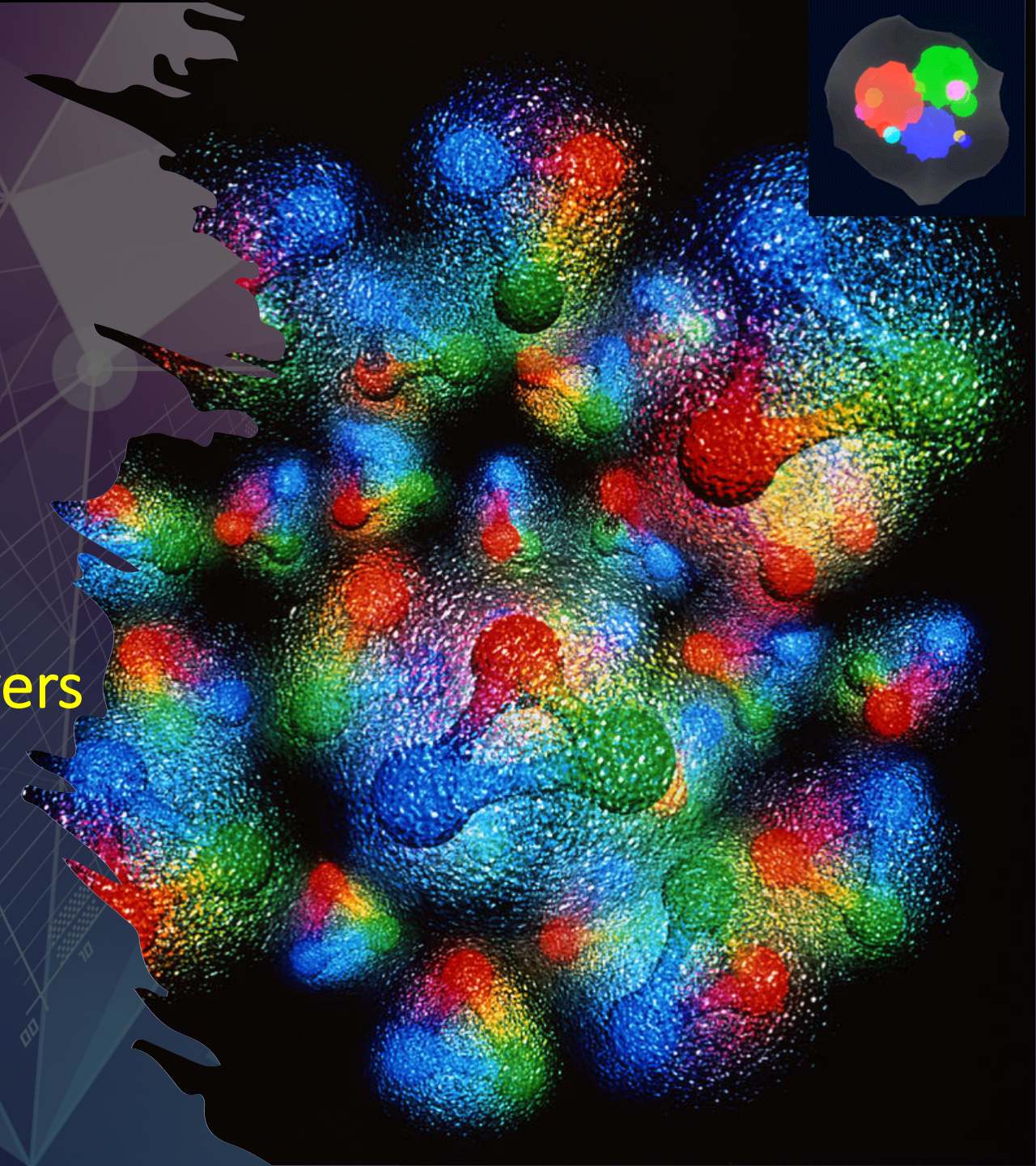


# The spin-directed momentum transfer model of polarized hyperon production

G. Goldstein, S. Liuti and D. Sivers  
(presented by S.Liuti)



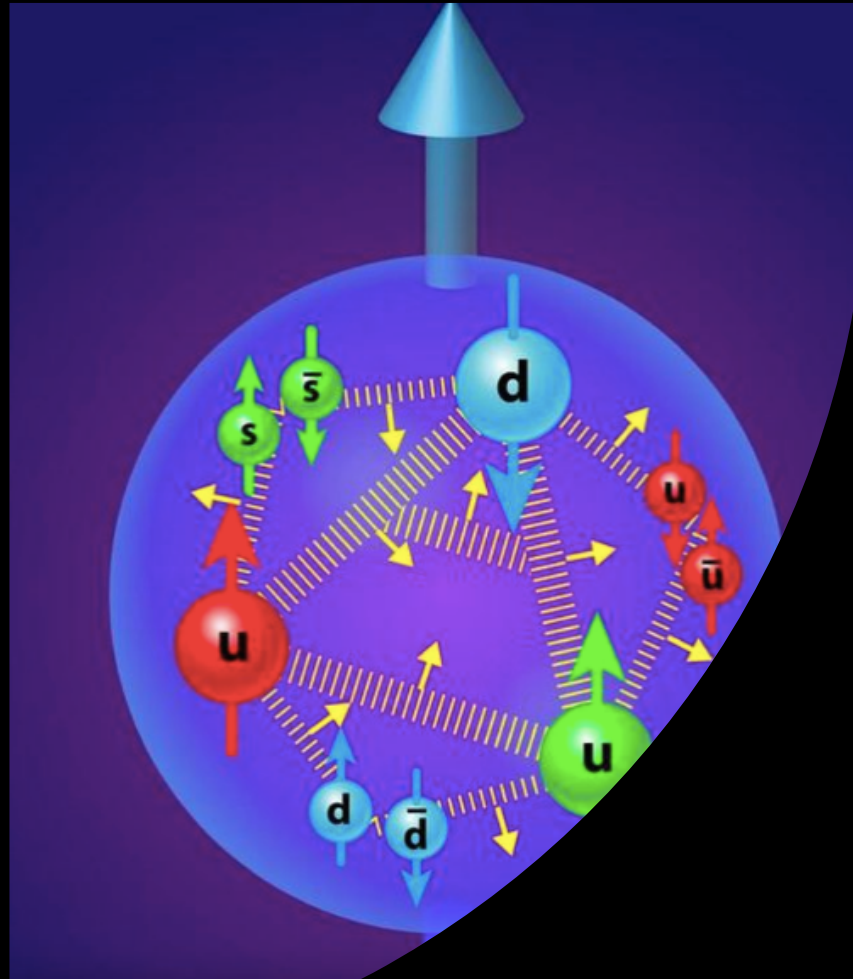
Collaborators: Gary Goldstein and Dennis Sivers

UVA group: Joshua Bautista\*, Brandon Kriesten\*\*, Zaki Panjsheeri\*

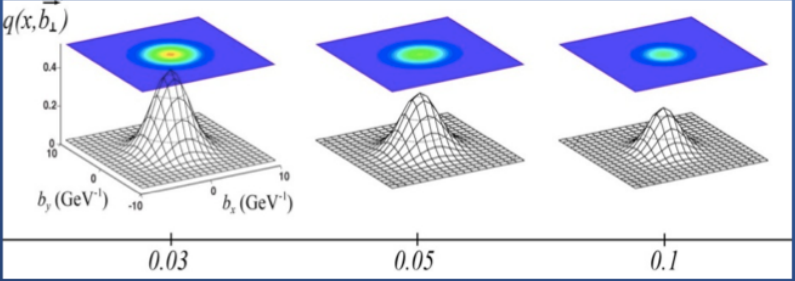
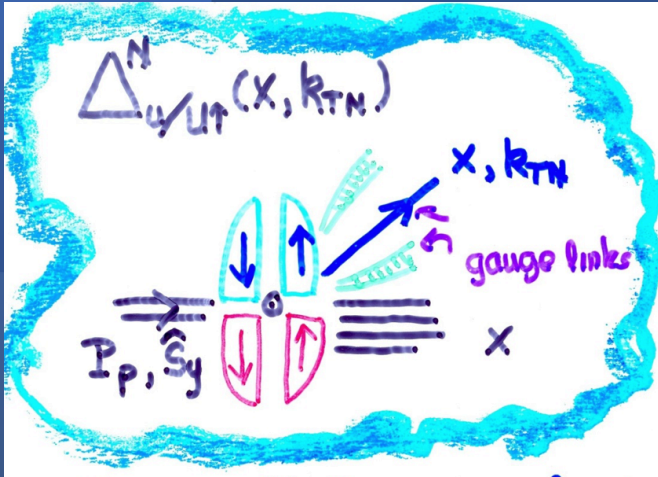
\*University of Virginia

\*\*SURA/Jefferson Lab

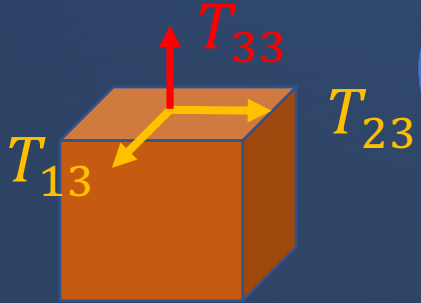
- 
- The proton spin remains a mystery
  - The flavor composition of the nucleon form factor remains a mystery
  - The perspective of the EIC has opened more possibilities for studying different types exclusive experiments beyond DVCS-type from proton targets
  - These new channels will allow us to understand nucleon spin is divided between gluons and orbital angular momentum



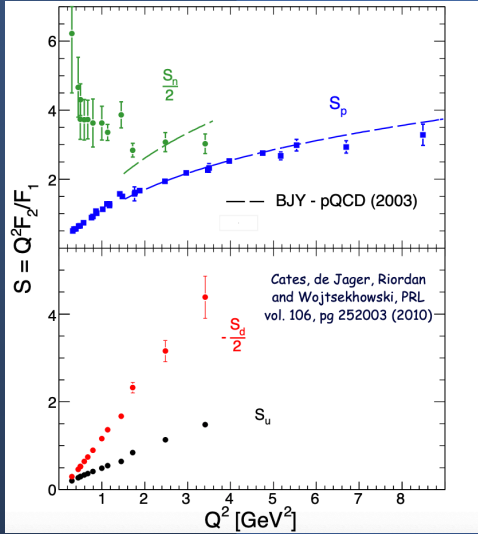
- How is hadronic spin generated as a collective phenomenon from the proton/neutron constituents?
- **What is the role of gluons?**
- How do we test this dynamics?



GPDs



Energy Momentum Tensor



Flavor form factors

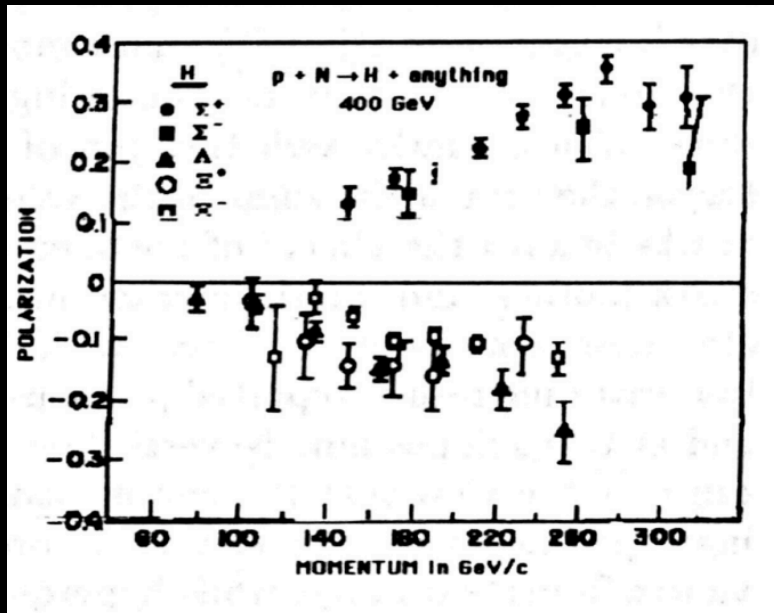
This means quark and gluon distributions

# Transverse Polarized Hyperon Production



Experiments on hyperon production performed during more than three decades have shown that the hyperon polarization is significant

From BNL, Fermilab ....



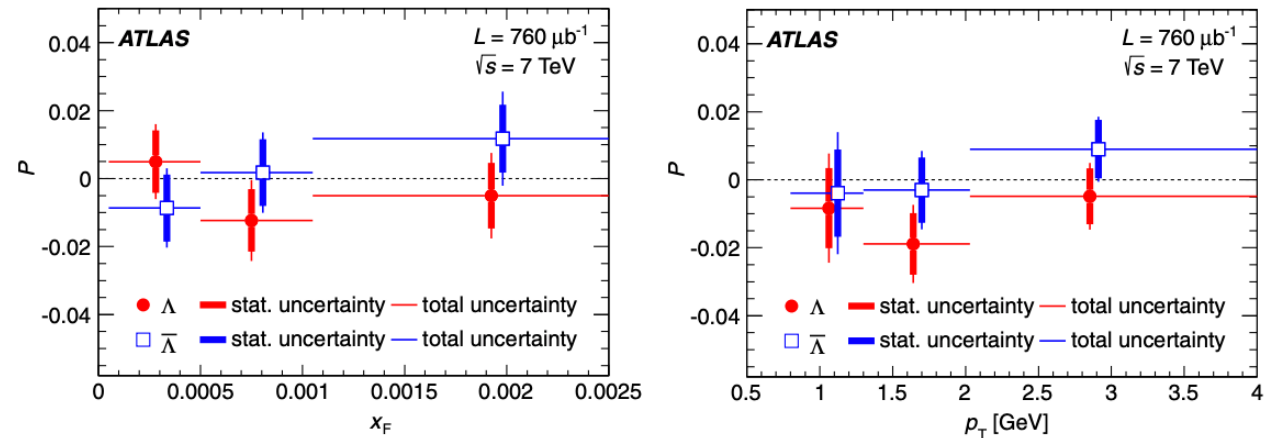
$P_H$

$H = \Lambda, \Sigma, \Xi, \dots$

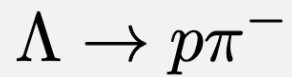
To the LHC ....

G. Aad *et al.*

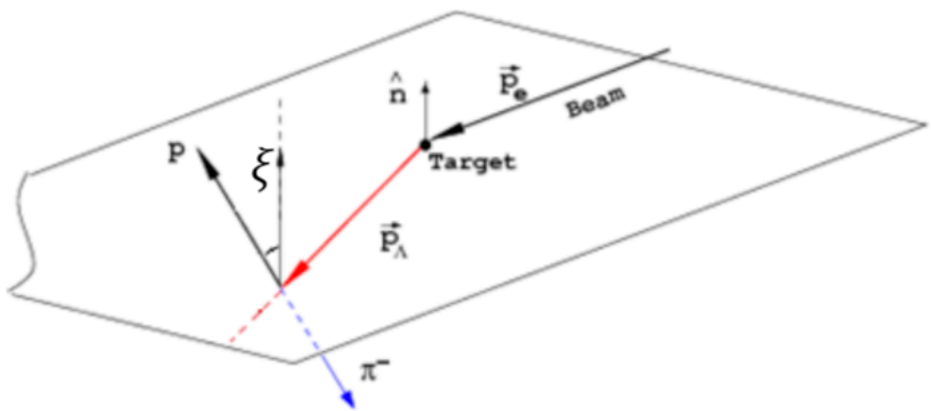
PHYSICAL REVIEW D **91**, 032004 (2015)



$$P_\Lambda = -0.010 \pm 0.005$$



Due to their self-analyzing weak decay, the polarization of  $\Lambda$  hyperons  $P_\Lambda$  can be accessed by inspecting the angular distribution of the protons produced in the decay



$$\frac{1}{N} \frac{dN}{d \cos \xi} = 1 + \alpha_\Lambda P_\Lambda \cos \xi$$

SSA

$$P_\Lambda = \frac{d\sigma^{lp \rightarrow l' \Lambda^\uparrow X} - d\sigma^{lp \rightarrow l' \Lambda^\downarrow X}}{d\sigma^{lp \rightarrow l' \Lambda^\uparrow X} + d\sigma^{lp \rightarrow l' \Lambda^\downarrow X}}$$

$$P_{\Lambda} = \frac{\Im m(\mathcal{M}_{nonflip}^* \mathcal{M}_{flip})}{\sigma_{unpol}}$$

- 1) The flip amp is different from 0
- 2) The non-flip amp has opposite **phase** to the flip one

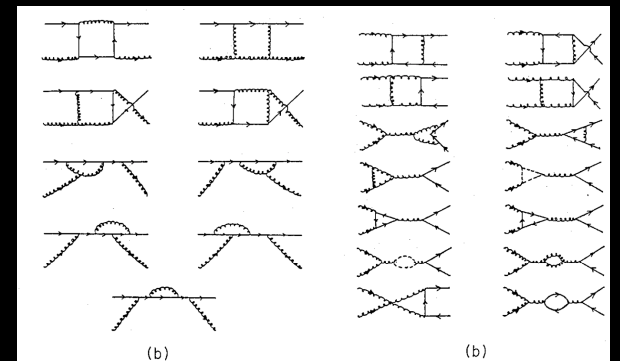
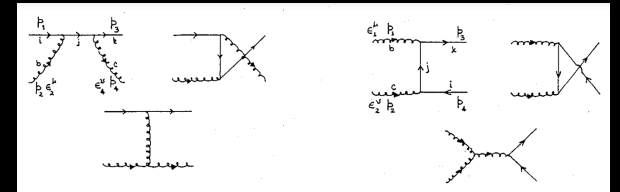
Which mechanisms can generate the quantum interference that can produce transverse polarization?



# 1. The perturbative approach

Dharmaratna, Goldstein '90's

- Hard scattering of partons
- gluon fusion produces a polarized s quark

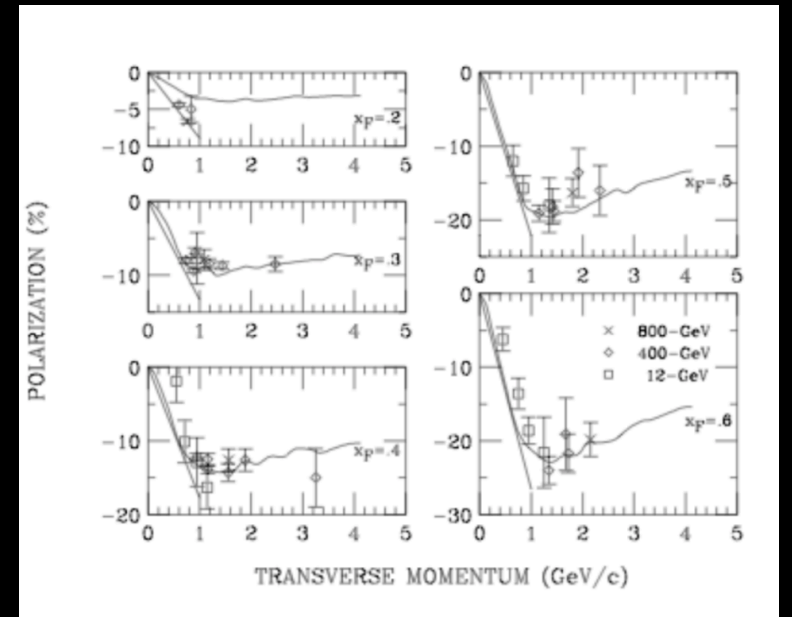


The polarization depends linearly on the quark mass  $\rightarrow$  its value increases for heavy quarks

$$\mathcal{P} = \alpha_s \frac{m(p^2 - k^2 \cos^2 \theta)}{24k \sin \theta D} \left( (N_1 + N_2)Y_+ + (N_1 - N_2)Y_- + N_3 \ln \frac{p - k}{p + k} + N_4 + 18k^3 \sin^2 \theta \cos \theta (\Sigma_1 + \Sigma_2) \right)$$

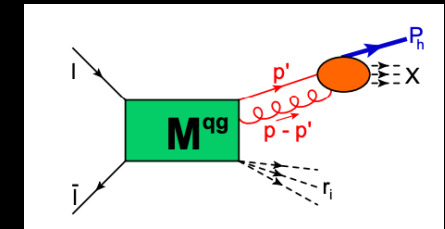
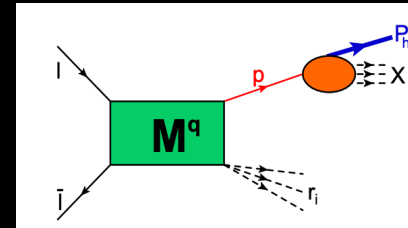
To make quantitative predictions the model needs to be improved with a semiclassical recombination mechanism

- 1)  $x_F$  of quark is smaller than experimental one but gets boosted by rescaling from non perturbative mechanism
- 2) It works in the assumption of TMD factorization
- 3) It contributes to the **current fragmentation** region



## 2. Twist three interference from quark current fragmentation

Pitonyak et al., 2019



- 1) A phase difference is introduced by the twist three amplitude for the current fragmentation
- 2) It works in the assumption of TMD factorization
- 3) It contributes to the **current fragmentation** region

# Kane Pumplin Repko (KPR) Factorization

In the limit  $m_q \rightarrow 0$  there exists within QCD perturbation theory a symmetry that is strongly broken in the nonperturbative sector of the theory.

The existence of this symmetry ensures that all parity-even single spin asymmetries in hard-scattering processes can be absorbed into the transverse-momentum dependence of the **fragmentation functions**.

**D. Sivers**

$$A_N = \alpha_S \frac{m_q}{\sqrt{s}} f_{TN}(\theta_{CM}, k_{TN})$$

G. Kane, J. Pumplin and W. Repko, PRL 41 (1978)

$$k_{TN} = \mathbf{k}_C \cdot (\mathbf{s} \times \mathbf{p}_A)$$

Spin-directed momentum transfer

### 3. Target fragmentation through $^3P_0$ mechanism

Sivers (2011)

Goldstein, SL, Sivers (2023)

Based on KPR factorization: Spin flip dynamics generating final hadron spin is in the target fragmentation function (**fracture function**)

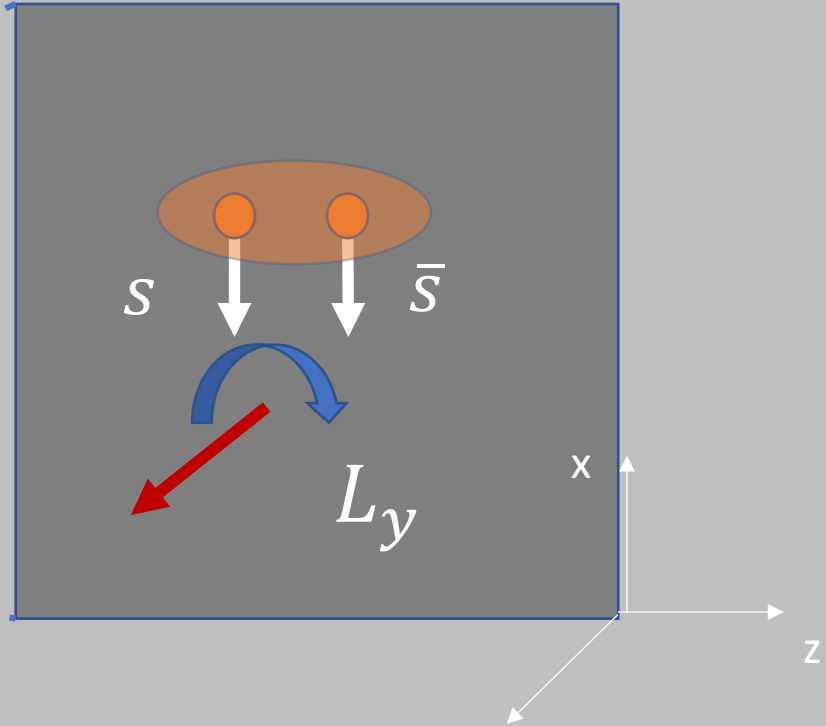
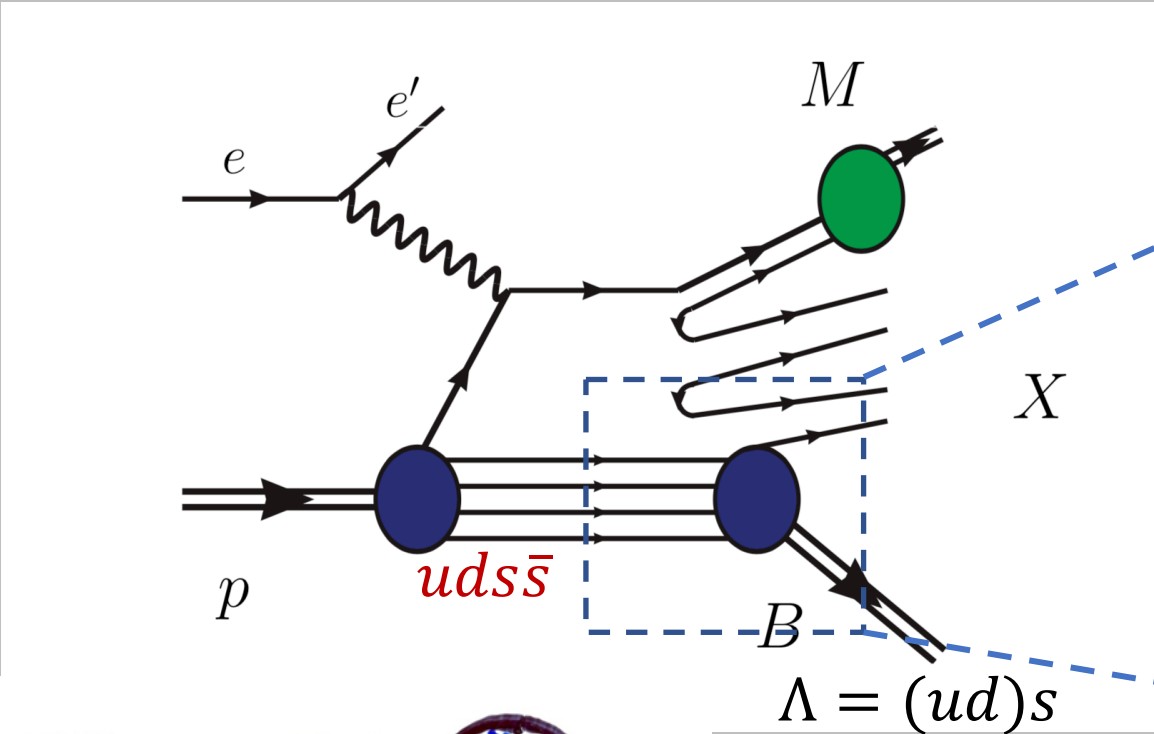
**Spin-directed momentum transfer framework** allows us to define kinematic variables involved in the underlying spin-orbit dynamics.

To produce a net hyperon transverse spin the quark **orbital angular momentum** between the initial and final states must differ by one unit.

**Four distinct spin-directed momentum transfer configurations/polarized fracfuns**

$$\begin{aligned} &\Delta M_{B,p^\uparrow}^q(x, k_{TN}; z, p_T^B) \\ &\Delta M_{B,p^\uparrow}^{q^\uparrow}(x, k_T; z, p_{TN}^B) \\ &\Delta M_{B^\uparrow,p^\uparrow}^q(x, k_T; z, p_{TN}^B) \\ &\Delta M_{B^\uparrow,p}^{q^\uparrow}(x, k_{TN}; z, p_T^B) \end{aligned}$$

# String breaking mechanism



D. Sivers

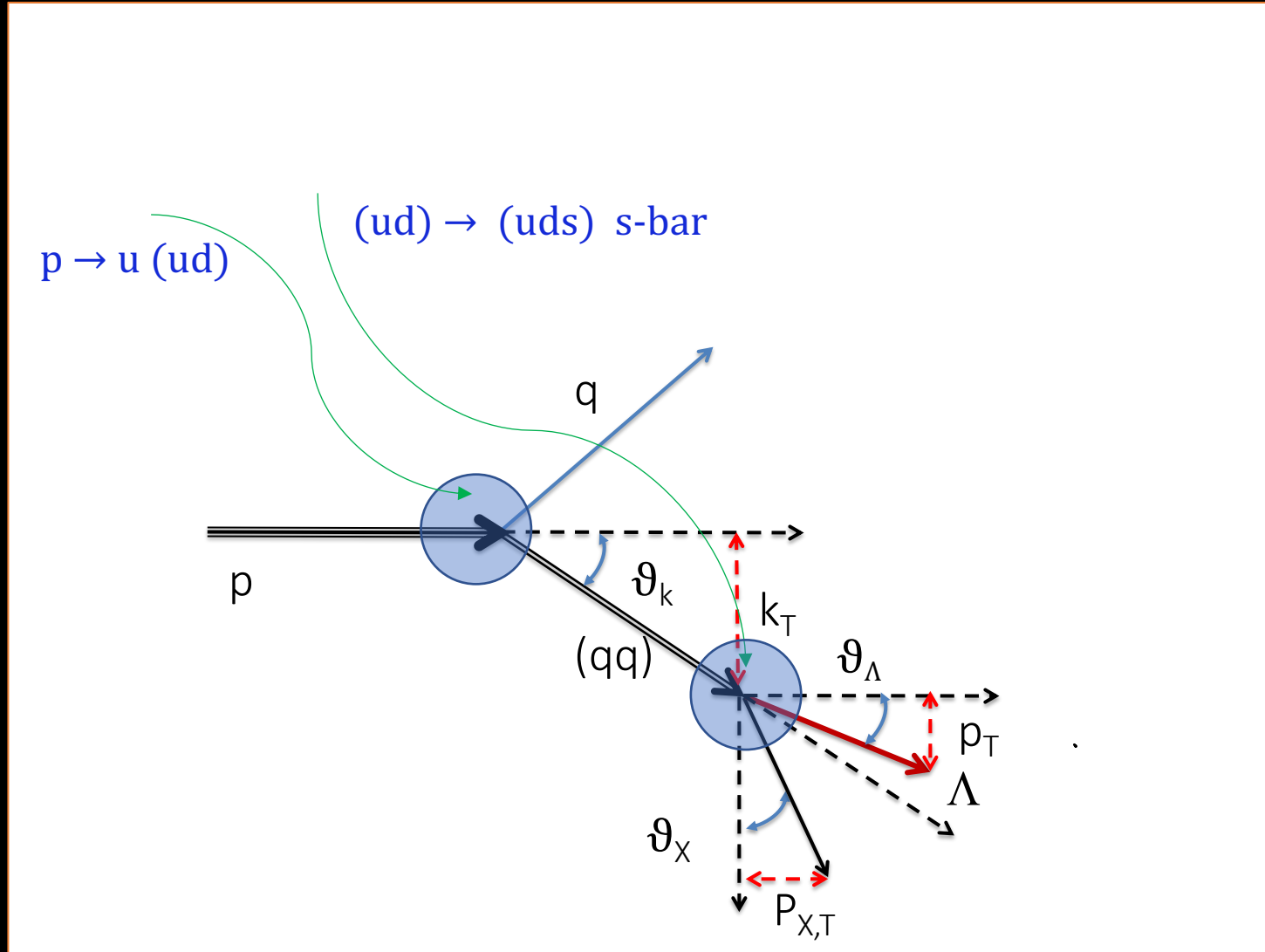
The  $(ud)$  diquark is produced with spin  $S = 0$

As the flux tube stretches an  $ss^-$ -pair is produced with transverse momentum,  $k_\perp$

The  $ss^-$  - pair have their spins aligned to balance the OAM

Rotational motion generates the breaking of the string

# Modeling the $^3P_0$ string breaking mechanism: Double Spectator Model



# Double Spectator Model

$$f_{\Lambda_H \Lambda_\gamma, \Lambda} = \sum_{\lambda} g_{\lambda}^{\Lambda_\gamma} \otimes \mathcal{F}_{\Lambda \lambda}^{\Lambda_H}(P, k)$$

Fracture function  
(see F. Benkmothar talk)

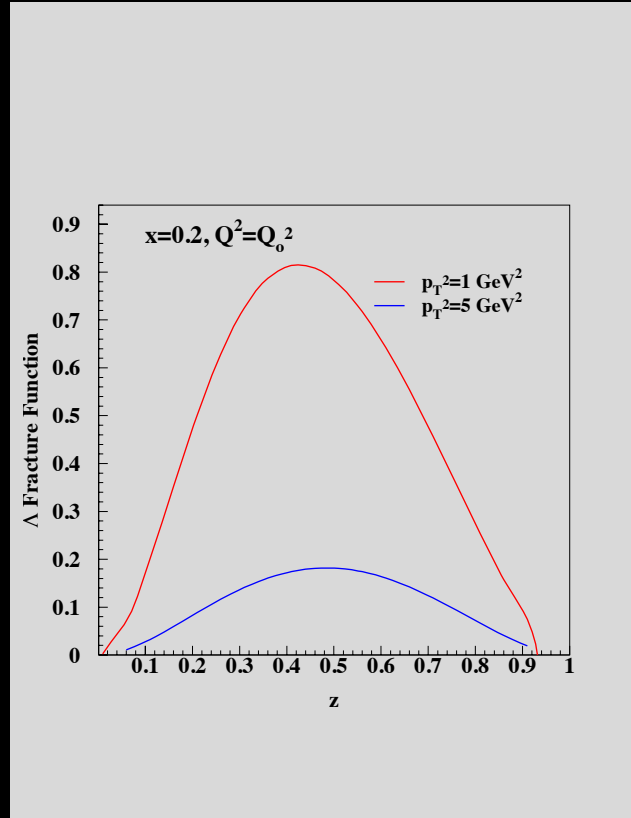
$$\mathcal{F}_{\Lambda, \lambda}^{\Lambda_H}(x, k_T, z, p_T, Q^2) = A_{\lambda, \Lambda} \sum_{\Lambda_X} B_{\Lambda_X}^{\Lambda_H}$$

$$A_{\Lambda, \lambda_q} = |\phi_{\lambda_q, \Lambda}(k, p)|^2$$

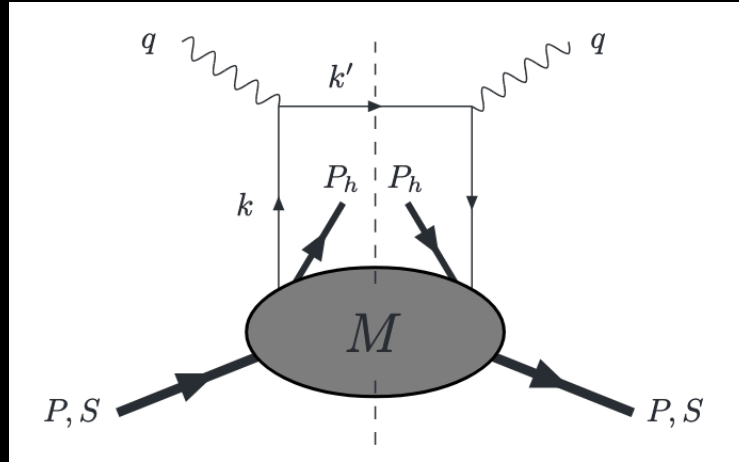
$$\phi_{\Lambda, \lambda_q}(k, p) = \Gamma(k) \frac{\bar{u}(k, \lambda_q) U(p, \Lambda)}{(k^2 - m^2)((p - k)^2 - M_{qq}^2)}$$

$$B_{\Lambda_X}^{\Lambda_B} = \tilde{\phi}_{\Lambda_X, \Lambda_B}^*(p_X, p_B) \tilde{\phi}_{\Lambda_X, \Lambda_B}(p_X, p_B)$$





# Target fragmentation cross section



$$\mathcal{M}^{[\gamma^-]} = \hat{u}_1 + \frac{\mathbf{P}_{2\perp} \times \mathbf{S}_\perp}{m_2} \hat{u}_{1T}^h + \frac{\mathbf{k}_\perp \times \mathbf{S}_\perp}{m_N} \hat{u}_{1T}^\perp + \frac{S_\parallel (\mathbf{k}_\perp \times \mathbf{P}_{2\perp})}{m_N m_2} \hat{u}_{1L}^{\perp h},$$

$$\mathcal{M}^{[\gamma^- \gamma_5]} = S_\parallel \hat{l}_{1L} + \frac{\mathbf{P}_{2\perp} \cdot \mathbf{S}_\perp}{m_2} \hat{l}_{1T}^h + \frac{\mathbf{k}_\perp \cdot \mathbf{S}_\perp}{m_N} \hat{l}_{1T}^\perp + \frac{\mathbf{k}_\perp \times \mathbf{P}_{2\perp}}{m_N m_2} \hat{l}_1^{\perp h},$$

$$\mathcal{M}^{[i\sigma^{i-} \gamma_5]} = S_\perp^i \hat{t}_{1T} + \frac{S_\parallel P_{2\perp}^i}{m_2} \hat{t}_{1L}^h + \frac{S_\parallel k_\perp^i}{m_N} \hat{t}_{1L}^\perp$$

## Quark polarization

	U	L	T
U	$\hat{u}_1$	$\hat{l}_1^{\perp h}$	$\hat{t}_1^h, \hat{t}_1^\perp$
L	$\hat{u}_{1L}^{\perp h}$	$\hat{l}_{1L}$	$\hat{t}_{1L}^h, \hat{t}_{1L}^\perp$
T	$\hat{u}_{1T}^h, \hat{u}_{1T}^\perp$	$\hat{l}_{1T}^h, \hat{l}_{1T}^\perp$	$\hat{t}_{1T}^h, \hat{t}_{1T}^{\perp h}$ $\hat{t}_{1T}^\perp, \hat{t}_{1T}^{\perp h}$

Nucleon polarization

M. Anselmino et al., Phys. Lett. B. 706 (2011), 46-52, [hep-ph] 1109.1132

# Conclusions

- The proton is a complex, multi-body environment, characterized by vortical motion of quarks and gluons
- We described a non perturbative QCD string breaking mechanism for polarized  $\Lambda$  production based on KPR factorization

## In progress...

- Full calculation of fracture functions cross section
- More to develop: predictions for other hyperon,  $\Sigma$ ,  $\Xi$ , ... polarizations and anti-Baryons
- Adding gluonic component to model