

Quarkonia as tools to study the multi-dimensional structure of the nucleon

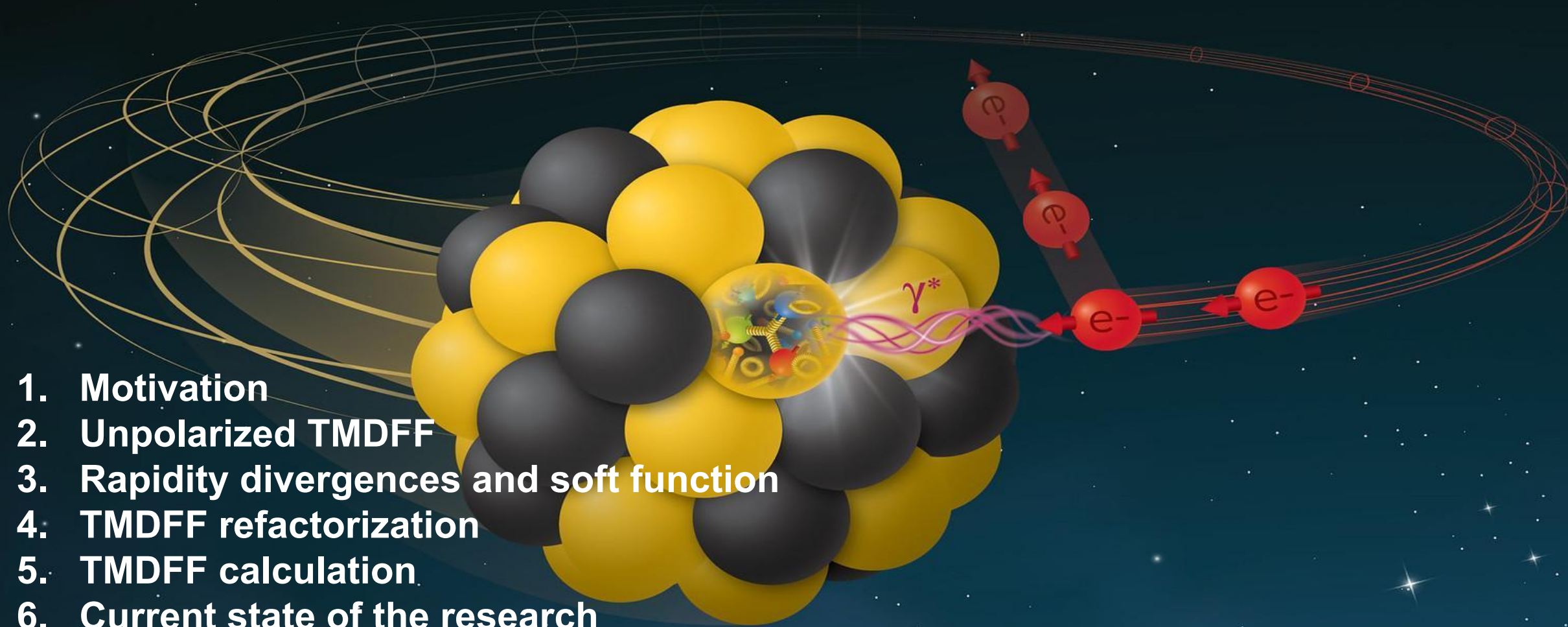
Samuel F. Romera

University of the Basque Country (EHU)

Hampton University Graduate Studies (HUGS) Program at
Jefferson Lab, 2023



Outline



1. Motivation
2. Unpolarized TMDFF
3. Rapidity divergences and soft function
4. TMDFF refactorization
5. TMDFF calculation
6. Current state of the research

Why quarkonia?

From an **experimental** point of view...

- Easy to produce and detect \longleftarrow
- Quarkonia production is one of the processes on which the EIC is most focused

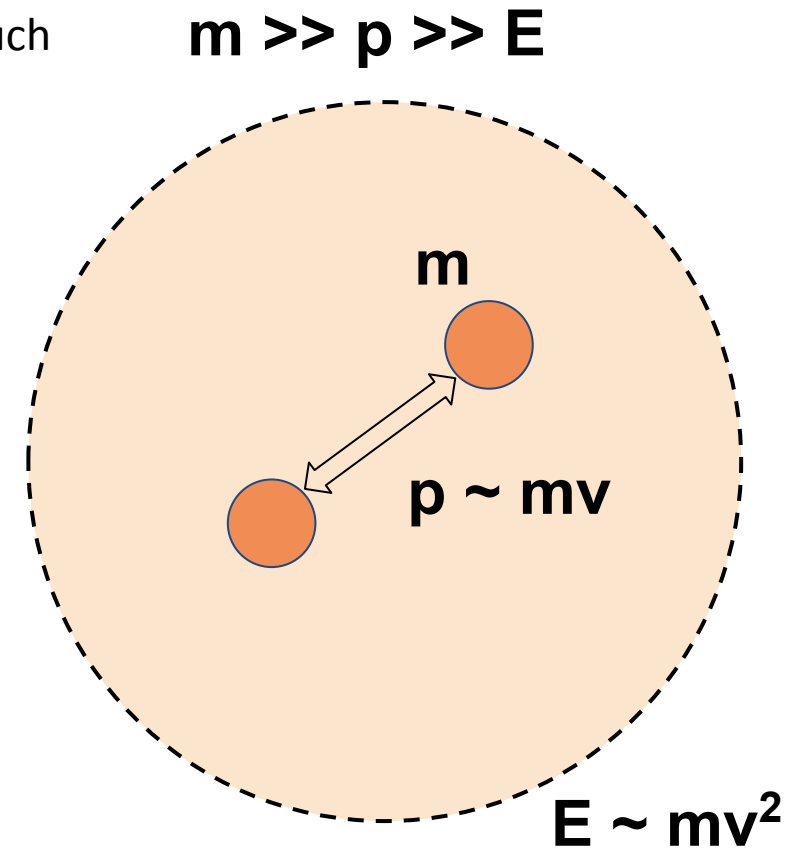
In particular quarkonium states such as J/ψ , $\psi(2S)$ and $\Upsilon(nS)$

From a **theoretical** point of view...

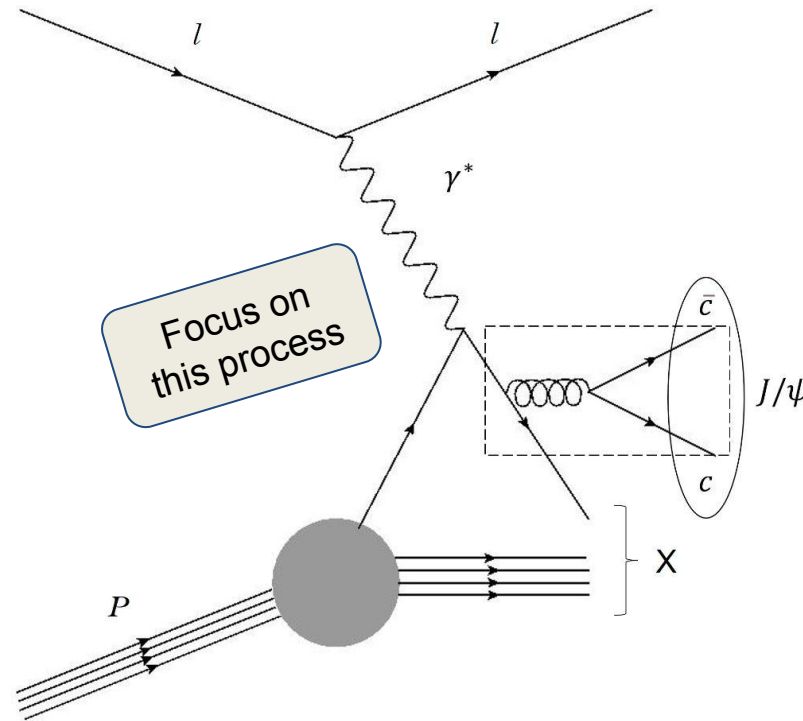
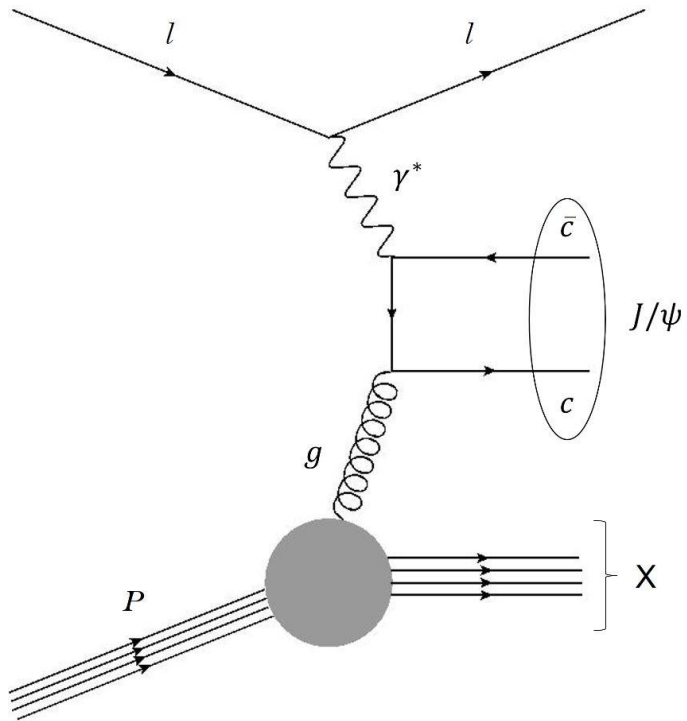
- Quarkonia are useful laboratories to study the interplay between pQCD and non-perturbative QCD
- There are much more physics yet to be understood and much more work yet to be done in quarkonium production



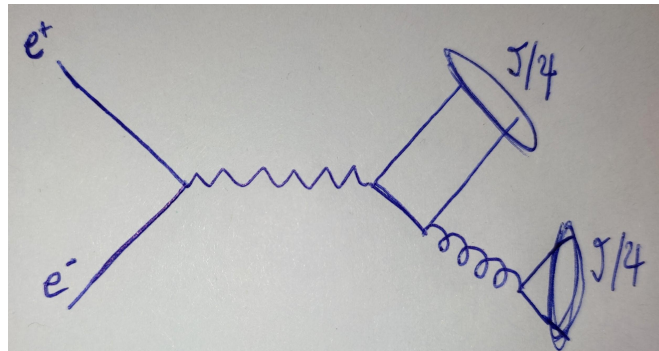
We'll see later!



Goal



- The main motivation is complete the J/ψ production in ep scattering
- The contribution by the light-quark fragmentation have been done in [arXiv:2007.05547](https://arxiv.org/abs/2007.05547) [hep-ph]
- Typical process carried by EIC: J/ψ production one of the most important



But we can also find the contribution of gluon TMDFF in other processes!

Unpolarized TMDFF

$$\tilde{f}_{q/\mathcal{N}}(x, b_T) = \int \frac{db^-}{4\pi} e^{-ixP^+b^-} \text{Tr} \left[\langle \mathcal{N} | \bar{\psi}_q(b) \gamma^+ \mathcal{W}(b, 0) \psi_q(0) | \mathcal{N} \rangle \right]$$

Thank you!

Patrick Barry's lecture...

How can we describe the process when a quark is produced in a hard interaction and then fragments into a detected hadron? TMD fragmentation function

$$\Delta_{g \rightarrow J/\psi}(z, \mathbf{b}_T) \sim \sum_X \int \frac{db^-}{4\pi} e^{-iP^+b^-/z} \langle 0 | T [\mathcal{B}_{n\perp}^\mu](b) | X, J/\psi \rangle \langle X, J/\psi | \bar{T} [\mathcal{B}_{n\perp\mu}](0) | 0 \rangle$$

$$D_{g \rightarrow J/\psi} \sim \boxed{Z_g \times R_g} \times \Delta_{g \rightarrow J/\psi}(z, \mathbf{b}_T)$$

Renormalization factors

Gluon field $\mathcal{B}_{n\perp}^\mu = \frac{1}{g} [\mathcal{W}_n^\dagger(b) i D_{n\perp}^\mu \mathcal{W}_n(b)]$

Overview of rapidity divergences and Soft Function (SF)

Cause

Our theory is boost invariant in the light-cone directions



Consequence

$$k^+ \rightarrow ak^+, k^- \rightarrow k^-/a, k^+k^-$$

$$a \rightarrow \infty \Rightarrow \text{Divergence}$$



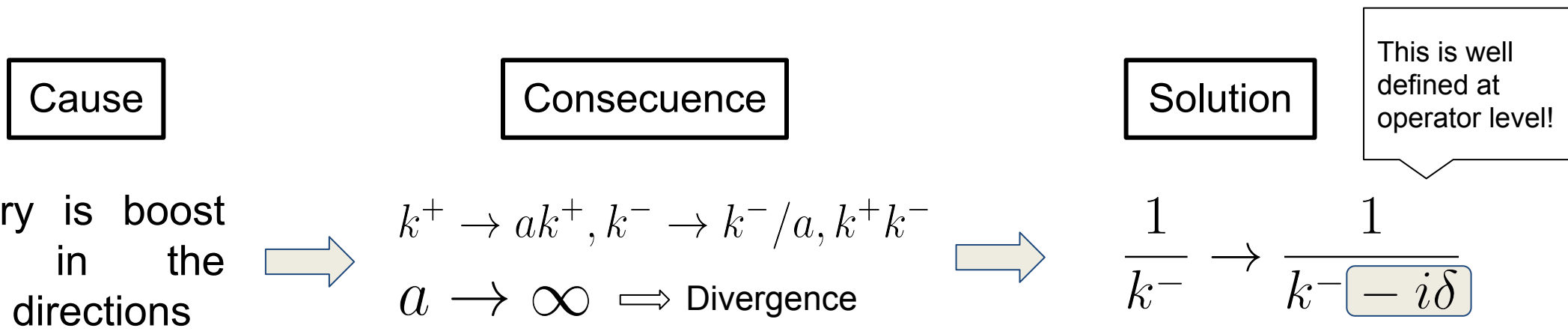
Solution

$$\frac{1}{k^-} \rightarrow \frac{1}{k^- - i\delta}$$

This is well defined at operator level!

Now the rapidity divergences are described by the regulator. We obtain something like that... $\ln\left(\frac{\delta^2}{Q^2}\right)$

Overview of rapidity divergences and Soft Function (SF)



Now the rapidity divergences are described by the regulator. We obtain something like that... $\ln\left(\frac{\delta^2}{Q^2}\right)$

At this point, the Soft Function enters in the game to cancel all the spurious rapidity divergences.

$$S(\mathbf{b}_\perp) = \frac{Tr}{N_c} \langle 0 | [S_n^\dagger S_{\bar{n}}] (0^+, 0^-, \mathbf{b}_\perp) [S_{\bar{n}}^\dagger S_n] (0^+, 0^-, \mathbf{0}) | 0 \rangle$$

Gluon TMDFF refactorization

Quarkonia, bound states of heavy quarks: massive quarks \rightarrow perturbative QCD. The heavy quarks are not relativistic: $v \ll 1 \rightarrow$ QCD turns non-perturbative

We have two scales strongly separated describing the quarkonium production.

Non-relativistic QCD (NRQCD)

NRQCD

$$\sigma^H = \sum_n \sigma_{Q\bar{Q}(n)} \langle 0 | \mathcal{O}^H(n) | 0 \rangle$$

$n = 2S+1 L_J^{[col.]}$

Can be organized in powers of v

gTMDFF in NRQCD

$$D_{g \rightarrow J/\psi}(z, \mathbf{b}_T) = \sum_n d_{g \rightarrow Q\bar{Q}(n)}(z, \mathbf{b}_T) \langle 0 | \mathcal{O}^{J/\psi}(n) | 0 \rangle$$

SDC

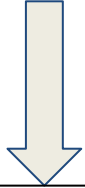
LDME \rightarrow Universals

$$\langle \mathcal{O}_n^{J/\psi} \rangle = \langle 0 | \chi^\dagger \mathcal{K}_n \psi a_{J/\psi}^\dagger a_{J/\psi} \psi^\dagger \mathcal{K}'_n \chi | 0 \rangle$$

$$D_{g \rightarrow J/\psi}(z, \mathbf{b}_T) = \sum_n d_{g \rightarrow Q\bar{Q}(n)}(z, \mathbf{b}_T) \langle 0 | \mathcal{O}^{J/\psi}(n) | 0 \rangle$$

TMDFF calculation

$$D_{g \rightarrow J/\psi}(z, \mathbf{b}_T) = \sum_n d_{g \rightarrow Q\bar{Q}(n)}(z, \mathbf{b}_T) \langle 0 | \mathcal{O}^{J/\psi}(n) | 0 \rangle$$



- QCD calculation at next to leading order in α_s
- Expansion around the threshold ($q = 0$)

$$\bar{u}(p)v(\bar{p}) = -2 \xi^\dagger(\mathbf{q} \cdot \boldsymbol{\sigma})\eta,$$

$$\bar{u}(p)\gamma^\mu v(\bar{p}) = L^\mu_j \left(\boxed{2E_q \xi^\dagger \sigma^j \eta} - \frac{2}{E_q + m_c} q^j \xi^\dagger(\mathbf{q} \cdot \boldsymbol{\sigma})\eta \right),$$

$$\begin{aligned} \bar{u}(p)(\gamma^\mu \gamma^\nu - \gamma^\nu \gamma^\mu)v(\bar{p}) &= (P^\mu L^\nu_j - P^\nu L^\mu_j) \left(\frac{2m_c}{E_q} \xi^\dagger \sigma^j \eta + \frac{2}{E_q(E_q + m_c)} q^j \xi^\dagger(\mathbf{q} \cdot \boldsymbol{\sigma})\eta \right) \\ &\quad + L^\mu_j L^\nu_k \xi^\dagger \{[\sigma^j, \sigma^k], \mathbf{q} \cdot \boldsymbol{\sigma}\} \eta, \end{aligned}$$

$$\bar{u}(p)(\gamma^\mu \gamma^\nu \gamma^\lambda - \gamma^\lambda \gamma^\nu \gamma^\mu)v(\bar{p})$$

$$\begin{aligned} &= L^\mu_i L^\nu_j L^\lambda_k \left(-E_q \xi^\dagger \{[\sigma^i, \sigma^j], \sigma^k\} \eta + \frac{q^i}{E_q + m_c} \xi^\dagger \{[\sigma^j, \sigma^k], \mathbf{q} \cdot \boldsymbol{\sigma}\} \eta \right. \\ &\quad \left. + \frac{q^j}{E_q + m_c} \xi^\dagger \{[\sigma^k, \sigma^i], \mathbf{q} \cdot \boldsymbol{\sigma}\} \eta + \frac{q^k}{E_q + m_c} \xi^\dagger \{[\sigma^i, \sigma^j], \mathbf{q} \cdot \boldsymbol{\sigma}\} \eta \right) \\ &\quad - \frac{2}{E_q} (P^\mu L^\nu_i L^\lambda_j + L^\mu_i L^\nu_j P^\lambda + L^\mu_j P^\nu L^\lambda_i) (\xi^\dagger q^i \sigma^j \eta - \xi^\dagger q^j \sigma^i \eta). \end{aligned}$$

TMDF calculation

$$D_{g \rightarrow J/\psi}(z, \mathbf{b}_T) = \sum_n d_{g \rightarrow Q\bar{Q}(n)}(z, \mathbf{b}_T) \langle 0 | \mathcal{O}^{J/\psi}(n) | 0 \rangle$$



Perturbative
NRQCD

$$\langle \chi^\dagger \sigma^j T^a \psi \mathcal{P}_{c\bar{c}', c\bar{c}} \psi^\dagger \sigma^i T^a \chi \rangle = 4m_c^2 \eta'^\dagger \sigma^j T^a \xi' \xi^\dagger \sigma^i T^a \eta$$

$$\langle \chi^\dagger \psi \mathcal{P}_{c\bar{c}', c\bar{c}} \psi^\dagger \chi \rangle = 4m_c^2 \eta'^\dagger \xi' \xi^\dagger \eta,$$

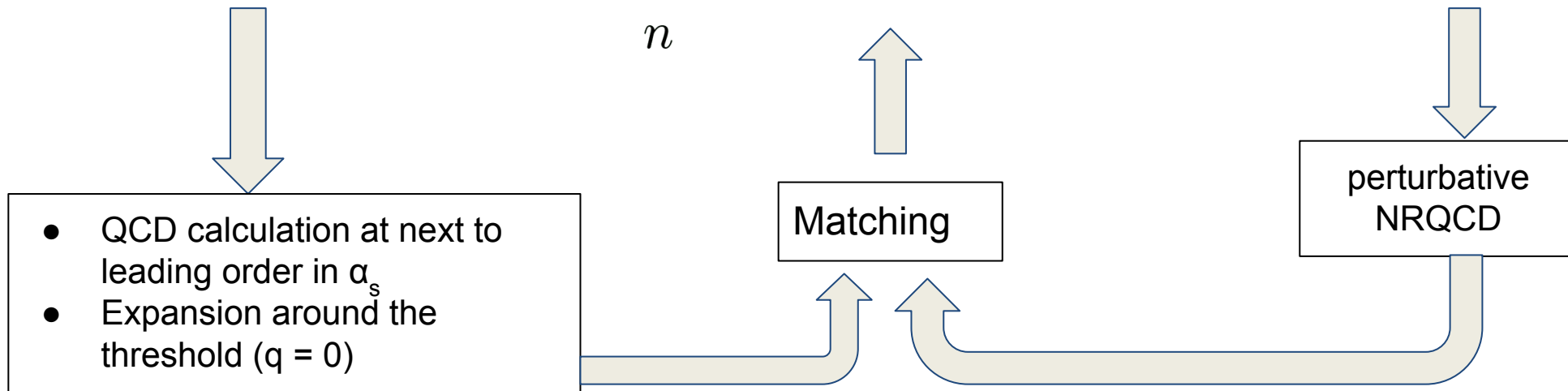
$$\langle \chi^\dagger T^a \psi \mathcal{P}_{c\bar{c}', c\bar{c}} \psi^\dagger T^a \chi \rangle = 4m_c^2 \eta'^\dagger T^a \xi' \xi^\dagger T^a \eta,$$

$$\langle \chi^\dagger (-\frac{i}{2} \overleftrightarrow{D}^m) T^a \psi \mathcal{P}_{c\bar{c}', c\bar{c}} \psi^\dagger (-\frac{i}{2} \overleftrightarrow{D}^n) T^a \chi \rangle = 4m_c^2 q'^m q^n \eta'^\dagger T^a \xi' \xi^\dagger T^a \eta,$$

$$\langle \chi^\dagger (-\frac{i}{2} \overleftrightarrow{D}^m) \sigma^i T^a \psi \mathcal{P}_{c\bar{c}', c\bar{c}} \psi^\dagger (-\frac{i}{2} \overleftrightarrow{D}^n) \sigma^j T^a \chi \rangle = 4m_c^2 q'^m q^n \eta'^\dagger \sigma^i T^a \xi' \xi^\dagger \sigma^j T^a \eta$$

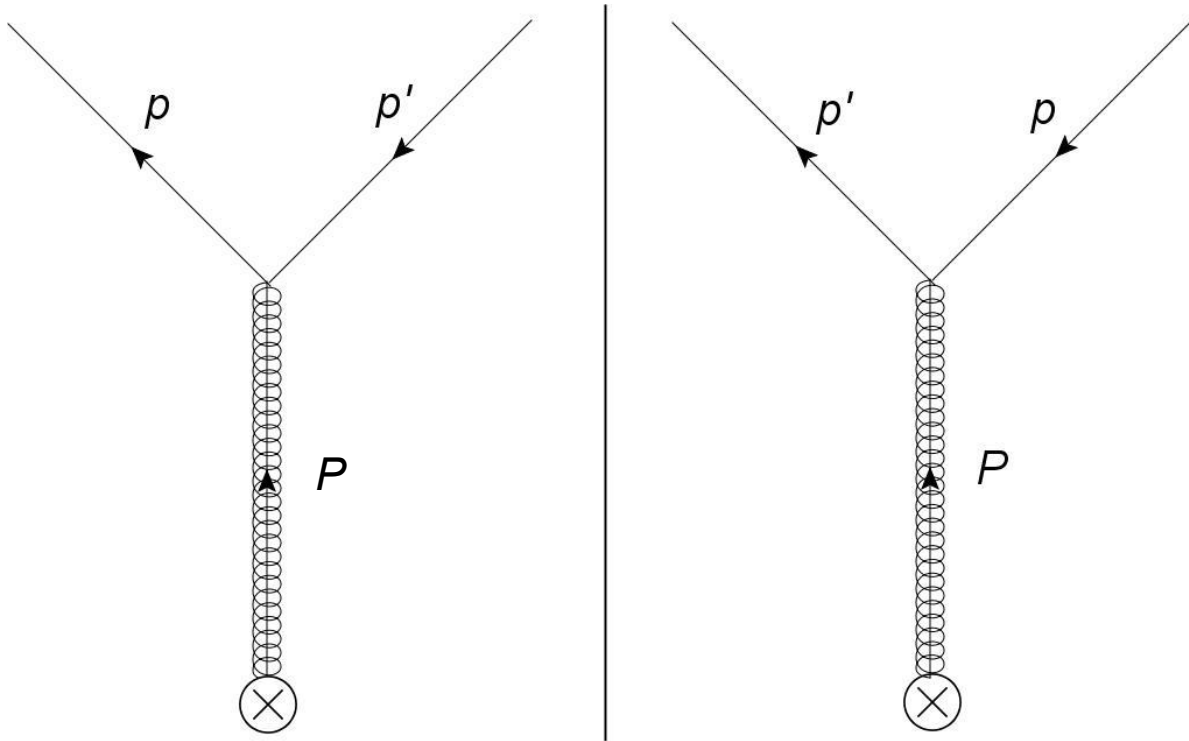
TMDFF calculation

$$D_{g \rightarrow J/\psi}(z, \mathbf{b}_T) = \sum_n d_{g \rightarrow Q\bar{Q}(n)}(z, \mathbf{b}_T) \langle 0 | \mathcal{O}^{J/\psi}(n) | 0 \rangle$$



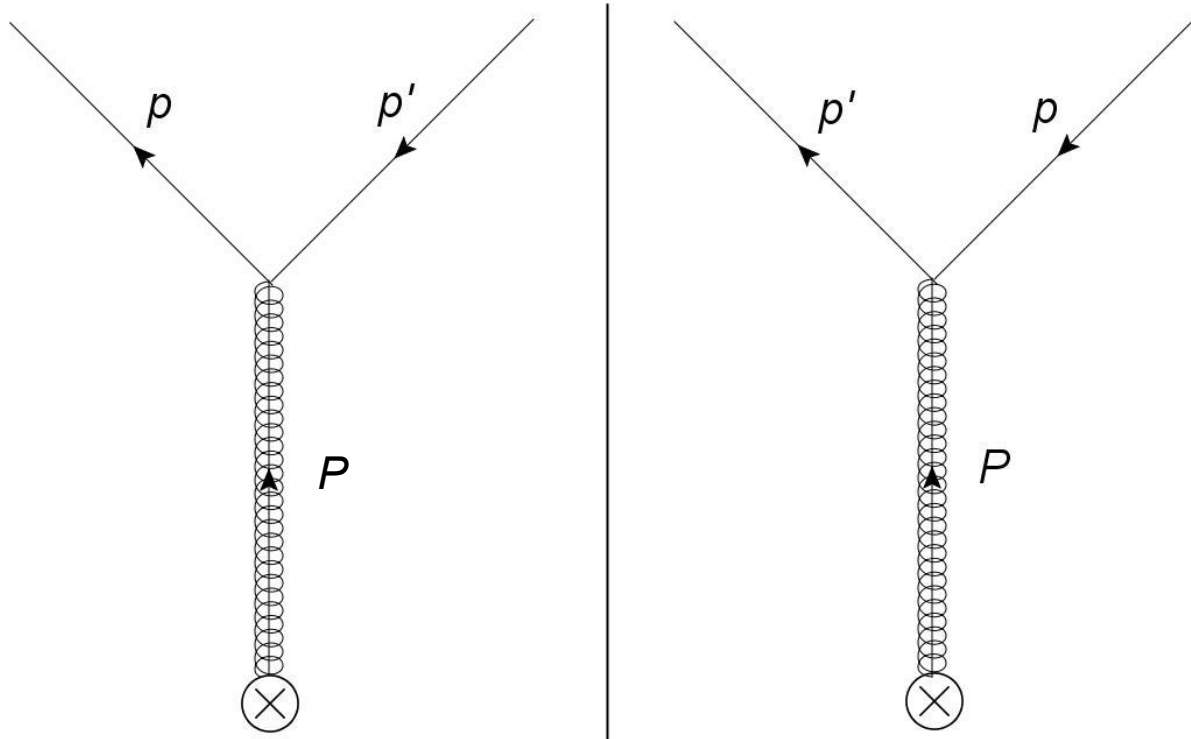
Feynman Diagrams

Leading order

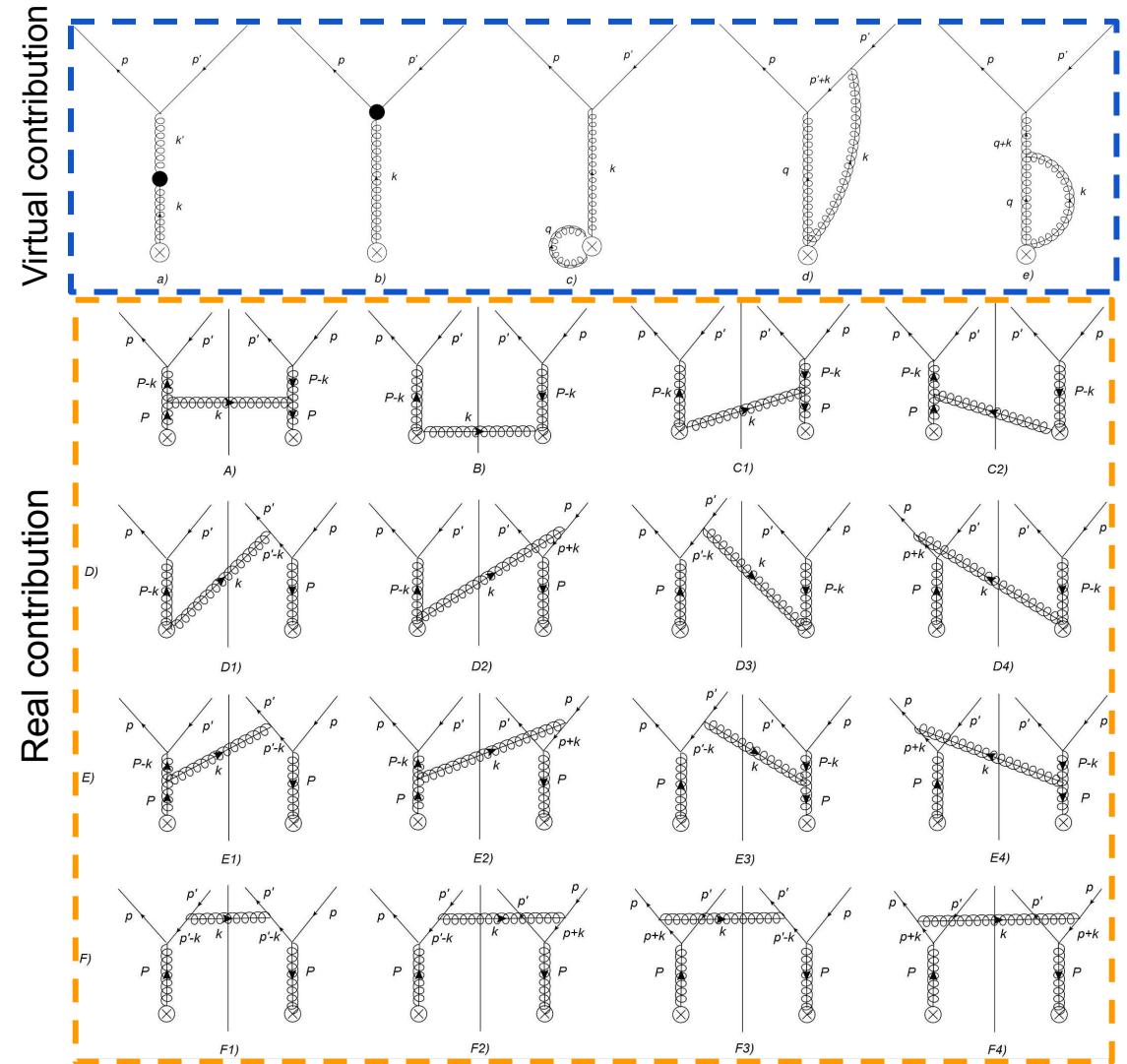


Feynman Diagrams

Leading order



Next to leading order



Some results

Virtual SDC

$$d^{a,b,c,d,e}(z, \mathbf{b}_\perp) = d^{LO}(z, \mathbf{b}_\perp) \frac{\alpha_s C_A}{2\pi} \left[\frac{1}{\epsilon_{UV}} \left(\frac{\beta_0}{2C_A} + \ln \frac{\delta^{+2}}{P^{+2}} \right) - \frac{1}{\epsilon_{IR}} \right.$$

$$+ 2 \ln^2 \frac{\delta^+}{P^+} + 2 \ln \frac{\delta^+}{P^+} \ln \frac{\mu^2}{4m_c^2} + \ln \frac{\mu^2}{m_c^2}$$

$$\left. - 4 \ln^2 2 + \frac{10}{3} \ln 2 + 2 - \pi^2 + \frac{123 - 10n_f}{27} \right]$$

Virtual SF

$$S_v = \frac{\alpha_s C_A}{2\pi} \left[\frac{-2}{\epsilon_{UV}^2} + \frac{2}{\epsilon_{UV}} \ln \frac{\delta^{+2} \zeta}{(P^+)^2 \mu^2} - \ln^2 \frac{(\delta^+)^2}{\mu^2} - \frac{\pi^2}{2} \right]$$

[arXiv:1502.05354](https://arxiv.org/abs/1502.05354)
[hep-ph]

$$\frac{R_g d^{a,b,c,d,e}(z, \mathbf{b}_\perp)}{d^{LO}(z, \mathbf{b}_\perp)} = \frac{\alpha_s C_A}{2\pi} \left[\frac{1}{\epsilon_{UV}^2} + \frac{1}{\epsilon_{UV}} \left(\frac{\beta_0}{2C_A} + \ln \frac{\mu^2}{\zeta} \right) - \frac{1}{\epsilon_{IR}} \right.$$

$$+ \frac{1}{2} \ln^2 \frac{\delta^{+2}}{\mu^2} + 2 \ln^2 \frac{\delta^+}{P^+} + 2 \ln \frac{\delta^+}{P^+} \ln \frac{\mu^2}{4m_c^2} + \ln \frac{\mu^2}{m_c^2}$$

$$\left. - \frac{3\pi^2}{4} - 4 \ln^2 2 + \frac{10}{3} \ln 2 + 2 + \frac{123 - 10n_f}{27} \right]$$

Real contribution!

Current state of the research

Two papers in progress...

Gluon TMDFF for
 J/ψ production

In collaboration with Ignazio Scimemi (Complutense University of Madrid - IPARCOS) and Miguel G. Echevarría (EHU)

J/ψ TMD shape function

In collaboration with Pieter Tael (University of Antwerp) and Miguel G. Echevarría (EHU)

Papers finished
before August!

...or so I hope