

Structure of hypernuclei:

*What will happen by the coupling of a  $\Lambda$  particle?*

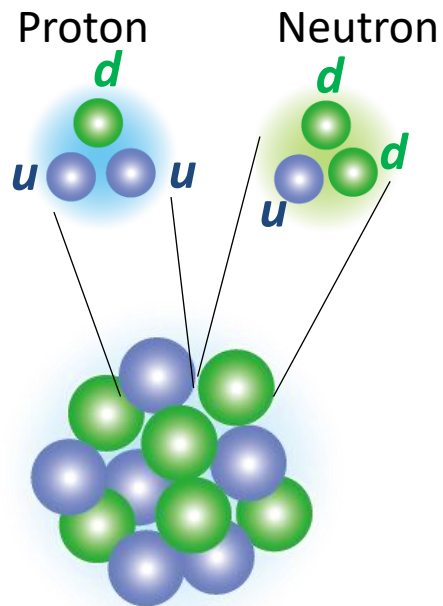
Masahiro Isaka, Hosei University

# Hypernuclei

## ◆ Normal nuclei

- Nucleons

- protons & neutrons

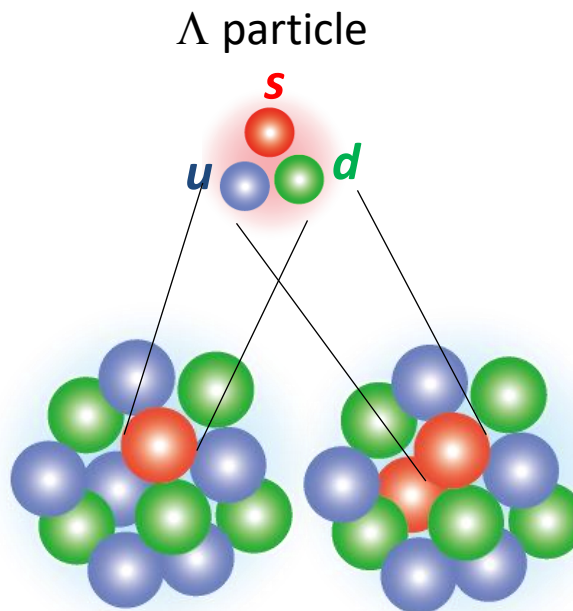


(Normal) nuclei

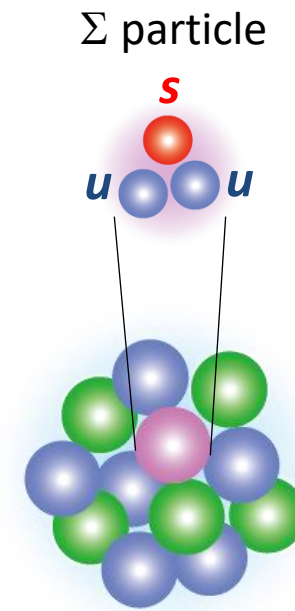
## ◆ Hypernuclei

- Nucleons and hyperon(s) ( $\Lambda$ ,  $\Sigma$ ,  $\Xi$ )
- Several kinds of hypernuclei exist

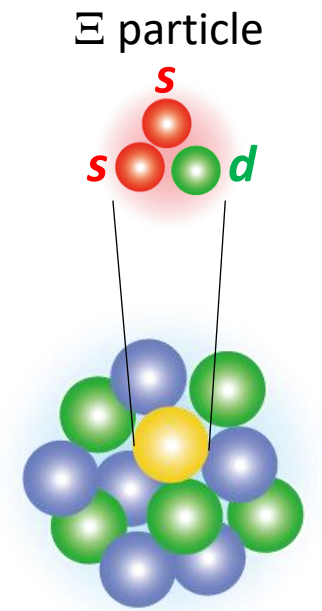
→ Today's talk: single  $\Lambda$  hypernuclei



$\Lambda$  hypernuclei  $\Lambda\Lambda$  hypernuclei



$\Sigma$  hypernuclei

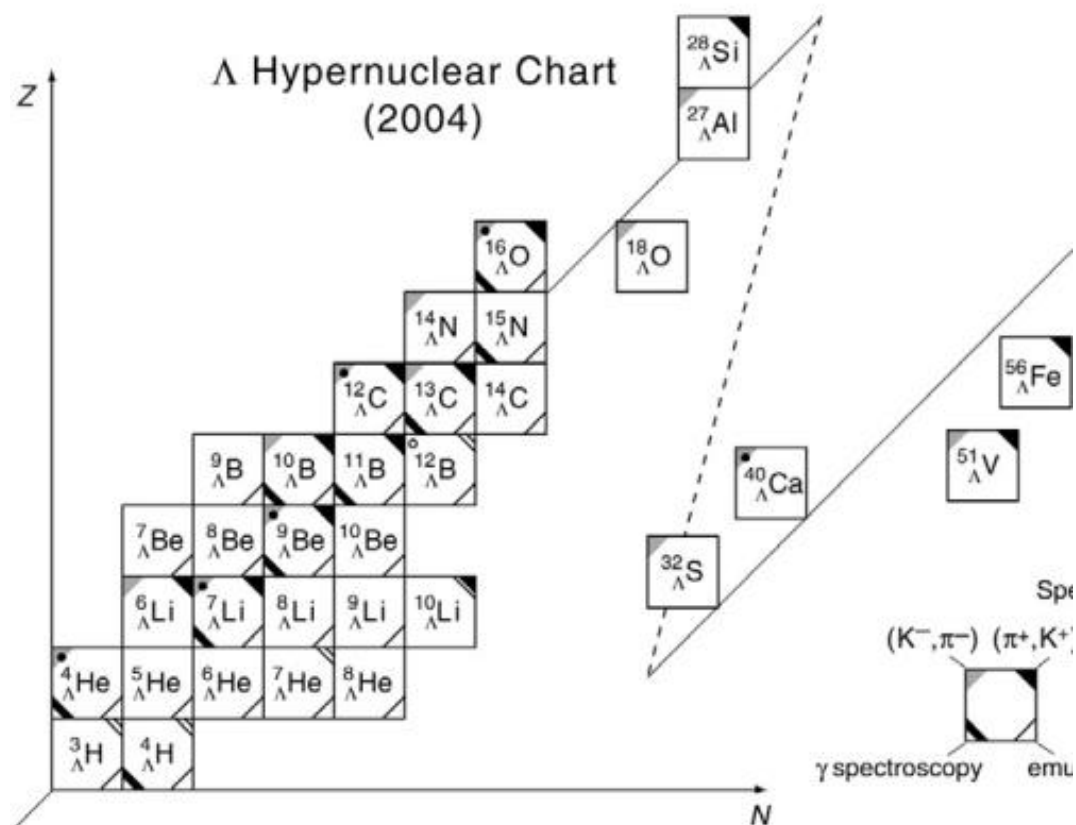


$\Xi$  hypernuclei

# Current situation of $\Lambda$ hypernuclear studies

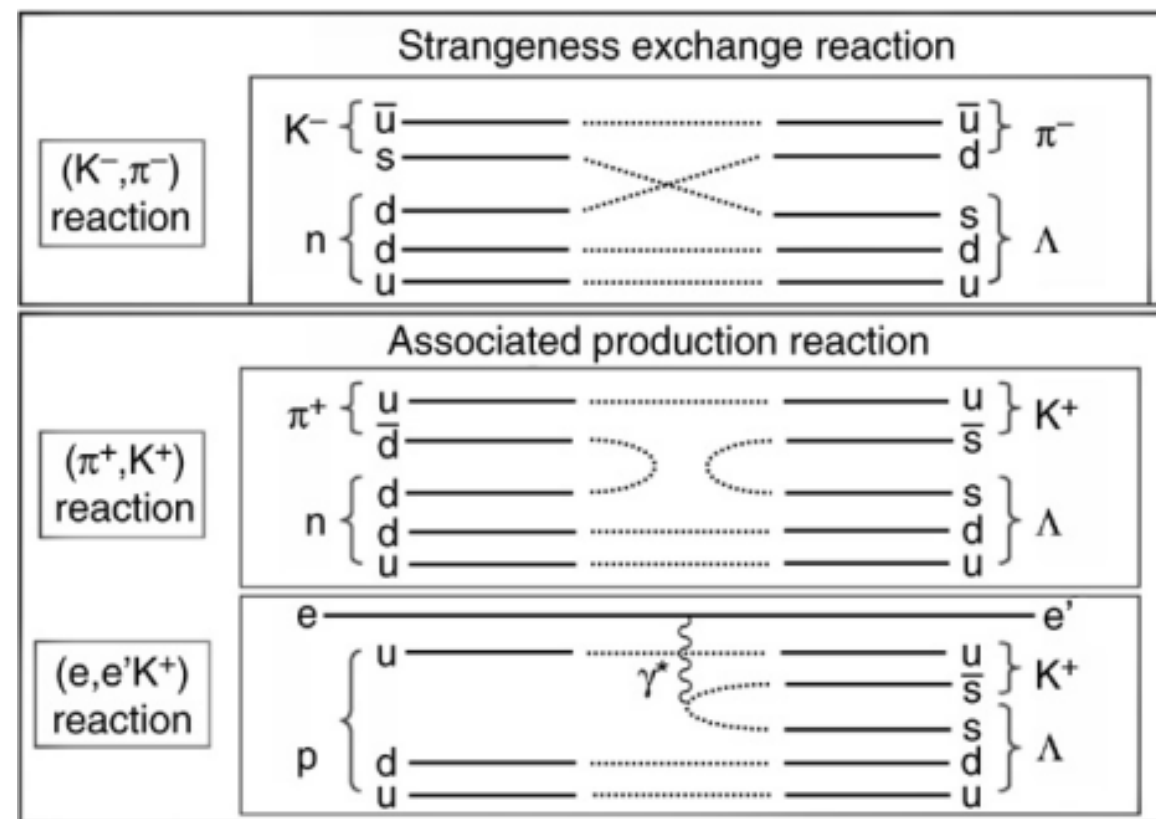
Hashimoto & Tamura, PNP **57** (2006), 564.

## Hypernuclei observed so far



Now,  $\Lambda$  hypernuclear chart is extended more!

## Production process of $\Lambda$ hypernuclei

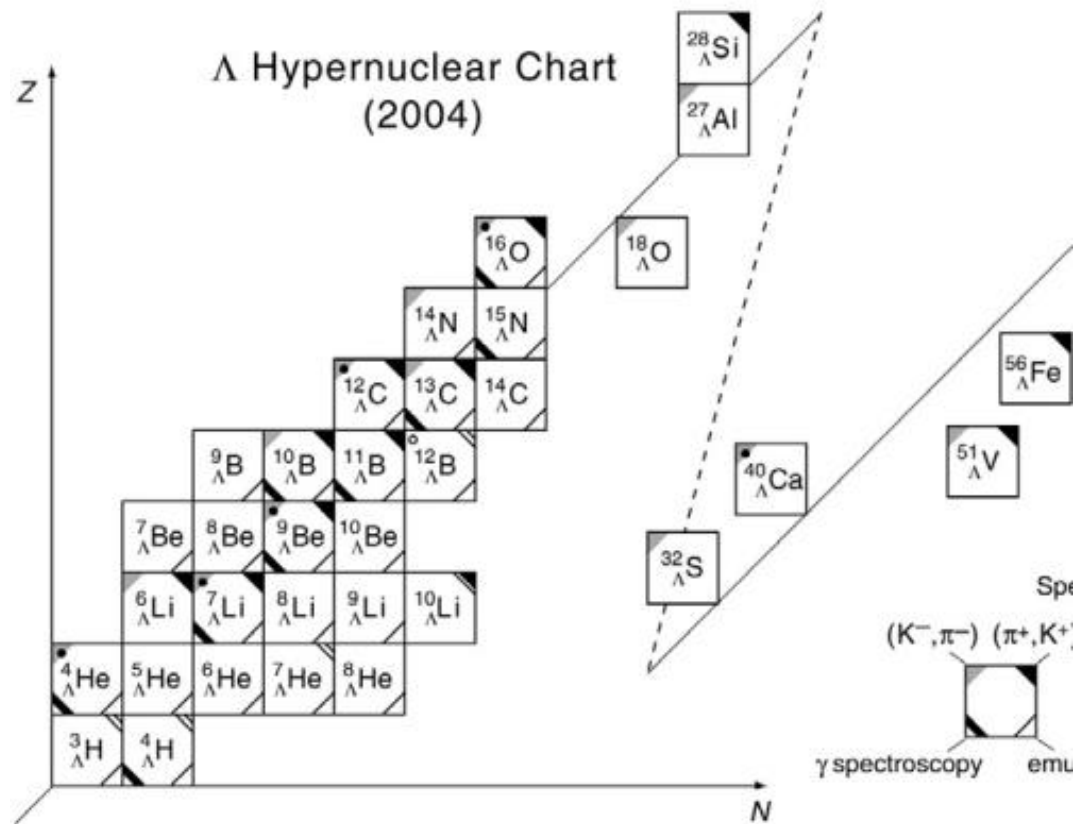


and  $\Lambda$  hypernuclear production by heavy ion collisions, etc.

# Current situation of $\Lambda$ hypernuclear studies

Hashimoto & Tamura, PNP **57** (2006), 564.

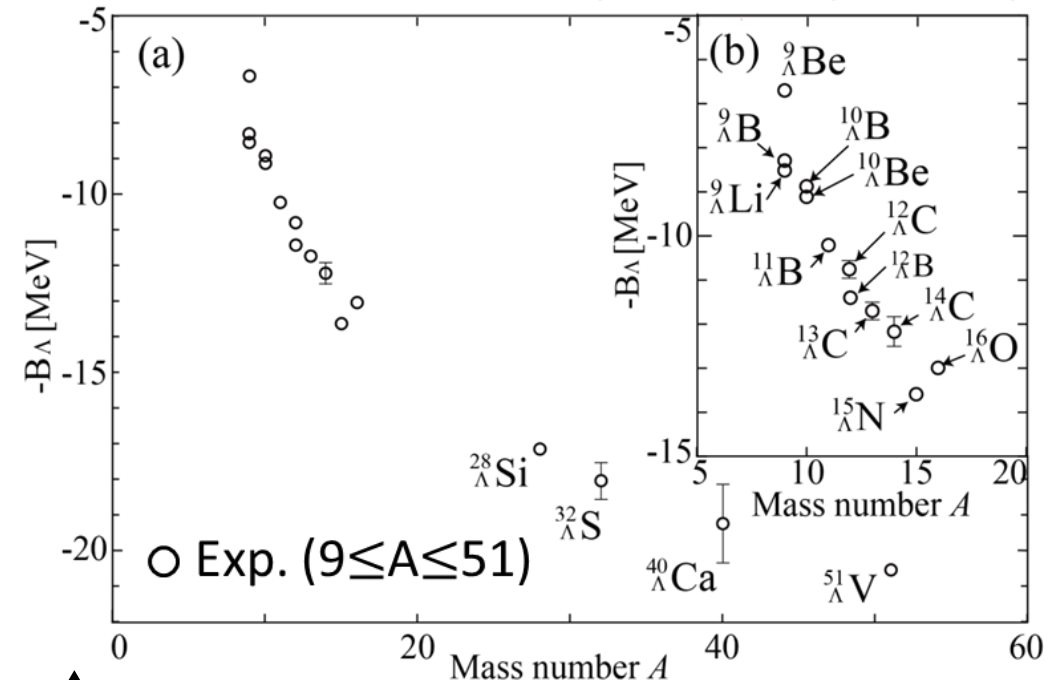
## Hypernuclei observed so far



Now,  $\Lambda$  hypernuclear chart is extended more!

## Observed data: $\Lambda$ binding energy $B_{\Lambda}$

A most basic and important quantity



$E$  ↑

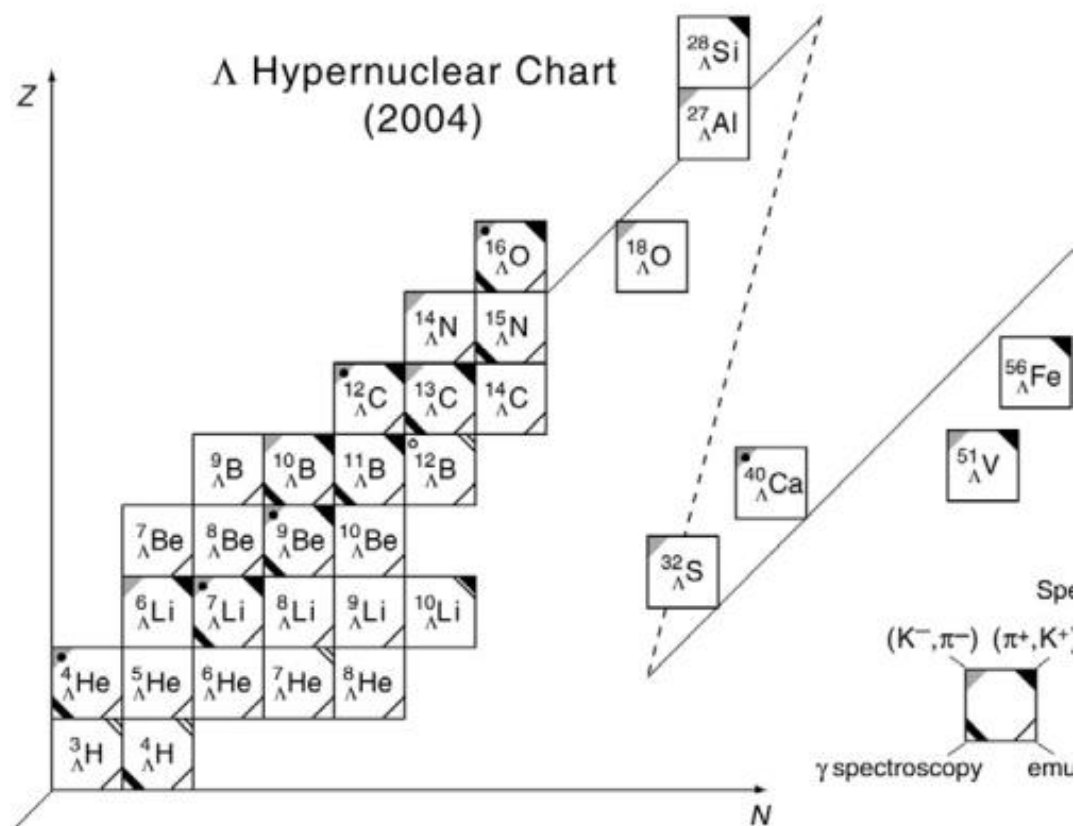
Energy of the core nuclei  ${}^A Z$

$B_{\Lambda}$ : energy gain by a  $\Lambda$  particle

E of corresponding hypernucleus  ${}^{A+1}_{\Lambda} Z$

# Current situation of $\Lambda$ hypernuclear studies

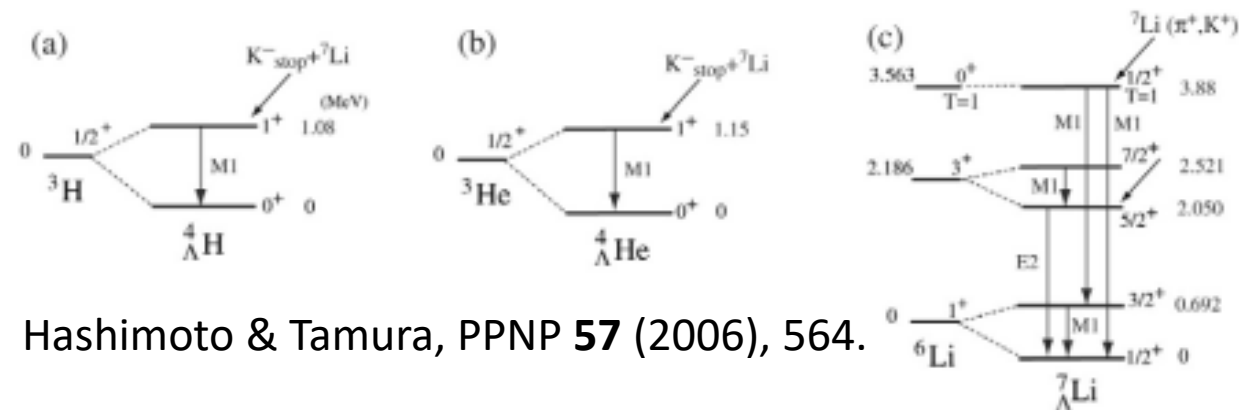
## Hypernuclei observed so far



Now,  $\Lambda$  hypernuclear chart is extended more!

## Observed data: Excited states

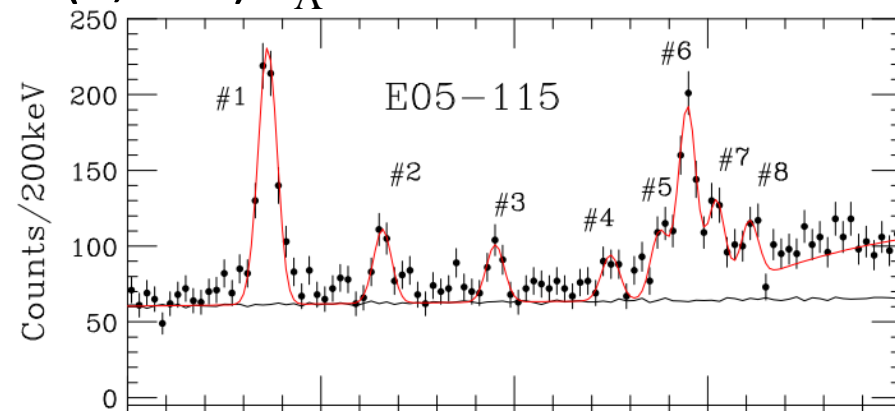
### $\gamma$ -ray spectroscopy data for bound states



Hashimoto & Tamura, PPNP **57** (2006), 564.

### $(e, e'K^+)$ reaction spectroscopy at JLab: not only bound but also unbound states

e.g.  ${}^{12}\text{C}(e, e'K^+){}^{12}_{\Lambda}\text{B}$  L. Tang, et al., PRC **90**, 034320(2014)



# Grand challenges of hypernuclear physics

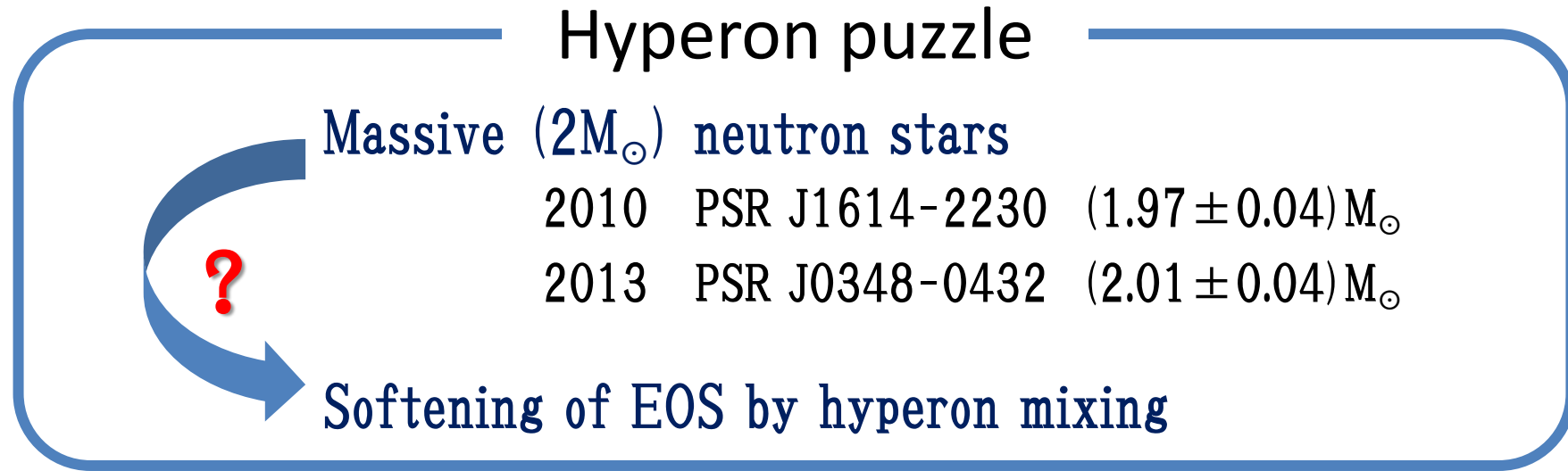
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## **Interaction:** To understand baryon-baryon interaction

- 2 body interaction between baryons (Y: hyperon, N: nucleon): YN, YY interactions
  - Studied through hypernuclei due to difficulty of YN scattering exp.
- Many-body interaction
  - Important issue in recent studies

## **Structure:** To understand many-body system of nucleons and hyperon

# Many-body force and “Hyperon puzzle” in neutron star



***How do we resolve?***

**Baryon many(three)-body force**

If strong repulsion exists acting in hyperonic channels,  
EOS of neutron star matter becomes stiff

**Important issue in hypernuclear physics:  
to reveal effects of hyperonic many-body force in  $\Lambda$  hypernuclei**

# Grand challenges of hypernuclear physics

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## **Interaction: To understand baryon-baryon interaction**

- 2 body interaction between baryons (Y: hyperon, N: nucleon): YN, YY interactions  
→ Studied through hypernuclei due to difficulty of YN scattering exp.
- Many-body interaction  
→ Important issue in recent studies

## **Structure: To understand many-body system of nucleons and hyperon**

- Addition of hyperon(s) shows new aspects of nuclear structure

e.g.) structure change by hyperon(s)

- No Pauli exclusion between N and Y
- YN interaction is different from nuclear force

} Hyperon as an impurity  
in hypernuclei



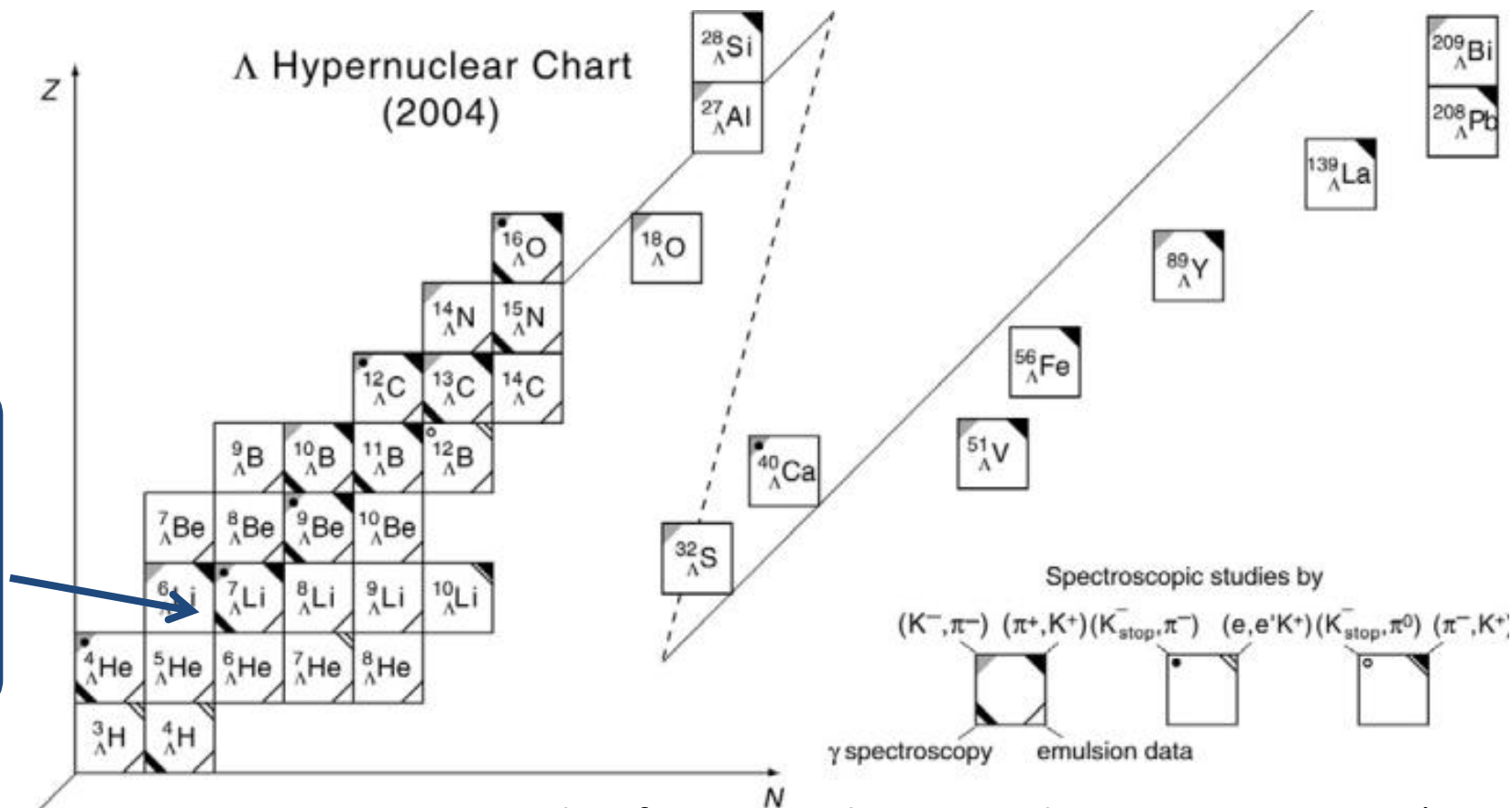
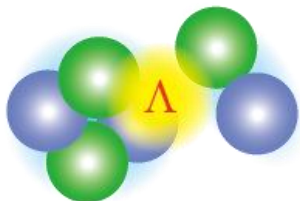
# Structure of $\Lambda$ hypernuclei

## ◆ $\Lambda$ hypernuclei observed so far

- Concentrated in light  $\Lambda$  hypernuclei
- Most have well-developed cluster structure

Light  $\Lambda$  hypernuclei

Developed cluster



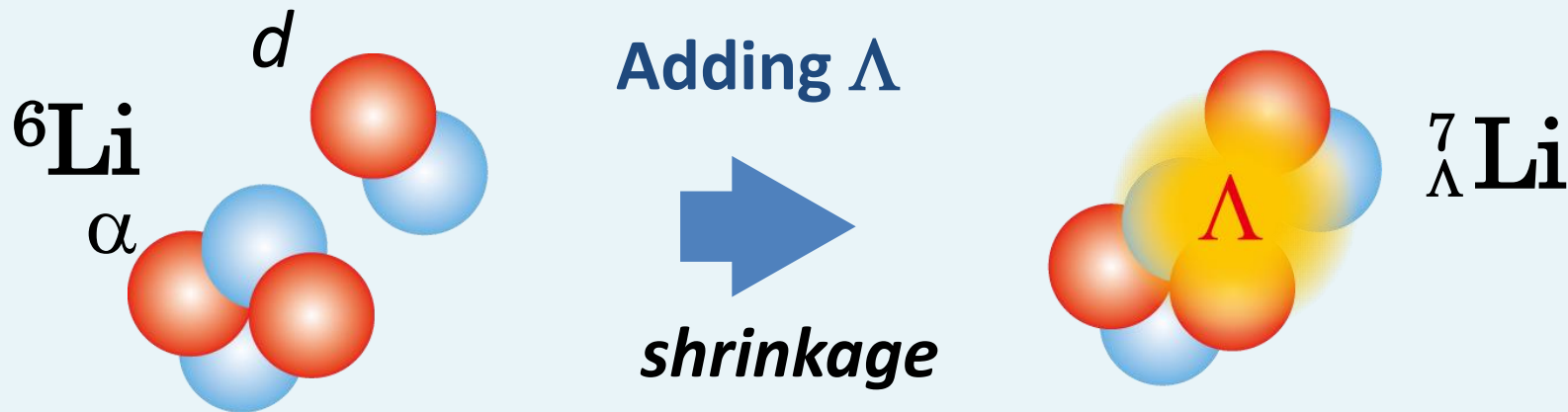
Taken from O. Hashimoto and H. Tamura, PPNP **57**(2006),564.

# Cluster structure in light hypernuclei

## ◆ Famous example of “impurity effects”

Example:  ${}^7_{\Lambda}\text{Li}$

Motoba, *et al.*, PTP**70**,189 (1983)  
Hiyama, *et al.*, PRC**59** (1999), 2351.  
Tanida, *et al.*, PRL**86** (2001), 1982.

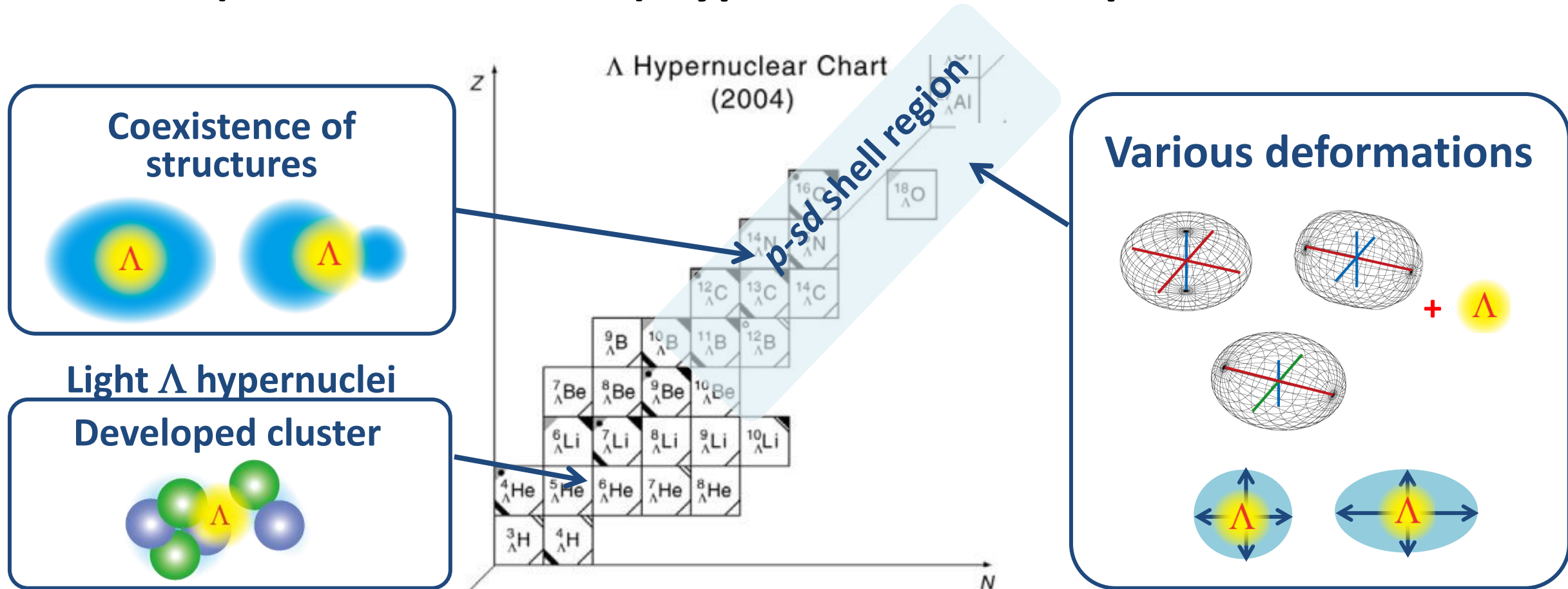


- $\Lambda$  reduces inter-cluster distance b/w  $\alpha + d$  of the core nucleus  ${}^6\text{Li}$
- Confirmed through  $B(E2)$  reduction

# Toward heavier and exotic $\Lambda$ hypernuclei

## ◆ Experiments at JLab, J-PARC, ... *etc.*

- Heavier(*sd*-shell and more) hypernuclei can be produced



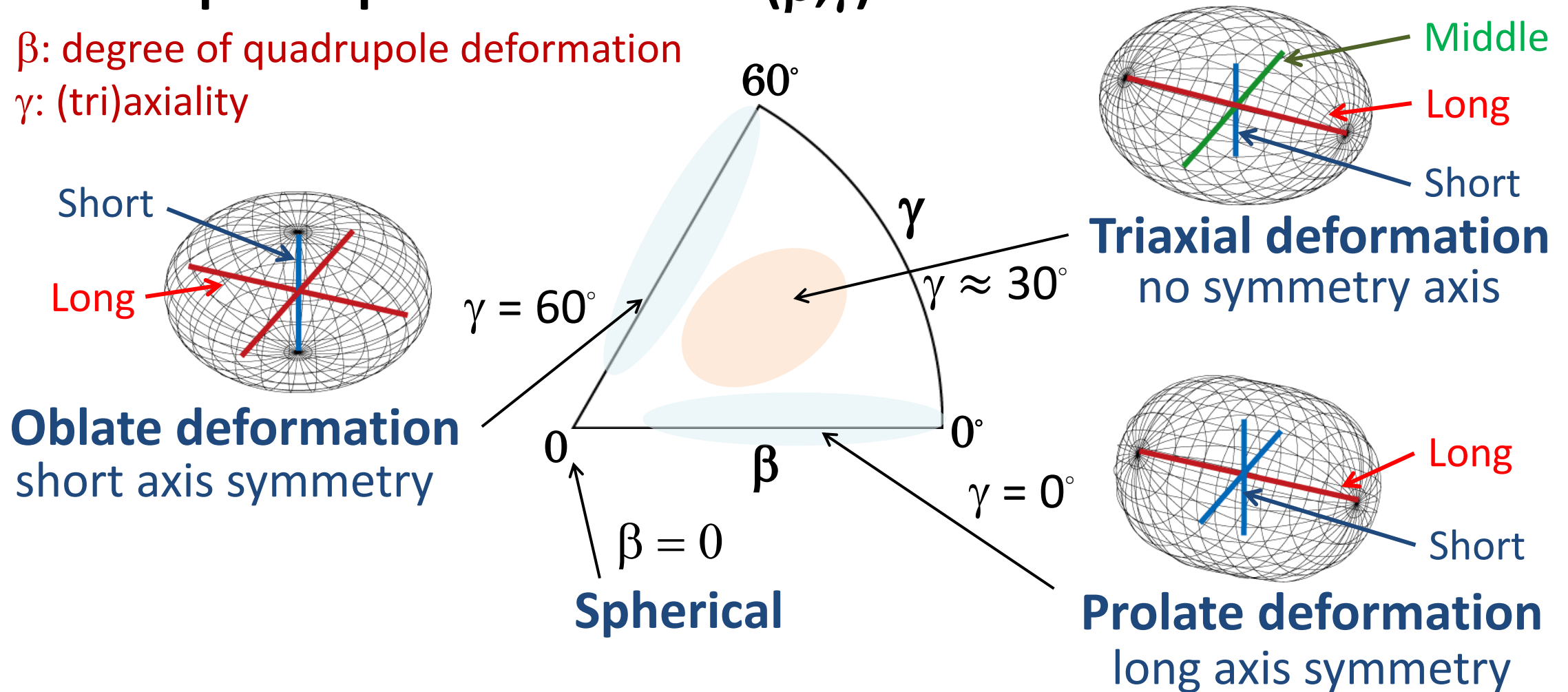
**What will happen if a  $\Lambda$  is coupled to nuclei with various structures ?**

# Deformation of nuclei

## ◆ Most of nuclei are deformed except for magic nuclei

### ● Nuclear quadrupole deformation ( $\beta, \gamma$ )

- $\beta$ : degree of quadrupole deformation
- $\gamma$ : (tri)axiality

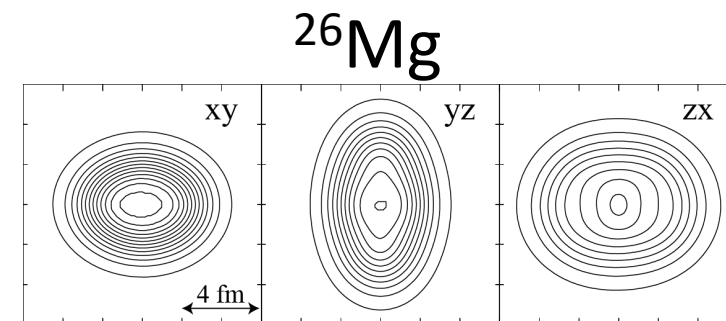
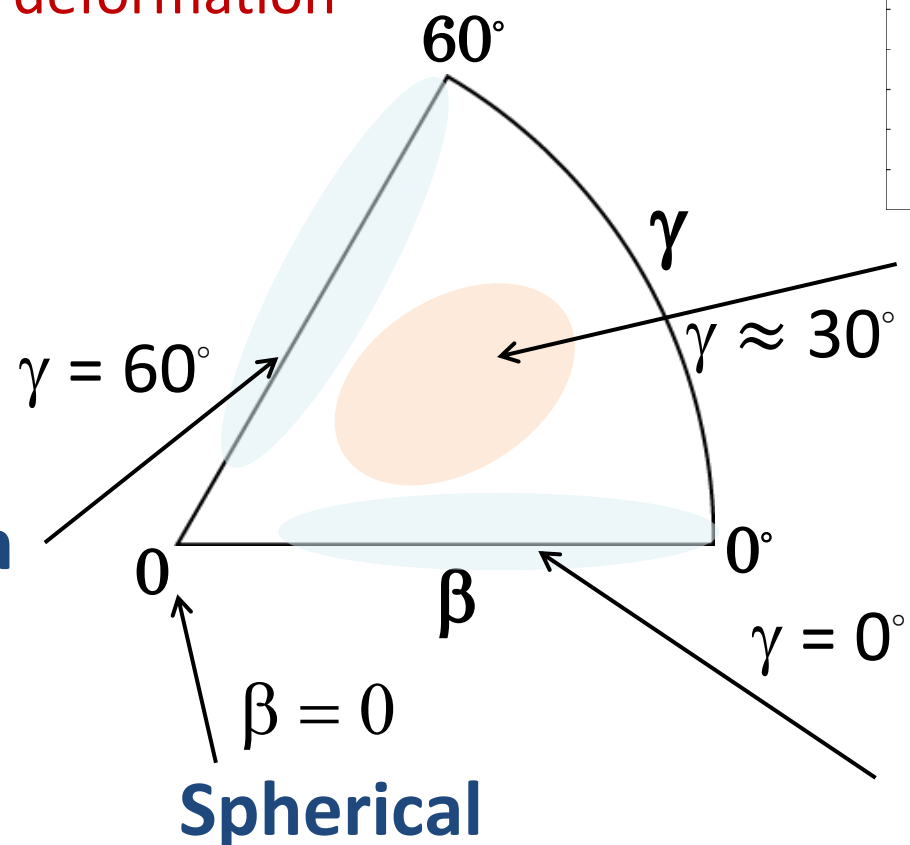
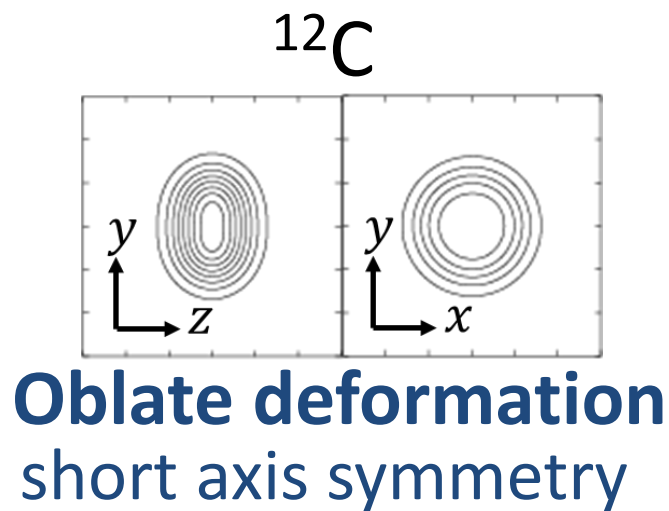


# Deformation of nuclei

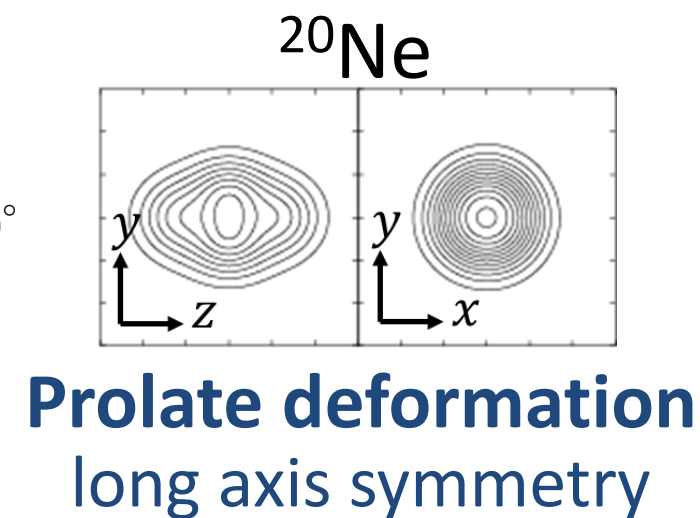
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**Triaxial deformation**  
no symmetry axis



# What will happen if a $\Lambda$ particle is coupled to nuclei ?

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- **Dynamical changes of nuclear structure**
  - Changes of cluster structure
  - Deformation changes
- **Sensitivity of the  $\Lambda$  binding energy  $B_\Lambda$  on nuclear structure**
  - Dependence of  $B_\Lambda$  on nuclear deformation
  - Mass number  $A$  dependence & many-body force effects
- **Coupling of  $\Lambda$  particle in  $p$  orbit to clustering/deformed core nuclei**
  - Genuine hypernuclear states
  - Possibility to probe nuclear deformation using  $\Lambda$  particle

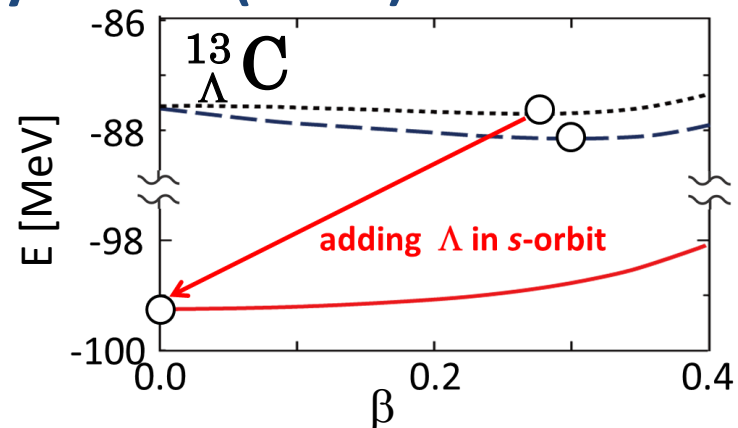


# Deformation change by $\Lambda$ particle

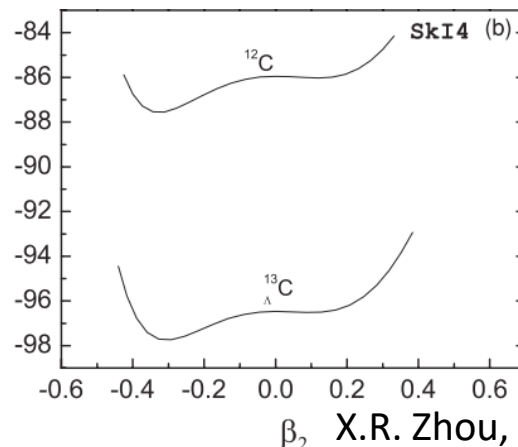
**Many authors predict that  $\Lambda$  in s-orbit reduces nuclear deformation**

**Antisymmetrized molecular**

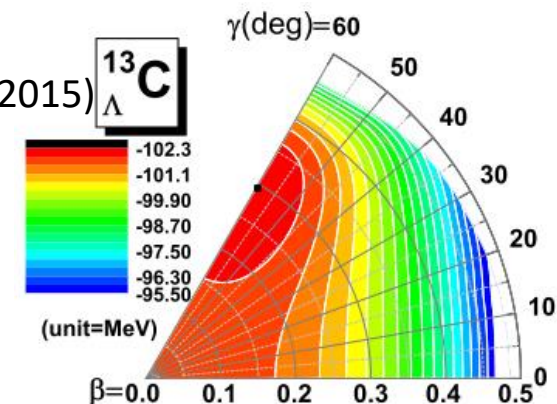
**Dynamics (AMD)** M.I, et al., PRC**83**, 044323(2011)



**Skyrme-Hartree-Fock (SHF)**

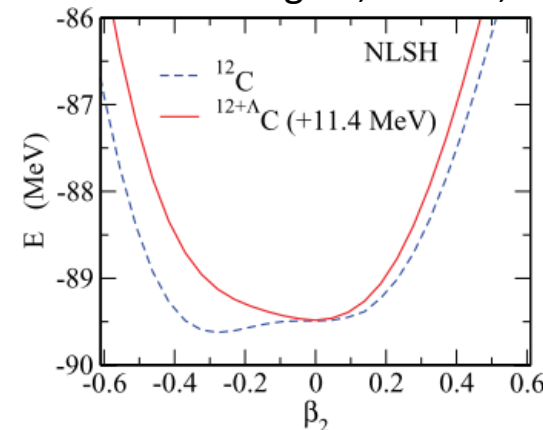


J.W. Cui, et al,  
PRC**91**,054306(2015)

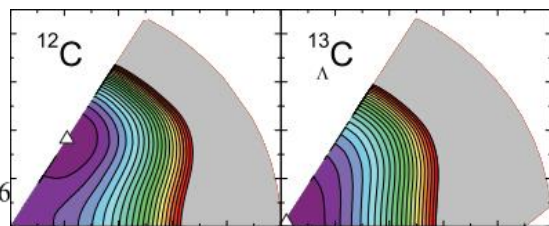


**Relativistic mean-field (RMF)**

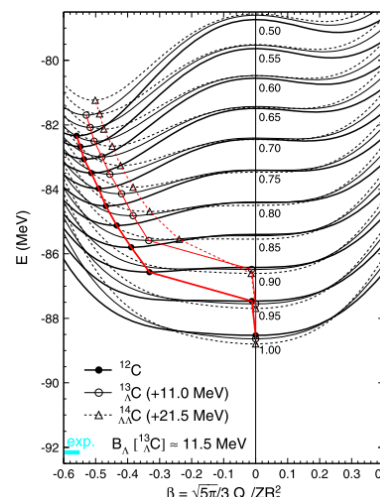
Win and Hagino, PR C**78**, 054311(2008)



B.N. Lu, et al., PRC**84**, 014328 (2011)



**RMF & SHF**



**Deformations/level structure with beyond-mean-field**

J.W. Cui, X.R. Zhou, H.J. Schulze,  
PRC**91**,054306('15)

H. Mei, K. Hagino, J.M. Yao, T. Motoba,  
PRC**91**, 064305(2015); **97**, 064318(2018)

H. J. Schulze, et al.,  
PTP**123**, 569('10)

*... and so on*

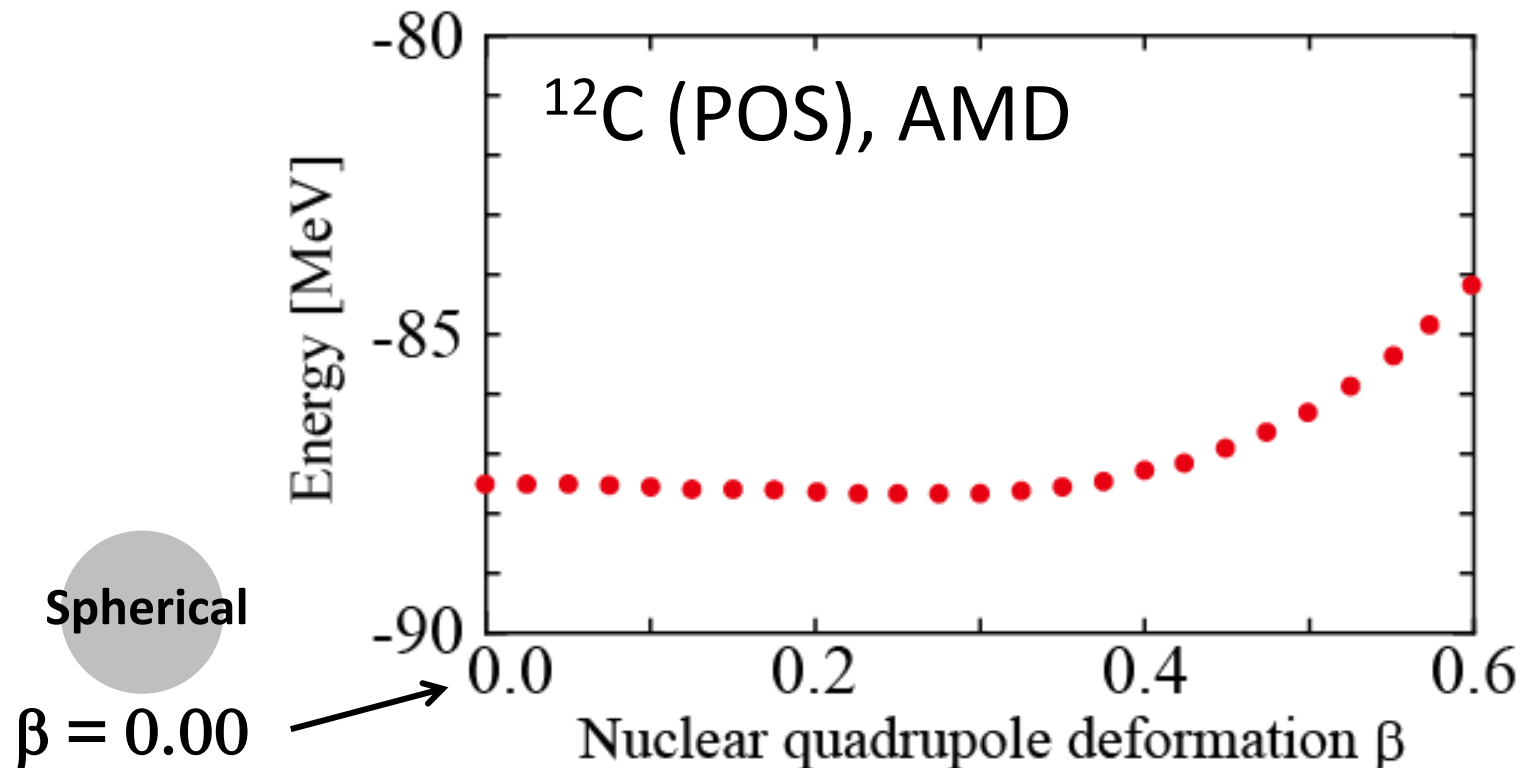
# Deformation change by $\Lambda$ particle

## ◆How to analyze from energy surface

Example:  $^{12}\text{C}$  with AMD(antisymmetrized molecular dynamics)

- Energy variation at each  $\beta$  (and  $\gamma$ )  $\rightarrow$  energy curve as a function of  $\beta$
- Energy minimum at  $(\beta, \gamma)$

M.I, et al., PRC83, 044323(2011)





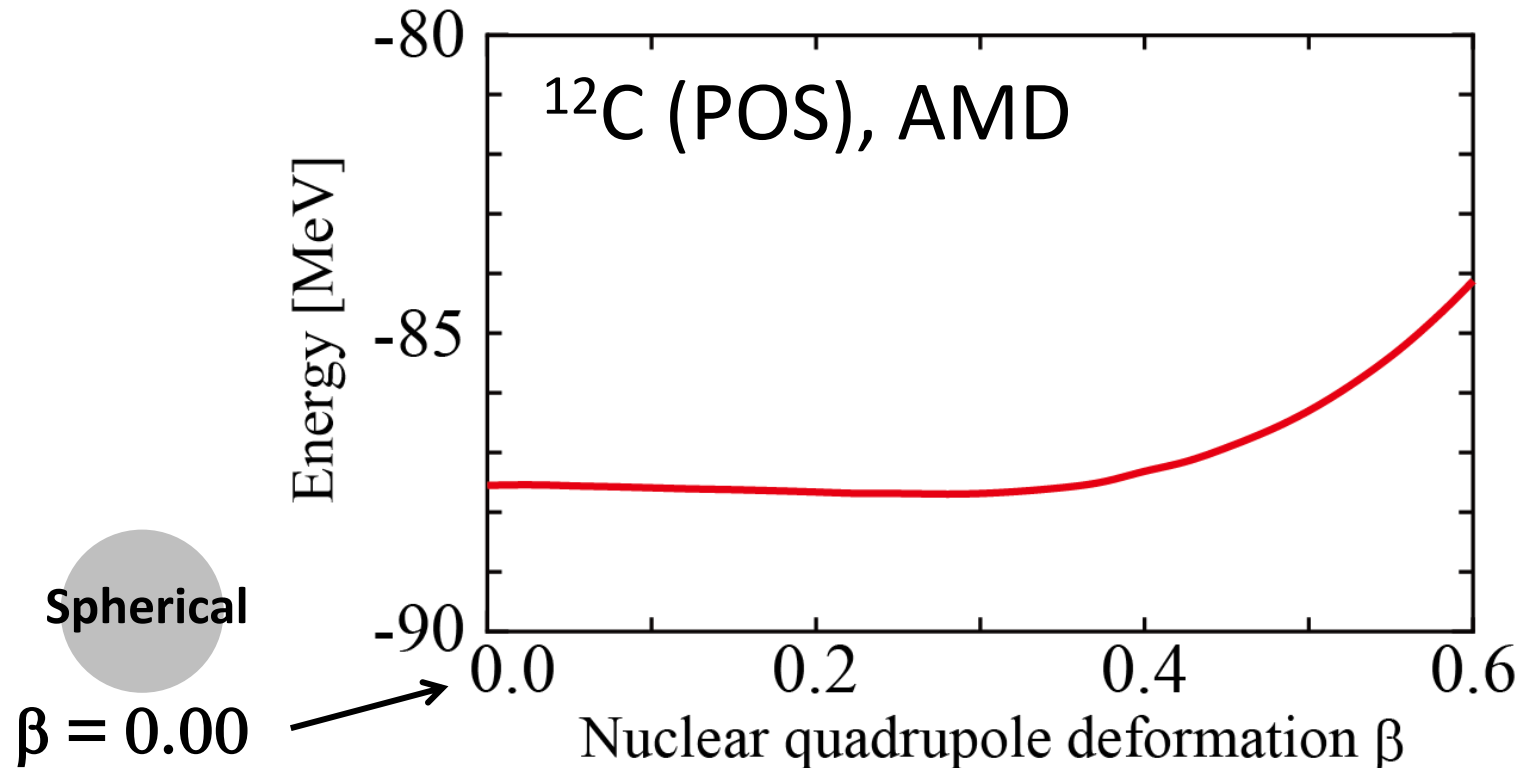
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M.I, et al., PRC83, 044323(2011)



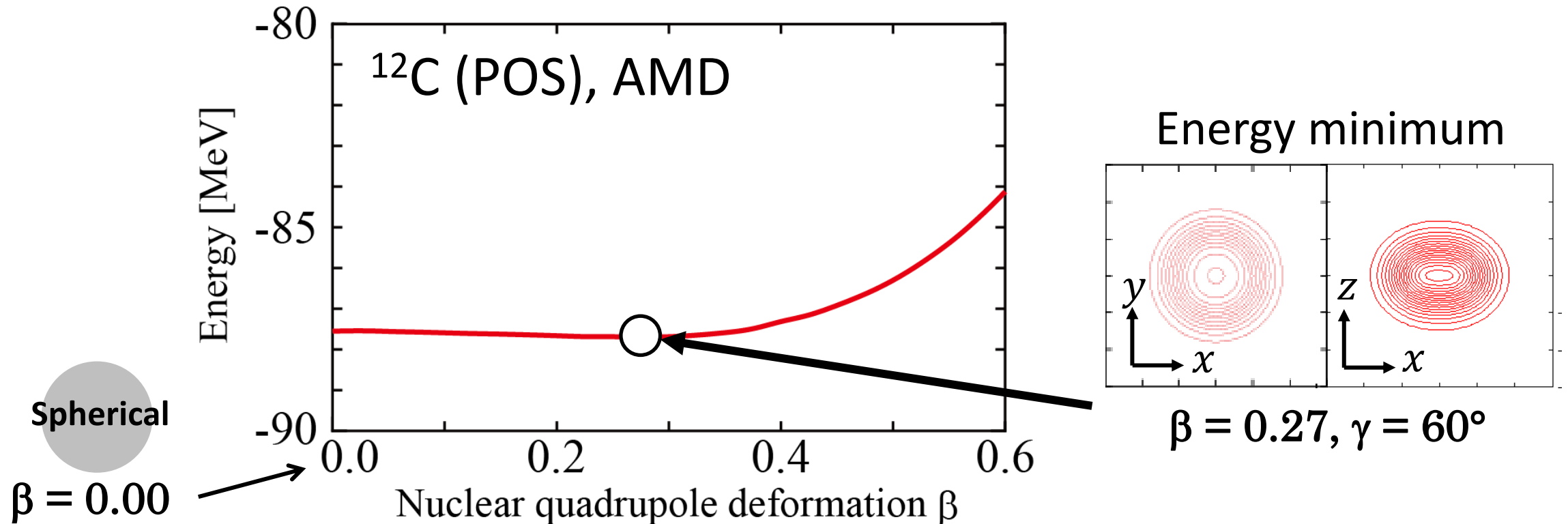
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M.I, et al., PRC83, 044323(2011)



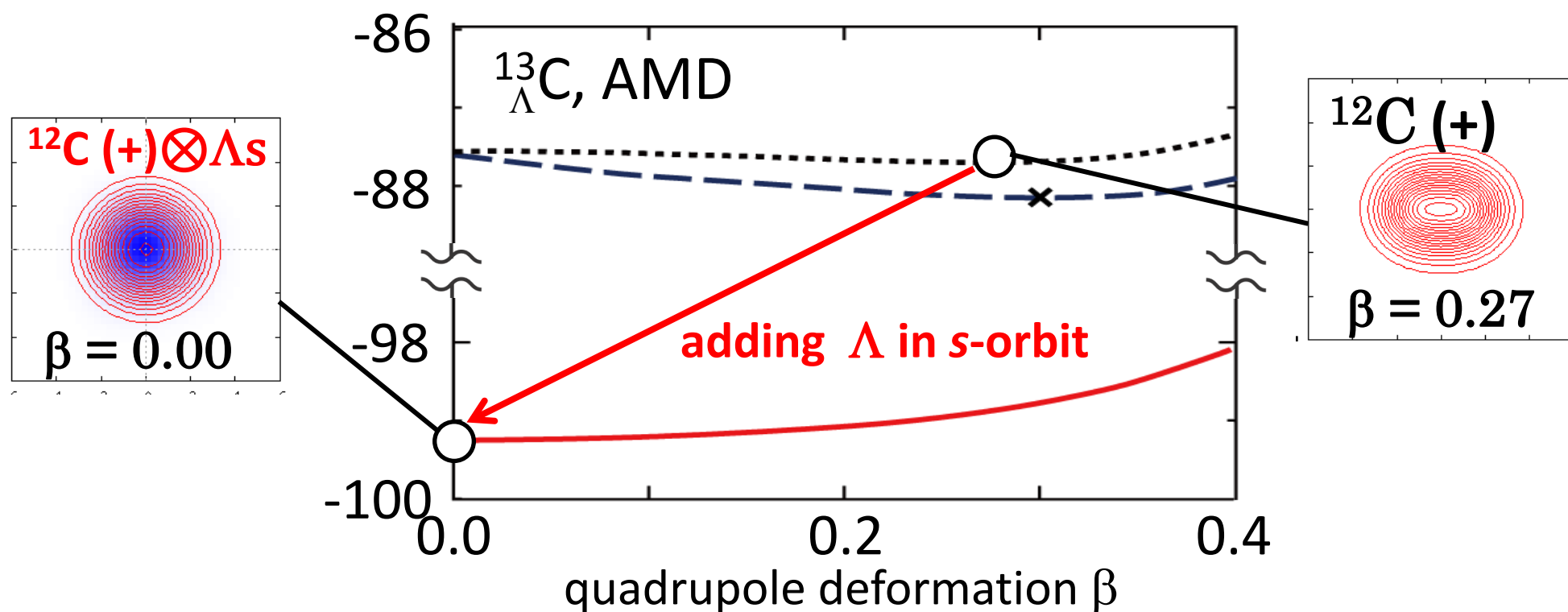
# Deformation change by $\Lambda$ in s-orbit

## ◆ How to analyze from energy surface

M.I, et al., PRC83, 044323(2011)

Example:  $^{12}\text{C}$  with AMD(antisymmetrized molecular dynamics)

- Energy variation of hypernucleus  
→ energy minimum moves to smaller  $\beta$  with  $\Lambda$  in s-orbit

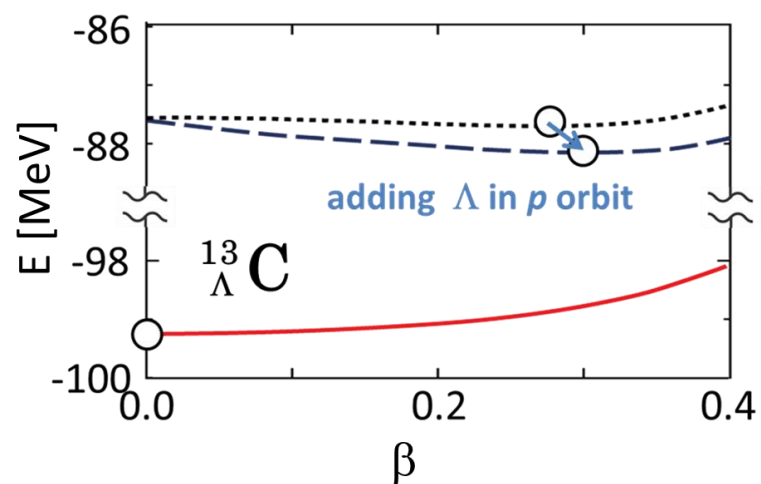


# Deformation change by $\Lambda$ particle

## $\Lambda$ in $p$ -orbit enhances nuclear deformation

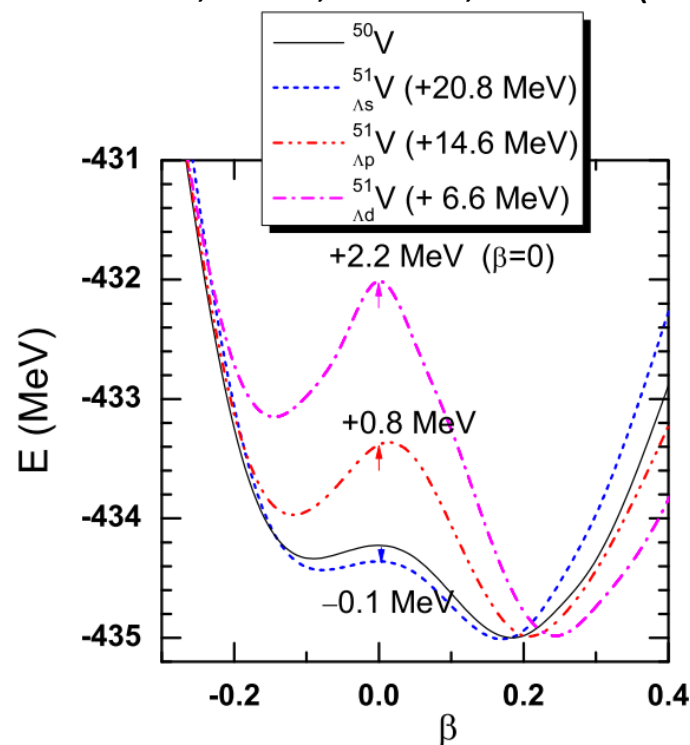
### Antisymmetrized molecular dynamics (AMD)

M.I, et al., PRC**83**, 044323(2011)



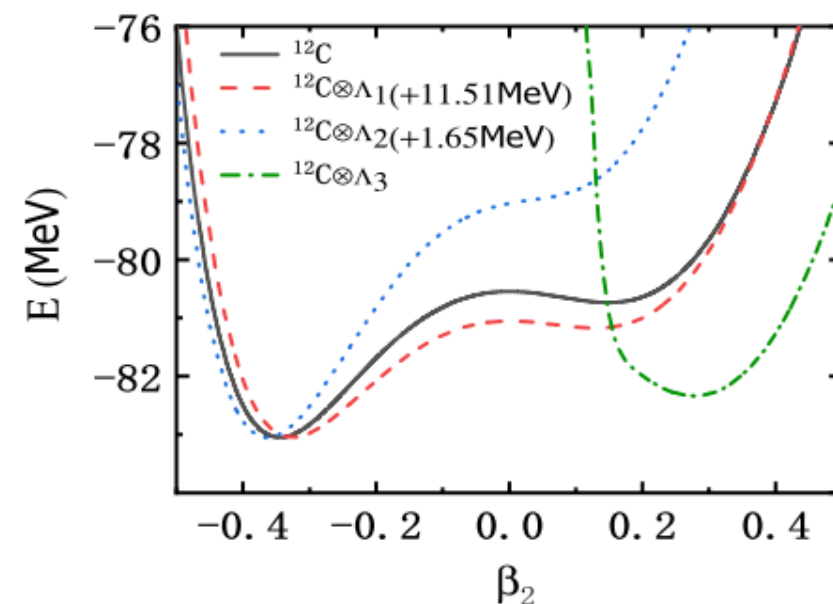
### Triaxially deformed relativistic mean-field

W. X. Xue, et al., PRC**91**, 024327(2015)



### Deformed Skyrme-Hartree-Fock (DSHF)

Bi-Cheng Fang et al., EPJA**56**,11(2020)



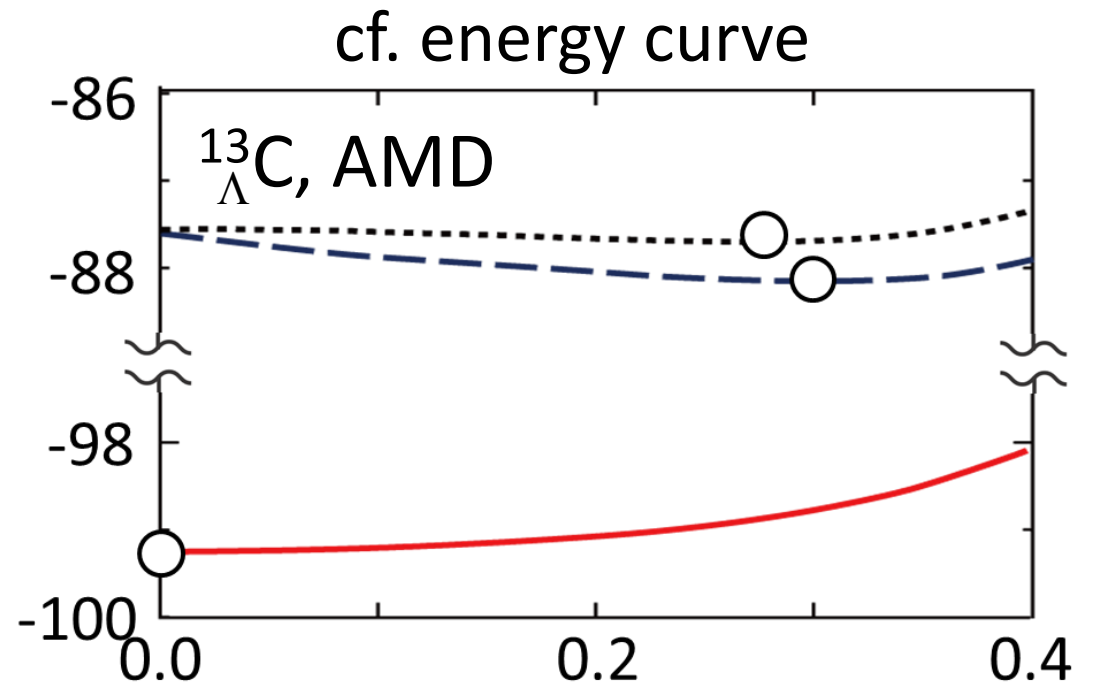
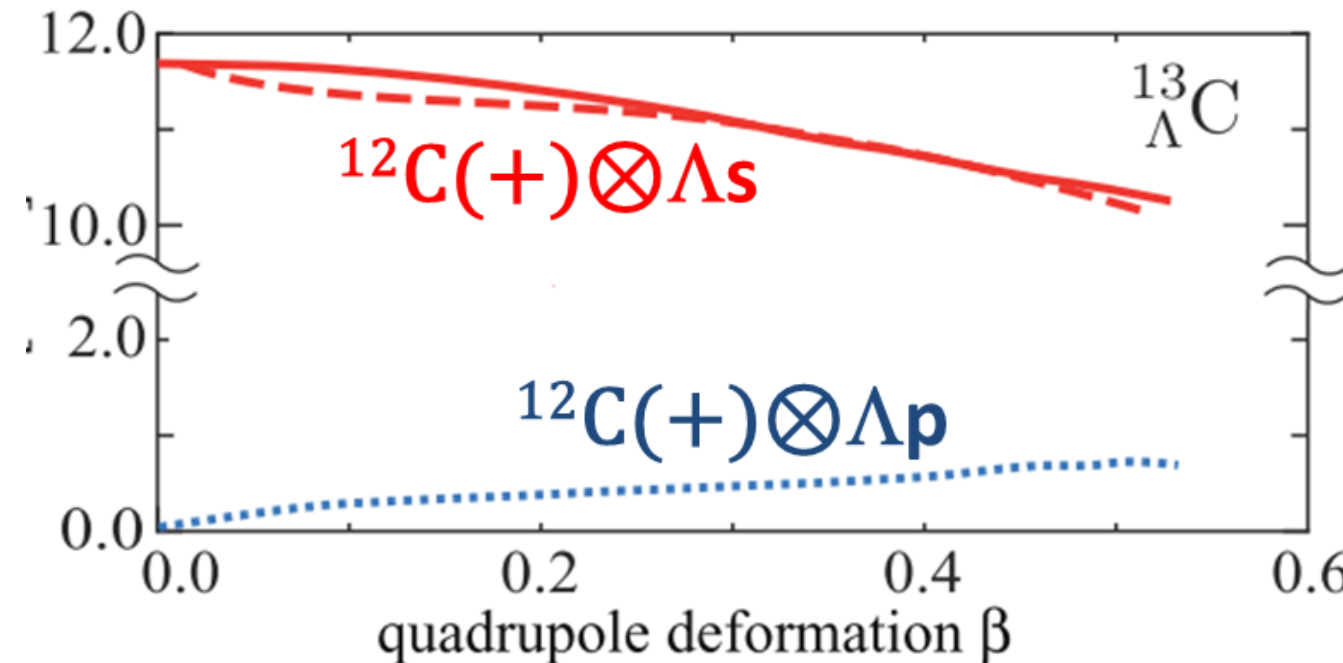
Deformation change of  $\Lambda$  in higher orbits such as  $d$ -orbit is also predicted by several papers:  
W. X. Xue, et al., PRC**91**, 024327(2015), X. Y. Wu, et al., PRC**95**, 034309(2017)

# Why does $\Lambda$ change nuclear deformation?

- $\Lambda$  in  $s$  orbit is deeply bound at small  $\beta$ , while  $\Lambda$  in  $p$  orbit prefers deformation
- Competition b/w  $\Lambda$  binding energy and energy surface of core nucleus

$$\Lambda \text{ binding energy: } b_{\Lambda}(\beta) = E_{core}(\beta) - E_{\Lambda}(\beta)$$

$b_{\Lambda}$  [MeV]

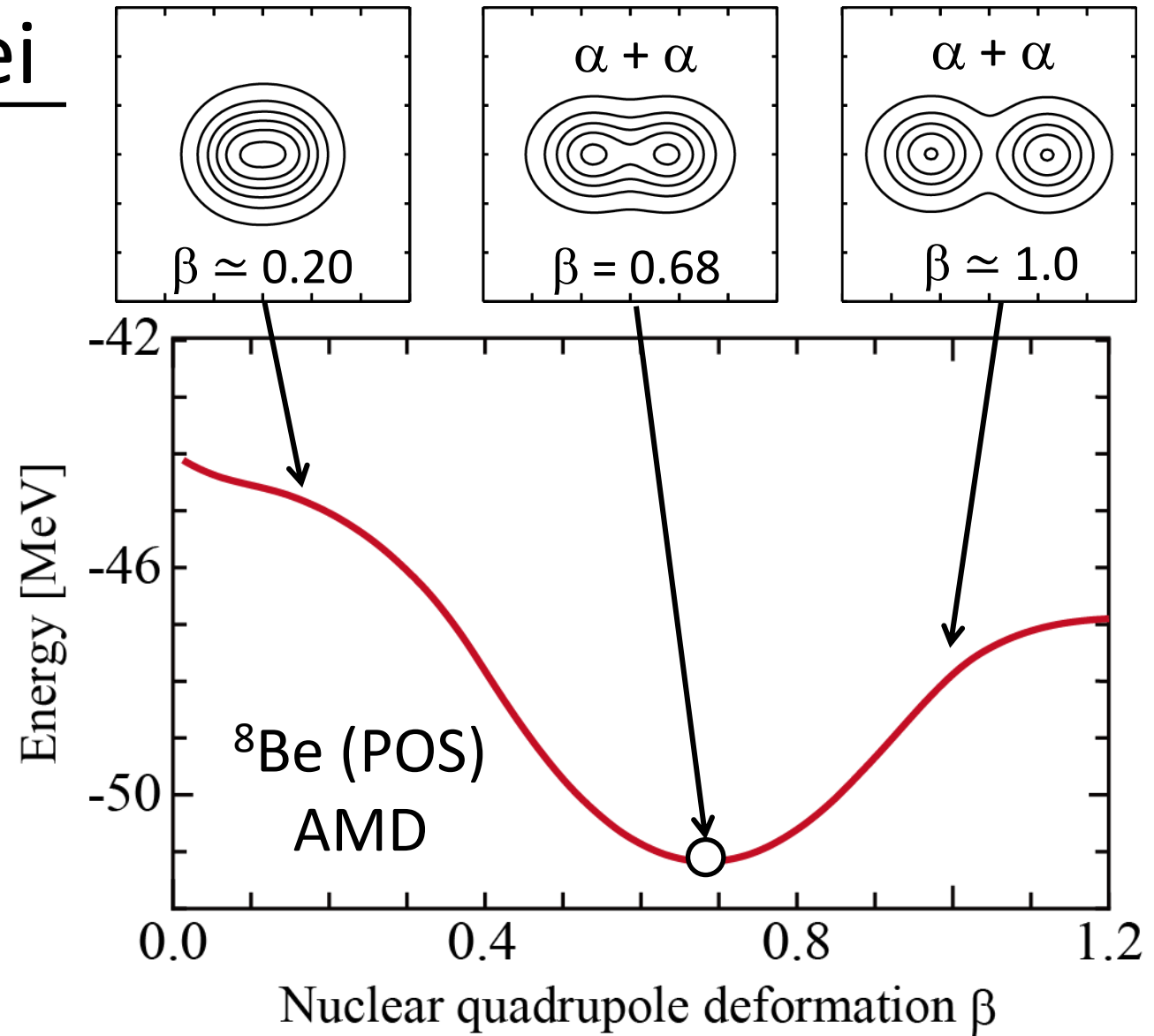
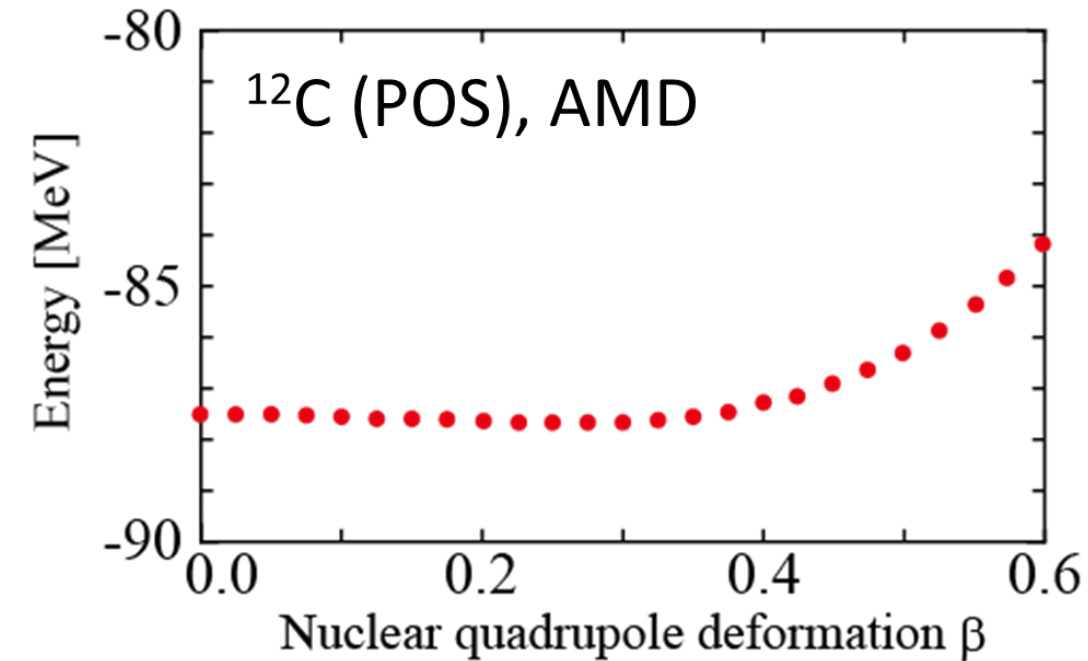


M.I., et al., PRC83, 044323(2011)

“binding energy of  $\Lambda$ ” vs. “energy surface of the core nuclei”

# Energy surface of core nuclei

Shape of energy surface is related to nuclear structure

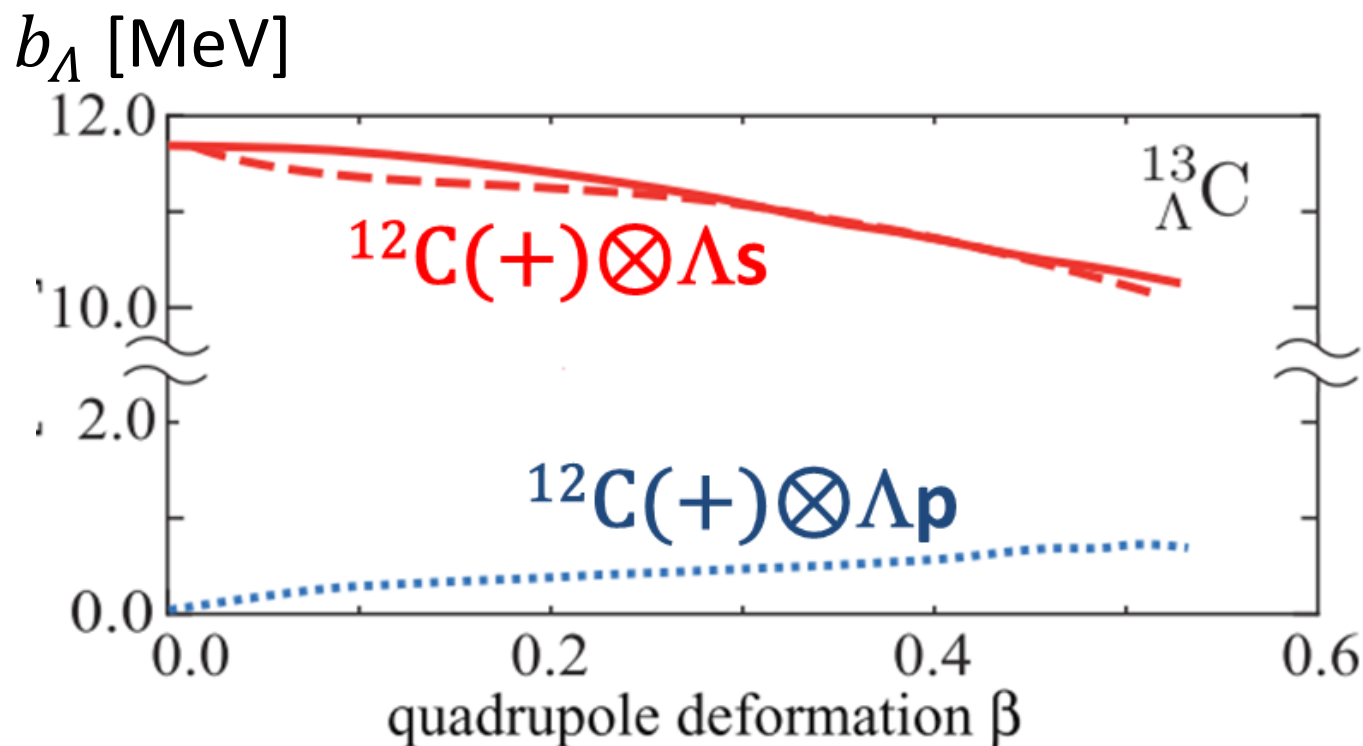



Energy increases when  $2\alpha$  clustering is reduced/enhanced from the minimum

# Deformation change by $\Lambda$ particle


**“Overlap between  $\Lambda$  and N” is the key!**

$\Lambda$  in s-orbit ( $p$ -orbit) is deeply bound with smaller  $\beta$  (larger  $\beta$ ) due to larger overlap between  $\Lambda$  and nucleons



$\Lambda$ in s orbit		Small $\beta$	Large $\beta$
Overlap b/w $\Lambda$ & N		Large	Small
$\Lambda$ N attraction		Large	Small

$\Lambda$ in $p$ orbit		Small $\beta$	Large $\beta$
Overlap b/w $\Lambda$ & N		Small	Large
$\Lambda$ N attraction		Small	Large

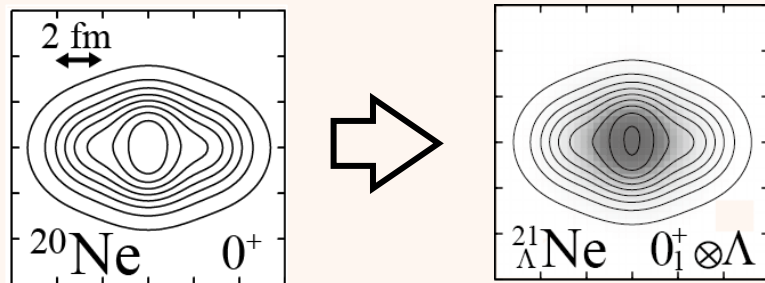
# Structure dependence of “impurity effects”

## ◆ Example: $^{21}_{\Lambda}\text{Ne}$ (prediction by AMD calc)

M. Isaka, et al., PRC83, 054304(2011)

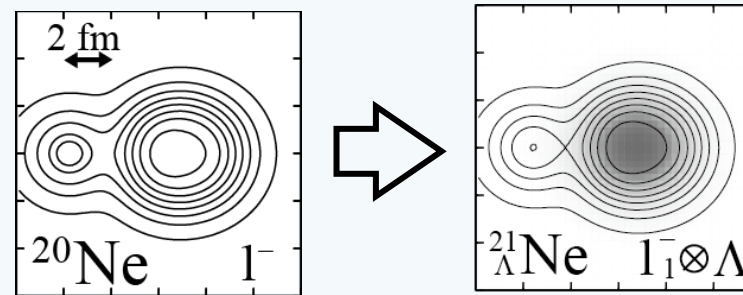
- Shrinkage/deformation change are larger in  $\alpha + ^{16}\text{O} + \Lambda$  cluster states, which appears as difference in intra-band B(E2) reduction

### Ground band

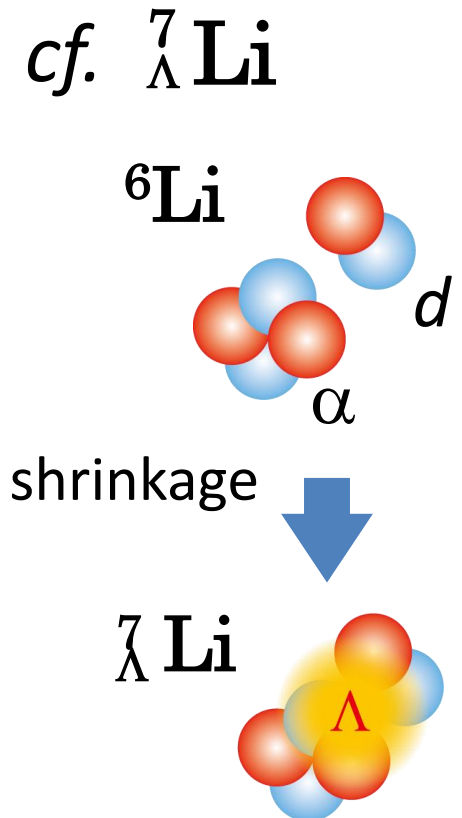


$^{20}\text{Ne}$ $K^\pi=0^+$	$r_{\text{RMS}}(\text{fm})$	$^{21}_{\Lambda}\text{Ne}$ $0^+ \otimes \Lambda s$	$r_{\text{RMS}}(\text{fm})$	$\Delta r_{\text{RMS}}(\text{fm})$
0+	2.97	(1/2)+	2.92	-0.05
2+	2.96	(3/2)+ (5/2)+	2.91	-0.05
4+	2.93	(7/2)+ (9/2)+	2.87	-0.06
6+	2.87	(11/2)+ (13/2)+	2.81	-0.05
			2.83	-0.04

### $K^\pi = 0^- (\alpha + ^{16}\text{O})$ band



$^{20}\text{Ne}$ $K^\pi=0^-$	$r_{\text{RMS}}(\text{fm})$	$^{21}_{\Lambda}\text{Ne}$ $0^- \otimes \Lambda s$	$r_{\text{RMS}}(\text{fm})$	$\Delta r_{\text{RMS}}(\text{fm})$
1-	3.27	(1/2)-	3.15	-0.11
3-	3.24	(3/2)- (5/2)- (7/2)-	3.15	-0.11
5-	3.23	(9/2)- (11/2)-	3.13	-0.11
7-	3.23	(13/2)- (15/2)-	3.14	-0.10
			3.11	-0.12
			3.11	-0.13
			3.06	-0.17
			3.05	-0.18





# What will happen if a $\Lambda$ particle is coupled to nuclei ?

---

- **Dynamical changes of nuclear structure**
  - Changes of cluster structure
  - Deformation changes
- **Sensitivity of the  $\Lambda$  binding energy  $B_\Lambda$  on nuclear structure**
  - Dependence of  $B_\Lambda$  on nuclear deformation
  - Mass number  $A$  dependence & many-body force effects
- **Coupling of  $\Lambda$  particle in  $p$  orbit to clustering/deformed core nuclei**
  - Genuine hypernuclear states
  - Possibility to probe nuclear deformation using  $\Lambda$  particle

# Dependence of $B_\Lambda$ on nuclear deformation

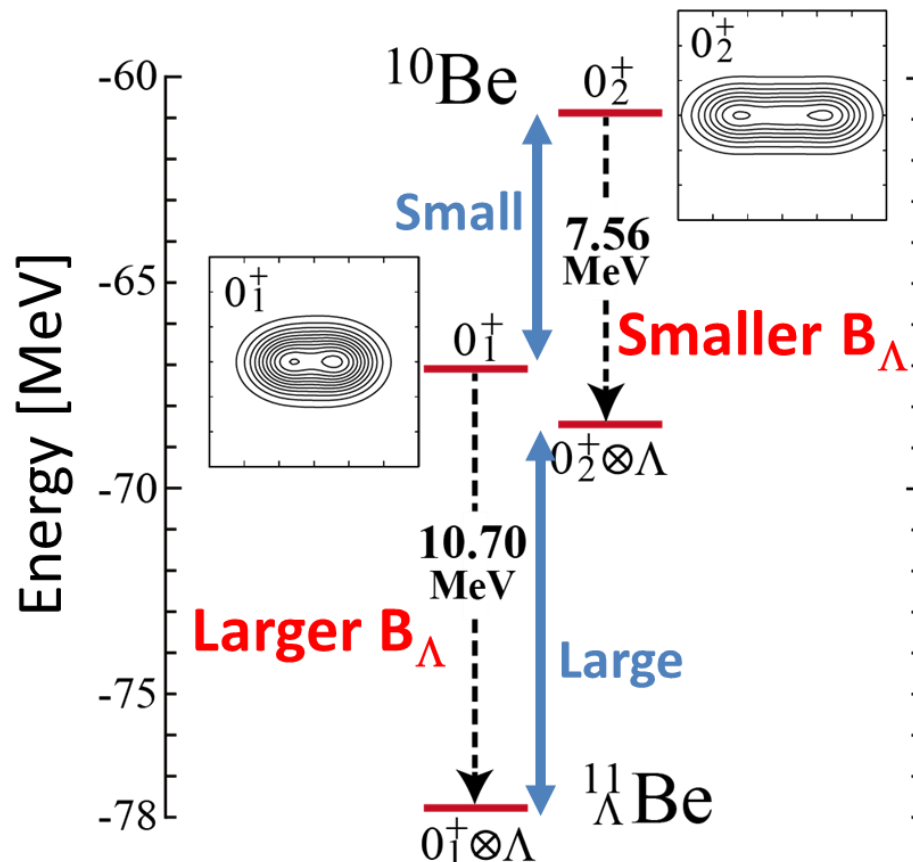
**$B_\Lambda$  is smaller in largely deformed state than less deformed state**

**Example:**

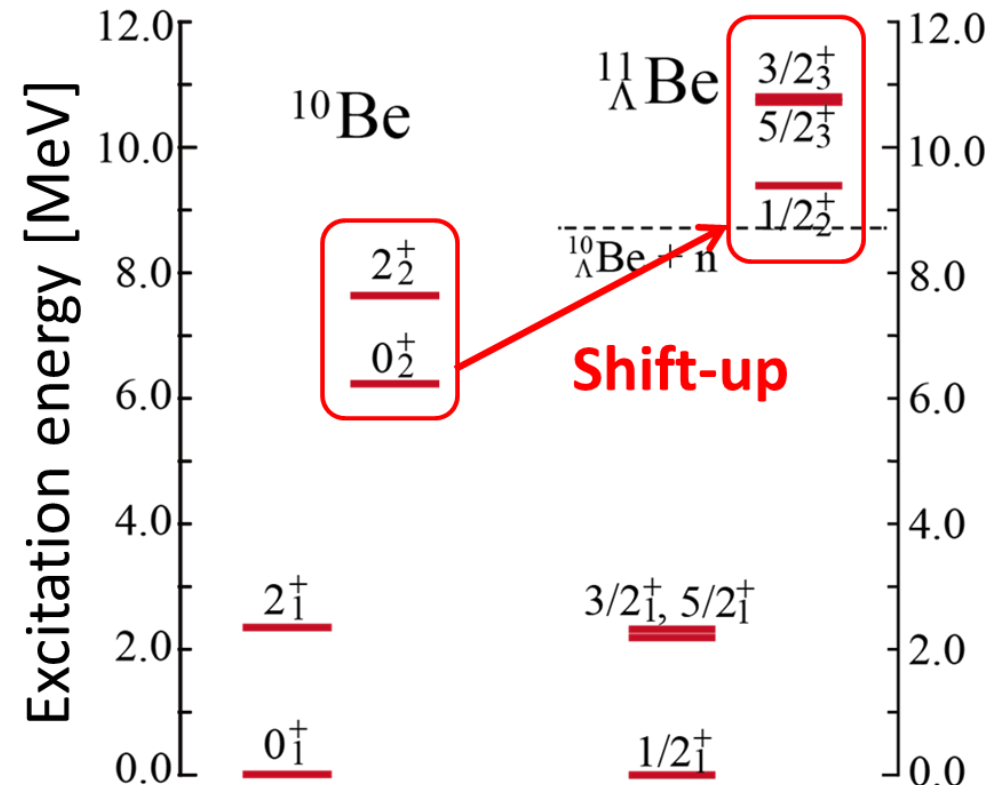


AMD calc

M.I. and Kimura,  
PRC92 (2015)



**If different deformation coexist, shift-up/down in excitation spectra**



# Dependence of $B_\Lambda$ on nuclear **structure**

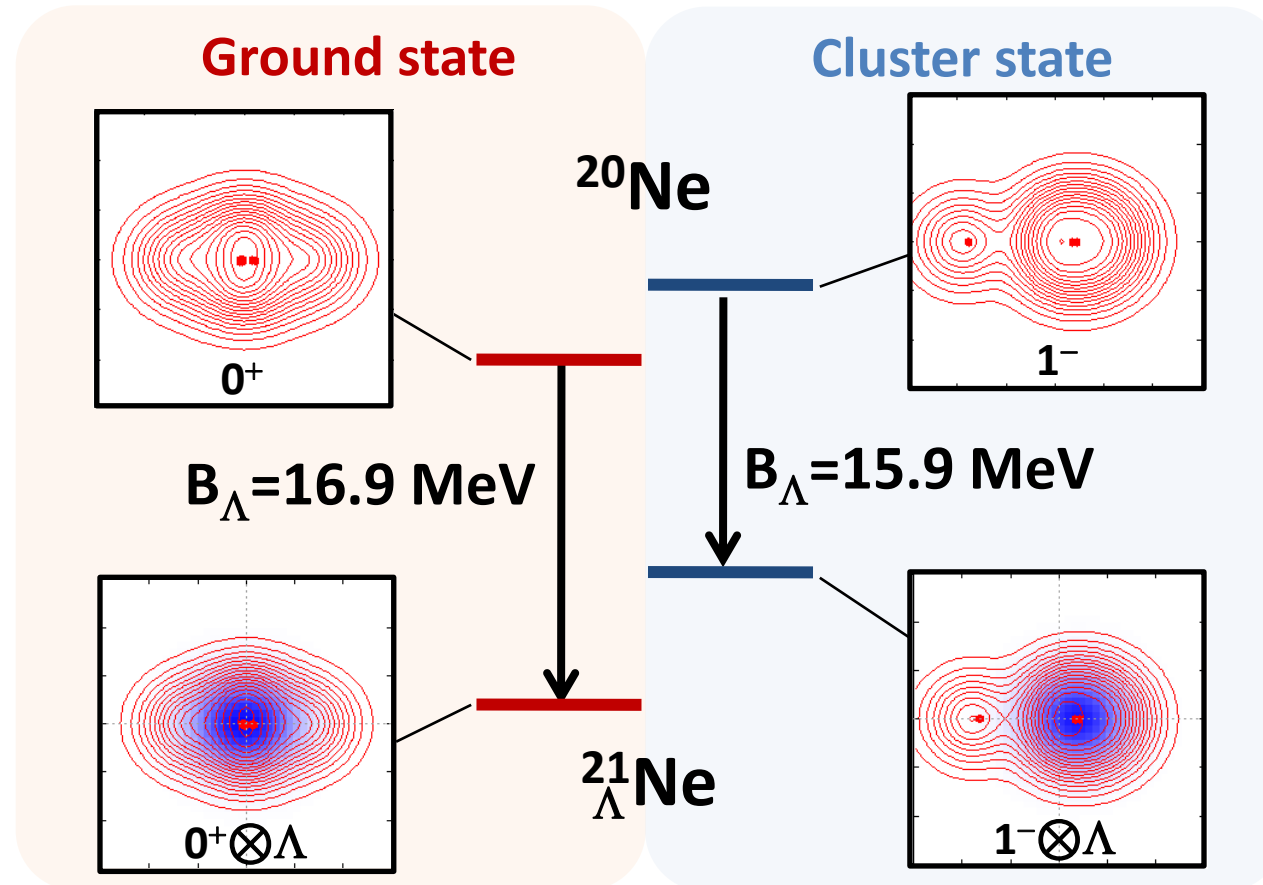
**$B_\Lambda$  is smaller in cluster states than mean-field like states**

**Example:**



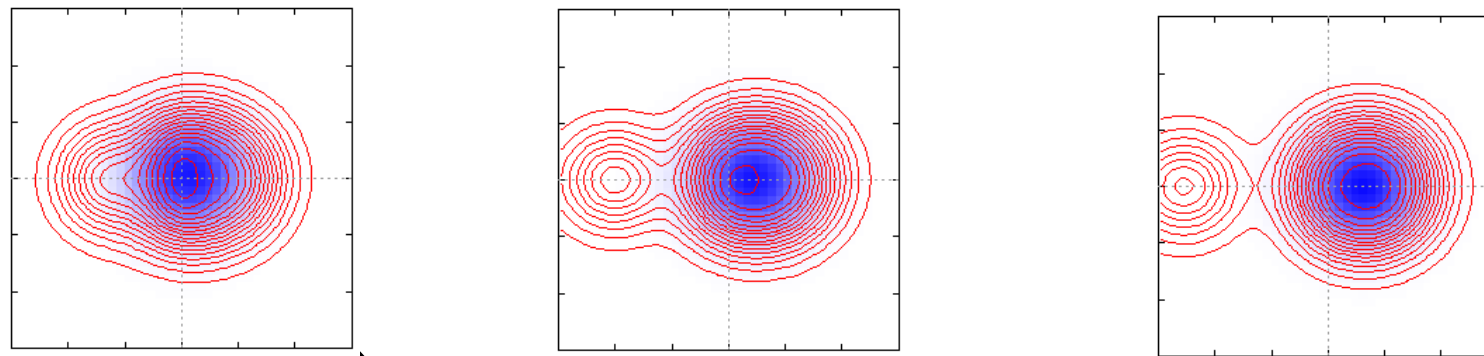
- Different structures coexist near the ground state
- $\Lambda$  in s orbit changes them, but the difference remains
- $\Lambda$  is localized around  $^{16}\text{O}$  in  $\alpha + ^{16}\text{O} + \Lambda$  state  $\rightarrow$  difference of  $B_\Lambda$

AMD calc  
M. Isaka, et al.,  
PRC83, 054304(2011)

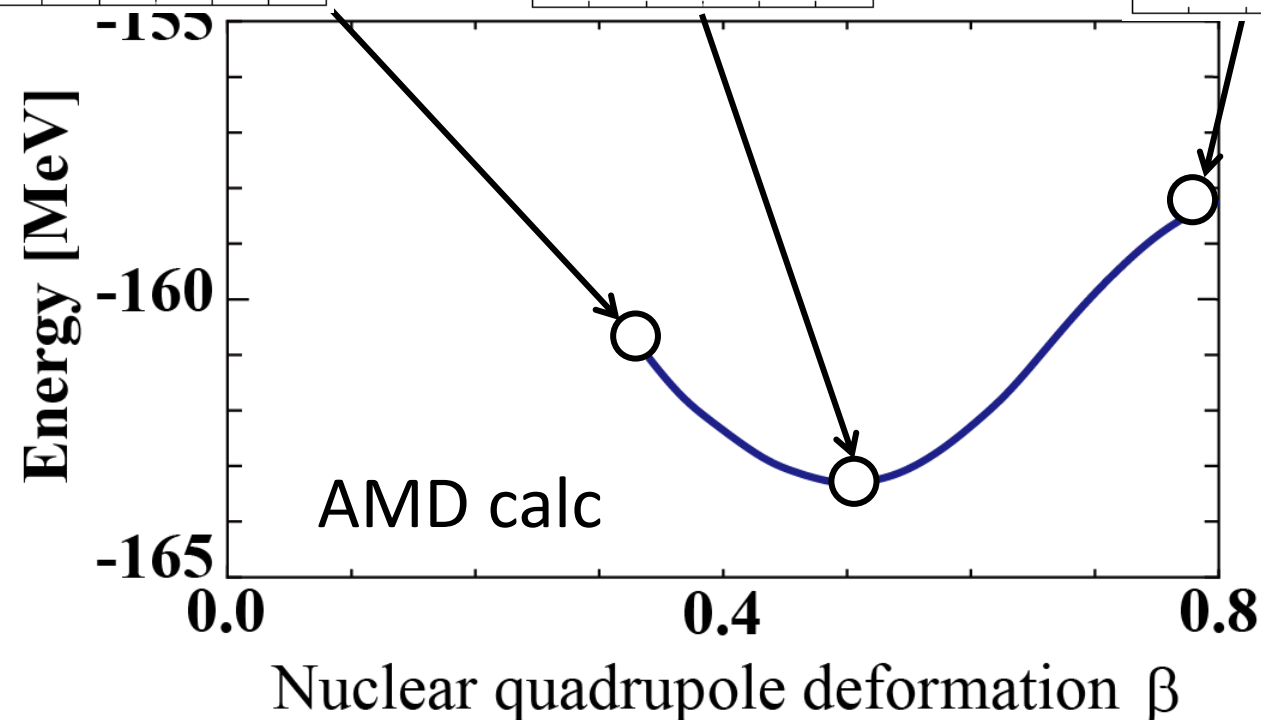


# Dependence of $B_\Lambda$ on nuclear **structure**

*Which cluster does a  $\Lambda$  particle prefer in  $\alpha + {}^{16}\text{O} + \Lambda$  state?*



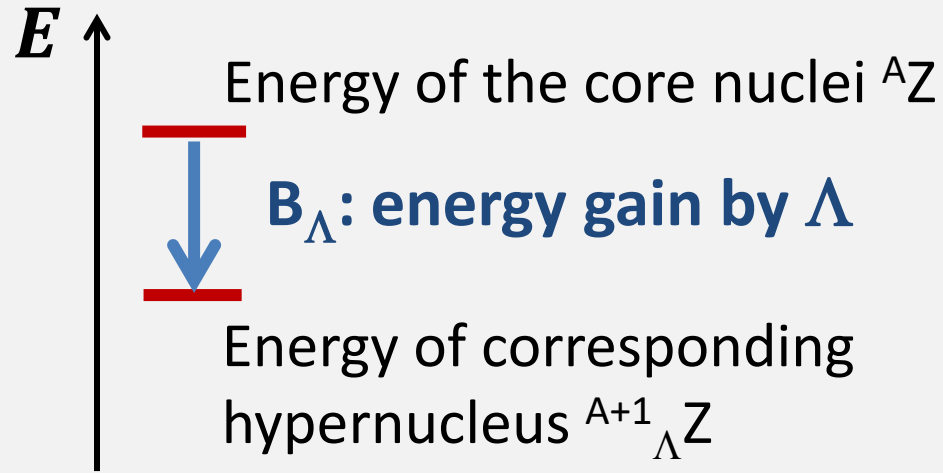
**“ $\alpha + {}^{17}_\Lambda\text{O}$  like”**  
 $\Lambda$  prefers the  ${}^{16}\text{O}$  cluster



Since  $B_\Lambda({}^{21}_\Lambda\text{Ne}) > B_\Lambda({}^{17}_\Lambda\text{O})$ ,  
 $B_\Lambda$  in  $\alpha + {}^{16}\text{O} + \Lambda$  states is smaller  
than the ground state in  ${}^{21}_\Lambda\text{Ne}$

# A-dependence of $B_\Lambda$ and many-body force effects

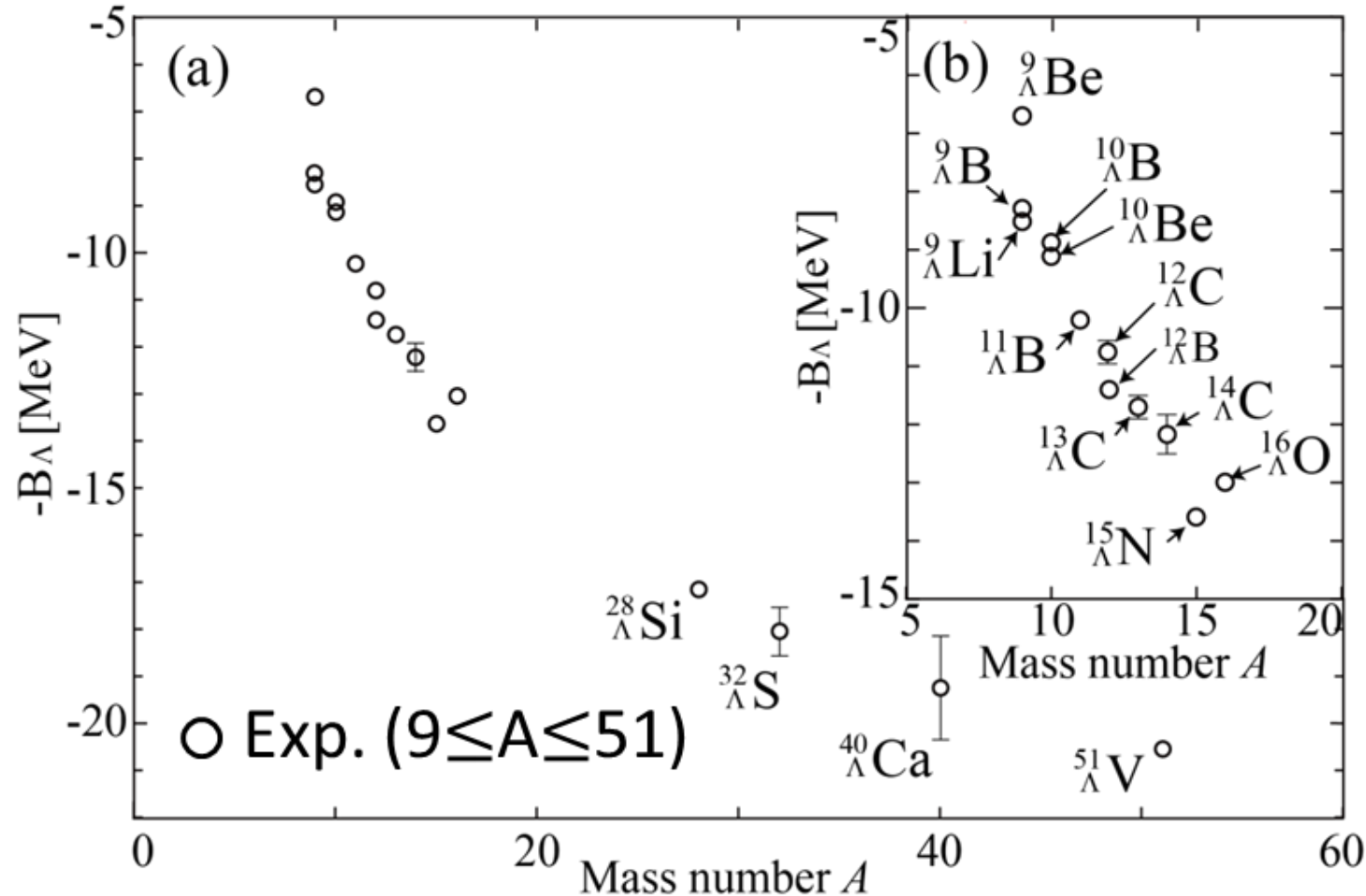
## $B_\Lambda$ : a most basic quantity



To describe  $A$  dependence of  $B_\Lambda$ ,

- Structure of hypernuclei
- YN and YNN interactions should be essential

Exp. data: Hashimoto & Tamura, PPNP57,564(2006) and references therein, Tang, *et. al.*, PRC90,034320(2014).



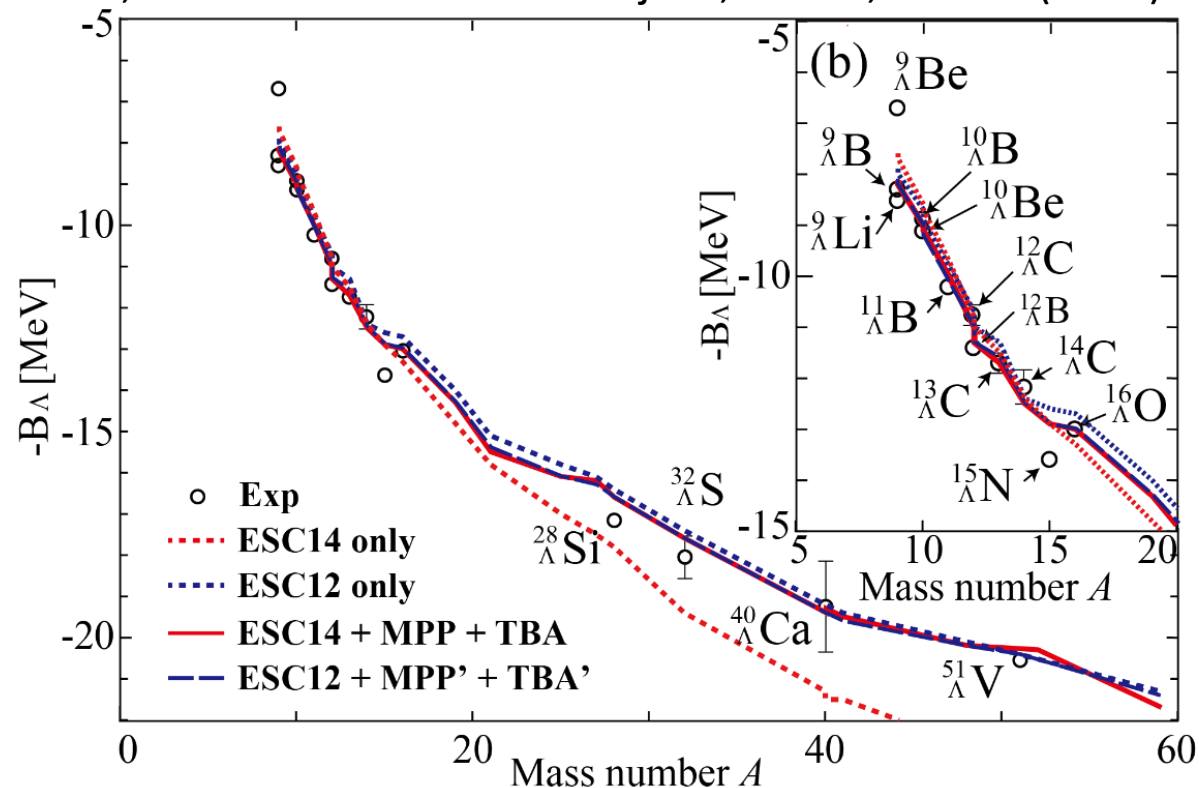
By describing structure of hypernuclei properly using appropriate  $\Lambda N$  interaction, many-body force (YNN) effects can be seen in  $A$  dep. of  $B_\Lambda$

# Some examples of theoretical attempts

Energy difference b/w  $B_{\Lambda}^{\text{cal}}$  with  $\Lambda N$  force and  $B_{\Lambda}^{\text{exp}}$  is a room for many-body force

## AMD w/ meson exch. $\Lambda N$ force (Nijmegen)

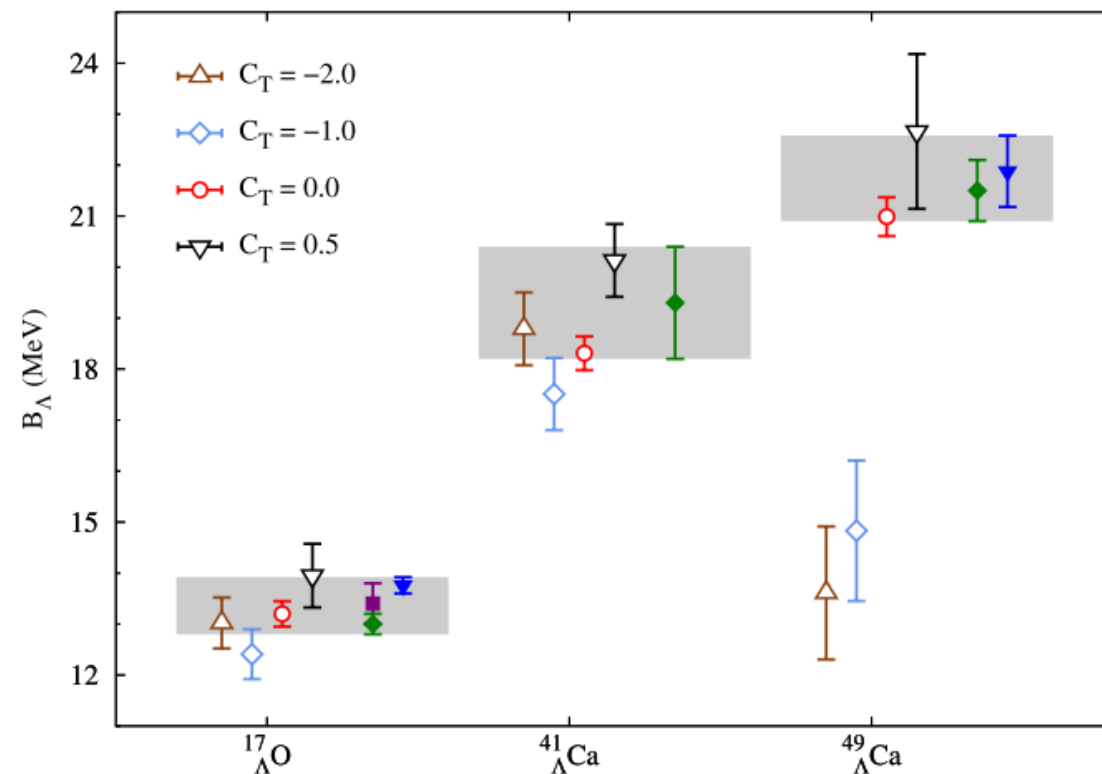
M. I, Y. Yamamoto and Th.A. Rijken, PRC95, 044308(2017)



- Taking into account hypernuclear deformation
- Ambiguity of  $\Lambda N$  potential vs. strength of  $\Lambda NN$  force

## auxiliary field diffusion Monte Carlo (AFDMC)

D. Lonardone et al., arXiv:1711.07521v3 (2018)



- Isospin dependence of  $\Lambda NN$  force
- Related to future JLab exp. (E12-15-008)

# Isospin-dependence of $\Lambda$ NN force

## ◆auxiliary field diffusion Monte Carlo (AFDMC) calculation

D. Lonardone et al., arXiv:1711.07521v3 (2018)

- The authors consider isospin dependence of  $\Lambda$ NN force related to charge symmetry breaking (CSB) in hypernuclei
- Isospin dependence is examined by artificially varying a parameter  $C_T$

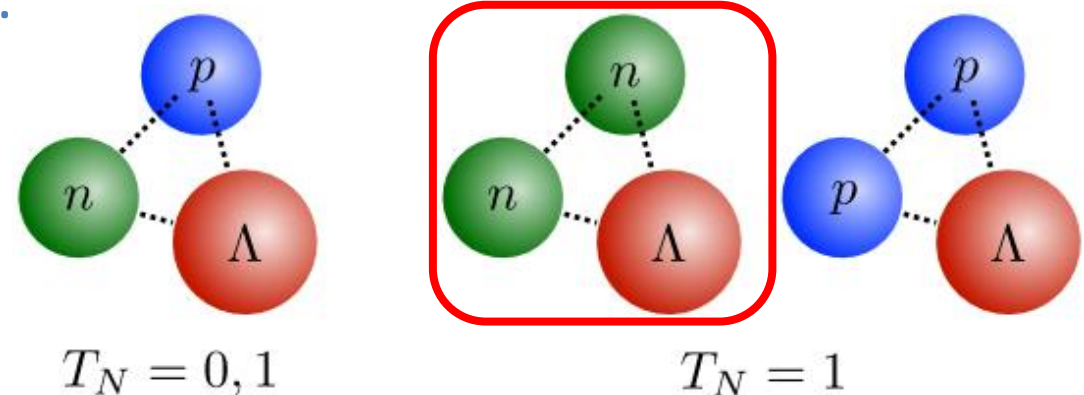
$$v_{\lambda ij}^T \boldsymbol{\tau}_i \cdot \boldsymbol{\tau}_j = -3 v_{\lambda ij}^T \mathcal{P}_{ij}^{T_N=0} + (1 + C_T) v_{\lambda ij}^T \mathcal{P}_{ij}^{T_N=1}$$

$C_T = 0$ : original  $\Lambda$ NN corr. to  $\Lambda$ - $\Sigma^0$  mixing mech.

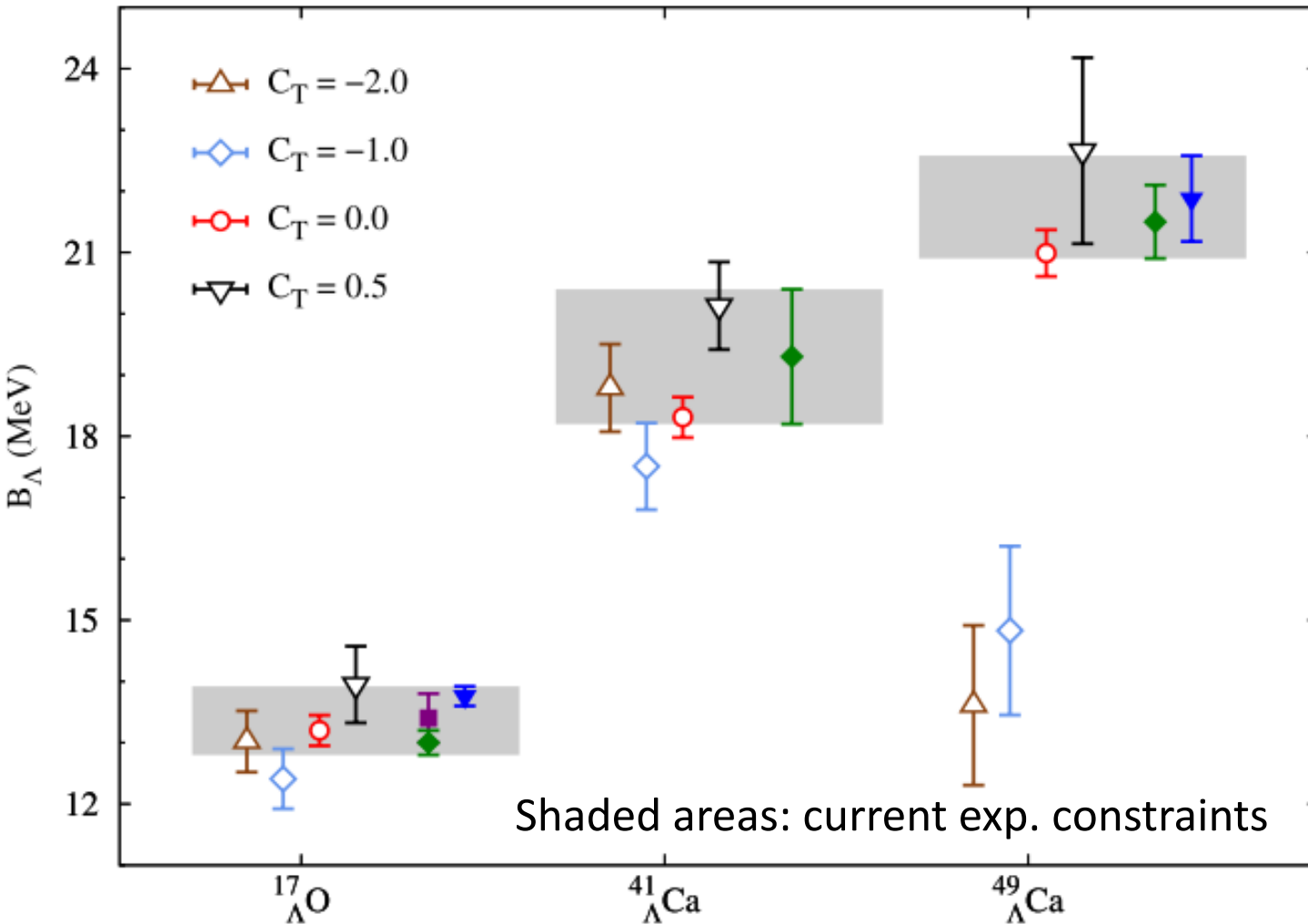
$C_T = -1$ :  $T_N = 1$  component is off

$C_T = -2$ : changing sign of  $T_N = 1$  comp.

Most important in n-rich systems



# Isospin-dependence of $\Lambda$ NN force



D. Lonardone et al., arXiv:1711.07521v3 (2018)

$$v_{\lambda ij}^T \boldsymbol{\tau}_i \cdot \boldsymbol{\tau}_j : \\ = -3 v_{\lambda ij}^T \mathcal{P}_{ij}^{T_N=0} + (1 + C_T) v_{\lambda ij}^T \mathcal{P}_{ij}^{T_N=1}$$

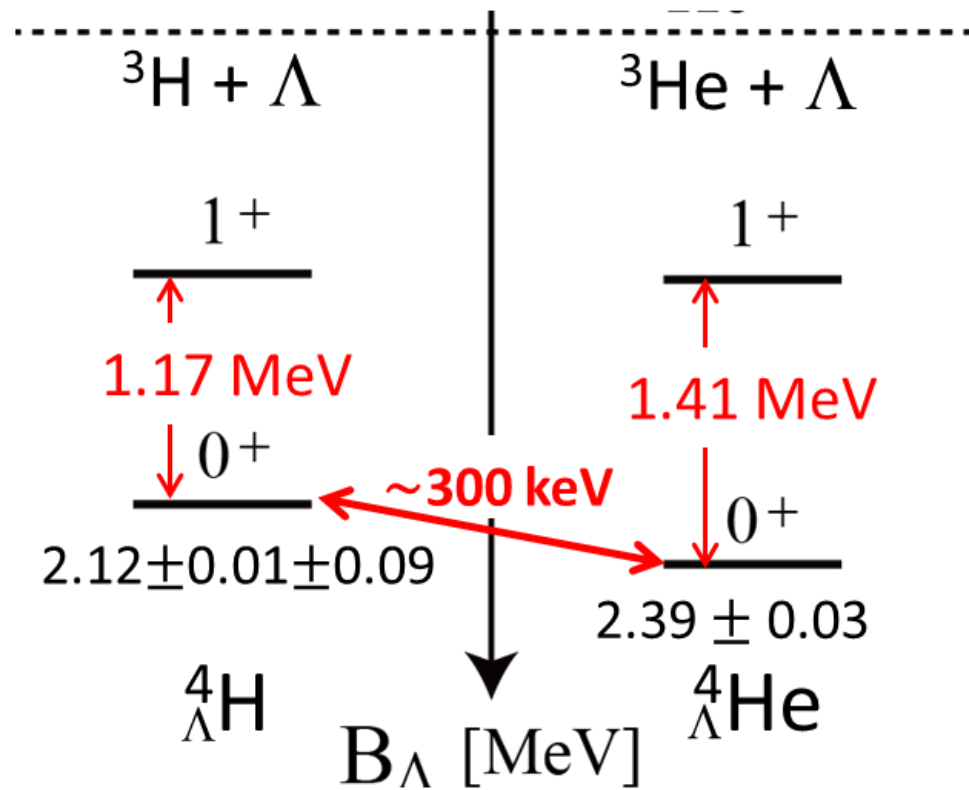
- Isospin dep. largely affects  $B_\Lambda$  in heavier  $Z \neq N$  hypernuclei
- Exp. data are not sufficiently accurate in  $40 \lesssim A \lesssim 50$

**$C_T$  can be determined by comparing with future JLab exp. (E12-15-008), which provides us new insight on many-body force**



# Related topic: charge symmetry breaking (CSB)

## ◆ CSB in $A=4$ $\Lambda$ hypernuclei



### *cf.* Charge symmetry in nuclei

- Binding energy, energy levels, ... etc. are almost the same in mirror nuclei except for Coulomb effects
- Charge symmetry of nuclear force

In  $^4_\Lambda\text{H}$  and  $^4_\Lambda\text{He}$

- Large difference in  $B_\Lambda$  and  $E_x(1^+)$
- $\Lambda p$  and  $\Lambda n$  interactions are different?
- Many-body effects with hyperon mixing?

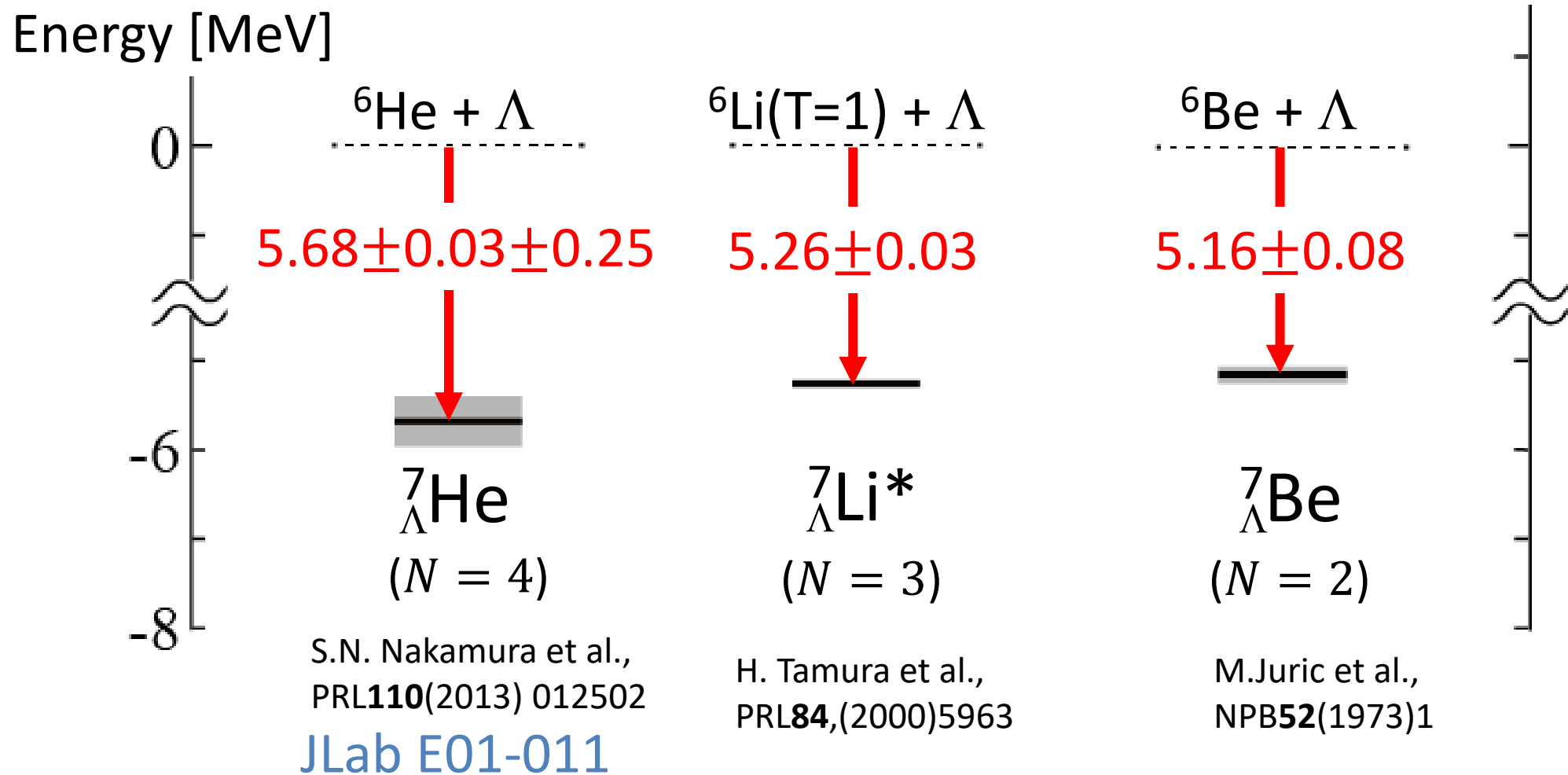
A.Esser et al., Phys. Rev. Lett.**114**,232501(2015)

T.O. Yamamoto, et al., PRL**115**, 222501('15)

# Related topic: charge symmetry breaking (CSB)

## ◆ CSB in $A=7$ hypernuclear isotriplet ( $T=1$ )

- Different trend from  $A=4$   $\Lambda$  hypernuclei: larger  $B_\Lambda$  as  $N$  increases



# What will happen if a $\Lambda$ particle is coupled to nuclei ?

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- **Dynamical changes of nuclear structure**

- Changes of cluster structure
- Deformation changes

- **Sensitivity of the  $\Lambda$  binding energy  $B_\Lambda$  on nuclear structure**

- Dependence of  $B_\Lambda$  on nuclear deformation
- Mass number  $A$  dependence & many-body force effects

- **Coupling of  $\Lambda$  particle in  $p$  orbit to clustering/deformed core nuclei**

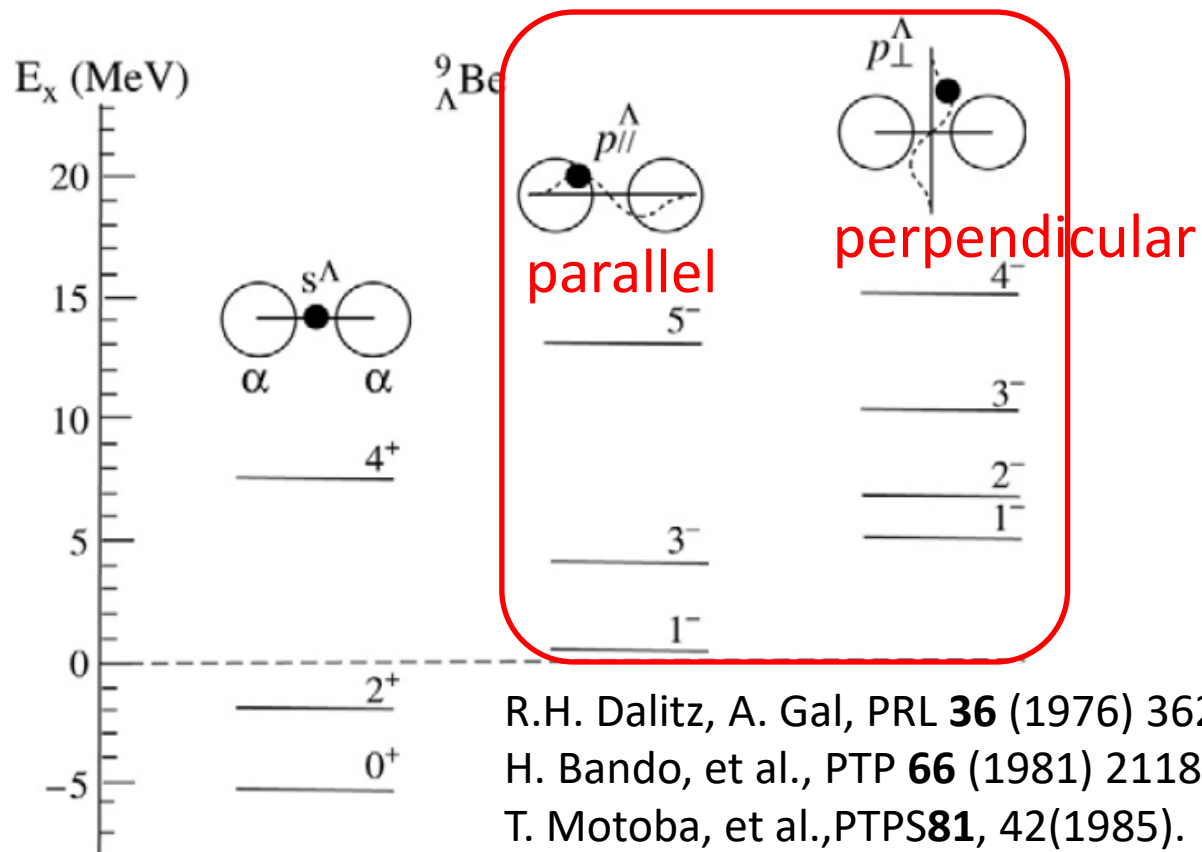
- Genuine hypernuclear states
- Possibility to probe nuclear deformation using  $\Lambda$  particle

# Genuine hypernuclear states in ${}^9_{\Lambda}\text{Be}$

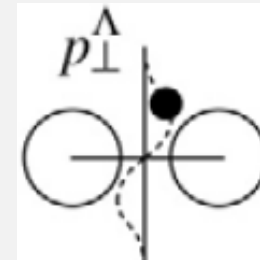
${}^9_{\Lambda}\text{Be}$ : axially symmetric  $2\alpha$  clustering

Two rotational bands as  $p$ -states  $\left\{ \begin{array}{l} \bullet \text{ Anisotropic } p \text{ orbit of } \Lambda \text{ hyperon} \\ \bullet \text{ Axial symmetry of } 2\alpha \text{ clustering} \end{array} \right.$

$\rightarrow$   $p$ -orbit parallel to/perpendicular to the  $2\alpha$  clustering

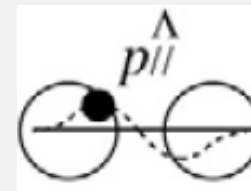


$p$  states in  ${}^9_{\Lambda}\text{Be}$



$\doteq$   ${}^9\text{Be}$

“ ${}^9\text{Be}$  analog states”



Forbidden for n in  ${}^9\text{Be}$   
due to Pauli principle

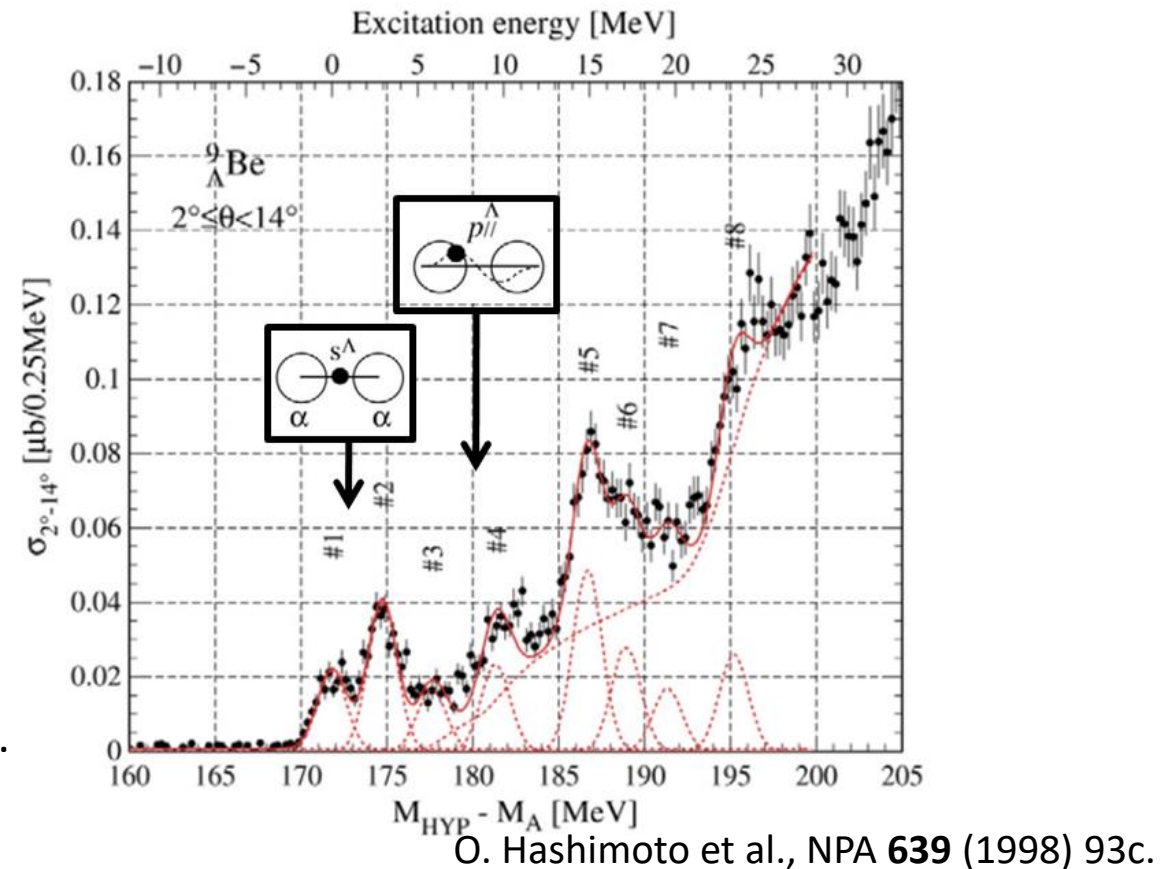
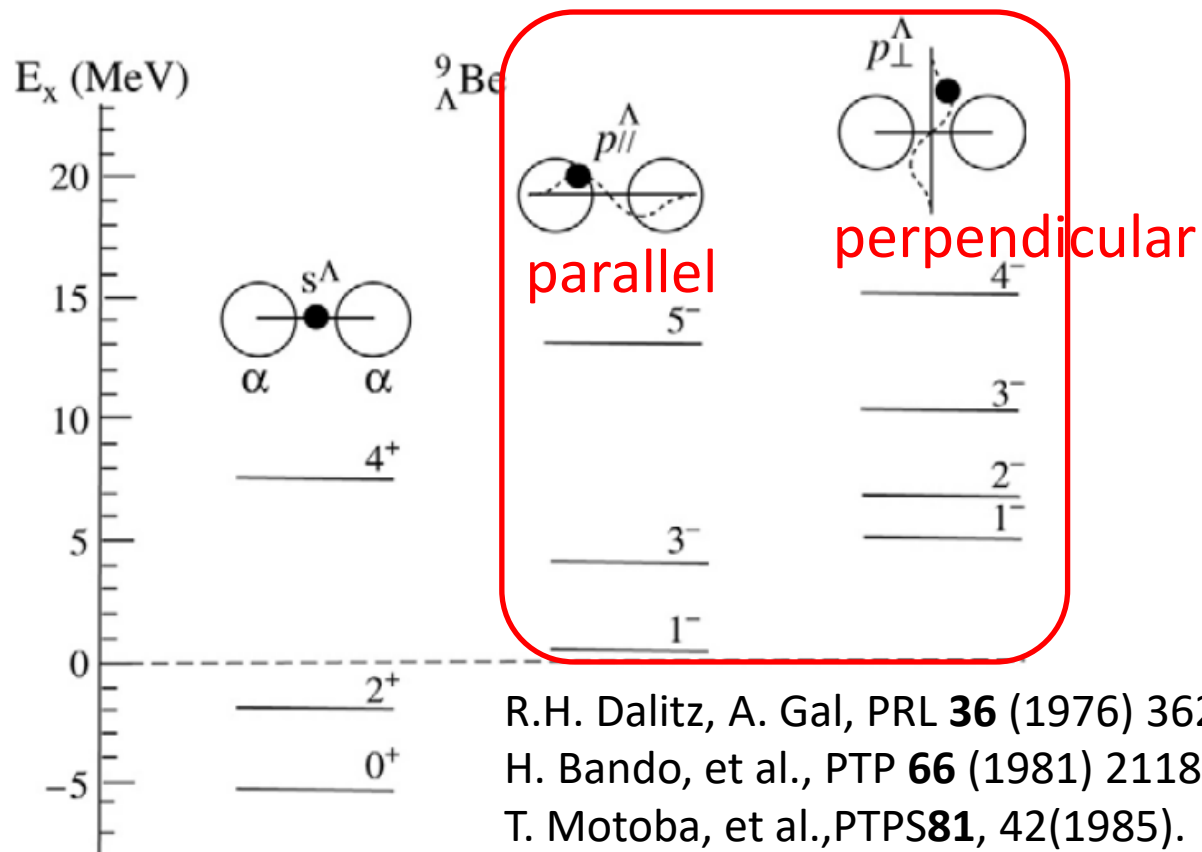
“genuine hypernuclear states”

# Genuine hypernuclear states in ${}^9_{\Lambda}\text{Be}$

${}^9_{\Lambda}\text{Be}$ : axially symmetric  $2\alpha$  clustering

Two rotational bands as  $p$ -states  $\left\{ \begin{array}{l} \bullet \text{ Anisotropic } p \text{ orbit of } \Lambda \text{ hyperon} \\ \bullet \text{ Axial symmetry of } 2\alpha \text{ clustering} \end{array} \right.$

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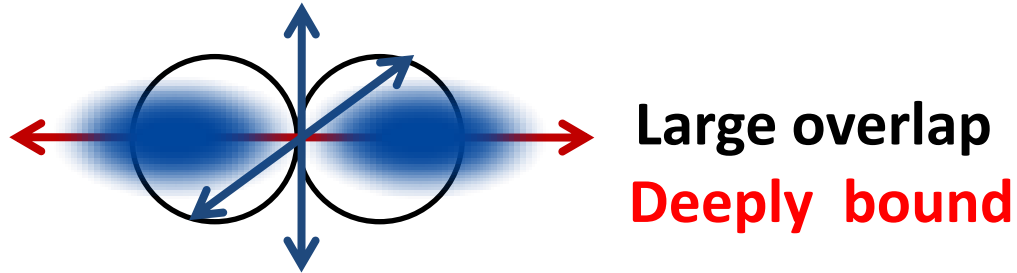


# Split of $p$ -state in ${}^9_{\Lambda}\text{Be}$

${}^9_{\Lambda}\text{Be}$ : axially symmetric  $2\alpha$  clustering

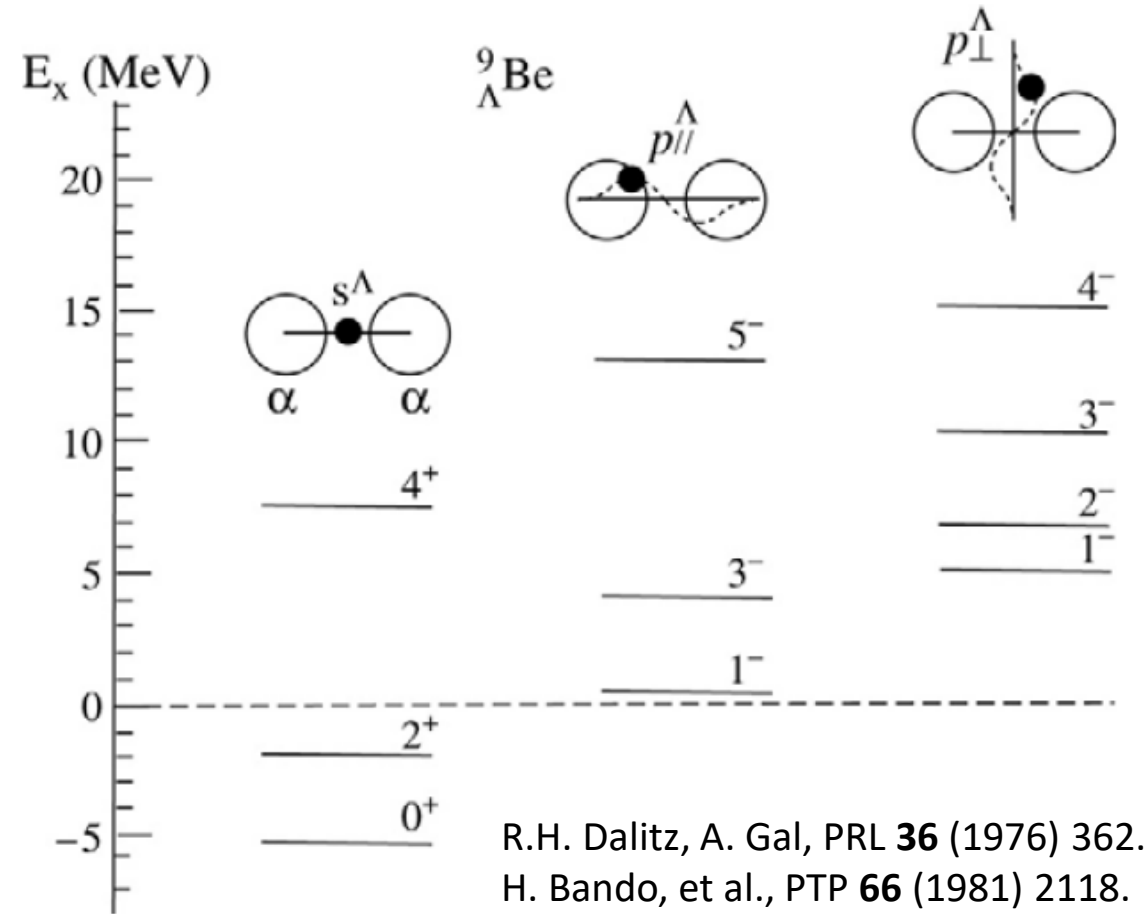
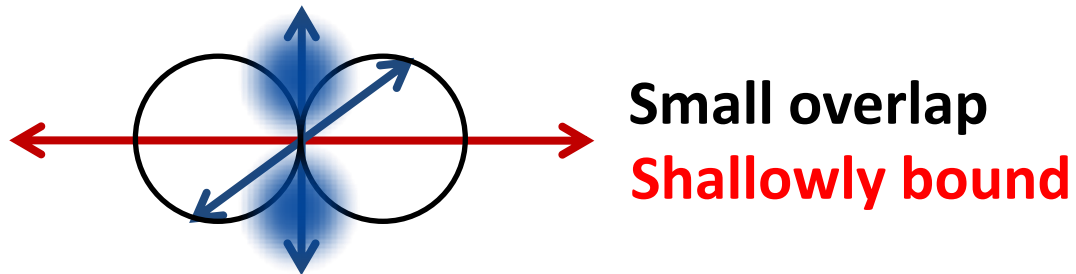
Genuine hypernuclear states:

$p$  orbit parallel to  $2\alpha$  (long axis)



${}^9\text{Be}$  analog states:

$p$  orbit perpendicular to  $2\alpha$  (short axes)



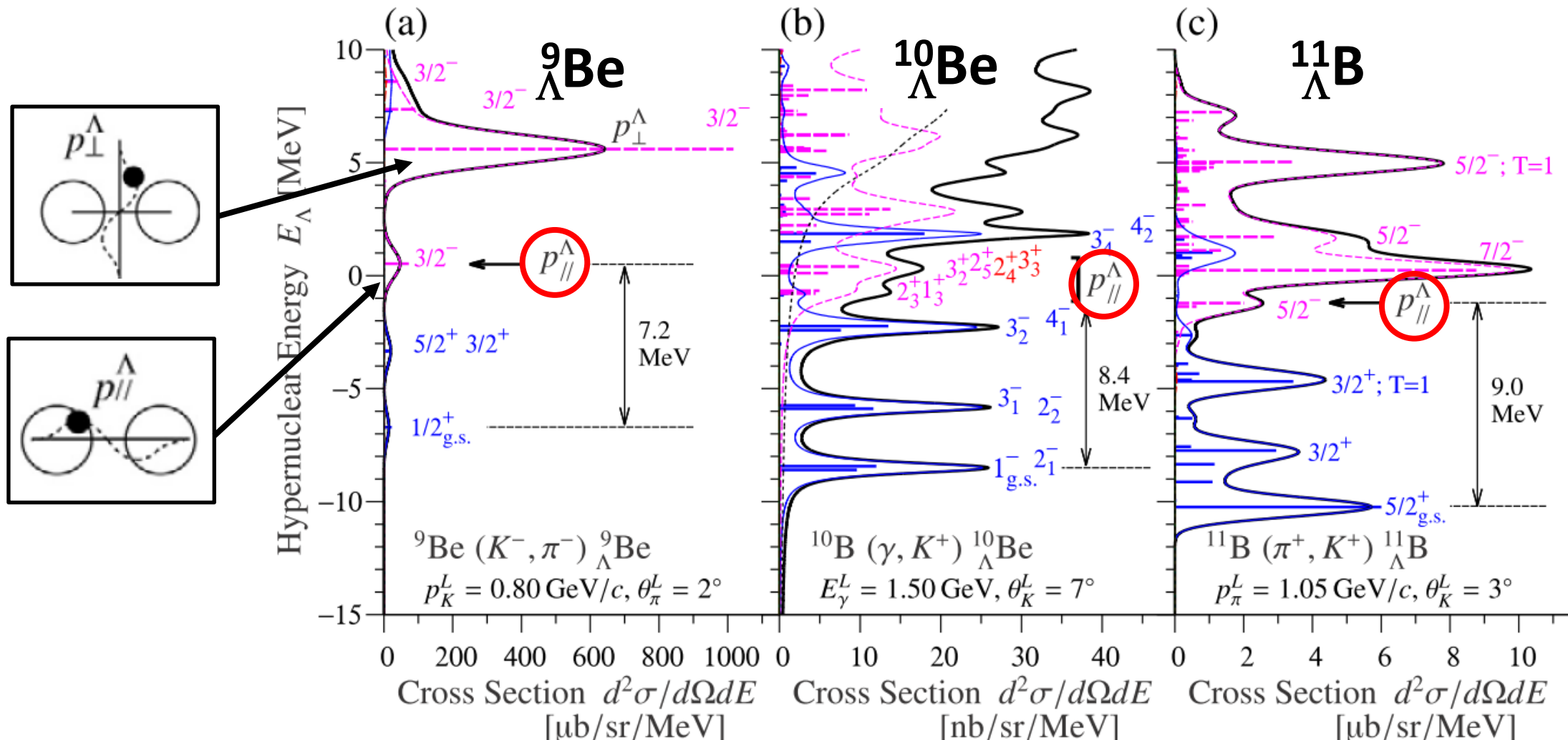
R.H. Dalitz, A. Gal, PRL **36** (1976) 362.  
H. Bando, et al., PTP **66** (1981) 2118.  
T. Motoba, et al.,PTPS**81**, 42(1985).

**$p$ -states splits into 2 bands depending on the direction of  $p$ -orbits**

# Genuine hypernuclear states in the other hypernuclei

Genuine hypernuclear states are predicted not only in  ${}^9_{\Lambda}\text{Be}$  but  ${}^{10}_{\Lambda}\text{Be}$  &  ${}^{11}_{\Lambda}\text{Be}$

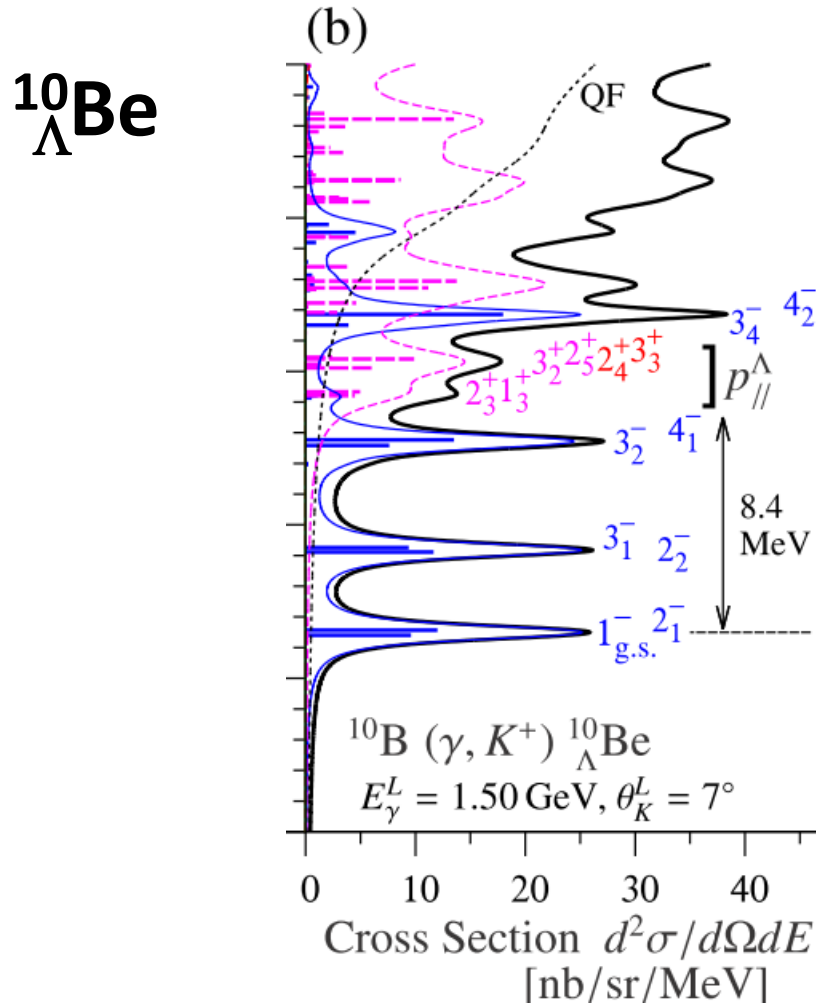
Shell model + DWIA calc. by Umeya et al., EPJ Web of Conference **271**, 01010(2022)



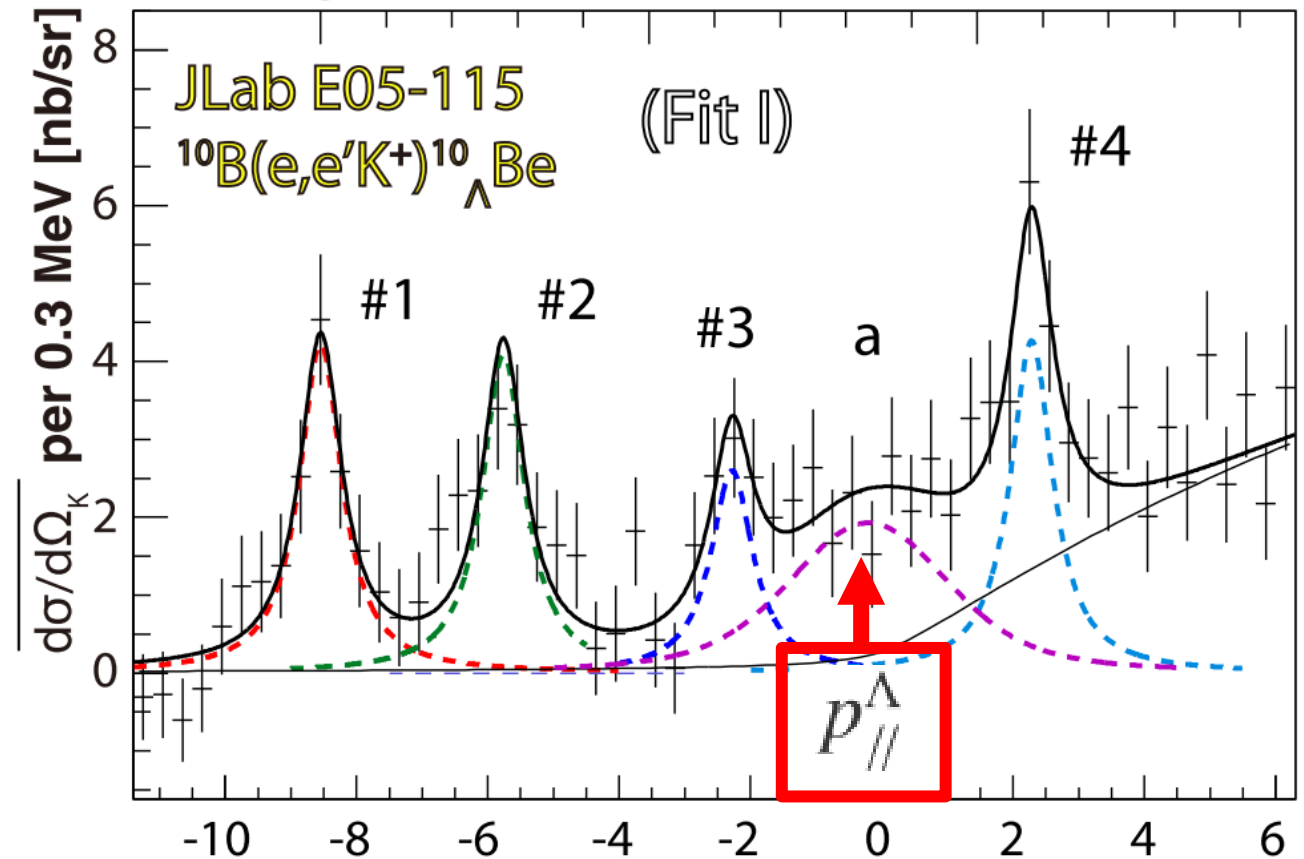


# Genuine hypernuclear states in the other hypernuclei

Shell model + DWIA by Umeya



T. Gogami et al., PRC**93**, 034314(2016)



Bump structure “a” can be interpreted as genuine hypernuclear states

**Coupling of  $\Lambda$  in  $p$  orbit to symmetric cluster structure causes splitting of  $p$  states**



# Coupling of $\Lambda$ in $p$ orbit to triaxially deformed nuclei

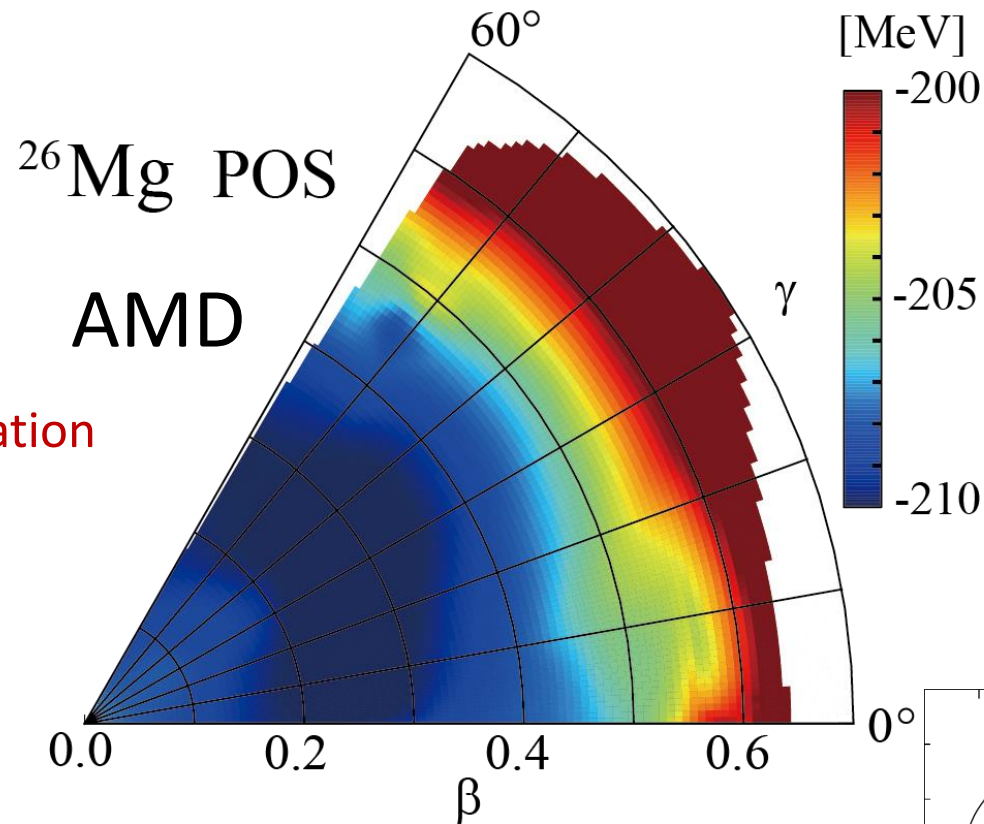
- $^{26}\text{Mg}$
- Shell gap in Nilsson diagram:  $Z=12$  (prolate) vs.  $N=14$  (oblate) → **triaxial**
  - $\beta, \gamma$ -softness is discussed by several authors

Terasaki et al. NPA**621**(1997)

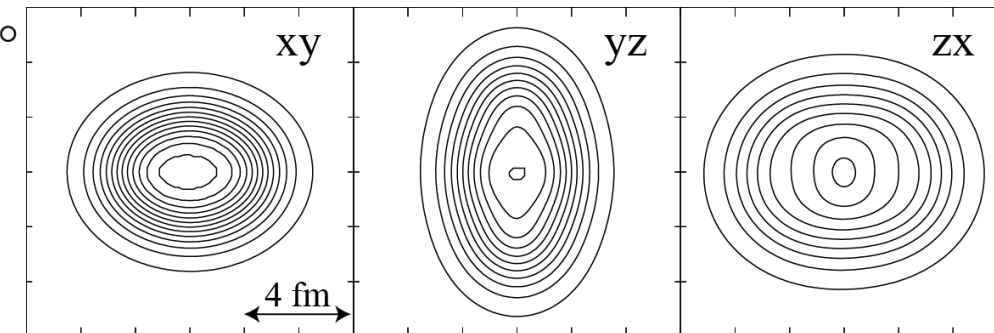
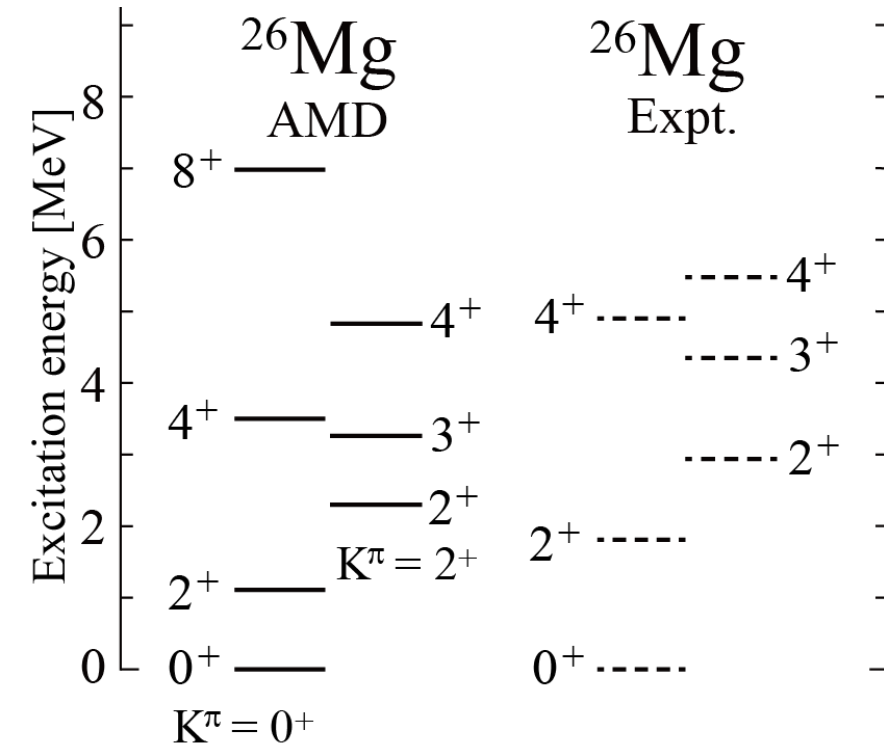
Rodriguez-Guzman et al. NPA**709** (2002)

Peru et al PRC**77** (2008)

Hinohara, Kanada-En'yo PRC**83** (2011)



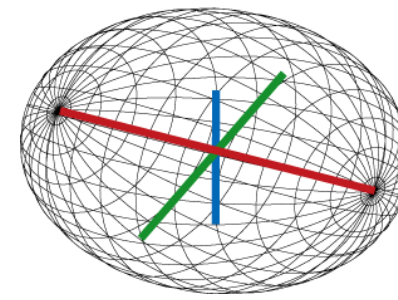
$\beta$ : degree of deformation  
 $\gamma$ : (tri)axiality



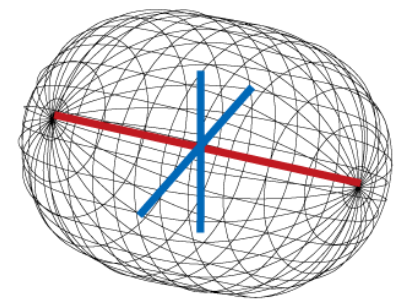
**What will happen if  $\Lambda$  in  $p$  orbit is coupled to triaxially deformed  $^{26}\text{Mg}$  ?**

# Triaxial deformation

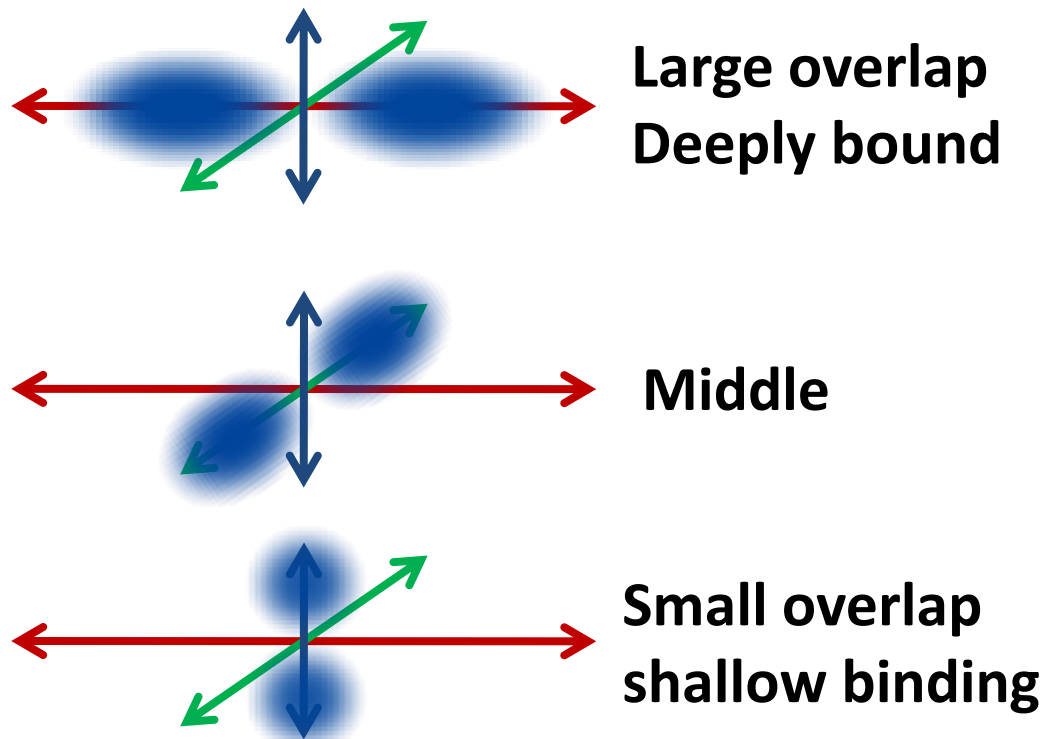
If nucleus is triaxially deformed,  
*p*-states can split into 3 different state



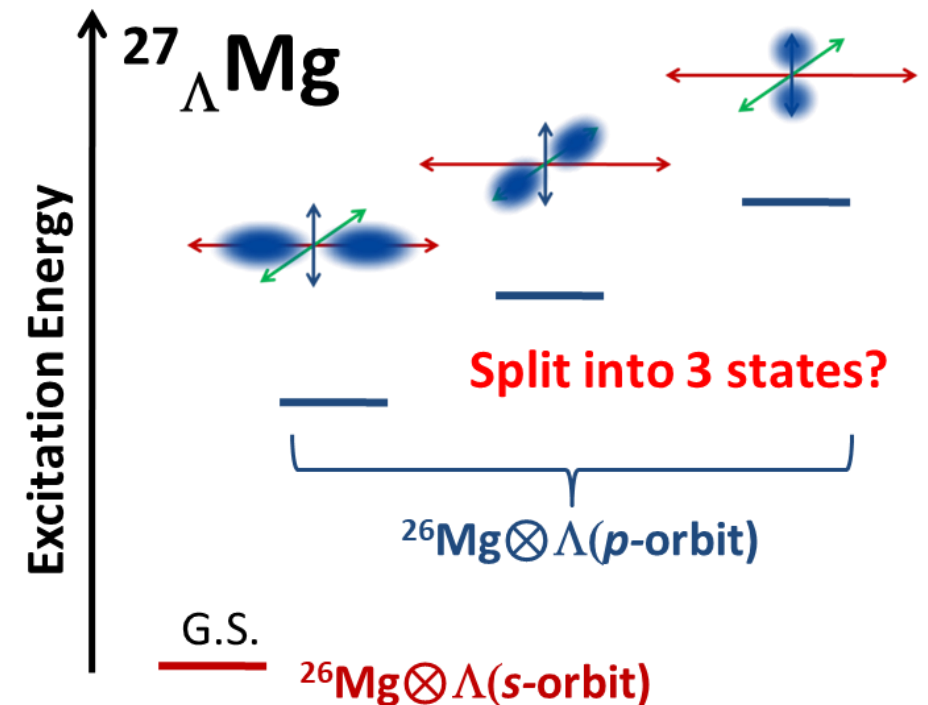
Triaxial deformation



Prolate deformation



Candidate: Mg hypernuclei

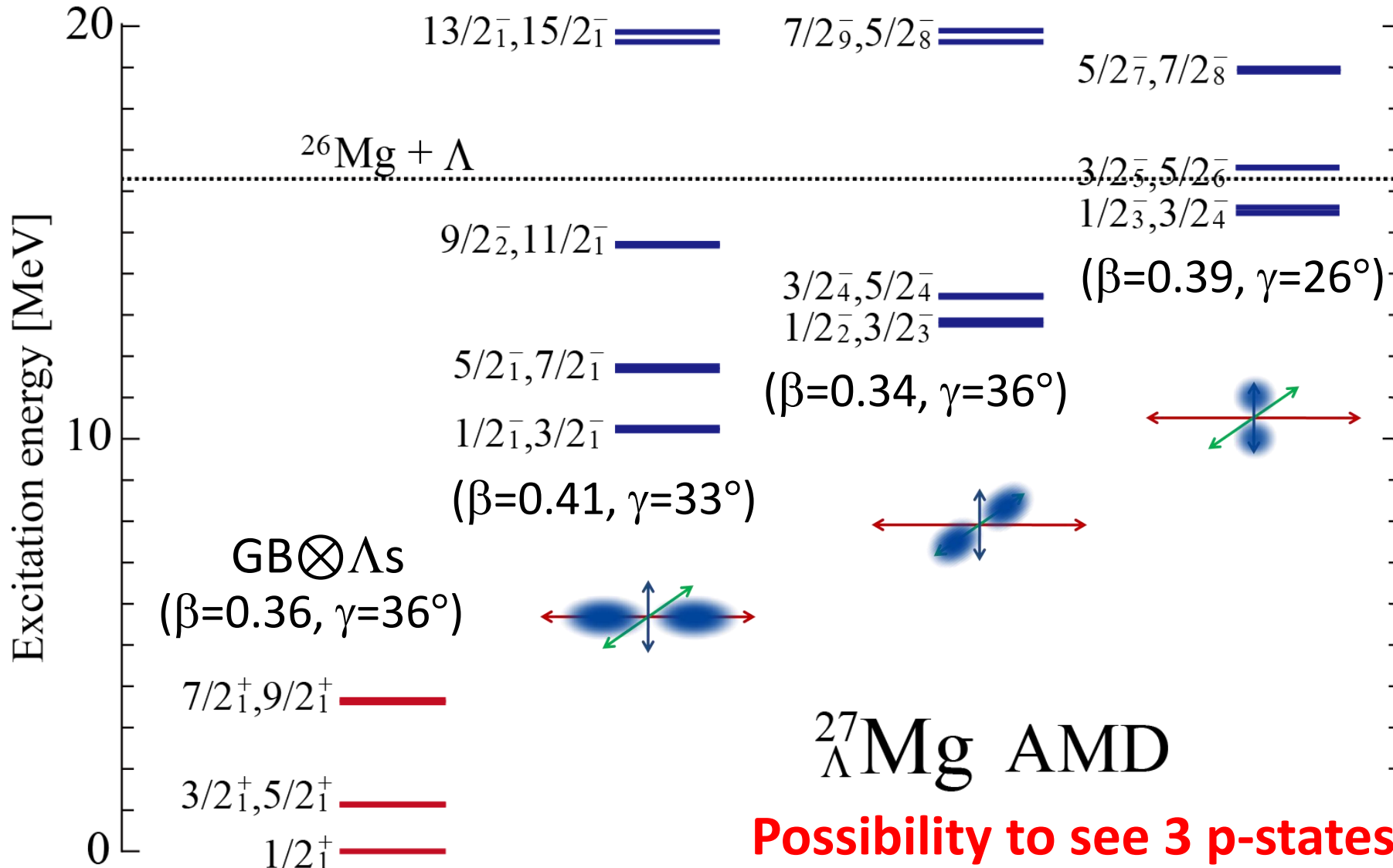


Observing the 3 different *p*-states is strong evidence of triaxial deformation

# Results: $^{27}_{\Lambda}\text{Mg}$

M.I., et al., PRC**87**, 021304(R) (2013)  
and calculation ongoing

- 3 bands are obtained by  $\Lambda$  in  $p$ -orbit → Splitting of the  $p$  states



Possibility to see 3 p-states by  $^{27}\text{Al}(e, e'K+)^{27}_{\Lambda}\text{Mg}$

# Summary

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In  $\Lambda$  hypernuclei, by the coupling of  $\Lambda$ , we can expect following phenomena:

## Dynamical changes of nuclear structure

- Changes of cluster structure
  - Deformation changes
- **Structure changes depending on core structure**

## Sensitivity of $B_\Lambda$ on nuclear structure

- $B_\Lambda$  depending on nuclear deformation
  - Both hyperonic interactions and structure of hypernuclei are important in  $A$ -dep. of  $B_\Lambda$
- **Many-body effects in  $\Lambda$  hypernuclei**      future exp. at JLab (E12-15-008)

## Coupling of $\Lambda$ particle to clustering/deformed core nuclei

- Genuine hypernuclear states not only in  $^9_\Lambda\text{Be}$  but also  $^{10}_\Lambda\text{Be}$       JLab E05-115
- Possibility to probe nuclear triaxial deformation using  $\Lambda$  in  $^{27}_\Lambda\text{Mg}$