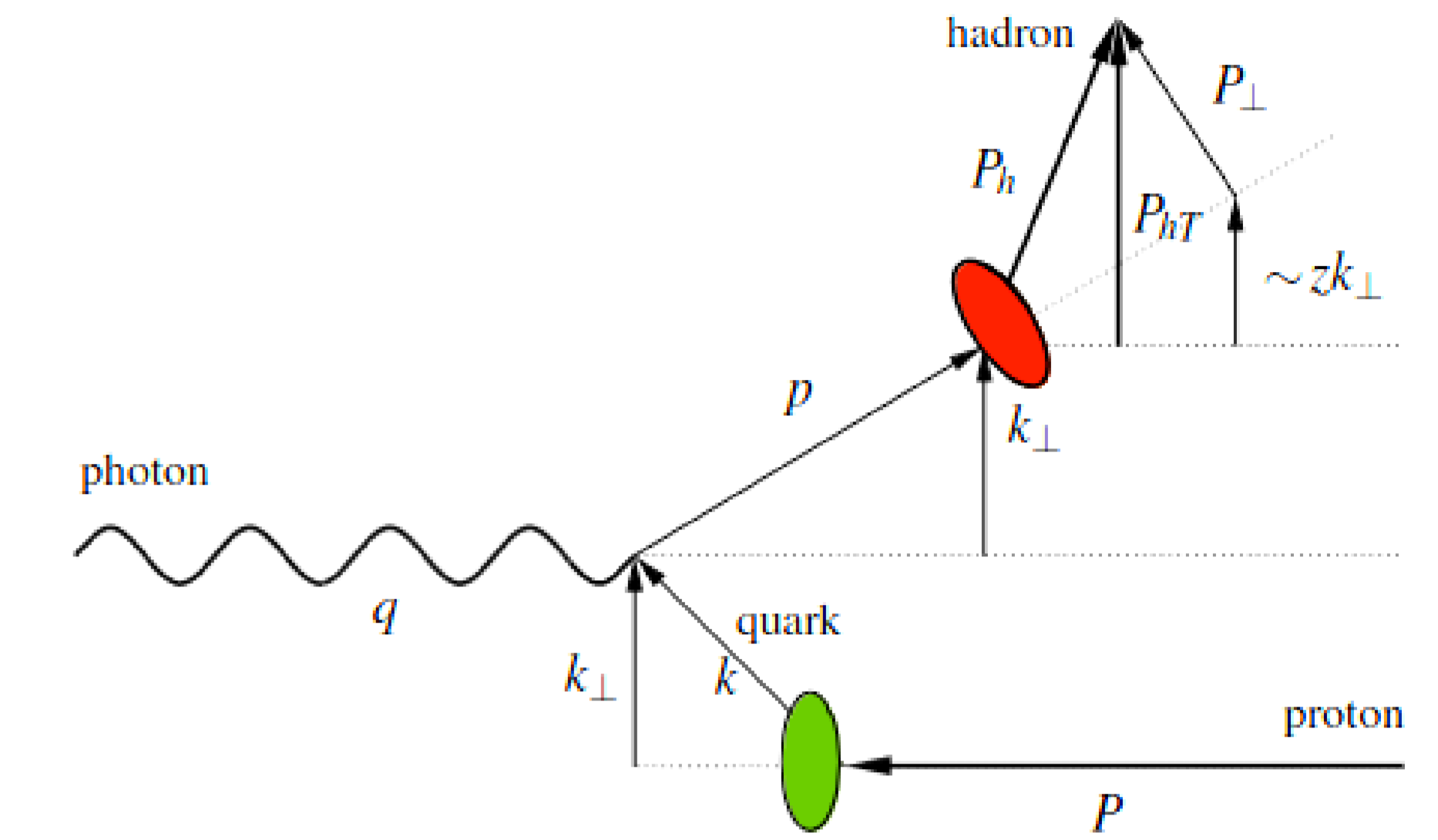
Asymmetries are good observables in **polarized processes**

$$A_{LL} = \frac{F_{LL}}{F_{UU,T}}$$

where

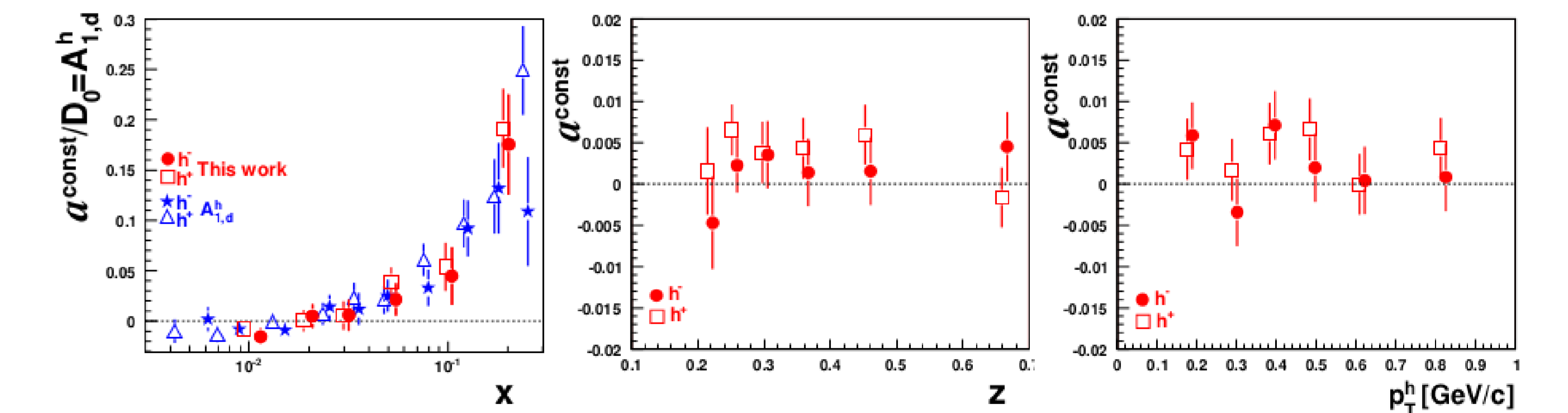
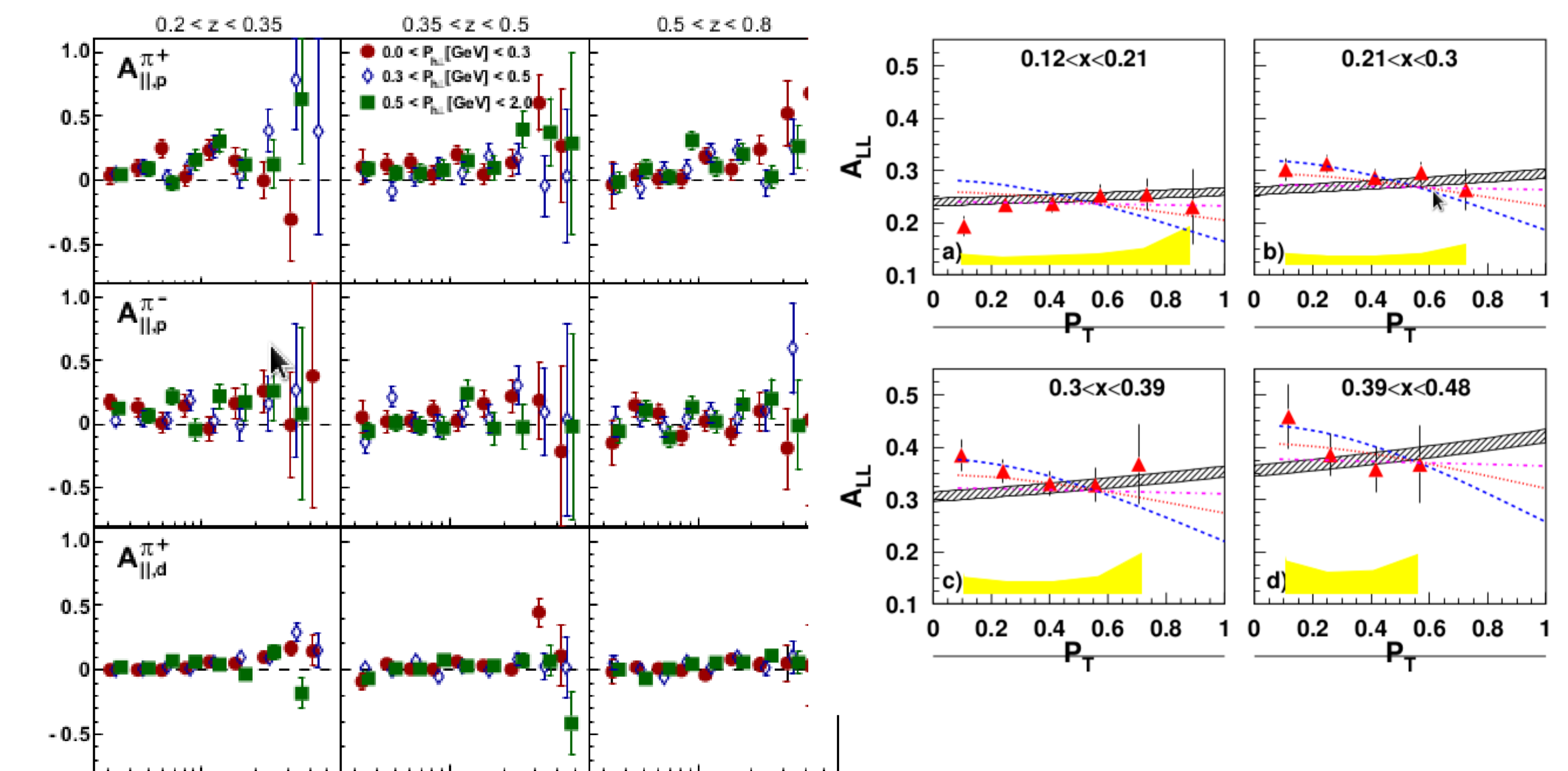
$$F_{LL} = \mathcal{C}[g_{1L}D_1] \quad F_{UU,T} = \mathcal{C}[f_1D_1]$$

## Structure of TMDs



$$F_{LL}(x, z, P_{hT}^2, Q^2) = x \sum_a \mathcal{H}_{LL}^a(Q^2, \mu^2) \int d^2k_{\perp} d^2P_{\perp} g_1^a(x, k_{\perp}^2; \mu^2) D_1^a(z, P_{\perp}^2; \mu^2) \delta^{(2)}(zk_{\perp} - P_{hT} + P_{\perp})$$

## Asymmetry: measurements



..and more complete data are coming from CLAS12!

## References

A. Bacchetta et al., *Unpolarized transverse momentum distributions from a global fit of Drell-Yan and semi-inclusive deep inelastic scattering data*, JHEP 10 (2022), arXiv:2206.07598v2 [hep-ph]

A. Bacchetta, M. Diehl, K. Goeke, A. Metz, P.J. Mulders and M. Schlegel, *Semi-inclusive deep inelastic scattering at small transverse momentum*, JHEP 02 (2007) 093 [hep-ph/0611265]

M. G. Alekseev et al. [COMPASS], arXiv:1007.1562 [hep-ex]

S. Jawalkar et al. [CLAS], arXiv:1709.10054 [nucl-ex]

		Quark polarization		
		Unpolarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 = \odot$	*	naive T-odd $h_1^\perp = \uparrow - \downarrow$ Boer-Mulders
	L	*	$g_1 = \rightarrow - \leftarrow$ helicity	$h_{1L}^\perp = \rightarrow - \leftarrow$ transversity
	T	naive T-odd $f_{1T}^\perp = \odot - \ominus$ Sivers	$g_{1T} = \uparrow - \downarrow$	$h_1 = \uparrow - \downarrow$ $h_{1T}^\perp = \rightarrow - \leftarrow$

		Quark polarization		
		Unpolarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$D_1 = \odot$		naive T-odd $H_1^\perp = \odot - \ominus$ Collins
	L		$G_{1L} = \rightarrow - \leftarrow$	$H_{1L}^\perp = \rightarrow - \leftarrow$
	T	naive T-odd $D_{1T}^\perp = \odot - \ominus$	$G_{1T} = \uparrow - \downarrow$	$H_1 = \uparrow - \downarrow$ $H_{1T}^\perp = \rightarrow - \leftarrow$