

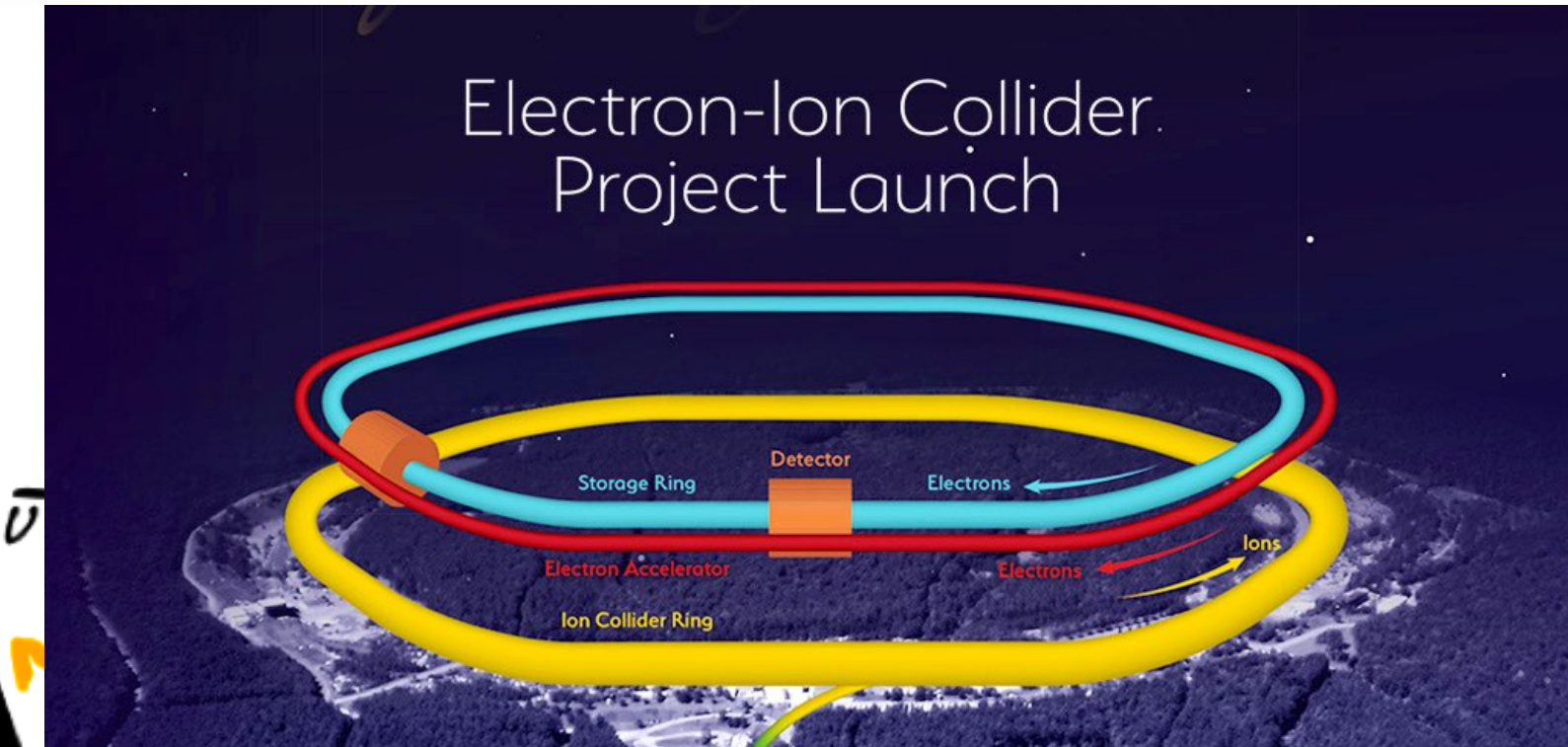
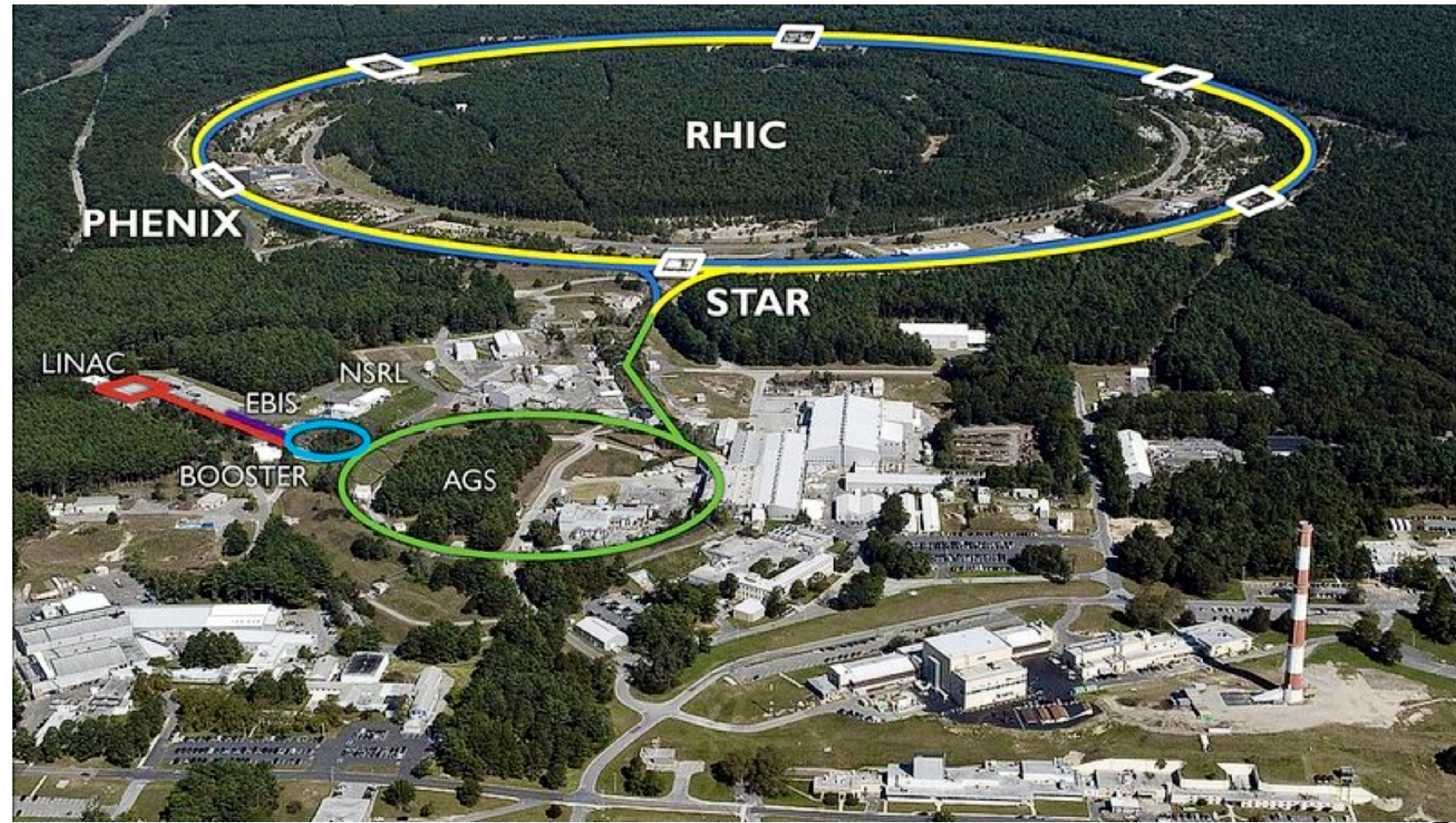


Lattice QCD: A Journey from RHIC to EIC

Swagato Mukherjee

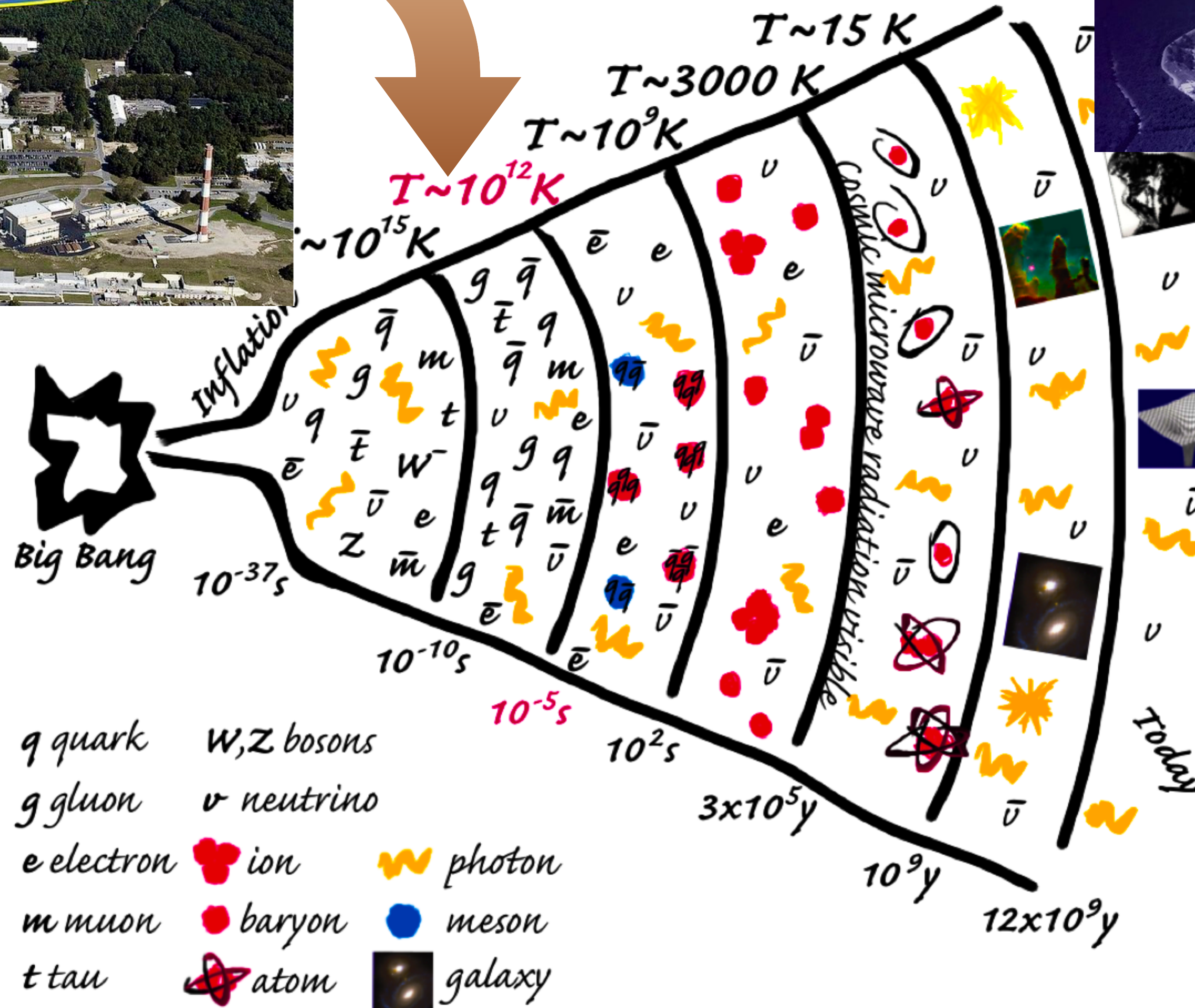
April 2023, GHP, Minneapolis, MN

my 14-year long journey across 14 billion years



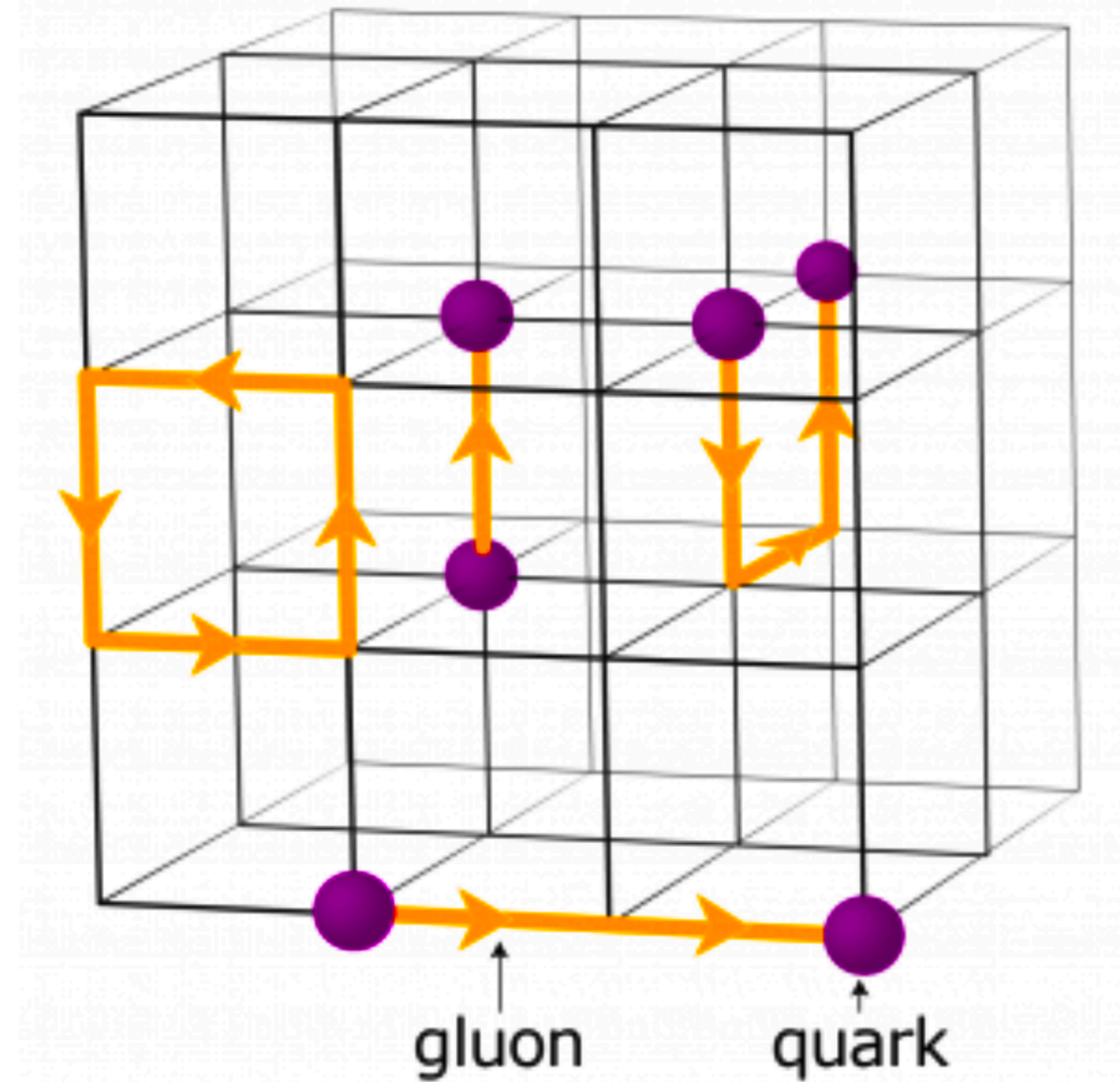
RHIC

EIC



the time machine — lattice QCD

QCD in its full glory: non-perturbative, gauge invariant regularization of QCD path integral



almost all results shown in this talk are for physical quark masses and in the continuum limit

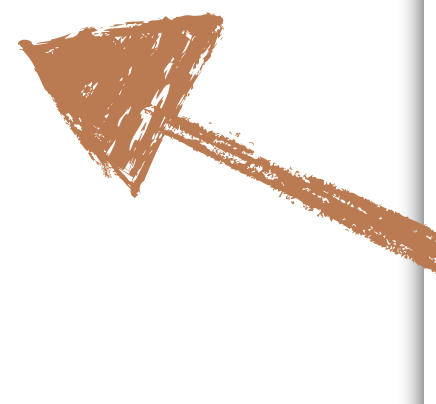


... swimming through quark gluon plasma ...

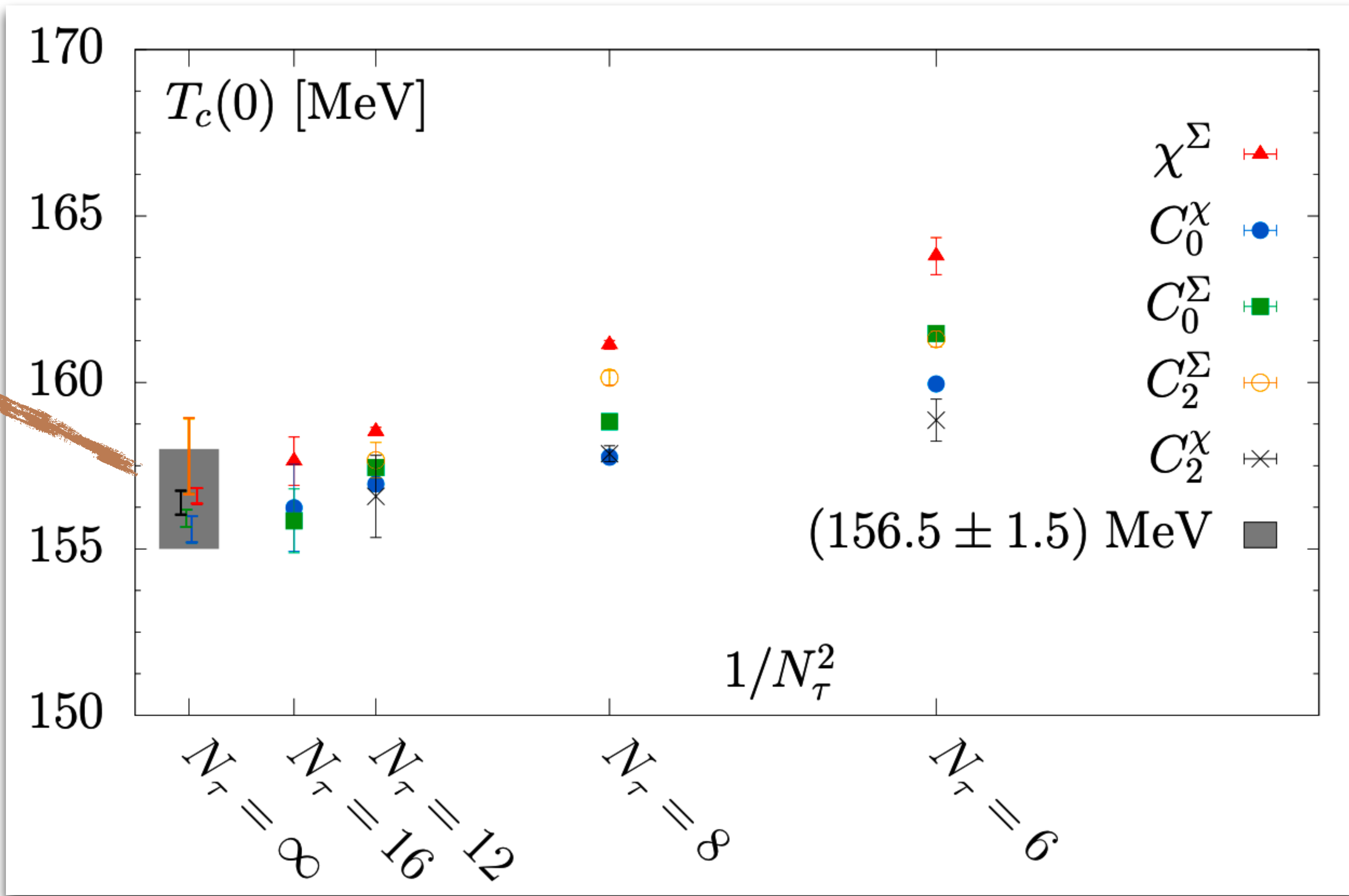


pseudocritical temperature

$T_c = 156.5 \pm 1.5 \text{ MeV}$



studying multiple observables related to punctuation of order parameter



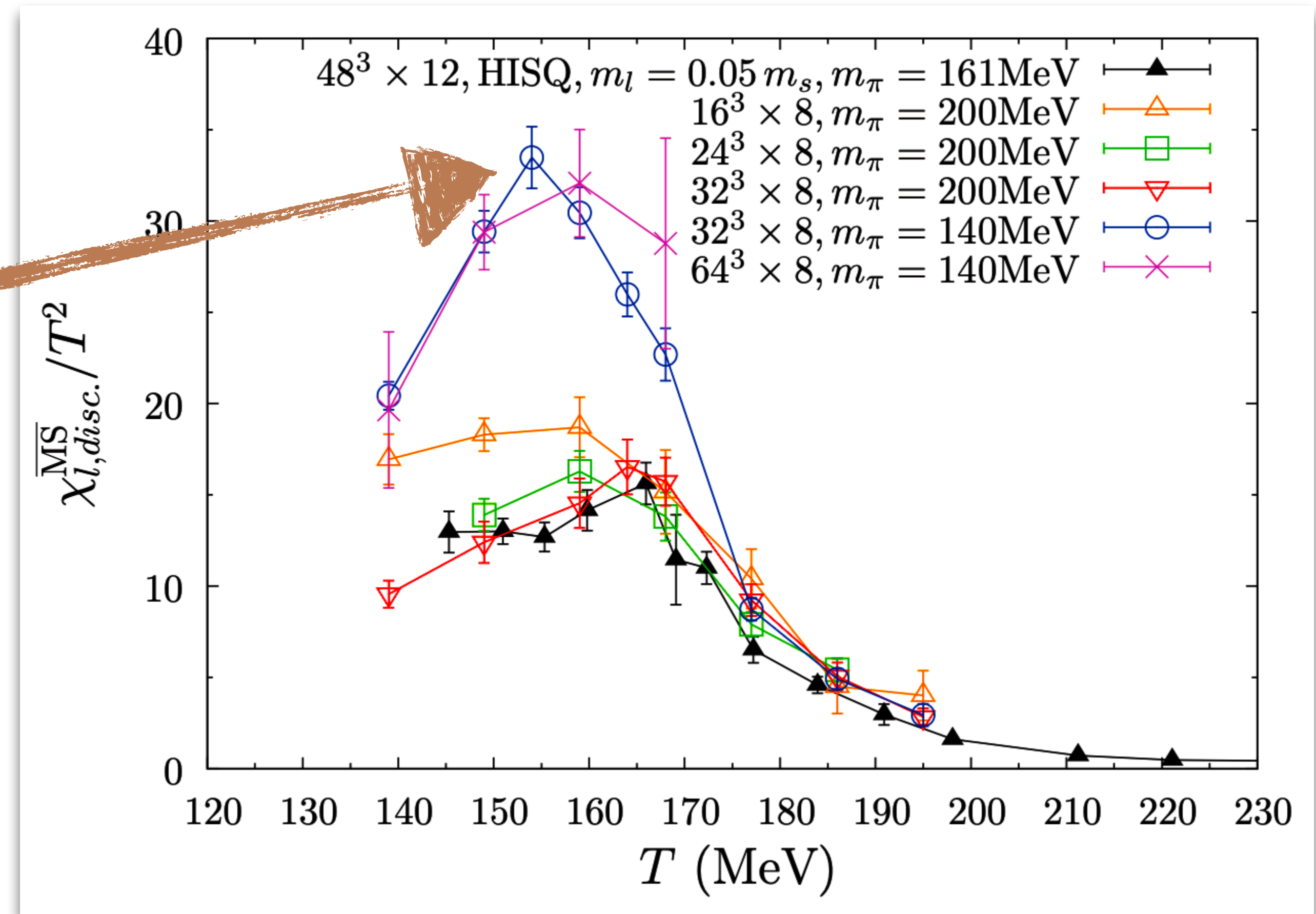
HotQCD, Phys. Lett. B 795,15 (2019)

HotQCD, Phys. Rev. D 85, 054503 (2012)

QCD transition is a crossover

no volume scaling observed
in fluctuation of the order
parameter

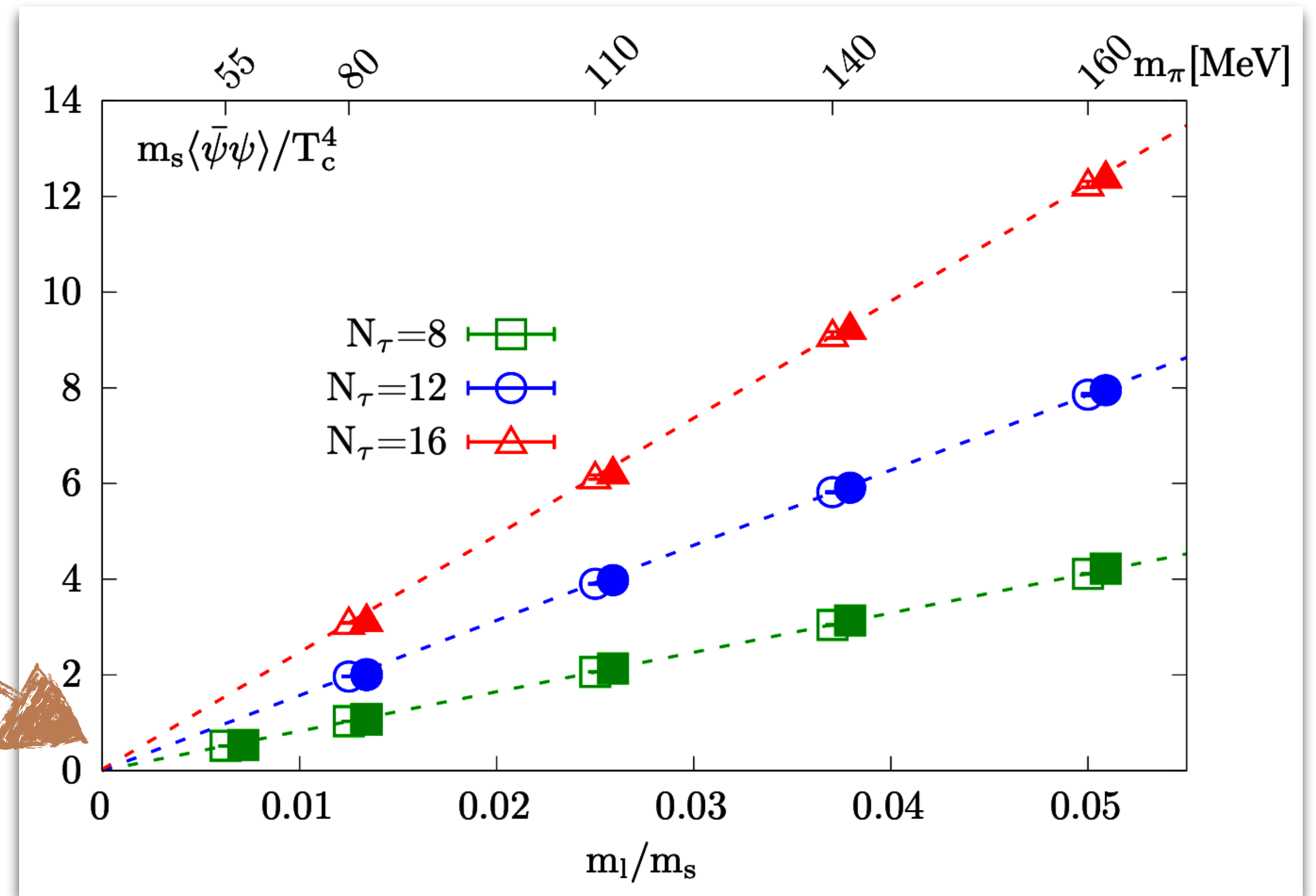
using chiral fermions



HotQCD, Phys. Rev. Lett. 113, 8, 082001(2014)

chiral symmetry is restored in the chiral limit ...

order parameter with
vanishing light quark
mass

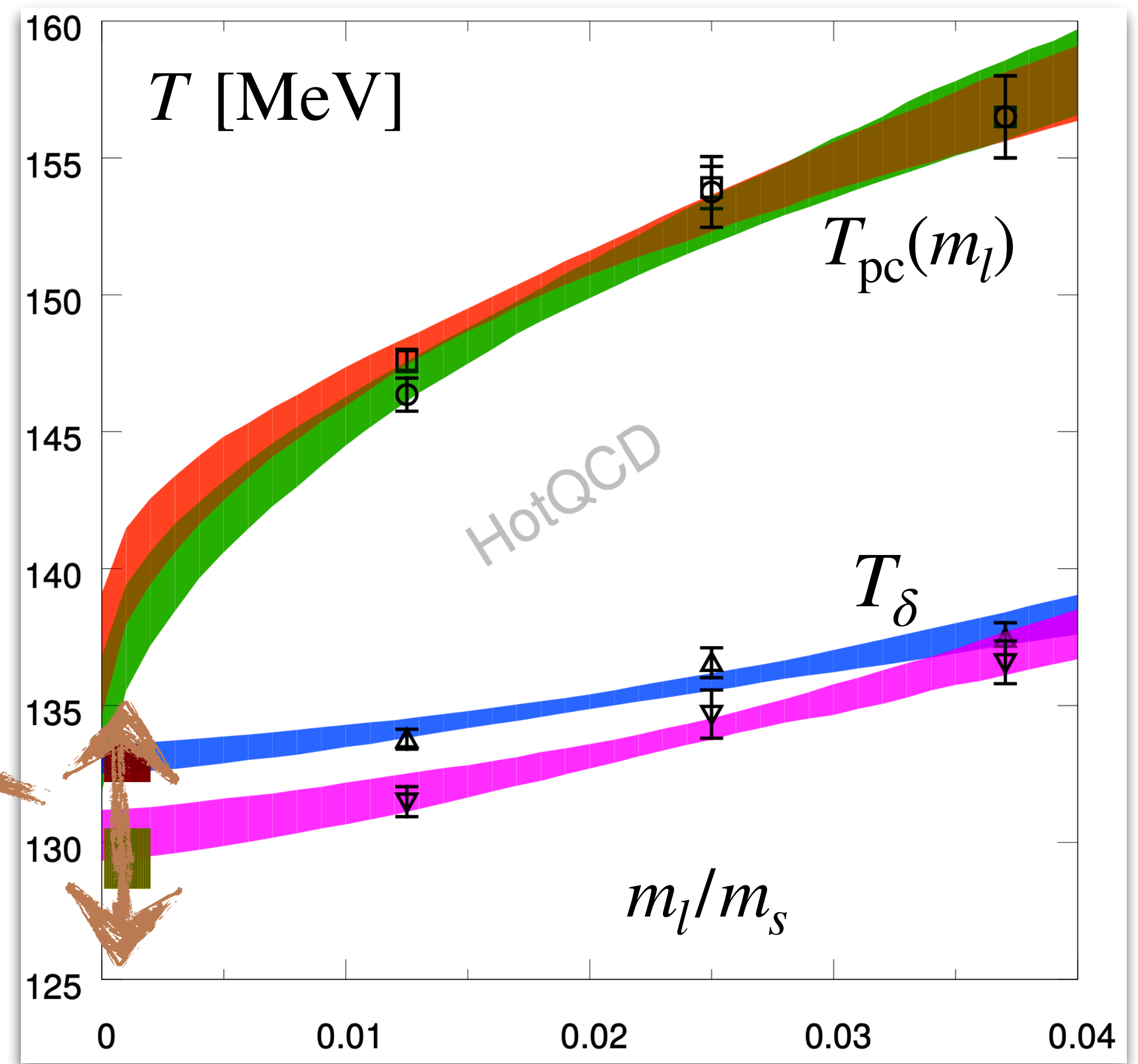


HotQCD, Phys. Rev. Lett. 126, 8, 082001 (2021)

chiral symmetry is restored in the chiral limit ...

second order transition in the chiral limit

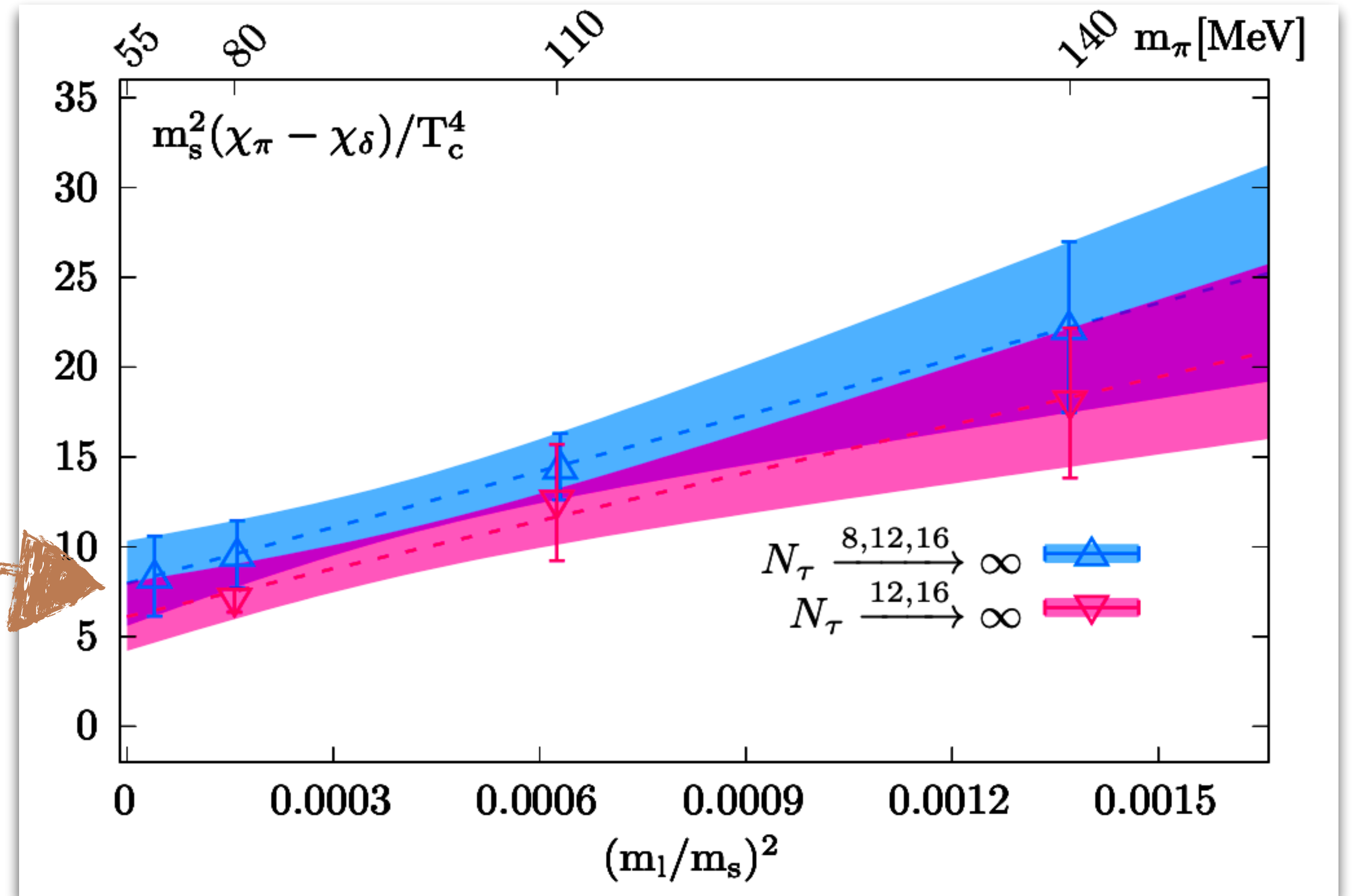
$$T_c^0 = 132^{+3}_{-6} \text{ MeV}$$



HotQCD, Phys. Rev. Lett. 123, 6, 062002 (2019)

... but anomalous axial symmetry remains broken

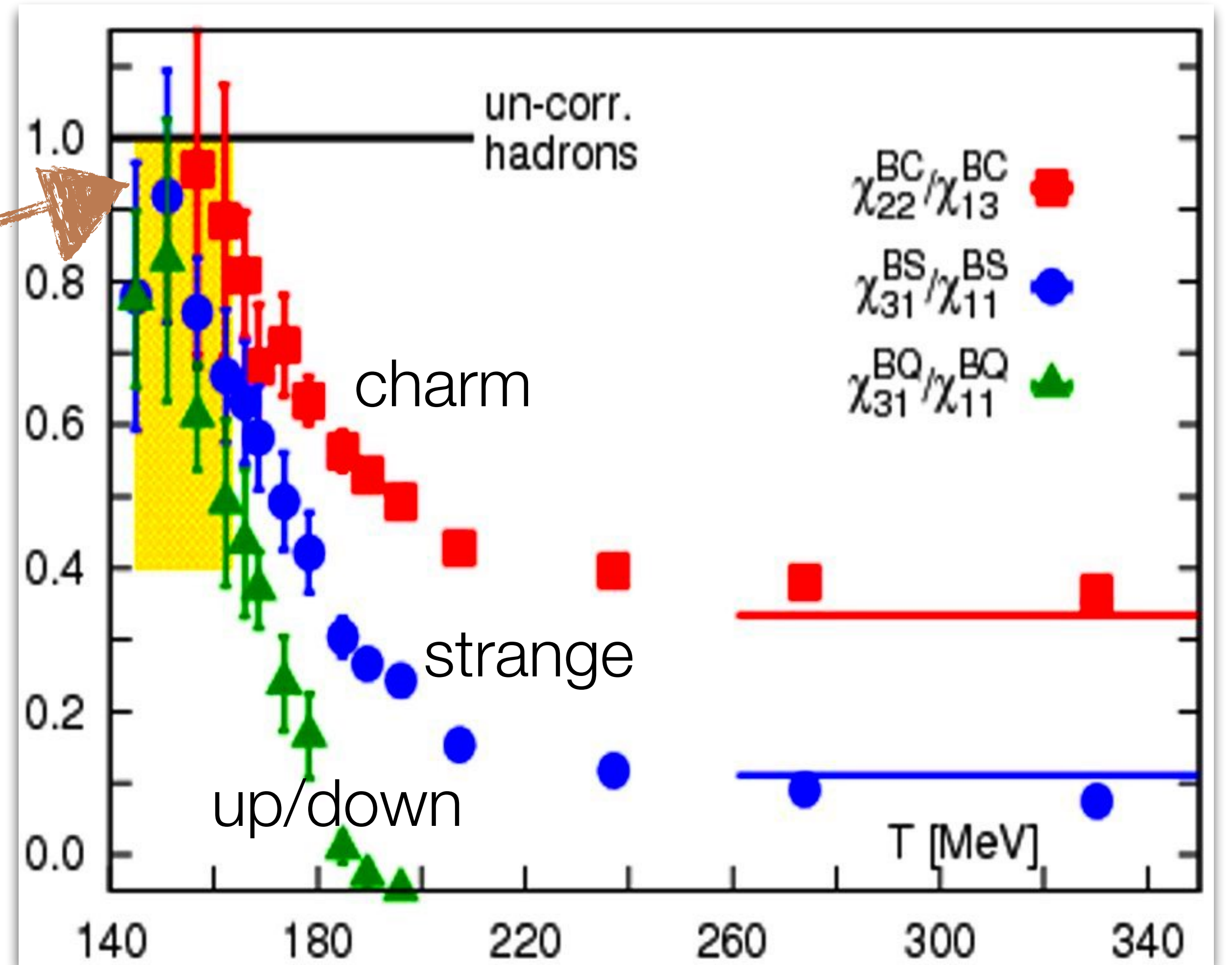
measures of axial symmetry breaking do not vanish in the chiral limit



HotQCD, Phys. Rev. Lett. 126, 8, 082001 (2021)

deconfinement – QGP

liberation of degrees of freedom carrying fractional charges around $T \approx 155$ MeV



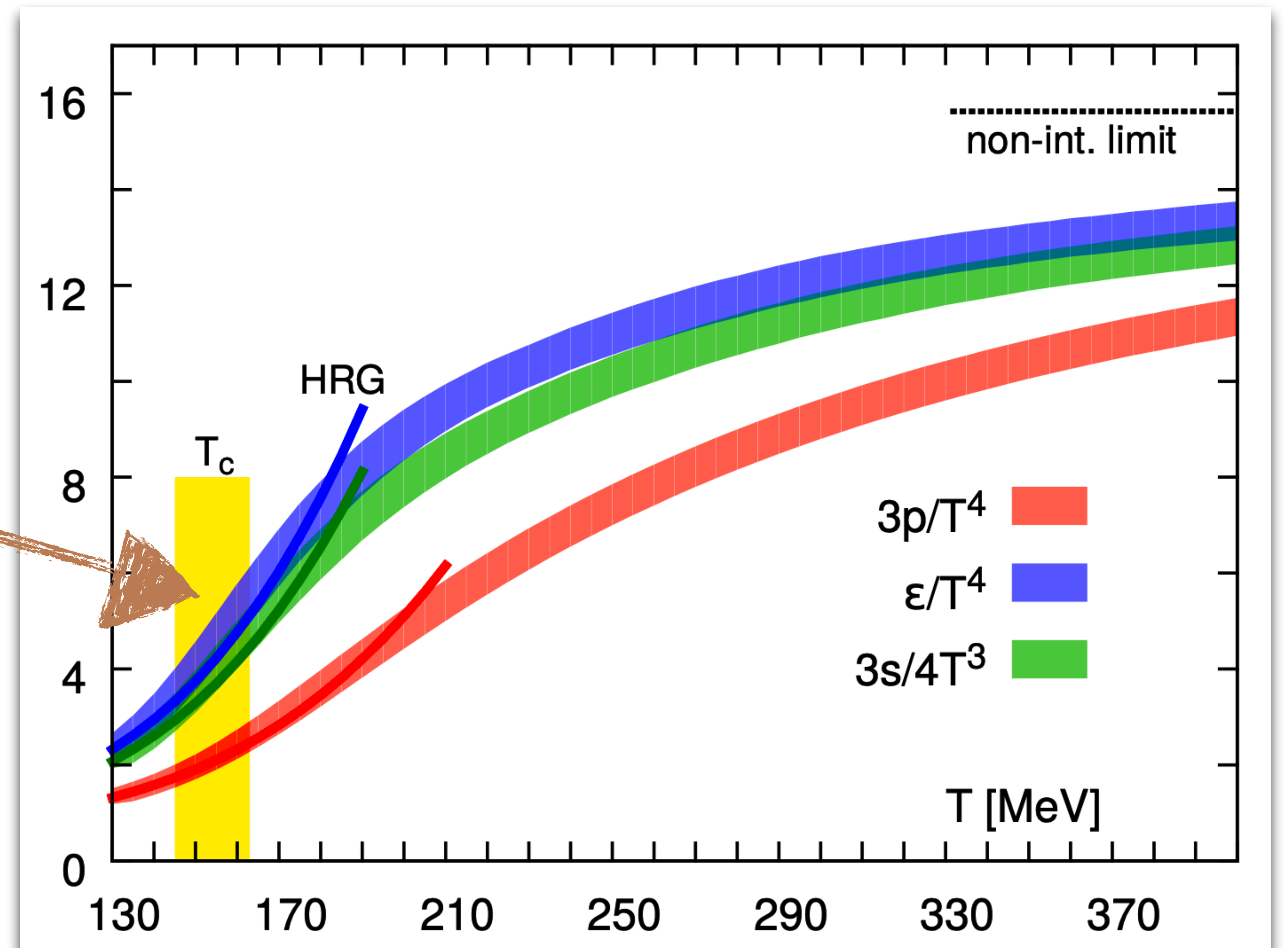
HotQCD, Phys. Rev. Lett. 111, 082301 (2013)

HotQCD, Phys. Lett. B 737, 210 (2014) 10

deconfinement – QGP

equation of state

rapidly rising pressure, energy
& entropy density around
 $T \approx 155$ MeV



HotQCD, Phys. Rev. D 90, 094503 (2014)

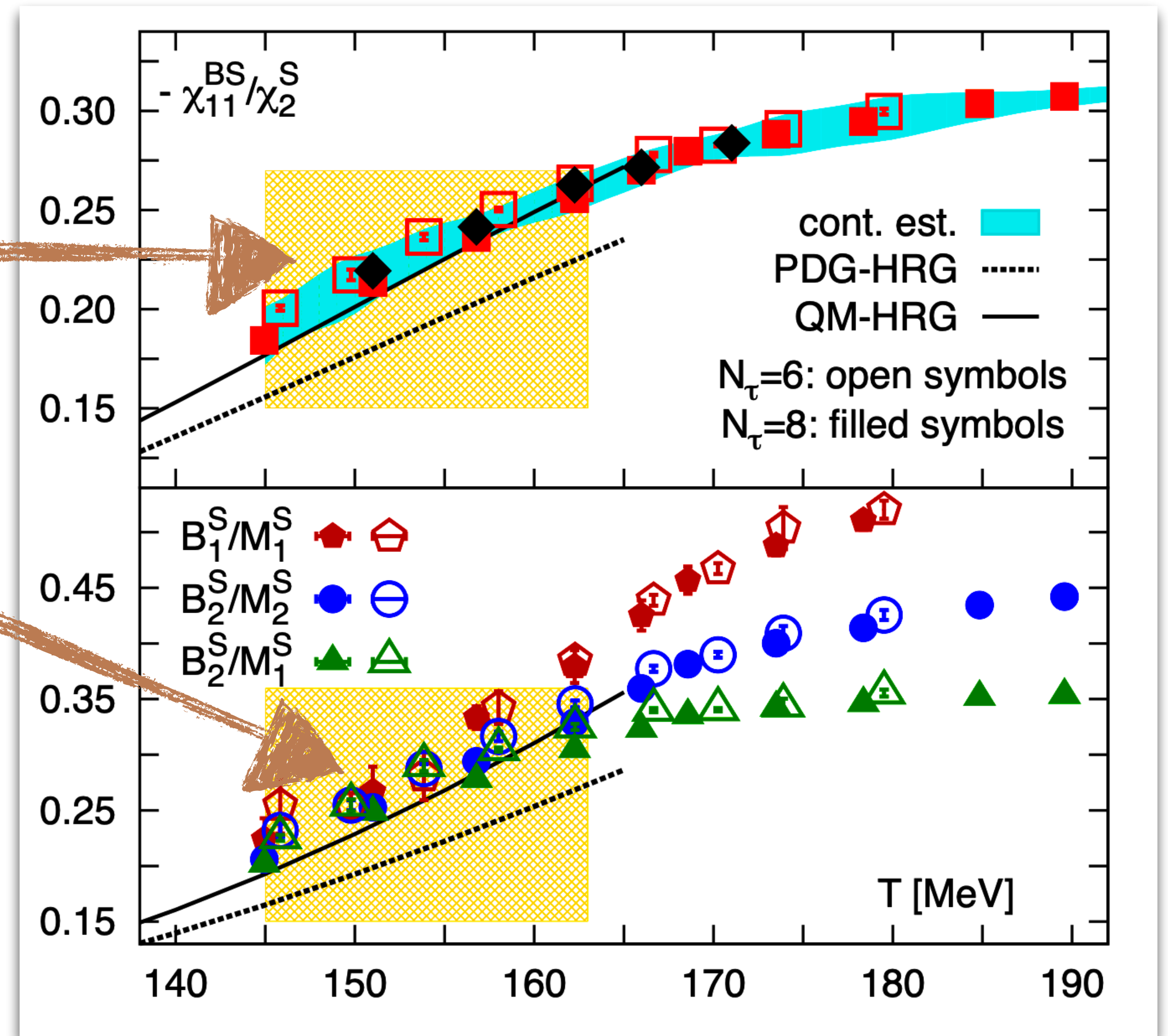
HotQCD, Phys. Rev. D 81, 054504 (2010)

signatures of yet unobserved hadrons

partial pressures of strange baryons can be described only by accommodating yet unobserved strange baryons

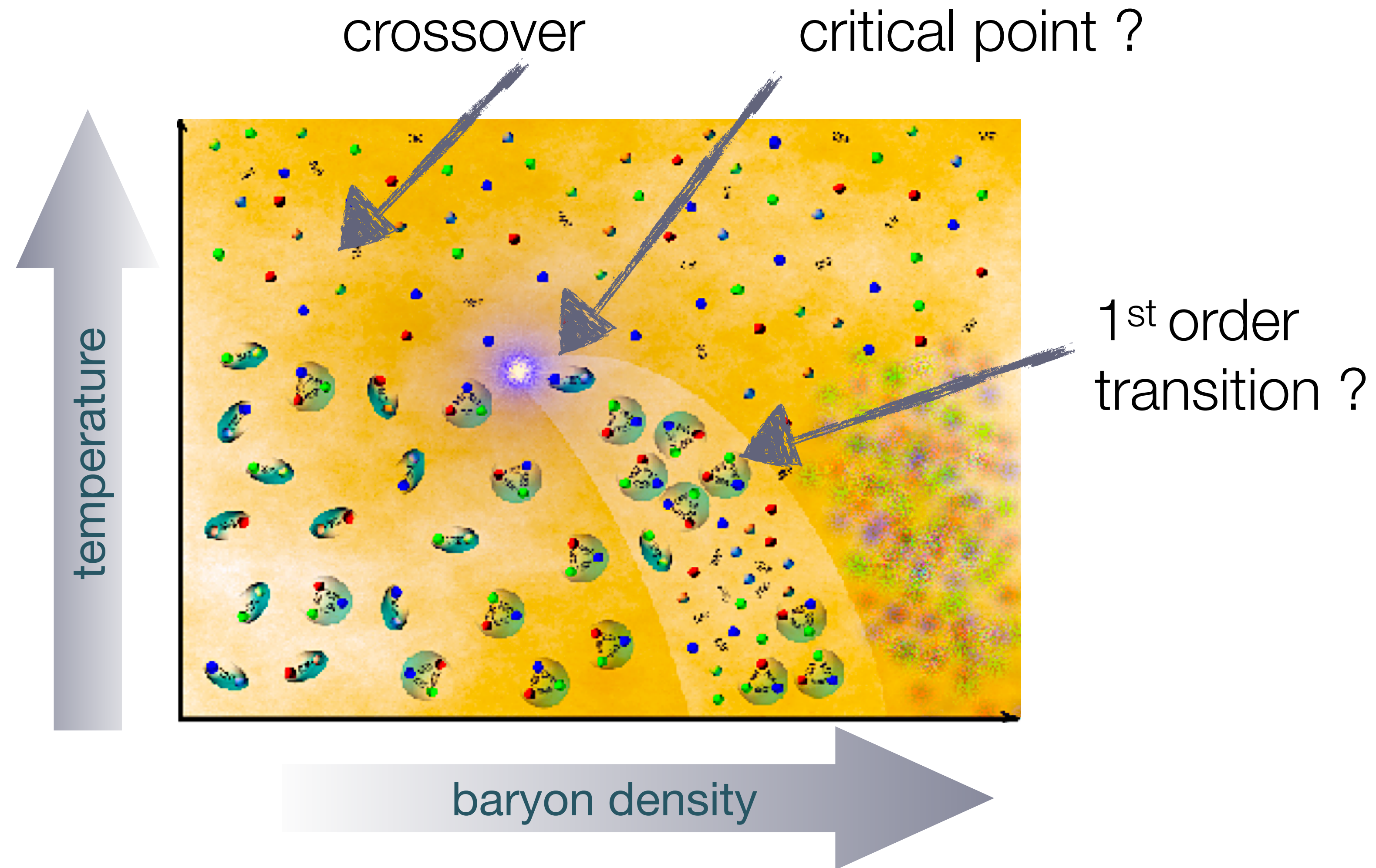
same for charm baryons

HotQCD, Phys. Lett. B 737, 210 (2014)



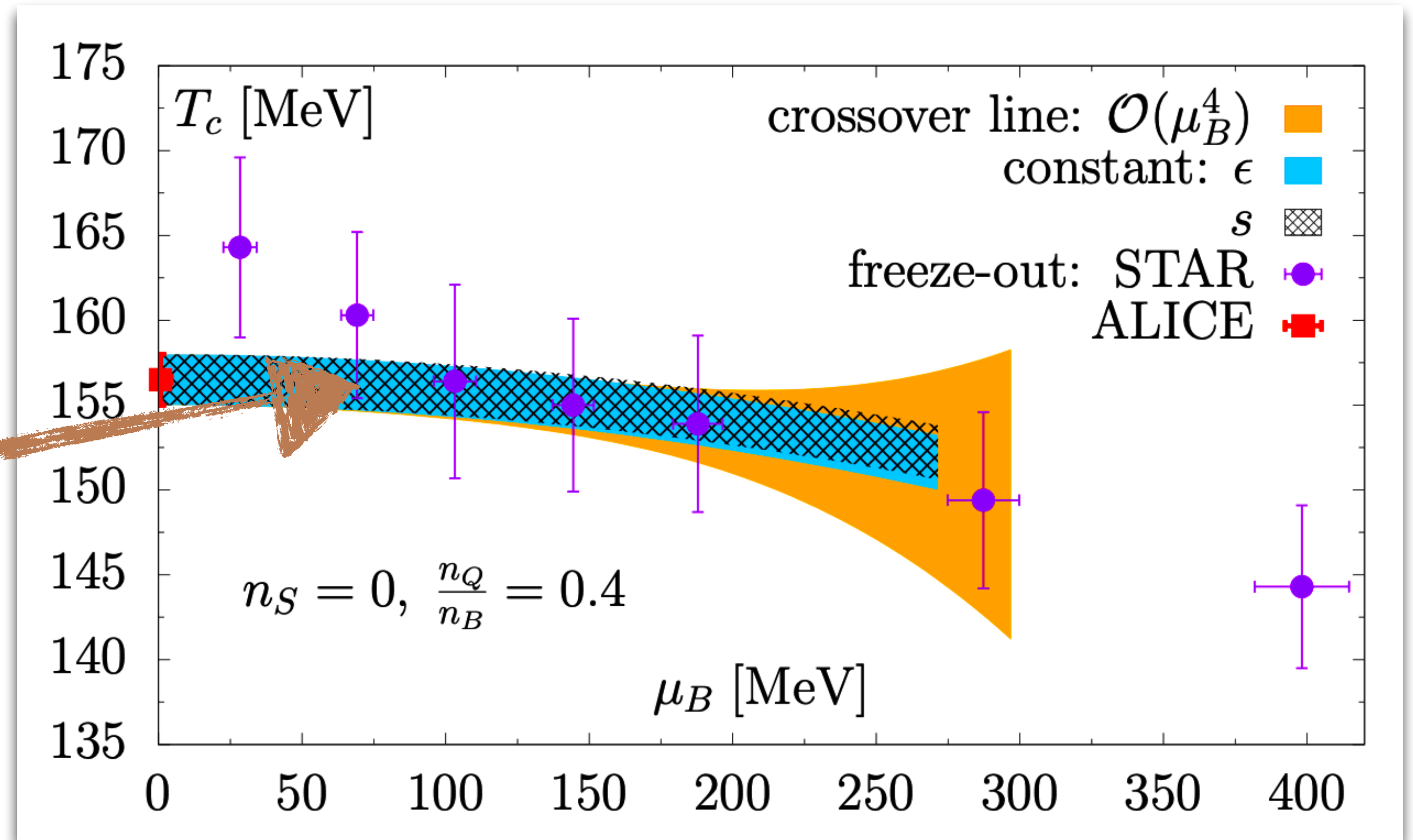
HotQCD, Phys. Rev. Lett. 113, 7, 072001 (2014)

... wondering through QCD phase diagram ...



QCD phase boundary in $T - \mu_B$ plane

freeze-out (hadronization)
at RHIC BES takes place
along the QCD phase
boundary

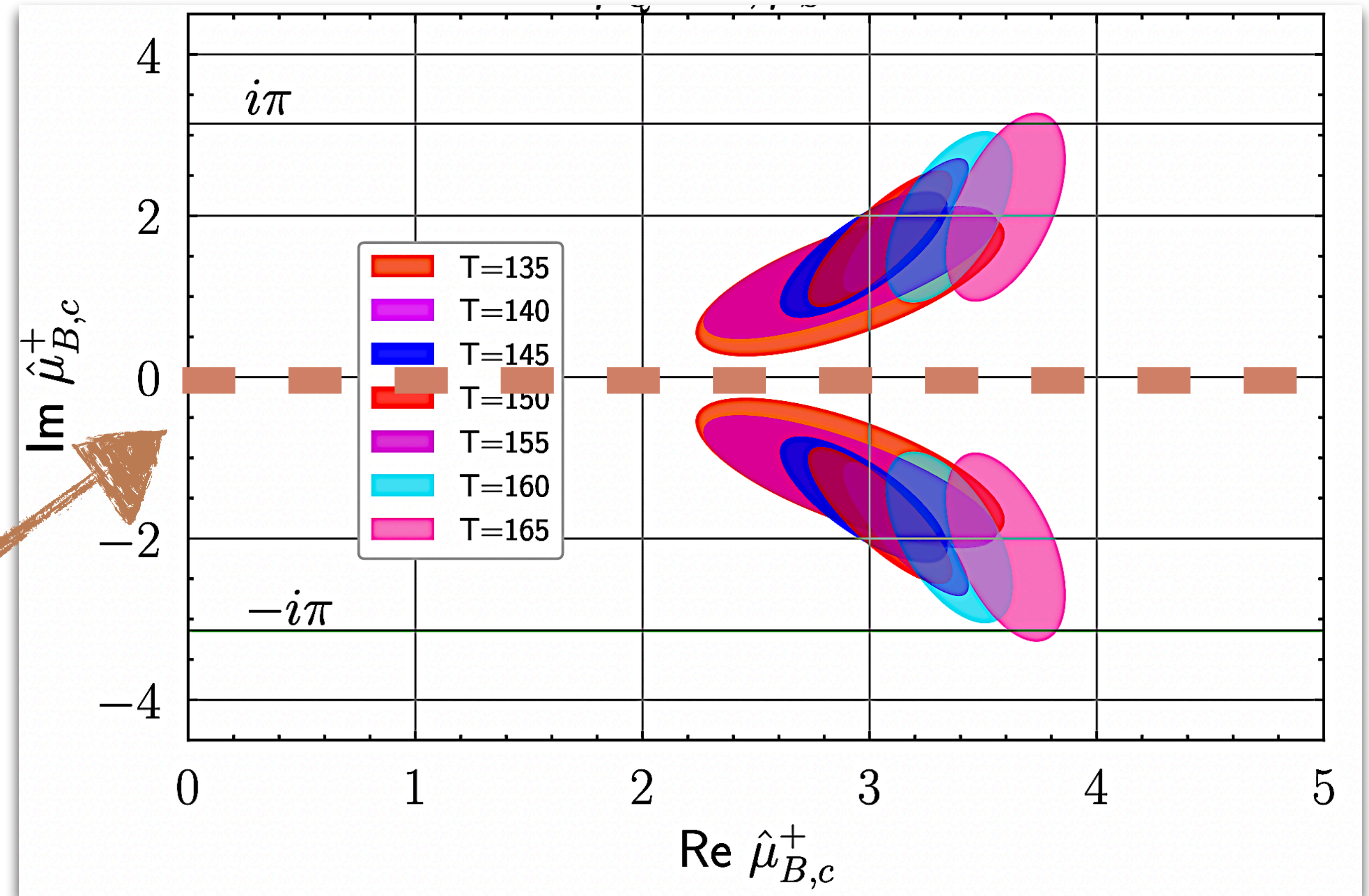


HotQCD, Phys. Lett. B 795,15 (2019)

where is QCD critical point ?

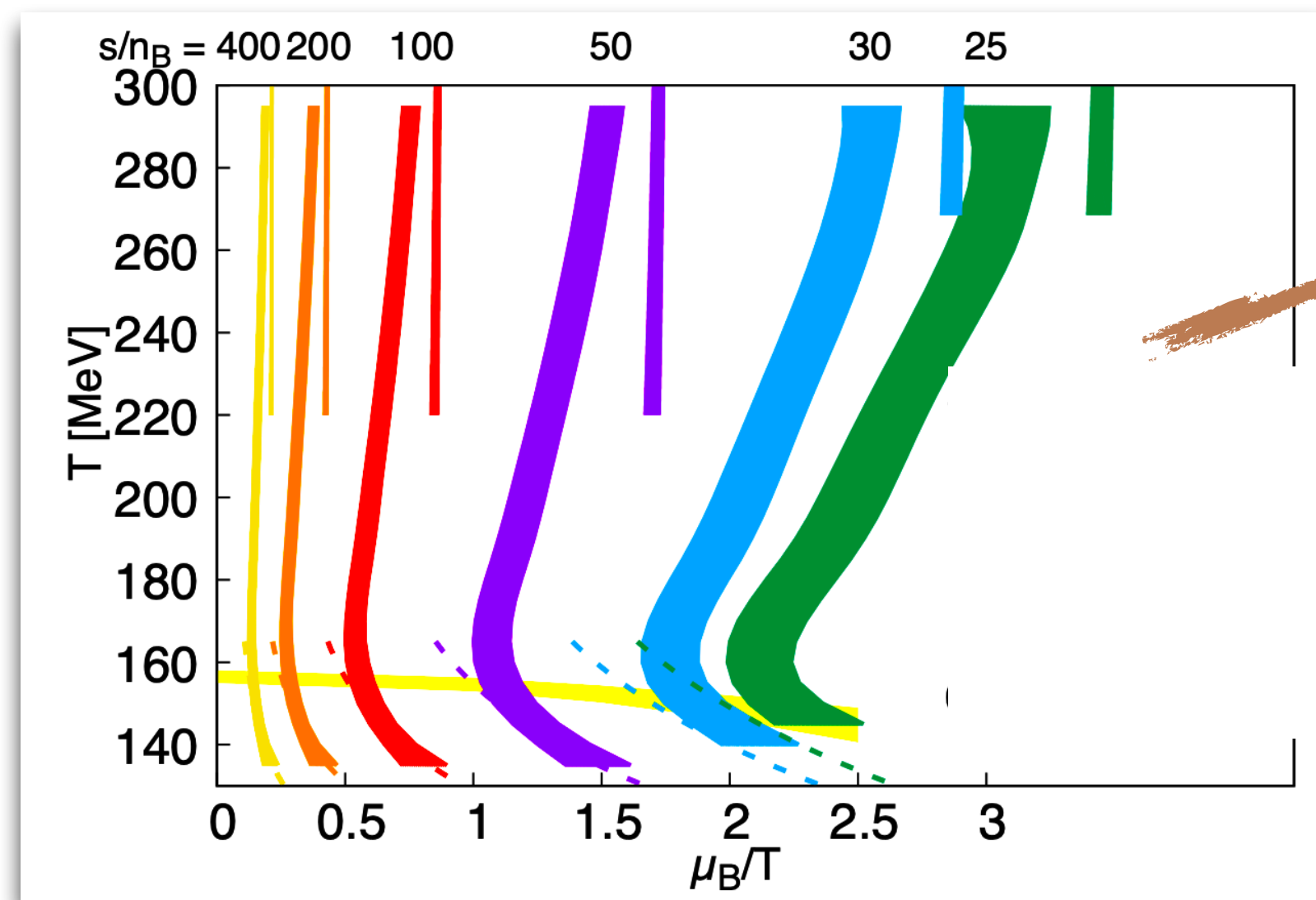
looking for zeros of QCD partition function in complex- μ_B plane

no hint of critical point for temperatures > 135 MeV

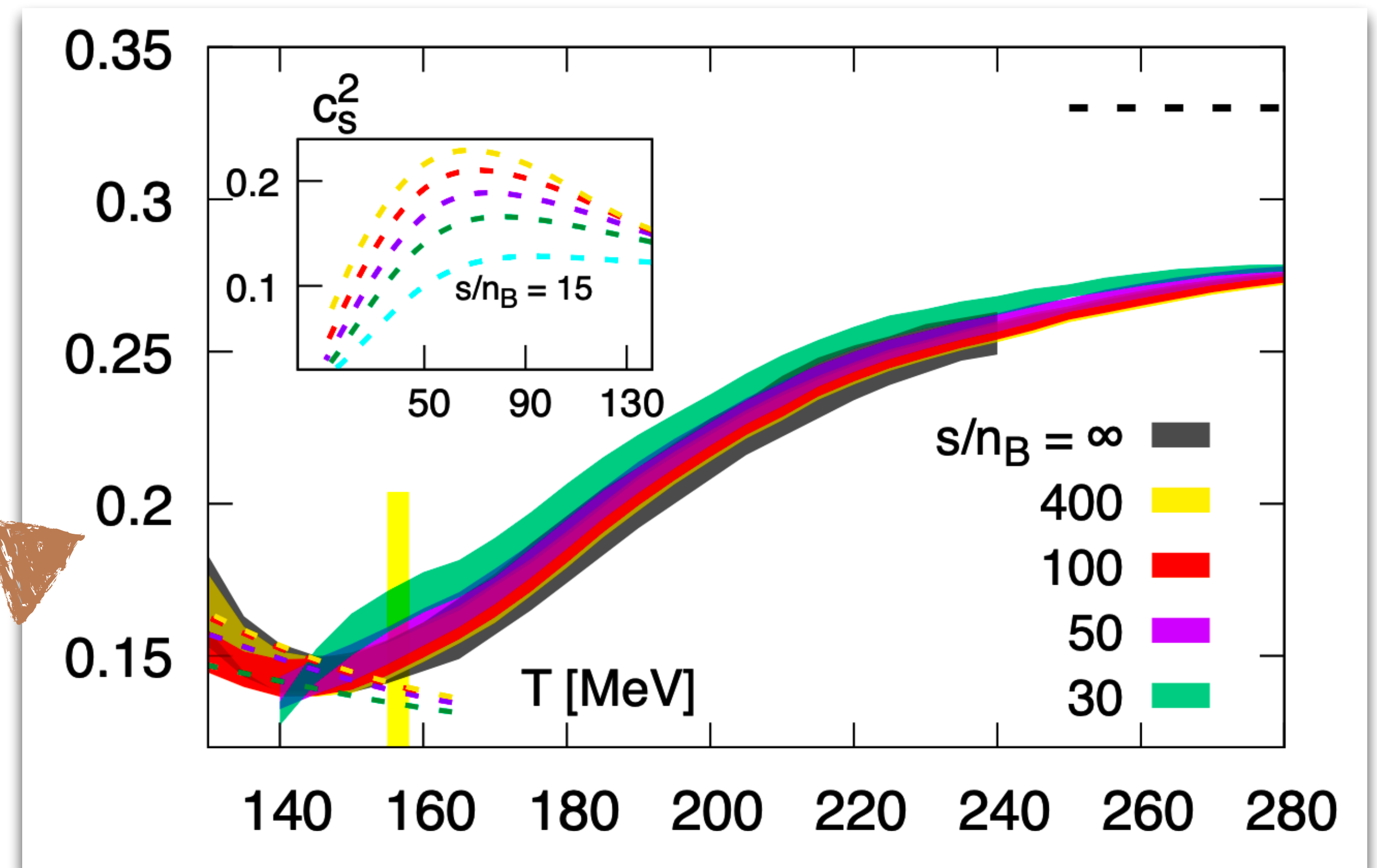


equation of state at $\mu_B > 0$

along the trajectories of
ideal hydrodynamic
evolution in the $T - \mu_B$
plane



speed of sound



HotQCD, 2212.09043

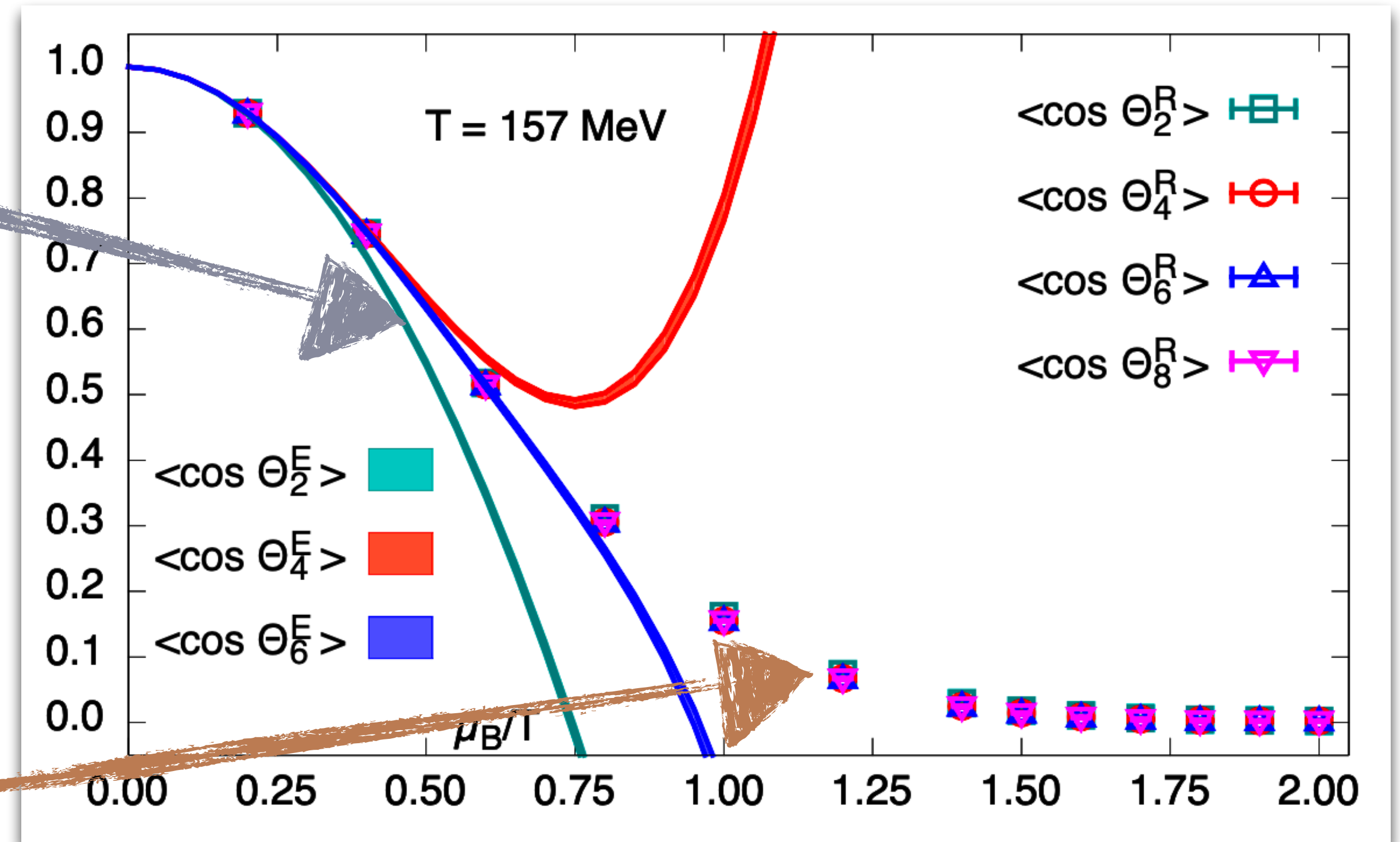
HotQCD, Phys. Rev. D 95, 054504 (2017)

QCD equation of state at $\mu_B > 0$

average phase of the
QCD partition function

very slow convergence of Taylor
series in powers of μ_B

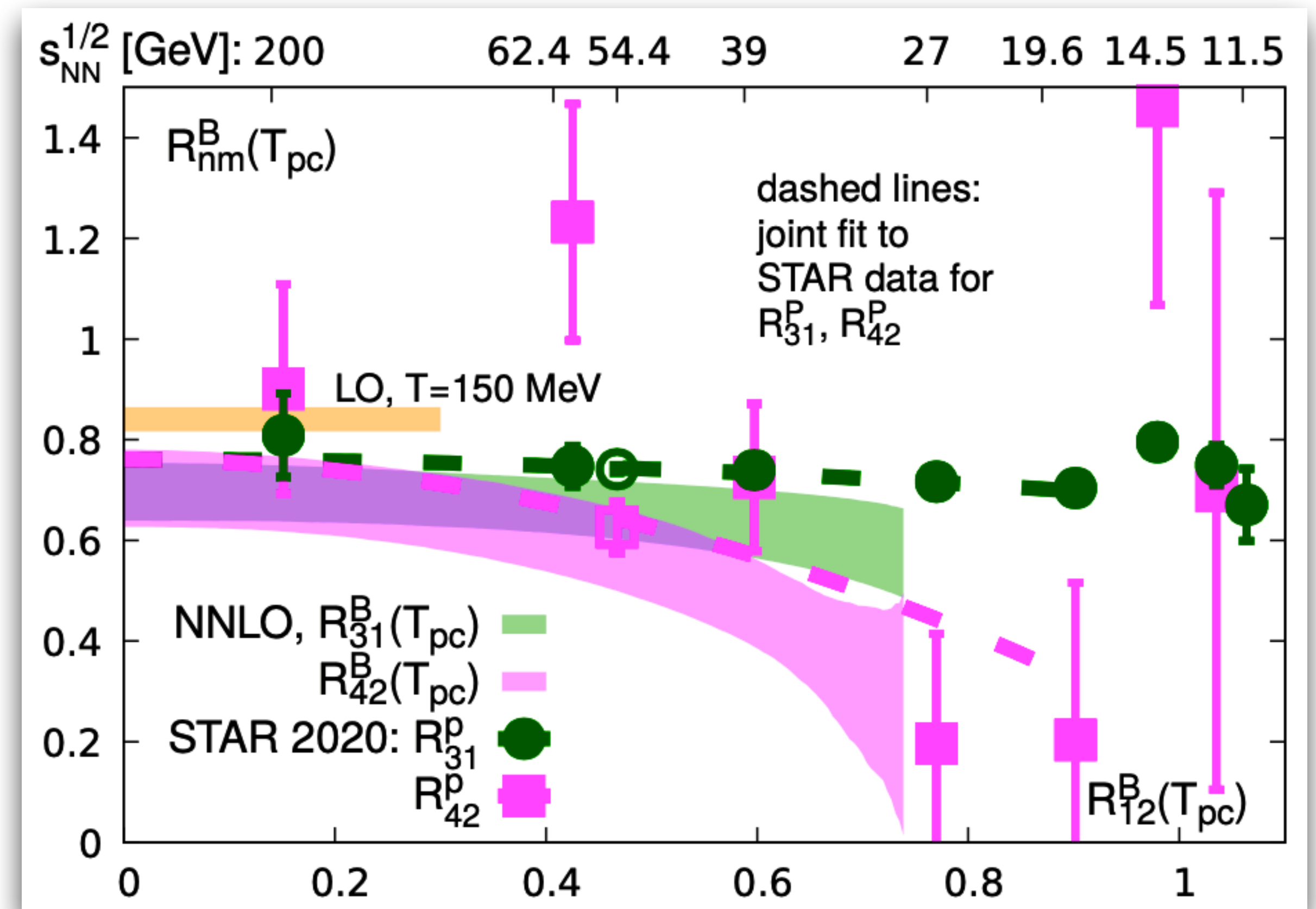
resumming up to 8-point
correlation functions of baryon
current to all orders in μ_B



lattice meets experiments: from $T - \mu_B$ to \sqrt{s}

lattice QCD: ratios of cumulants of net baryon fluctuations along $T - \mu_B$ phase boundary

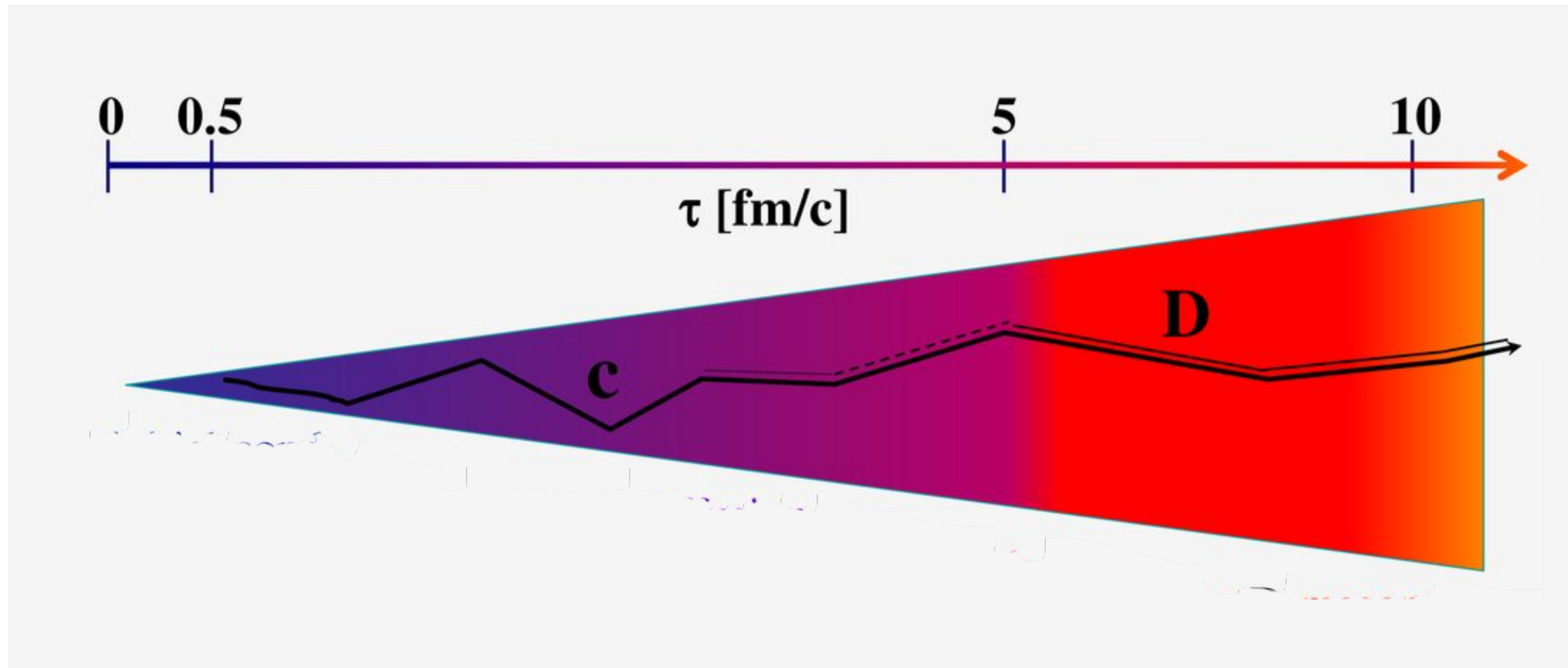
STAR: ratios of cumulants of net proton fluctuations for different collision energies (\sqrt{s})



HotQCD, Phys. Rev. Lett. 109, 192302 (2012)

HotQCD, Phys. Rev. D 101, 7, 074502 (2020)

... dancing with heavy quarks in QGP ...

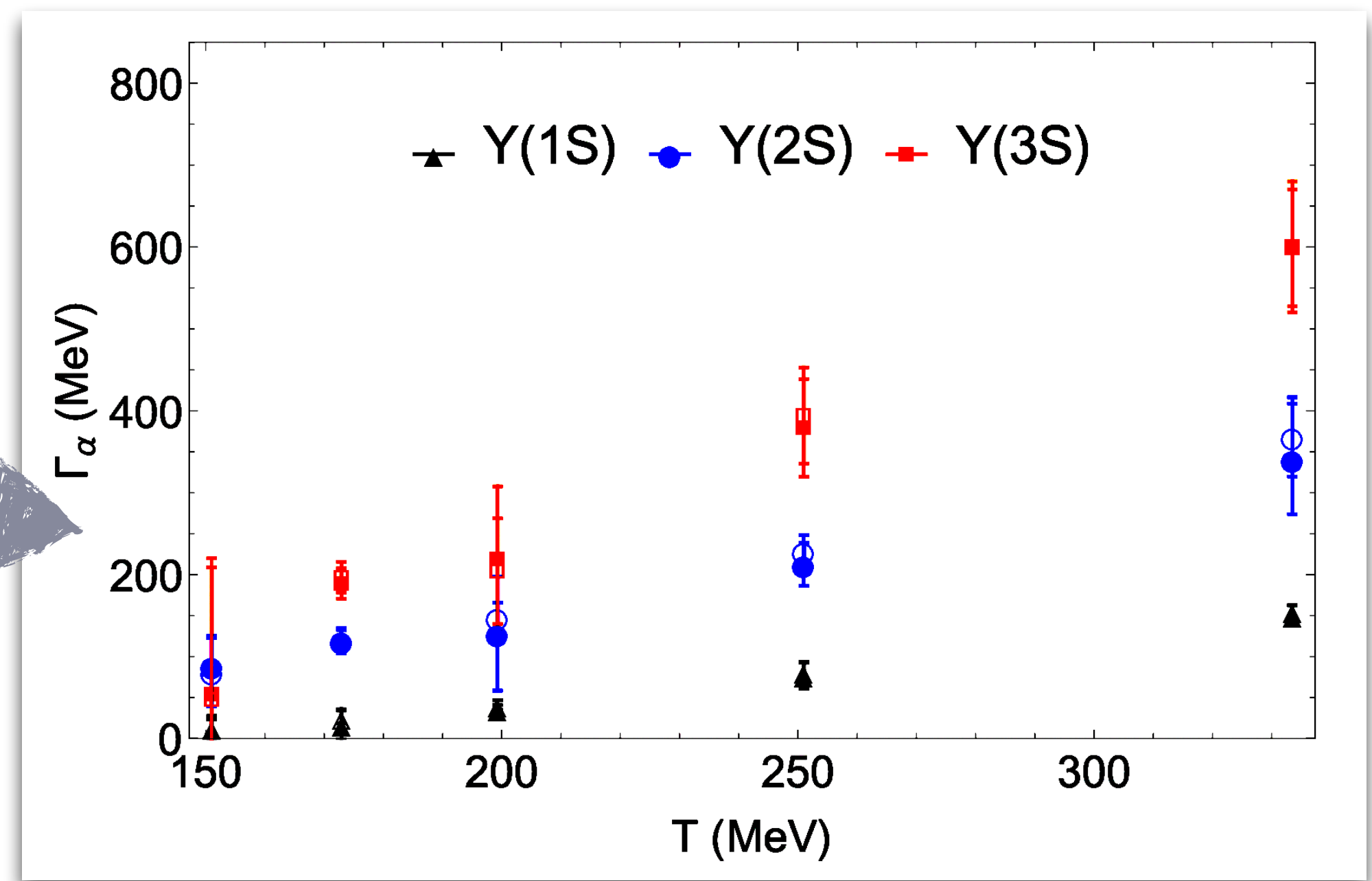
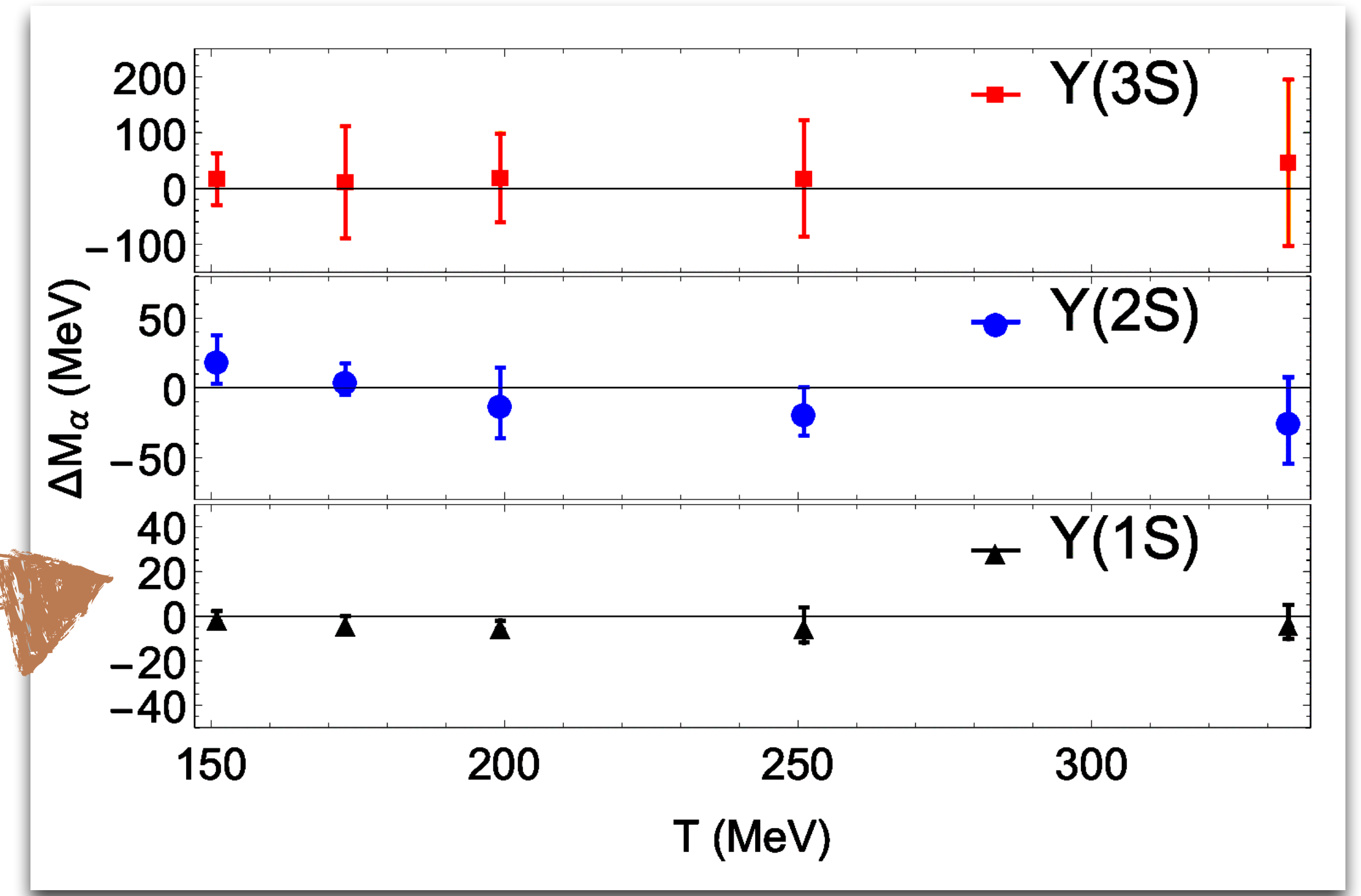


excited bottomonia in QGP

first lattice QCD results for up to 3S & 2P bottomonia in QGP

small thermal mass shift

large thermal widths



Larsen *et al.*, Phys. Lett. B 800, 135119 (2020)

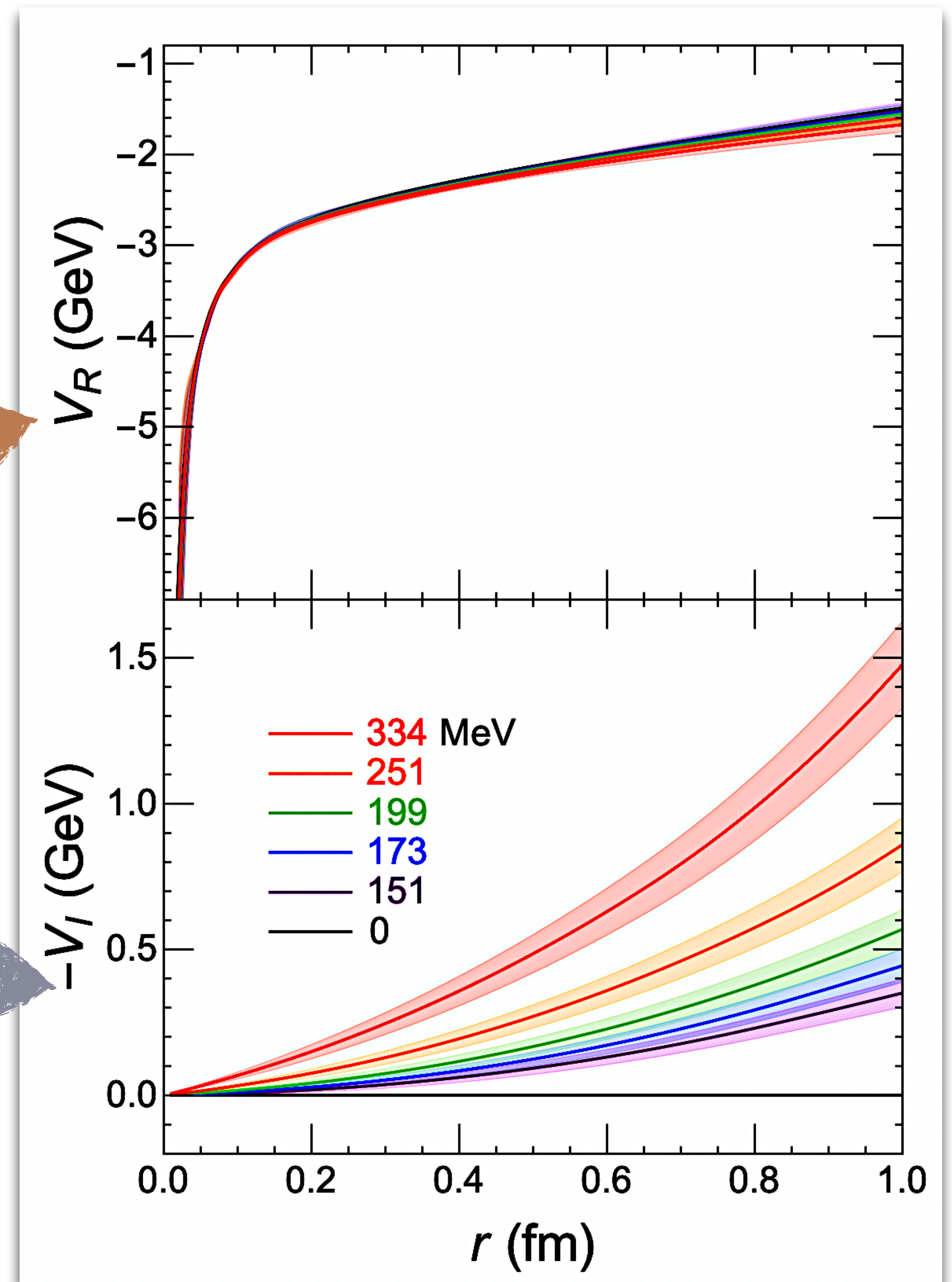
heavy quark interactions in QGP

heavy quark potential

from neural network fits to in-medium
bottomonia masses and widths

real part: almost no
temperature dependence

imaginary part: strong
temperature dependence



Shi et al., Phys. Rev. D 105, 1, 014017 (2022)

heavy quark interactions in QGP

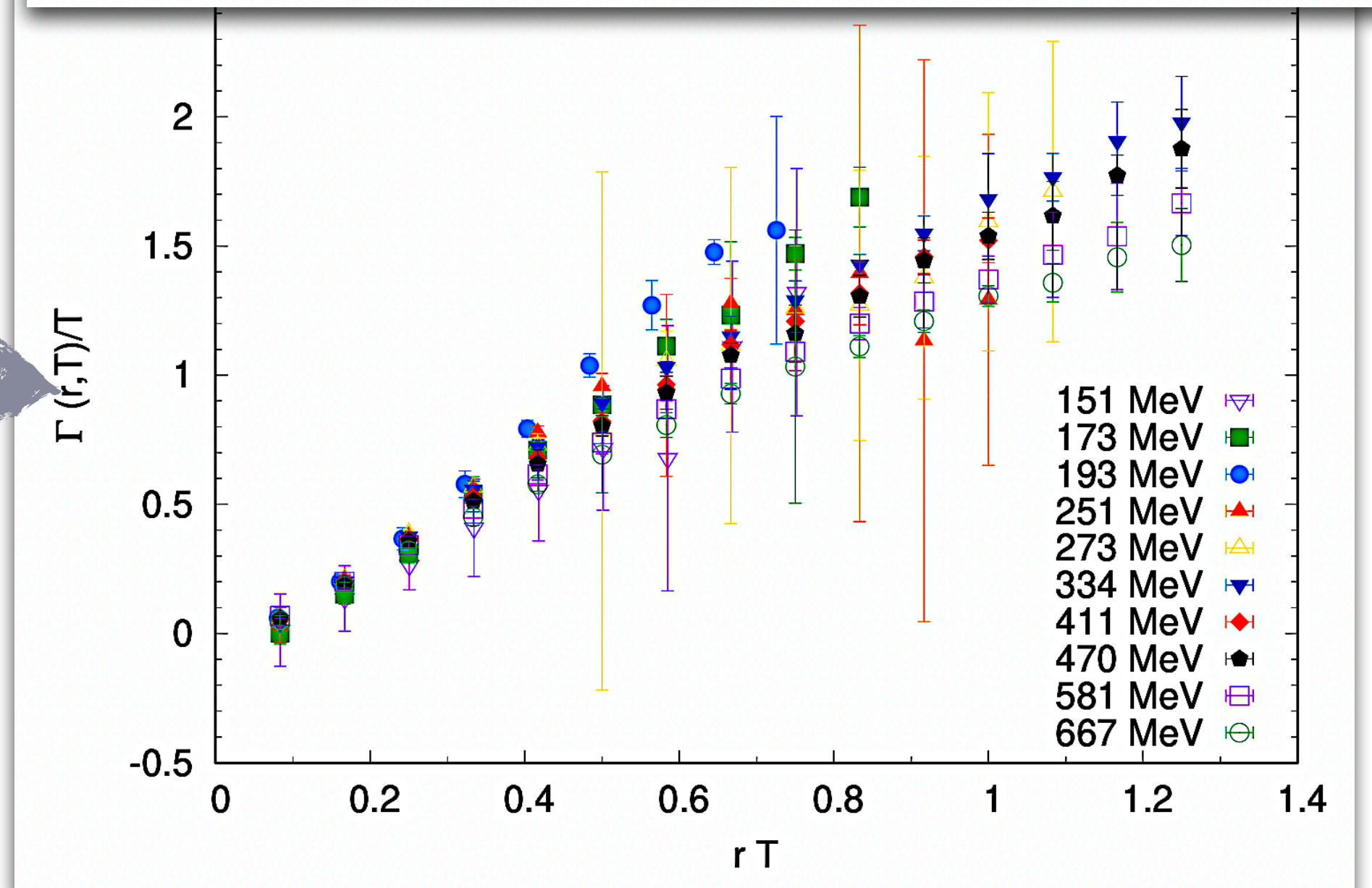
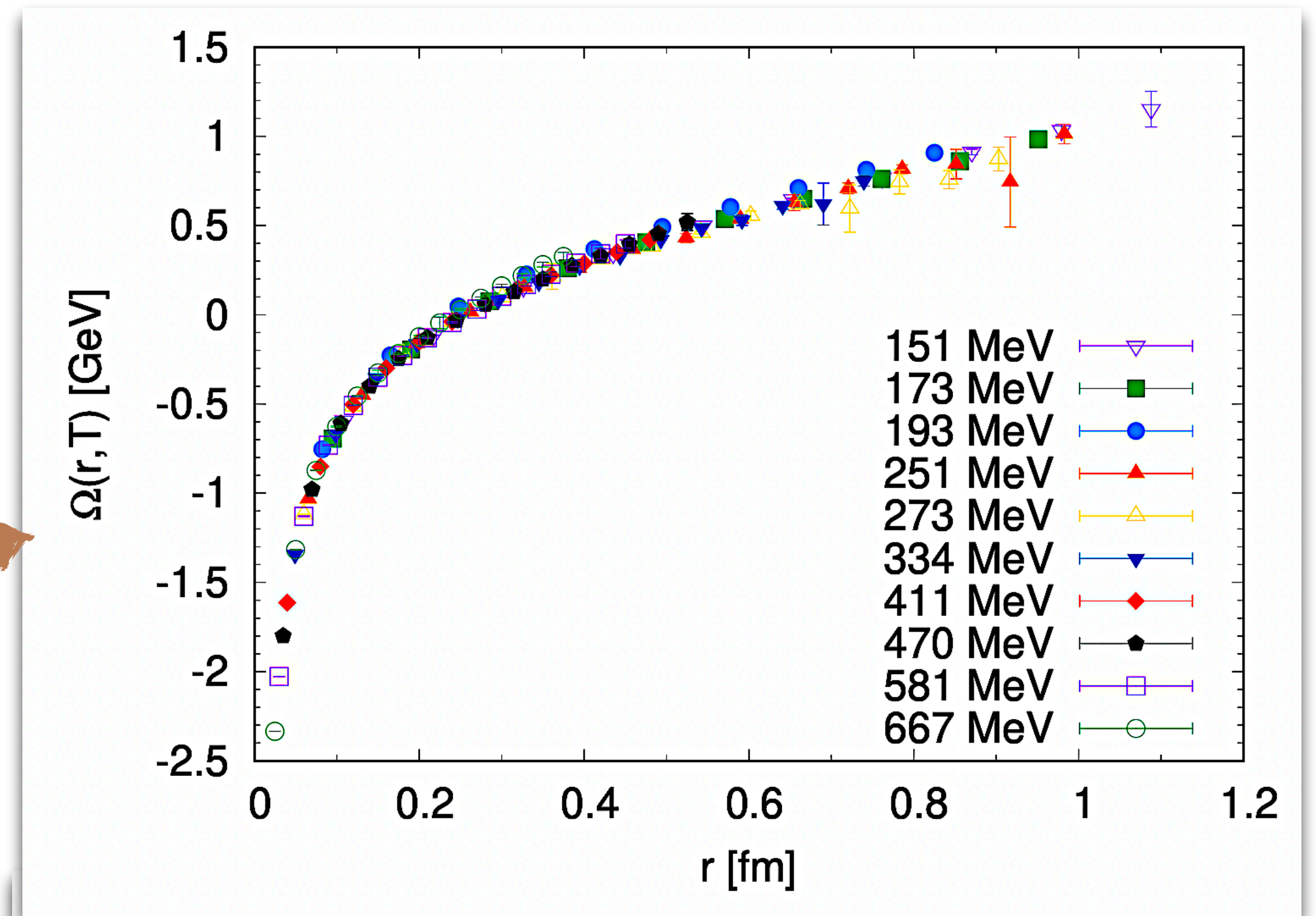
static quark potential

where is color screening ?

real part: almost no
temperature dependence

imaginary part: strong
temperature dependence

HotQCD, Phys. Rev. D 105, 054513 (2022)



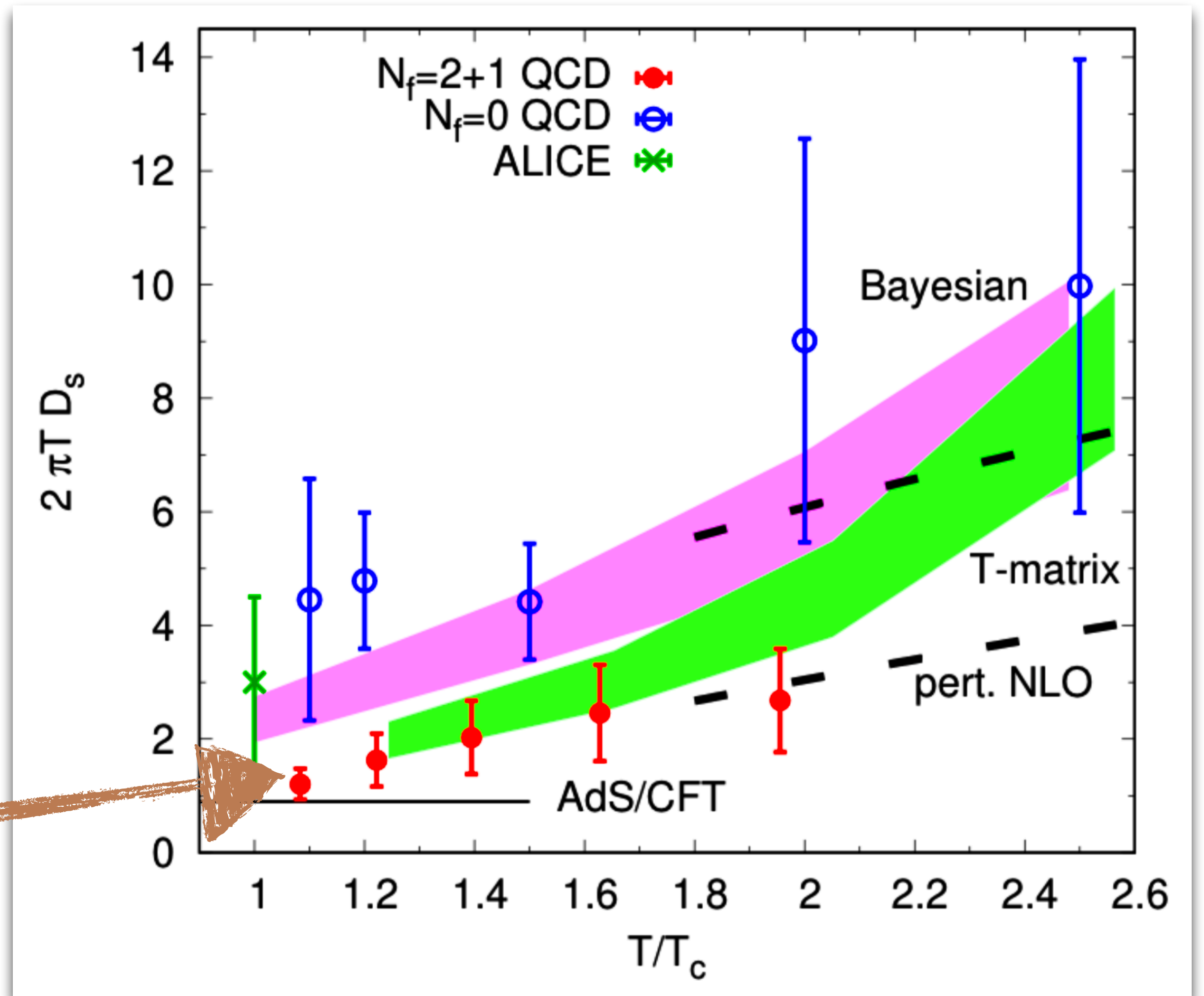
heavy quark diffusion

first lattice QCD calculations of transport coefficient including light dynamical fermions

HotQCD, 2302.08501

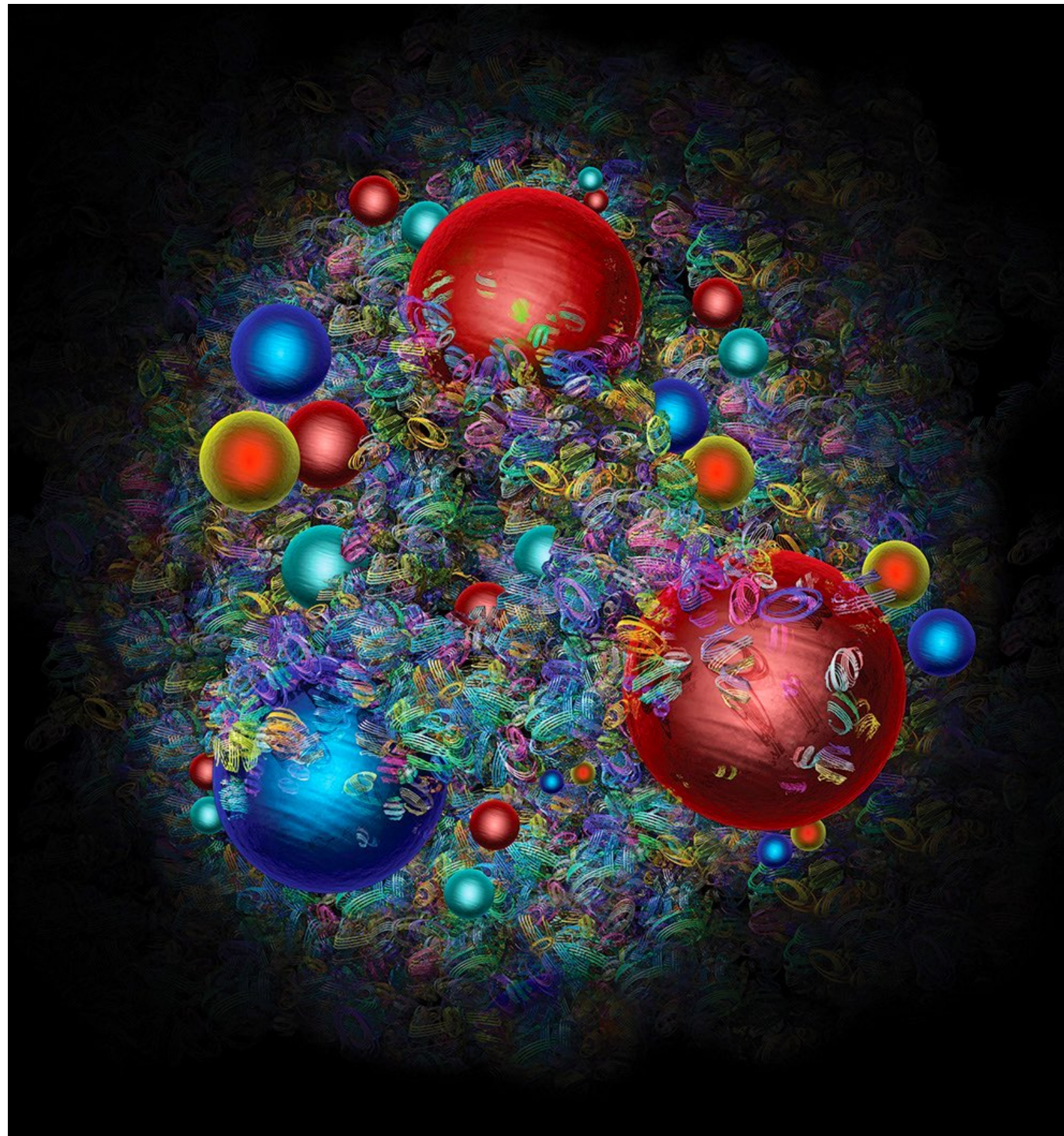
very fast equilibration of heavy quark in QGP

static quark spatial diffusion coefficient



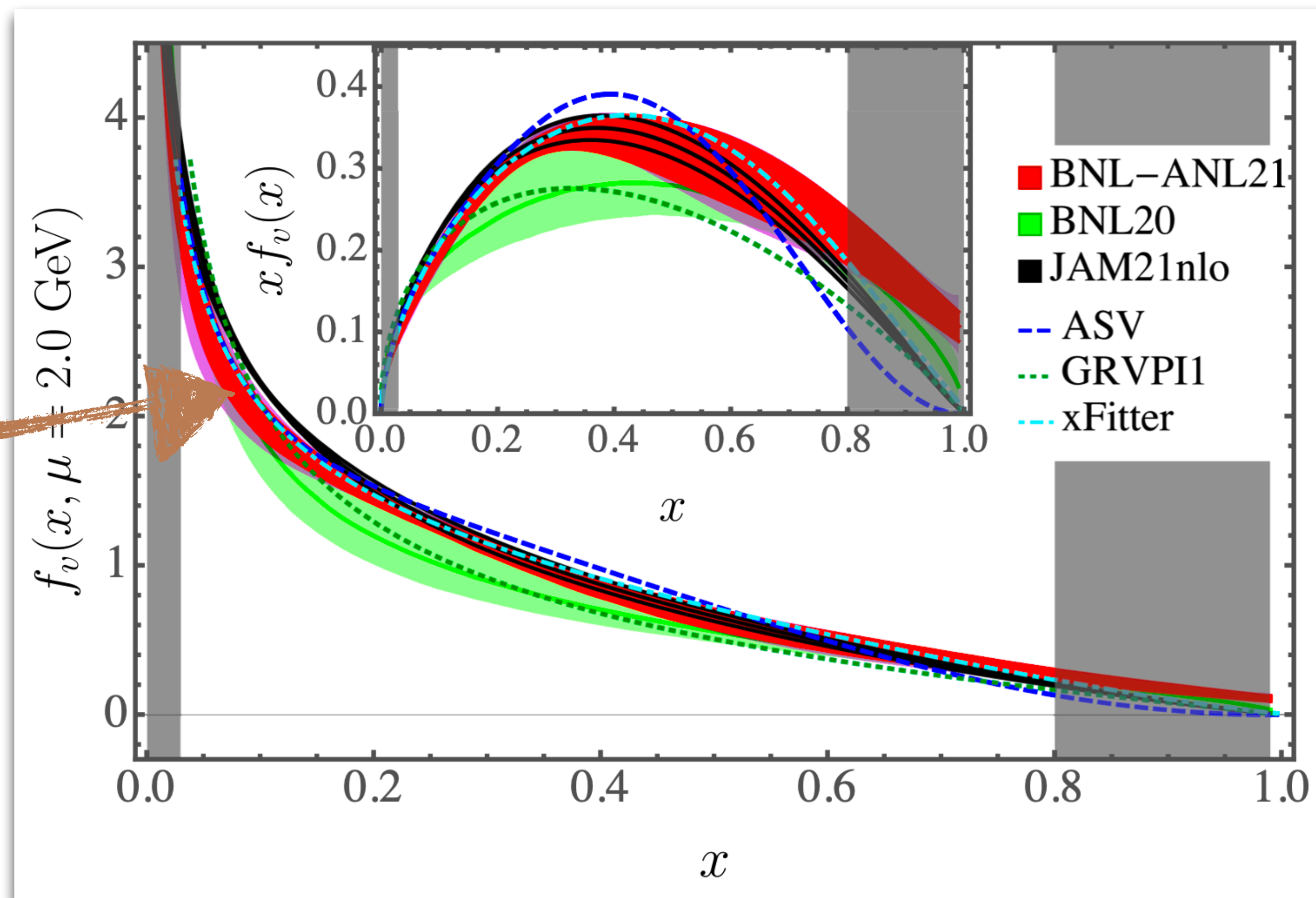
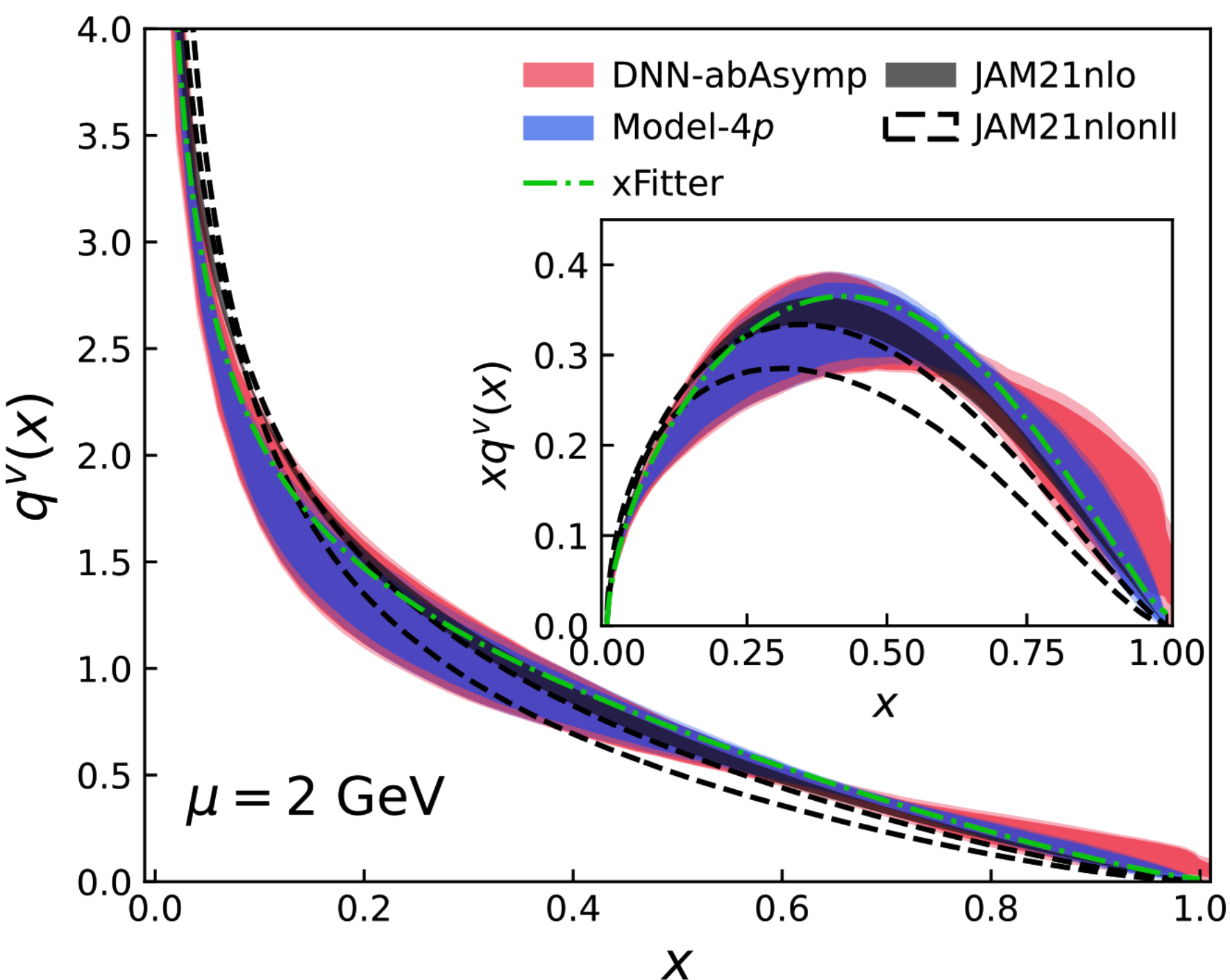
... 14 billion year later ...

... lost in hadronic heartland ...



peeking inside pion

first lattice QCD calculations of valance quark distribution of pion at NNLO

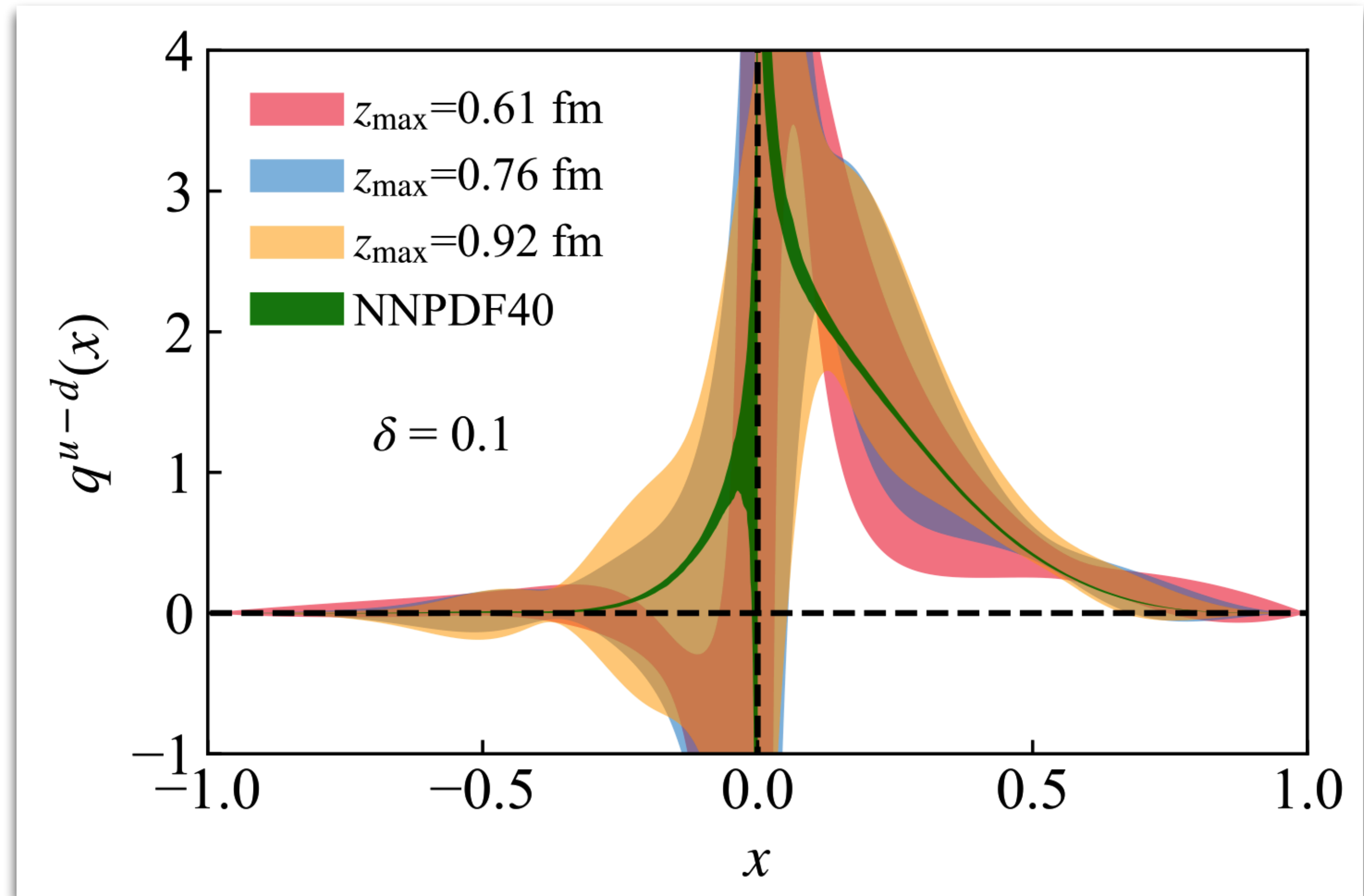


Gao *et al.*, Phys. Rev. Lett. 128, 14, 142003 (2022)

Gao *et al.*, Phys. Rev. D 106, 11, 114510 (2022)

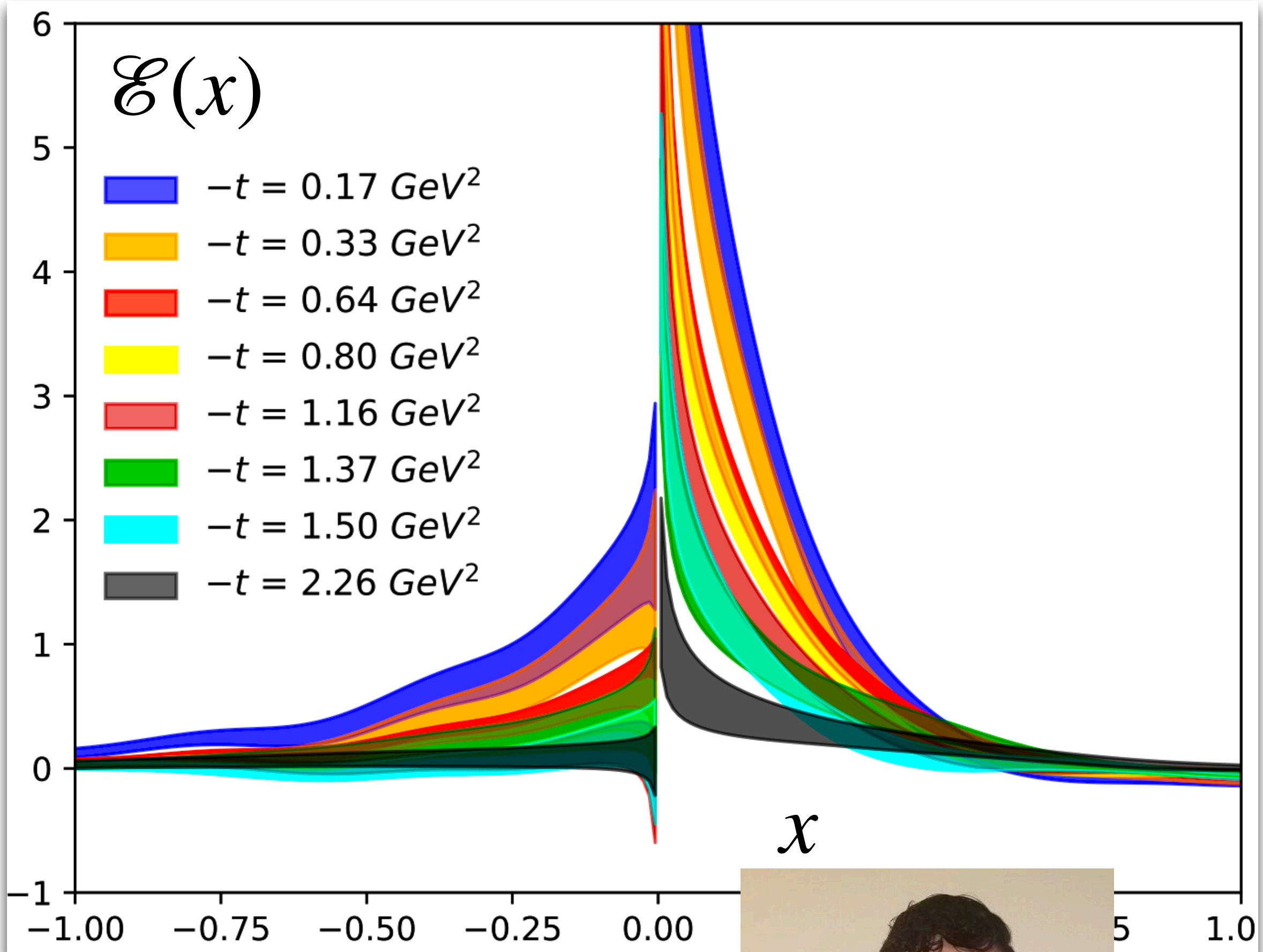
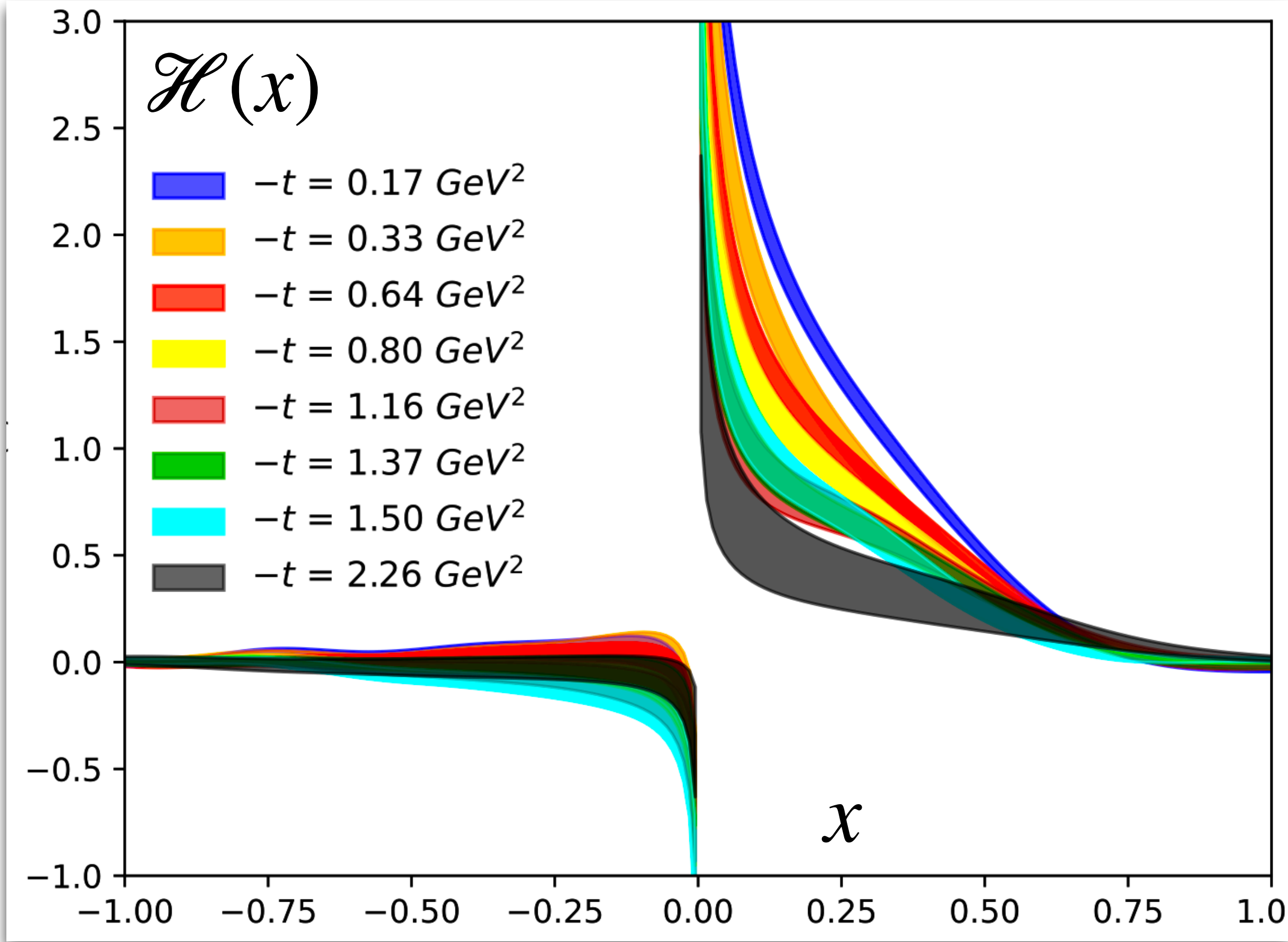
diving into proton

first lattice QCD calculations
with physical quark masses
of valance quark distribution
of proton at NNLO



Gao *et al.*, Phys. Rev. D, to be published (2212.12569)

2+1 dimensional image of proton ...

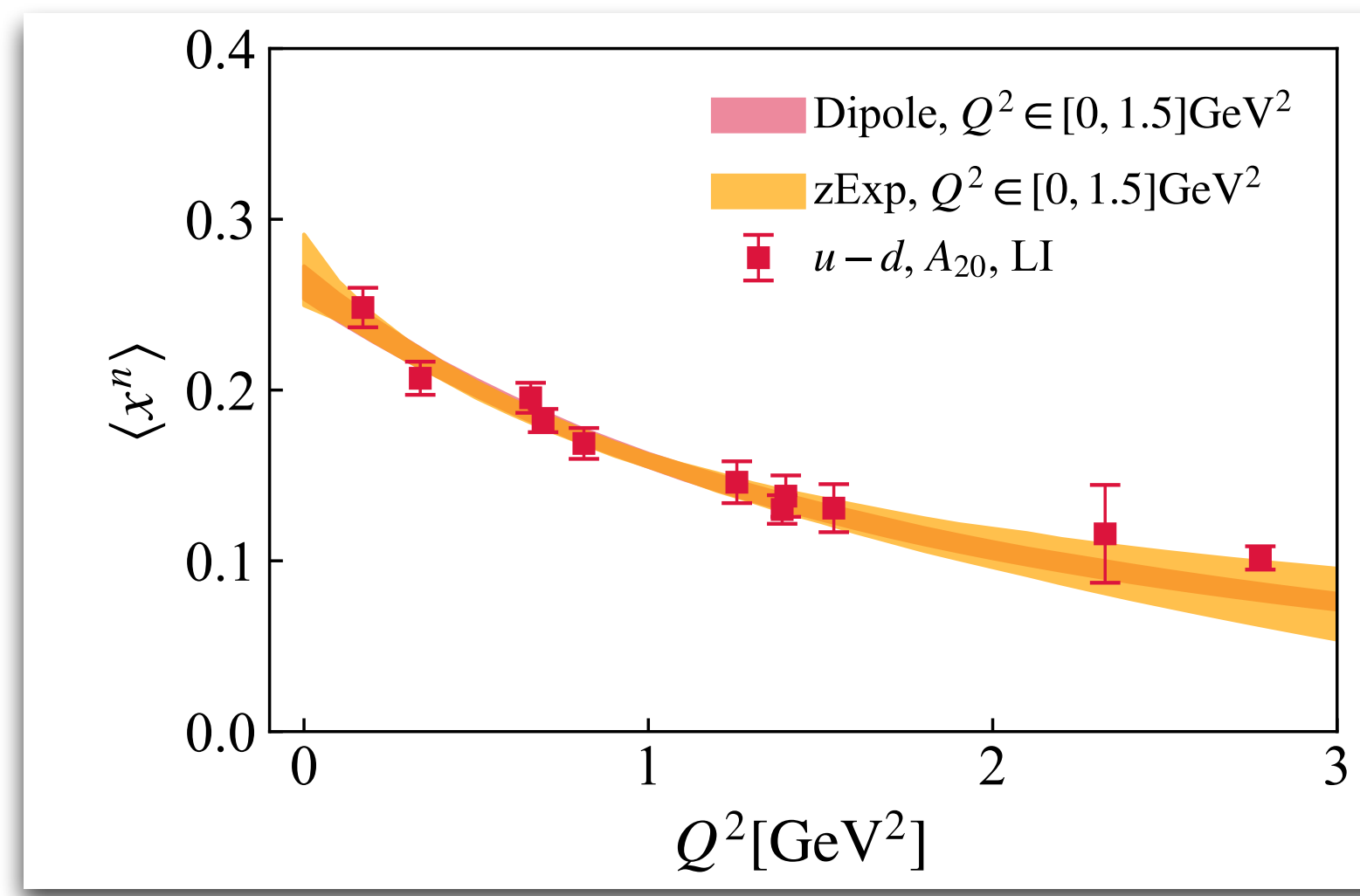
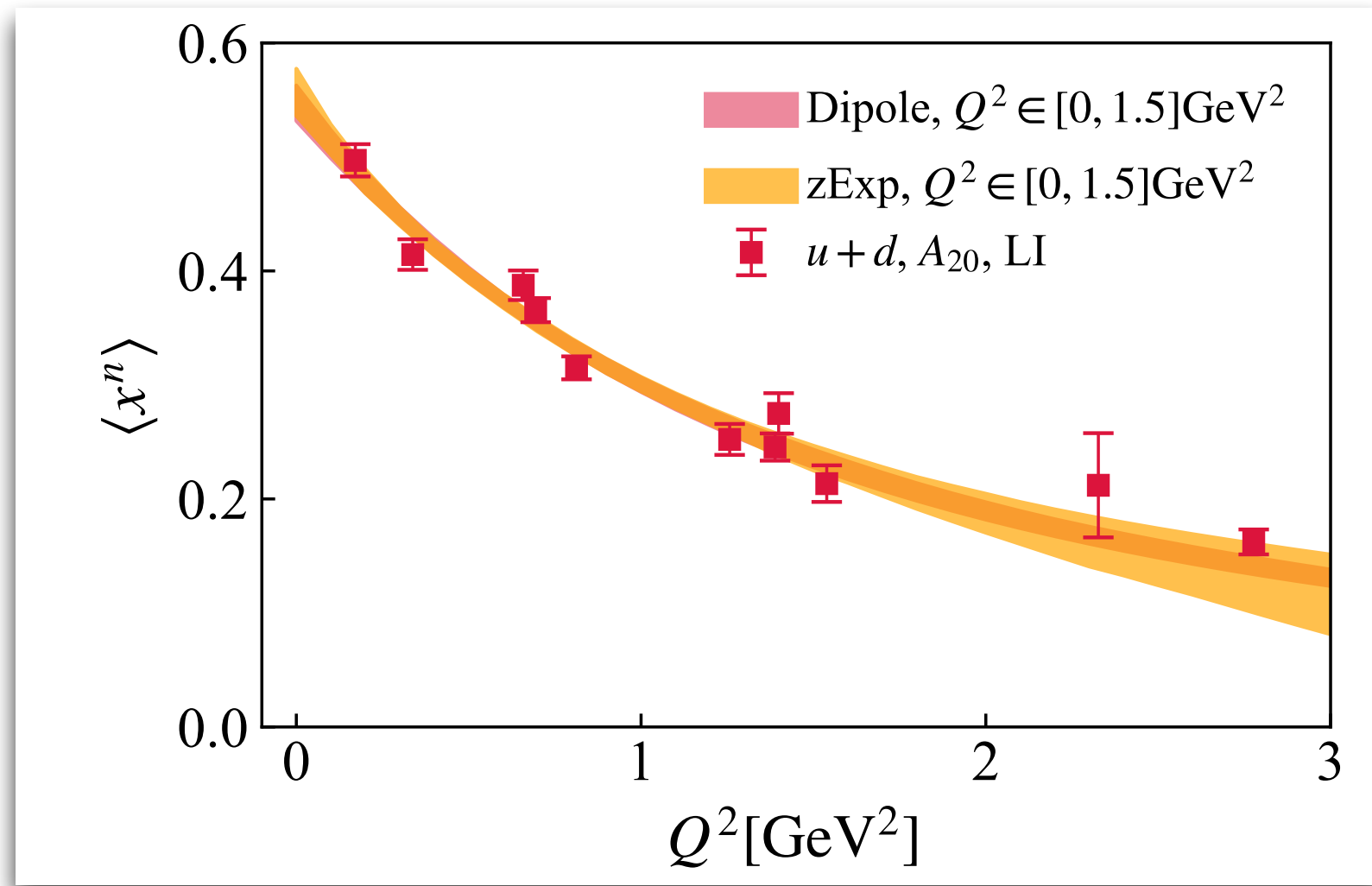


Bhattacharya *et al.*, Phys. Rev. D 106, 1, 114512 (2022)

talk: Joshua Miller



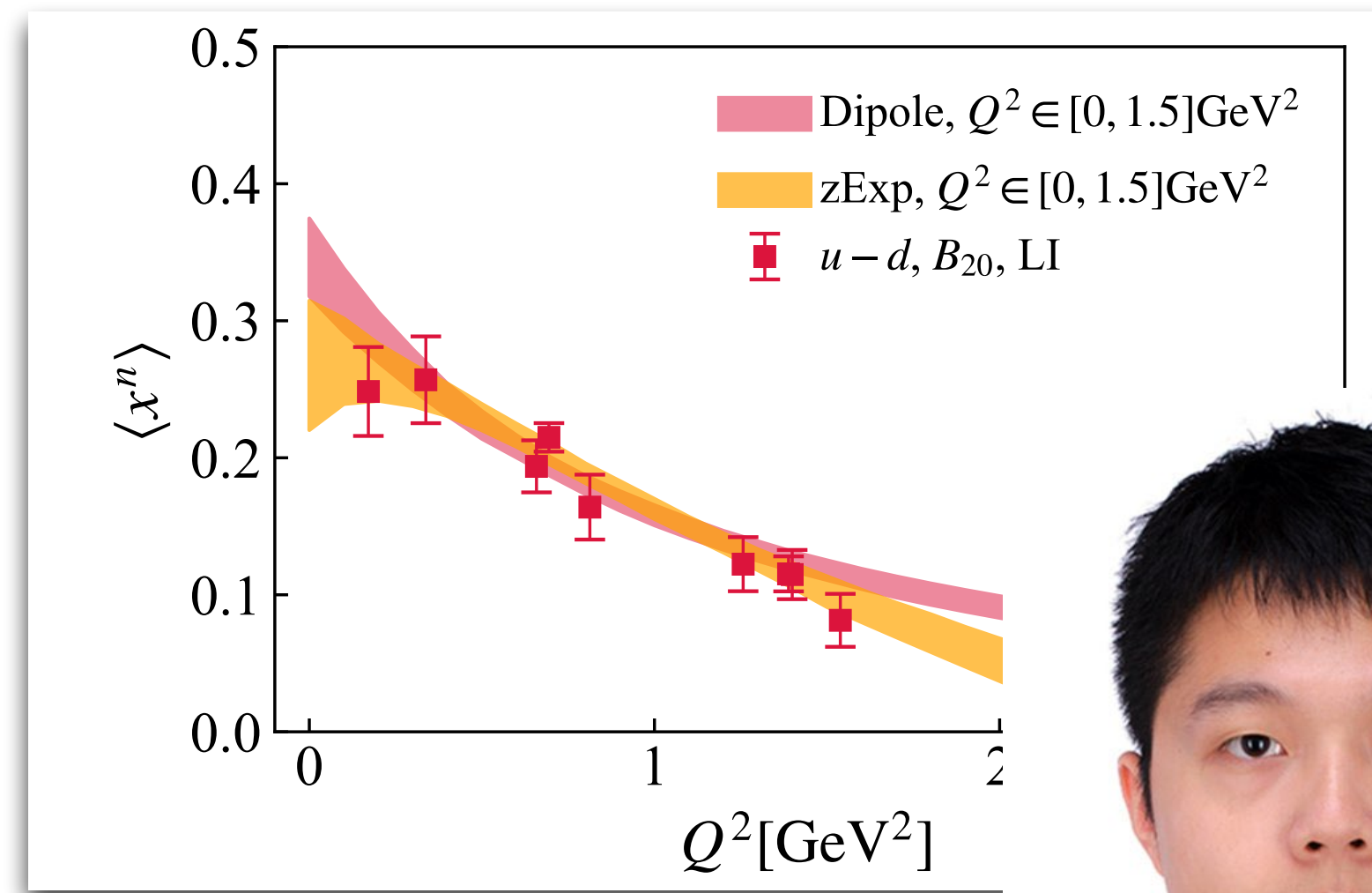
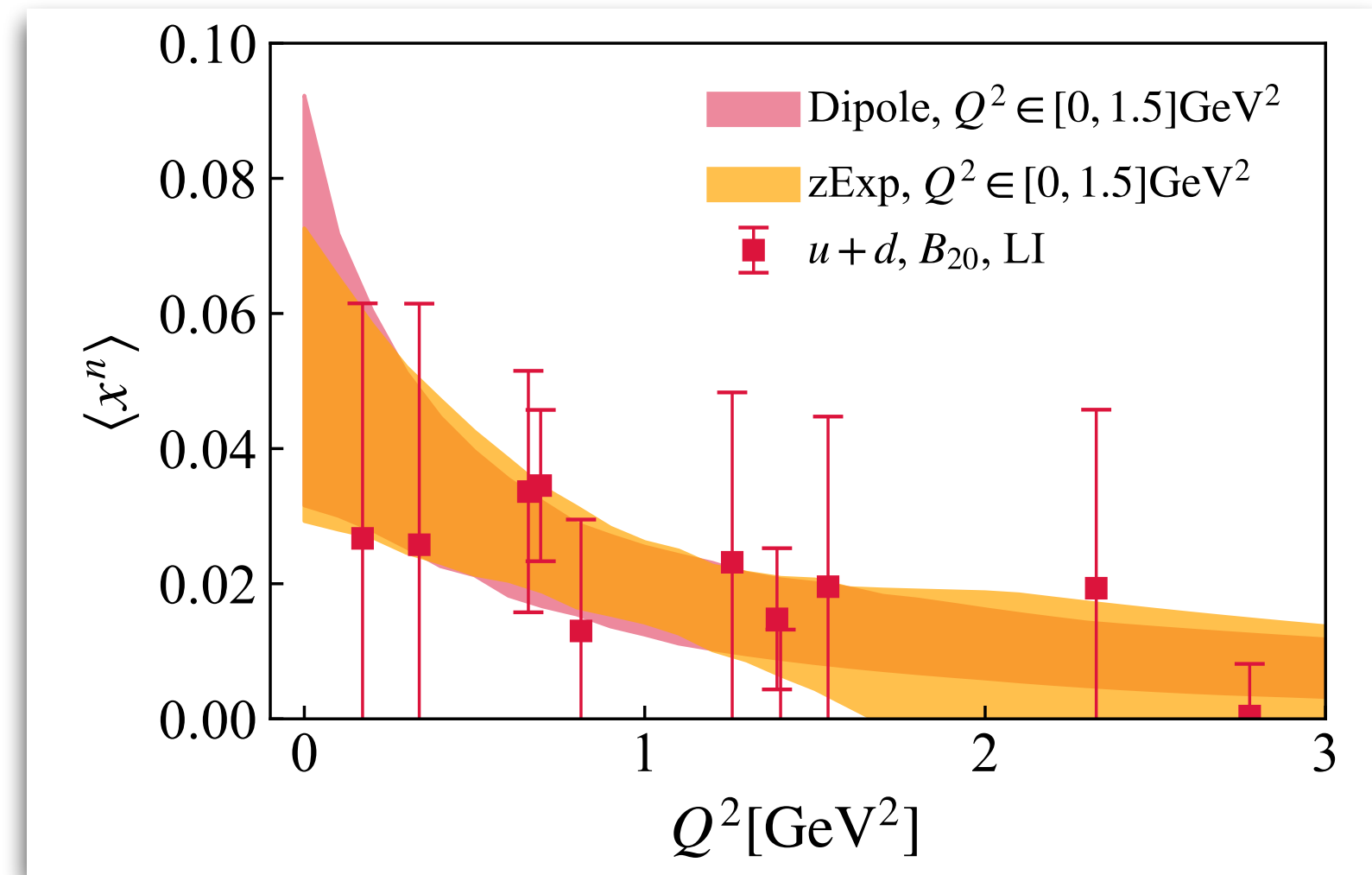
... from GPD to quark angular contribution to proton spin



$$J^{u-d} = 0.267(27)(39)$$

$$J^{u+d} = 0.301(14)(02)$$

$$J^q = \frac{1}{2}(A_{20}^q(0) + B_{20}^q(0))$$



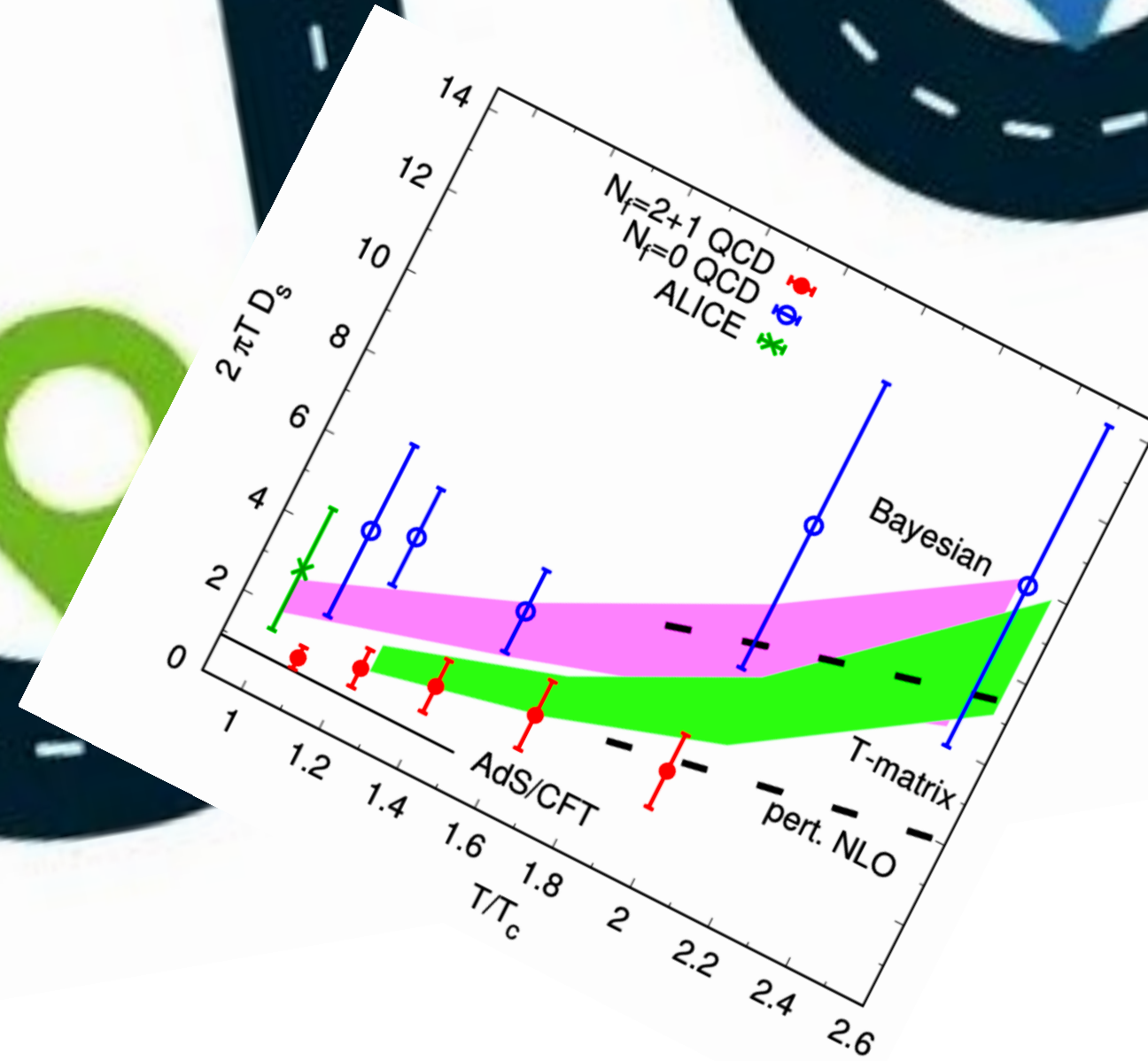
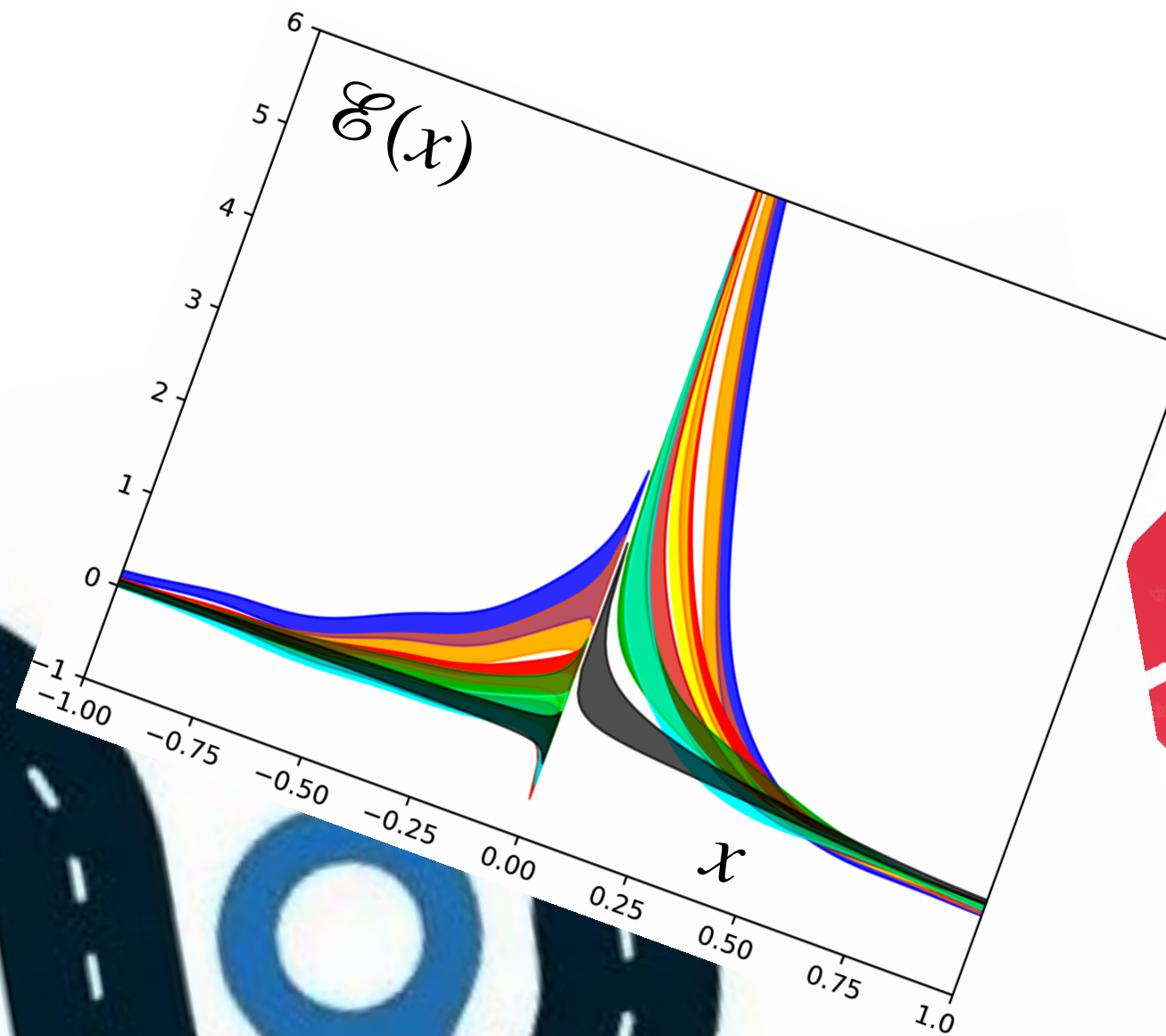
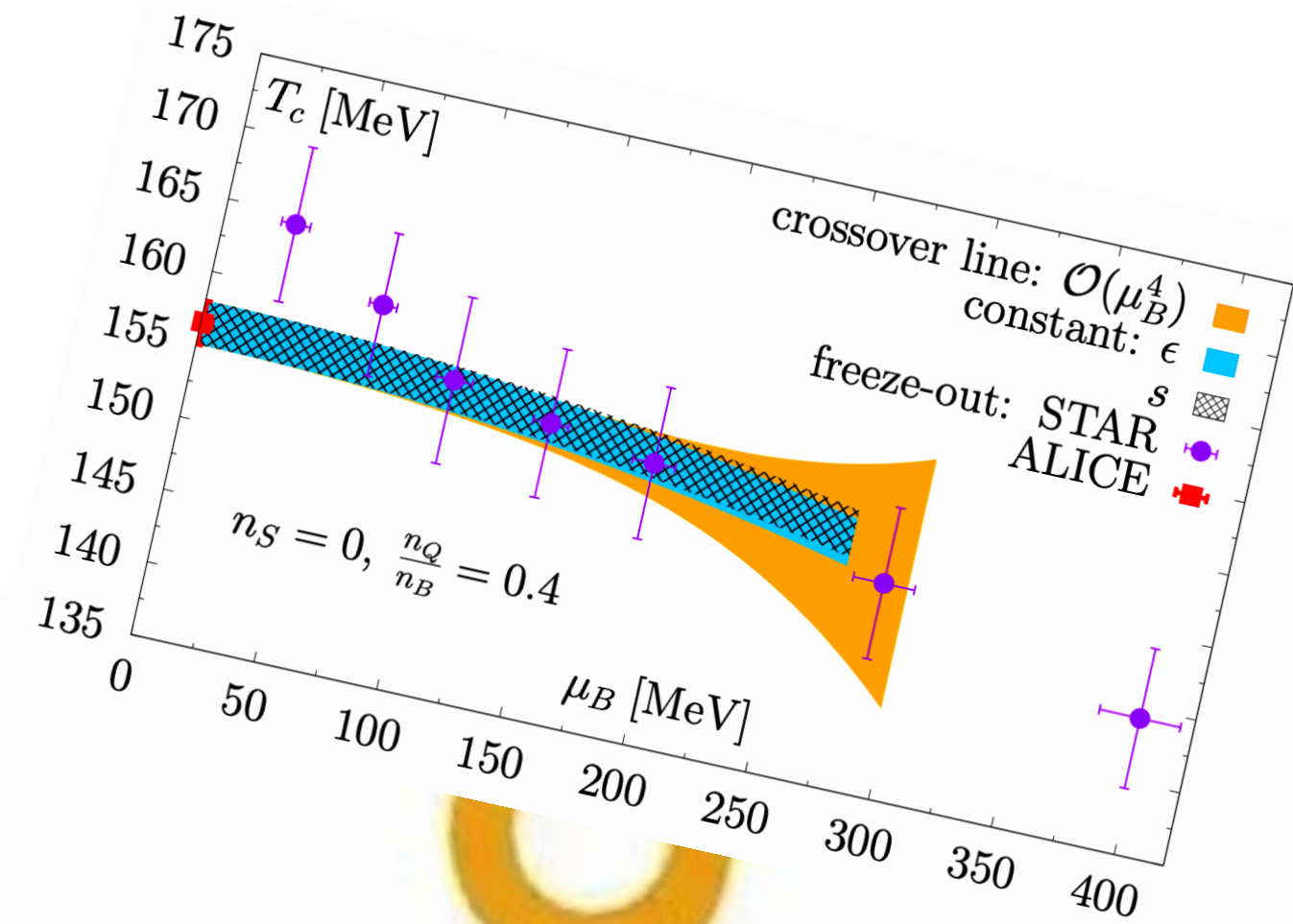
$$\int_{-1}^1 dx x^n H^q(x, \xi = 0, t) = A_{n+1,0}^q(t)$$

$$\int_{-1}^1 dx x^n E^q(x, \xi = 0, t) = B_{n+1,0}^q(t)$$



talk: Xiang Gao

at the end ...



... beginning of new journey

