



Polarized electron/positron beams for rare isotope studies at FRIB

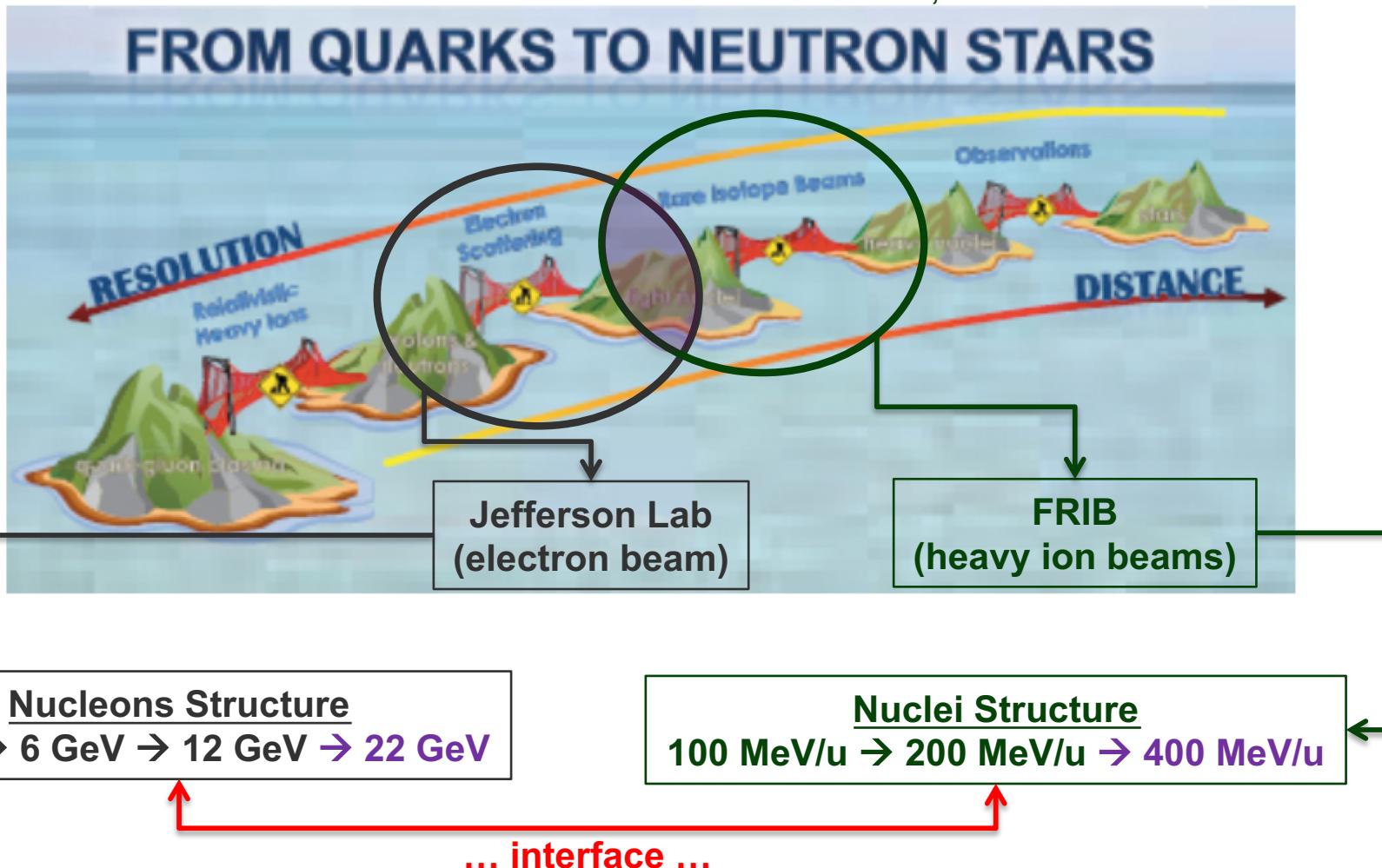
Paul Guèye, Ambar Rodriguez Alicea, Peter Ostroumov, Jeremy Rebenstock, Maya Wallach

MICHIGAN STATE
UNIVERSITY

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ENERGY | Office of
Science

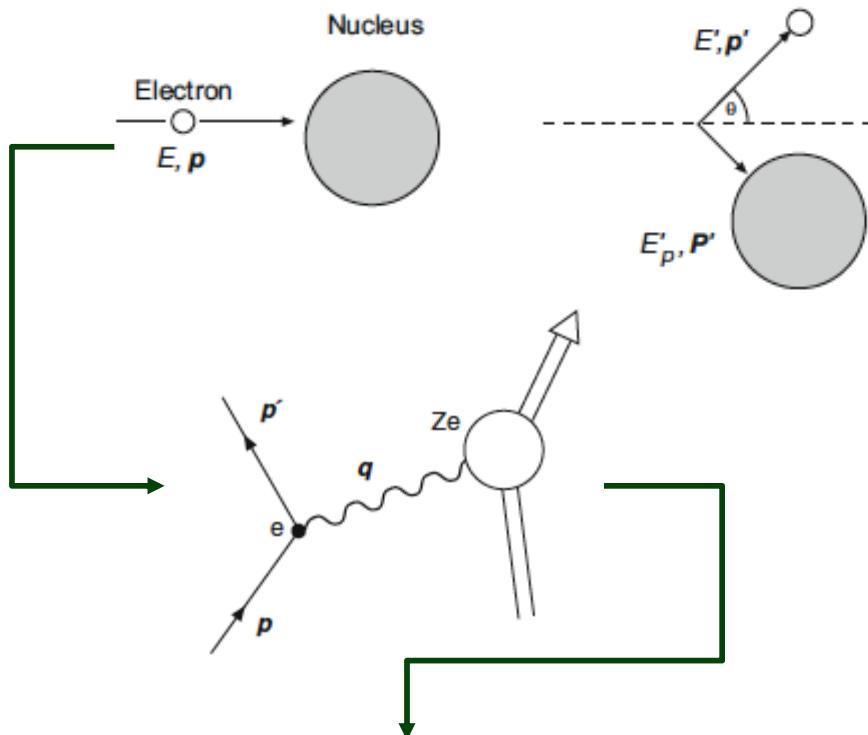
Two Probes + One Target = Scaling

The National Academies of Science, 2013



Peaking Inside Closed Rooms

Bogdan Povh, Klaus Rith, Christoph Scholz, Frank Zetsche • Werner Rodejohann
Particles and Nuclei: An Introduction to the Physical Concepts



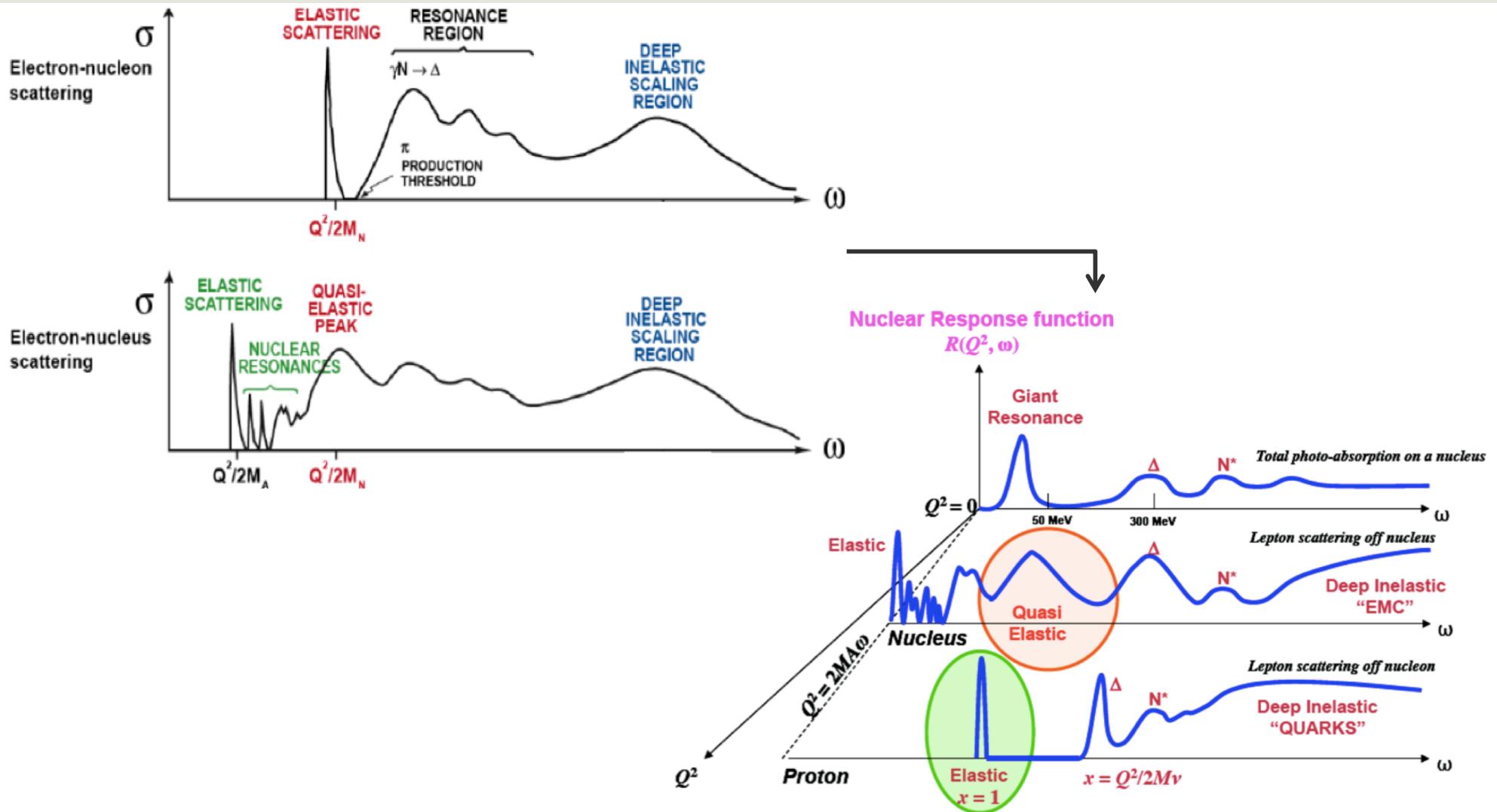
“seeing” inside a closed room



https://www.123rf.com/photo_38569847_stock-vector-house-cartoon-interior-cartoon-living-rooms-with-furniture-.html

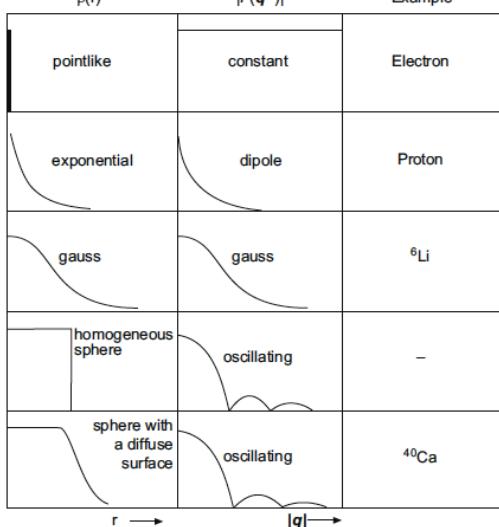
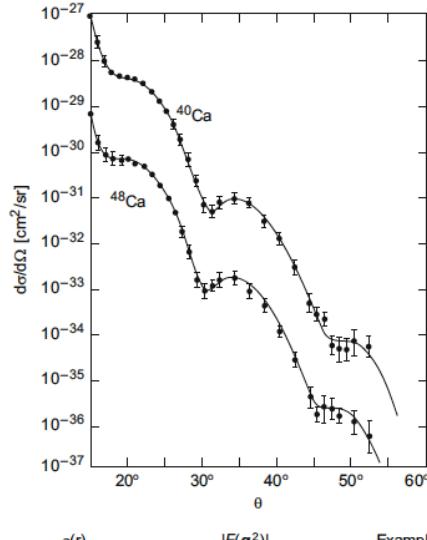
$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{exp.}} = \left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}}^* \cdot |F(q^2)|^2$$

Electron Scattering Experiments



Nuclear Radii

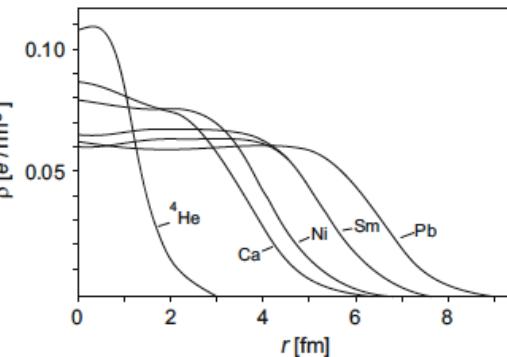
Bogdan Povh, Klaus Rith, Christoph Scholz, Frank Zetsche • Werner Rodejohann
Particles and Nuclei: An Introduction to the Physical Concepts



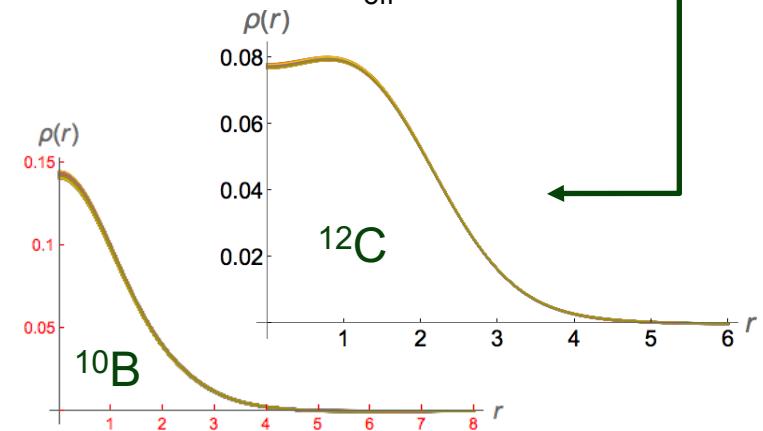
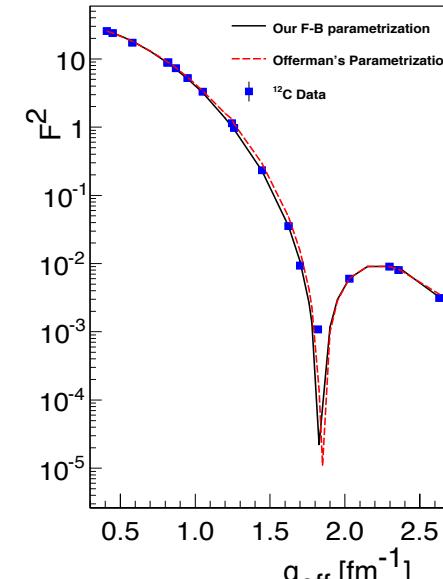
$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega} \right)_{Mott} | F(Q^2) |^2$$

$$F_p(q^2) = \frac{1}{4\pi} \int d^3r j_0(qr) \rho_p(r)$$

$$\begin{aligned} ZF_p &= 4\pi \int_0^\infty \rho_p r^2 dr \\ &= \sum_{\nu=1}^{\infty} (-1)^{\nu+1} \frac{4\pi R_p}{q_\nu^2} a_\nu \end{aligned}$$



A. Kabir, PhD Thesis (2019)
Low Energy Deuteron Experiment, JLab



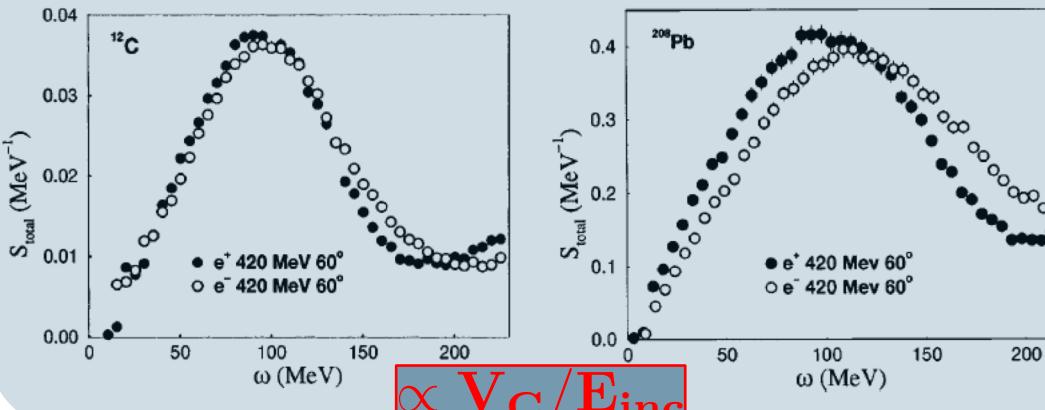
Facility for Rare Isotope Beams

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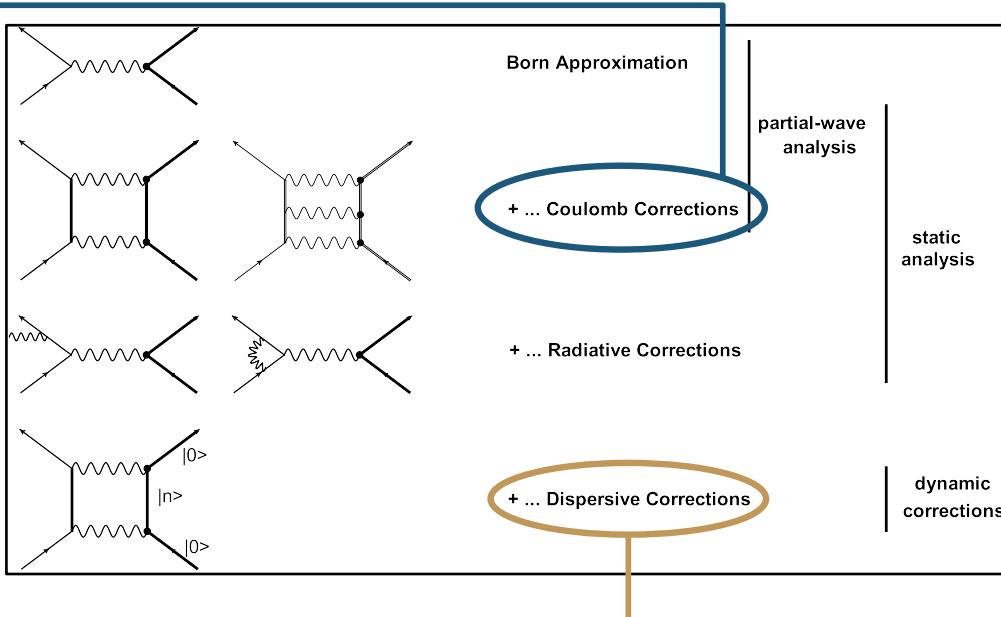
What About Higher Orders Corrections?

P. Guèye et al., Phys. Rev., C60 , 044308 (1999)

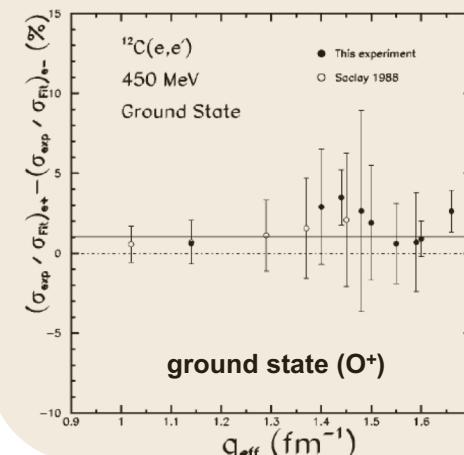


Effective Momentum Approximation (EMA)

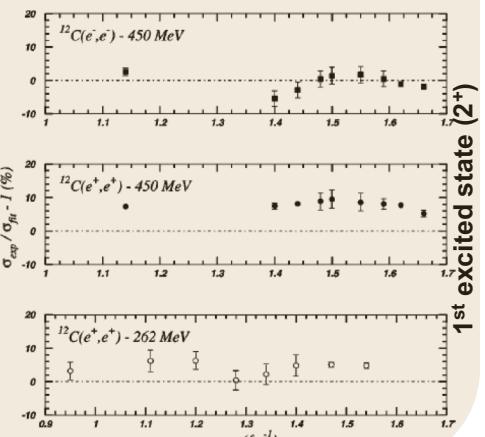
Nucleus	$\langle r^2 \rangle_{\text{exp}}^{1/2}$ (fm)	$ V_C(0) $ (MeV)	$ V_C $ (MeV)	$ V_C _{\text{fit}}$ (MeV)
^{12}C	2.464	4.6	3.3	3.1 ± 0.25
^{40}Ca	3.450	10.5	7.9	7.4 ± 0.6
^{48}Ca	3.451	10.4	7.9	7.4 ± 0.6
^{56}Fe	3.714	12.5	9.5	8.9 ± 0.7
^{90}Zr	4.258	16.7	12.8	11.9 ± 0.9
^{154}Gd	5.124	21.8	16.9	15.9 ± 1.2
^{208}Pb	5.503	25.9	20.1	18.9 ± 1.5



P. Guèye et al., Phys. Rev., C57 , 2107 (1998)



P. Guèye et al., Phys. Rev., C63, 051303(R) (2001)



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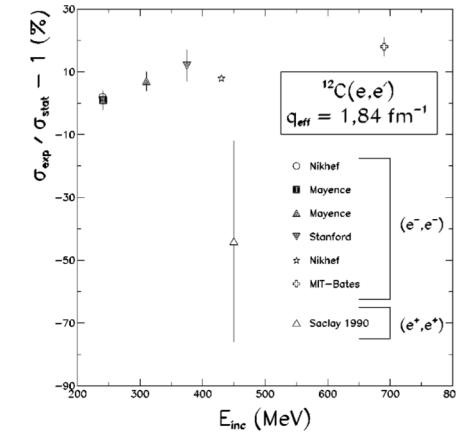
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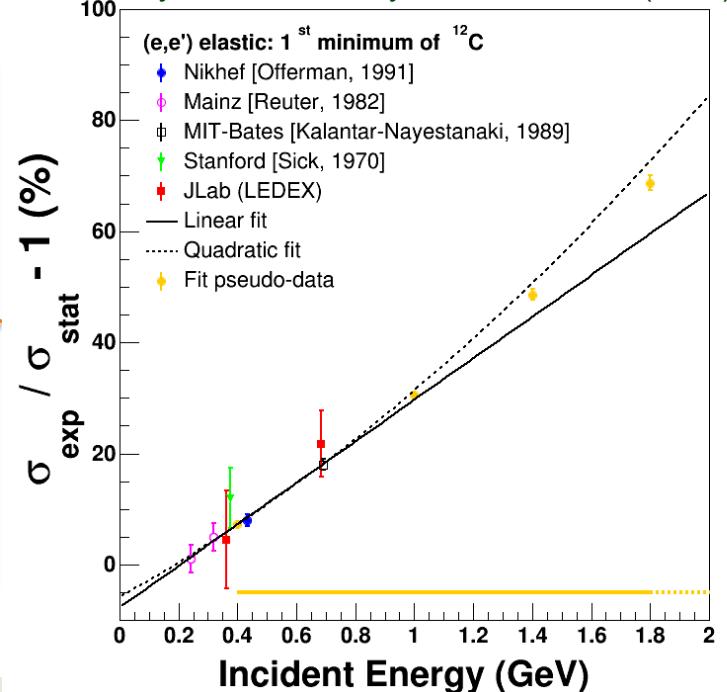
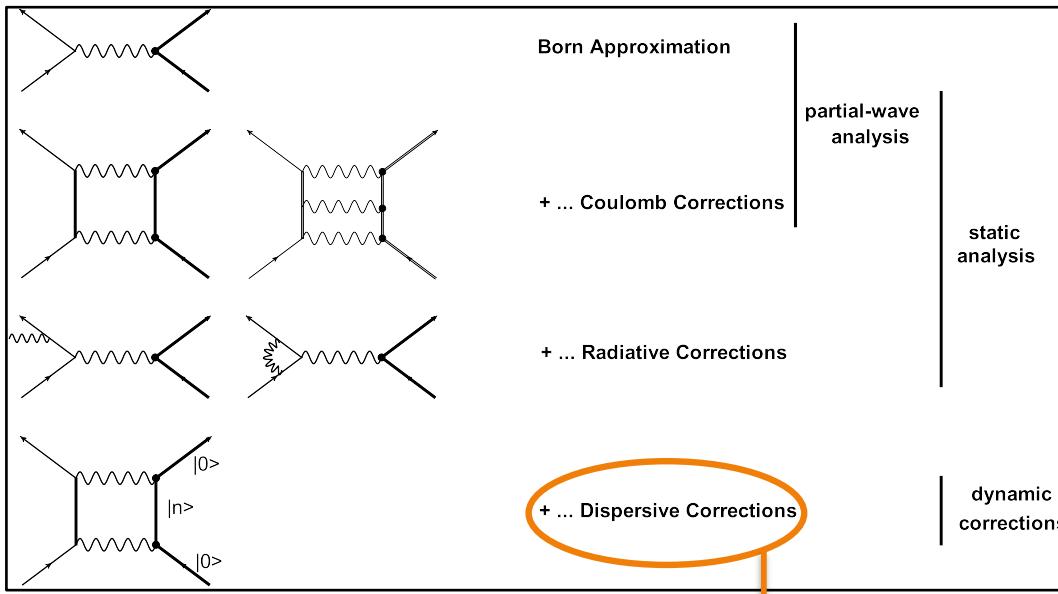
Are Dispersive Effects really negligible?

P. Gueye et al., PRC57, 5 (1998)

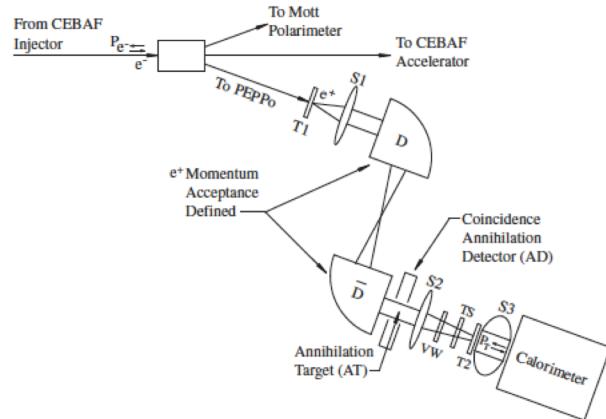
- **Coulomb Corrections**
 - ✓ PWBA → DWBA
 - ✓ Effective Momentum Approximation
- **Dispersive Effects**
 - ✓ Energy Dependence?



P. Gueye et al., Eur. Phys. Jour. A56:126 (2020)

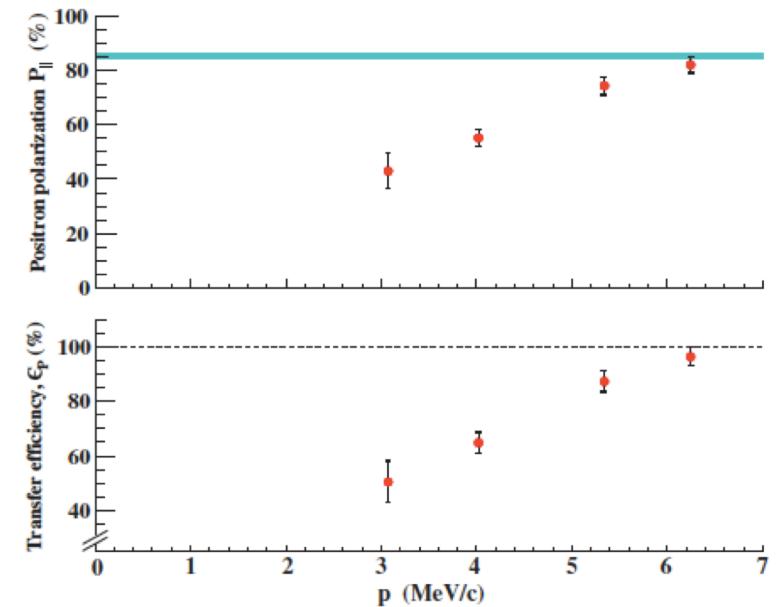
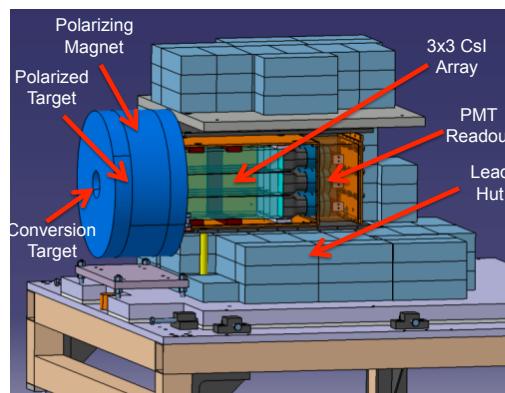
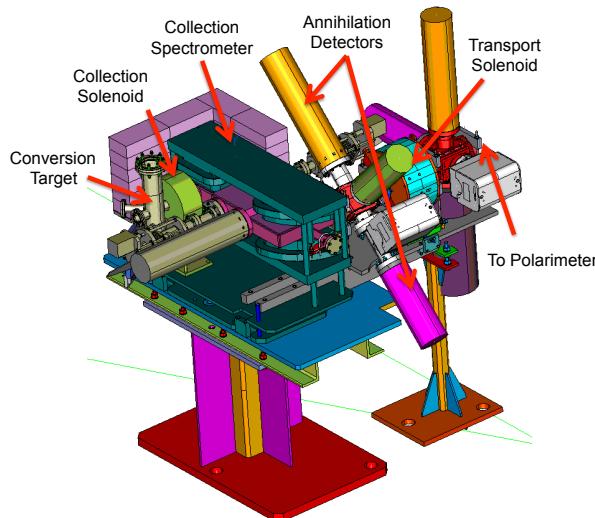
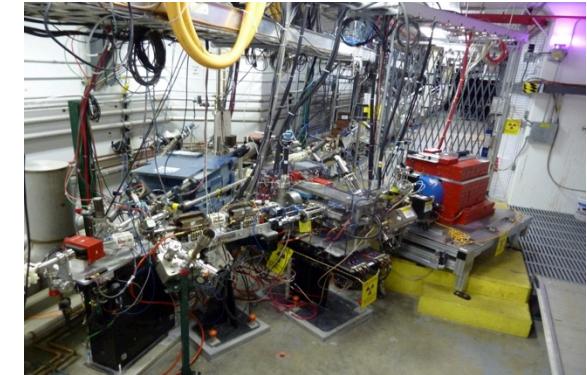


Polarized Positron Beams – 20 years later! (... possible scheme for the EIC)



Polarized Electrons for Polarized Positrons
D. Abbott *et al.*, PRL 116, 214801 (2016)

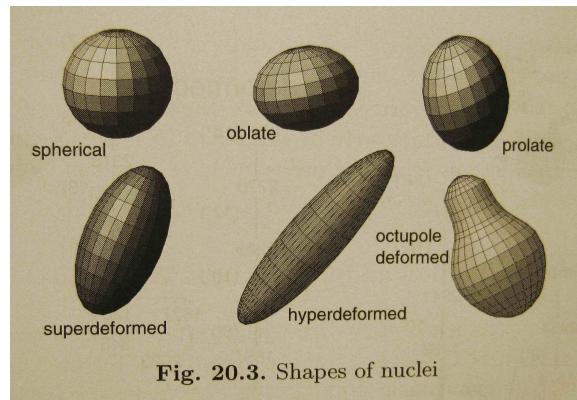
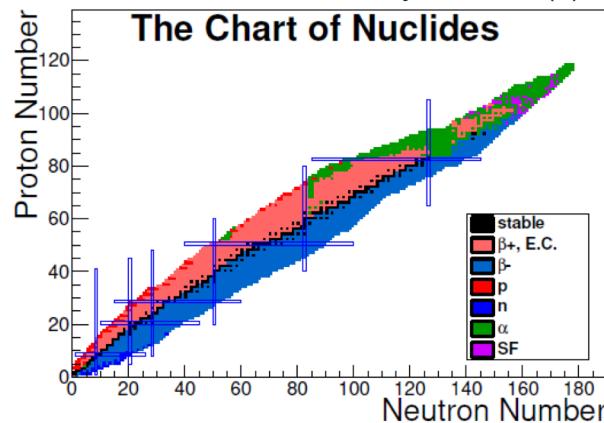
- Experiment in the CEBAF injector
- Highly polarized positrons
- 80% @ 6.5 MeV
- R&D for EIC
- Last PhD @ HU (A. Adeyemi, 2016)



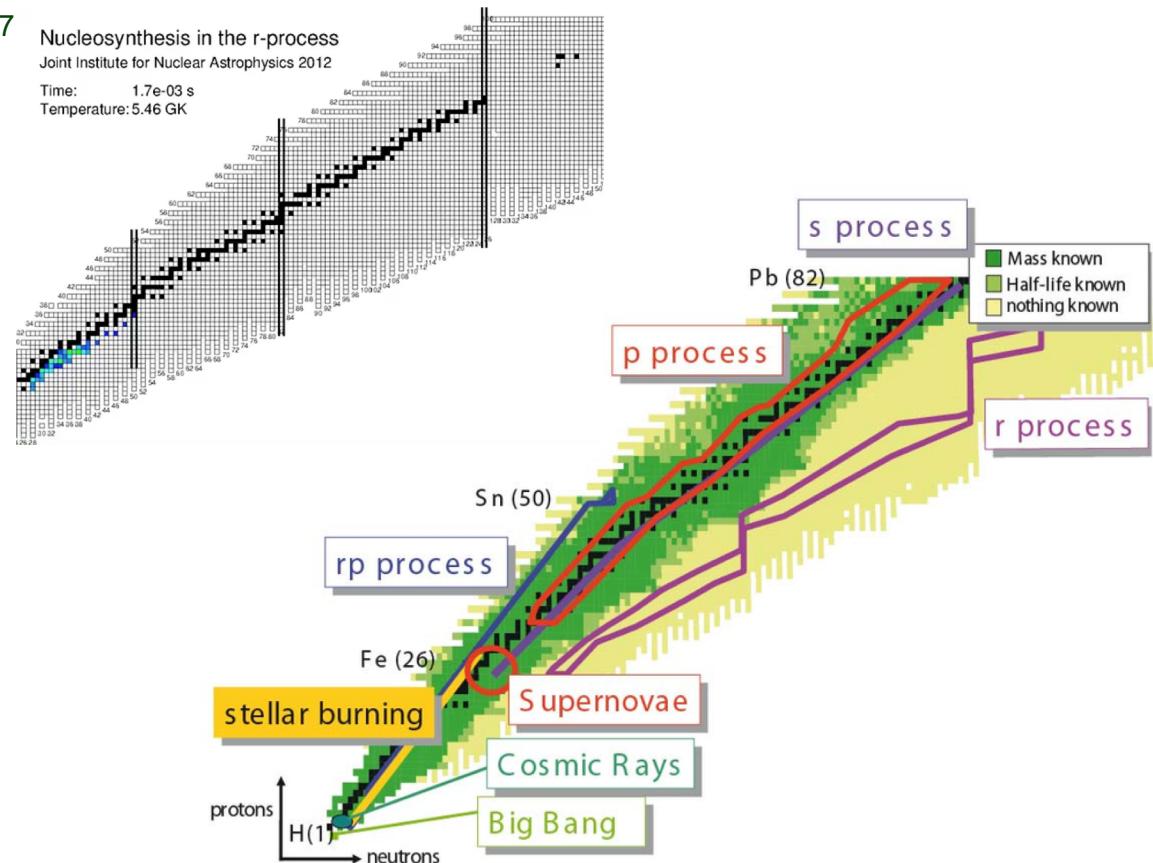
Nuclear Chart & Nuclear Science

- 288 stables ($T_{1/2} >$ life of the solar system)
- 3,308 different isotopes have been discovered (status in 2019)

Data from G. Audi, et al., Chinese Physics C, 41(3),030001, 2017



<https://web2.ph.utexas.edu/~coker2/index.files/rotation2.htm>



K.Blaum et al., Contemporary Physics, 51, 149-175 2010

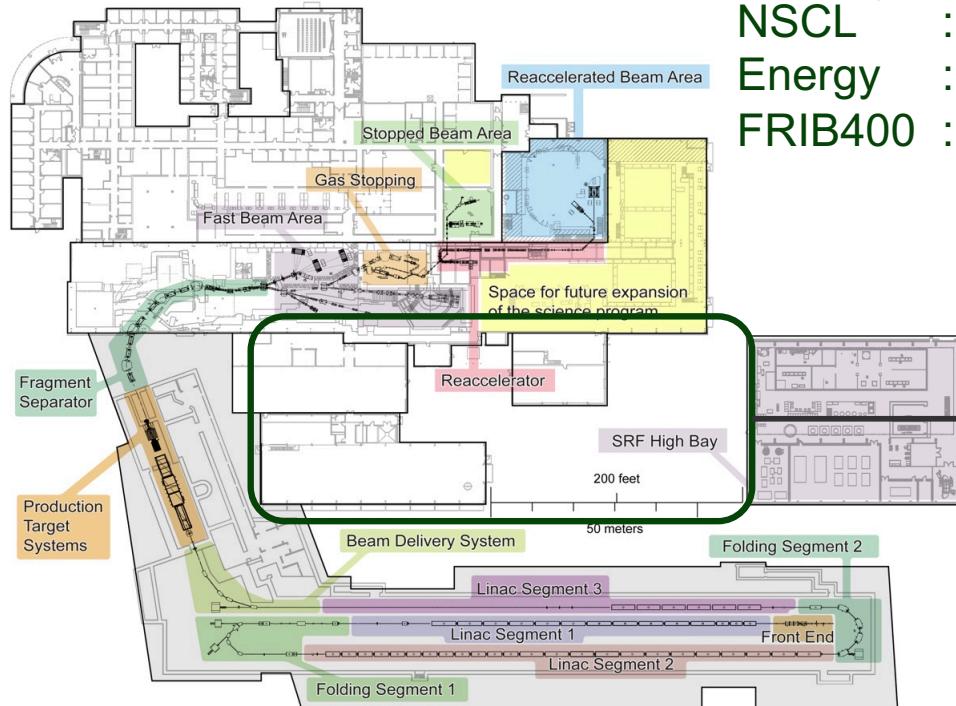
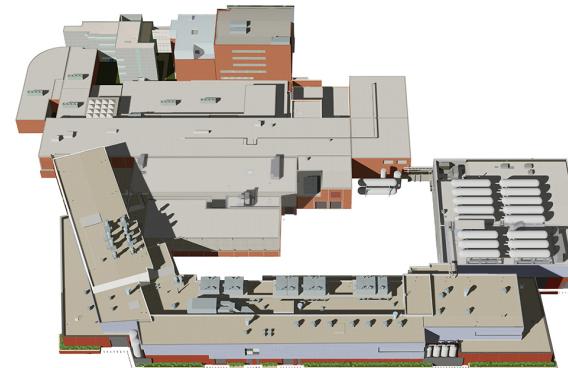
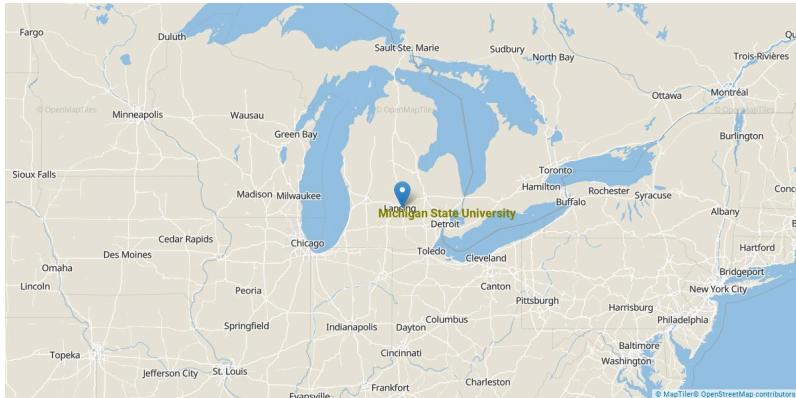


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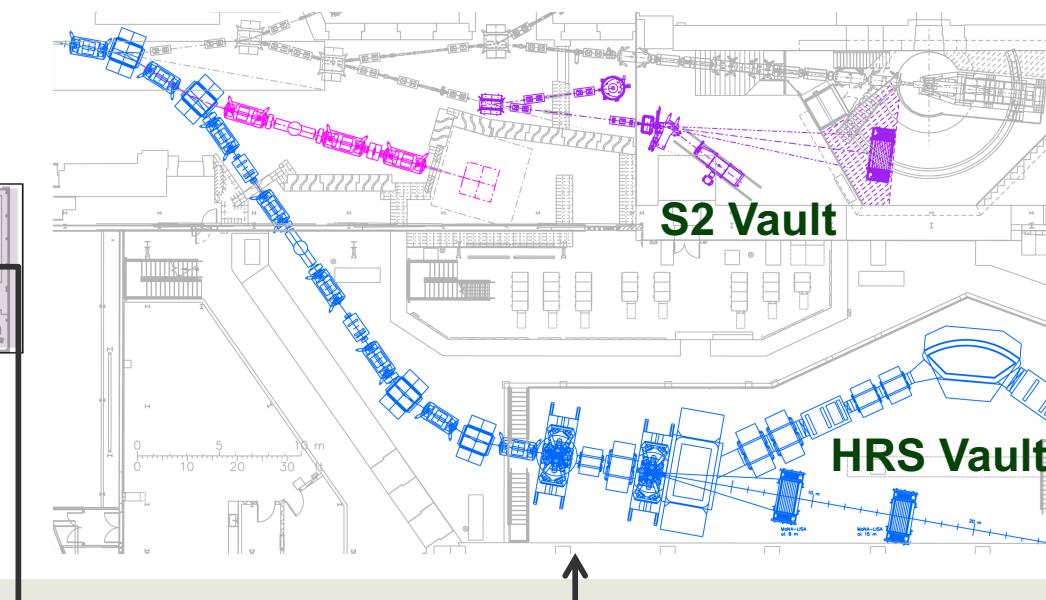
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Facility for Rare Isotope Beams

(www.frib.msu.edu; start: May 10, 2022)



NSCL : ~100 MeV/u
Energy : ~200 MeV/u
FRIB400 : ~400 MeV/u



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Invariant Mass Technique (MoNA Collaboration – S2 Vault)

$$\left. \begin{array}{l} E_U = E_F + E_n \\ \vec{P}_U = \vec{P}_F + \vec{P}_n \end{array} \right\} \rightarrow M_U = \sqrt{E_U^2 - \vec{P}_U^2}$$

$$E_{\text{decay}} = \sqrt{M_F^2 + M_n^2 + 2(E_F E_n - \vec{P}_F \vec{P}_n)} - (M_F + M_n)$$

electronics racks

MoNA-LISA neutron detector array

charged particle detector chamber

Sweeper dipole magnet

magnetic field shield

CAESAR gamma detector array beam tracking detector chambers

neutrons

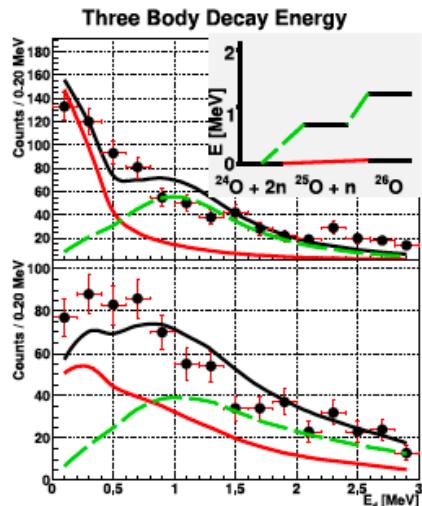
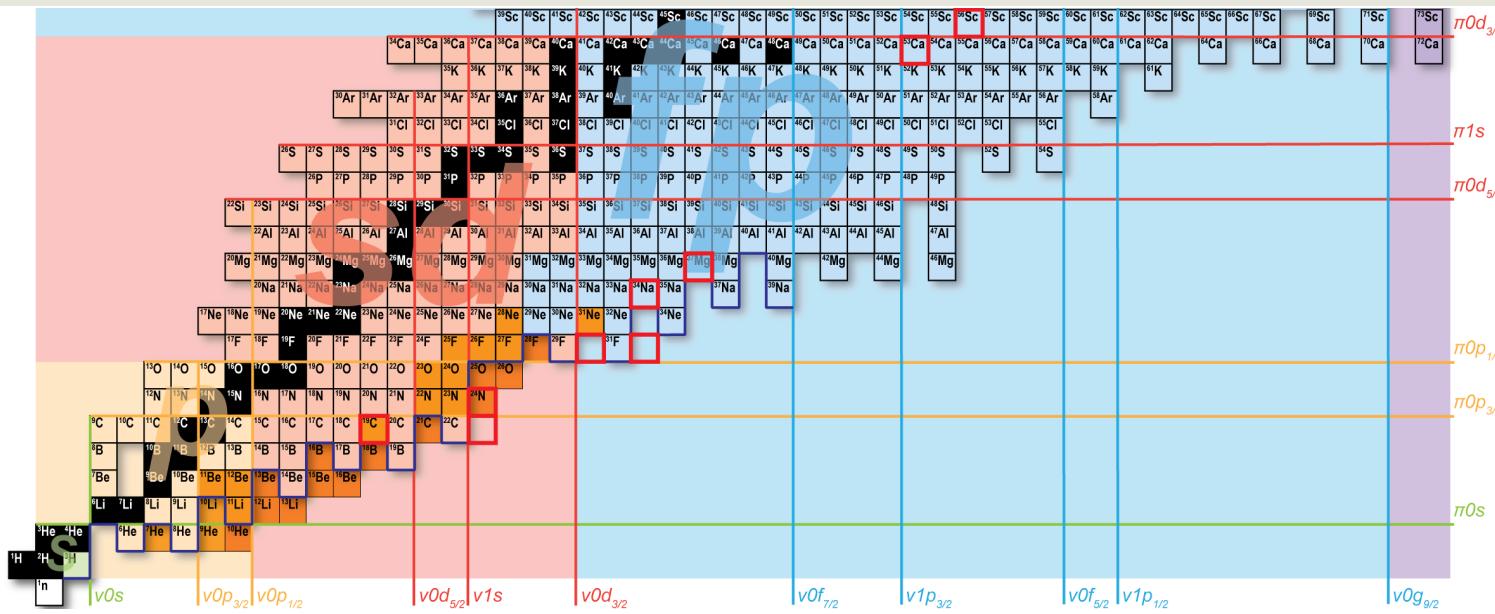
charged fragments

beam

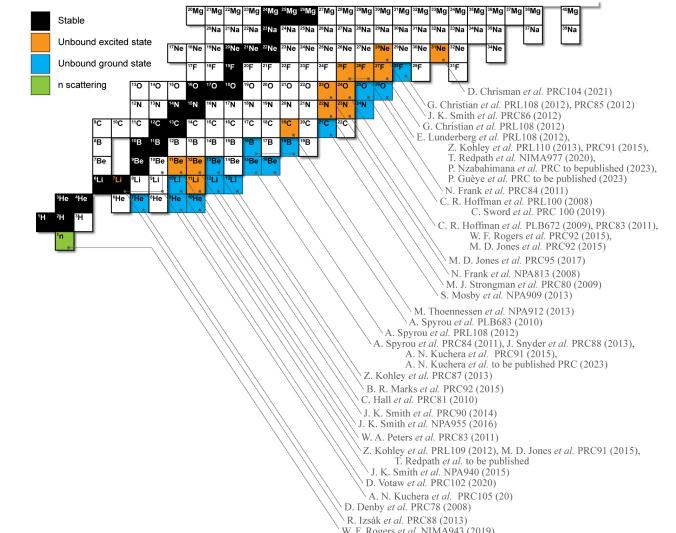
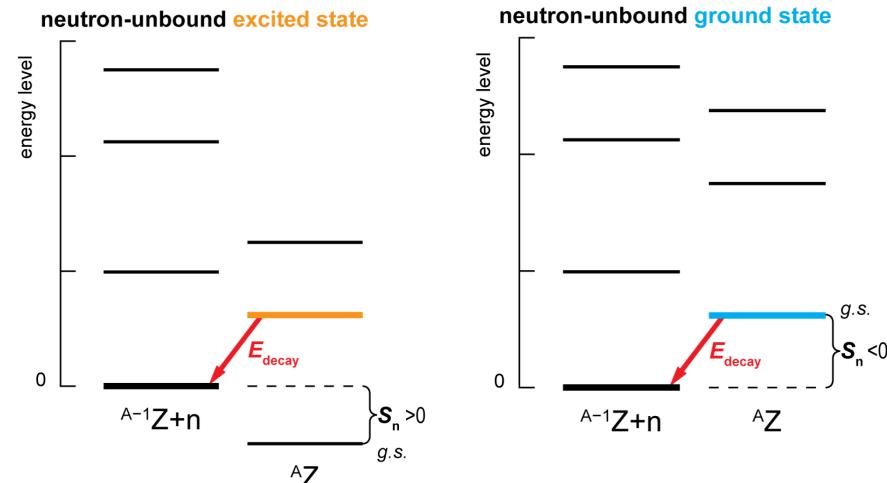
breakup target



Neutron Unbound States



T. Redpath et al, NIMA, 97, 164284 (2020)



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Nuclear Tomography? What About Polarization?

Meson electro-production

$$\begin{aligned} \frac{d\sigma_v}{d\Omega_\eta} = & \frac{|\mathbf{k}|}{k_\gamma^{cm}} P_\alpha P_\beta \{ R_T^{\beta\alpha} + \varepsilon_L R_L^{\beta\alpha} \\ & + [2\varepsilon_L(1+\varepsilon)]^{1/2} ({}^cR_{TL}^{\beta\alpha} \cos \phi_\eta + {}^sR_{TL}^{\beta\alpha} \sin \phi_\eta) \\ & + \varepsilon ({}^cR_{TT}^{\beta\alpha} \cos 2\phi_\eta + {}^sR_{TT}^{\beta\alpha} \sin 2\phi_\eta) \\ & + h [2\varepsilon_L(1-\varepsilon)]^{1/2} ({}^cR_{TL'}^{\beta\alpha} \cos \phi_\eta + {}^sR_{TL'}^{\beta\alpha} \sin \phi_\eta) \\ & + h(1-\varepsilon^2)^{1/2} R_{TT'}^{\beta\alpha} \}, \end{aligned} \quad (12)$$

G. Knöchlein, D. Drechsel, L. Tiator
Z. Phys. A352, 327-343 (1995)

**3D nucleon tomography!!
(DVCs, parton distributions ...)**

Table 1. Polarization observables in pseudoscalar meson electroproduction. A star denotes a response function which does not vanish but is identical to another response function via a relation in App. A

		Target			Recoil			Target + Recoil								
β	$-$	$-$	$-$	x'	y'	z'	x'	x'	x'	y'	y'	y'	z'	z'	z'	
α	$-$	x	y	z	$-$	$-$	x	y	z	x	y	z	x	y	z	
T	R_T^{00}	0	R_T^{0y}	0	0	$R_T^{y'0}$	0	$R_T^{x'x}$	0	$R_T^{x'z}$	0	*	0	$R_T^{z'x}$	0	$R_T^{z'z}$
L	R_L	0	R_L^{0y}	0	0	*	0	$R_L^{x'x}$	0	$R_L^{x'z}$	0	*	0	*	0	*
cTL	${}^cR_{TL}^{00}$	0	${}^cR_{TL}^{0y}$	0	0	*	0	${}^cR_{TL}^{x'x}$	0	*	0	*	0	${}^cR_{TL}^{z'x}$	0	*
sTL	0	${}^sR_{TL}^{0x}$	0	${}^sR_{TL}^{0z}$	${}^sR_{TL}^{x'0}$	0	${}^sR_{TL}^{z'0}$	0	*	0	*	0	*	0	*	0
cTT	${}^cR_{TT}^{00}$	0	*	0	0	*	0	*	0	*	0	*	0	*	0	*
sTT	0	${}^sR_{TT}^{0x}$	0	${}^sR_{TT}^{0z}$	${}^sR_{TT}^{x'0}$	0	${}^sR_{TT}^{z'0}$	0	*	0	*	0	*	0	*	0
${}^cTL'$	0	${}^cR_{TL'}^{0x}$	0	${}^cR_{TL'}^{0z}$	${}^cR_{TL'}^{x'0}$	0	${}^cR_{TL'}^{z'0}$	0	*	0	*	0	*	0	*	0
${}^sTL'$	${}^sR_{TL'}^{00}$	0	${}^sR_{TL'}^{0y}$	0	0	*	0	${}^sR_{TL'}^{x'x}$	0	*	0	*	0	${}^sR_{TL'}^{z'x}$	0	*
TT'	0	$R_{TT'}^{0x}$	0	$R_{TT'}^{0z}$	$R_{TT'}^{x'0}$	0	$R_{TT'}^{z'0}$	0	*	0	*	0	*	0	*	0



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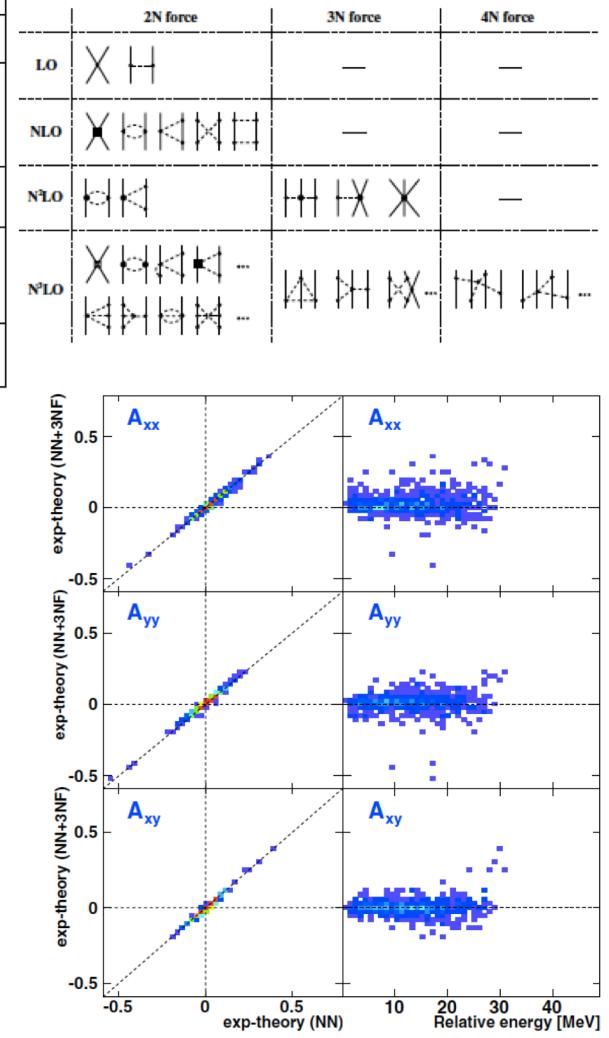
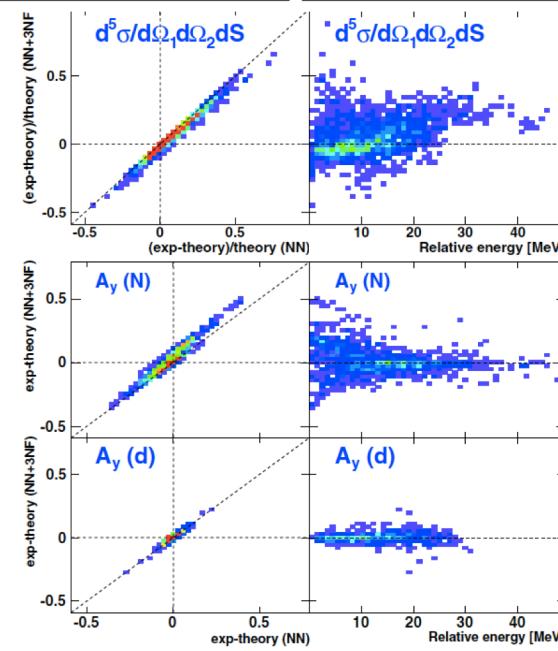
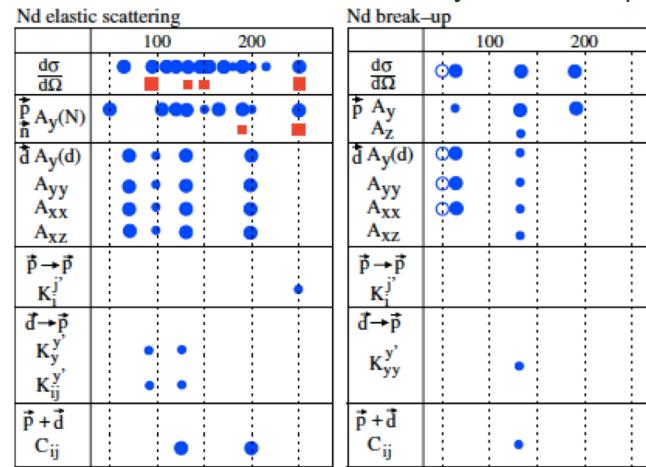
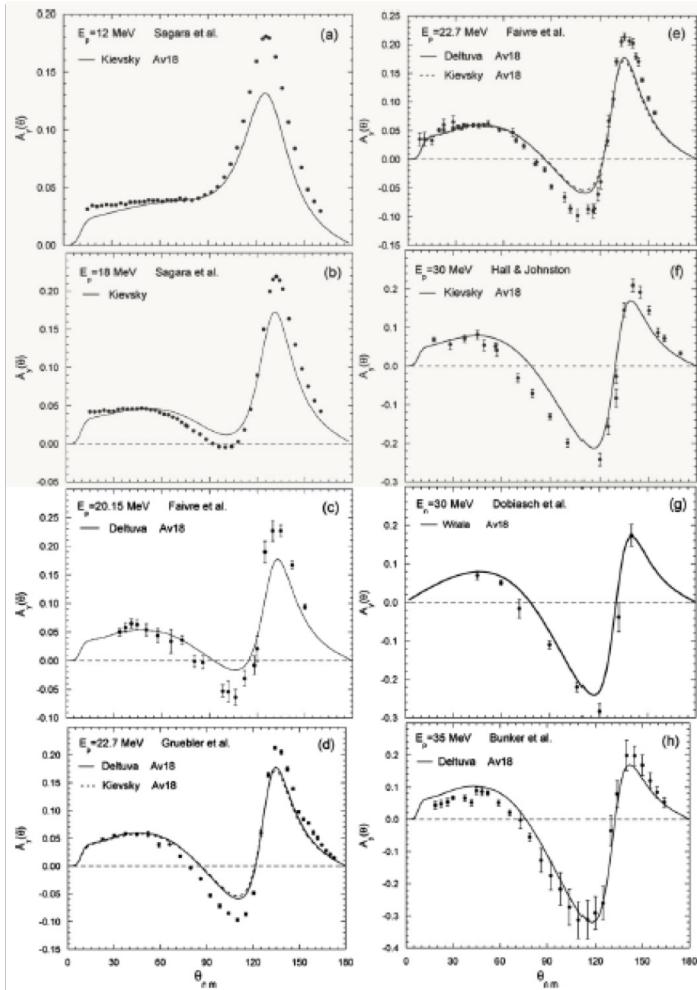
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Polarization in Heavy Ion Physics

Kalantar-Nayestanaki, Rep. Prog. Phys. **75** (2012) 016301

Torow et al., Nucl. Part. Phys. **35** (2008) 125104



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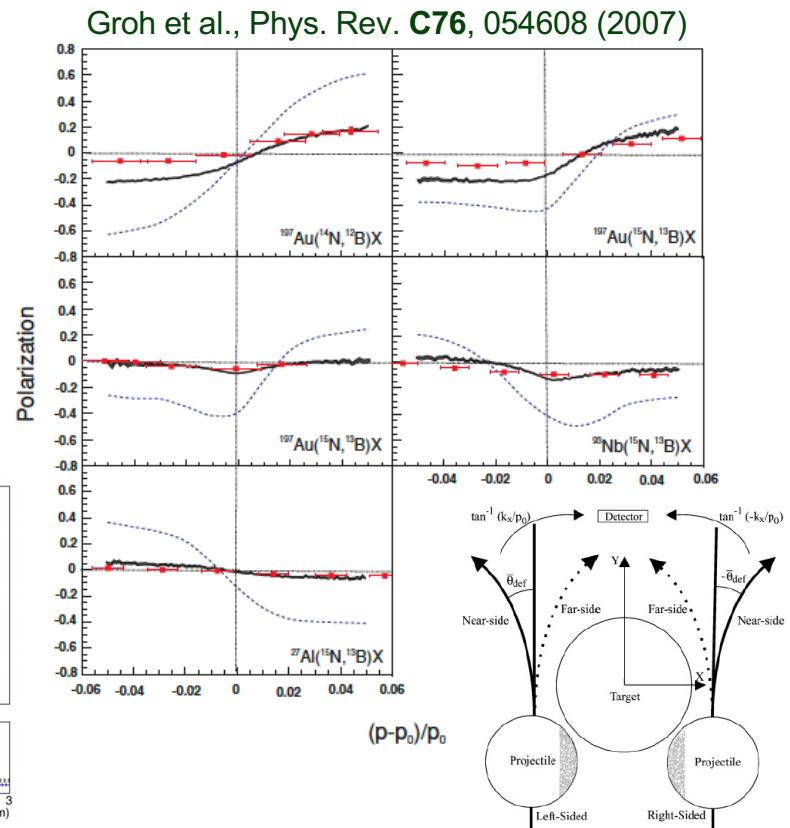
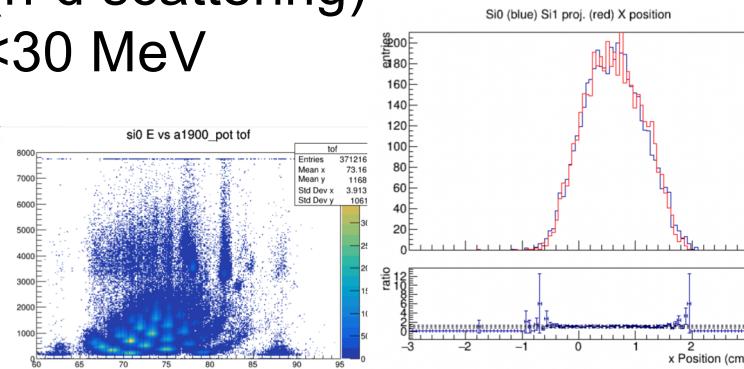
Measuring Polarized Observables with Rare Isotopes



Georgia Votta, GS

- Another “nitch” @ FRIB
- Beam polarization
 - ✓ Spin polarization in nuclear structure: allows spin-parity assignment
 - ✓ Fragments are (always) longitudinally polarized:
spin flip (polarized $^{31}\text{Na}^+$ @ TRIUMF)
 - ✓ Some good references
 - P. Mantica *et al.*, Phys. Rev. **C55**, 2501 (1997)
 - D. Hoff *et al.*, PRL **119**, 232501 (2017);
 - D. Hoff *et al.*, Phys. Rev. **C97**, 054605 (2018)

- Target polarization
 - ✓ The A_y puzzle (n-d scattering)
 - ✓ Asymmetry in <30 MeV
 - ✓ 3N & 4N forces





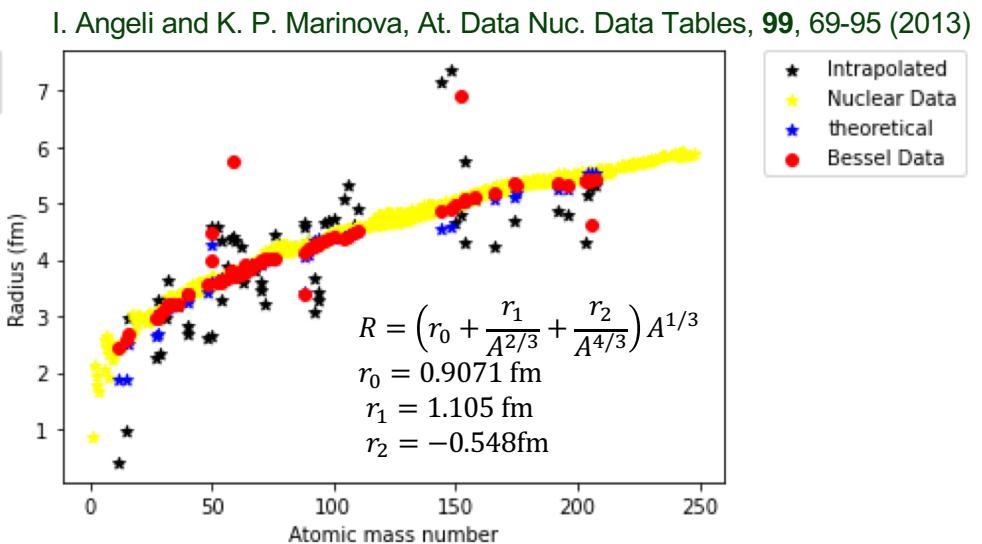
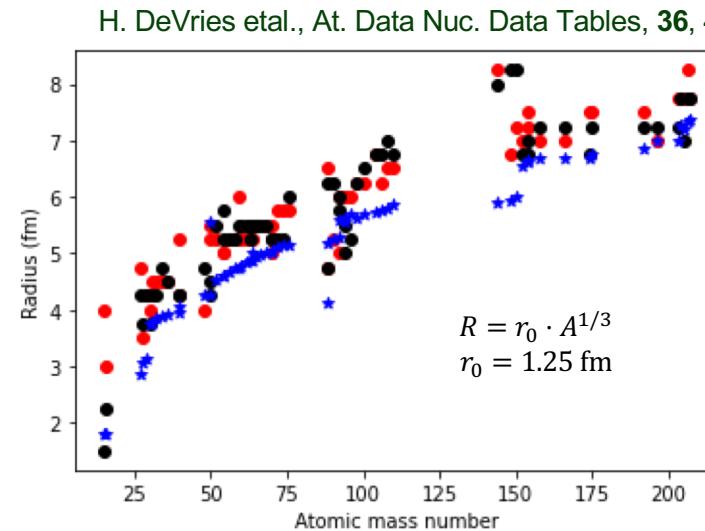
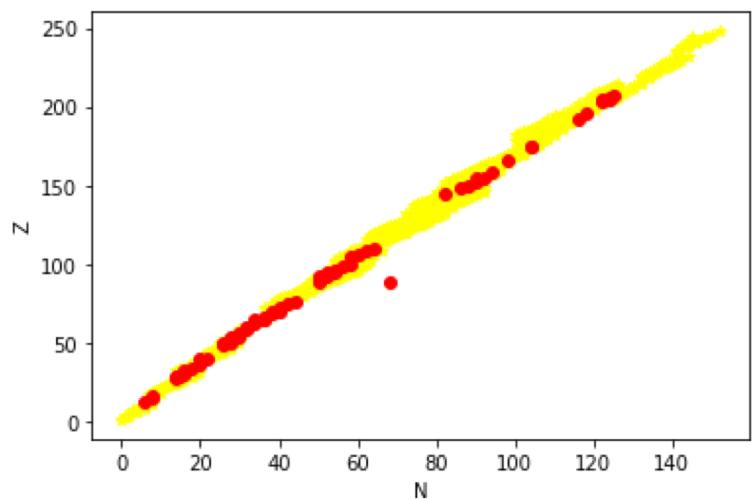
Rare Isotope Nuclear Radii

Ambar Rodriguez Alicea, GS

$$\rho(r) = \frac{1}{2\pi^2} \int F(q) \frac{\sin(qr')}{qr'} q^2 dq \\ = \sum a_\nu j_0(qr)$$

$$F(q) = 4\pi \int \frac{\sin(qr'/\hbar)}{qr'/\hbar} r'^2 \rho(r') dr'$$

$$j_0(qr) = \sum_{n=0}^{\infty} \frac{(-1)^n (qr)^{2n}}{2^{2n} (n!)^2} = \frac{\sin(qr)}{qr}$$

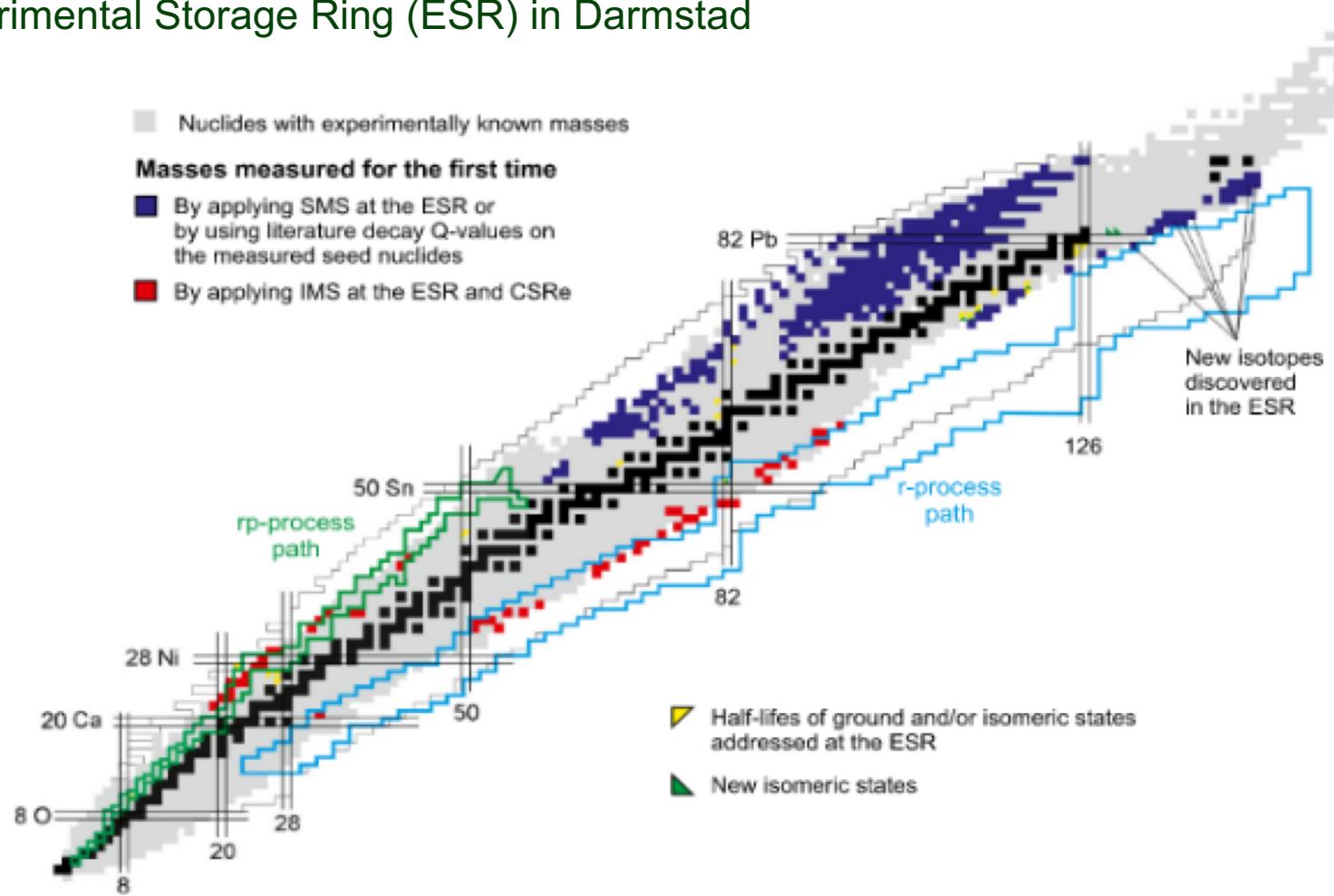


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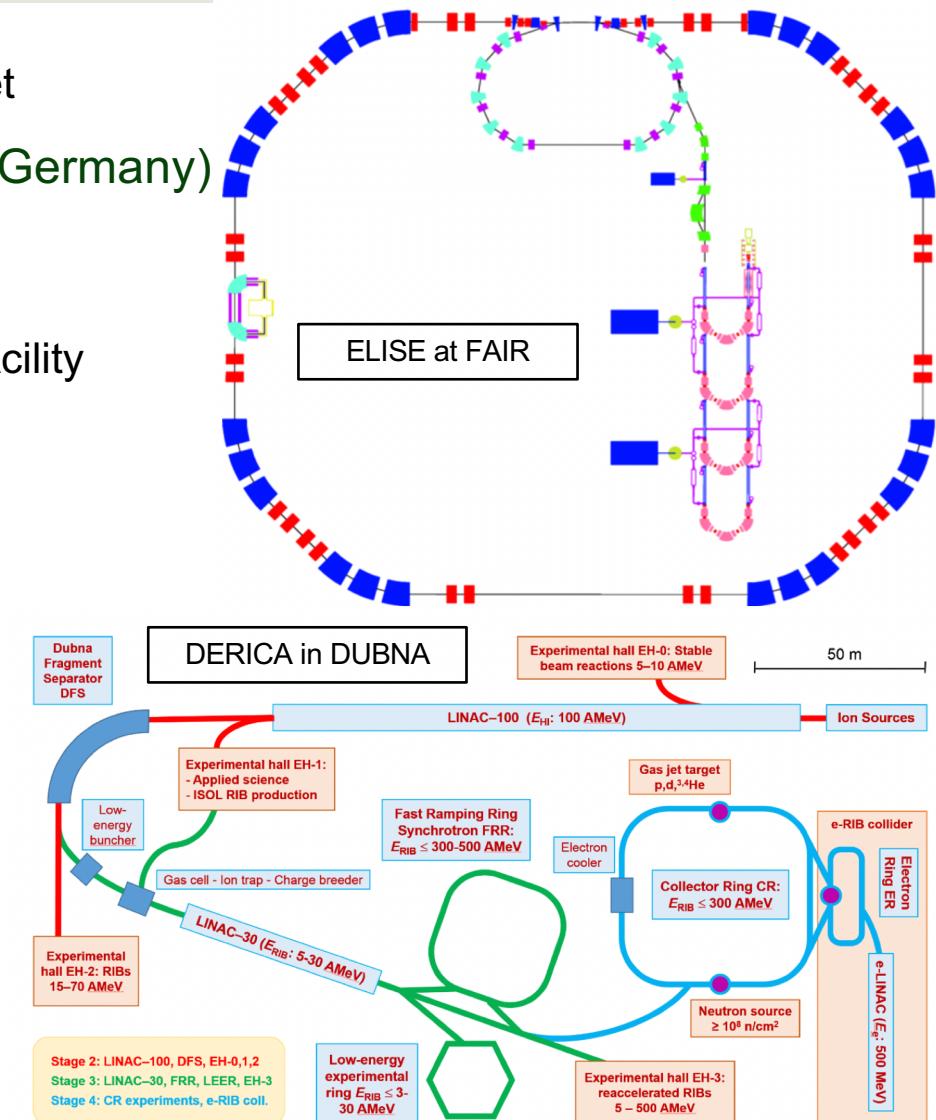
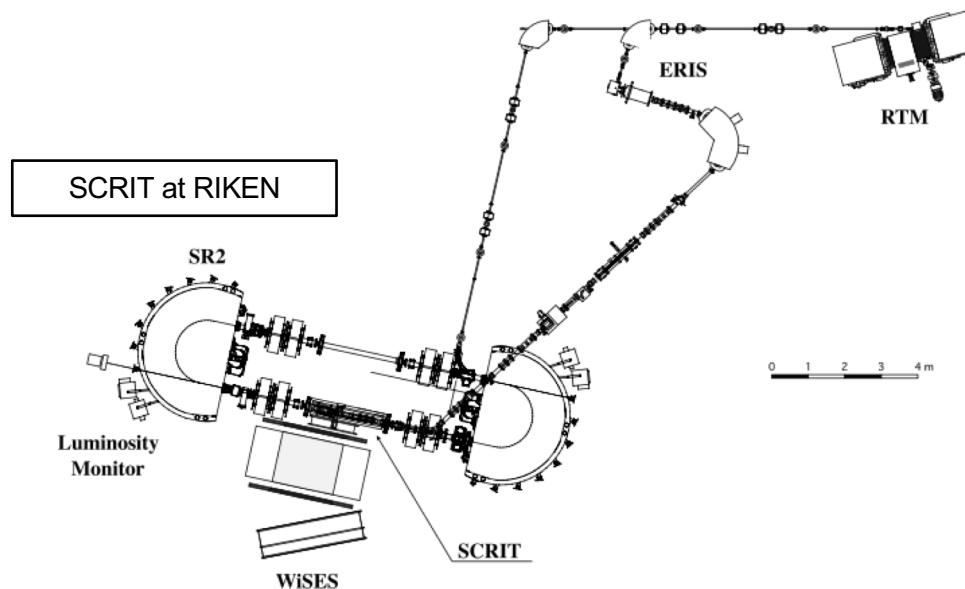
Storage Rings & Mass Measurements

Experimental Storage Ring (ESR) in Darmstad



Electron/Rare Isotopes Systems

- RI-RIKEN
 - SCRIT: Self-Confining Radioactive isotope Ion Target
- Facility for Anti-proton and Ion Research (FAIR, Germany)
 - ELISE: Electron-Ion Scattering in a Storage Ring
- DUBNA (Russia)
 - DERICA: Dubna Electron-Radioactive Ion Collider fAcility



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Workshop, Symposia & Brainstorming Sessions

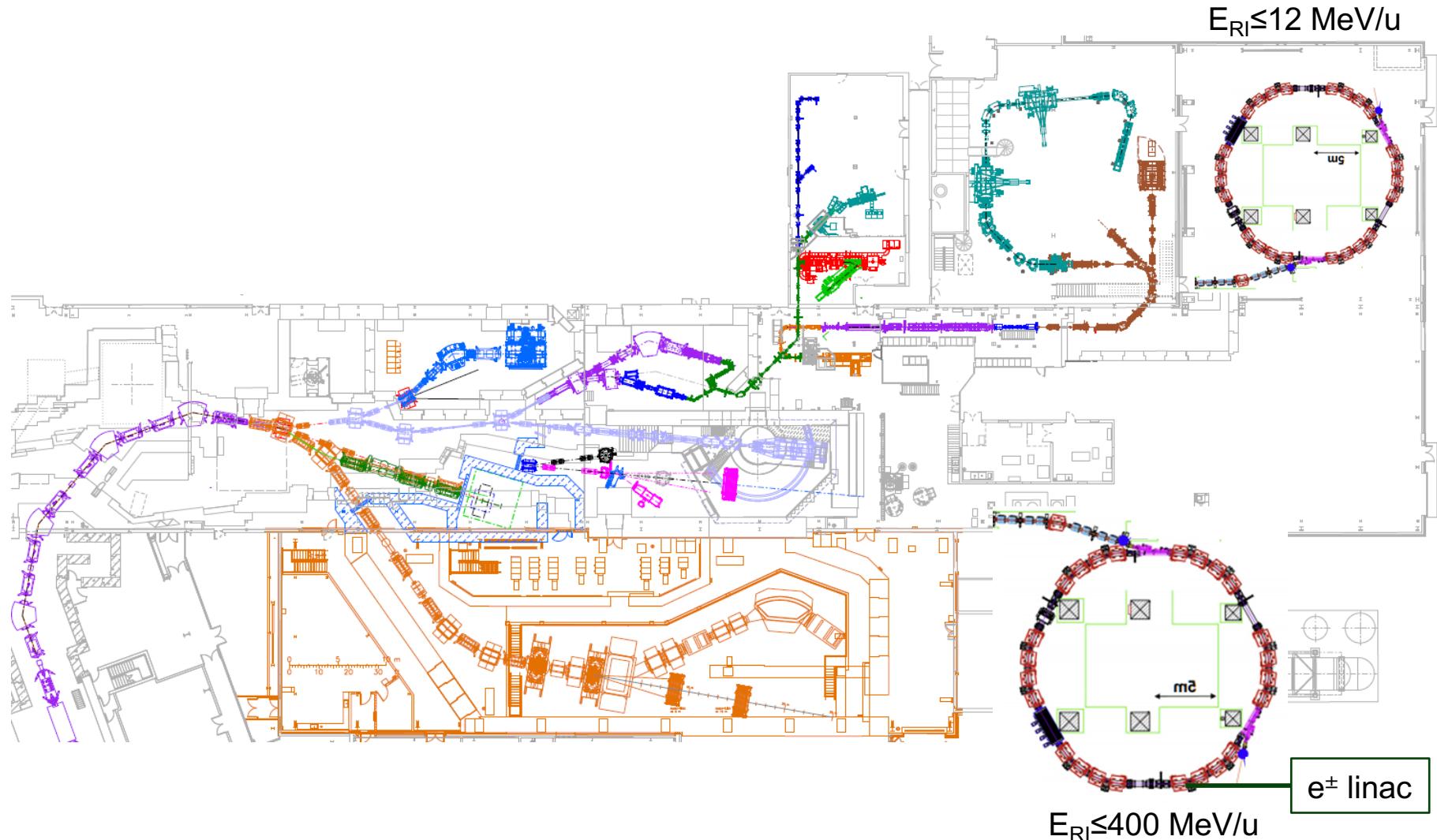
- North America Storage Rings & Neutron Captures Workshop
 - June 28-30, 2021
 - <https://meetings.triumf.ca/event/235/overview>
- APS/DNP symposia
 - October 11-14, 2021
 - JF: Advances and Opportunities in Polarized Targets and Beams I
 - KF: Advances and Opportunities in Polarized Targets and Beams II
- Brainstorming sessions
 - March 25, 2022 and April 01, 2022
 - Participants: Eric Voutier, Toshimi Suda, Maya Wallach, Claude Marchand, Dominic Marchand, Joe Grames, Kei Minamisono, Peter Ostroumov, Kent Paschke, Ryan Richards, Or Hen, Alain Lapierre, Michael Kohl



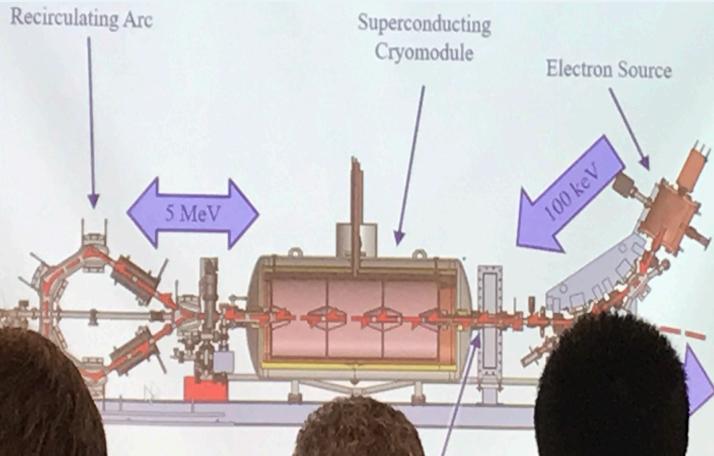
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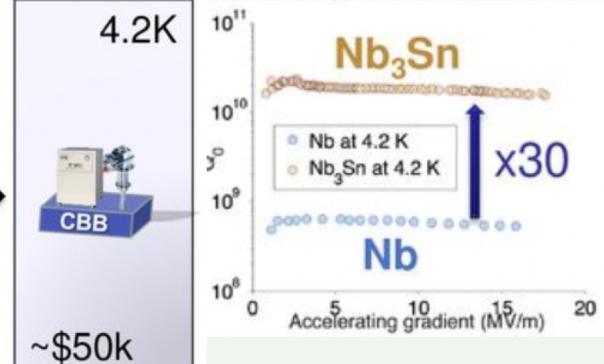
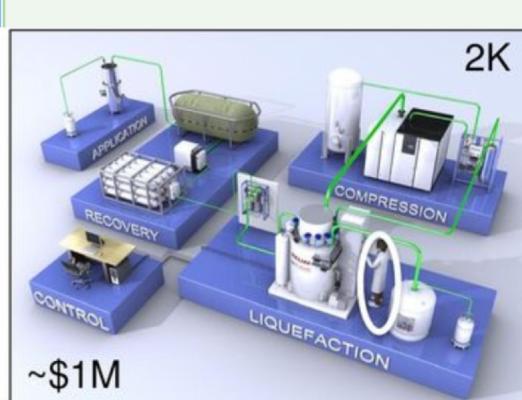
e^\pm -FRIB Concept



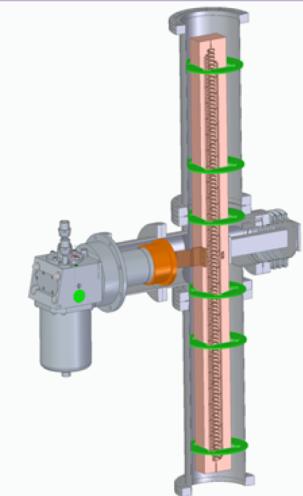
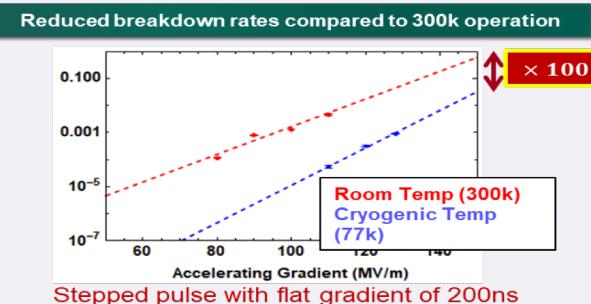
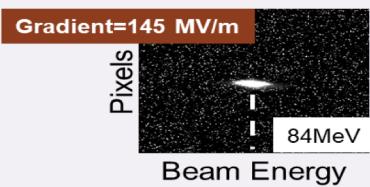
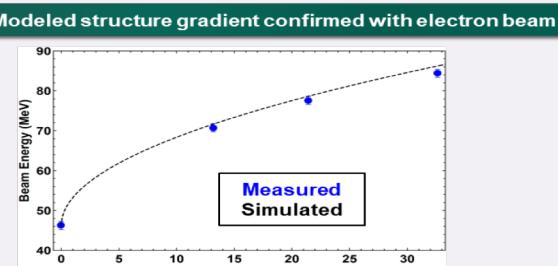
Electron/Positron Beam



Niowave, <https://www.niowaveinc.com>



Center for Bright Beams @ Cornell, <https://cbb.cornell.edu>



Very High Energy Electron (VHEE) @ SLAC

- Compact, ultrahigh efficiency/gradient linac
- Operating at 80K, require 20 MW to reach 100 MV/m
- Typical linacs would have required about 200 MW of power



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RI Storage Rings – 1

Experimental Storage Ring, ESR at GSI

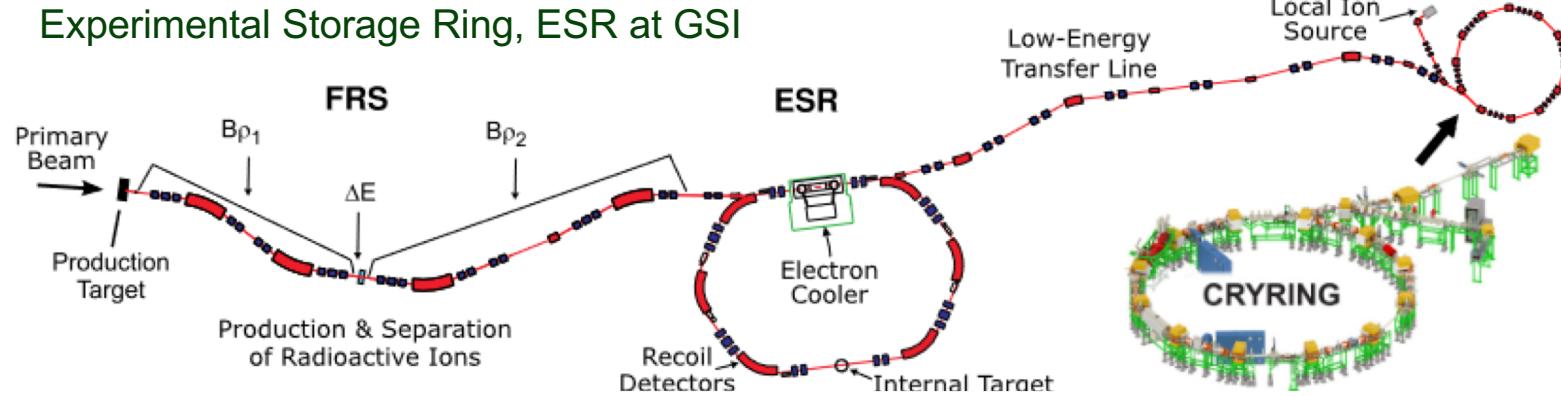


Table 1
Active heavy ion storage rings.

	ESR	CSRc	R3
Circumference	108.4 m	128.8 m	60.4 m
Maximum magnetic rigidity	10 Tm	9.4 Tm	5.0 Tm
Maximum beam energy U^{q2+}	550 MeV/u	500 MeV/u	
Minimum beam energy U^{q2+}	3 MeV/u	5 MeV/u	
Beam cooling			
Internal target			
Standard mode			
Tunes Q_x/Q_y			
Acceptance (Λ_x/Λ_y) [$\mu\text{m}/(\delta p/p)$] [%]			
Isochronous mode			
Tunes Q_x/Q_y			
Acceptance (Λ_x/Λ_y) [$\mu\text{m}/(\delta p/p)$] [%]			
Transition energy γ_t			

The figure consists of three maps. The left map shows Europe with a red dot indicating the location of the ESR in Germany. The middle map shows Asia with a red dot indicating the location of the CSRc in China. The right map shows Japan with a red dot indicating the location of the R3 ring in Japan.



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RI Storage Ring – 2



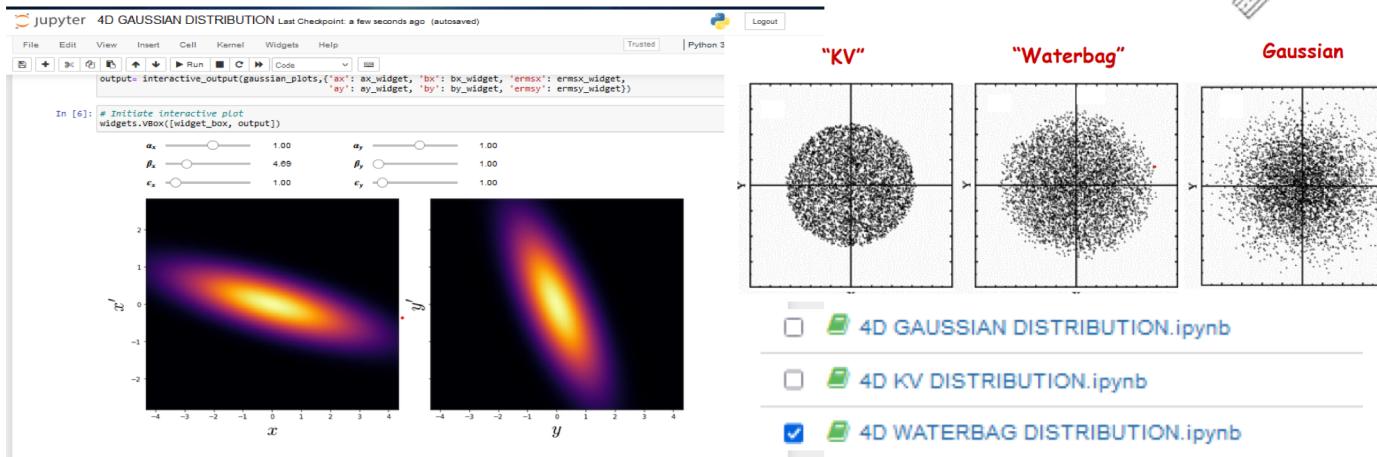
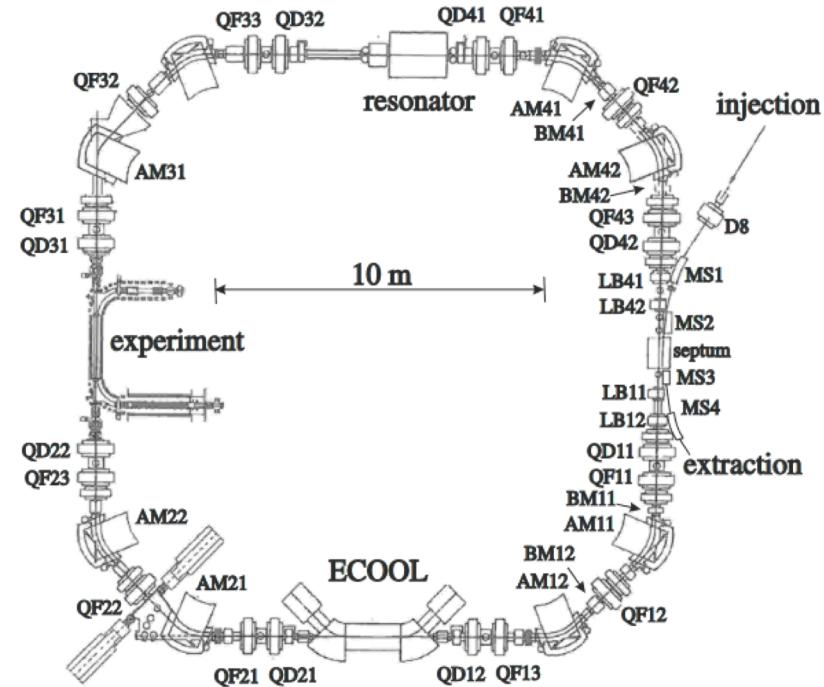
Jeremy Rebenstock, UG

▪ Components

- Injection
- Beam cooling system (phase space)
- Internal target
- Magnet system
- Detector (diagnostics) system
- Extraction

▪ Test Storage Ring facility (model)

- Heidelberg, 1988-2012



The equation for the beam ellipse, with our Twiss parameterization can be written as:

$$\epsilon = \gamma x^2 + 2\alpha xx' + \beta x'^2$$

$$\gamma = \frac{1 + \alpha^2}{\beta}$$



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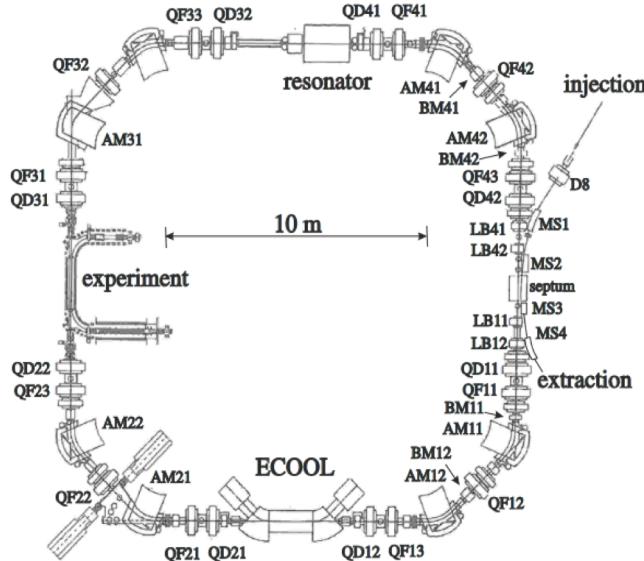
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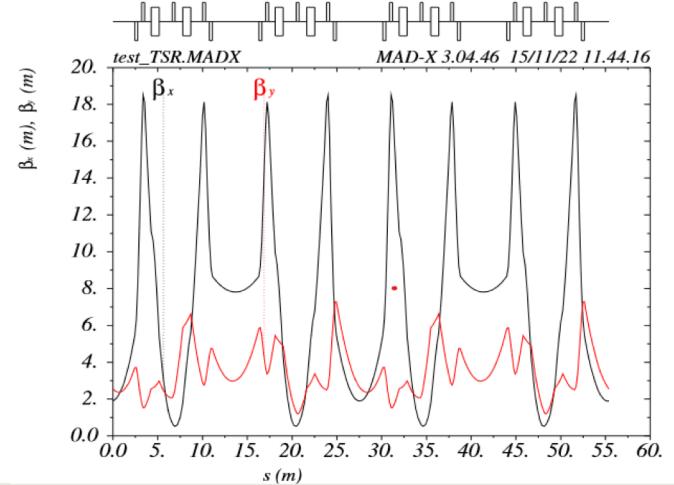
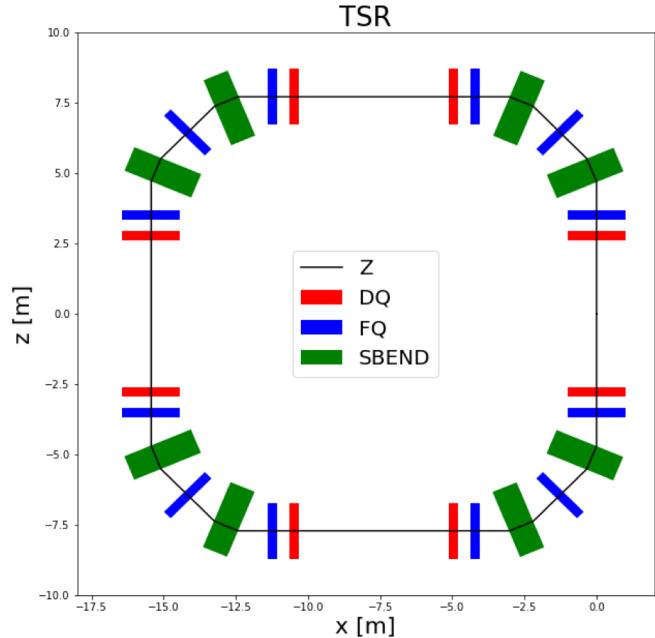
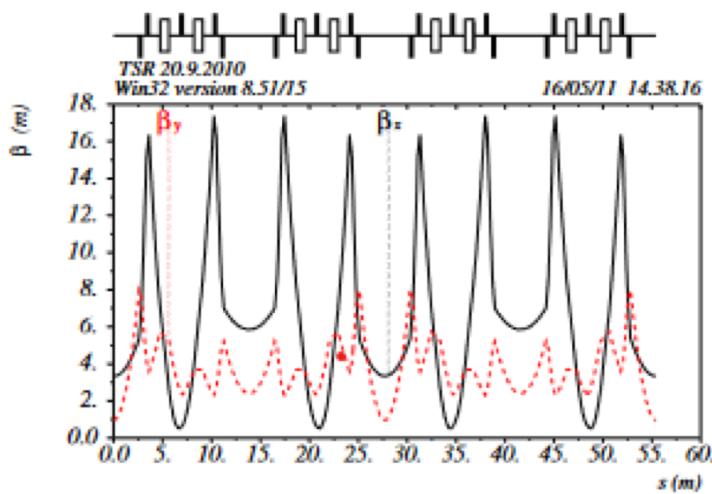
RI Storage Ring – 3



Jeremy Rebenstock, UG



Cristhian Gonzalez Ortiz
PhD advisor: P. Ostroumov
TSR Conceptual report



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Some numbers ...

T. Suda and H. Simon, Prog. Part. Nucl. Phys., **96**, 1-31 (2017)

Table 2

Estimate for the required luminosities for different studies in colliding beam kinematics. It is assumed that the maximum running time shall not exceed four weeks.

Reaction	Deduced quantity	Target nuclei	Luminosity \mathcal{L} $\text{cm}^{-2} \text{s}^{-1}$
Elastic scattering at small q	r.m.s. charge radii	Light Medium	10^{24}
First minimum in elastic form-factor	Density distribution with 2 parameters	Light Medium Heavy	10^{28} 10^{26} 10^{24}
Second minimum in elastic form-factor	Density distribution with 3 parameters	Medium Heavy	10^{29} 10^{26}
Giant resonances	Position, width, strength, decays	Medium Heavy	10^{28} 10^{28}
Quasi-elastic scattering	Spectroscopic factors, spectral function, momentum distributions	Light	10^{29}

	Ee (GeV)	Beam Current	Ne/s	Target (/cm²)	L (/cm²/s)
Hofstadter (1950)	0.15	1 nA	6.25×10^{09}	6.0×10^{19}	3.8×10^{29}
JLab	12	100 μA	6.25×10^{14}	6.0×10^{22}	3.8×10^{37}
SCRIT (20 W)	0.15-0.30	300 mA	1.88×10^{18}	1.0×10^{09}	1.9×10^{27}
FRIB (400 kW)	0.2	300 mA	1.88×10^{18}	1.9×10^{13}	3.6×10^{31}
positrons (1/100)	0.2	3 mA	1.88×10^{16}	1.9×10^{13}	3.6×10^{29}



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Next Steps

■ Science case

- Rare isotope rings: nuclear astrophysics community
- e^\pm scattering: electron scattering community
- Goals
 - » Series of (hybrid) workshops: summer 2023, fall 2023, spring 2024
 - » Science report: spring/summer 2024

■ Accelerator systems

- RI ring systems: White Paper NSAC LRP
 - » Nuc. Astro. team is working on it
- e^\pm system
 - » Interest: Cornell, BNL ... Jlab



NSAC Long Range Plan Town Hall Meeting on Nuclear Structure, Reactions, and Astrophysics
November 14–16, 2022, Argonne National Laboratory

Thank You!

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