

Hadronic parity violation in few-nucleon systems

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Beyond the Standard Model Physics at low energies

- New physics at scale Λ_{new}
- For $E \ll \Lambda_{\text{new}}$: effective interactions with SM degrees of freedom
- Possible structure: local 4-quark operators $\bar{q}O_1q \bar{q}O_2q$
- In nuclei: isolate through symmetry violations
- QCD nonperturbative: manifestation of 4-quark operators at hadronic level?

Parity-violating NN interactions

- Hadronic manifestation of PV quark interactions
- Interplay of weak and nonperturbative strong interactions
- Range of weak interactions ~ 0.002 fm
 - Sensitive to quark-quark correlations inside nucleon
 - “Inside-out” probe: no need to go to high energies
- Relative strength in NN case $\sim G_F m_\pi^2 \approx 10^{-7}$
- Isolate through pseudoscalar observables $(\vec{p} \cdot \vec{\sigma})$

Parity-violating NN interactions

- Meson-exchange model
 - Single-meson exchange (π, ρ, ω, \dots)
 - One parity-conserving, one parity-violating meson-nucleon coupling
 - Most common framework: DDH model
- Effective field theories
 - Pionless: 5 PV low-energy constants (LECs) at LO
 - Chiral: 1 PV πN LEC at LO, 5 PV NN LECs at NLO
- Experimental constraints very difficult to obtain

Parity-violating NN interactions in pionless EFT

- At very low energies: pion exchange not resolved
- Only nucleons as explicit degrees of freedom
- Contact terms with increasing number of derivatives
- Parity determined by orbital angular momentum L : $(-1)^L$
- Simplest PV interaction: $L \rightarrow L \pm 1$
- Leading order: five S - P transition operators with corresponding LECs

Parity-violating NN interactions in pionless EFT

- Five independent LECs at LO

$$C^{(3S_1-1P_1)}, \quad C_{(\Delta I=0)}^{(1S_0-3P_0)}, \quad C_{(\Delta I=1)}^{(1S_0-3P_0)}, \quad C_{(\Delta I=2)}^{(1S_0-3P_0)}, \quad C^{(3S_1-3P_1)}$$

- Parametrize short-distance details
- Determine from
 - Underlying theory → nonperturbative QCD
 - Experimental results → suite of high-precision measurements in few-nucleon systems at very low energies
- Additional theoretical constraints?

Large- N_c QCD

- Generalize QCD from $SU(3)$ to $SU(N_c)$ gauge theory

- Taken with $g^2 N_c$ fixed

- Systematic expansion in $1/N_c$

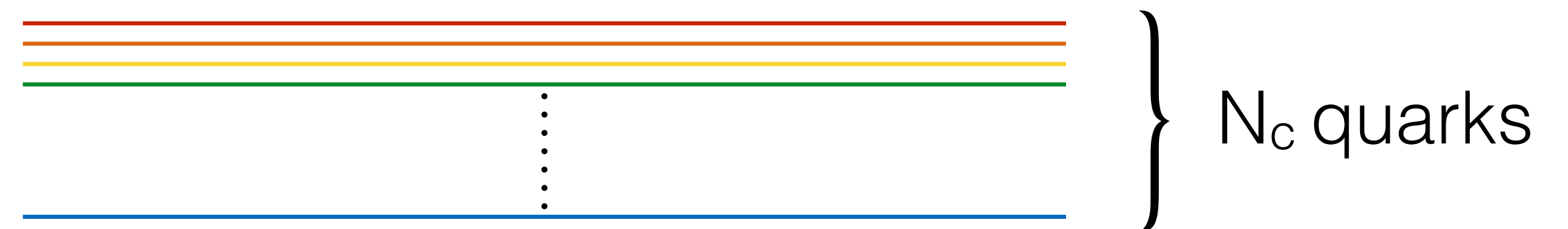
- Phenomenologically successful

- Baryons

 - Bound states of N_c quarks

 - Baryon mass $M \sim O(N_c)$

 - Spectrum: (contracted) $SU(4)$ spin-flavor symmetry $u \uparrow, u \downarrow, d \uparrow, d \downarrow$



Large- N_c QCD and effective field theories

- Capture some nonperturbative QCD aspects
- Based on symmetries
- Expansions in small parameter ($1/N_c$ vs p/Λ)
- Complementary
- Individually successful

Combine to obtain double expansion

- Large- N_c scaling of low-energy coefficients of EFT

Caveats

- Nuclear matter forms classical crystal for $N_c \rightarrow \infty$?
 - Assume that symmetries of NN interactions do not change
- Nucleon and Δ degenerate in large- N_c limit
- Δ plays important role in meson-baryon interactions
 - Ignore intermediate Δ states

PV NN interactions in the large- N_c expansion

- General operators structure

▸ LO in N_c [$\mathcal{O}(N_c)$]

$$\mathbf{p}_- \cdot (\vec{\sigma}_1 \times \vec{\sigma}_2) \vec{\tau}_1 \cdot \vec{\tau}_2$$

$$\mathbf{p}_- \cdot (\vec{\sigma}_1 \times \vec{\sigma}_2) [\tau_1 \tau_2]_2^{zz}$$

▸ NLO in N_c [$\mathcal{O}(N_c^0)$]

$$\mathbf{p}_+ \cdot (\vec{\sigma}_1 \tau_1^3 - \vec{\sigma}_2 \tau_2^3)$$

$$\mathbf{p}_- \cdot (\vec{\sigma}_1 + \vec{\sigma}_2) (\vec{\tau}_1 \times \vec{\tau}_2)^3$$

$$\mathbf{p}_- \cdot (\vec{\sigma}_1 \times \vec{\sigma}_2) (\vec{\tau}_1 + \vec{\tau}_2)^3$$

$$[(\mathbf{p}_+ \times \mathbf{p}_-) \cdot \vec{\sigma}_1 \mathbf{p}_- \cdot \vec{\sigma}_2 + (\mathbf{p}_+ \times \mathbf{p}_-) \cdot \vec{\sigma}_2 \mathbf{p}_- \cdot \vec{\sigma}_1] (\vec{\tau}_1 \times \vec{\tau}_2)^3$$

PV NN interactions in the large- N_c expansion

- ▶ NNLO in N_c [$\mathcal{O}(N_c)$]

$$\mathbf{p}_- \cdot (\vec{\sigma}_1 \times \vec{\sigma}_2)$$

$$\mathbf{p}_+^2 \mathbf{p}_- \cdot (\vec{\sigma}_1 \times \vec{\sigma}_2) \vec{\tau}_1 \cdot \vec{\tau}_2$$

$$\mathbf{p}_+ \cdot (\vec{\sigma}_1 - \vec{\sigma}_2)$$

$$\mathbf{p}_+ \cdot (\vec{\sigma}_1 - \vec{\sigma}_2) \vec{\tau}_1 \cdot \vec{\tau}_2$$

$$\mathbf{p}_+ \cdot (\vec{\sigma}_1 - \vec{\sigma}_2) [\tau_1 \tau_2]_2^{zz}$$

$$\mathbf{p}_+^2 \mathbf{p}_- \cdot (\vec{\sigma}_1 \times \vec{\sigma}_2) [\tau_1 \tau_2]_2^{zz}$$

- Can be multiplied by independent functions $U_i(\vec{p}_-^2) \sim \mathcal{O}(1)$

Parity violation in pionless EFT + large N_c

- Large- N_c scaling of PV couplings

$$\begin{aligned} \mathcal{C}^{(^3S_1-^1P_1)} &\sim N_c & \mathcal{C}_{(\Delta I=0)}^{(^1S_0-^3P_0)} &\sim N_c & \mathcal{C}_{(\Delta I=2)}^{(^1S_0-^3P_0)} &\sim N_c \sin^2 \theta_W \\ \mathcal{C}_{(\Delta I=1)}^{(^1S_0-^3P_0)} &\sim N_c^0 \sin^2 \theta_W & \mathcal{C}^{(^3S_1-^3P_1)} &\sim N_c^0 \sin^2 \theta_W \end{aligned}$$

- Different terms related by Fierz transformations

$$\mathcal{C}^{(^3S_1-^1P_1)} = 3 \mathcal{C}_{(\Delta I=0)}^{(^1S_0-^3P_0)} [1 + O(1/N_c^2)]$$

- Only **two independent** terms at leading order in $1/N_c$ expansion
- Isotensor term is LO in combined expansion
- Isovector term (pion exchange) is NLO



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A New Paradigm for Hadronic Parity Nonconservation and Its Experimental Implications

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Comparison with experiment

- Longitudinal asymmetry in $\vec{p}p$ scattering and induced polarization in $np \rightarrow d\vec{\gamma}$:

$$\mathcal{C}^{(^3S_1-^1P_1)} / \mathcal{C} = (-1.1 \pm 1.0) \times 10^{-10} \text{ MeV}^{-1}$$

$$\mathcal{C}_{(\Delta I=2)}^{(^1S_0-^3P_0)} / \mathcal{C} = (7.4 \pm 2.3) \times 10^{-11} \text{ MeV}^{-1}$$

- γ -ray asymmetry in $\vec{n}p \rightarrow d\gamma$:

$$\mathcal{C}^{(^3S_1-^3P_1)} / \mathcal{C} = (-4.6 \pm 2.1) \times 10^{-11} \text{ MeV}^{-1}$$

- Large errors
- Not inconsistent with large- N_c expectation

Conclusions

- Parity violation in pionless EFT and large N_c
 - Two LECs at LO in combined pionless EFT and large- N_c expansion
 - Relationship between isoscalar LECs
 - Isovector LEC only NLO in large- N_c expansion, but isotensor LEC is LO
 - Trends, not predictions
 - Only upper limits on size
 - Other scales can impact relative sizes
 - LECs strongly depend on renormalization point
 - Renormalization-point dependence driven by S -wave interaction