The unique opportunities of PVDIS at high energy

Mark Dalton Acknowledgements to Nobuo Sato for huge contributions



Importance of Stran

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- Proton spin puzzle: a larş
 The polarized strange c
 the strange distribution
- TMDs: Theoretical descri central to the JLab 12 pro distribution.

PV

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interactions at high energies s which absorb or produce a lepton system of huge mass Q^2 such that the ratio Q^2/s is finite. An observable class of processes meeting this requirement is production of massive lepton pairs in hadron-hadron (dcollisions,⁵ viz.,

$$p+p \rightarrow (\mu^+\mu^-) + \cdots$$

Our remarks apply equally to any colliding pair to the let such as (pp), $(\overline{p}p)$, (πp) , (πp) , (πp) and o final reptons sive since $(\mu^+\mu^-)$, $(e\overline{e})$, $(\mu\nu)$, and (ψ) havior at large χ hereing

What is going on here can be best illustrated in a center-of-mass frame. If a massive state with $Q^2 \sim s$ emerges from one of the colliding protons A or B as in Fig. 1(a), it is impossible to satisfy both energy and momentum conservation in the



can

FIG. 1. one of the case it is partons of parton-an

 $\operatorname{uc-}_{antiparte}$ of the matrix $(d - \overline{a})$ tiparte Viewed f (1) Splitting function say, ann to the left ons sive since in of mass

 $Q^2 = x$

where r

Factorization

hard, calculable part



different processes probe different parts of the distributions

PDF Global Analysis

PDF: probability densities of the longitudinal momentum fraction x of quarks and gluons relative to their parent hadron momentum

- Experimental data (~ 4k data points)
- QCD factorization: PDF (same for PPDF, FF)
- Global analysis

global analyses differ in: input data, parameterization $\eta_{(x,Q^2)}^{1}$ treatment of heavy quarks, value of Ω s, experimental errors treatment, theoretical error estimation.

Fits have been done for 3 decades but the strange remains poorly determined

Momentum sum rule is important constraint



Process	Reaction	Subprocess	PDFs probed	x
letter de letter	$\ell^{\pm}\{p,n\} \to \ell^{\pm} + X$	$\gamma^* q \to q$	q, ar q, g	$x\gtrsim 0.01$
	$\ell^{\pm} n/p \to \ell^{\pm} + X$	$\gamma^* d/u o d/u$	d/u	$x\gtrsim 0.01$
	$ u(\bar{\nu})N \to \mu^-(\mu^+) + X $	$W^*q \to q'$	q, ar q	$0.01 \lesssim x \lesssim 0.5$
	$\nu N \to \mu^- \mu^+ + X$	$W^*s \to c$	S	$0.01 \lesssim x \lesssim 0.2$
	$\bar{\nu}N \to \mu^+\mu^- + X$	$W^*\bar{s} \to \bar{c}$	\overline{s}	$0.01 \lesssim x \lesssim 0.2$
	$e^{\pm}p \to e^{\pm} + X$	$\gamma^*q \to q$	g,q,ar q	$0.0001 \lesssim x \lesssim 0.1$
	$e^+p \to \bar{\nu} + X$	$W^+\{d,s\} \to \{u,c\}$	d,s	$x\gtrsim 0.01$
	$e^{\pm}p \to e^{\pm}c\bar{c} + X$	$\gamma^* c \to c, \gamma^* g \to c\bar{c}$	c,g	$0.0001 \lesssim x \lesssim 0.1$
	$e^{\pm}p \to jet(s) + X$	$\gamma^*g \to q\bar{q}$	g	$0.01 \lesssim x \lesssim 0.1$
	$\overrightarrow{\ell}^{\pm} \{ \overrightarrow{p}, \overrightarrow{d}, \overrightarrow{n} \} \to \ell^{\pm} + X$	$\gamma^* q \to q$	$\Delta q + \Delta ar q$, Δg	$0.003 \lesssim x \lesssim 0.8$
N ₂ N ₁ pp	$pp \to \mu^+ \mu^- + X$	$uar{u}, dar{d} o \gamma^*$	$ar{q}$	$0.015 \lesssim x \lesssim 0.35$
	$pn/pp \to \mu^+\mu^- + X$	$(u\bar{d})/(u\bar{u}) o \gamma^*$	$ar{d}/ar{u}$	$0.015 \lesssim x \lesssim 0.35$
	$p\bar{p}(pp) \rightarrow jet(s) + X$	gg, qg, qq ightarrow 2jets	g,q	$0.005 \lesssim x \lesssim 0.5$
	$p\bar{p} \to (W^{\pm} \to \ell^{\pm} \nu) + X$	$ud \to W^+$, $\bar{u}\bar{d} \to W^-$	$u,d,ar{u},ar{d}$	$x\gtrsim 0.05$
	$pp \to (W^{\pm} \to \ell^{\pm} \nu) + X$	$u ar{d} o W^+$, $d ar{u} o W^-$	$u,d,ar{u},ar{d},(g)$	$x\gtrsim 0.001$
	$p\bar{p}(pp) \to (Z \to \ell^+ \ell^-) + X$	$uu, dd(uar{u}, dar{d}) ightarrow Z$	u,d(g)	$x\gtrsim 0.001$
	$pp \to (W+c) + X$	$gs \to W^- c, g\bar{s} \to W^+ \bar{c}$	$s, ar{s}$	$x \sim 0.01$
	$pp \to tt + X$	gg ightarrow tt	g	$x \sim 0.01$
	$\overrightarrow{n} \rightarrow W^{\pm} + X$	$u_I \bar{d}_P \to W^+ d_I \bar{u}_P \to W^-$	$\Delta u \ \Delta \bar{u} \ \Delta d \ \Delta \bar{d}$	$0.05 \le x \le 0.4$
	$\overrightarrow{p} \overrightarrow{p} \rightarrow \pi + X$	$gg \rightarrow qg, qg \rightarrow qg$	Δg	$0.05 \stackrel{\sim}{_\sim} x \stackrel{\sim}{_\sim} 0.4$
ll	\rightarrow + (\rightarrow \rightarrow) + (\rightarrow +)	* ~ ` ~	$\Delta u \ \Delta \bar{u} \ \Delta d \ \Delta \bar{d}$	0.005 < ~ < 0.5
	$\ell = \{ p, a \} \to \ell = h + X$	$\gamma q \rightarrow q$	Δg	$0.003 \gtrsim x \gtrsim 0.5$
N SIDIS	$\vec{\ell}^{\pm}\{\vec{p},\vec{d}\} \to \ell^{\pm}D + X$	$\gamma^*g \to c\bar{c}$	Δg	$0.06 \lesssim x \lesssim 0.2$

Nucleons in the initial state: Parton Distribution Functions

Emanuele R. Nocera (Oxford)

FFs and Global QCD Fits

June 11 2017 13 / 48

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Strange PDF

Various extractions differ by more than and order of magnitude Many fits not focusing on strange contribution

Results depend on parameterizations and assumptions in global fit.



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Knowledge of PDFs



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Parity violating electron scattering

usually small

/

$$A_{PV} = \frac{G_F Q^2}{4\sqrt{2}\pi\alpha_{\text{em}}} \left[a_2(x) + \frac{1 - (1 - y)^2}{1 + (1 - y)^2} a_3(x) \right]$$

$$a_{2}(x) = -2 g_{A}^{e} \frac{F_{2}^{\gamma Z}(x)}{F_{2}^{\gamma}(x)} = \frac{2 \sum_{q} e_{q} g_{V}^{q} q^{+}(x)}{\sum_{q} e_{q}^{2} q^{+}(x)}$$

$$F_{2}^{\gamma Z}(x) \qquad \sum_{q} e_{q} g_{A}^{q} q^{-}(x)$$

$$a_3(x) = -2g_V^e \frac{F_3^{\gamma L}(x)}{F_2^{\gamma}(x)} = -4g_V^e \frac{\sum_q e_q g_A q^{-\gamma}(x)}{\sum_q e_q^2 q^{+\gamma}(x)}$$

$$g_V^u = \frac{1}{2} - \frac{4}{3} \sin^2 \theta_W$$
$$g_V^d = -\frac{1}{2} + \frac{2}{3} \sin^2 \theta_W$$
$$g_A^e = -\frac{1}{2}$$
$$g_V^e = -\frac{1}{2} + 2 \sin^2 \theta_W$$
$$g_A^u = -g_A^d = \frac{1}{2}$$
$$q^+(x) = q(x) + \bar{q}(x)$$
$$q^-(x) = q(x) - \bar{q}(x)$$

Brady, PHYSICAL REVIEW D 84, 074008 (2011)

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Isolating the strange quarks

Parity-violating DIS allows strange contribution to be isolated, when combined with e.m. p and n DIS data at low/intermediate x

At leading order in QCD

$$F_{2}^{\gamma p} = \frac{4}{9}x(u+\bar{u}) + \frac{1}{9}x(d+\bar{d}+s+\bar{s}) + \cdots$$

$$F_{2}^{\gamma n} = \frac{4}{9}x(d+\bar{d}) + \frac{1}{9}x(u+\bar{u}+s+\bar{s}) + \cdots$$

$$F_{2}^{\gamma Z,p} = (\frac{1}{3} - \frac{8}{9}\sin^{2}\theta_{W})x(u+\bar{u})$$

$$+(\frac{1}{6} - \frac{2}{9}\sin^{2}\theta_{W})x(d+\bar{d}+s+\bar{s}) + \cdots$$

$$\approx \frac{1}{9}x(u+\bar{u}+d+\bar{d}+s+\bar{s}) + \cdots$$
(for $\sin^{2}\theta_{W} \approx \frac{1}{4}$)
$$F_{2}^{\gamma Z,p} = (\frac{1}{3} - \frac{8}{9}\sin^{2}\theta_{W})x(d+\bar{d}+s+\bar{s}) + \cdots$$

$$K_{2}^{\gamma Z,p} = (\frac{1}{3} - \frac{1}{9}\sin^{2}\theta_{W})x(d+\bar{d}+s+\bar{s}) + \cdots$$

$$K_{2}^{\gamma Z,p} = (\frac{1}{3} - \frac{1}{9}\sin^{2}\theta$$

3 equations with 3 unknowns

$$s + \bar{s} = 3(5F_2^{\gamma Z p} - F_2^{\gamma p} - F_2^{\gamma n})$$

in JAM framework.

Parity violating electron scattering

$$A_{PV} = \frac{G_F Q^2}{4\sqrt{2}\pi\alpha_{\rm em}} \left[a_2(x) + \frac{1 - (1 - y)^2}{1 + (1 - y)^2} a_3(x) \right]$$

Beam energy dependence gives access to $a_3(x)$



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PVDIS at high energy | JLab 22 GeV Workshop | 24 January 2023

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Not a new idea!

SLAC-PROPOSAL E149 bis May 4, 1993

DIS-PARITY

Parity Violation in Deep Inelastic Electron Scattering

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ABSTRACT

We propose to measure the parity-violating interference between electromagnetic and weak coupling in deep-inelastic electron scattering from hydrogen and deuterium. The primary physics goals are to test the electro-weak Standard Model and to determine the flavor content of the quark sea. The experiment will determine the Z-quark coupling to an equivalent precision in $\sin^2(\theta_w)$ of approximately 0.003 (1.2%), and will be quite sensitive to the possible existence of certain types of new physics beyond the Standard Model. The experiment will also measure two new combinations of quark distribution functions with errors of about 0.05. The deuterium case is particularly interesting because it directly isolates the strange sea quark contributions.

The experiment consists of scattering polarized electrons from long liquid hydrogen and deuterium targets in End Station A. The scattered electrons are detected in two large solid angle magnetic spectrometers centered around scattering angles of 4.25 and 6.5 degrees for an incident beam energy of 29.1 GeV. At E = 48.6 GeV, the 6.5° spectrometer is moved to 2.75°. The spectrometers are entirely made up of existing magnets. The detectors consist primarily of lead glass arrays that can be made from the surplus ASP pieces. The beam is required to have a polarization of 80% or more, and high current $(4 \times 10^{11} e^{-1})$ per beam pulse at 120 Hz non-SLED, 2×10^{11} SLED). The combination of higher beam polarization, higher beam energy, and larger spectrometer solid angle provides a factor of eight greater sensitivity than the original SLAC experiment on deuterium. The addition of hydrogen allows more sensitive studies of the flavor content of the quark sea.

A total of twelve weeks of running time is requested for an initial run at 29.1 GeV, of which one third will be used for checkout, calibrations, and background studies, and two thirds for data taking. A subsequent run of ten weeks at 48.6



GeV is also requested. In a relatively modest running time, this experiment will provide significantly improved precision on electroweak physics at low Q^2 and on the flavor content of the quark sea. We emphasize that this experiment can only - be undertaken at SLAC.

Experimental Overview

Beam: 40 uA @ 22 GeV. 40 cm liquid H₂ or D₂ 100 days of running

Luminosity (instantaneous) 4.3x10³⁸ cm⁻²s⁻²

Luminosity (integrated) 4.5x10⁶ fb⁻¹

cf EIC Yellow Report proton ~ 100 fb⁻¹ deuteron ~ 10 fb⁻¹



SoLID spectrometer and detectors optimized for PVDIS experiment: high luminosity, small systematics, polarimetry, beam quality and control

The only place in the world where this is possible.

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SoLID 22 GeV



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m JLab}$ 22/

Corrections and Interpretation

Target Mass Corrections



Implemented in some PDF extractions

Nuclear Corrections

Higher Twist Corrections

non-pert. parton correlations



Implemented in some PDF extractions

binding, Fermi motion, off-shellness calculations are largely modeldependent

Final State Interactions



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QED

Recent recognition of this issue. New formalism to include effects in global fit.

Interpretation - QED effects

QED "corrections" depend on the input hadronic tensor Not possible to construct model-independent QED RC corrections Need to include QED in global analysis

Size of RCs quickly become extremely large at small x



Liu, Melnitchouk, Qiu, Sato, JHEP11(2021)157 Liu, Melnitchouk, Qiu, Sato, PHYSICAL REVIEW D 104, 094033 (2021)

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Interpretation - QED effects

RCs quickly become extremely large at small x, resonance region begins at higher x. Limits range of useful data Higher beam energy and higher Q² help a lot.



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Summary

Strange PDF is the least well known

Parity Violating DIS provides a unique new combination of PDFs and has much greater sensitivity to the strange

SoLID @ 22 GeV with 100 days can provide a very strong constraint on $s + \overline{s}$.

The interpretation of data is easier at higher energies and Q².

Extras



Connection to $\sin^2 \theta_W$



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Accessing Strange Quarks

 di-muon production in neutrino-nucleus scattering 	0.01 < x < 0.2
 W and Z rapidity distributions 	x >~ 0.001
 LHC W+c production 	x ~ 0.01
 Semi-inclusive K production: not included in global fits (fragmentation) 	0.02 < x < 0.6
 Parity Violating electron scattering 	0.1 < x < 0.5

It's still not clear whether the strange sea is as big as the up and down sea $(k \sim 1)$ or half as big $(k \sim 0.5)$

 $k = \frac{2s}{(\bar{u} + \bar{d})}$

Neutrino nucleus

di-muon production in neutrinonucleus scattering tags strangeness



- initial state nuclear modifications of the PDFs themselves (partly under control—nuclear PDF fits)
- final state interactions
 - medium-induced gluon bremsstrahlung;
 - final state suppression of charm production (measured at RHIC significantly larger than perturbative calculations for QGP)

Neutrino nucleus

Neutrino and electron scattering data from iron do not agree below x~0.1 Sensitive to different nuclear effects?



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Proton-proton

ATLAS fits get k \approx I using collider only data



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Proton F₂



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SoLID acceptance

SOLID acceptance @ 6.6 GeV



SOLID acceptance @ 11 GeV

SOLID acceptance @ 22 GeV



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