Exclusive charmonium production and gluonic structure

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Near threshold J/ψ production based on gluonic GPDs



Figure: Photoproduction of J/ψ probing gluon GPDs.

• the differential cross section for J/ψ photoproduction in terms of gluon GPDs is given by [Guo, Ji, and Liu, 2021]

$$\frac{d\sigma}{dt} = \frac{\alpha_{\rm EM} e_Q^2}{4\left(W^2 - M_N^2\right)^2} \frac{(16\pi\alpha_S)^2}{3M_V^3} |\psi_{\rm NR}(\mathbf{0})|^2 |G(t,\xi)|^2$$

Near threshold J/ψ production based on gluonic GPDs

• where the function $G(t,\xi)$ is written in terms of gluon GPDs $F_g(x,\xi,t)$ as

$$G(t,\xi) = \frac{1}{2\xi} \int_{-1}^{1} \mathrm{d}x \mathcal{A}(x,\xi) \mathcal{F}_{g}(x,\xi,t)$$

where the hard kernel $\mathcal{A}(x,\xi)$ reads

$$\mathcal{A}(x,\xi) \equiv \frac{1}{x+\xi-i0} - \frac{1}{x-\xi+i0}$$

Near threshold J/ψ production based on gluonic GPDs

• after summing/averaging over the final and initial proton spin

$$|G(t,\xi)|^{2} = \frac{1}{\xi^{4}} \left\{ \left(1 - \frac{t}{4M_{N}^{2}} \right) E^{2} - 2E(H+E) + \left(1 - \xi^{2} \right) (H+E)^{2} \right\}$$

where we have defined

$$\int_0^1 dx H_g(x,\xi,t) = A(t) + \xi^2 D(t) \equiv H ,$$

$$\int_0^1 dx E_g(x,\xi,t) = B(t) - \xi^2 D(t) \equiv E$$

• Witten diagram showing the couplings of the graviton with bulk Dirac fermions and vector mesons [K.M., and Zahed, 2019]



Figure: Witten diagram.

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- the main elements of the Witten diagram are composed of:
- 1/ the bulk-to-boundary propagator of the vector mesons (or virtual photons for space-like momenta $q^2=-Q^2$) as

$$V(q,z) = \mathcal{V}(q=iQ,z) = g_5 \sum_n \frac{F_n \phi_n(z)}{Q^2 + m_n^2},$$

where $\phi_n(m_n, z)$, m_n , $f_n \equiv -F_n/m_n$, and $g_5 \sim \frac{1}{\sqrt{N_c}}$ are the bulk wave function, mass, decay constant, and hadronic coupling constant of each meson resonances, respectively

2/ the bulk-to-boundary propagator of the spin-2 glueballs (for space-like momenta $k^2=-{\cal K}^2$)

$$h(k,z) = \mathcal{H}(K,z) = \sqrt{2}\kappa \sum_{n} \frac{F_n \psi_n(z)}{K^2 + m_n^2},$$

where $\psi_n(m_n, z)$, m_n , $f_n \equiv -F_n/m_n$, and $\kappa \sim \frac{1}{N_c}$ are the bulk wave function, mass, decay constant, and hadronic coupling constant of each glueball resonances

3/ the bulk-to-bulk propagators of the vector meson and glueball resonances

$$G_V(q',z,z')=\sum_n \frac{\phi_n(z)\phi_n(z')}{q'^2-m_n^2},$$

and

$$G_h(k,z,z') = \sum_n \frac{\psi_n(z)\psi_n(z')}{k^2 - m_n^2}$$

• the normalized bulk wave function of one of the vector meson resonances $\phi_{n=0} \equiv \phi_V$ takes the form

$$\phi_V = c_V z J(M_V z) = \frac{f_V}{M_V} \times M_V z J(M_V z)$$

where $J(M_V z)$ is a special function that depends on the details of the holographic model • the decay constant f_V , for a meson at rest, defined as

$$< 0|J_{V,i}|V_j> = f_V M_V \delta_{ij}$$

is calculable in a given holographic model to QCD, and can be extracted experimentally from the leptonic width as

$$\Gamma(V \to \ell^+ \ell^-) = \frac{4\pi}{3} \alpha_{QED}^2 e_V^2 \frac{f_V^2}{M_V}$$

where e_V is the electric charge of the constituent quarks of the vector meson, and for $V = (J/\psi, \Upsilon)$: $e_V = (2/3, 1/3)$, $M_V = (3.097, 9.460)$ GeV and $e_V f_V = (270, 238)$ MeV

Kinematics of the $\gamma^* p \rightarrow V p$ Process

ullet the differential cross section for the photoproduction process $\gamma^* p o V p$ is given by

$$rac{d\sigma}{dt} = rac{e^2}{64\pi s |\mathbf{q}_{\gamma}|^2} |\mathcal{A}_{\gamma*p
ightarrow Vp}(s,t)|^2 \,,$$

where

$$|\mathsf{q}_{\gamma}| = rac{1}{2\sqrt{s}}\sqrt{s^2 - 2(-Q^2 + m_N^2)s + (-Q^2 - m_N^2)^2} \,,$$

with $s = W^2 = (p_1 + q_1)^2$, and $t = (p_1 - p_2)^2 = (q_1 - q_2)^2$ where $q_{1,2}$ are the four-vectors of the virtual photon and vector meson, respectively

J/ψ photoproduction near threshold: a probe to gluonic gravitational form factors of proton

 for heavy vector meson photoproduction the bulk-to-bulk propagators of glueball resonances factorizes as

$$G_h(k,z,z') = \sum_n \frac{\psi_n(z'\to 0)\psi_n(z)}{k^2 - m_n^2} = \frac{\psi_n(z'\to 0)}{\sqrt{2\kappa}F_n} \times \sum_n \frac{\sqrt{2\kappa}F_n\psi_n(z)}{k^2 - m_n^2}$$



Figure: Factorization of Witten diagram.

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The differential cross section for photoproduction of J/ψ near threshold

• after summing over the photon, J/ψ , and proton polarizations and spin, the differential cross section for photoproduction of J/ψ near threshold is [K.M., Zahed, 2019-2022]

$$egin{aligned} \left(rac{d\sigma}{dt}
ight) &= \mathcal{N}^2 imes \left[\mathcal{A}(t) + \xi^2 D(t)
ight]^2 \ & imes rac{1}{\mathcal{A}^2(0)} imes rac{1}{128\pi(s-m_N^2)^2} imes \mathcal{F}(s,t,M_V,m_N) imes \left(1-rac{t}{4m_N^2}
ight) \end{aligned}$$

 \bullet the normalization factor ${\cal N}$ is defined as

$$\mathcal{N}^2 \equiv e^2 imes \left(rac{f_V}{M_V}
ight)^2 imes \mathbb{V}^2_{hAA} imes 2\kappa^2 imes A^2(0) = 7.768 \, \mathrm{GeV}^{-4}$$

• we also have $F(s,t)\sim s^4\sim 1/\xi^4$ with $rac{d\sigma}{dt}\sim s^2 imes \left({\it A}({\it K})+{\it D}({\it K})/s^2
ight)$

Pentaquarks at LHCb

• The masses and total widths of the three charm pentaquark states recently reported by LHCb [Aaij:2019] are

$$\begin{array}{rcl} m_{P_c} &=& 4311.9 \pm 0.7 \; {\rm MeV} & \Gamma_{P_c} = 9.8 \pm 2.7 \; {\rm MeV} \\ m_{P_c} &=& 4440.3 \pm 1.3 \; {\rm MeV} & \Gamma_{P_c} = 20.6 \pm 4.9 \; {\rm MeV} \\ m_{P_c} &=& 4457.3 \pm 0.6 \; {\rm MeV} & \Gamma_{P_c} = 6.4 \pm 2.0 \; {\rm MeV} \end{array}$$

• the bulk interaction vertex between the Dirac fermion ψ and the pentaquark field Ψ is [Liu, K.M., Nowak, and, Zahed, 2021]

$$\eta_X \int dz d^4 x \sqrt{|g(z)|} e^{-\phi_N(z)} \sum_{\xi=1,2} \bar{\Psi}_{\xi} e^M_A e^N_B \sigma^{AB} F_{MN} \psi_{\xi} ,$$



Figure: The s- and t-channel contribution to near threshold photoproduction of J/ψ .

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Figure: s-channel contribution to the photo-production cross section for $V = J/\psi$ versus \sqrt{s} , showing the three charm pentaquarks.



Figure: Total cross section for $V = J/\psi$ photo-production: the blue-solid curve is the t-channel contribution from [Mamo:2019], the red-solid curve is the sum of t- and s-channel contribution showing the three holographic pentaquarks times $N_s = 2.0 \times 10^6$ to make them visible, and the data are from GlueX [GlueX:2019].

The differential cross section for photoproduction of J/ψ near threshold: comparison to experiment



Figure: The $A(k^2)$ form factor extracted from the two-dimensional cross section data of $J/\psi - 007$ collaboration (Duran, et al., 2022) using our holographic QCD approach (orange dash-dot curve) and the GPD+VMD approach [Guo, et al., 2021] (green solid curve), compared to the latest lattice calculation [Pefkou, et al., 2021] (blue dotted curve).

The differential cross section for photoproduction of J/ψ near threshold: comparison to experiment



Figure: The $D(k^2) = 4C(k^2)$ form factor extracted from the two-dimensional cross section data of $J/\psi - 007$ collaboration (Duran, et al., 2022) using our holographic QCD approach (orange dash-dot curve) and the GPD+VMD approach [Guo, et al., 2021] (green solid curve), compared to the latest lattice calculation [Pefkou, et al., 2021] (blue dotted curve).

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The differential cross section for photoproduction of J/ψ near threshold: comparison to experiment

Table: Gluonic GFF parameters and the corresponding proton mass and scalar radii, determined by $J/\psi - 007$ collaboration through a two-dimensional fit following the holographic QCD and GPD+VMD approach, compared to the latest lattice results.

Гheoretica∣ approach	$\chi^2/{\rm n.d.f}$	m_A (GeV ²)	m_C (GeV ²)	$C_g(0)$	$\sqrt{\langle r_m^2 angle}$ (fm)	$\sqrt{\langle r_s^2 angle}$ (fm)
Holographic QCD	0.925	$1.575 {\pm} 0.059$	$1.12{\pm}0.21$	-0.45 ± 0.132	$0.755{\pm}0.035$	$1.069{\pm}0.056$
GPD + VMD	0.924	2.71±0.19	$1.28\pm$ 0.50	$\textbf{-0.20}\pm0.11$	0.472±0.042	$0.695{\pm}0.071$
Lattice		$1.641{\pm}~0.043$	1.07 ± 0.12	-0.483± 0.133	0.7464±0.025	$1.073 {\pm} 0.066$

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Photoproduction of J/ψ from near threshold to far from threshold



Figure: The total cross section for J/ψ photoproduction.

Thank You!