

Consistent large transverse momentum matching in TMDs with “Hadron Structure Oriented” approach

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GHP Workshop 2023

Based on

- **The resolution to the problem of consistent large transverse momentum in TMDs**

([arXiv:2303.04921](https://arxiv.org/abs/2303.04921))

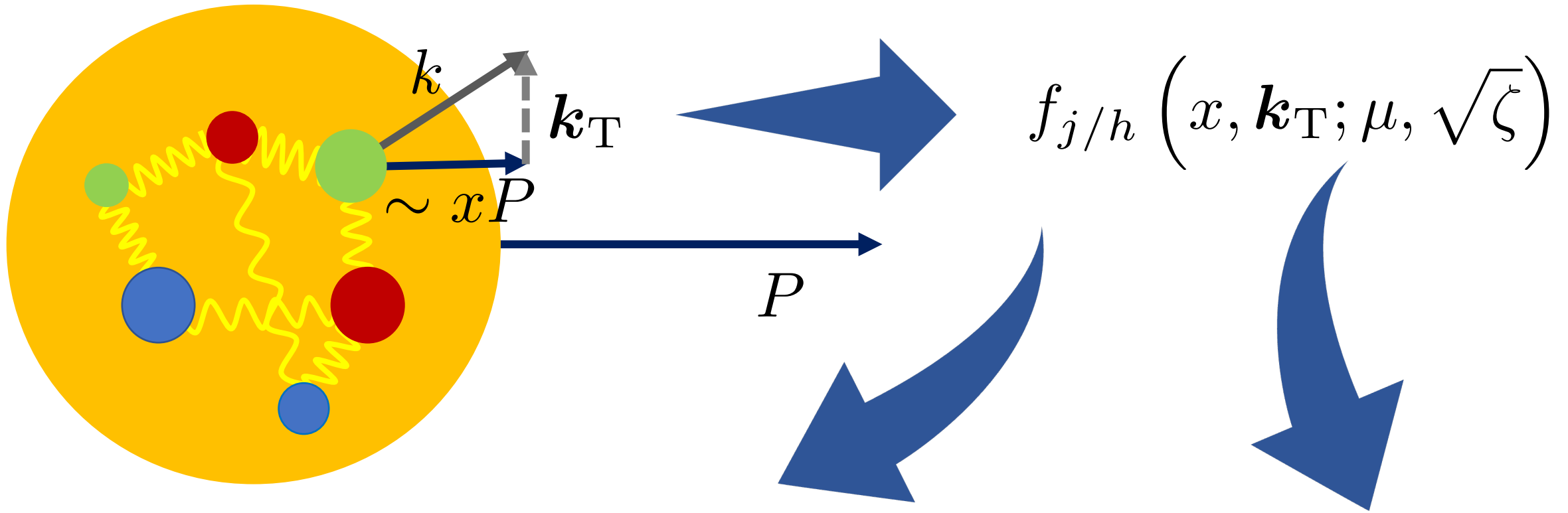
(J. O. Gonzalez-Hernandez, T. Rainaldi, T. C. Rogers)

- **Combining nonperturbative transverse momentum dependence with TMD evolution**

([PhysRevD.106.034002](https://arxiv.org/abs/2303.04921))

(J. O. Gonzalez-Hernandez, T. C. Rogers, N. Sato)

Why TMDs ?



Hadronic structure
(nonperturbative effects)

Predictions
(evolution & universality)

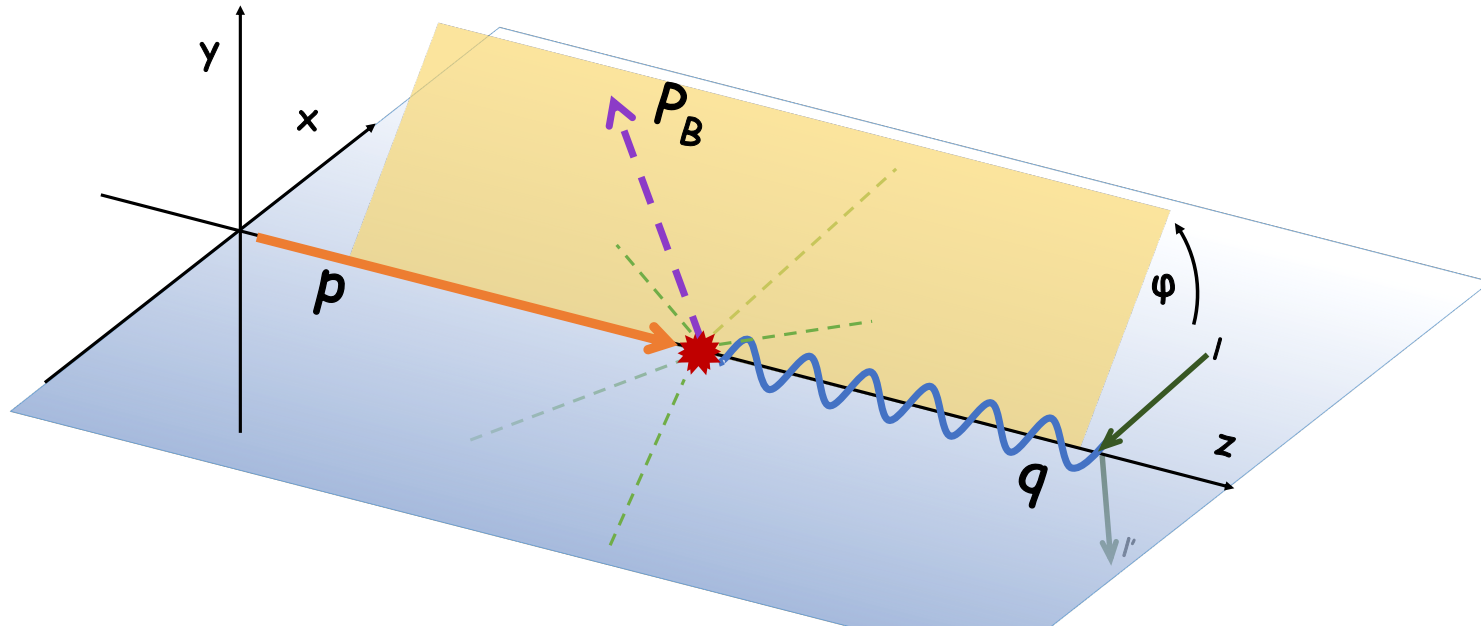
SIDIS

$$\frac{d\sigma}{dx dy dz dq_T^2} = \underbrace{W_{\text{SIDIS}}}_{q_T \ll Q} + \underbrace{Y_{\text{SIDIS}}}_{q_T \sim Q} + \mathcal{O}\left(\frac{m^2}{Q^2}\right)$$

FO collinear factorization

FO_{SIDIS} – ASY_{SIDIS}

$$H \int d^2 \mathbf{k}_{1T} d^2 \mathbf{k}_{2T} f_{j/p} \left(x, \mathbf{k}_{1T}; \mu, \sqrt{\zeta} \right) D_{h/j} \left(z, z \mathbf{k}_{2T}; \mu, \sqrt{\zeta} \right) \delta^{(2)} \left(\mathbf{q}_T + \mathbf{k}_{1T} - \mathbf{k}_{2T} \right)$$



Conventional approach :

$$H \int d^2 \mathbf{k}_{1T} d^2 \mathbf{k}_{2T} f_{j/p} \left(x, \mathbf{k}_{1T}; \mu, \sqrt{\zeta} \right) D_{h/j} \left(z, z \mathbf{k}_{2T}; \mu, \sqrt{\zeta} \right) \delta^{(2)} \left(\mathbf{q}_T + \mathbf{k}_{1T} - \mathbf{k}_{2T} \right)$$



Fourier Transform

$$H \int \frac{d^2 \mathbf{b}_T}{(2\pi)^2} e^{-i \mathbf{b}_T \cdot \mathbf{q}_T} \tilde{f}_{j/p} \left(x, \mathbf{b}_T; \mu, \sqrt{\zeta} \right) \tilde{D}_{h/j} \left(z, \mathbf{b}_T; \mu, \sqrt{\zeta} \right)$$

Solve evolution equations relating input scale with SIDIS scale

$$\frac{\partial \ln \tilde{f}_{j/p}(x, \mathbf{b}_T; \mu, \sqrt{\zeta})}{\partial \ln \sqrt{\zeta}} = \tilde{K}(\mathbf{b}_T; \mu)$$

$$\frac{d \ln \tilde{f}_{j/p}(x, \mathbf{b}_T; \mu, \sqrt{\zeta})}{d \ln \mu} = \gamma(\alpha_S(\mu); \mu/\sqrt{\zeta})$$

$$\frac{d\tilde{K}(\mathbf{b}_T; \mu)}{d \ln \mu} = -\gamma_K(\alpha_S(\mu))$$

Same for FF

$$\begin{aligned} \mu &= \sqrt{\zeta} \\ \mu_0 &= \sqrt{\zeta_0} \end{aligned}$$

$$\begin{aligned} \tilde{f}_{j/p}(x, \mathbf{b}_T; \mu, \sqrt{\zeta}) &= \tilde{f}_{j/p}(x, \mathbf{b}_T; \mu_0, \sqrt{\zeta_0}) \times \\ &\times \exp \left\{ \int_{\mu_0}^{\mu} \frac{d\mu'}{\mu'} \left[\gamma(\alpha_S(\mu'); 1) - \ln \left(\frac{\sqrt{\zeta}}{\mu'} \right) \gamma_K(\alpha_S(\mu')) \right] + \ln \left(\frac{\sqrt{\zeta}}{\sqrt{\zeta_0}} \tilde{K}(\mathbf{b}_T; \mu_0) \right) \right\} \end{aligned}$$

Separate $b_T < b_{max}$ & $b_T > b_{max}$ regions with a b_* prescription

$$\tilde{f}_{j/p}(x; \mathbf{b}_T; \mu_Q, Q) = \tilde{f}_{j/p}(x; \mathbf{b}_*; \mu_Q, Q) \underbrace{\frac{\tilde{f}_{j/p}(x; \mathbf{b}_T; \mu_Q, Q)}{\tilde{f}_{j/p}(x; \mathbf{b}_*; \mu_Q, Q)}}_{\exp\{-g_{j/p}(x, \mathbf{b}_T)\}}$$

Same for FF

Perturbatively calculable with fixed order collinear factorization

Nonperturbative

$$g_K(\mathbf{b}_T) \equiv \tilde{K}(\mathbf{b}_*; \mu) - \tilde{K}(\mathbf{b}_T; \mu)$$

Choose ansatzes for g functions

$$g_{j/p}(x, \mathbf{b}_T) = \frac{1}{4} M_F^2 b_T^2$$

$$g_{h/j}(z, \mathbf{b}_T) = \frac{1}{4z^2} M_D^2 b_T^2$$

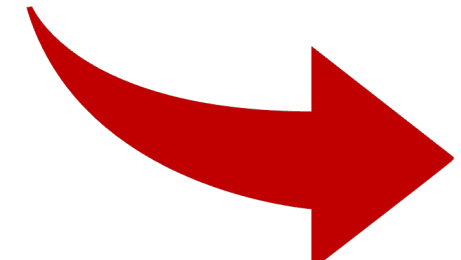
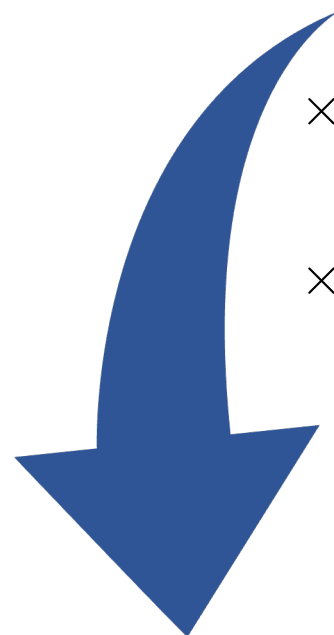
$$g_K(\mathbf{b}_T) = \frac{g_2}{2M_K^2} \ln(1 + M_K^2 b_T^2)$$

$$g_K(\mathbf{b}_T) = \frac{1}{2} M_K^2 b_T^2$$

Relate μ_{b_*} with input scale Q_0 and get OPE expansion

$$\mu_{b_*} \equiv \frac{2e^{-\gamma_E}}{b_T}$$

$$\tilde{f}_{j/p}(x; \mathbf{b}_T; \mu_Q, Q) = \tilde{f}_{j/p}^{\text{OPE}}(x; \mathbf{b}_*; \mu_{b_*}, \mu_{b_*}) \times \exp \left\{ \int_{\mu_{b_*}}^{\mu_Q} \frac{d\mu'}{\mu'} \left[\gamma(\alpha_S(\mu'); 1) - \ln \left(\frac{Q}{\mu'} \right) \gamma_K(\alpha_S(\mu')) \right] + \ln \left(\frac{Q}{\mu_{b_*}} \right) \tilde{K}(\mathbf{b}_*; \mu_{b_*}) \right\} \times \exp \left\{ -g_{j/p}(x, \mathbf{b}_T) - g_K(\mathbf{b}_T) \ln \left(\frac{Q}{Q_0} \right) \right\}$$



Nonperturbative

Perturbatively calculable



Drop this

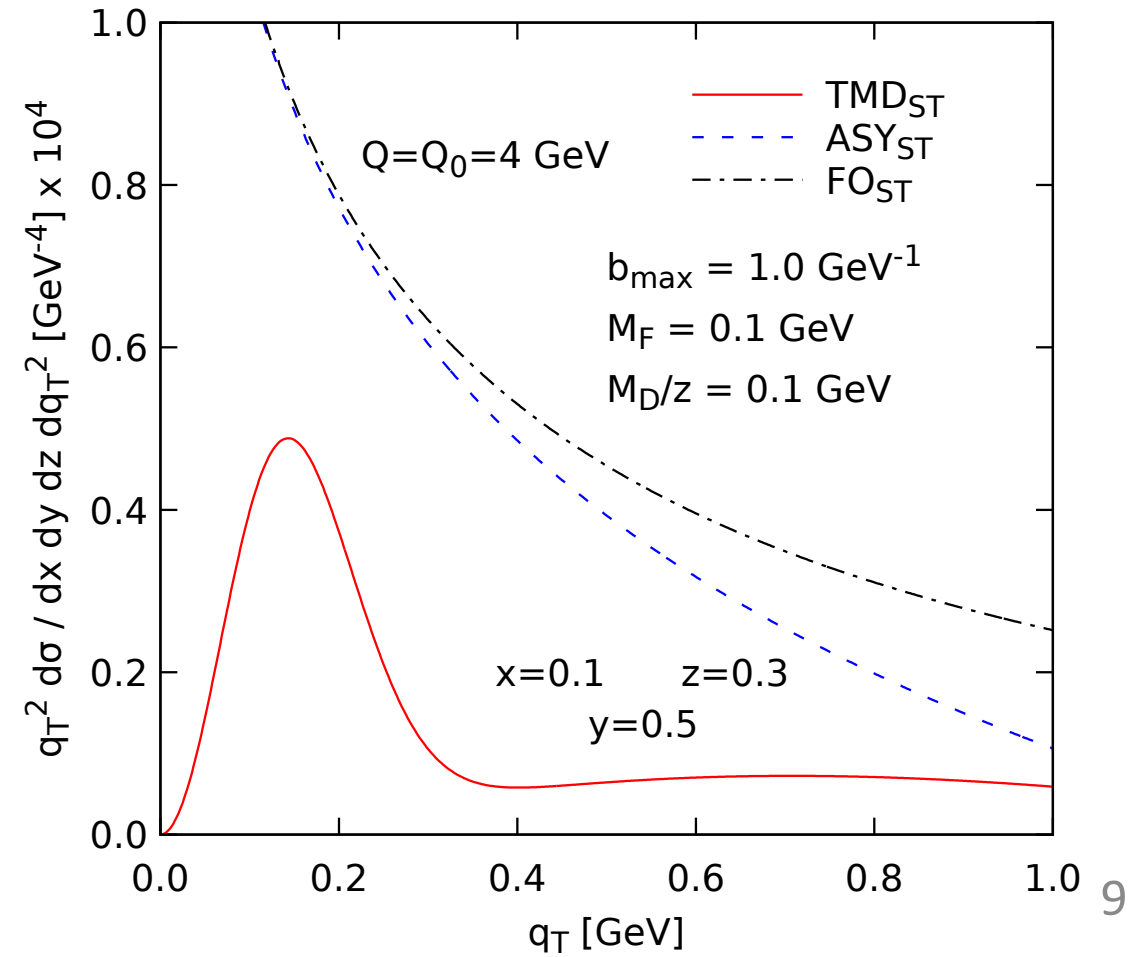
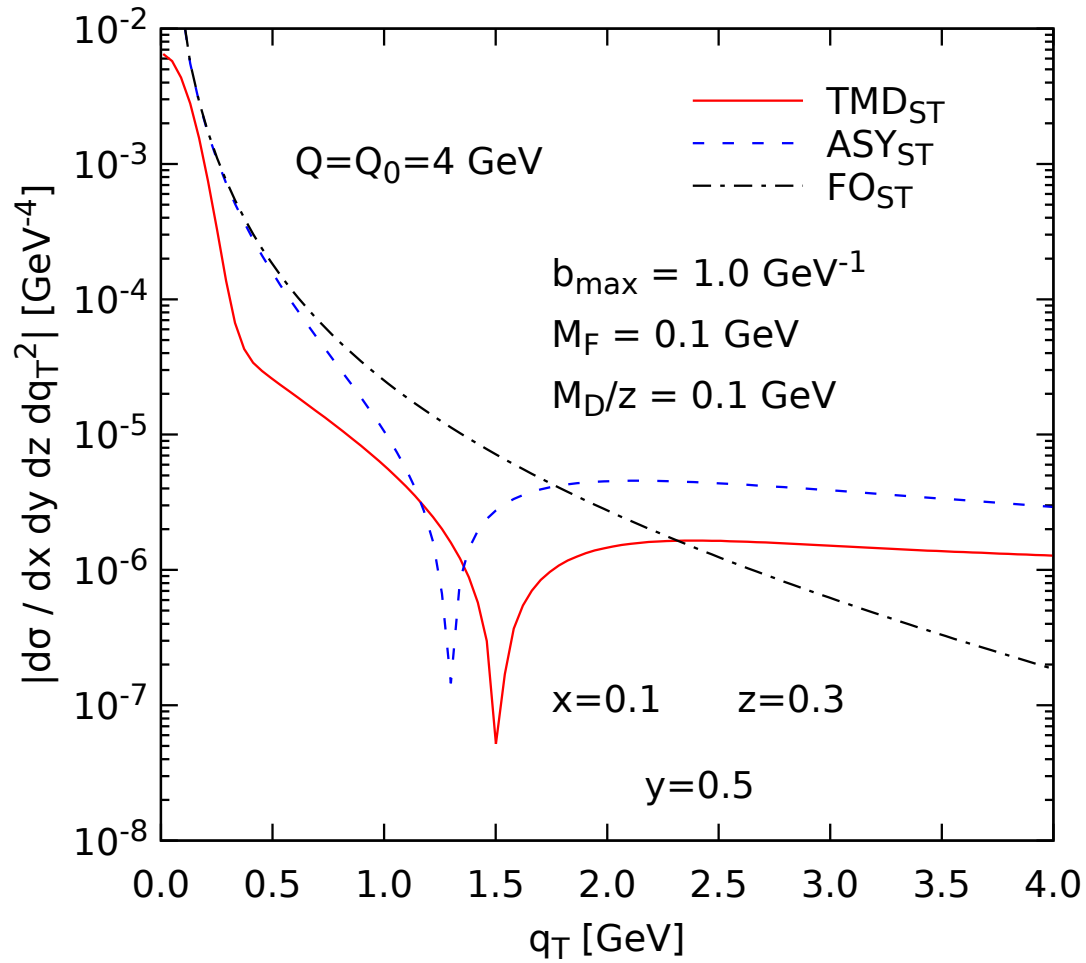
$$\tilde{f}_{j/p}^{\text{OPE}}(x, \mathbf{b}_*; \mu_{b_*}, \mu_{b_*}) = \tilde{C}_{j/j'}(x/\xi, \mathbf{b}_*; \mu_{b_*}, \mu_{b_*}) \otimes \tilde{f}_{j'/p}(\xi; \mu_{b_*}) + \mathcal{O}(m_{\text{max}}^2)$$

Fixed order collinear factorization

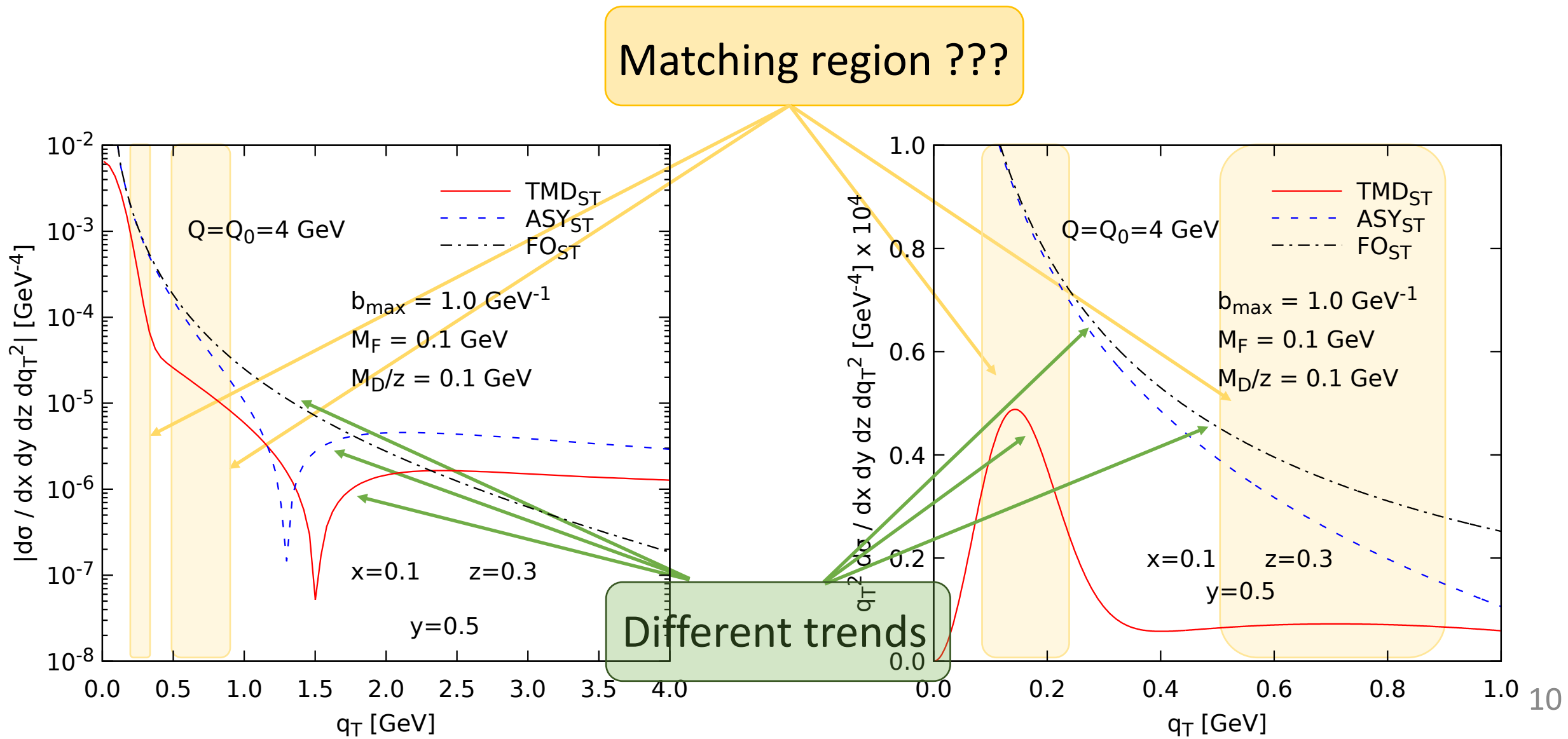
Same for FF

Conventional approach results for SIDIS

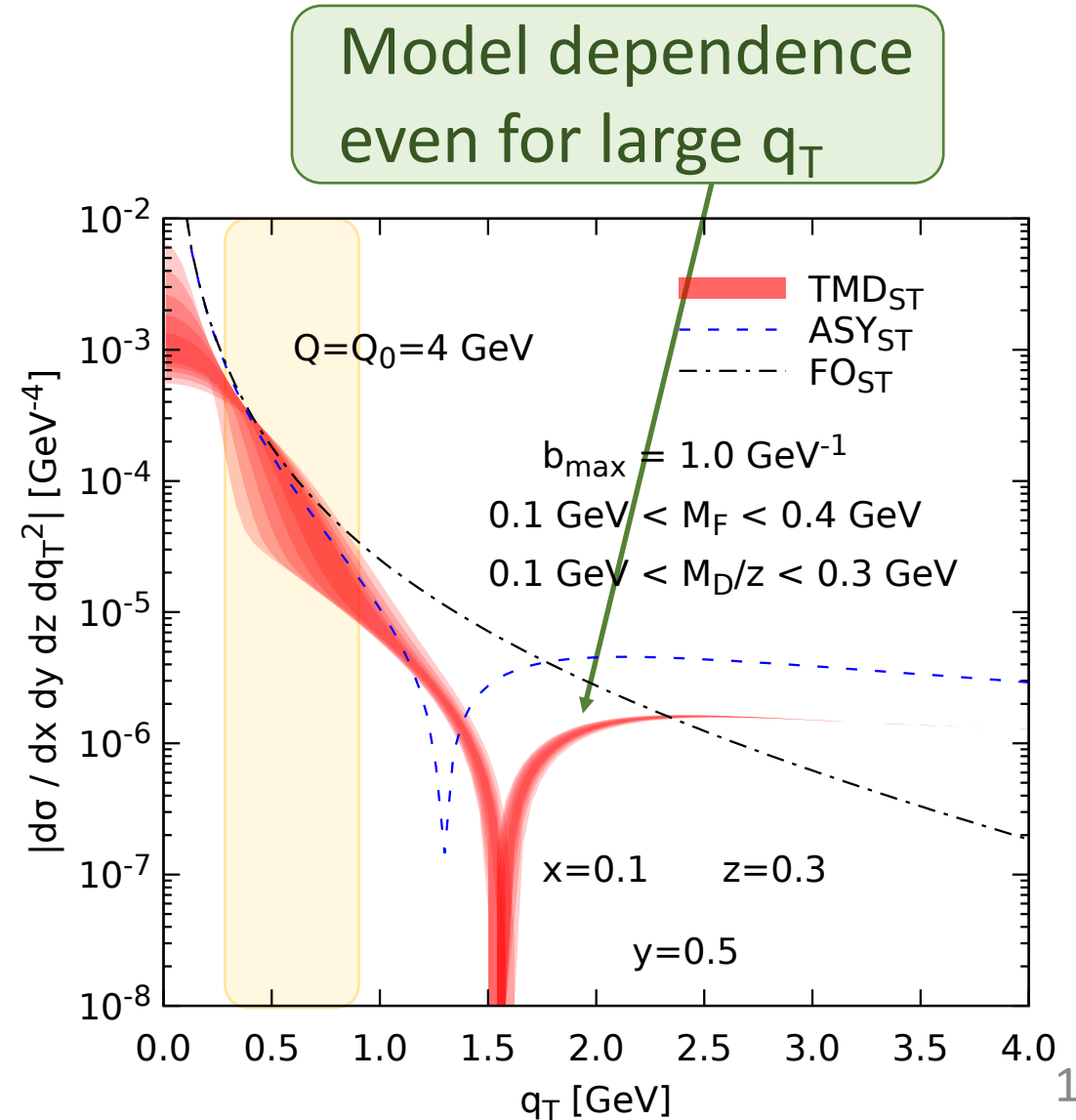
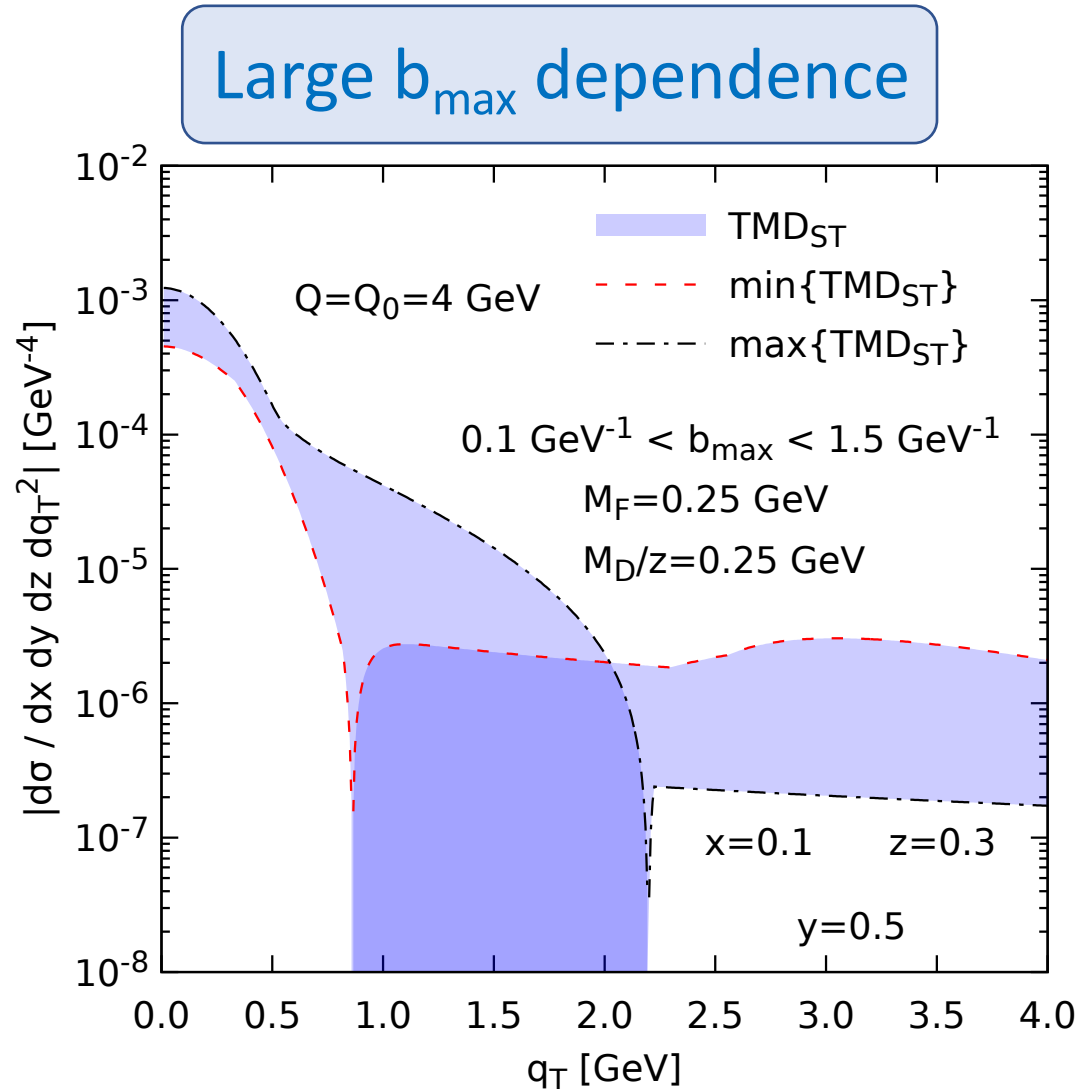
$$\Lambda_{\text{QCD}} \ll q_T \ll Q_0$$



Conventional approach results for SIDIS



(Other) Issues with conventional approach



Hadron Structure Oriented approach (HSO)

- Fixes TMDs parametrization at **input scale** Q_0
- Uses uniquely determined TMDs for **all transverse momenta**
- Interpolates **perturbative (large k_T)** and **nonperturbative (small k_T)** TM regions
- Can **swap NP models** easily
- Explicit (approximate) **probability interpretation**

No need for b_* prescription !

TMD PDF & FF HSO parametrization at input scale

Fixed order collinear factorization

$\mathcal{O}(\alpha_S)$

$$f_{j/p}(x, \mathbf{k}_T; \mu_{Q_0}, Q_0^2) = \frac{1}{2\pi} \frac{1}{k_T^2 + m_{f_{j/p}}^2} \left[A_{j/p}^f(x; \mu_Q) + B_{j/p}^f(x; \mu_Q) \ln \left(\frac{Q_0^2}{k_T^2 + m_{f_{j/p}}^2} \right) \right] + \frac{1}{2\pi} \frac{A_{j/p}^{f,g}(x; \mu_Q)}{k_T^2 + m_{f_{g/p}}^2} + C_{j/p}^f f_{\text{core},j,p}(x, \mathbf{k}_T; Q_0^2)$$

Such that

Small k_T model

NP parameters

$$f_{j/p}^c(x; \mu_Q) \equiv 2\pi \int_0^{k_c} dk_T k_T f_{j/p}(x, \mathbf{k}_T; \mu_Q, \sqrt{\zeta}) = f_{j/p}(x; \mu_Q) + \Delta_{j/p}(x; \mu_Q, k_c) + \text{p.s.}$$

TMD PDF & FF HSO parametrization at input scale

Fixed order collinear factorization

$\mathcal{O}(\alpha_S)$

$$D_{h/j}(z, z\mathbf{k}_T; \mu_{Q_0}, Q_0^2) = \frac{1}{2\pi z^2} \frac{1}{k_T^2 + m_{D_{h/j}}^2} \left[A_{h/j}^D(z; \mu_Q) + B_{h/j}^D(z; \mu_Q) \ln \left(\frac{Q_0^2}{k_T^2 + m_{D_{h/j}}^2} \right) \right] + \frac{1}{2\pi z^2} \frac{A_{h/g}^{D,g}(z; \mu_Q)}{k_T^2 + m_{D_{h/g}}^2} + C_{h/j}^D D_{\text{core},h/j}(z, z\mathbf{k}_T; Q_0^2)$$

Such that

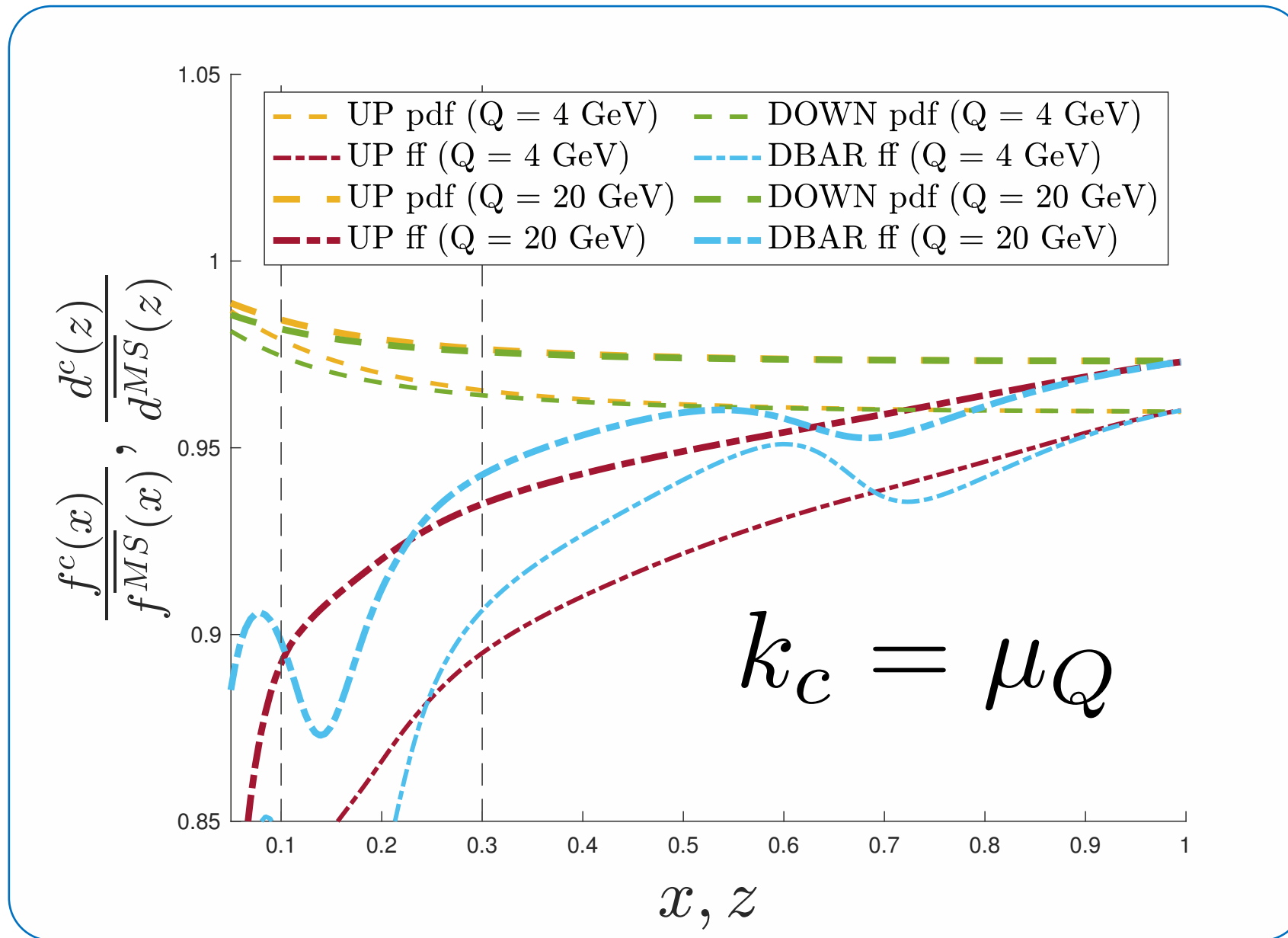
Small k_T model

NP parameters

$$d_{h/j}^c(z; \mu_Q) \equiv 2\pi z^2 \int_0^{k_c} dk_T k_T D_{h/j}(z, z\mathbf{k}_T; \mu_Q, \sqrt{\zeta})$$

$$= d_{h/j}(z; \mu_Q) + \Delta_{h/j}(z; \mu_Q, k_c) + \text{p.s.}$$

From $\overline{\text{MS}}$ to Cutoff schemes



Choose “core” models (examples)

$$f_{\text{core},i/p}^{\text{Gauss}}(x, \mathbf{k}_T; Q_0^2) = \frac{e^{-k_T^2/M_F^2}}{\pi M_F^2}$$

$$D_{\text{core},h/j}^{\text{Gauss}}(z, z\mathbf{k}_T; Q_0^2) = \frac{e^{-z^2 k_T^2/M_D^2}}{\pi M_D^2}$$

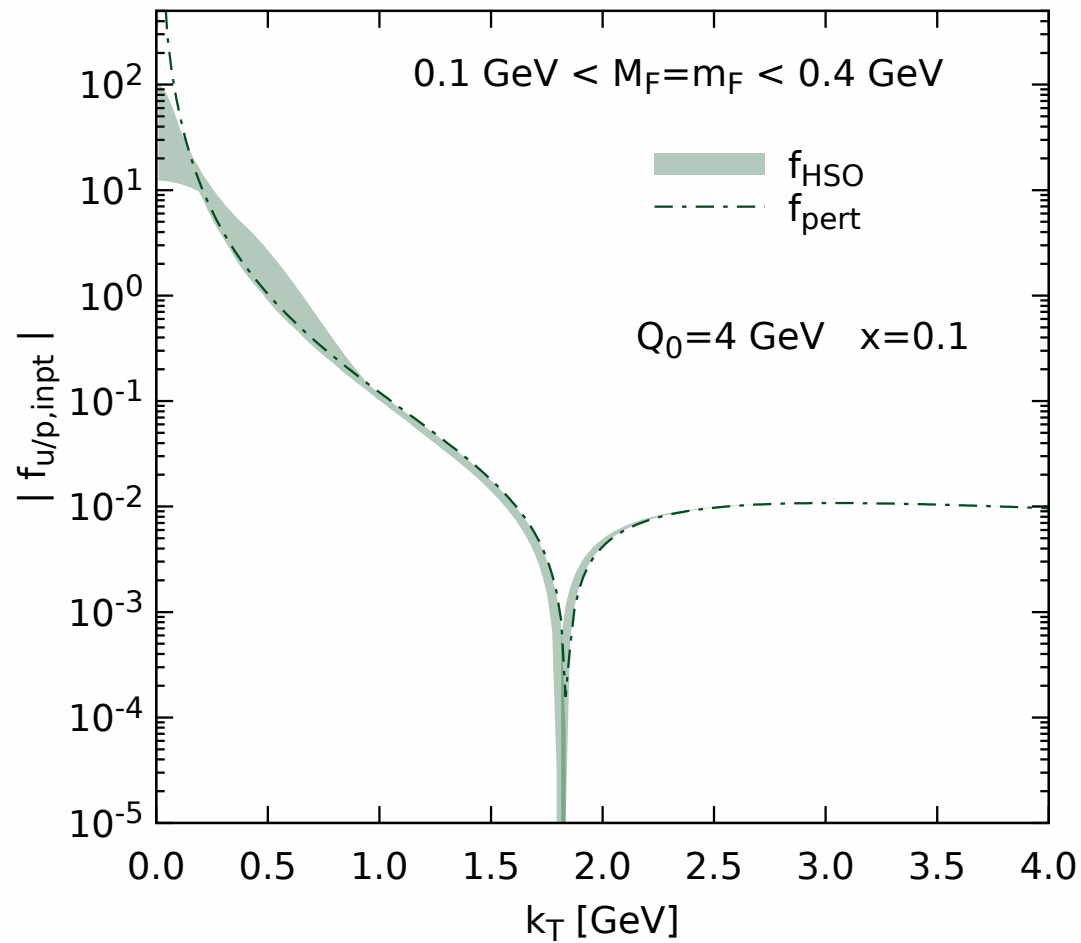
Gaussian “core” models

Spectator-like “core” models

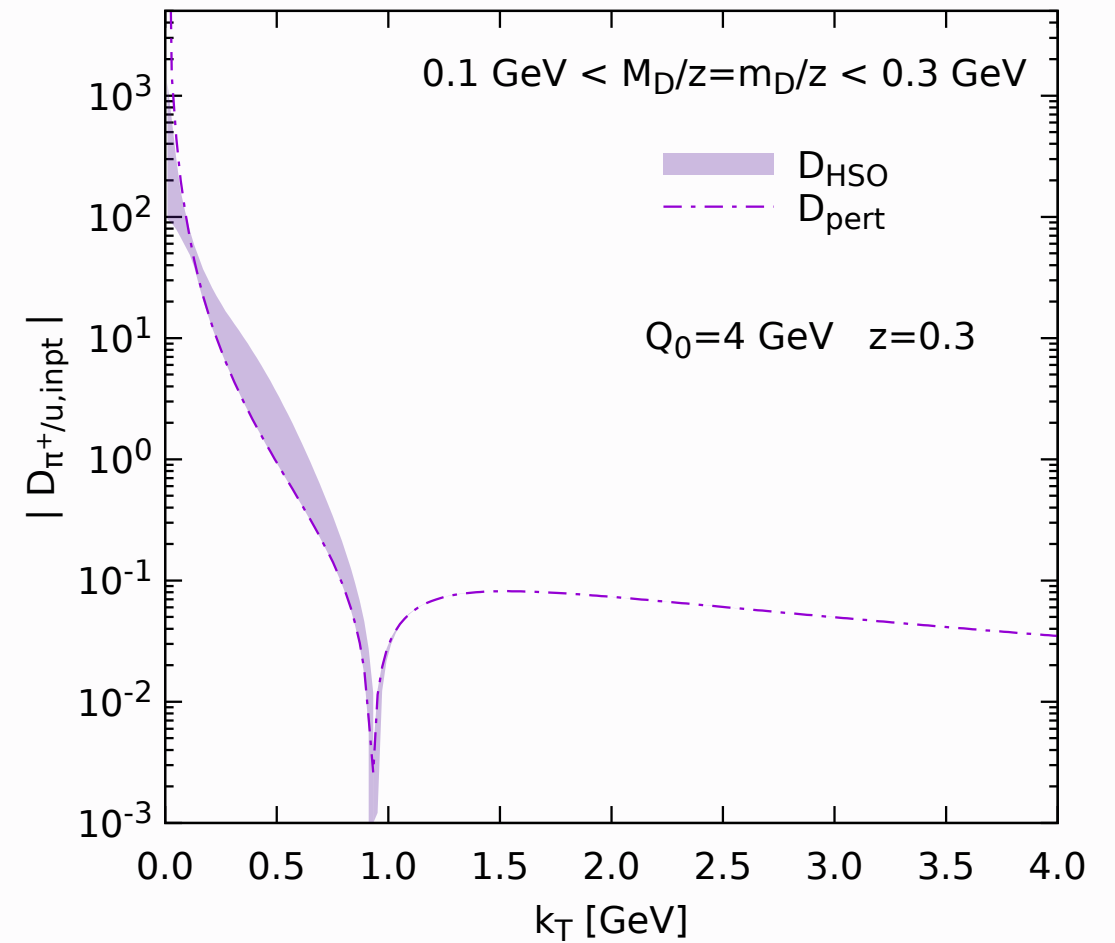
$$f_{\text{core},j/p}^{\text{Spect}}(x, \mathbf{k}_T; Q_0^2) = \frac{6M_{0F}^6}{\pi(2M_F^2 + M_{0F}^2)} \frac{M_F^2 + k_T^2}{(M_{0F}^2 + k_T^2)^4}$$

$$D_{\text{core},h/j}^{\text{Spect}}(z, z\mathbf{k}_T; Q_0^2) = \frac{2M_{0D}^4}{\pi(M_D^2 + M_{0D}^2)} \frac{M_D^2 + z^2 k_T^2}{(M_{0D}^2 + z^2 k_T^2)^3}$$

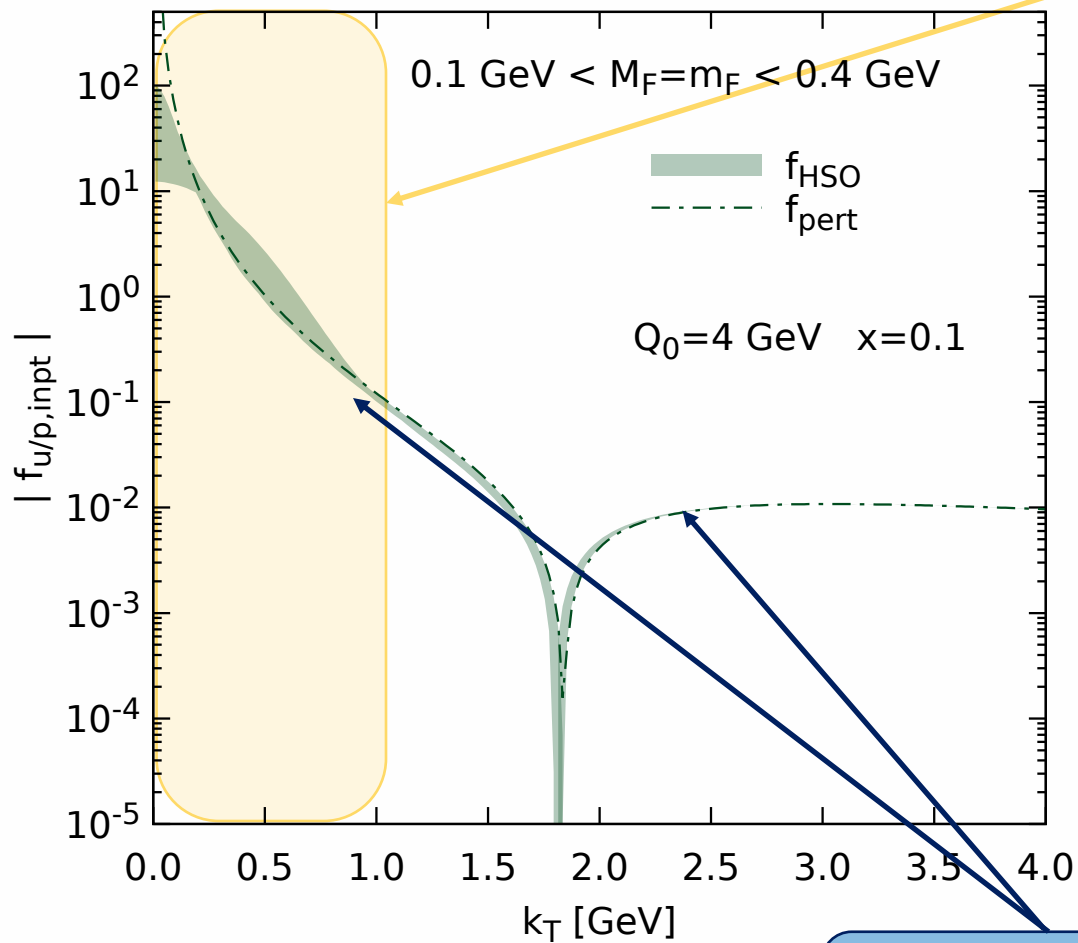
Up-quark from Proton TMD pdf



π^+ from Up-quark TMD ff

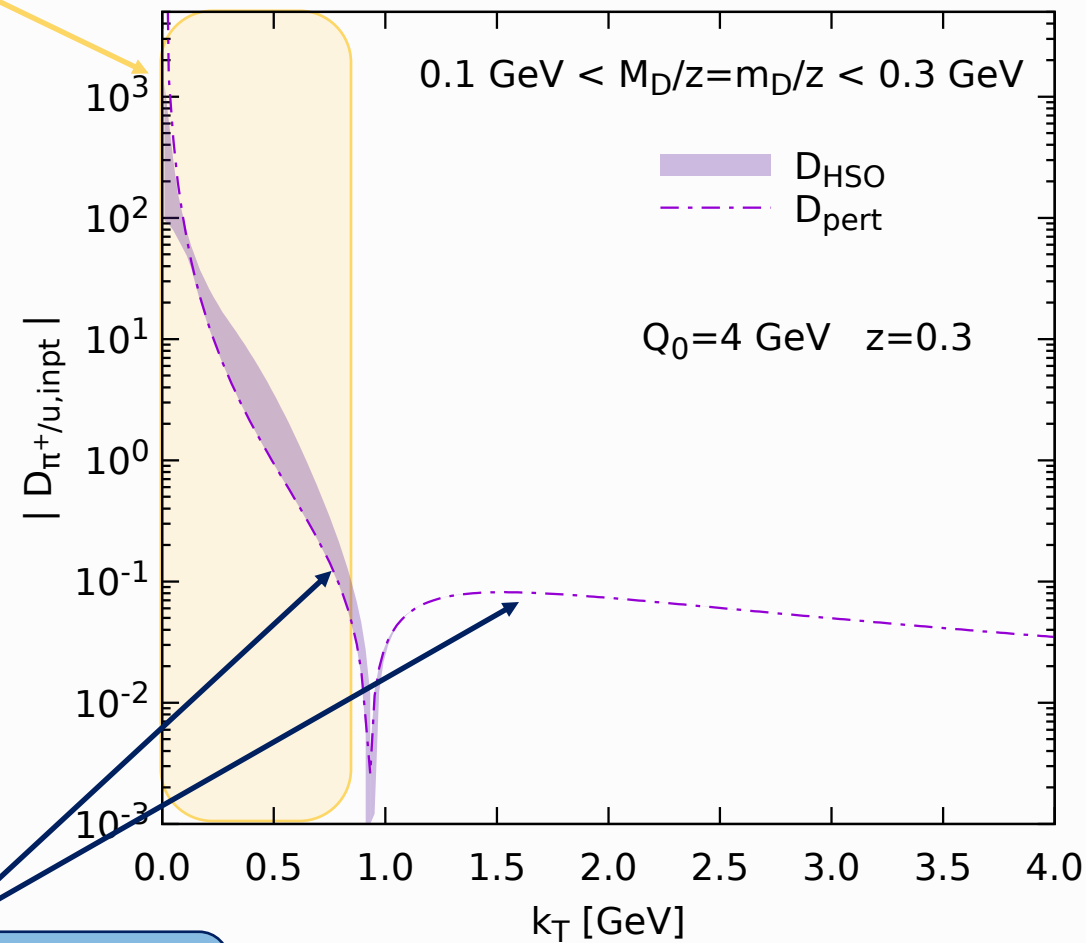


Up-quark from Proton TMD pdf



Model dependence

π^+ from Up-quark TMD ff



Smooth transition

Asymptotic term

HSO

Conventional

$$ASY_{\text{HSO}} = \lim_{\frac{q_{\text{T}}}{Q} \rightarrow \sim 1, \frac{m^2}{Q^2} \rightarrow 0} W_{\text{HSO}}$$

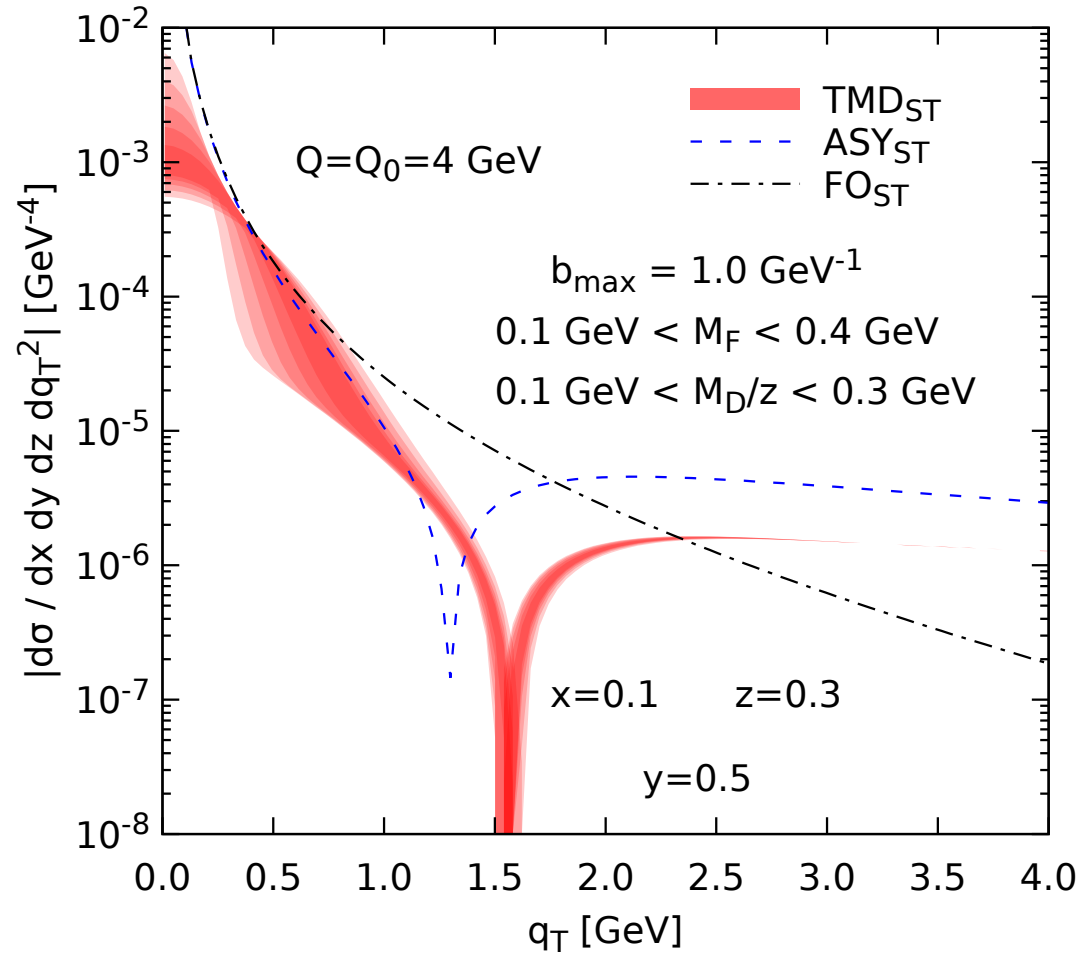
$$ASY_{\text{ST}} = \lim_{\frac{q_{\text{T}}}{Q} \rightarrow 0, \frac{m^2}{Q^2} \rightarrow 0} FO_{\text{ST}}$$



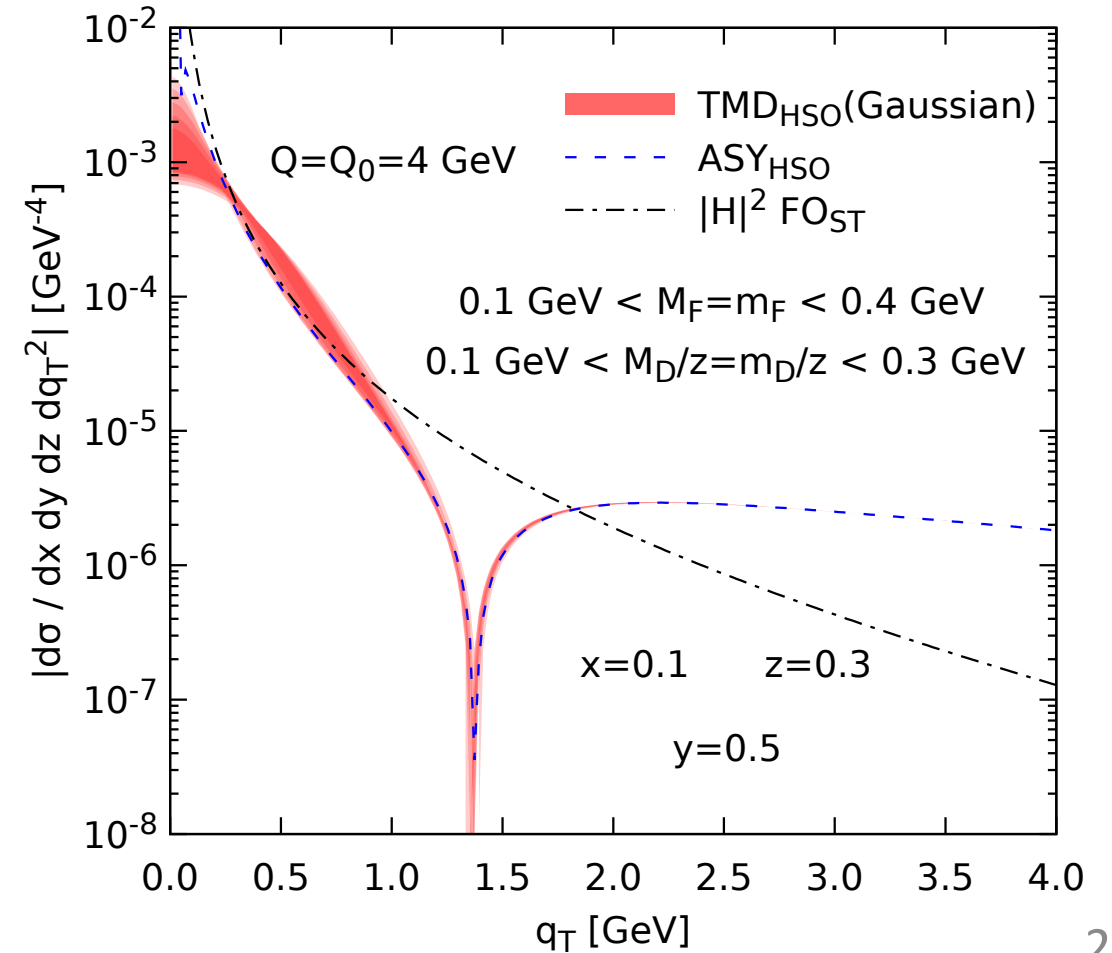
$$\begin{aligned}
 [f, D]_{\text{ASY}} = & D^{\text{pert}}(z, z\mathbf{q}_{\text{T}}; \mu_Q, Q^2) f^c(x; \mu_Q) + \frac{1}{z^2} f^{\text{pert}}(x, -\mathbf{q}_{\text{T}}; \mu_Q, Q^2) d^c(z; \mu_Q) \\
 & + \int d^2\mathbf{k}_{\text{T}} \left\{ f^{\text{pert}}(x, \mathbf{k}_{\text{T}} - \mathbf{q}_{\text{T}}/2; \mu_Q, Q^2) D^{\text{pert}}(z, z(\mathbf{k}_{\text{T}} + \mathbf{q}_{\text{T}}/2); \mu_Q, Q^2) \right. \\
 & - D^{\text{pert}}(z, z\mathbf{q}_{\text{T}}; \mu_Q, Q^2) f^{\text{pert}}(x, \mathbf{k}_{\text{T}} - \mathbf{q}_{\text{T}}/2; \mu_Q, Q^2) \Theta(\mu_Q - |\mathbf{k}_{\text{T}} - \mathbf{q}_{\text{T}}/2|) \\
 & \left. - D^{\text{pert}}(z, z(\mathbf{k}_{\text{T}} + \mathbf{q}_{\text{T}}/2); \mu_Q, Q^2) f^{\text{pert}}(x, -\mathbf{q}_{\text{T}}; \mu_Q, Q^2) \Theta(\mu_Q - |\mathbf{k}_{\text{T}} + \mathbf{q}_{\text{T}}/2|) \right\}
 \end{aligned}$$

Conventional vs HSO - SIDIS cross section

Conventional

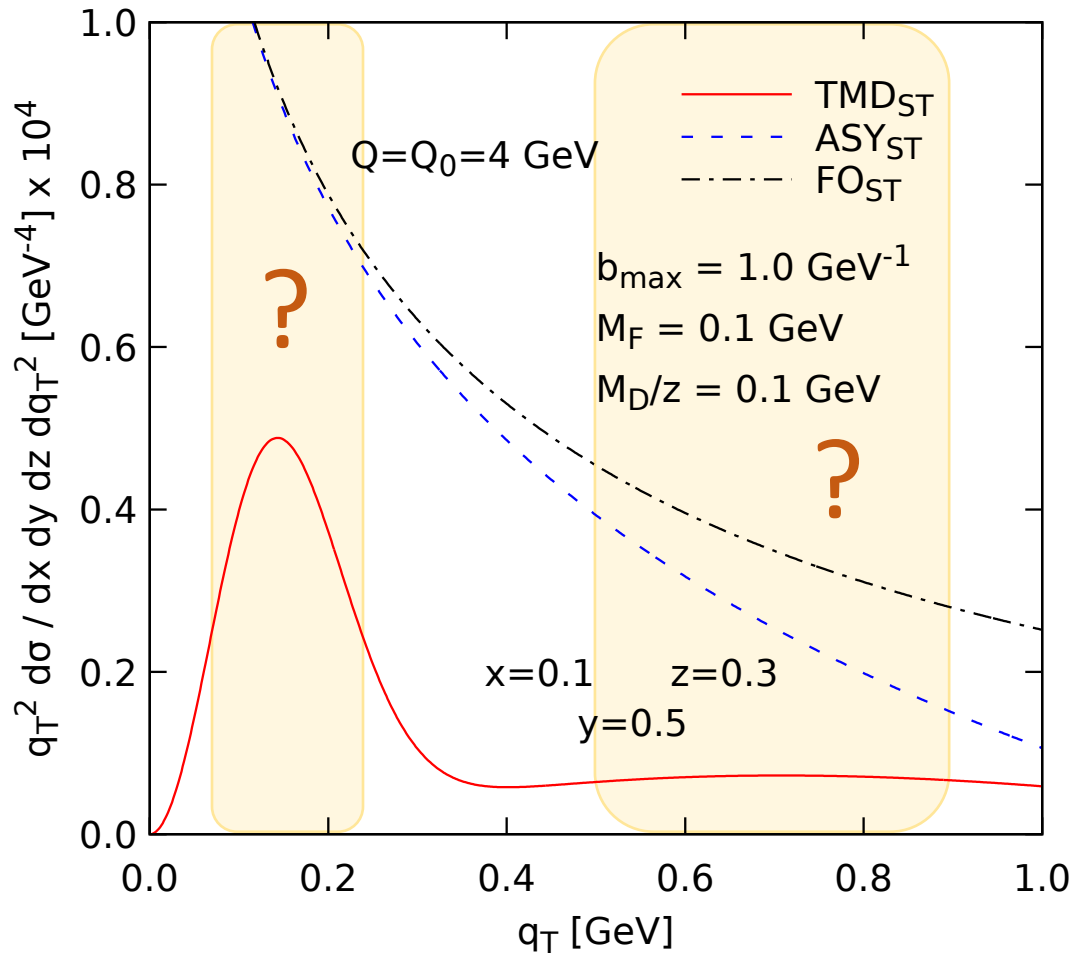


HSO (Gaussian)

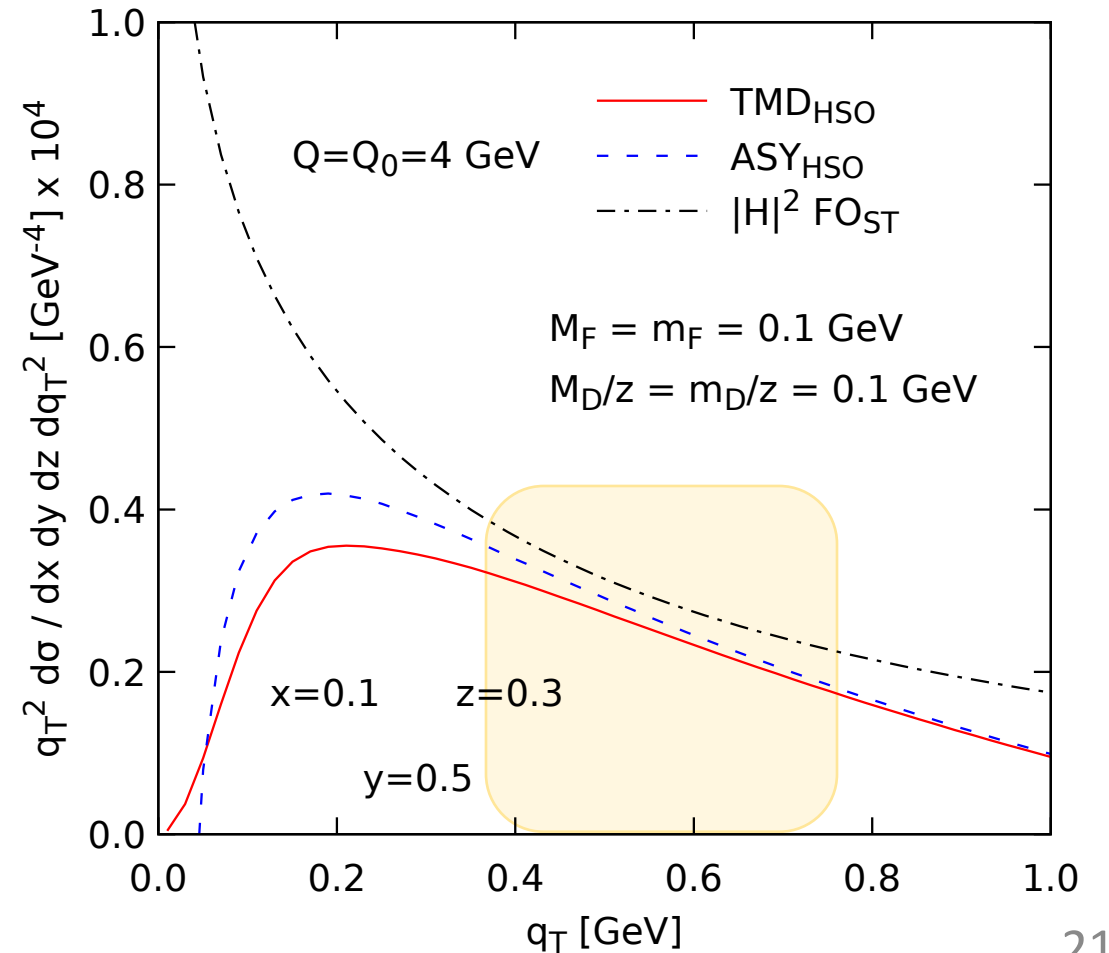


Conventional vs HSO - SIDIS cross section

Conventional



HSO (Gaussian)



Summary

- Consistent TMD parametrization for large TM at input scale
- No need of b_{\max}
- Improved TM behavior in matching region

NEXT:

- Check with data (SIDIS, DY, DIA, ...)
- Add higher orders
- Incorporate NP calculations (lattice, EFT, ...)

Thank you

Backup slides

Coefficients (TMD pdf)

$$A_{i/p}^f(x; \mu_{Q_0}) \equiv \sum_{ii'} \delta_{i'i} \frac{\alpha_s(\mu_{Q_0})}{\pi} \left\{ [(P_{i'i} \otimes f_{i'/p})(x; \mu_{Q_0})] - \frac{3C_F}{2} f_{i'/p}(x; \mu_{Q_0}) \right\},$$

$$B_{i/p}^f(x; \mu_{Q_0}) \equiv \sum_{i'i} \delta_{i'i} \frac{\alpha_s(\mu_{Q_0}) C_F}{\pi} f_{i'/p}(x; \mu_{Q_0}),$$

$$A_{i/p}^{f,g}(x; \mu_{Q_0}) \equiv \frac{\alpha_s(\mu_{Q_0})}{\pi} [(P_{ig} \otimes f_{g/p})(x; \mu_{Q_0})],$$

$$C_{i/p}^f \equiv \frac{1}{N_{i/p}^f} \left[f_{i/p}(x; \mu_{Q_0}) - A_{i/p}^f(x; \mu_{Q_0}) \ln \left(\frac{\mu_{Q_0}}{m_{f_{i,p}}} \right) - B_{i/p}^f(x; \mu_{Q_0}) \ln \left(\frac{\mu_{Q_0}}{m_{f_{i,p}}} \right) \ln \left(\frac{Q_0^2}{\mu_{Q_0} m_{f_{i,p}}} \right), \right. \\ \left. - A_{i/p}^{f,g}(x; \mu_{Q_0}) \ln \left(\frac{\mu_{Q_0}}{m_{f_{g,p}}} \right) + \frac{\alpha_s(\mu_{Q_0})}{2\pi} \left\{ \sum_{ii'} \delta_{i'i} [C_{\Delta}^{i/i'} \otimes f_{i'/p}](x; \mu_{Q_0}) + [C_{\Delta}^{i/g} \otimes f_{g/p}](x; \mu_{Q_0}) \right\} \right].$$

Coefficients (TMD ff)

$$A_{h/j}^D(z; \mu_{Q_0}) \equiv \sum_{jj'} \delta_{j'j} \frac{\alpha_s(\mu_{Q_0})}{\pi} \left\{ [(P_{jj'} \otimes d_{h/j'})(z; \mu_{Q_0})] - \frac{3C_F}{2} d_{h/j'}(z; \mu_{Q_0}) \right\},$$

$$B_{h/j}^D(z; \mu_{Q_0}) \equiv \sum_{jj'} \delta_{j'j} \frac{\alpha_s(\mu_{Q_0}) C_F}{\pi} d_{h/j'}(z; \mu_{Q_0}),$$

$$A_{h/j}^{D,g}(z; \mu_{Q_0}) \equiv \frac{\alpha_s(\mu_{Q_0})}{\pi} [(P_{gj} \otimes d_{h/g})(z; \mu_{Q_0})],$$

$$C_{h/j}^D \equiv \frac{1}{N_{h/j}^D} \left[d_{h/j}(z; \mu_{Q_0}) - A_{h/j}^D(z; \mu_{Q_0}) \ln \left(\frac{\mu_{Q_0}}{m_{D_{h,j}}} \right) - B_{h/j}^D(z; \mu_{Q_0}) \ln \left(\frac{\mu_{Q_0}}{m_{D_{h,j}}} \right) \ln \left(\frac{Q_0^2}{\mu_{Q_0} m_{D_{h,j}}} \right), \right. \\ \left. - A_{h/j}^{D,g}(z; \mu_{Q_0}) \ln \left(\frac{\mu_{Q_0}}{m_{D_{h,g}}} \right) + \frac{\alpha_s(\mu_{Q_0})}{2\pi} \left\{ \sum_{jj'} \delta_{j'j} [C_{\Delta}^{j'/j} \otimes d_{h/j'}](z; \mu_{Q_0}) + [C_{\Delta}^{g/j} \otimes d_{h/g}](z; \mu_{Q_0}) \right\} \right].$$

b_T space TMD distributions at input scale

$$\begin{aligned}
 z^2 \tilde{D}_{\text{inpt},h/j}(z, \mathbf{b}_T; \mu_{Q_0}, Q_0^2) &= K_0 (b_T m_{D_{h,j}}) \left[A_{h/j}^D(z; \mu_{Q_0}) + B_{h/j}^D(z; \mu_{Q_0}) \ln \left(\frac{b_T Q_0^2 e^{\gamma_E}}{2m_{D_{h,j}}} \right) \right] \\
 &\quad + K_0 (b_T m_{D_{h,g}}) A_{h/j}^{D,g}(z; \mu_{Q_0}) \\
 &\quad + C_{h/j}^D z^2 \tilde{D}_{\text{core},h/j}(z, \mathbf{b}_T; Q_0^2),
 \end{aligned}$$

$$\begin{aligned}
 \tilde{f}_{\text{inpt},i/p}(x, \mathbf{b}_T; \mu_{Q_0}, Q_0^2) &= K_0 (b_T m_{f_{i,p}}) \left[A_{i/p}^f(x; \mu_{Q_0}) + B_{i/p}^f(x; \mu_{Q_0}) \ln \left(\frac{b_T Q_0^2 e^{\gamma_E}}{2m_{f_{i,p}}} \right) \right] \\
 &\quad + K_0 (b_T m_{f_{g,p}}) A_{g/p}^f(x; \mu_{Q_0}) \\
 &\quad + C_{i/p}^f \tilde{f}_{\text{core},i/p}(x, \mathbf{b}_T; Q_0^2),
 \end{aligned}$$

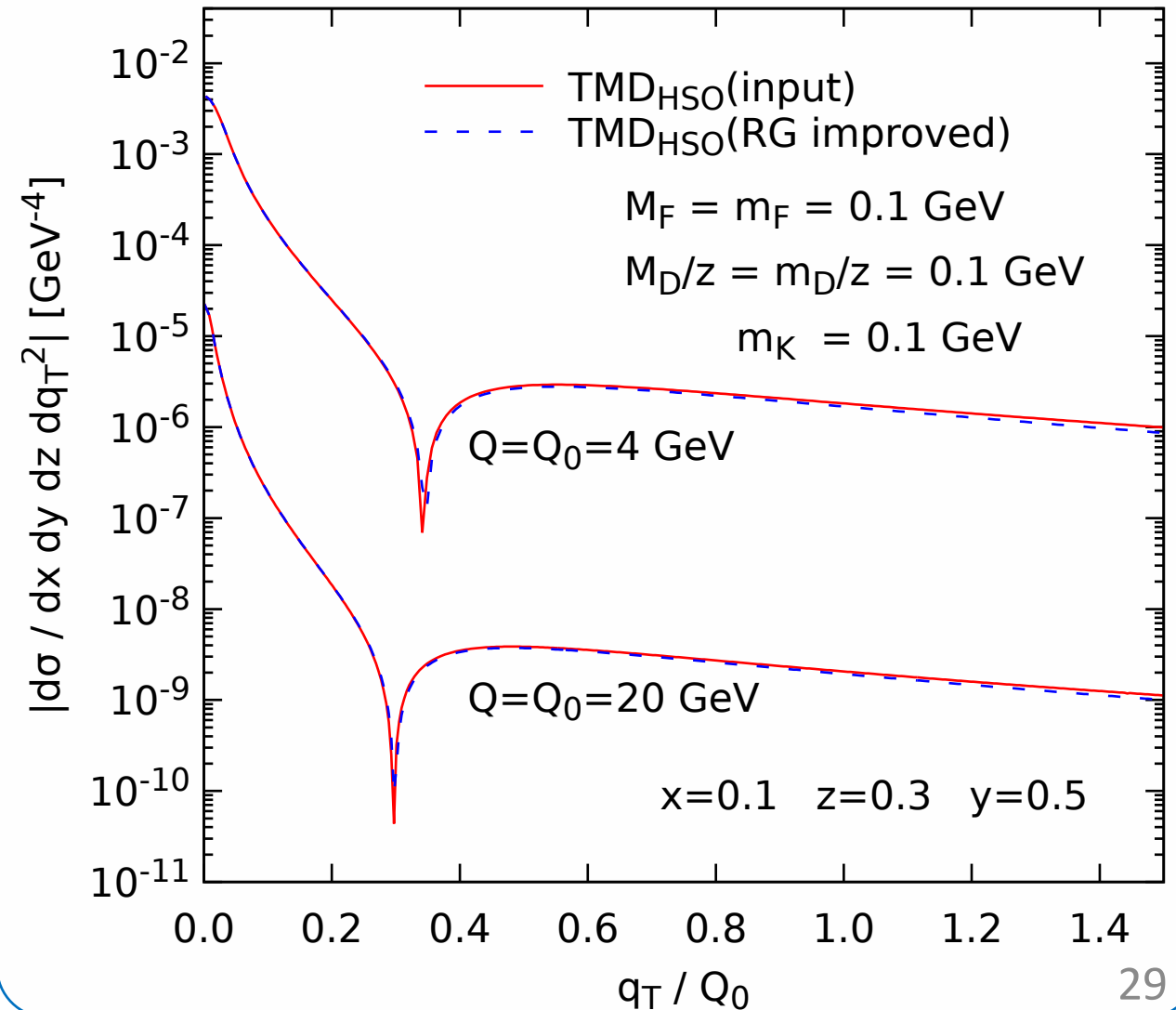
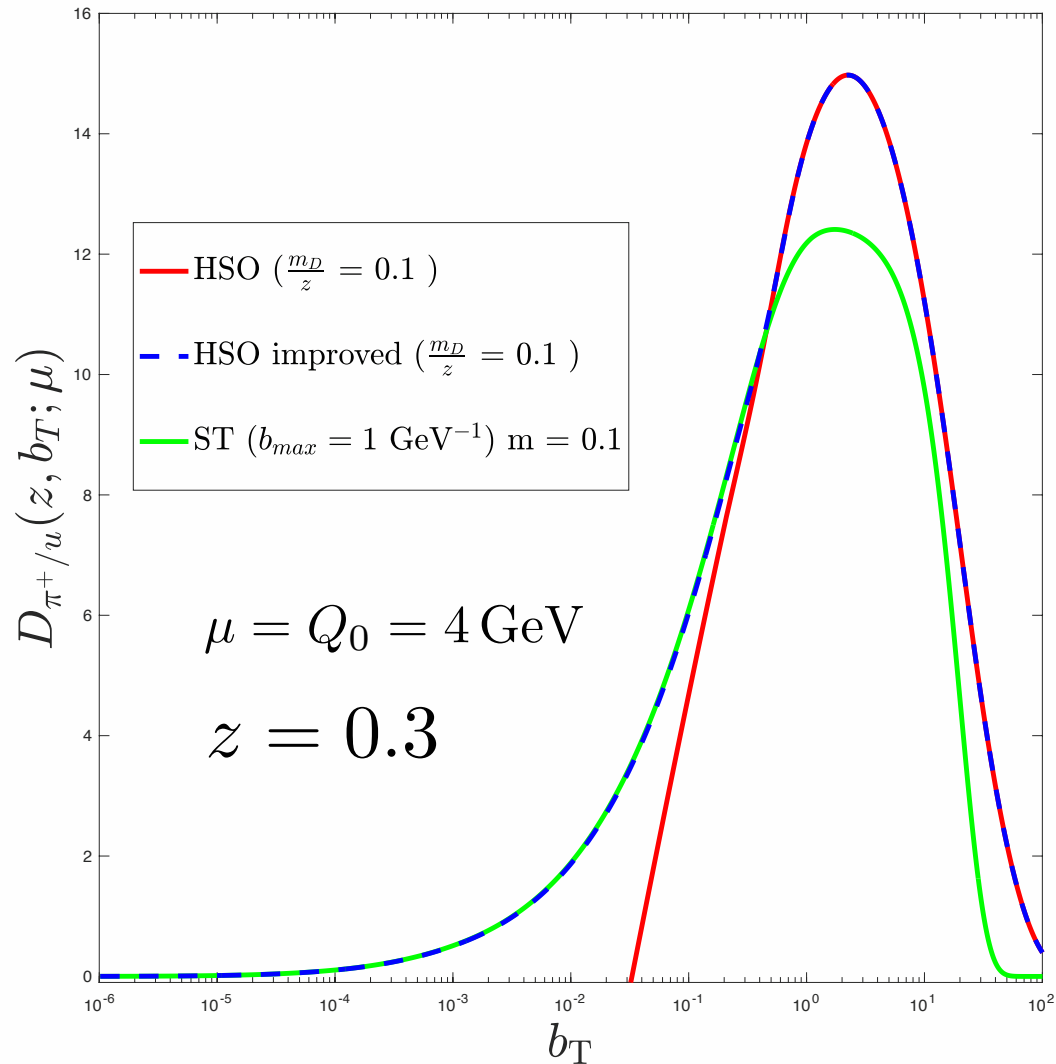
b_T – space RG improvement

$$\tilde{D}_{h/j}(z, \mathbf{b}_T; \mu_{Q_0}, Q_0^2) = \tilde{D}_{\text{inpt},h/j}(z, \mathbf{b}_T; \mu_{\bar{Q}_0}, \bar{Q}_0^2) E(\bar{Q}_0/Q_0, b_T)$$


$$\tilde{f}_{i/p}(x, \mathbf{b}_T; \mu_{Q_0}, Q_0^2) = \tilde{f}_{\text{inpt},i/p}(x, \mathbf{b}_T; \mu_{\bar{Q}_0}, \bar{Q}_0^2) E(\bar{Q}_0/Q_0, b_T).$$


$$E(\bar{Q}_0/Q_0, b_T) \equiv \exp \left\{ \int_{\mu_{\bar{Q}_0}}^{\mu_{Q_0}} \frac{d\mu'}{\mu'} \left[\gamma(\alpha_s(\mu'); 1) - \ln \frac{Q_0}{\mu'} \gamma_K(\alpha_s(\mu')) \right] + \ln \frac{Q_0}{\bar{Q}_0} \tilde{K}_{\text{inpt}}(b_T; \mu_{\bar{Q}_0}) \right\}$$

b_T – space RG improvement (an example)



Collins-Soper kernel

$$K_{\text{inpt}}^{(1)}(k_T; \mu_{Q_0}) = \frac{\alpha_s(\mu_{\overline{Q_0}}) C_F}{\pi^2} \frac{1}{k_T^2 + m_K^2} + C_K \delta^{(2)}(\mathbf{k}_T)$$


$$K_{\text{inpt}}^{(1)}(b_T; \mu_{Q_0}) = \frac{2\alpha_s(\mu_{\overline{Q_0}}) C_F}{\pi} K_0(b_T m_K) + C_K$$


$$\tilde{K}(b_T; \mu_{Q_0}) = \frac{2\alpha_s(\mu_{\overline{Q_0}}) C_F}{\pi} \left[K_0(b_T m_K) + \ln \left(\frac{m_K}{\mu_{\overline{Q_0}}} \right) \right] - \int_{\mu_{\overline{Q_0}}}^{\mu_{Q_0}} \frac{d\mu'}{\mu'} \gamma_K(\alpha_s(\mu'))$$