

Nuclear imaging with diffractive vector meson production in UPCs

SPIN Symposium

Joint GPD/Nuclear Session
Sep 26th, 2023

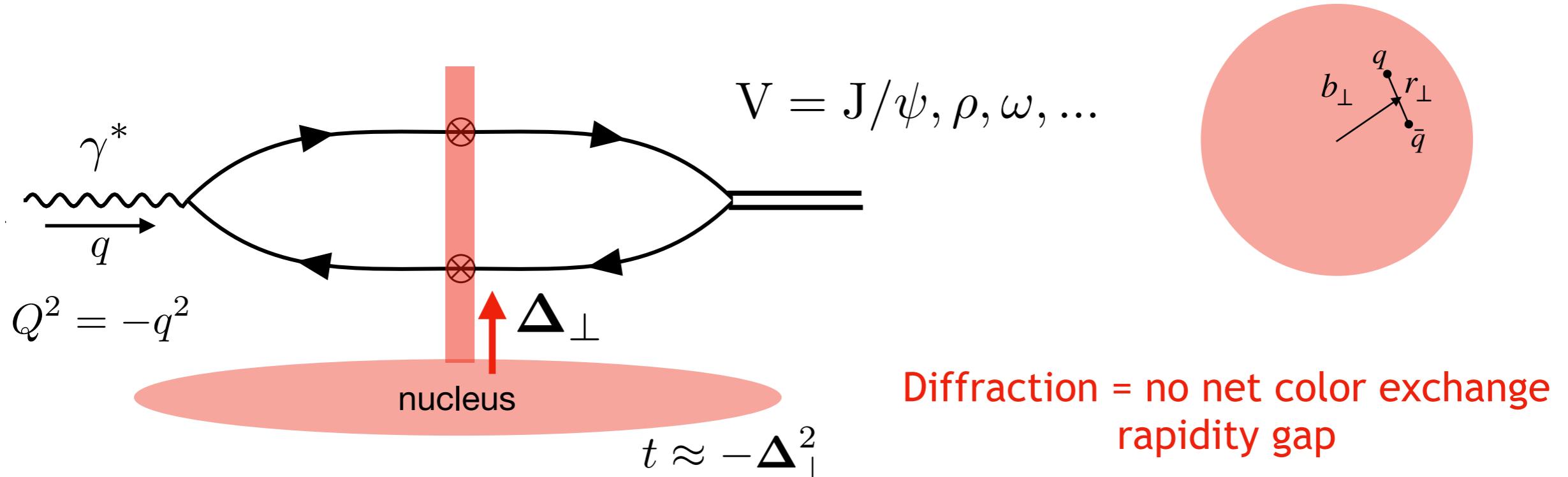


Farid Salazar

*Phys.Rev.D 106 (2022) 7, 074019. arXiv: 2207.03712. H. Mäntysaari, FS, B. Schenke.
Work in progress 2310.XXXX. HM, FS, BS, Chun Shen, Wenbin Zhao*

Diffraction, Saturation and Vector Meson Production

The dipole picture



Ψ_{γ^*} computed in QED, Ψ_V needs a non-perturbative modeling

$$\mathcal{A}(Q, \Delta_{\perp}) = \int_{\mathbf{b}_{\perp}, \mathbf{r}_{\perp}, z} e^{-i\Delta_{\perp} \cdot \mathbf{b}_{\perp}} \Psi_{\gamma}^*(Q, \mathbf{r}_{\perp}, z) D(\mathbf{r}_{\perp}, \mathbf{b}_{\perp}) \Psi_V(\mathbf{r}_{\perp}, z)$$

Interaction of the dipole with small- x background field, aka CGC, shock-wave...

Diffraction, Saturation and Vector Meson Production

Coherent and nuclear imaging

- Coherent (aka exclusive): Target remains intact

$$d\sigma_{coh}/d|t| \propto \langle \mathcal{A}^\dagger(\Delta_\perp) \rangle_x \langle \mathcal{A}(\Delta_\perp) \rangle_x$$

Simplest model

$$\langle D(\mathbf{r}_\perp, \mathbf{b}_\perp) \rangle_x \propto T(\mathbf{b}_\perp) \quad d\sigma_{coh} \propto \left| \int d^2 \mathbf{b}_\perp e^{-i \Delta_\perp \cdot \mathbf{b}_\perp} T(\mathbf{b}_\perp) \right|^2$$

Study Δ_\perp -dependence -> form factor of nuclei (Nuclear imaging)

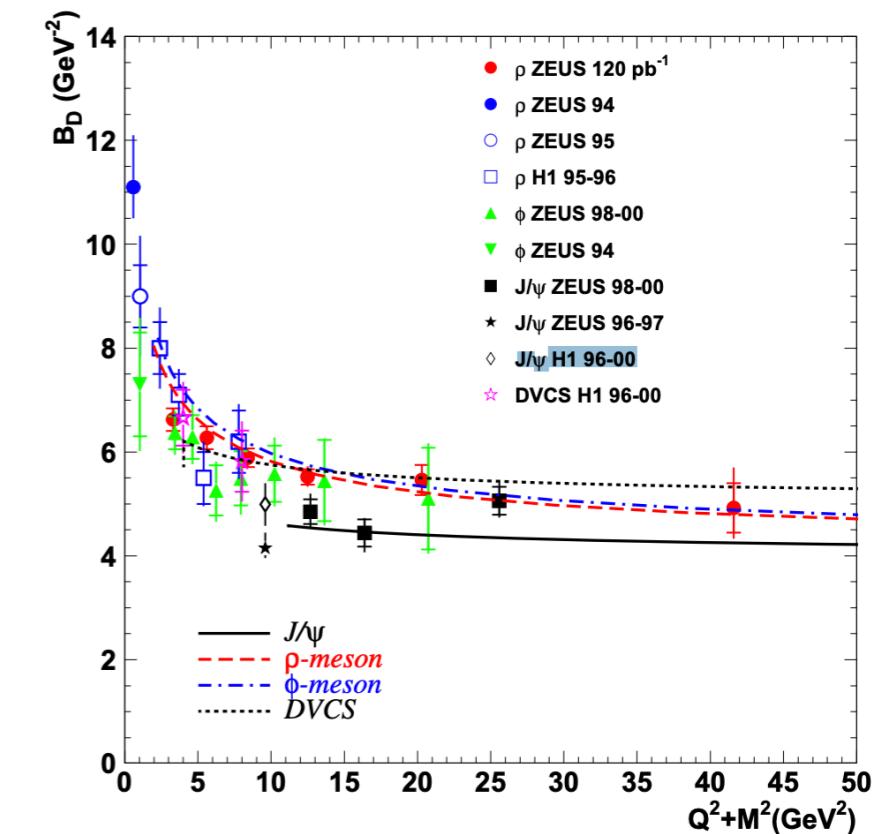
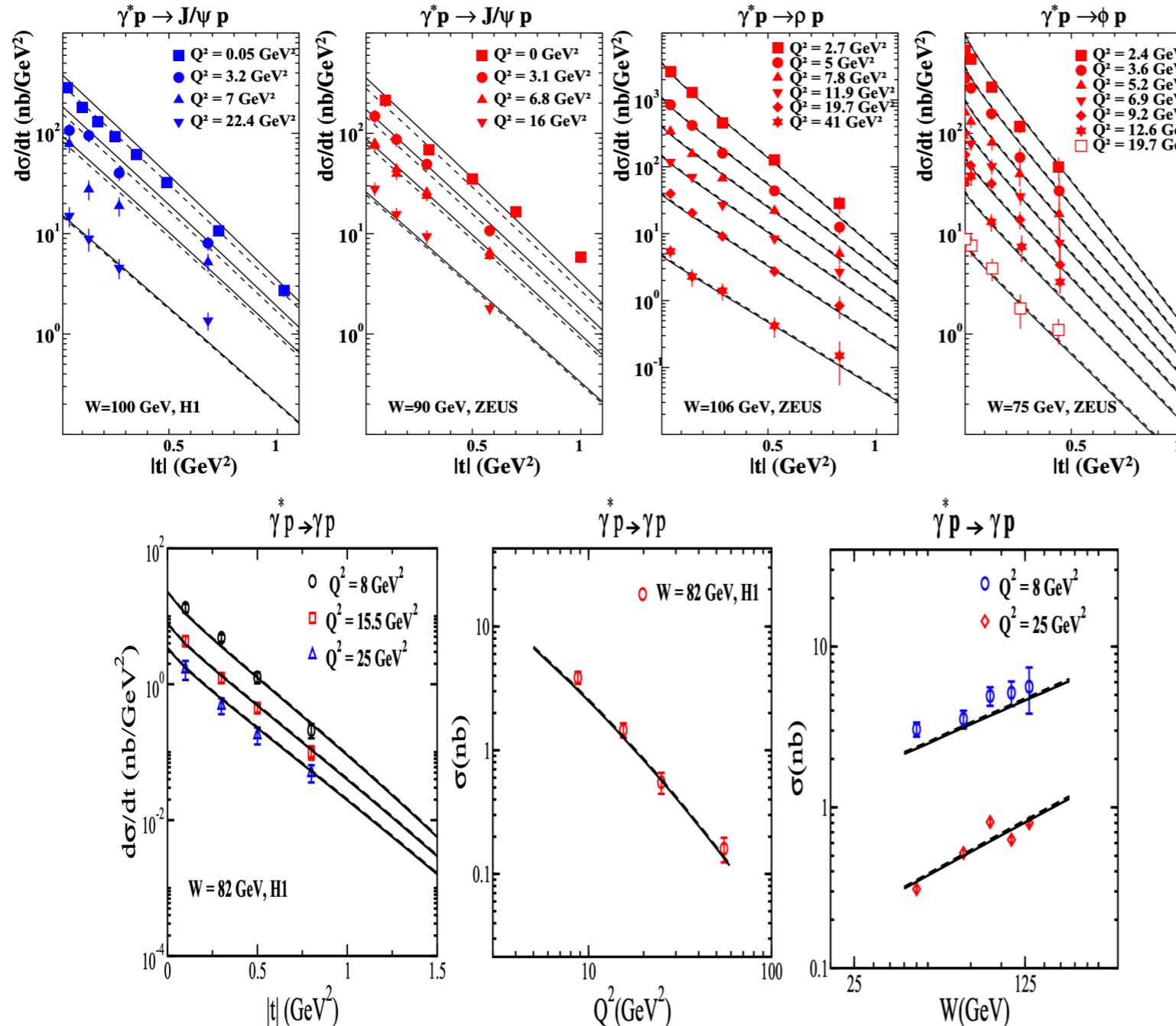
Saturation model (e.g. GBW or IPSAT) modify the profile

$$\langle D(\mathbf{r}_\perp, \mathbf{b}_\perp) \rangle_x = 1 - e^{-g(x, 1/\mathbf{r}_\perp^2) T(\mathbf{b}_\perp)}$$

Teaney and Kowalski (PRD 2003)

Diffraction, Saturation and Vector Meson Production

Comparison of HERA to saturation



Extract proton size from slope of exclusive spectra

Saturation model = IPSAT

Rezaeian, Siddikov, Klundert, Venugopalan (PRD 2013)

Diffraction, Saturation and Vector Meson Production

Incoherent and fluctuations

- Incoherent: Target breaks-up

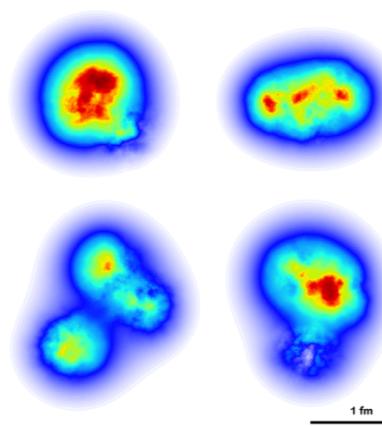
$$d\sigma_{incoh} \propto \langle \mathcal{A}^\dagger(\Delta_\perp) \mathcal{A}(\Delta_\perp) \rangle_x - \langle \mathcal{A}^\dagger(\Delta_\perp) \rangle_x \langle \mathcal{A}(\Delta_\perp) \rangle_x$$

Origin of fluctuations in proton:

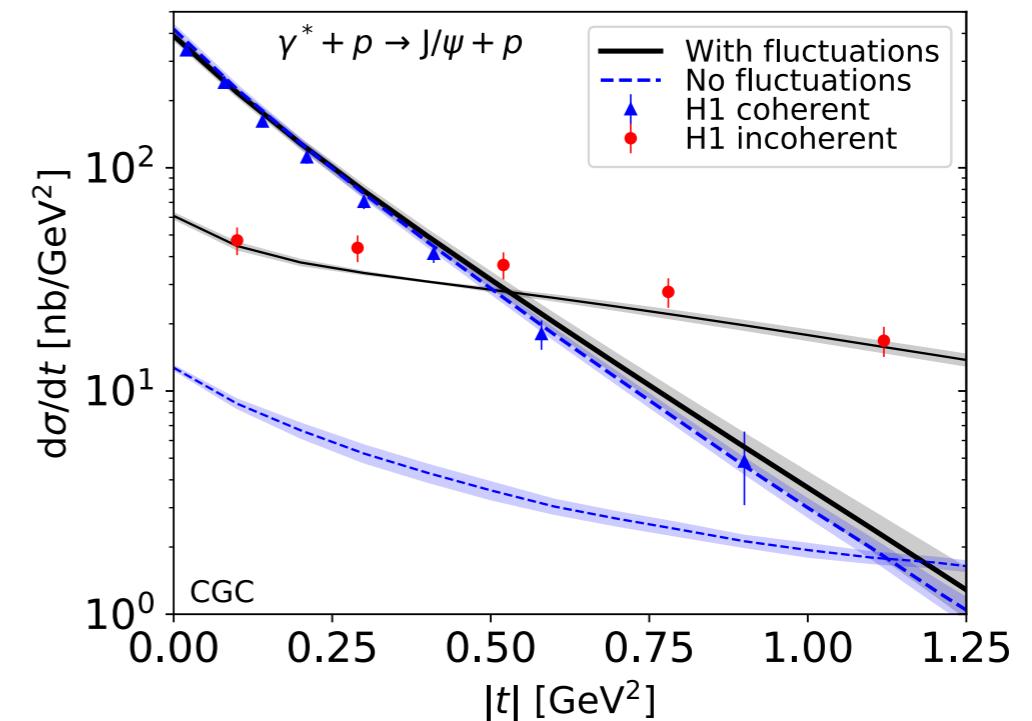
Color charge correlations in CGC (beyond IPSAT or GBW)

$$\langle D^\dagger(\mathbf{r}_\perp, \mathbf{b}_\perp) D(\mathbf{r}'_\perp, \mathbf{b}'_\perp) \rangle_x - \langle D^\dagger(\mathbf{r}_\perp, \mathbf{b}_\perp) \rangle_x \langle D(\mathbf{r}'_\perp, \mathbf{b}'_\perp) \rangle_x \sim \mathcal{O}(1/N_c^2)$$

Sub-nucleon fluctuations (hotspots)

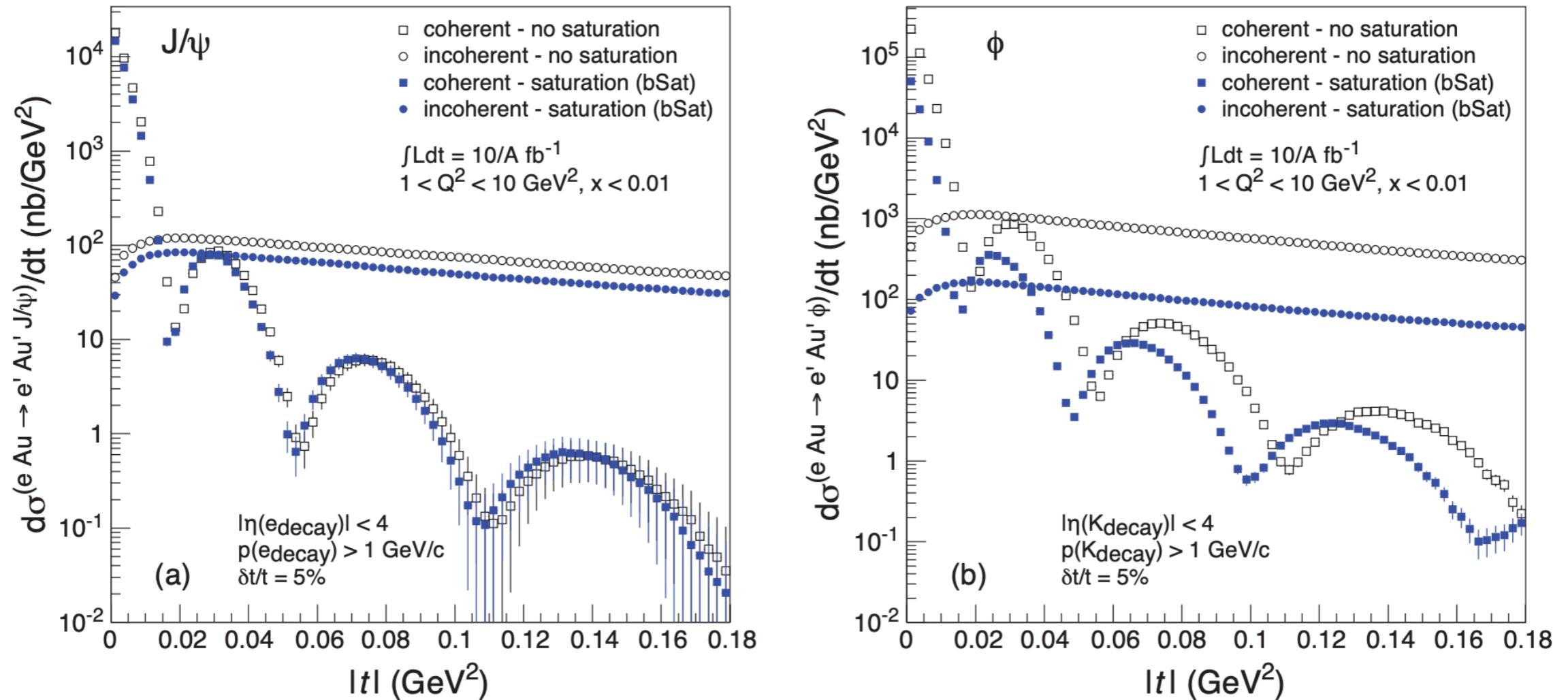


Mantysaari and Schenke (PRL 2016)



Diffraction, Saturation and Vector Meson Production

Expectations for the future EIC



Toll, Ullrich (PRC 2013)

Characteristic dips at \sim multiples of nuclear size (linearized model)

Saturation modifies $|t|$ spectrum, large modifications for ϕ less so for J/ψ

Incoherent production dominates $|t| > 0.02 \text{ GeV}^2$

Diffraction, Saturation and Vector Meson Production

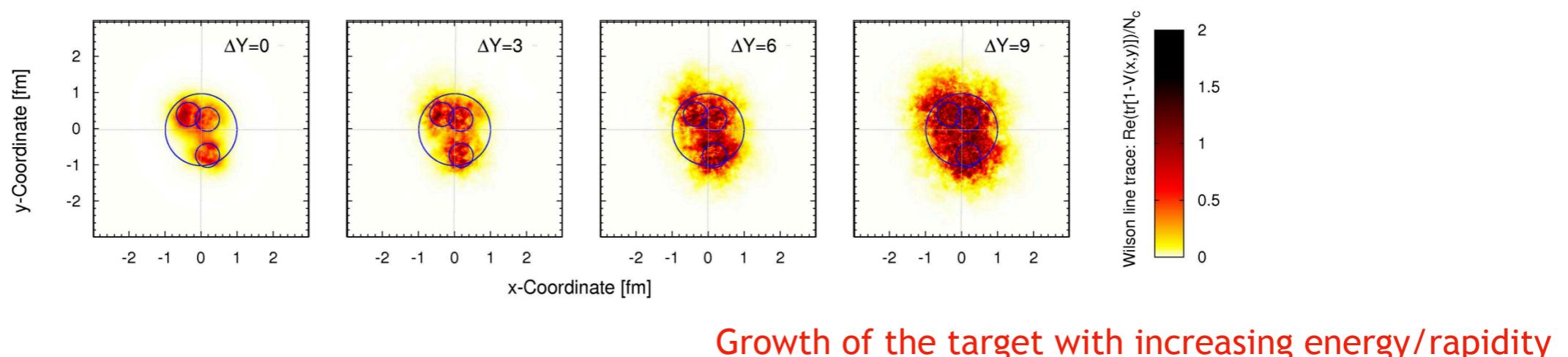
Our setup

- Initial condition: McLerran-Venugopalan model with nucleon substructure

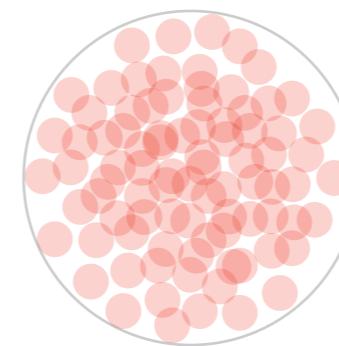
$$\langle \rho^a(\mathbf{x}_\perp) \rho^b(\mathbf{y}_\perp) \rangle \propto \delta^{ab} \delta^{(2)}(\mathbf{x}_\perp - \mathbf{y}_\perp) T_p(\mathbf{x}_\perp) \quad T_p(\mathbf{b}_\perp) = \frac{1}{2\pi B_q N_q} \sum_{i=1}^{N_q} p_i e^{-(\mathbf{b}_\perp - \mathbf{b}_{i\perp})^2/(2B_q)}$$

Parameters constrained by HERA data

- Small-x evolution (gluon emissions/recombinations) via JIMWLK



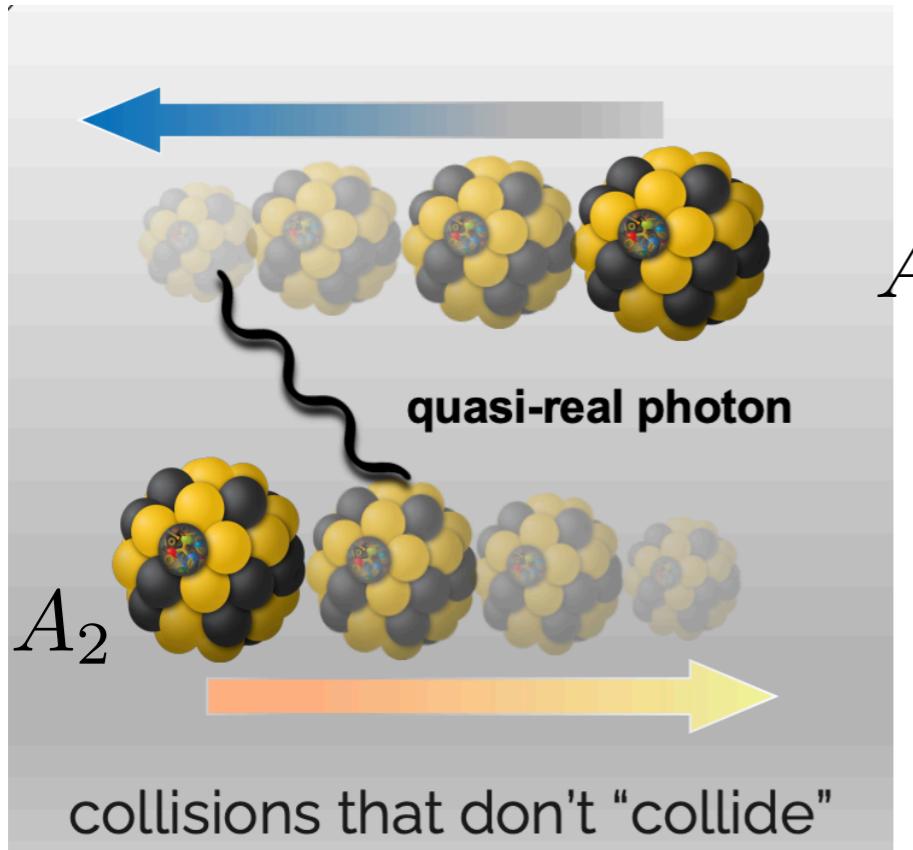
- For collisions with nuclei: we sample nucleons according to Woods-Saxon distribution



Connection to heavy ions see Mäntysaari, Schenke, Shen and Tribedy (PLB 2017)

Ultra-peripheral collisions

Photon-nucleus collisions with heavy-ions



Ion (electrically charged) serves as a photon source, the other ion serves as the nuclear target

Minimum distance between ions:

$$B_{min} \geq R_{A_1} + R_{A_2}$$

Credit: Kong Tu DIS 2023

Photon fluxes

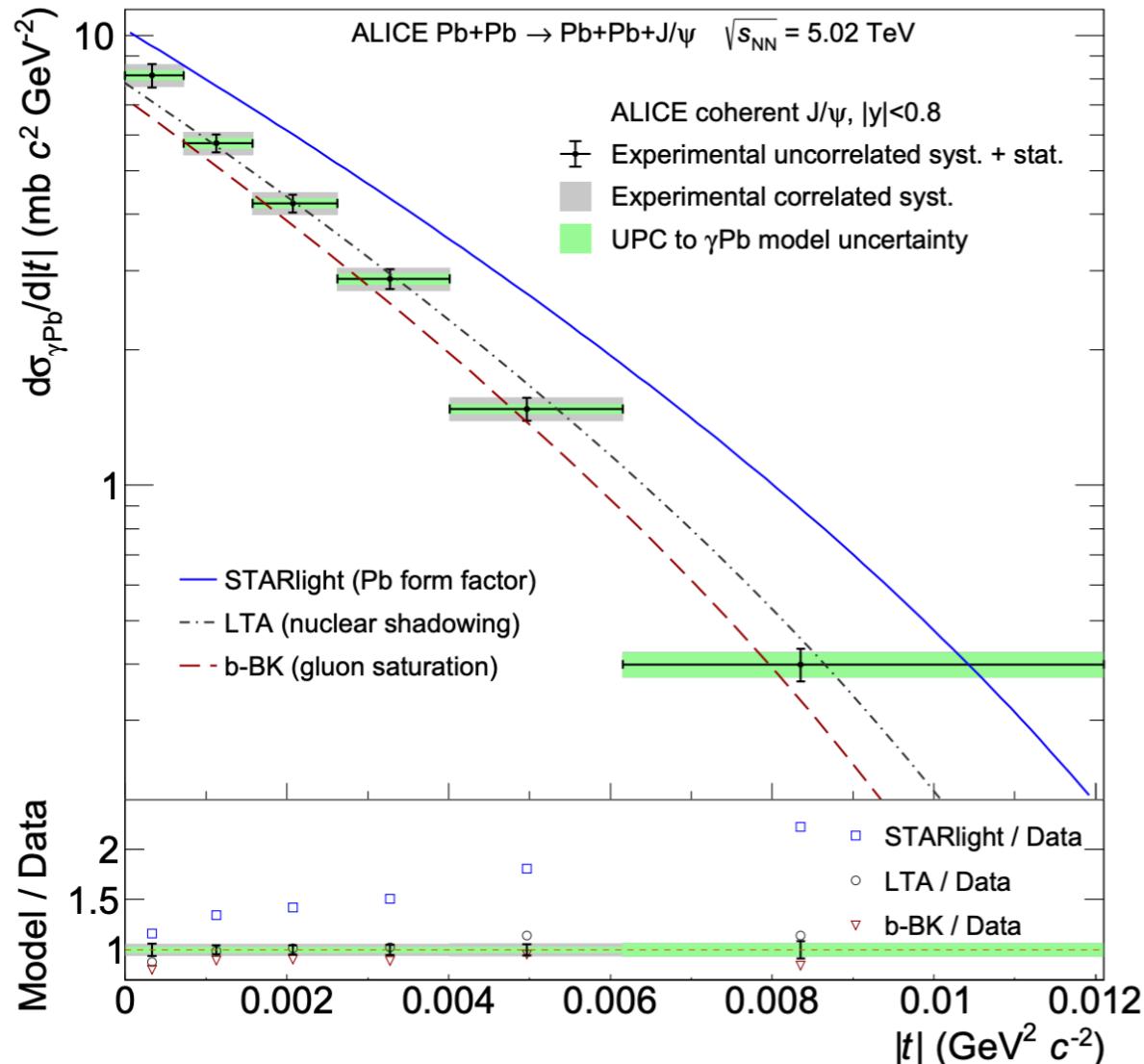
$$\frac{d\sigma^{A_1+A_2 \rightarrow A_1+A_2+V}}{d|t|dy} = N(x_1) \frac{d\sigma^{\gamma+A_2 \rightarrow A_2+V}}{d|t|dy} \Big|_{x_2} + N(x_2) \frac{d\sigma^{\gamma+A_1 \rightarrow A_1+V}}{d|t|dy} \Big|_{x_1}$$

$$x_1 = \frac{M_V e^{-y}}{\sqrt{s}}$$

$$x_2 = \frac{M_V e^y}{\sqrt{s}}$$

Ultra-peripheral collisions

Results from LHC for J/Ψ in γPb

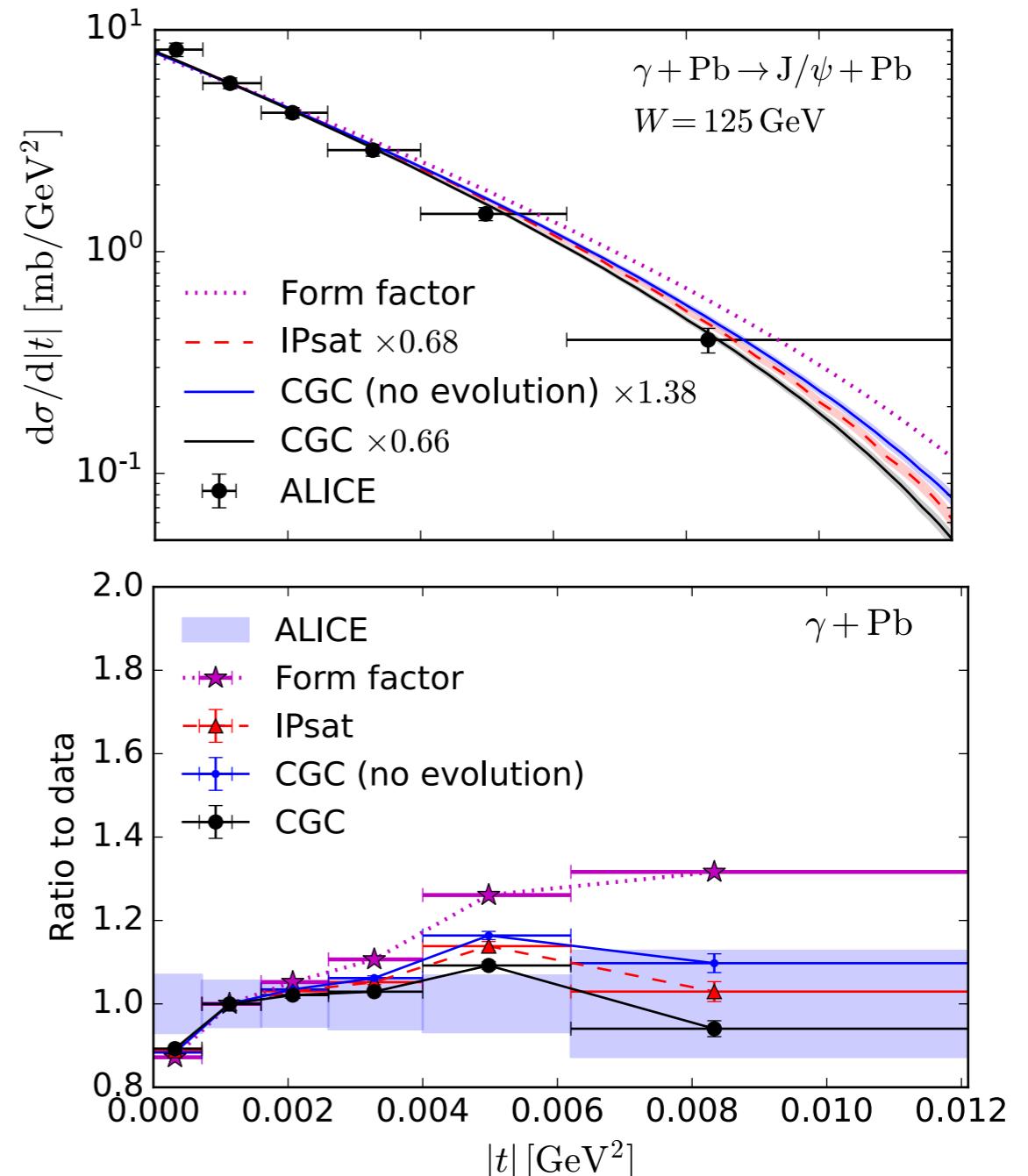


ALICE extracts (via unfolding) $\gamma^* Pb$ collision
from UPC (PLB 2021)

STARlight = FT of form factor (FT of Woods-Saxon)

LTA = Leading Twist Shadowing

b-BK= saturation framework



H. Mäntysaari, FS, B. Schenke (PRD 2022)

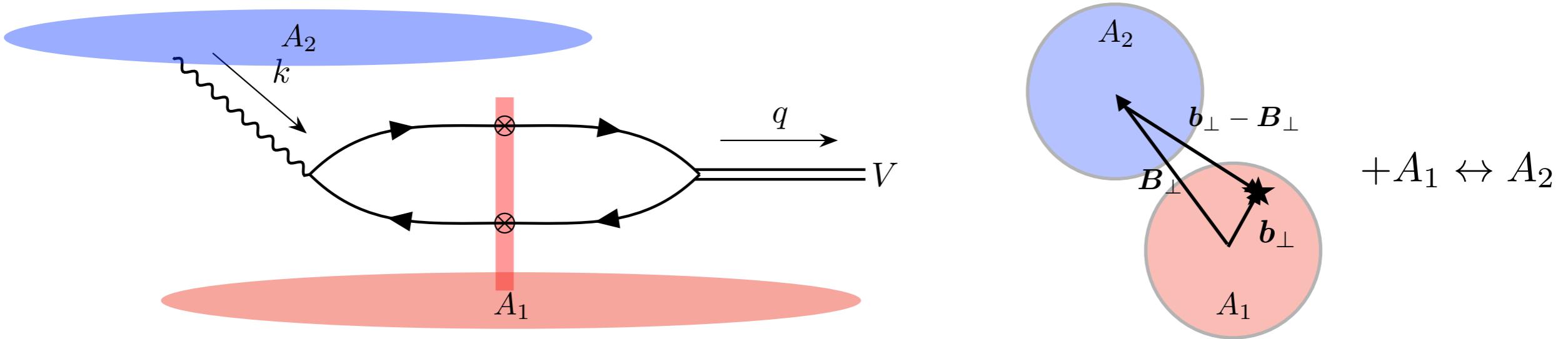
Ultra-peripheral collisions

Quantum interference and finite photon kT

“The emitters and targets can switch roles, and the two possibilities are indistinguishable, so [quantum] interference may occur” Klein, Nystrand (PRL 2000)

Photon carries small but non-zero transverse momentum kT

Joint impact parameter-TMD framework proposed by Xing, Zhang, Zhou, Zhou (JHEP 2020)



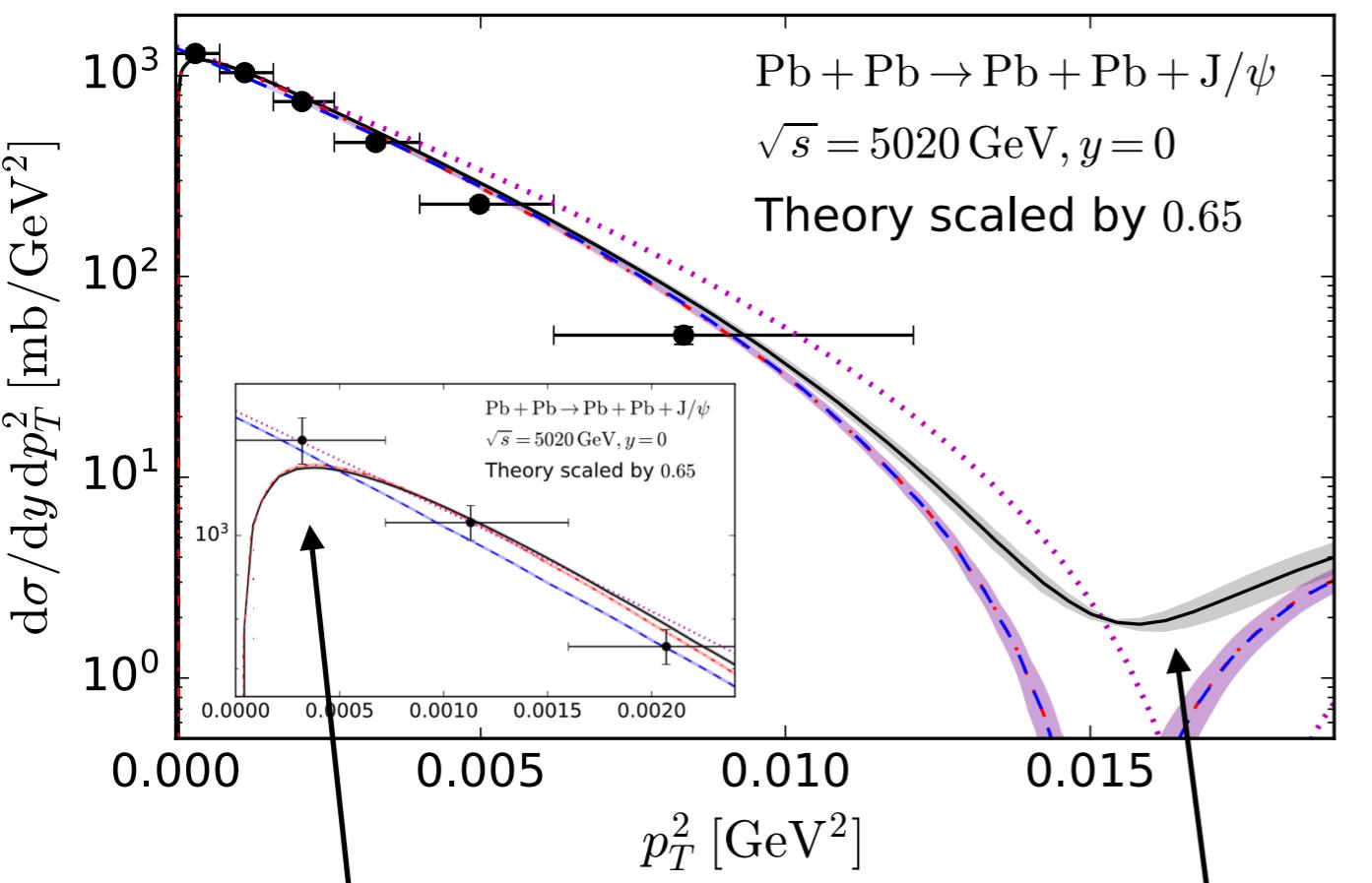
$$\frac{d\sigma^{A_1+A_2 \rightarrow V+A_1+A_2}}{dp_{\perp}^2 dy} = \frac{1}{4\pi} \int_{|B_{\perp}| > B_{\min}} d^2 B_{\perp} \langle \mathcal{M}_{\lambda}^{\dagger}(y, p_{\perp}, B_{\perp}) \rangle \langle \mathcal{M}_{\lambda}(y, p_{\perp}, B_{\perp}) \rangle$$

$$\mathcal{M}^{\lambda}(y, p_{\perp}, B_{\perp}) = \int_{b_{\perp}} e^{-ip_{\perp} \cdot b_{\perp}} \left[\langle \tilde{\mathcal{A}}(b_{\perp}) \rangle_{x_1} \tilde{\mathcal{F}}_{\perp}^{\alpha}(x_2, b_{\perp} - B_{\perp}) + \langle \tilde{\mathcal{A}}(b_{\perp} - B_{\perp}) \rangle_{x_2} \tilde{\mathcal{F}}_{\perp}^{\alpha}(x_1, b_{\perp}) \right] \epsilon_{\alpha}^{V\lambda,*}$$

Incorporates both quantum interference and non-zero photon kT

Ultra-peripheral collisions

Results from LHC for J/Ψ in $PbPb$



Photon kT washes away the
diffractive dip

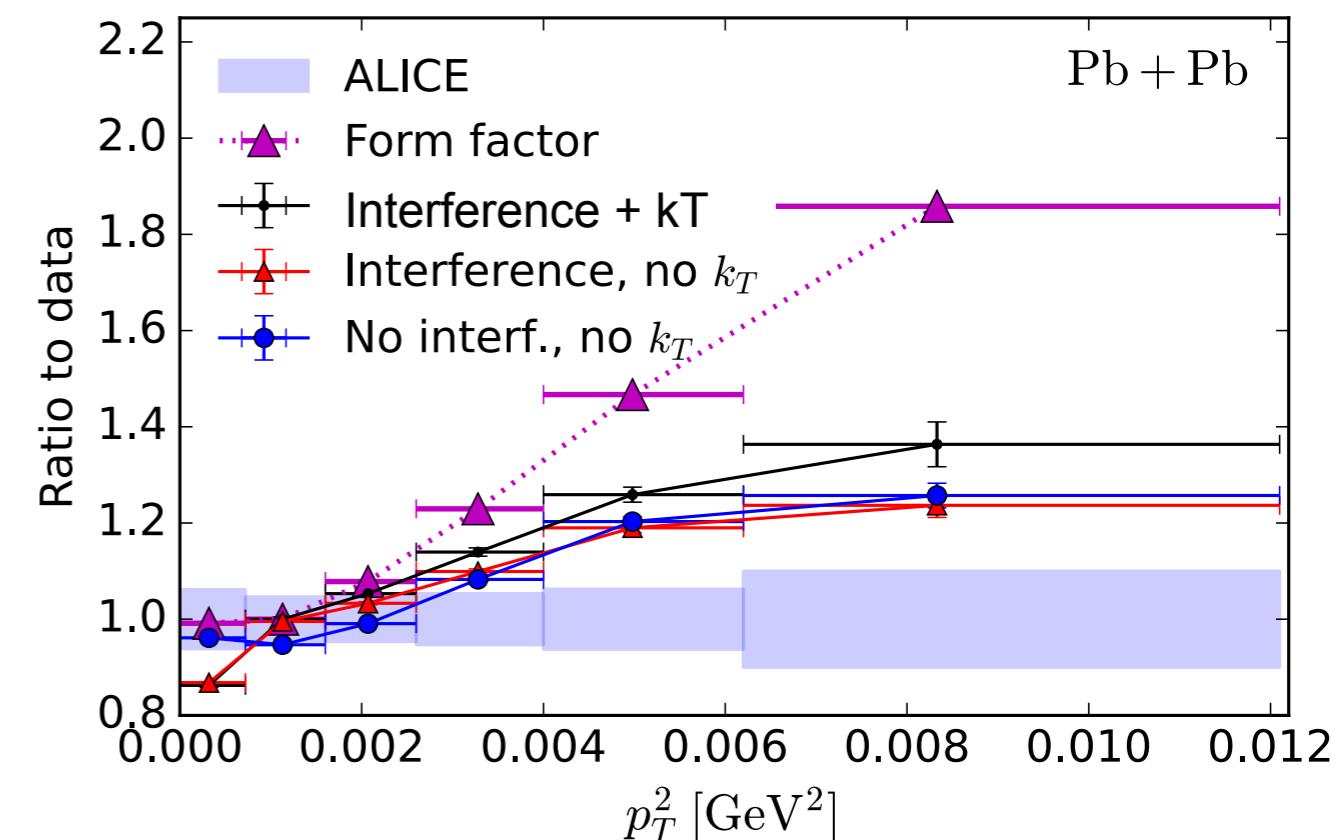
Interference at play

Cross-section vanishes at low p_\perp

Larger discrepancy between theory and data for $PbPb$ than in γPb .

Unfolding seems to be inconsistent with the joint-impact parameter TMD framework

All results computed with CGC

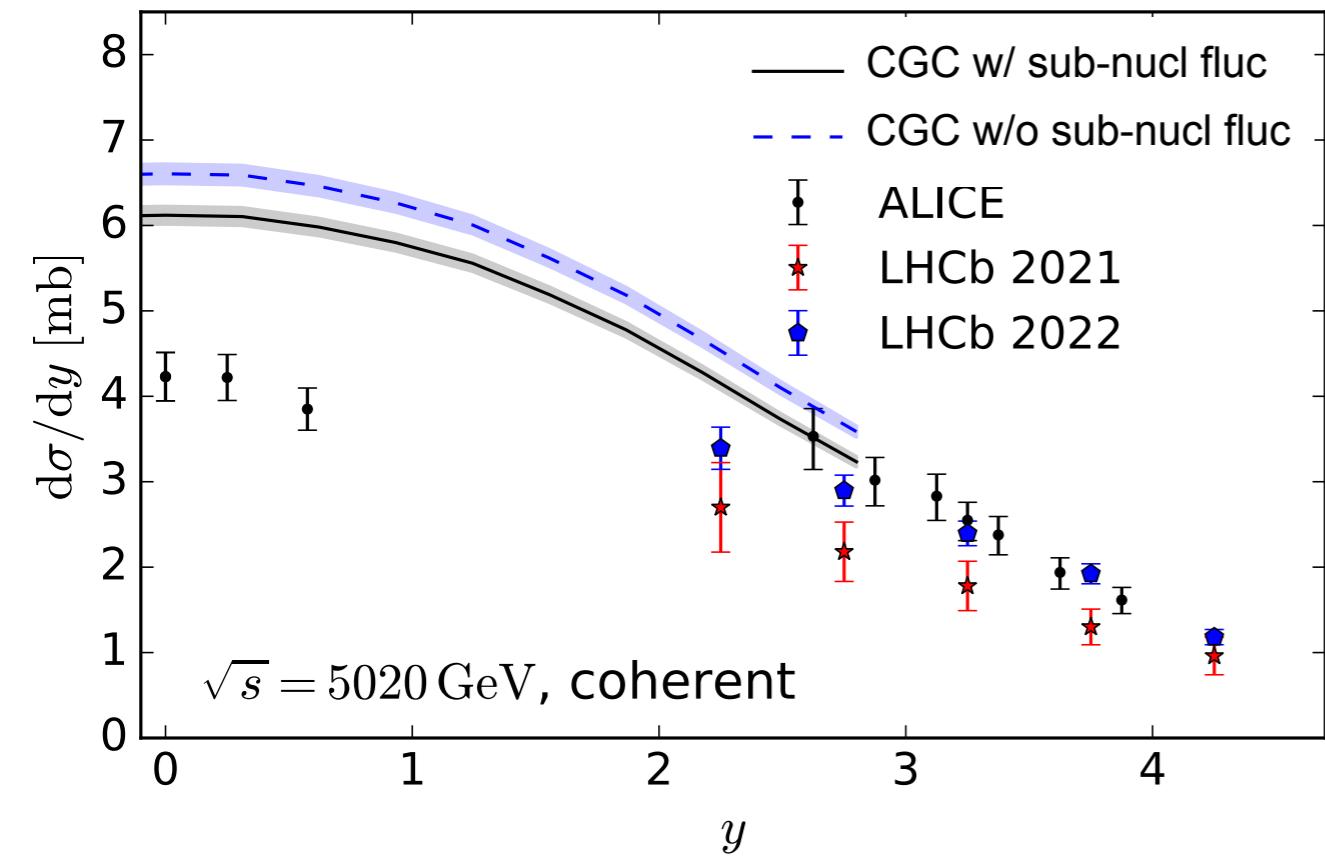


Theory: Photon kT makes pT
spectrum wider

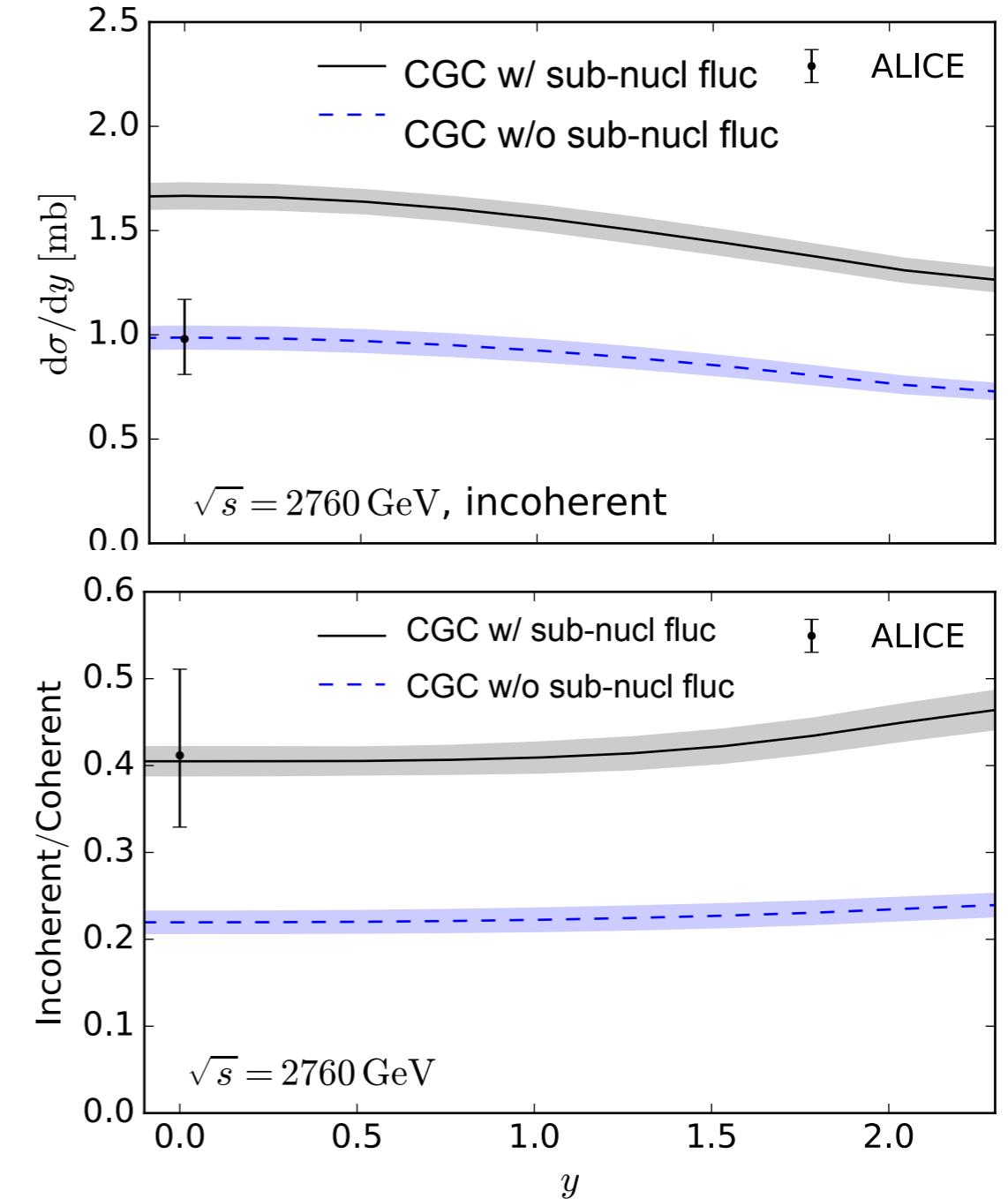
Data: Photon kT makes pT spectrum
steeper. Counterintuitive!

Ultra-peripheral collisions

Rapidity dependence and incoherent production



Normalization fixed by HERA ep data
Current setup has too little nuclear suppression



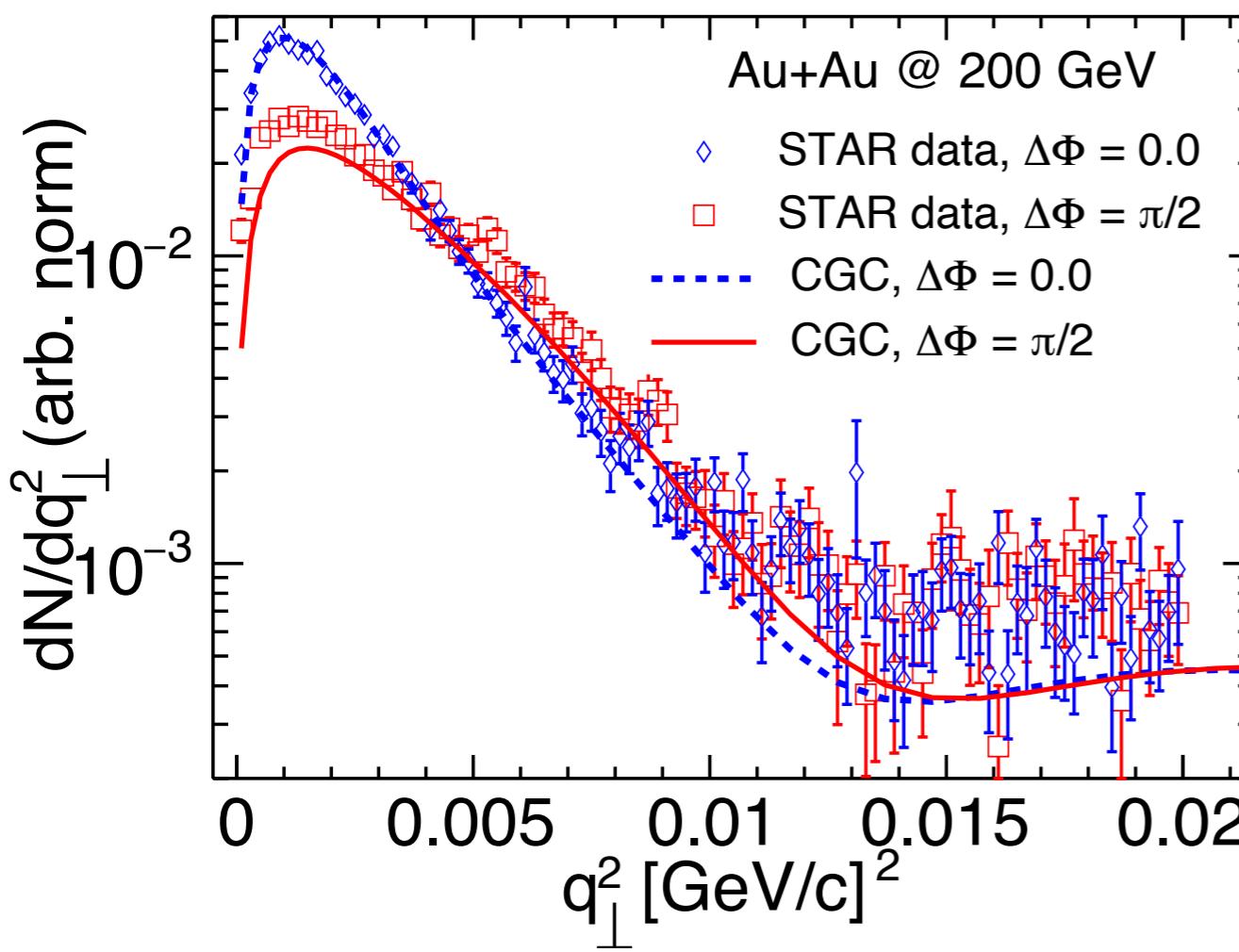
Normalization should cancel in ratio
coherent-to-incoherent. Data seems to
favor sub-nucleon fluctuations

Ultra-peripheral collisions

Azimuthal correlations at RHIC

Correlations in daughter particles $\rho \rightarrow \pi^+ \pi^-$

$$P_\perp = \frac{1}{2}(p_\perp^{\pi^+} - p_\perp^{\pi^-}) \quad q_\perp = p_\perp^{\pi^+} + p_\perp^{\pi^-}$$

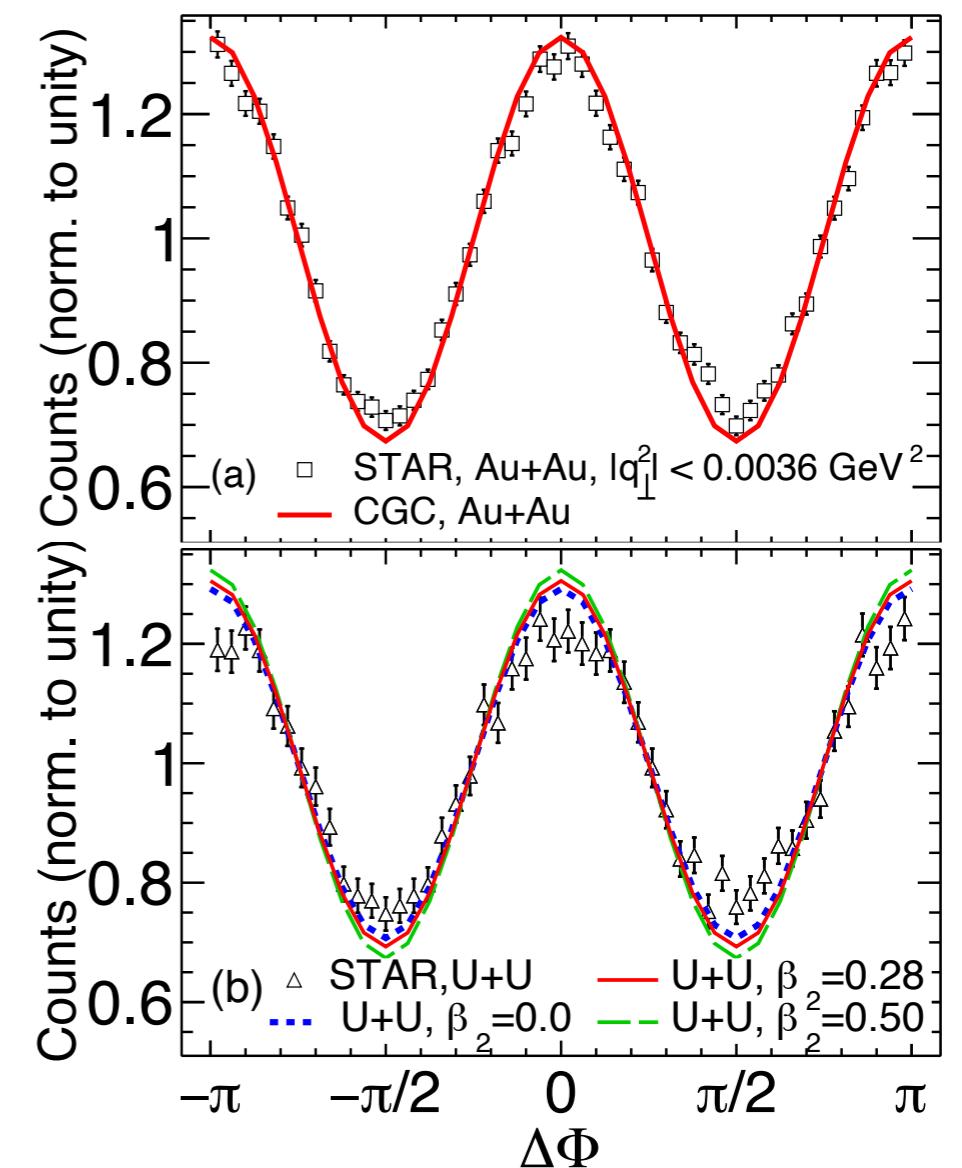


Correlations induced by interference and finite photon kT

Work in progress 2310.XXXX. HM, FS, BS, Chun Shen, Wenbin Zhao

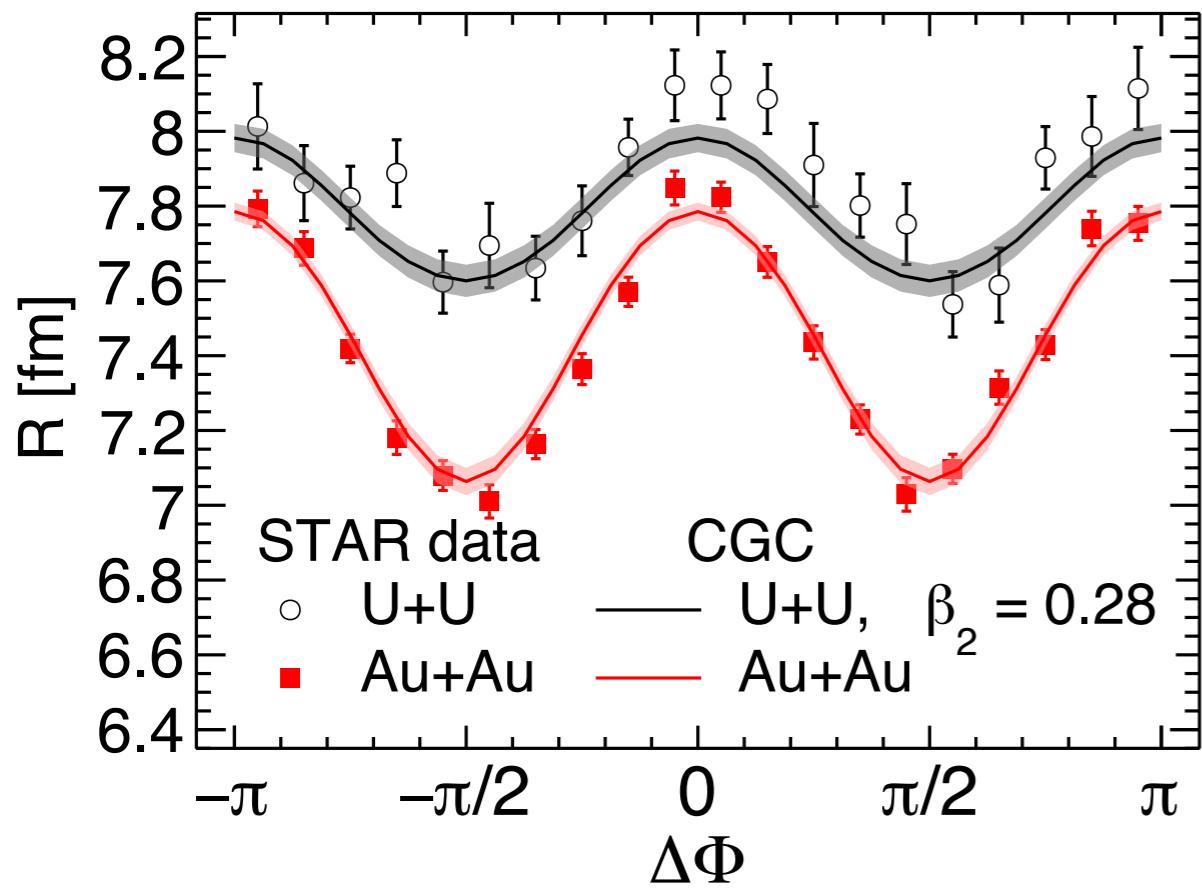
See more at Yajin's talk

$\Delta\Phi$ angle between P_\perp and q_\perp

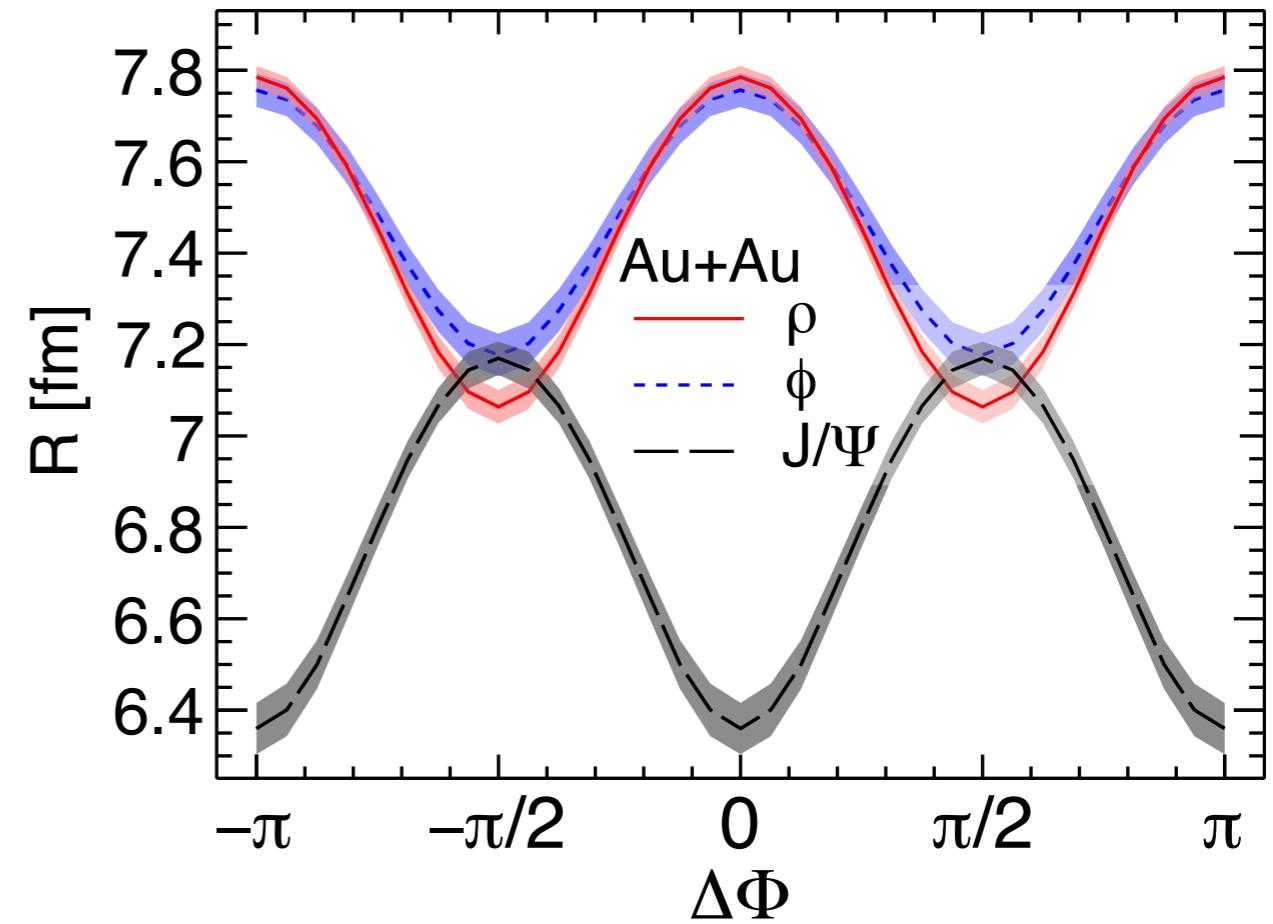


Ultra-peripheral collisions

Angle dependent nuclear radius



Nuclear radius (size) extracted from qT spectrum is angle-dependent!



Sign of correlation depends on spin of particle in the final state (spin 0 vs spin 1/2)

Summary

- Diffractive vector meson production in photon-nucleus collisions allows imaging of nuclei at HERA, EIC, and UPCs
- The saturation framework provides a reasonable description of the data (steeper spectrum than the form factor)
- In UPCs, one must take into account interference effects and non-zero photon kT
- Interference leads to interesting azimuthal correlations on the daughter particles