Status of the DVCS proposals

Positron WG Collaboration Meeting

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Introduction

- Two proposals using a positron beam were submitted to PAC48 (2020)
- ➤ One experiment proposed to use CLAS12 (Hall B) with a low current (≈45 nA) polarized beam
- > One experiment proposed to use NPS (Hall C) with a high current (\approx 5 μ A) unpolarized beam

PAC 48 SUMMARY OF RECOMMENDATIONS										
PR12-20-009	E. Voutier	Beam charge asymmetries for Deeply Virtual Compton Scattering on the proton at CLAS12	В	100			C2	4		
PR12-20-012	C. Munoz Camacho	Deeply Virtual Compton Scattering using a positron beam in Hall C	С	77			C2	4		

Add C2 definition here!

Motivation



Access in helicity-independent cross section

Access in helicity-dependent cross-section

Positron production and transport



Electrons

Dominated by damping in the LINACS

Dominated by synchrotron rad. in Arcs

Area	δp/p	ε _x	εγ	
	[x10 ⁻³]	[nm]	[nm]	
Chicane	0.5	4.00	4.00	
Arc 1	0.05	0.41	0.41	
Arc 2	0.03	0.26	0.23	
Arc 3	0.035	0.22	0.21	
Arc 4	0.044	0.21	0.24	
Arc 5	0.060	0.33	0.25	
Arc 6	0.090	0.58	0.31	
Arc 7	0.104	0.79	0.44	
Arc 8	0.133	1.21	0.57	
Arc 9	0.167	2.09	0.64	
Arc 10	0.194	2.97	0.95	
Hall D	0.18	2.70	1.03	

Area	δ p/ p	ε _x	ε _y	
	[x10 ⁻³]	[nm]	[nm]	
Chicane	10	500	500	
Arc 1	1	50	50	
Arc 2	0.53	26.8	26.6	
Arc 3	0.36	19	18.6	
Arc 4	0.27	14.5	13.8	
Arc 5	0.22	12	11.2	
Arc 6	0.19	10	9.5	
Arc 7	0.17	8.9	8.35	
Arc 8	Arc 8 0.16		7.38	
Arc 9	Arc 9 0.16		6.8	
MYAAT01	0.18	9.13	6.19	

Positrons

At 11 GeV, after Arc9, e+ beam size ~twice bigger than e- beam

Averaging εx and ε_y:



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TAC comments on positron

- The implementation of a multi-Hall, high current, high polarization positron beam at CEBAF raises multiple and complex challenges, as detailed in the TAC report
- If the PAC finds our physics program compelling, our collaboration is ready to engage with the Lab to investigate its feasibility.

TAC conclusion:

In conclusion, while a positron beam upgrade is a major upgrade which will require substantial accelerator physics development, a detailed cost and implementation plan, and expensive changes to the CEBAF accelerator, a multi-Hall positron beam capability could have great potential for a future JLAB 12-GeV science program.

Hall B positron experiment

CLAS12

V. Burkert et al. NIMA 959 (2020) 163419



- There is **no difference** bewteen **e**⁻ and **e**⁺ beam transport in Hall B beam line, nor in beam related detector background.

- Beam diagnostics are expected to operate similarly with e⁻ and e⁺ beam.

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Hall B beam time request

We are asking for a total of **100 days** of beam, operating **CLAS12** with **45 nA e**⁻ and **e**⁺ beams polarized at **60%**, distributed in :

		Beam parameters					Sol	Tor		
Purpose	Label	q	Not	E	I	λ	Target	Dol	Dol	Time
		(e)	Ivat.	(GeV)	(nA)	(%)		FOI.	FOI.	(h)
			Р							24
$en \rightarrow en$	Cal		1	2.2		0		+	+	24
$cp \rightarrow cp$	Uai.			2.2		0		+		24
					15		5 cm LH ₂			24
$ep \rightarrow ep\gamma$	Phy.		S		40			- 2		480
Background	Cal.		6	10.6		60		-		48
$ep \rightarrow ep\gamma$	Phy.							+	1	480
Background	Cal.							+		48
Commission	ning							+		72
$ep \rightarrow ep$	Cal			2.2		0		+		24
								10-70		24
Commission	ing	ng			15		$5~\mathrm{cm}$	-	_	72
$ep \rightarrow ep\gamma$	Phy.		5		40	60	LH_2			480
Background	Cal.			10.6		00		-		48
$ep \rightarrow ep\gamma$	Phy.							+	1	480
Background	Cal.							+	1	48
									Total	2400

- 80 days for physics data taking;
- 20 days for commissioning and calibration;

using lepton beams of different nature

2 days with the CEBAF e⁻ beam;

 \circ 46 days with the secondary e⁻ beam;

o 52 days with the secondary e⁺ beam;

and different energies

```
9 days at 2.2 GeV;
91 days at 10.6 GeV.
```

Hall B projections

> A sample of expected experimental data... 15 bins in $(x_B, Q^2) \times 6$ bins in t = 90 azimuthal dependences per observable (A_{UU}^C, A_{LU}^C) .



Hall B impact of positrons

 \geq The importance of positron beams for the determination of CFFs can be quantified in a model-dependent way depending on : the cross section model, the GPDs model, and the hypotheses of the fitting approach.

Observable	σ_{UU}	A_{LU}	A_{UL}	A_{LL}	A_{UU}^C	A_{LU}^C
Sytematics (%)	5	3	3	$3 \oplus 3$	3	3

=

Fitting of $\{\mathcal{H}, \tilde{\mathcal{H}}\}$ CFFs assuming model values for $\{\mathcal{E}, \tilde{\mathcal{E}}\}$ CFFs.





KM15 Δ Re H without / with positrons

DVCS with positrons in Hall C

$$|\mathcal{T}(\pm ep \to \pm ep\gamma)|^2 = |\mathcal{T}^{BH}|^2 + |\mathcal{T}^{DVCS}|^2 + \mathcal{I}$$
Opposite sign for e- & e+

✓ Precise determination of the absolute photon electro-production cross section

- ✓ Clean, model-independent separation of DVCS² and DVCS-BH interference
- ✓ More stringer constraints on CFFs by combining e⁻ & e⁺ data

<u>In a nutshell:</u>

- Same experimental configuration as approved experiment E12-13-010
- > Expected positron beam momentum spread comparable with current electron beam
- Positron beam size larger than current electron beam (twice bigger at 11 GeV according to current simulation)
- No additional systematic uncertainties expected due to the use of positrons 10/18

Hall C positron experiment using NPS

- 1080 PbWO₄ crystals
- 0.6 Tm sweeping magnet
- F250ADC sampling electronics
- Large opening angle beam pipe
- SHMS as carriage for rotation



0<u>-</u>

0.5

1.5

2.5 MM² GeV²



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PR12-20-012: Kinematic & beam time request



Separation of DVCS² and BH-DVCS interference

Projections based on the KM15 model (Kumericki and Mueller, 2015)



Impact on Compton Form Factors (CFFs) extraction



(factor of ~2 for HT and NLO)

PAC48 report

Hall B proposal

As a secondary e^+ beam has a much larger momentum dispersion and emittance than the primary e^- beam currently existing, one has the same $\delta p/p$ with a spot size 2-3 times larger at the target. Therefore, a target cell with 50% larger (15 mm diameter) entrance and exit windows is needed to avoid any interaction with the target structure frames. This difference in the beam parameters leads the collaboration to require equal statistics of electron and positron data using secondary e^+ and e^- beams.

Hall C proposal

Therefore, there are no concerns about the experimental setup, and the committee is confident that the described measurements can be performed.

Common to Hall B and Hall C proposals

Issues: The iTAC report concludes that "while a positron beam upgrade is a major upgrade, which will require substantial accelerator physics development, a detailed cost and implementation plan, and expensive changes to the CEBAF accelerator, a multi-Hall positron beam capability could have great potential for a future JLAB 12-GeV science program." The PAC fully agrees with this statement.

To clearly show what science would be left unaddressed without having a polarized positron beam available at JLab, the physics impact of positron beam experiments should be demonstrated more rigorously. This should include an impact study with positron beam pseudo-data and all other existing and anticipated future data with an electron beam. One would also like to see which amount of electron data alone would be needed to obtain comparable constraints on the CFFs. Furthermore, to guide the technological development of a positron source, it is important that the experiment determines clear performance requirements for the positron beam.

The PAC notes that proposal PR12-20-012 requires no electron data set with equal beam conditions for an extraction and separation of CFFs with high accuracy. It would be interesting if both groups could come to a common understanding on this issue.

EPJA paper (Pawel et al.)

Impact of a positron beam at JLab on an unbiased determination of DVCS Compton form factors

CLAS12 data (unpolarized target): e- & e+



Gepard (K. Kumericki's code)

🕋 gepard

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* Tool for studying the 3D quark and gluon distributions in the nucleon

View page source

Tool for studying the 3D quark and gluon distributions in the nucleon

Fitting theory to data

At the moment, only the standard least-squares fitting is implemented, and it uses the iminuit Python interface for the Minuit2 C++ library.

Gepard user creates a MinuitFitter object, where first argument is a dataset (collection of DataPoint objects), and second argument is a Theory object. Then fit() method of this MinuitFitter objects starts the fit (by calling migred optimizer of Minuit). Before this, user should release some parameters of the Theory because all parameters are by default fixed at the moment of the creation of Theory object.

So, the minimal example of fitting is:

```
>>> import gepard as g
>>> tass fittest(g.gpd.PuNNormGPD, g.cff.MellinBarnesCFF, g.dvcs.BMK):
... pass
>>> tha = FitTest()
>>> f = g.MinuitFitter(pts, th)
>>> f.release_parameters('ns', 'ms2', 'secs')
>>> f.fit()
```

Final values of chi-square and of parameters are available as

>>> f.minuit.fval								
8.411								
>>> f	.pr	int_parameters()						
ns	=	1.46 +- 0.45						
ms2	=	0.93 +- 0.07						
secs	-	-0.32 +- 0.01						

After successful fit of theory object th, user can access parameter uncertainties as th.parameters_errors dictionary, and full covariance matrix (inverse of the chi-square Hessian matrix) as th.covariance dictionary.

Covariance matrix can then be used to propagate uncertainty to prediction of observables, like this:

```
>>> th.predict(pts[0], uncertainty=True)
(13.25, 1.62)
```

where parameter-dependent form factors (such as CFFs) can also be "predicted", i. e., calculated together with their uncertainty:

>>> pt = g.DataPoint(x8=0.01, t=-0.2, Q2=10)
>>> th.predict(pt, observable='ImH', uncertainty=True)
(273.2, 25.8)

Summary and outlook

- > Both Hall B and Hall C proposals are C2 from PAC48 (3 years ago)
- To address the report, a better, more comprehensive way to show the impact of positrons for DVCS and CFF extraction needs to be presented:

Two possibilities:

- Extend EPJA paper to all DVCS data (present and future) using PARTONS
- Develop a global fit using the GEPARD code
- The Hall B proposal needs to clarify (or justify better) the need of electron new/additional electron data for the BCA measurement

BACK-UP

Systematic uncertainties

Source	pt-to-pt (%)	scale (%)
Acceptance	0.4	1.0
Electron PID	<0.1	<0.1
Efficiency	0.5	1.0
Electron tracking	0.1	0.5
Charge	0.5	1.0
Target thickness	0.2	0.5
Kinematics	0.4	<0.1
Exclusivity	1.0	2.0
π^0 subtraction	0.5	1.0
Radiative corrections	1.2	2.0
Total	1.8-1.9	3.4-3.5

The π^0 electroproduction cross section would be measured concurrently with DVCS with both electrons and positrons, and would allow to monitor the systematics of the e- and e+ runs

Impact on Compton Form Factors (CFFs) extraction



 ✓ Combined fit of all electron data from approved experiment E12-13-010

(helicity-dependent AND helicity-independent cross sections)

- $\checkmark\,$ Fits with and without the proposed positron data
- Fits include helicity-conserving CFFs, but also +1 helicity-flip CFFs ("HT") and +2 helicity-flip CFFs ("NLO")
- \checkmark Cross sections generated with CFFs values fitted to 6 GeV data

In order to extract the CFFs we exploit the combined

- Azimuthal dependence (ϕ)
- Beam-energy dependence
- Q²-dependence
- Helicity dependence (for E12-13-010 data)
- Beam-charge dependence
- of the DVCS cross section

Correlation coefficients

Correlations between different CFFs are significantly improved by a combined fit with positrons

$$|\rho_{i,j}| = \operatorname{cov}[\mathbb{F}_i, \mathbb{F}_j]/(\sigma_i \sigma_j)$$

Electrons & Positrons

Sm(Ĩ,) Sm(Ĥ_) 0.9 0.9 ℜe(Ĥ_) ℜe(Ĥ_) 0.8 0.8 Sm(H_) Sm(H_) **ℜe(H_**) %e(H_) 0.7 0.7 ଞm(ୖH_⊶) ິສm(ୖୖୄ୷) 0.6 0.6 **ℜe(ୖH_{₀+})** ઉદ(Ĥୁ) 0.5 0.5 ଞm(Hୁ) Sm(H_,) 0.4 0.4 େ(H_{o+}) **ℜe(H_{_1})** 0.3 0.3 ଞm(Ĥ₊₊) ଞm(Ĥ₊₊) Re(Ĥ₊₊ ℜe(Ĥ₊_) 0.2 0.2 Sm(H__) Sm(H_.) 0.1 0.1 େ(H₊₊) **ℜe(H₊₊)** n 0 ℜe(Ĥ_↓) (⁺⁺)m βe(H,,) $h(\widetilde{H}_{0+})$ 3m(H,,) 3m(Ĥ₊₊) îe(H₀,) Sm(H₀,) %e(Ĥ₀,) Sm(Ĥ₀,) 3ke(H__) 3m(H__) 3m(Ĥ_++) 3m(H,₁) βe(H__) 3m(Ĥ_,) βe(H̃₊₊) ßm(Ĥ__) βe(H₀,) %e(Ĥ₊₊) Ste(H₁₁ Sm(H HT NLO LT/LO $(t = -0.26 \text{ GeV}^2)$ Much better separation of H & Ht CFFs at LT/LO

(from -94% without positrons to -39% when electron and positrons are combined, in this t-bin)

Electrons only

$$\begin{split} \sigma(ep \to ep\gamma) &= \underbrace{|BH|^2}_{\text{Known to} \sim 1\%} + \underbrace{\mathcal{I}(BH \cdot DVCS)}_{\text{Linear combination of GPDs}} + \underbrace{|DVCS|^2}_{\text{Bilinear combination of GPDs}} \\ \mathcal{I} \propto 1/y^3 &= (k/\nu)^3, \\ \left|\mathcal{T}^{DVCS}\right|^2 \propto 1/y^2 &= (k/\nu)^2 \end{split}$$



 φ -dependence provides 5 independent observables:

 ${\sim}1$, ${\sim}\cos\varphi, {\sim}\sin\varphi$, ${\sim}\cos(2\varphi), {\sim}\sin(2\varphi)$

• Cross section measured at 2 beam energies and constant Q^2 , x_B , t



Leading-twist and LO simultaneous fit of both beam energies (dashed line) does not reproduce the data
 Light-cone axis in the (q,q') plane (Braun et al.): II++, II++, E++, E++, E++





- Using only helicity-conserving CFFs ("LT/LO") the fit of both beam energies (dashed line) does not reproduce the data
- Including helicity-flip CFFs, either single-helicity flip ("HT") or double-helicity flip ("NLO") satisfactorily reproduce the angular dependence (blue solid line)

DVCS² and \mathcal{I} (DVCS·BH) separated in NLO and higher-twist scenarios



• DVCS² & *I* significantly different in each scenario

 Sizeable DVCS² contribution in the higher-twist scenario in the helicity-dependent cross section

Nature Commun. 8, 1408 (2017)

DVCS process: leading twist ambiguity

- DVCS defines a preferred axis: light-cone axis
- At finite Q^2 and non-zero t, there is an ambiguity:
 - **1** Belitsky et al. ("BKM", 2002–2010): light-cone axis in plane (q, P)
 - 2 Braun et al. ("BMP", 2014): light-cone axis in plane (q,q')easier to account for kin. corrections $\sim O(M^2/Q^2)$, $\sim O(t/Q^2)$

$$\begin{aligned} \mathcal{F}_{++} &= & \mathbb{F}_{++} + \frac{\chi}{2} \left[\mathbb{F}_{++} + \mathbb{F}_{-+} \right] - \chi_0 \mathbb{F}_{0+} \\ \mathcal{F}_{-+} &= & \mathbb{F}_{-+} + \frac{\chi}{2} \left[\mathbb{F}_{++} + \mathbb{F}_{-+} \right] - \chi_0 \mathbb{F}_{0+} \\ \mathcal{F}_{0+} &= & -(1+\chi) \mathbb{F}_{0+} + \chi_0 \left[\mathbb{F}_{++} + \mathbb{F}_{-+} \right] \end{aligned} \right\} \xrightarrow{\mathbb{F}_{-+} = 0} \begin{cases} \mathcal{F}_{++} &= (1 + \frac{\chi}{2}) \mathbb{F}_{++} \\ \mathcal{F}_{-+} &= \frac{\chi}{2} \mathbb{F}_{++} \\ \mathcal{F}_{0+} &= \chi_0 \mathbb{F}_{++} \end{cases} \end{aligned}$$

(eg. $\chi_0 = 0.25$, $\chi = 0.06$ for $Q^2 = 2$ GeV², $x_B = 0.36$, t = -0.24 GeV²)

DVCS cross-section: $\varphi \& Q^2$

$$\mathcal{I} = \frac{i_0/Q^2 + i_1 \cos \varphi/Q + i_2 \cos 2\varphi/Q^2 + i_3 \cos 3\varphi/Q}{\mathcal{P}_1 \mathcal{P}_2}$$

$$\mathsf{DVCS}^2 = \frac{d_0/Q^2 + d_1 \cos \varphi/Q^3 + d_2 \cos 2\varphi/Q^4}{2}.$$

The product of the BH propagators reads:

$$\mathcal{P}_1 \mathcal{P}_2 = 1 + \frac{p_1}{Q} \cos \varphi + \frac{p_2}{Q^2} \cos 2\varphi.$$

Reducing to a common denominator ($\times \mathcal{P}_1 \mathcal{P}_2$), one obtains:

$$\mathcal{P}_{1}\mathcal{P}_{2}\mathcal{I} + \mathcal{P}_{1}\mathcal{P}_{2}\mathsf{DVCS}^{2} = \boxed{(i_{0} + d_{0})/Q^{2}} + d_{1}p_{1}/2/Q^{4} + p_{2}d_{2}/2/Q^{6} \\ + [i_{1}/Q + (p_{1}d_{0} + d_{1})/Q^{3} + (p_{1}d_{2} + p_{2}d_{1})/2/Q^{5}]\cos\varphi \\ + [i_{2}/Q^{2} + (p_{2}d_{0} + p_{1}d_{1}/2 + d_{2})/Q^{4}]\cos2\varphi \\ + [i_{3}/Q + (p_{1}d_{2} + p_{2}d_{1})/2/Q^{5}]\cos3\varphi \\ + [p_{2}d_{2}/4/Q^{6}]\cos4\varphi \,.$$

The $\mathcal I$ and DVCS² terms **mix at leading order in 1/\mathbf{Q}** in the φ expansion