

Overview of SIDIS Experiments in Halls C



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- Semi-Inclusive Deep Inelastic Scattering (SIDIS)
- Past, Present and Future of Hall C SIDIS
- Opinions

Preliminaries

- Thanks to the many people for sharing their slides from previous talks.
- My apologies in advance for any mis-statements or misrepresentations I may make; they are my mistakes, not yours!

Do parton distributions and fragmentation functions factorize at Jefferson Lab energies?

Flavor Decomposition of SIDIS

$$\frac{1}{\sigma_{(e,e')}} \frac{d\sigma}{dz} (ep \rightarrow hX) = \frac{\sum_{q} e_q^2 f_q(x) D_q^h(z)}{\sum_{q} e_q^2(x) f_q(x)}$$

 $f_q(x)$: parton distribution function $D_q^h(z)$: fragmentation function

- Leading-Order (LO) QCD
- after integration over $p_{h\perp}$ and φ_h
- NLO: gluon radiation mixes x and z dependences
- Target-Mass corrections at large z
- In(1-z) corrections at large z



$$M_x^2 = W'^2 \sim M^2 + Q^2 (1/x - 1)(1 - z)$$

With p_T and k_T dependences, some kind of convolution is necessary to obtain final $P_{h\perp}$

SIDIS in Hall C

Past

• E00-108: 6 GeV beam with HMS and SOS spectrometers $\pi^{+/-}$ (H,D)

Present

- E12-09-002 (CSV): 12 GeV with HMS and SHMS $\pi^{+/-}$ (D)
- E12-09-017 (pT-SIDIS) *π*^{+/- ,} K^{+/-} (H,D)

Future!

- E12-13-007 (*π*⁰-SIDIS) HMS and NPS + PR12-23-014 (R) (H,D)
- E12-06-104 (R-SIDIS) *π*^{+/-} (H,D)
- C12-15-006 Tagged DIS!

Past

Factorization Test: E00-108



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JLab E00-108 Results

T. Navasardyan et al., PRL 98 022001 (2007)

- Cross section/simulation based on factorization prediction
- Good Agreement at low z
- Delta Resonance at high z



T. Navasardyan et al., PRL 98 022001 (2007)

dơ/dΩ_edE_edzdP_t²d≬ (nb/GeV³/sr)

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JLab Measurement: E00-108



- E= 5.5, x=0.3, Q²=2.3
- Similar, but different slopes for H, D
- Using simple gaussian+Cahn model, combined data yields momentum widths of pdf and fragmentation functions

from Phys. Lett. B665 (2008) 20

Present

Charge Symmetry Violation (CSV) Experiment E12-09-002

Introduction

What is Charge symmetry?

Charge symmetry (CS) is a specific rotation in isospin space. It is the invariance with respect to rotation of π about the T2 axis. $\begin{bmatrix} H, P_{CS} \end{bmatrix} = 0 \qquad \qquad P_{CS} |d\rangle = |u\rangle$ $P_{CS} = \exp(i\pi T2) \qquad \qquad P_{CS} |u\rangle = -|d\rangle$

Low Energy: CS in nuclei

CS operator interchanges neutrons and protons

- pp and nn scattering lengths are nearly the same
- $M_n \simeq M_p$
- $B(n, {}^{3}He) \simeq B(p, {}^{3}H)$ and energy levels in other mirror nuclei are equal (to 1%)
- $m(^{3}He) \simeq m(^{3}H)$

After electromagnetic corrections CS respected down to \sim 1%

QCD: Quark level

- $u^p(x, Q^2) = d^n(x, Q^2)$ $d^p(x, Q^2) = u^n(x, Q^2)$
- Origin of CS violations:
 → Electromagnetic interaction

$$\rightarrow \delta m = m_d - m_u$$

Naively, one would expect CSV would be on the order of $(m_d - m_u)/\langle M \rangle$, where $\langle M \rangle$ is roughly $0.5 - 1.0 \text{ GeV} \rightarrow \text{CSV}$ effect about 1%



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Motivation

- Charge symmetry violation is an important ingredient for pushing the precision frontier in the partonic structure of the nucleon
- Charge symmetry is often assumed in extracting PDFs from data where the data is limited in sensitivity to CS violation
- The validity of charge symmetry is a necessary condition for many relations between structure functions and sum rules
- Flavor symmetry violation extraction $\bar{u}(x) \neq \bar{d}(x)$ relies on the implicit assumption of charge symmetry (in the sea quarks)
- Charge symmetry violation viable part of explanation for the anomalous value of the Weinberg angle extracted by NuTeV experiment
- CSV is related to our understanding of the flavor dependence of the quark masses (one of the key unsolved problems in Physics why is $m_d \sim m_u \neq m_s \neq m_c \neq m_b \neq m_t$)

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SIDIS Formalism

Charge Symmetry Violation

In the PDFs: $\delta d(x) = d^{p}(x) - u^{n}(x), \delta u(x) = u^{p}(x) - d^{n}(x).$ $CSV(x) = \delta d - \delta u$

In Fragmentation Functions

$$\delta D(z) = \frac{D_u^{\pi^+} - D_d^{\pi^-}}{D_u^{\pi^+}}$$

Leading order methodology for iso-scaler targets (Londergan, Pang, and Thomas PRD54(1996)3154)

$$R_{meas}^{D}(x,z) = \frac{4N^{D\pi^{-}}(x,z) - N^{D\pi^{+}}(x,z)}{N^{D\pi^{+}}(x,z) - N^{D\pi^{-}}(x,z)} = \frac{4R_{Y}(x,z) - 1}{1 - R_{Y}(x,z)}$$
(1)

where $N^{D\pi^{\pm}}(x,z)$ is the **measured yield** of π^{\pm} electroproduction on a deuterium target, $R_Y = N^{D\pi^-} / N^{D\pi^+}$ is the yield ratio, and we rely the following:

Factorization $N^{\mathrm{N}h} \propto \sum e_i^2 q_i^{\mathrm{N}}(x) D_i^h(z)$

Impulse Approximation

$$N^{D\pi^{\pm}}(x,z) = N^{p\pi^{\pm}}(x,z) + N^{n\pi^{\pm}}(x,z)$$

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CSV in the Valence Region

Leading order experimental analysis \rightarrow will need higher order global analysis

Londergan, Pang and Thomas PRD54(1996)3154

$$D(z) \mathbf{R}(x, z) + A(x)\mathbf{CSV}(x) + F(z)\delta D(z) = \mathbf{B}(x, z)$$

$$D(z) = \frac{1 - \Delta(z)}{1 + \Delta(z)}, \Delta(z) = \frac{D_u^{\pi^-}(z)}{D_u^{\pi^+}(z)}$$

$$CSV(x) = \delta d - \delta u$$

$$R(x, z) = \frac{5}{2} + R_{meas}^D$$

$$A(x) = \frac{-4}{3(u_v + d_v)}$$

$$F(z) = \frac{4 + \Delta(z)}{3(1 - \Delta^2(z))}$$

$$A(x) \text{ and } B(x, z) \text{ are computed from PDF and FF fits}$$

$$B(x, z) = \frac{5}{2} + R_{sea_s}^D(x, z) + R_{sea_s}^D(x, z) + R_{sea_s}^D(x)$$

$$R_{sea_s}^D(x, z) = \frac{\Delta_s(z)[s(x) + \overline{s}(x)]/(1 + \Delta(z))}{[u_v^p(x) + d_v^p(x)]}$$

$$\Delta_s(z) = \frac{D_s^-(z) + D_s^+(z)}{D_u^+(z)}$$



Extract simultaneously D(z) and CSV(x) from each (Q^2,x) setting



Experiment in Hall C – E12-09-002

Measurements: $D(e, e'\pi^+)$ and $D(e, e'\pi^-)$

Setup

- 11 GeV e^- beam
- 10 cm LD₂ target
- SHMS $\rightarrow \pi^{\pm}$, HMS $\rightarrow e'$

Each x setting has 4 z measurements $z_j = 0.4, 0.5, 0.6, 0.7$ $R_Y(x, z) = Y^{D\pi^-}(x, z)/Y^{D\pi}(x, z)$ $\boxed{R_{Meas}^D(x, z) = \frac{4R_Y(x, z) - 1}{1 - R_Y(x, y)}}$ $Q^2 = 3.5 \text{ GeV}^2 \rightarrow x = 0.30, 0.35, 0.40, 0.45$ $Q^2 = 5.1 \text{ GeV}^2 \rightarrow x = 0.45, 0.50, 0.55, 0.60$



Assuming $\delta D(z) = 0$, for each Q^2 we have 16 equations and 8 unknowns: $D(z_j)$ and $CSV(x_i)$ D(z) R(x, z) + A(x)CSV(x) = B(x, z)



 R_{meas}^D Results

$$D(z) R(x,z) + A(x)CSV(x) = B(x,z)$$
$$R(x,z) = \frac{5}{2} + R^{D}_{meas}(x,z)$$



$$R_{meas}^{D}(x, z) \text{ for } Q^{2} = 4 \text{ GeV}^{2}$$

bin projected on z axis.
$$R_{meas}^{D}(x, z) = \frac{4R_{Y}(x, z) - 1}{1 - R_{Y}(x, z)}$$

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Factorization

Charge Ratio Sum and Differences

$\sigma_p^{\pi^+} - \sigma_p^{\pi^-} - 4u_v(x) - d_v(x) - p^-$	$\frac{d_v}{d_v} = \frac{4 - 3R^-}{2R^-}$
$\frac{1}{\sigma_d^{\pi^+} - \sigma_d^{\pi^-}} = \frac{1}{3(u_v() + d_v(x))} = R$	$u_v = 3R^- + 1$



Ratios should not depend on z.



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Simultaneous Extraction of $\Delta(z)$ and CSV(x)

Fragmentation ratio and valence CSV parton distribution

Four parameter fit

$$\Delta(z) \equiv \frac{D_u^{\pi^-}(z)}{D_u^{\pi^+}(z)} = z^{\alpha}(1-z)^{\beta}$$
$$CSVx \equiv \delta d - \delta u = x^a(1-x)^b(x-c)$$
$$c \text{ is determined from the constraint: } \int_0^1 CSV(x)dx = 0$$

$$c = \frac{\int_0^1 x^{(a+1)} (1-x)^b}{\int_0^1 x^a (1-x)^b} = \frac{B(a+2,b+1)}{B(a+1,b+1)}, B(x,y) = \frac{\Gamma(x)\Gamma(y)}{\Gamma(x+y)}$$
$$R_{fit}^D(x,z) = \frac{B(x,z) - A(x)CSV(x)}{D(z)} - \frac{5}{2}$$

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Results after standard ρ background subtraction



The data points on the right are plotted using the extracted Δ , $R^{D}_{meas}(x,z)$, and the equation below



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Jefferson Lab Exp. E12-09-017: Precise measurements of (e,e' π^{\pm}) and (e,e'K[±]) cross sections at Semi-Inclusive Deep Inelastic Scattering (SIDIS) Kinematics

- Precise measurements to test the assumptions of factorization of SIDIS process at photon invariant momentum transfer $Q^2=q^2-\nu^2$, at moderate Bjorken $x = Q^2/2M\nu$ (*M* is proton mass)
- Allow exploration of assumptions of favored/disfavored fragmentation of different flavor quarks using ¹H and ²H targets
- Investigate possible target mass effects
- Investigate possible higher twist effects
- Complement SIDIS measurements in large open acceptance detector CLAS at Jefferson Lab Hall B

SIDIS cross section with Transverse Momentum-dependent Parton and Fragmentation Distributions (TMDs)

 $f_q(x,{f k}_ot)$: parton distribution function as function of intrinsic parton $k_ au$ $D^h_q(z,{f p}_ot)$: fragmentation function as a function of fragmentation $p_ au$

Cross section for SIDIS hadron of fractional energy z_h and transverse momentum P_T

$$\frac{d^5 \sigma^{\ell p \to \ell h X}}{dx_B \, dQ^2 \, dz_h \, d^2 \mathbf{P}_T} = \sum_q e_q^2 \int d^2 \mathbf{k}_\perp \, f_q(x, \mathbf{k}_\perp) \, \frac{2\pi \alpha^2}{x_B^2 s^2} \, \frac{\hat{s}^2 + \hat{u}^2}{Q^4}$$
$$\times D_q^h(z, \mathbf{p}_\perp) \, \frac{z}{z_h} \, \frac{x_B}{x} \left(1 + \frac{x_B^2}{x^2} \frac{k_\perp^2}{Q^2}\right)^{-1}$$

from Anselmino et al. (hep-ph/0412316v1)

Model for basic data interpretation

• Now perform k_{\perp} integration and keep terms order O(k \perp/Q) on previous cross section expression to get

$$\begin{aligned} \frac{d^5 \sigma^{\ell p \to \ell h X}}{dx_B \, dQ^2 \, dz_h \, d^2 \mathbf{P}_T} \simeq &\sum_q \frac{2\pi \alpha^2 e_q^2}{Q^4} \, f_q(x_B) \, D_q^h(z_h) \bigg[(1 + (1 - y)^2) \\ &-4 \, \frac{(2 - y)\sqrt{1 - y} \, \langle k_\perp^2 \rangle \, z_h}{\langle P_T^2 \rangle \, Q} \, \bigg[\frac{1}{\pi \langle P_T^2 \rangle} \, e^{-P_T^2/\langle P_T^2 \rangle} \, , \end{aligned}$$

$$\begin{aligned} &\text{where } \langle P_T^2 \rangle = \langle p_T^2 \rangle + z_h^2 \langle k_\perp^2 \rangle \, .\end{aligned}$$

- At each (x, Q²,z) fit multiplicities with $y = M_0 b e^{-bP_T^2} (1 + AP_T cos(\phi))$
- Fit parameters are M_0 , A, and $b = 1/\langle P_T^2 \rangle$
- Results very preliminary

Example of fits of parameterization to multiplicities

 π + x=0.31 Q2 = 3.1 z = 0.33 N=151 χ^2 /dof = 1.9



Fit Results for $y = M_0 b e^{-bP_T^2} (1 + AP_T cos(\phi))$ Parameters



Curves are DSS frag. functions with CTEQ5 pdfs. The solid circles are with DVM ρ subtraction, while the crosses are with no ρ subtraction. The lefthand arrow shows the z value for (W')**2=3 GeV, while the righthand arrow is for (W')**2 = 2.5 GeV**2.

Fit Results for $y = M_0 b e^{-bP_T^2} (1 + AP_T cos(\phi))$ Parameters



The curve is of the form $\langle P_T^2 \rangle = \langle p_t^2 \rangle$ + $z^2 \langle k_{\perp}^2 \rangle$ where $\langle p_t^2 \rangle$ is the width from fragmentation (0.2 GeV² here) for both favored and unfavored FF), and $\langle k_{\perp}^2 \rangle$ is the intrinsic q transverse momentum width (also 0.2 GeV²) for both u and d quarks.

Fit Results for $y = M_0 b e^{-bP_T^2} (1 + AP_T cos(\phi))$ Parameters



<u>Generally</u> the results increase with increasing z, and are significantly greater than 0 for π -, but closer to zero for π +. The Cahn effect would predict negative values of A, in contradiction with the data.

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Future

Experiment E12-13-007: π^0 SIDIS



Transverse momentum widths of quarks with different flavor (and polarization) can be different



 $P_{T} = p_{t} + z k_{t} + O(k_{t}^{2}/Q^{2})$

E12-13-007 goal: Measure the basic SIDIS cross sections of π° production off the proton, including a map of the P_T dependence (P_T ~ Λ < 0.5 GeV), to validate^(*) flavor decomposition and the k_T dependence of (unpolarized) up and down quarks

(*) Can only be done using spectrometer setup capable of %-type measurements (an essential ingredient of the global SIDIS program!)

Requires new ~25 msr Neutral-Particle Spectrometer

Advantages of (e,e' π°) beyond (e,e' $\pi^{+/-}$)

- Many experimental and theoretical advantages to validate understanding of SIDIS with neutral pions
- **Can verify:** $\sigma^{\pi^{0}}(x,z) = \frac{1}{2} (\sigma^{\pi^{+}}(x,z) + \sigma^{\pi^{-}}(x,z))$
- **D** Confirms understanding of flavor decomposition/ k_{T} dependence

PAC: "the cross sections are such basic tests of the understanding of SIDIS at 11 GeV kinematics that they will play a critical role in establishing the entire SIDIS program of studying the partonic structure of the nucleon."

Experiment E12-13-007: π^0 SIDIS



 $P_{h\perp}$ Coverage of SIDIS experiments

(e,e'π[±]) with SHMS E12-09-017

(e,e'π⁰) with NPS E12-13-007



PR12-23-014: SIDIS basic (e,e' π^0) cross sections





New Run Group addition proposal to PAC51 (P. Bosted, E. Kinney, H. Mkrtchyan, V. Tadevosyan, R. Ent, T. Horn, et al.)

Measure $R_{LT} = \sigma_L / \sigma_T$, the ratios of d/u cross sections, the transverse momentum dependence of the cross section, and the spin-independent and beam-spin-dependent modulations of the cross section

Projections for R_{LT} SIDIS as function of p_T and z



Physics goals are driven by the need to more fully understand the production processes that enter SIDIS for better understanding of the 3D nucleon structure

• Dynamic and target higher twist, deep-exclusive processes, VM, CSV ¹³

Measurement of the ratio $R = \sigma_L / \sigma_T$ in Semi-Inclusive Deep Inelastic Scattering

E12-06-104, Spokespersons: P. Bosted, R. Ent, E. Kinney, and H. Mkrtchyan

- This experiment will make precise measurements of R in charged π and K SIDIS on H and D targets as a function of Q^2 , fractional hadron momentum z, and hadron transverse momentum p_T
- Standard technique to measure R: Vary the virtual photon polarization ε by using different incident beam energies and electron scattering angles, while keeping the Q^2 , x, z, and p_T constant. Will use the two magnetic spectrometers in Hall C.

 $\varepsilon = \left[1 + 2\left(\frac{Q^2}{4M^2x^2}\right)\tan\frac{\theta^2}{2}\right]^{-1} \qquad \sigma = \Gamma(\sigma_{\rm T} + \varepsilon\sigma_{\rm L} + \varepsilon\cos(2\phi)\sigma_{\rm TT} + [\varepsilon(\varepsilon+1)/2]^{1/2}\cos(\phi)\sigma_{\rm LT})$

$R = \sigma_L / \sigma_T$ is a basic aspect of the photon-parton interaction



 Almost no experimental knowledge of R in SIDIS

Projections for E12-06-104 vs existing Cornell Data (projections assume $R_{SIDIS} = R_{DIS}$) Comparable 1.6% systematic uncertainties not indicated





Experiment E12-06-104: R SIDIS

- An essential measurement in understanding SIDIS in LO factorized form at these energies $\sim \sum_{q}^{\pi} \sim \sum_{q}^{2} q(x) D_{q}^{\pi}(z)$ quark
 - Previous JLab cross section experiments experiments suggest this factorized picture is valid at JLab energies at appropriate final hadronic state energies
- We will be able to test many common assumptions used in SIDIS analyses:

$$R_{SIDIS} = R_{DIS}?$$

$$R_{SIDIS}^{\pi^{+}} = R_{SIDIS}^{\pi^{-}}?$$

$$R_{SIDIS}^{H} = R_{SIDIS}^{D}?$$

$$R_{SIDIS}^{\pi^{+}} = R_{SIDIS}^{K^{+}}?$$

$$R_{SIDIS}^{K^{+}} = R_{SIDIS}^{K^{-}}?$$

- Important for determining spin structure function g_1^h (need term $(1 + \varepsilon R)$ to get g_1^h/F_1^h from A_{\parallel}^h)
- At low z, expect DIS Q² behavior ($\sim 1/Q^2$), but as z \rightarrow 1, expect Deep-Exclusive Q² behavior ($\sim Q^2$)
- Completely unknown p_T behavior, which might impact on TMD analyses

Hall C Kinematic Reach

HMS + SHMS (or NPS) Accessible Phase Space for SIDIS



The Tagged Deep Inelastic Scattering (TDIS) Experiment

Knowledge of meson structure is critical to a complete understanding of the emergence hadron mass.



Lack of meson targets \Rightarrow No direct measurement

TDIS is a direct measurement of the mesonic content of nucleons and extraction of the pion's F₂ structure functions, by deep inelastic scattering off the virtual-meson cloud.

The well established spectator tagging used to access the "meson cloud" target.

C1 conditionally approved with A- rating for 27 PAC days (up for jeopardy review @ PAC 51)

Two Additional Run-group Experiments Approved Run-group: TDIS, kaon-TDIS and neutron-TDIS

Many-fold increased interest in the technique and the science goal - over 50 publications with more than 1200 citations (e.g. LRP white paper & EIC yellow report).

Significant theory progress - global QCD analysis including leading neutron HERA data

Substantial experimental progress - mTPC based recoil detector, streaming readout, simulations validated with BoNUS12 data, high rate & high occupancy tracking algorithm



TDIS will be a pioneering experiment and a necessary first step for future experiments at the EIC and 22 GeV JLab.

Strong Hall C Capabilities

- Hall C is at the high luminosity frontier, precisely measuring small cross sections and performing longitudinal/transverse separations are our bread and butter
- The ability to rapidly switch between different targets can greatly reduce systematic uncertainties, and a wide choice of nuclei is possible
- Do not fail to appreciate the ability to set a spectrometer at a precise angle; alignment of large collider detectors takes years can take years (e.g. Fermilab W mass result)
- High luminosity measurements with polarized beams and targets are well established capabilities at Jefferson Lab
- Flexibility to change experimental layout (spectrometers, calorimeters, compact photon source) allows for optimized measurements with precision

Concluding Opinions

- The general capabilities of the Hall allow a broad SIDIS program already at 12 GeV; the focus is on precise measurements of small cross sections and the ability to perform longitudinal/transverse separations.
- The first 12 GeV experiments have produced new results!
- The high precision Hall C experiments will complement those of Hall B.
- As we explore SIDIS at 12 GeV, we will move beyond the nucleon as a collection of free point objects! These are the signatures of the emergent physical structure we wish to understand!!!

Backup Slides

Experiment E12-09-002

Kinematic Coverage

Charge Symmetry Violating Quark Distributions via Precise Measurement of π^+/π^- Ratios in Semi-inclusive Deep Inelastic Scattering.



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New Theoretical Guidance

Mapping the Kinematical Regimes of Semi-Inclusive Deep Inelastic Scattering, M. Boglione, A. Dotson, L. Gamberg, S. Gordon, J.O. Gonzalez-Hernandez, A. Prokudin, T.C. Rogers, and N. Sato, (2019), ArXiv:1904.12882



Example: Study of kinematic deviation between Nachtmann and Bjorken x

What are the tools available?



Neutral Pion Spectrometer



HMS and SHMS: Base Spectrometers

•HMS and SHMS that can reach 6+ and 11+ GeV/c with 10% momentum bites (single and coincidence mode)

0.1% reconstruction
 efficiency easily
 demonstrated

•Excellent PID

•Cerenkov and lead glass shower counters

•Can go up to 5.5 and 10.5 deg with sub-mrad pointing accuracy



Experimental Hall C

