

A continuum Schwinger method for pion GPDs

Observable implications at future electron ion colliders

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Generalized parton distributions (GPDs)

(GPD) – Generalized parton distributions:

Non-local quark and gluon operators, evaluated between hadron states in non-forward kinematics and projected onto the light front.

[Fortsch.Phys.:42(1994)101]

[Phys.Lett.B:380(1996)417]

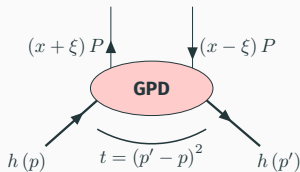
[Phys.Rev.D:55(1197)7114]

Example: Leading twist chiral-even GPDs of a spinless hadron.

$$H_{\pi}^q(x, \xi, t) = \frac{1}{2} \int \frac{d\lambda}{2\pi} e^{i\lambda x} \langle \pi(p') | \bar{\psi}^q(-\lambda n/2) \gamma^{\mu} \psi^q(\lambda n/2) | \pi(p) \rangle n_{\mu},$$

$$H_{\pi}^g(x, \xi, t) = \int \frac{d\lambda}{2\pi} e^{i\lambda x} \langle \pi(p') | G^{\mu\alpha}(-\lambda n/2) G_{\alpha}^{\nu}(\lambda n/2) | \pi(p) \rangle n_{\mu} n_{\nu}$$

GPDs: Definition and properties

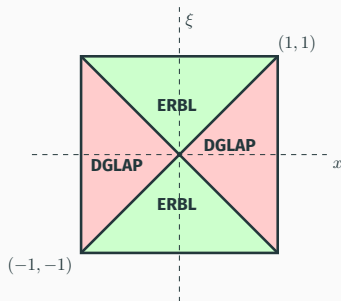


- x : Momentum fraction of P .
 ξ : Fraction of momentum longitudinally transferred.
 t : Momentum transfer.

Kinematics:

[Phys.Rept:388(2003)41]

- **DGLAP** ($|x| > |\xi|$):
Emits/takes a quark ($x > 0$)
or antiquark ($x < 0$).
- **ERBL**: ($|x| < |\xi|$):
Emits pair quark-antiquark.



GPDs: Definition and properties

Probabilistic interpretation:

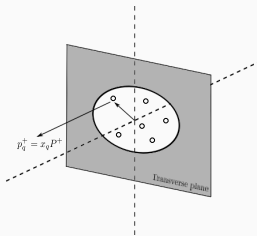
Probability amplitude of finding a parton at a given position in transverse plane carrying a momentum fraction “ x ” of the hadron’s averaged light-cone momentum.

[Phys.Rev.D:66(2002)119903]

Striking properties:

- PDFs as forward limit.
- Electromagnetic and gravitational FFs as Mellin moments.
- Parametrize DVCS amplitudes through CFFs.

[e.g. Phys.Rev.D:59(1999)07009, Phys.Rev.D:56(1997)2982]



“3D picture” of hadrons

GPDs: Definition and properties

- **Support:**

[Phys.Lett.B:428(1998)359]

$$(x, \xi) \in [-1, 1] \otimes [-1, 1]$$

Analyticity/Causality

- **Polynomiality:** Order- m Mellin moments are degree- $(m+1)$ polynomials in ξ .

[X.Ji-JPG:1181(24)1998, A.Radyushkin-PLB:81(449)1999]

$$\int_{-1}^1 dx x^m H(x, \xi, t) = \sum_{\substack{k=0 \\ k \text{ even}}}^{m+1} c_k^{(m)}(t) \xi^k$$

Lorentz invariance

- **Positivity:**

[P.V.Pobylitsa-PRD:114015(65)2002, B.Pire et al.-EPJC:103(8)1999]

$$|H^q(x, \xi, t=0)| \leq \sqrt{q \left(\frac{x+\xi}{1+\xi} \right) q \left(\frac{x-\xi}{1-\xi} \right)}, \quad |x| \geq \xi$$

Positivity of Hilbert space norm

- **Low energy soft-pion theorem**

[M.V.Polyakov-NPB:231(555)1999, C.Mezrag et al.-PLB:190(741)2015]

PCAC/Axial-Vector WTI

GPD modeling: State of the art

1. Overlap representation

[Nucl.Phys.B596(2001)33]

Based on LFWFs, $\Psi^q(x, k_{\perp}^2)$

} Polynomiality ?
Positivity ✓

2. Double Distribution representation

[Fortsch.Phys.:42(1994)101, hep-ph/0101225]

Relying on Radon transform, \mathcal{R}

} Polynomiality ✓
Positivity ?

Problem: Different modeling strategies and different issues

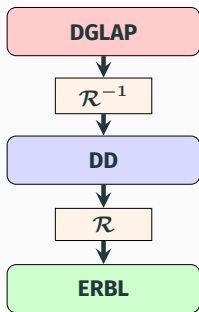
Solution

Covariant extension: Given a DGLAP GPD, the covariant extension allows for computing the corresponding ERBL GPD such that polynomiality is satisfied. [Eur.Phys.J.C:77(2017)12,906]

Covariant extension

Covariant extension: Given a DGLAP GPD, the covariant extension allows for computing the corresponding ERBL GPD such that polynomiality is satisfied. [Eur.Phys.J.C:77(2017)12,906]

$$H(x, \xi) = \mathcal{R} [h(\beta, \alpha)] + \frac{1}{|\xi|} D^+ \left(\frac{x}{\xi} \right) + \text{sign}(\xi) D^- \left(\frac{x}{\xi} \right)$$



1. Build positive DGLAP GPD
2. Covariant extension: ERBL GPD

GPD properties			
Support [Diehl-PLB(1998)]	✓	Positivity [Poly.-PRD(2002), Pire-EPJC(1999)]	✓
Polynomiality [Ji-JPG(1998), Radyu.-PLB(1999)]	✓	Soft-pion [Poly.-NPB(1999), Mezr.-PLB(2015)]	✓

3. Soft pion theorem: fix $D^\pm(\alpha, 0)$

GPD modeling: DGLAP region (I)

1. Overlap representation [Nucl.Phys.B596(2001)33]

$$H^q(x, \xi, t)|_{|x| \geq \xi} = \int \frac{d^2 k_{\perp}}{16\pi^3} \Psi^{q*}(x_-, k_{\perp}^2, -) \Psi^q(x_+, k_{\perp}^2, +)$$

2. Assume factorization of the LFWF

$$\Psi^q(x, k_{\perp}^2) \propto \varphi(x) \phi(k_{\perp}^2)$$

↓ (Overlap rep.)

$$H^q(x, \xi, t)|_{|x| \geq \xi} = \sqrt{q \left(\frac{x-\xi}{1-\xi} \right) q \left(\frac{x+\xi}{1+\xi} \right)} \Phi(x, \xi, t)$$

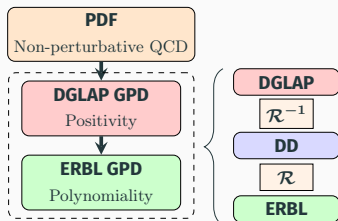
↓ ($t=0$)

$$H^q(x, \xi, 0)|_{|x| \geq \xi} = \sqrt{q \left(\frac{x-\xi}{1-\xi} \right) q \left(\frac{x+\xi}{1+\xi} \right)}$$

[Phys. Lett. B 815 (2021) 136158]

[Chin. Phys. C 46 (2022) 1, 013105]

[Phys. Rev. D 105 (2022) 9, 094012]



PDFs are the sole ingredient needed to build positive DGLAP GPD

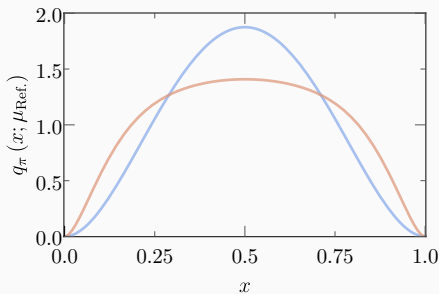
GPD modeling: DGLAP region (II)

1. Under certain PTIR, chiral symmetry allows to factorize LFWF:

$$\Psi_{\pi}^{\lambda_1 \lambda_2}(x, k_{\perp}^2) = \sqrt{q_{\pi}(x)} \frac{i^{\lambda_1 \lambda_2} M^2}{(k_{\perp}^2 + M^2)^2}$$

2. Positive pion GPD [Phys. Rev. D 105 (2022) 9, 094012]

$$H_{\pi}^q(x, \xi, t)|_{\text{DGLAP}} = \frac{\sqrt{q_{\pi}(x_-) q_{\pi}(x_+)}}{(1+z^2)^2} \left[3 + \frac{1-2z}{1+z} \frac{\operatorname{arctanh}\left(\sqrt{\frac{z}{1+z}}\right)}{\sqrt{\frac{z}{1+z}}} \right]$$
$$z = -t(1-x)^2 / 4M^2(1-\xi^2)$$



Two models:

- Algebraic model

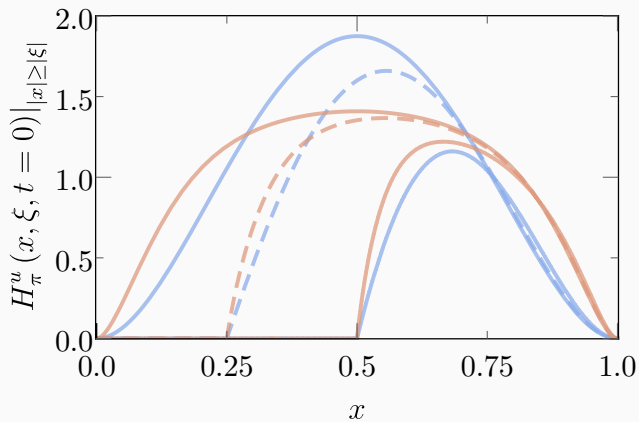
$$q_{\pi}(x) = 30x^2(1-x)^2$$

- Realistic model (DSE)

[Phys.Rev.D:101(2020)5,054014]

$$q_{\pi}(x) = \mathcal{N}_q x^2(1-x)^2 \times \left[1 + \gamma x(1-x) + \rho \sqrt{x(1-x)} \right]$$

GPD modeling: DGLAP region (III)



GPD modeling: ERBL region (I)

Covariant extension:

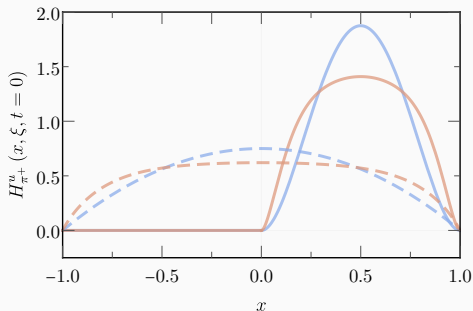
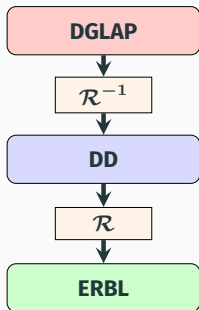
$$H^q(x, \xi, t) = \mathcal{R}[h(\beta, \alpha, t)] + \frac{1}{|\xi|} D^+ \left(\frac{x}{\xi}, t \right) + \text{sgn}(\xi) D^- \left(\frac{x}{\xi}, t \right)$$

Fix D-terms with soft pion theorem:

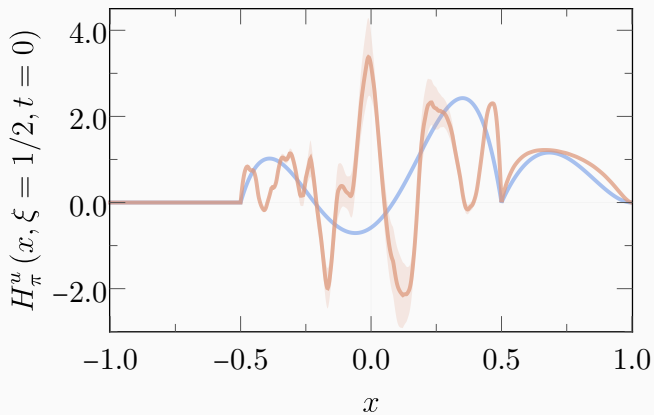
[Nucl.Phys.B:555(1999)231, Phys.Lett.B:741(2015)190]

$$H_{\pi}^{I=0}(x, \xi, t) \Big|_{\xi=1, t=0} = H_{\pi}(x, \xi, t) - H_{\pi}(-x, \xi, t) \Big|_{\xi=1, t=0} = 0$$

$$H_{\pi}^{I=1}(x, \xi, t) \Big|_{\xi=1, t=0} = H_{\pi}(x, \xi, t) + H_{\pi}(-x, \xi, t) \Big|_{\xi=1, t=0} = \varphi_{\pi} \left(\frac{1+x}{2} \right)$$



GPD modeling: ERBL region (II)



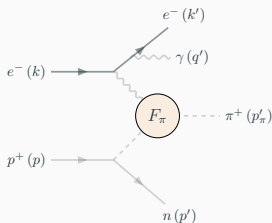
Phenomenology of pion GPDs: Sullivan process (I)

Sullivan process [Phys.Rev.D:5(1972)1732]

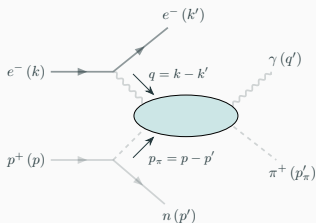
Deep inelastic electron-proton scattering with πn fixed final states.

One-pion-exchange approximation: [Eur.Phys.J.C:58(2008)179]

- $|t|^{Max.} = 0.6 \text{ GeV}^2$
 - Factorization: $\sigma_L^{\gamma^*} \gg \sigma_{\perp}^{\gamma^*}$
- } Met at EIC [arXiv[phys.ins-det]2103.05419]

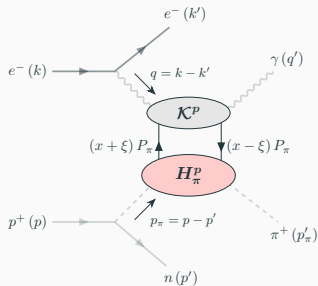
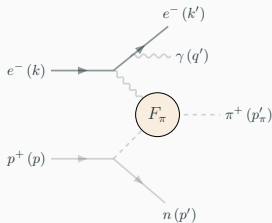


Employed for EFFs
[Phys.Rev.C:78(2008)045203]



Provide access to pion GPDs

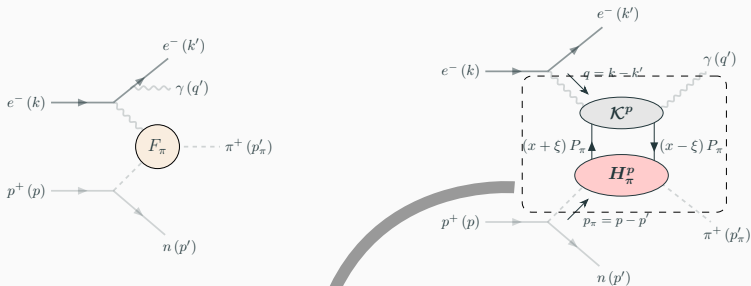
Phenomenology of pion GPDs: Sullivan process (II)



$$\mathcal{M}_{\text{Sullivan}} = i\sqrt{2}g_{\pi NN}\bar{u}_\sigma(p')\gamma_5 u_\sigma(p)\frac{i}{t - m_\pi^2}F(t; \Lambda)\mathcal{M}_{e\pi \rightarrow e\gamma\pi}$$

$$\mathcal{M}_{e\pi \rightarrow e\gamma\pi} = \mathcal{M}_{\text{BH}} + \mathcal{M}_{\text{DVCS}} \Rightarrow |\mathcal{M}_{e\pi \rightarrow e\gamma\pi}|^2 = |\mathcal{M}_{\text{DVCS}}|^2 + |\mathcal{M}_{\text{BH}}|^2 \pm \mathcal{I}(\lambda)$$

Phenomenology of pion GPDs: Sullivan process (II)



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DVCS amplitudes are parametrized by hadron GPDs through CFFs.

[Phys.Rev.D:55(1997)7114]

$$\mathcal{M}_{\text{DVCS}} \propto \mathcal{H}_\pi(\xi, t; Q^2) = \sum_{p=q,g} \int_{-1}^1 \frac{dx}{\xi} \mathcal{K}^p\left(\frac{x}{\xi}, \frac{Q^2}{\mu_F^2}, \alpha_S(\mu_F^2)\right) H_\pi^p(x, \xi, t; \mu_F^2)$$

Phenomenology of pion GPDs: The path towards DVCS

1. Start from state-of-the-art **models** for pion GPDs.
2. Leading order **scale evolution** for GPDs.
3. Convolution with coefficient functions: **CFFs**.

GPD Modeling

[Phys.Lett.B 815 (2021) 136158]

[Chin. Phys. C 46 (2022) 1, 013105]

[Phys. Rev. D 105 (2022) 9, 094012]

QCD evolution

[Eur. Phys. J. A 58 (2022) 1, 10]

[Phys. Rev. D 105 (2022)9, L091502]

[Eur. Phys. J. C 82 (2022) 10, 888]

LO/NLO CFFs

[Phys. Rev. D 105 (2022) 9, 094012]

[Phys. Rev. Lett. 128 (2022) 20,

202501]

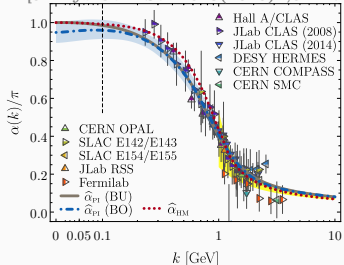
Phenomenological study of pion's structure
Phys. Rev. Lett. 128 (2022) 20, 202501

Scale evolution: Hadron scale

Reference scale: $\mu_{\text{Ref.}} = 331 \text{ MeV}$

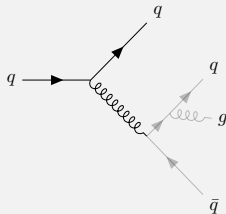
No gluons nor sea quarks

[J.Phys.Conf.Ser.1643(2020)1,012177]



Parton splitting (DGLAP)

[Nucl.Phys.B:126(1977)298]

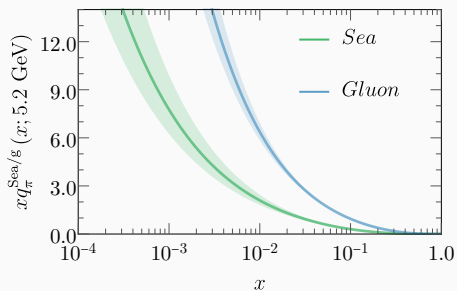
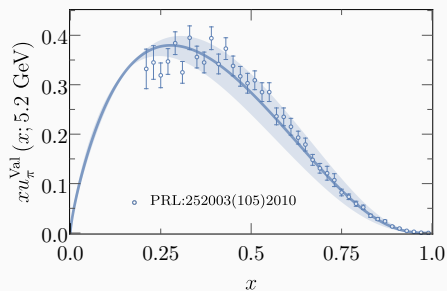


Renormalization group equations for GPD operators yield:

[Phys.Rev.D:55(1997)7114, Eur.Phys.J.C:82(2022)10,888]

$$\frac{dH^{(\pm)}(x, \xi, t; \mu^2)}{d \log \mu^2} = \frac{\alpha_s(\mu^2)}{4\pi} \int_x^\infty \frac{dy}{y} \mathcal{P}^{(\pm)}(y, \kappa, \alpha_s(\mu^2)) H^{(\pm)}\left(\frac{x}{y}, \xi, t; \mu^2\right)$$

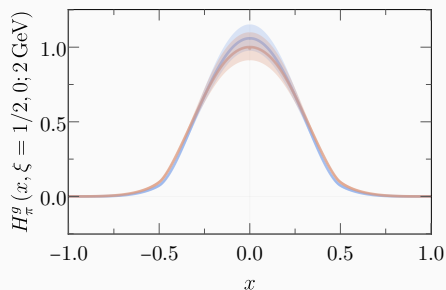
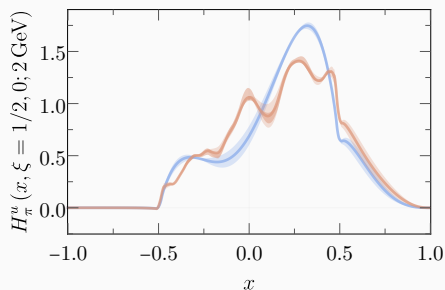
Scale evolution: PDFs and GPDs



- Non-zero gluon distribution generated by scale evolution.

Scale evolution: PDFs and GPDs

[Phys. Rev. D 105 (2022) 9, 094012]



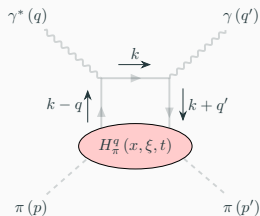
- Non-zero gluon distribution generated by scale evolution.
- Uncertainty band narrowed.
- Continuity along $x = \xi$ lines.

Compton form factors (I)

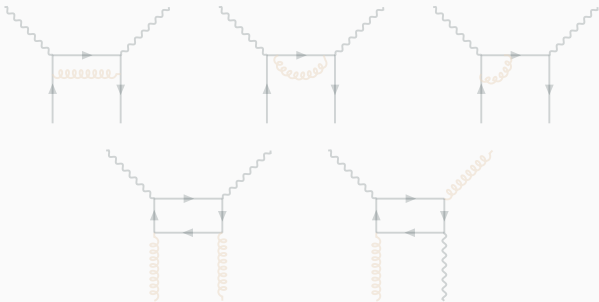
$$\mathcal{H}_\pi^p(\xi, t; Q^2) = \int_{-1}^1 \frac{dx}{\xi} \mathcal{K}^p\left(\frac{x}{\xi}, \frac{Q^2}{\mu_F^2}, \alpha_s(\mu_F^2)\right) H_\pi^p(x, \xi, t; \mu_F^2)$$



Leading order

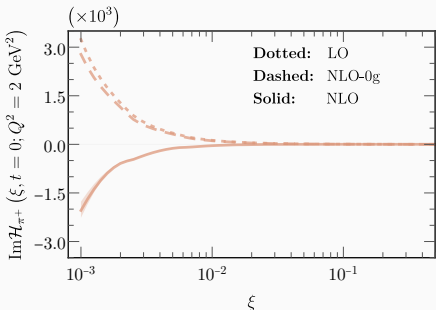
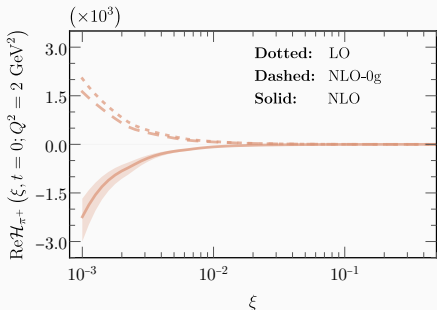


Next to leading order



Compton Form Factors (II)

[Phys. Rev. D 105 (2022) 9, 094012]

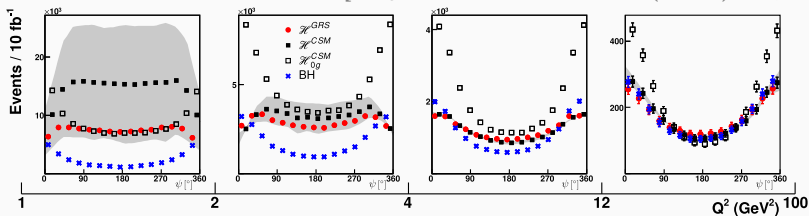


- Small effect of NLO corrections to quark amplitudes.
- Dominant effect of gluons.

Gluon dominance makes essential at least NLO accuracy in any phenomenological analysis of DVCS at an EIC.

Phenomenology of pion GPDs: EIC event-rates

[Phys. Rev. Lett. 128 (2022) 20, 202501]



The cross-section can be written as:

[Phys.Rev.D:79(2009)014017s]

$$|\mathcal{M}_{e\pi \rightarrow e\gamma\pi}|^2 = C_{\text{BH}} F_{\pi}^2(t_{\pi}) \pm \frac{F_{\pi}}{Q} (C_{\text{int}}(\psi) \text{Re}\mathcal{H}_{\pi} + \lambda S_{\text{Int}}(\psi) \text{Im}\mathcal{H}_{\pi}) + \frac{1}{Q^2} C_{\text{DVCS}} |\mathcal{H}_{\pi}|^2$$

$$|\mathcal{H}_{\pi}|^2 = \text{Re}^2(\mathcal{H}_{\pi}) + \text{Im}^2(\mathcal{H}_{\pi}) \sim \text{Re}^2(|\mathcal{H}_{\pi}^{q,\text{LO}}| - |\mathcal{H}_{\pi}^g|) + \text{Im}^2(|\mathcal{H}_{\pi}^{q,\text{LO}}| - |\mathcal{H}_{\pi}^g|)$$

Low energy, $Q^2 < 2 \text{ GeV}^2$

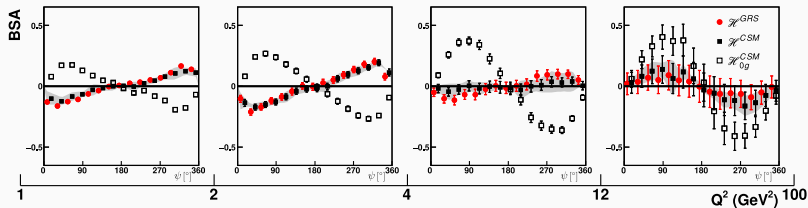
Intermediate energies

$$|\mathcal{H}_{\pi}|_{Q^2 < 2 \text{ GeV}^2}^2 \sim |\mathcal{H}_{\pi}^g|_{Q^2 < 2 \text{ GeV}^2}^2 > |\mathcal{H}_{\pi}^{q,\text{LO}} - \mathcal{H}_{\pi}^g|_{2 < Q^2 < 12 \text{ GeV}^2}^2 \sim |\mathcal{H}_{\pi}|_{2 < Q^2 < 12 \text{ GeV}^2}^2$$

Quark-gluon “interference”: Modulates expected number of events

Phenomenology of pion GPDs: EIC beam-spin asymmetries

[Phys. Rev. Lett. 128 (2022) 20, 202501]



$$\mathcal{A}(\psi) = \frac{\sigma^\uparrow - \sigma^\downarrow}{\sigma^\uparrow + \sigma^\downarrow} \sim s_{\text{Int}}^1 \text{Im} \mathcal{H}_\pi \sin \psi$$

$$\text{Im}(\mathcal{H}_\pi) = \text{Im}(\mathcal{H}_\pi^{q,\text{LO}}) + \text{Im}(\mathcal{H}_\pi^{q,\text{NLO}}) + \text{Im}(\mathcal{H}_\pi^g) \sim \text{Im}|\mathcal{H}_\pi^{q,\text{LO}}| - \text{Im}|\mathcal{H}_\pi^g|$$

Sign inversion: **Smoking gun** for gluon dominance

Summary and perspectives

1. *Can we obtain “theoretically complete” pion GPDs?* **YES!**

- Positivity: “Factorization hypothesis”
- Polynomiality: Covariant extension
- Support and soft-pion theorem
- Continuity along $x = \xi$

2. *What can we learn from them about the real world?*

- Hadron scale hypothesis
- GPDs might be accessible in future experiments
- Crucial role of gluons within pions
 - Modulate expected statistics
 - Manifest themselves through beam-spin asymmetries

Future work

- “Incomplete” covariant extension: Trigger Lattice efforts
- Refine evolution
- Neutral pions: Direct window to GPDs?