

An extraction of the Sivers, and the Boer-Mulders functions in $SU(3)$ with DNNs



Ishara Fernando & Dustin Keller
September 25-29, 2023



**25TH INTERNATIONAL
SPIN PHYSICS
SYMPOSIUM**

September 24 – 29, 2023
Durham Convention Center
Durham, NC, USA

Based on the recent publication (September 2023)
<https://doi.org/10.1103/PhysRevD.108.054007>



U.S. DEPARTMENT OF
ENERGY

Office of
Science

This work is supported by DOE contract DE-FG02-96ER40950

Outline

- A brief introduction to TMDs
- Sivers asymmetry from SIDIS
- Generalization of $\mathcal{N}_q(x)$ for $SU(3)_{\text{flavor}}$
- Deep Neural Network (DNN) method
- Testing with pseudo-data
- DNN Fits & Results for Sivers function
- DNN Boer-Mulders function extraction
- Summary and Outlook

TMD PDFs

$$\Phi(x, k_T; S) = \int \frac{d\xi^- d\xi_T}{(2\pi)^3} e^{ik \cdot \xi} \langle P, S | \bar{\psi}(0) \mathcal{U}_{[0, \xi]} \psi(\xi) | P, S \rangle |_{\xi^+ = 0}$$

At leading-twist, the Quark correlator can be decomposed into 8 components (6 T - even and 2 T -odd terms)

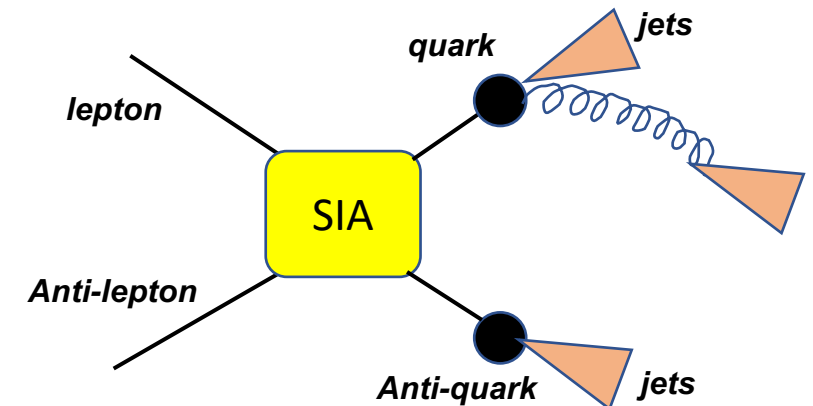
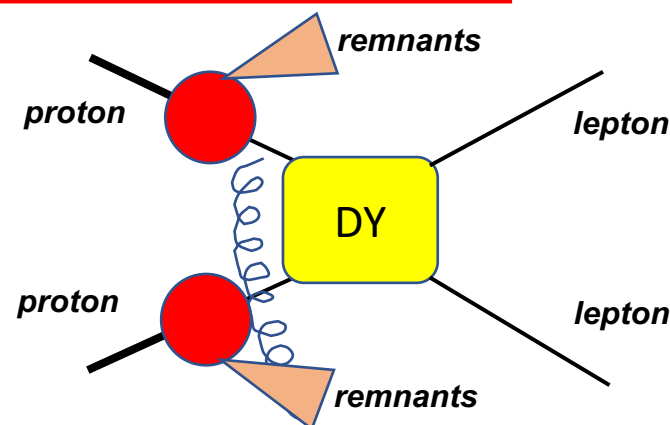
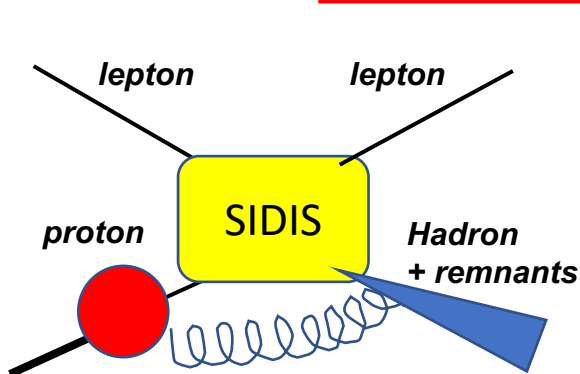
$$\begin{aligned} \Phi(x, k_T, P, S) = & f_1(x, k_T^2) \frac{\not{P}}{2} + \frac{h_{1T}(x, k_T^2)}{4} \gamma_5 [\not{S}_T, \not{P}] + \frac{S_L}{2} g_{1L}(x, k_T^2) \gamma_5 \not{P} + \frac{k_T \cdot S_T}{2M} g_{1T}(x, k_T^2) \gamma_5 \not{P} \\ & + S_L h_{1L}^\perp(x, k_T^2) \gamma_5 \frac{[k_T, \not{P}]}{4M} + \frac{k_T \cdot S_T}{2M} h_{1T}^\perp(x, k_T^2) \gamma_5 \frac{[k_T, \not{P}]}{4M} \end{aligned}$$

T-even

$$+ i h_1^\perp(x, k_T^2) \frac{[k_T, \not{P}]}{4M} - \frac{\epsilon_T^{k_T S_T}}{4M} f_{1T}^\perp(x, k_T^2) \not{P}$$

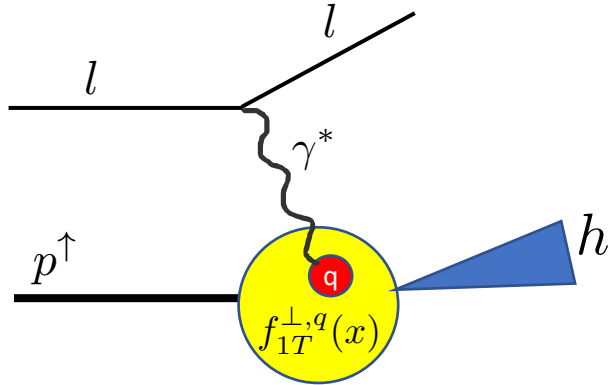
T-odd

		Quark Polarization		
		U	L	T
Nucleon Polarization	U	$f_1 = \odot$	N/A	$h_1^\perp = \odot - \odot$ Boer-Mulders
	L	N/A	$g_{1L} = \odot - \odot$ Helicity	$h_{1L}^\perp = \odot - \odot$
	T	$f_{1T}^\perp = \odot - \odot$ Sivers	$g_{1T}^\perp = \odot - \odot$	$h_1 = \odot - \odot$ $h_{1T}^\perp = \odot - \odot$ Transversity



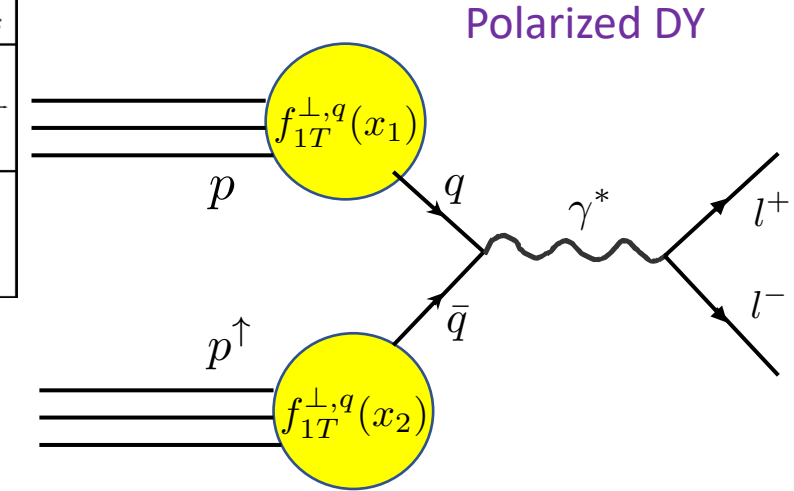
TMD PDFs

Polarized Semi Inclusive DIS



		Quark Polarization		
		U	L	T
Nucleon Polarization	U	$f_1 = \odot$	N/A	$h_1^\perp = \odot - \odot$ Boer-Mulders
	L	N/A	$g_{1L} = \odot - \odot$ Helicity	$h_{1L}^\perp = \odot - \odot$
	T	$f_{1T}^\perp = \odot - \odot$ Sivers	$g_{1T}^\perp = \odot - \odot$	$h_1 = \odot - \odot$ $h_{1T}^\perp = \odot - \odot$ Transversity

* For these two processes
TMD factorization is proven



$$\frac{d\sigma_{SIDIS}^{LO}}{dx dy dz dp_T^2 d\phi_h d\psi} = \left[\frac{\alpha}{xy Q^2} \frac{y^2}{2(1-\epsilon)} \left(1 + \frac{y^2}{2x} \right) \right] \times (F_{UU,T} + \epsilon F_{UU,L}) \left\{ 1 + \cos 2\phi_h \left(\epsilon A_{UU}^{\cos 2\phi_h} \right) + S_T \left[\sin(\phi_h - \phi_s) \left(A_{UT}^{\sin(\phi_h - \phi_s)} \right) + \sin(\phi_h + \phi_s) \left(\epsilon A_{UT}^{\sin(\phi_h + \phi_s)} \right) + \sin(3\phi_h - \phi_s) \left(\epsilon A_{UT}^{\sin(3\phi_h - \phi_s)} \right) \right] \right\}$$

$$\frac{d\sigma^{LO}}{d\Omega} = \frac{\alpha_{em}^2}{F_q} F_v^1 \left\{ 1 + \cos^2 \theta + \sin^2 \theta \cos 2\phi_{CS} A_U^{\cos 2\phi_{CS}} + S_T \left[(1 + \cos^2 \theta) \sin \phi_s A_T^{\sin \phi_s} + \sin^2 \theta \left(\sin(2\phi_{CS} + \phi_s) A_T^{\sin(2\phi_{CS} + \phi_s)} + \sin(2\phi_{CS} - \phi_s) A_T^{\sin(2\phi_{CS} - \phi_s)} \right) \right] \right\}$$

$$A_{UU}^{\cos 2\phi_h} \propto h_1^{\perp q} \otimes H_{1q}^{\perp h}$$

BM * CF

$$A_{UT}^{\sin(\phi_h - \phi_s)} \propto f_{1T}^{\perp q} \otimes D_{1q}^h$$

Sivers * FF

$$A_{UT}^{\sin(\phi_h + \phi_s)} \propto h_1^q \otimes H_{1q}^{\perp h}$$

Transv * CF

$$A_{UT}^{\sin(3\phi_h - \phi_s)} \propto h_{1T}^{\perp q} \otimes H_{1q}^{\perp h}$$

Pretz * CF

$$\left. h_1^{\perp q} \right|_{SIDIS} = - \left. h_1^{\perp q} \right|_{DY}$$

$$\left. f_{1T}^{\perp q} \right|_{SIDIS} = - \left. f_{1T}^{\perp q} \right|_{DY}$$

$$\left. h_1^q \right|_{SIDIS} = \left. h_1^q \right|_{DY}$$

$$\left. h_{1T}^{\perp q} \right|_{SIDIS} = \left. h_{1T}^{\perp q} \right|_{DY}$$

$$A_T^{\cos 2\phi_{CS}} \propto h_1^{\perp q} \otimes h_1^{\perp q}$$

BM * BM

$$A_T^{\sin \phi_s} \propto f_1^q \otimes f_{1T}^{\perp q}$$

PDF * Sivers

$$A_T^{\sin(2\phi_{CS} - \phi_s)} \propto h_1^{\perp q} \otimes h_1^q$$

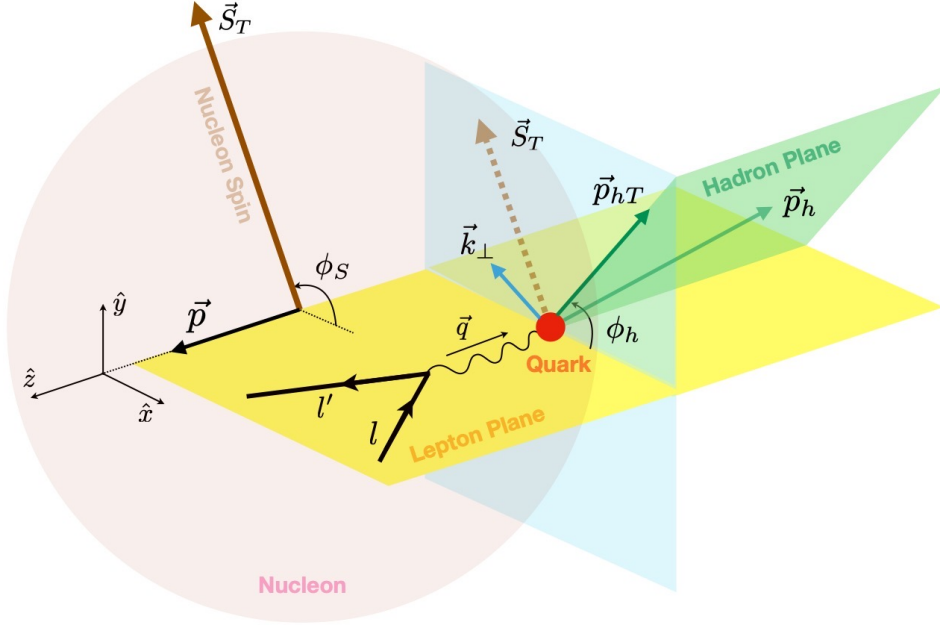
BM * Transv

$$A_T^{\sin(2\phi_{CS} + \phi_s)} \propto h_1^{\perp q} \otimes h_{1T}^{\perp q}$$

BM * Pretz

Sivers Asymmetry from SIDIS

$$\frac{d^5\sigma^{lp\rightarrow lhX}}{dx dQ^2 dz d^2p_\perp} = \sum_q e_q^2 \int d^2\mathbf{k}_\perp \left(\frac{2\pi\alpha^2}{x^2 s^2} \frac{\hat{s}^2 + \hat{u}^2}{Q^4} \right) \times \hat{f}_{q/p^\uparrow}(x, k_\perp) D_{h/q}(z, p_\perp) + \mathcal{O}(k_\perp/Q) ,$$



$$\hat{f}_{q/p^\uparrow}(x, k_\perp) = f_{q/p}(x, k_\perp) + \frac{1}{2} \Delta^N f_{q/p^\uparrow}(x, k_\perp) \vec{S}_T \cdot (\hat{p} \times \hat{k}_\perp)$$

$$\Delta^N f_{q/p^\uparrow}(x, k_\perp) = 2\mathcal{N}_q(x) h(k_\perp) f_{q/p}(x, k_\perp)$$

Anselmino et al. (2017)

Single Spin Asymmetry (Sivers Asymmetry)

$$\langle p_\perp^2 \rangle = 0.12 \pm 0.01 \text{ GeV}^2$$

$$\langle k_\perp^2 \rangle = 0.57 \pm 0.08 \text{ GeV}^2$$

$$A_{UT}^{\sin(\phi_h - \phi_S)}(x, y, z, p_{hT}) = \frac{d\sigma^{l^\uparrow p \rightarrow hlX} - d\sigma^{l^\downarrow p \rightarrow hlX}}{d\sigma^{l^\uparrow p \rightarrow hlX} + d\sigma^{l^\downarrow p \rightarrow hlX}} \equiv \frac{d\sigma^\uparrow - d\sigma^\downarrow}{d\sigma^\uparrow + d\sigma^\downarrow}$$

$$\mathcal{A}_0(z, p_{hT}, m_1)$$

$$= \frac{\sqrt{2} e z p_{hT}}{m_1} \frac{[z^2 \langle k_\perp^2 \rangle + \langle p_\perp^2 \rangle] \langle k_S^2 \rangle^2}{[z^2 \langle k_S^2 \rangle + \langle p_\perp^2 \rangle]^2 \langle k_\perp^2 \rangle} \times \exp \left[- \frac{p_{hT}^2 z^2 (\langle k_S^2 \rangle - \langle k_\perp^2 \rangle)}{(z^2 \langle k_S^2 \rangle + \langle p_\perp^2 \rangle) (z^2 \langle k_\perp^2 \rangle + \langle p_\perp^2 \rangle)} \right]$$

$$\langle k_S^2 \rangle = \frac{m_1 \langle k_\perp^2 \rangle}{m_1^2 + \langle k_\perp^2 \rangle}$$

$$A_{UT}^{\sin(\phi_h - \phi_S)}(x, z, p_{hT}) = \mathcal{A}_0(z, p_{hT}, m_1) \left(\frac{\sum_q \mathcal{N}_q(x) e_q^2 f_q(x) D_{h/q}(z)}{\sum_q e_q^2 f_q(x) D_{h/q}(z)} \right)$$

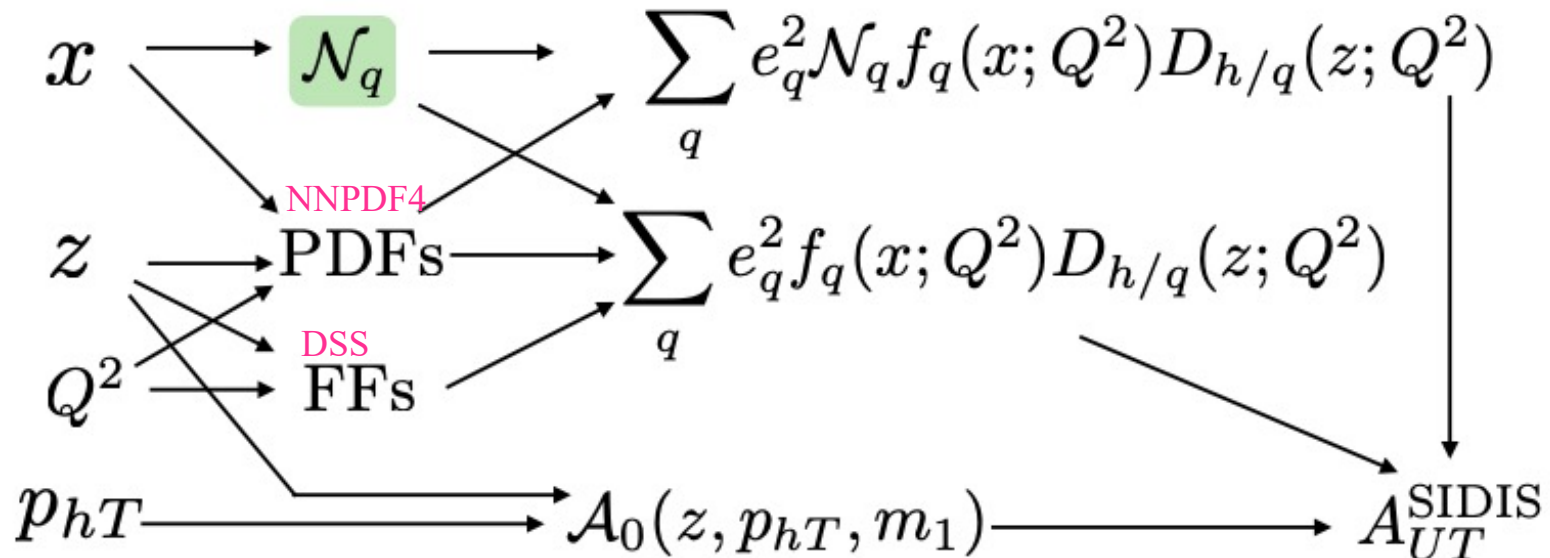
$$\mathcal{N}_q(x) = N_q x^{\alpha_q} (1-x)^{\beta_q} \frac{(\alpha_q + \beta_q)^{(\alpha_q + \beta_q)}}{\alpha_q^{\alpha_q} \beta_q^{\beta_q}}$$

$$\mathcal{N}_{\bar{q}}(x) = N_{\bar{q}}$$

DNN Approach

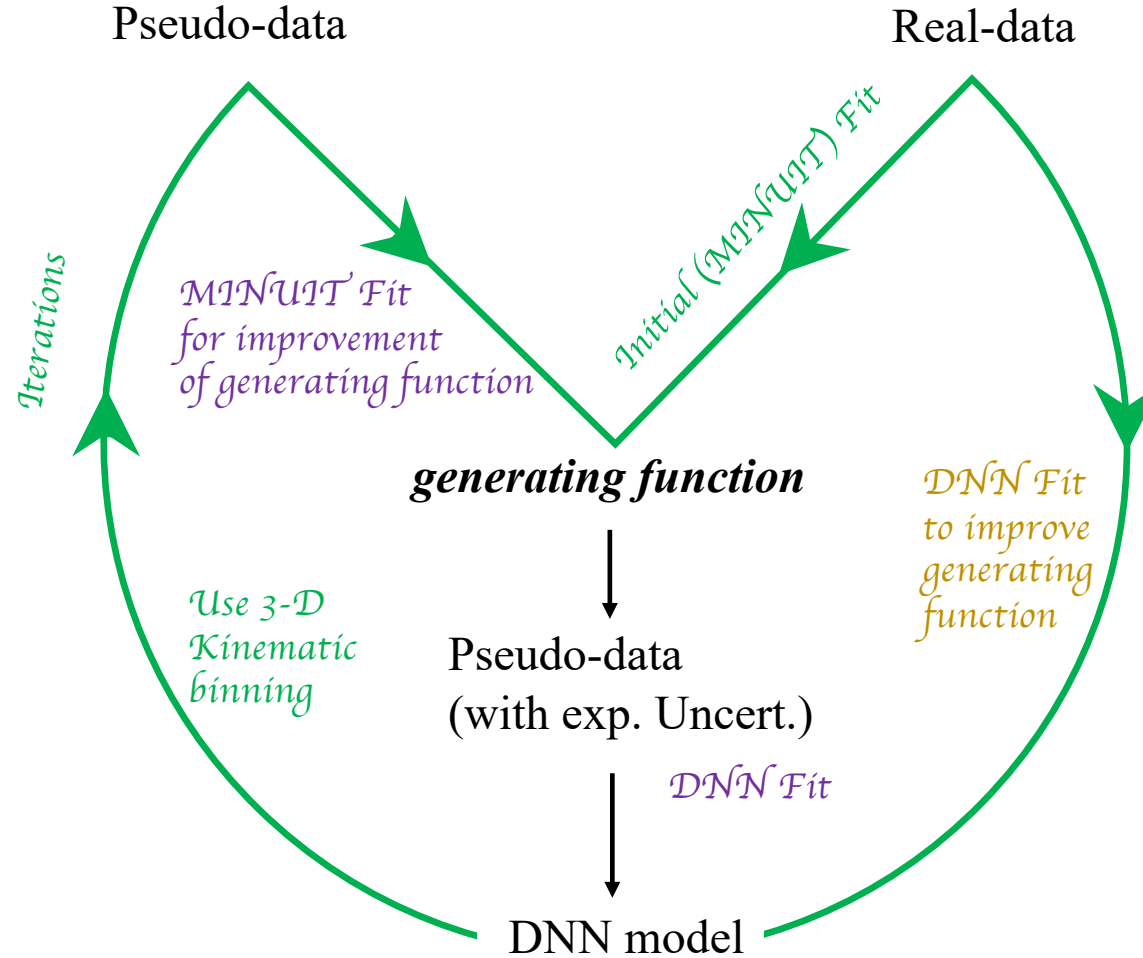
$$A_{UT}^{\sin(\phi_h - \phi_S)}(x, z, p_{hT}) = \mathcal{A}_0(z, p_{hT}, m_1) \left(\frac{\sum_q \mathcal{N}_q(x) e_q^2 f_q(x) D_{h/q}(z)}{\sum_q e_q^2 f_q(x) D_{h/q}(z)} \right)$$

- The exceptional capacity of DNN to be ideal for function approximation (Universal Approximation Theorem).
- Each quark flavor q is independently handled by a separate $\mathcal{N}_q(x)$.
- The only input to each $\mathcal{N}_q(x)$ is x .
- Statistical & Systematic uncertainties from the experimental data are combined in quadrature; then propagated using bootstrap method by generating replicas.
- Systematic uncertainty in method is evaluated with variations in generating function.



The “DNN Method” for extracting TMDs

I. P. Fernando and D. Keller
Phys. Rev. D.108.054007 (2023)



- Systematic study for both DNN models were performed separately using various generating functions.

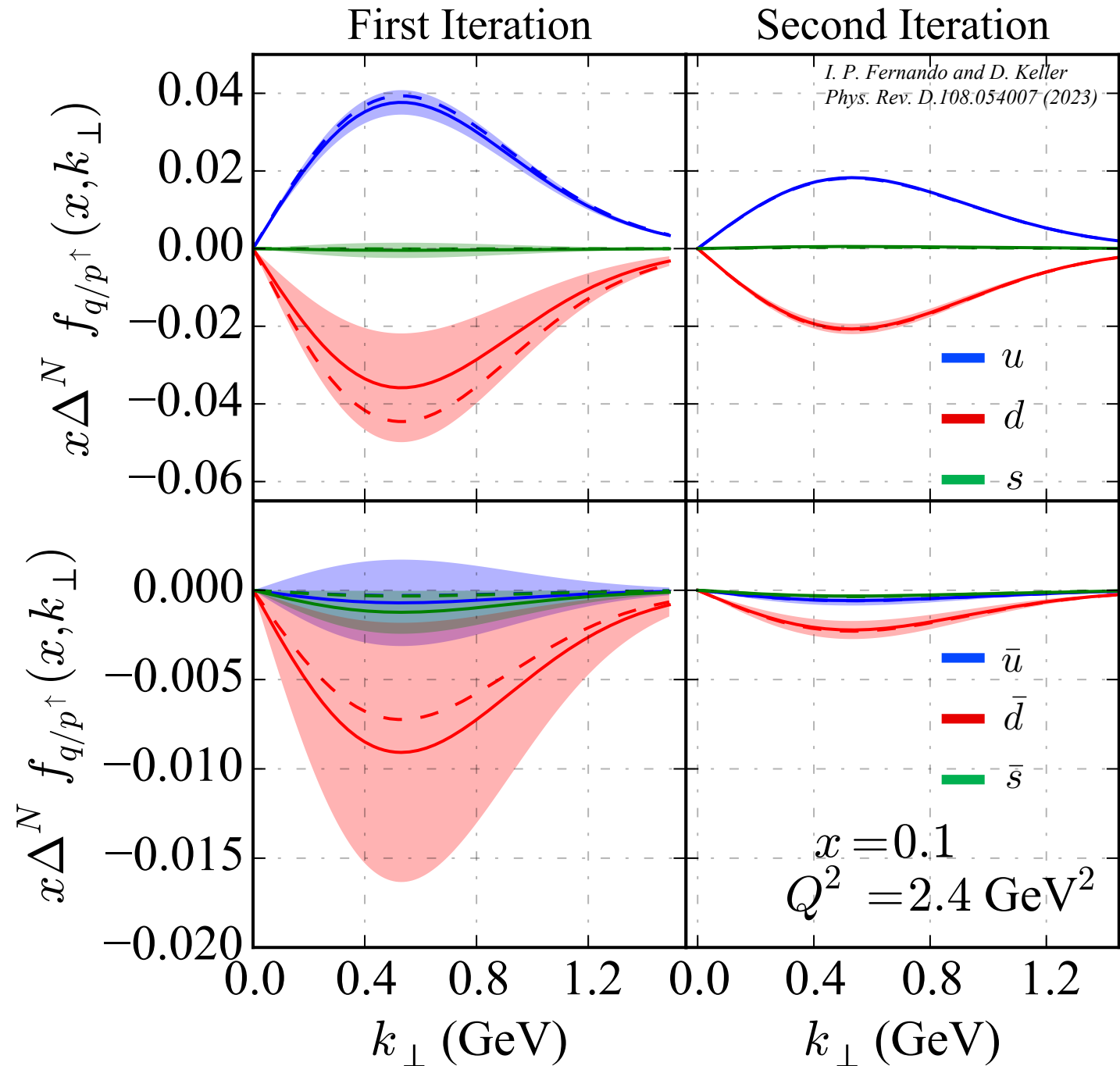
- We trained two separate models for “proton” and “neutron” (deuteron)
- To take full advantage of the information provided by the model testing in the previous slide, the steps from method testing with pseudo-data are performed again separately for proton and deuteron SIDIS data.

$$\epsilon_q(x, k_{\perp}) = \left(1 - \frac{|\Delta^N f_{q/p^{\uparrow}}^{(\text{true})} - \Delta^N f_{q/p^{\uparrow}}^{(\text{mean})}|}{\Delta^N f_{q/p^{\uparrow}}^{(\text{true})}} \right) \times 100\%$$

$$\sigma_q(x, k_{\perp}) = \sqrt{\frac{\sum_i \left(\Delta^N f_{q/p^{\uparrow}}^{(i)} - \Delta^N f_{q/p^{\uparrow}}^{(\text{mean})} \right)^2}{N}}$$

DNN Method testing with Pseudo data

- Dashed lines represent the **generating function** in each iteration.
- A comparison:
Improving the **generating function**
Fine-tuning the hyperparameters
- Solid-lines and the band represent the mean and 68% CL with 1000 replicas of the DNN model.



Data Selection

Proton DNN
model

Deuteron DNN
model

Projections from
Deuteron DNN model

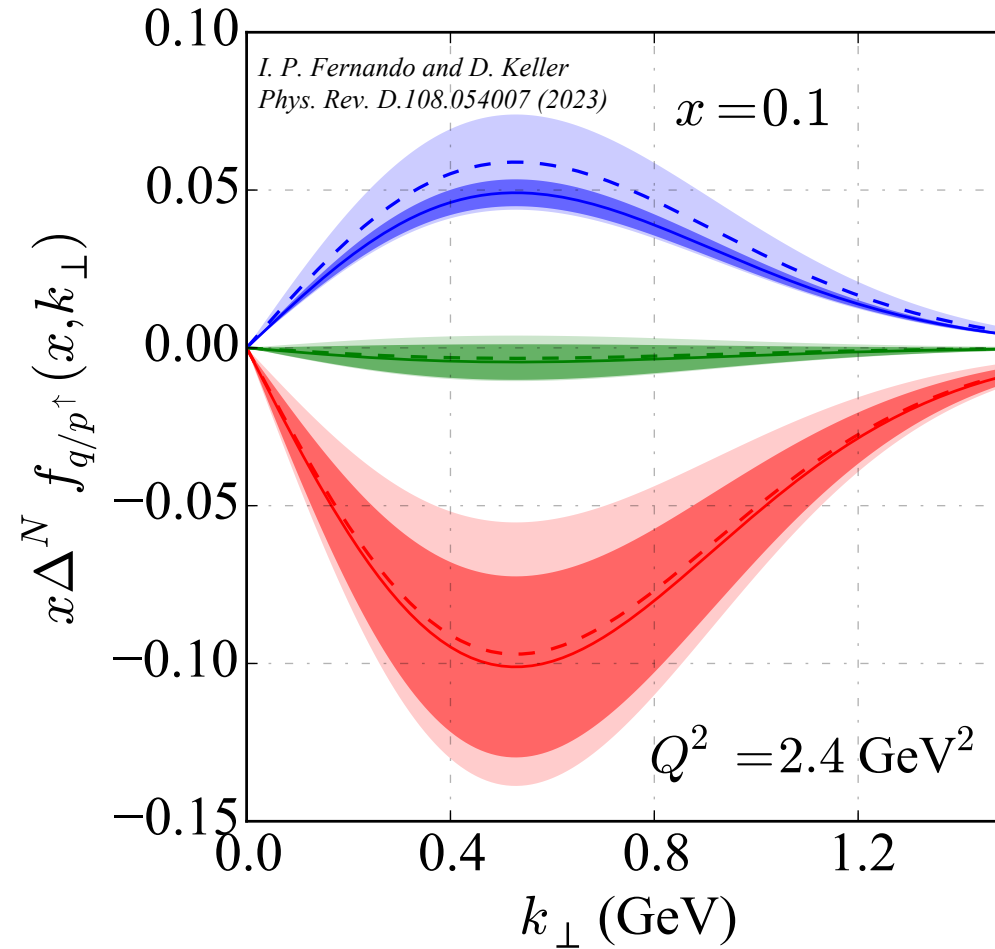
Dataset	Kinematic coverage	Reaction	Data points
HERMES2009 (SIDIS) [53]	$0.023 < x < 0.4$	$p^\uparrow + \gamma^* \rightarrow \pi^+$	21
	$0.2 < z < 0.7$	$p^\uparrow + \gamma^* \rightarrow \pi^-$	21
	$0.1 < p_{hT} < 0.9$	$p^\uparrow + \gamma^* \rightarrow \pi^0$	21
	$Q^2 > 1 \text{ GeV}^2$	$p^\uparrow + \gamma^* \rightarrow K^+$	21
		$p^\uparrow + \gamma^* \rightarrow K^-$	21
HERMES2020 (SIDIS) [55]	$0.023 < x < 0.6$	$p^\uparrow + \gamma^* \rightarrow \pi^+$	27, 64
	$0.2 < z < 0.7$	$p^\uparrow + \gamma^* \rightarrow \pi^-$	27, 64
	$0.1 < p_{hT} < 0.9$	$p^\uparrow + \gamma^* \rightarrow \pi^0$	27
	$Q^2 > 1 \text{ GeV}^2$	$p^\uparrow + \gamma^* \rightarrow K^+$	27, 64
		$p^\uparrow + \gamma^* \rightarrow K^-$	27, 64
COMPASS2015 (SIDIS) [54]	$0.006 < x < 0.28$	$p^\uparrow + \gamma^* \rightarrow \pi^+$	26
	$0.2 < z < 0.8$	$p^\uparrow + \gamma^* \rightarrow \pi^-$	26
	$0.15 < p_{hT} < 1.5$	$p^\uparrow + \gamma^* \rightarrow K^+$	26
	$Q^2 > 1 \text{ GeV}^2$	$p^\uparrow + \gamma^* \rightarrow K^-$	26
COMPASS2009 (SIDIS) [49]	$0.006 < x < 0.28$	$d^\uparrow + \gamma^* \rightarrow \pi^+$	26
	$0.2 < z < 0.8$	$d^\uparrow + \gamma^* \rightarrow \pi^-$	26
	$0.15 < p_{hT} < 1.5$	$d^\uparrow + \gamma^* \rightarrow K^+$	26
	$Q^2 > 1 \text{ GeV}^2$	$d^\uparrow + \gamma^* \rightarrow K^-$	26
JLAB2011 (SIDIS) [52]	$0.156 < x < 0.396$	$^3\text{He}^\uparrow + \gamma^* \rightarrow \pi^+$	4
	$0.50 < z < 0.58$	$^3\text{He}^\uparrow + \gamma^* \rightarrow \pi^-$	4
	$0.24 < p_{hT} < 0.43$ $1.3 < Q^2 < 2.7$		
COMPASS2017 (DY) [50]	$0.1 < x_N < 0.25$	$p^\uparrow + \pi^- \rightarrow l^+ l^- X$	15
	$0.3 < x_\pi < 0.7$		
	$4.3 < Q_M < 8.5$		
	$0.6 < q_T < 1.9$		

HERMES2020
3D binned data

Projections from
Proton DNN model

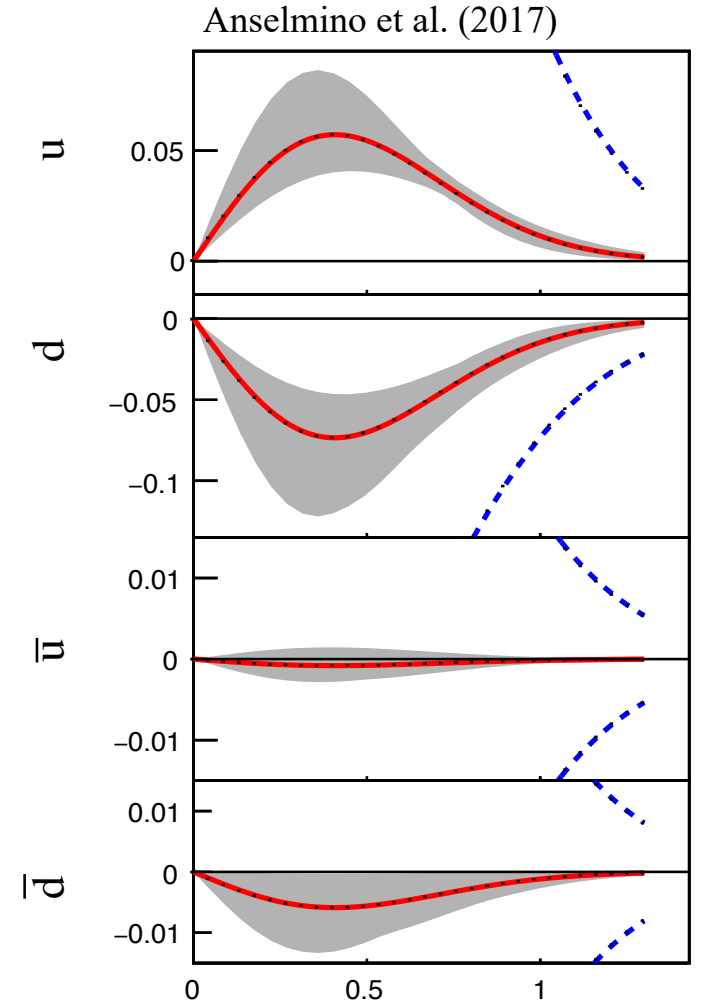
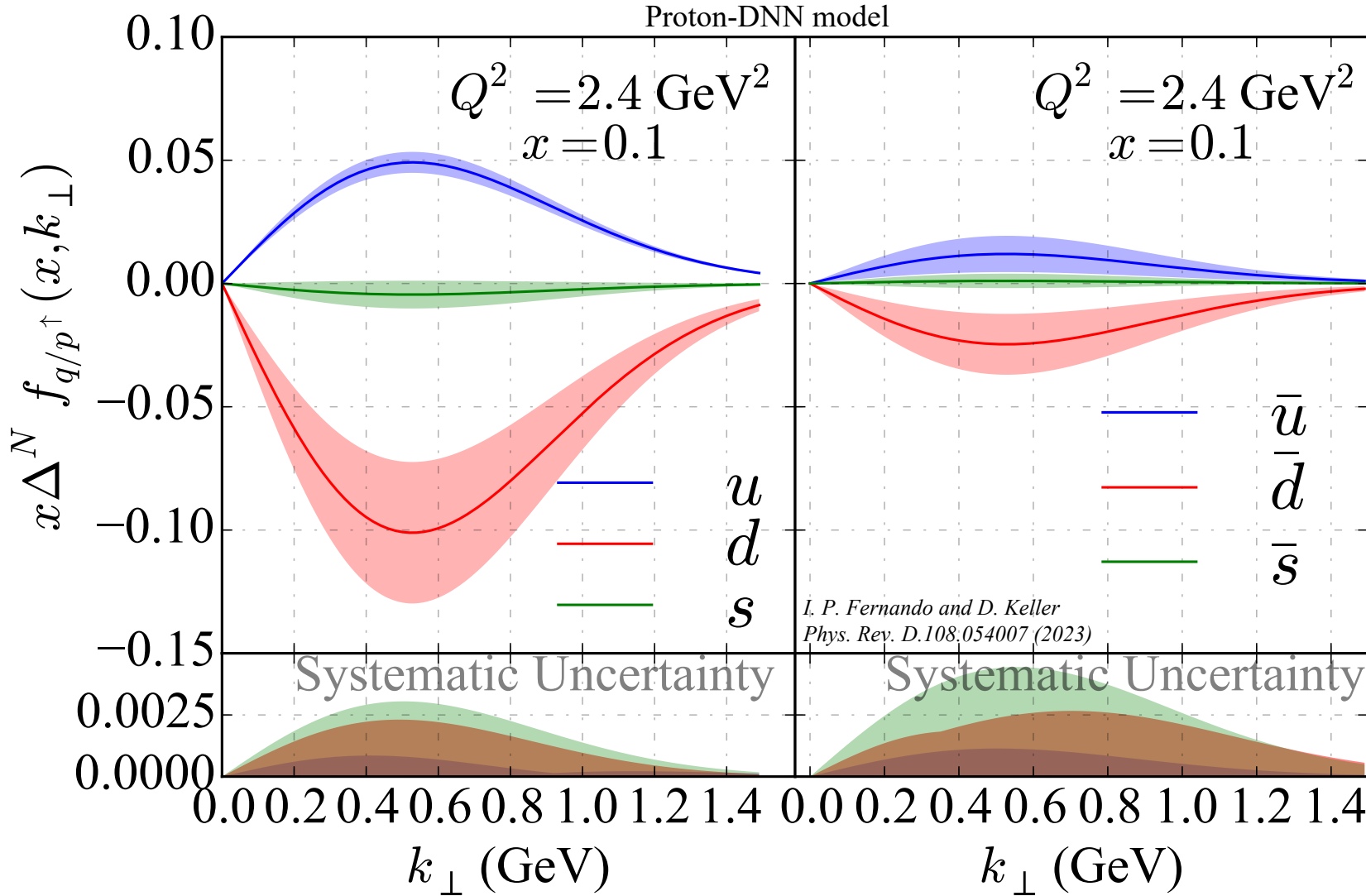
$$\Delta^N f_{q/p^\uparrow}(x, k_\perp)|_{\text{SIDIS}} = - \Delta^N f_{q/p^\uparrow}(x, k_\perp)|_{\text{DY}}$$

DNN Method: With Real data (Quality of the extraction)



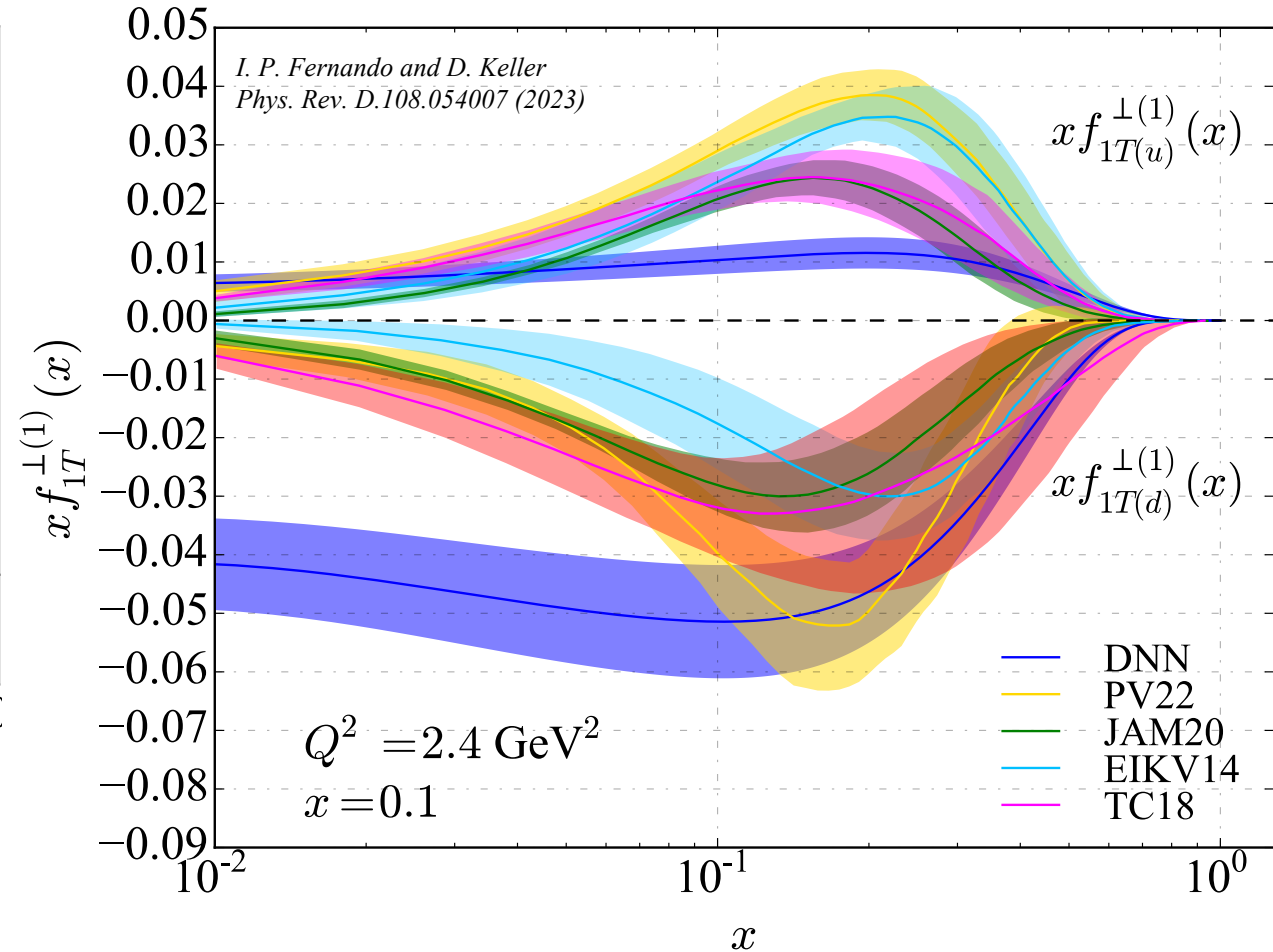
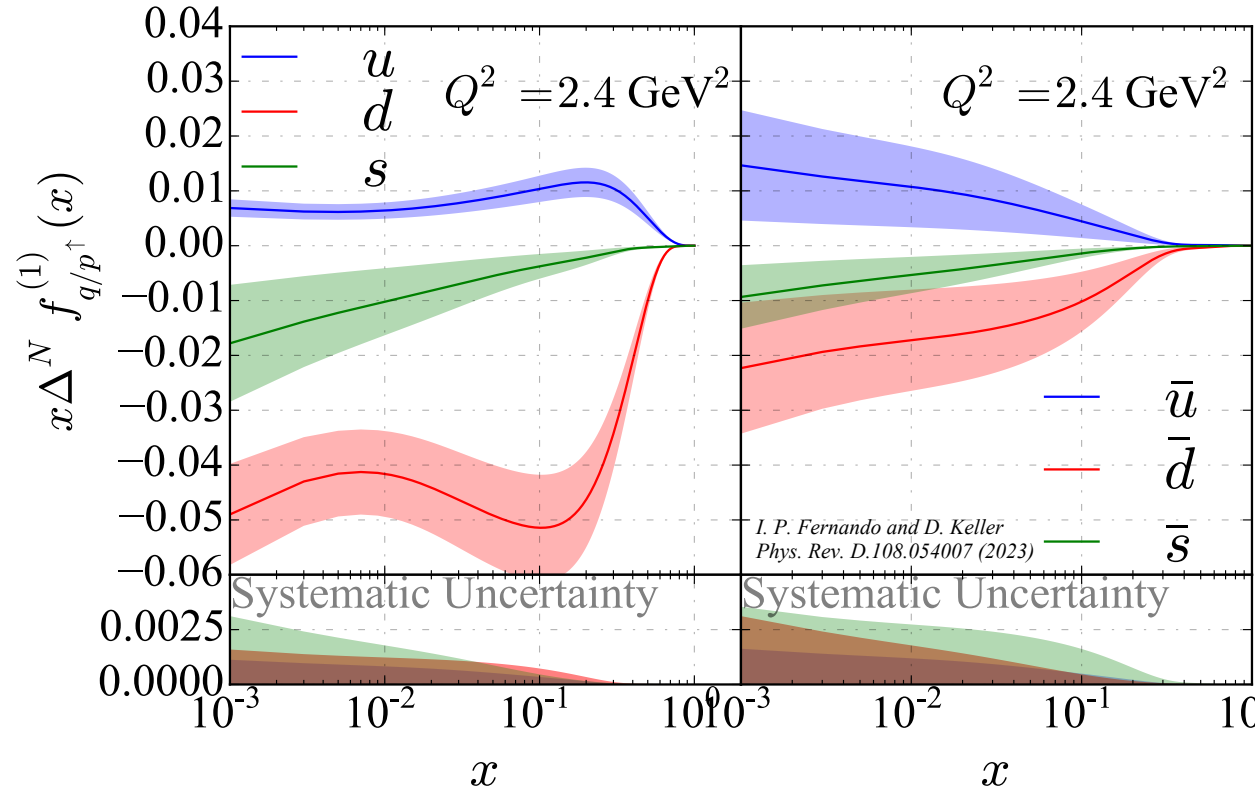
The qualitative improvement of the extracted Sivers functions for u (blue), d (red), and s (green) quarks at $x = 0.1$ and $Q^2 = 2.4 \text{ GeV}^2$ using the optimized proton-DNN model at the Second Iteration (solid-lines with dark-colored error bands with 68% CL), compared to the First Iteration (dashed-lines with light-colored error bands with 68% CL)

Sivers functions from the “Proton” DNN Model



Sivers 1st moments from the “Proton” Model

$$\Delta^N f_{q/p^\uparrow}^{(1)}(x) = \int d^2 k_\perp \frac{k_\perp}{4m_p} \Delta^N f_{q/p^\uparrow}(x, k_\perp) = -f_{1T}^{\perp(1)q}(x)$$

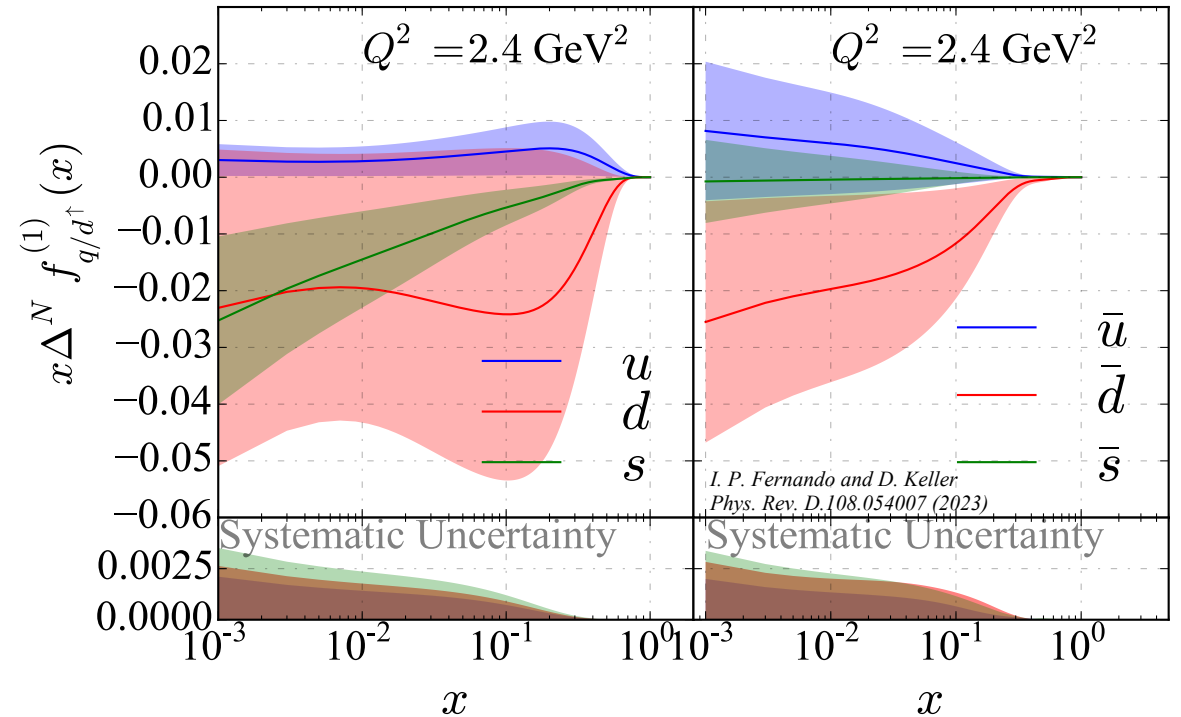
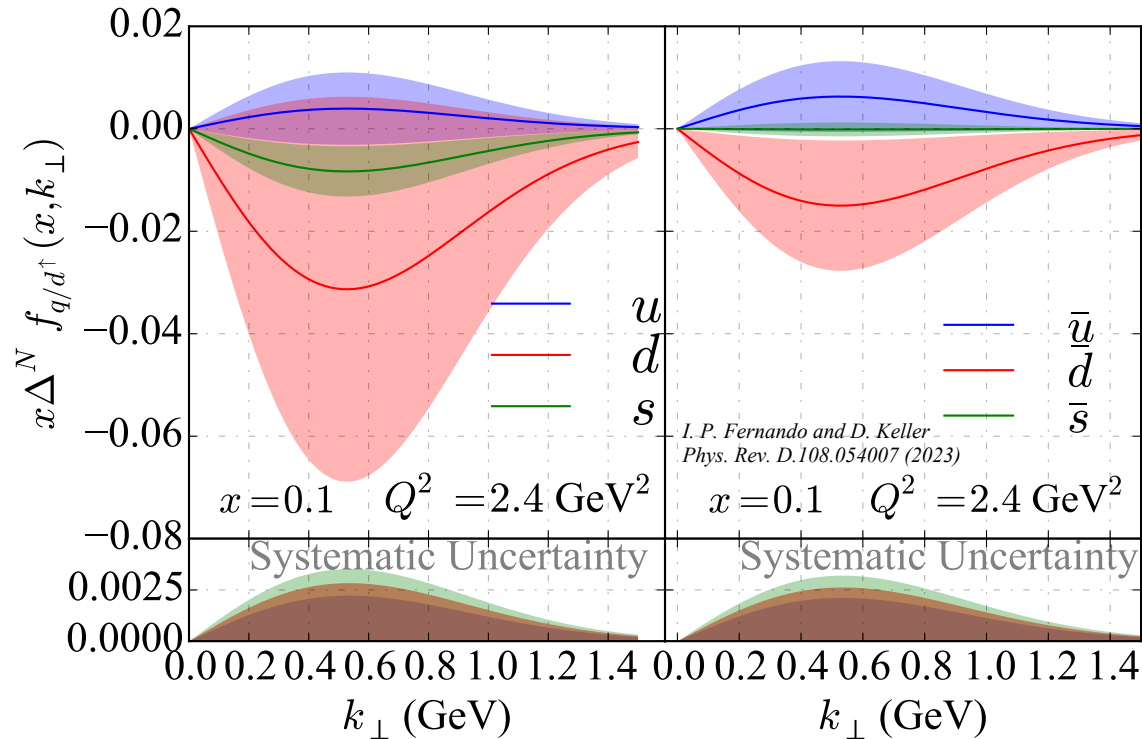


DNN Method: Results from the “Deuteron” Model

- Trained on COMPASS 2009 SIDIS data with Deuteron target.
- Did not imposed iso-spin symmetric conditions, or data cuts.

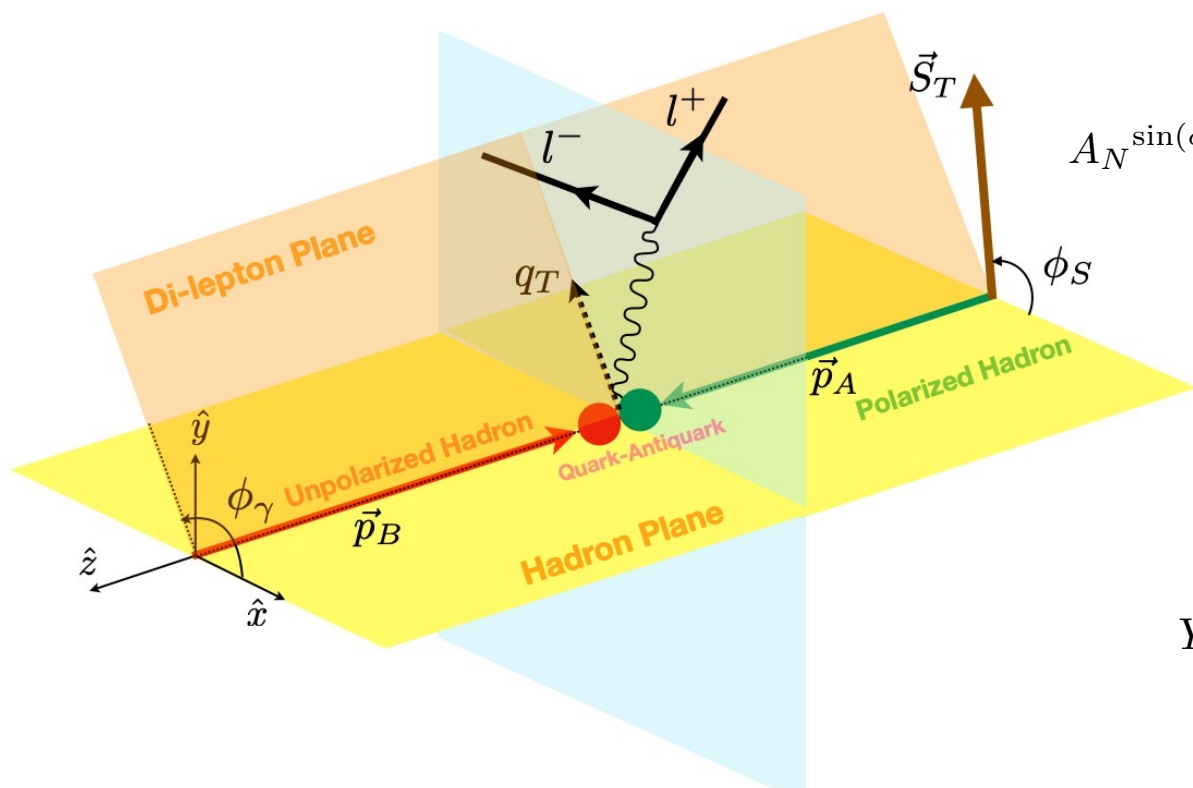
$$\cancel{f_{1T,u\leftarrow d}^\perp = f_{1T,d\leftarrow d}^\perp = \frac{f_{1T,u\leftarrow p}^\perp + f_{1T,d\leftarrow p}^\perp}{2}}$$

$$\Delta^N f_{q/p^\uparrow}^{(1)}(x) = \int d^2 k_\perp \frac{k_\perp}{4m_p} \Delta^N f_{q/p^\uparrow}(x, k_\perp) = -f_{1T}^{\perp(1)q}(x)$$



DNN Model Projections: DY

Based on Anselmino et al. (2017)



$$A_N^{\sin(\phi_\gamma - \phi_S)}(x_F, M, q_T) = \mathcal{B}_0(q_T, m_1) \frac{\sum_q \frac{e_q^2}{x_1 + x_2} \mathcal{N}_q(x_1) f_{q/A}(x_1) f_{\bar{q}/B}(x_2)}{\sum_q \frac{e_q^2}{x_1 + x_2} f_{q/A}(x_1) f_{\bar{q}/B}(x_2)}$$

$$\mathcal{B}_0(q_T, m_1) = \frac{q_T \sqrt{2e}}{m_1} \frac{Y_1(q_T, k_S, k_{\perp 2})}{Y_2(q_T, k_{\perp 1}, k_{\perp 2})}$$

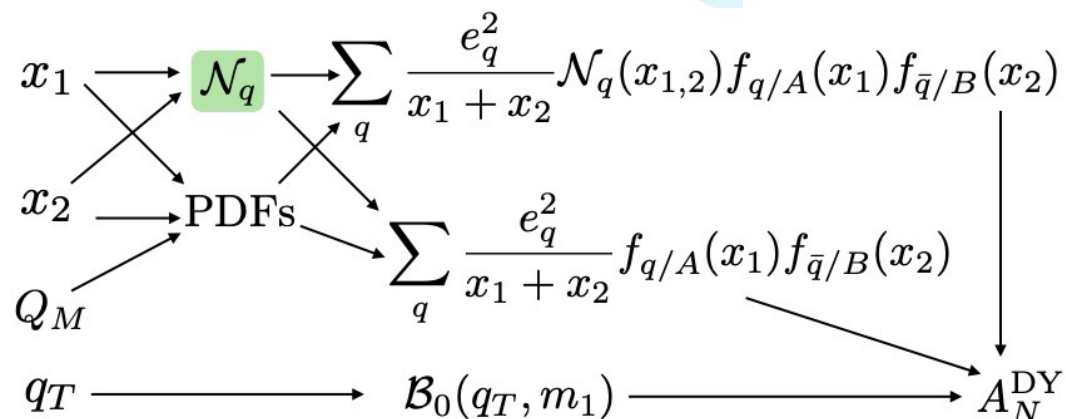
$$Y_1(q_T, k_S, k_{\perp 2}) = \left(\frac{\langle k_S^2 \rangle^2}{\langle k_{\perp 2}^2 \rangle (\langle k_S^2 \rangle + \langle k_{\perp 2}^2 \rangle)^2} \right) \times \exp \left(\frac{-q_T^2}{\langle k_S^2 \rangle + \langle k_{\perp 2}^2 \rangle} \right)$$

$$Y_2(q_T, k_{\perp 1}, k_{\perp 2}) = \left(\frac{1}{\langle k_{\perp 1}^2 \rangle + \langle k_{\perp 2}^2 \rangle} \right) \times \exp \left(\frac{-q_T^2}{\langle k_{\perp 1}^2 \rangle + \langle k_{\perp 2}^2 \rangle} \right)$$

$$\frac{1}{\langle k_S^2 \rangle} = \frac{1}{m_1^2} + \frac{1}{\langle k_{\perp 1}^2 \rangle}$$

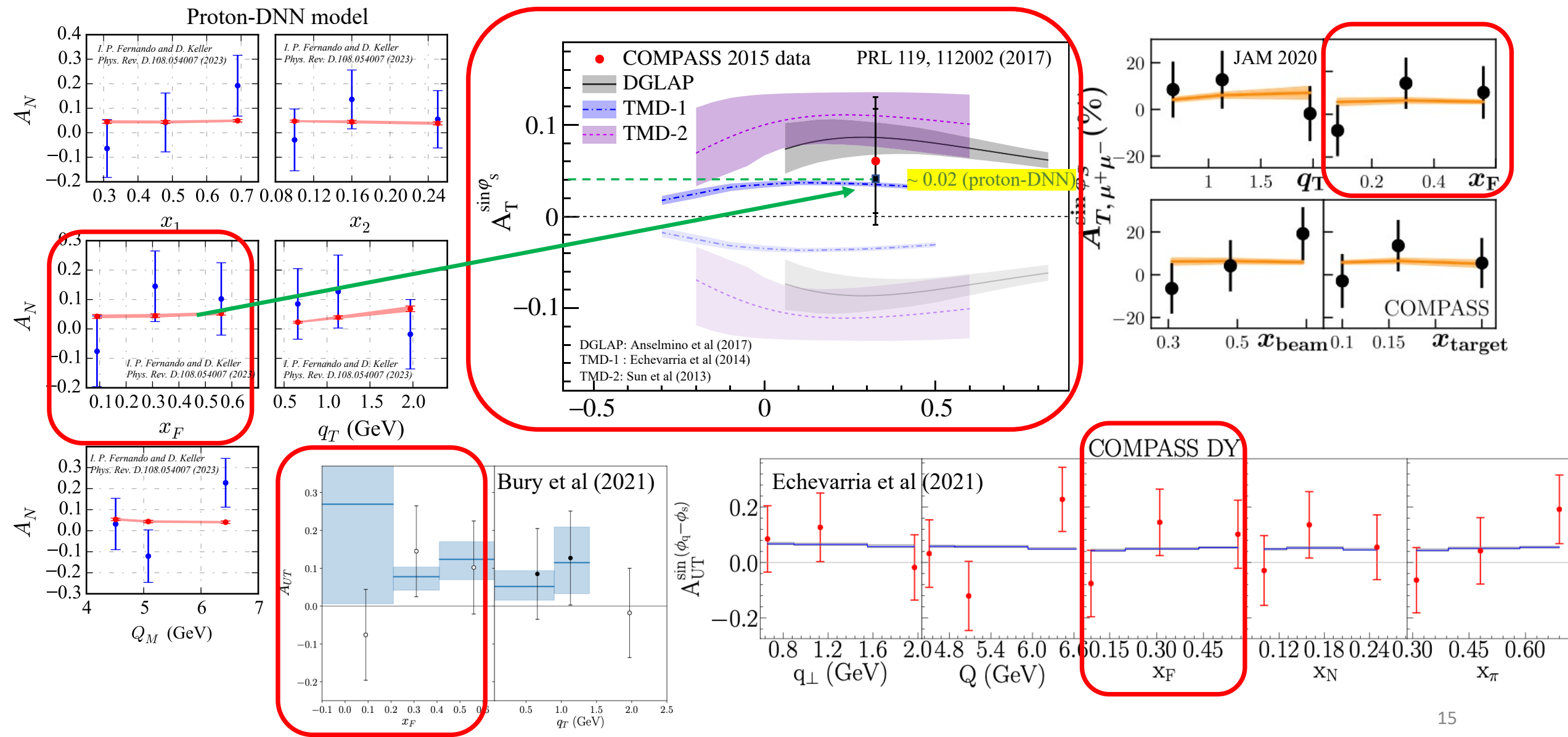
$$\langle k_{\perp 1}^2 \rangle = \langle k_{\perp 2}^2 \rangle = \langle k_{\perp}^2 \rangle = 0.25 \text{ GeV}^2$$

$$\Delta^N f_{q/p^\uparrow}(x, k_\perp)|_{\text{SIDIS}} = -\Delta^N f_{q/p^\uparrow}(x, k_\perp)|_{\text{DY}}$$



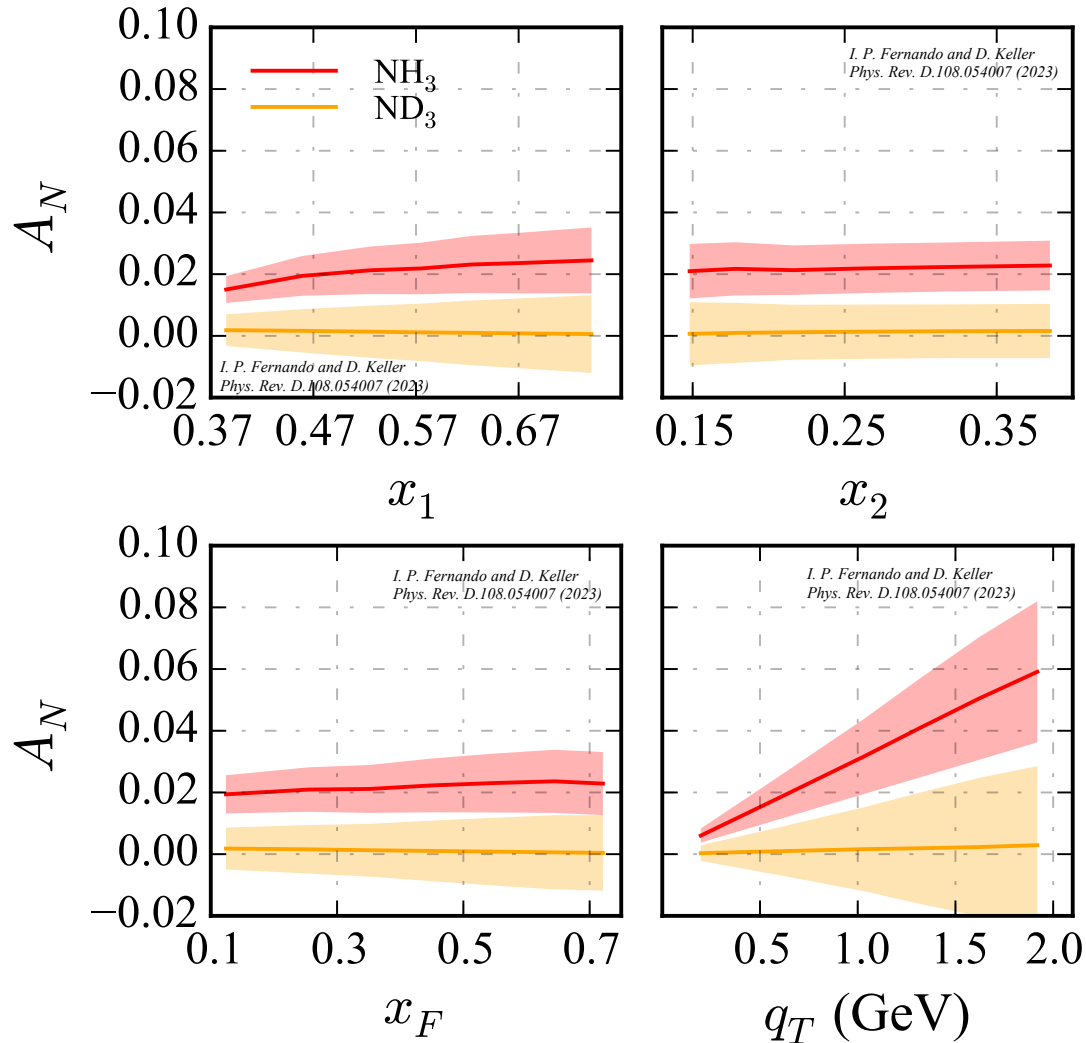
DNN Model Projections: DY

COMPASS 2017 DY Projections



DNN Model Projections: DY @ SpinQuest

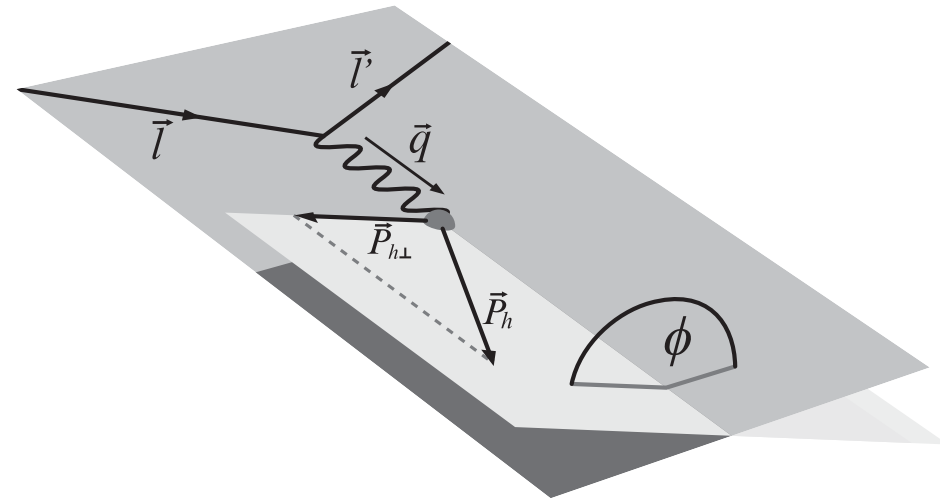
DNN Models



- SpinQuest (E1039) experiment at Fermilab is aiming to extract the Sivers function for the light-sea quarks.
- Unpolarized 120 GeV proton beam with polarized proton and deuteron targets (separately).
- Proton-DNN model predictions (Red)
Deuteron-DNN model predictions (Orange)

Boer-Mulders (BM) function

Boer Mulders function describes the polarized quark distributions inside an unpolarized hadron.



Based on the framework by Barone et al., PRD 91, 074019 (2015)

$$\frac{d\sigma}{dx_B dy dz_h dp_{hT} d\phi} = \frac{\pi\alpha^2}{Q^2 x_B y} \left((1 + (1 - y)^2) F_{UU} + 2(2 - y) \sqrt{1 - y} F_{UU}^{\cos\phi} \cos\phi + 2(1 - y) F_{UU}^{\cos 2\phi} \cos 2\phi \right)$$

Twist-3

$$F_{UU}^{\cos\phi, h}|_{\text{BM}} = \sum_q e_q^2 x \int d^2\mathbf{k}_\perp \frac{k_\perp p_{hT}^2 - z(\mathbf{k}_\perp \cdot \mathbf{p}_{hT})}{Q p_{hT}^2} \frac{k_\perp}{m_p} h_1^\perp(x, k_\perp) \frac{2p_\perp}{z m_h} H_1^\perp(z, p_\perp) + F_{UU}^{\cos\phi}|_{\text{Cahn}} = -2 \sum_q e_q^2 x \int d^2\mathbf{k}_\perp \frac{(\mathbf{k}_\perp \cdot \mathbf{p}_{hT})}{Q p_{hT}} f_q(x, k_\perp) D_q(z, p_\perp)$$

Twist-4

$$F_{UU}^{\cos 2\phi, h}|_{\text{BM}} = - \sum_q e_q^2 x \int d^2\mathbf{k}_\perp \frac{p_{hT}(\mathbf{k}_\perp \cdot \mathbf{p}_{hT}) + z(\mathbf{k}_\perp^2 \mathbf{p}_{hT} - 2(\mathbf{k}_\perp \cdot \mathbf{p}_{hT}))}{2k_\perp p_{hT}^2} \frac{k_\perp}{m_p} h_1^\perp(x, k_\perp) \frac{2p_\perp}{z m_h} H_1^\perp(z, p_\perp) + F_{UU}^{\cos\phi}|_{\text{Cahn}} = 2 \sum_q e_q^2 x \int d^2\mathbf{k}_\perp \frac{2(\mathbf{k}_\perp \cdot \mathbf{p}_{hT}) - \mathbf{k}_\perp^2 \mathbf{p}_{hT}^2}{Q^2 p_{hT}^2} f_q(x, k_\perp) D_q(z, p_\perp)$$

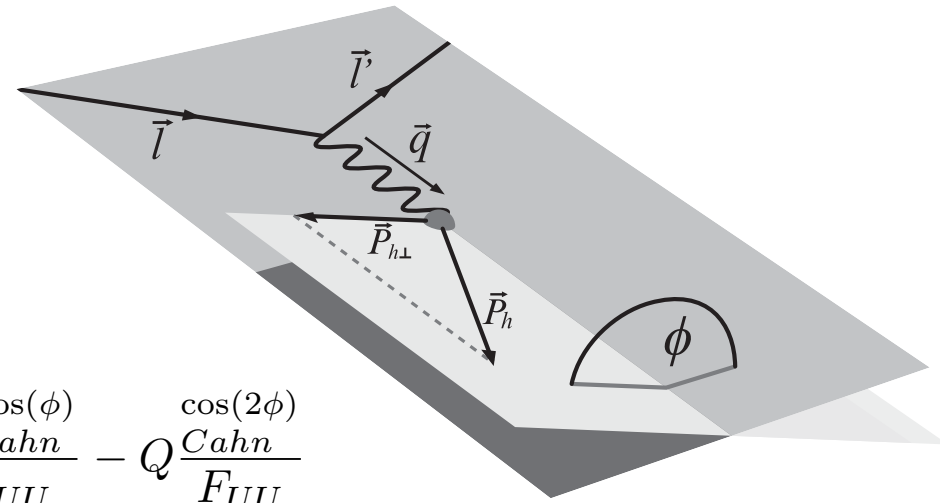
Boer-Mulders (BM) function

$$A^{\cos(\phi)} = A_{BM}^{\cos(\phi)} + A_{Cahn}^{\cos(\phi)}$$

$$A^{\cos(2\phi)} = A_{BM}^{\cos(2\phi)} + A_{Cahn}^{\cos(2\phi)}$$

An exploratory rearrangement

$$A^{\cos(\phi)} - Q A^{\cos(2\phi)} = \frac{F_{BM}^{\cos(\phi)}}{F_{UU}} - Q \frac{F_{BM}^{\cos(2\phi)}}{F_{UU}} + \frac{F_{Cahn}^{\cos(\phi)}}{F_{UU}} - Q \frac{F_{Cahn}^{\cos(2\phi)}}{F_{UU}}$$



$$h_1^\perp(x, k_\perp) = \mathcal{N}_q(x) h(k_\perp) f_q(x, k_\perp; Q^2)$$

$$\mathcal{N}_q(x) = N_q x^\alpha (1-x)^\beta \frac{(\alpha + \beta)^{(\alpha + \beta)}}{\alpha^\alpha \beta^\beta}$$

$$f_q(x, k_\perp; Q^2) = f_q(x; Q^2) \frac{e^{-k_\perp^2 / \langle k_\perp^2 \rangle}}{\pi \langle k_\perp^2 \rangle}$$

$$h(k_\perp) = \sqrt{2e} \frac{k_\perp}{m_{BM}} e^{-k_\perp^2 / m_{BM}^2}$$

$$\Delta D_{h/q\uparrow}(z, p_\perp) = \mathcal{N}_q^C(z) h^C(k_\perp) D_{h/q}(z, p_\perp; Q^2)$$

$$\mathcal{N}_q^C(x) = N_q^C z^\gamma (1-z)^\delta \frac{(\gamma + \delta)^{(\gamma + \delta)}}{\gamma^\gamma \delta^\delta}$$

$$D_{h/q}(z, p_\perp; Q^2) = D_{h/q}(z; Q^2) \frac{e^{-p_\perp^2 / \langle p_\perp^2 \rangle}}{\pi \langle p_\perp^2 \rangle}$$

$$h^C(p_\perp) = \sqrt{2e} \frac{p_\perp}{m_C} e^{-p_\perp^2 / m_C^2}$$

From the framework by Barone et al., PRD 91, 074019 (2015)

$$\begin{aligned} \frac{F_{BM}^{\cos(\phi)}}{F_{UU}} - Q \frac{F_{BM}^{\cos(2\phi)}}{F_{UU}} &= \frac{2(2-y)}{[1+(1-y)^2]} \left(\frac{e p_{hT}}{m_{BM} m_c Q} \right) \frac{\langle p_{hT}^2 \rangle}{\langle p_{hT}^2 \rangle_{BM}^3} \left(\frac{\sum_q e^q \mathcal{N}_q(x) f_{q/p}(x) \mathcal{N}_q^C(x) D_{h/q}(z)}{\sum_q e^q f_{q/p}(x) D_{h/q}(z)} \right) \exp \left(\frac{p_{hT}^2}{\langle p_{hT}^2 \rangle} - \frac{p_{hT}^2}{\langle p_{hT}^2 \rangle_{BM}} \right) \\ &\quad \left(\frac{2\sqrt{1-y}}{\langle p_{hT}^2 \rangle_{BM}} \frac{\langle k_\perp^2 \rangle_{BM}^2 \langle p_\perp^2 \rangle_C^2}{\langle k_\perp^2 \rangle^2 \langle p_\perp^2 \rangle^2} (z^2 \langle k_\perp^2 \rangle_{BM}^2 \langle (p_{hT}^2 - \langle p_{hT}^2 \rangle_{BM}^2) + \langle p_\perp^2 \rangle_C^2 \langle p_{hT}^2 \rangle_{BM}^2) + Q^2 p_{hT} \frac{z \langle k_\perp^2 \rangle_{BM}^2 \langle p_\perp^2 \rangle_C^2}{\langle k_\perp^2 \rangle^2 \langle p_\perp^2 \rangle^2} \right) \end{aligned}$$

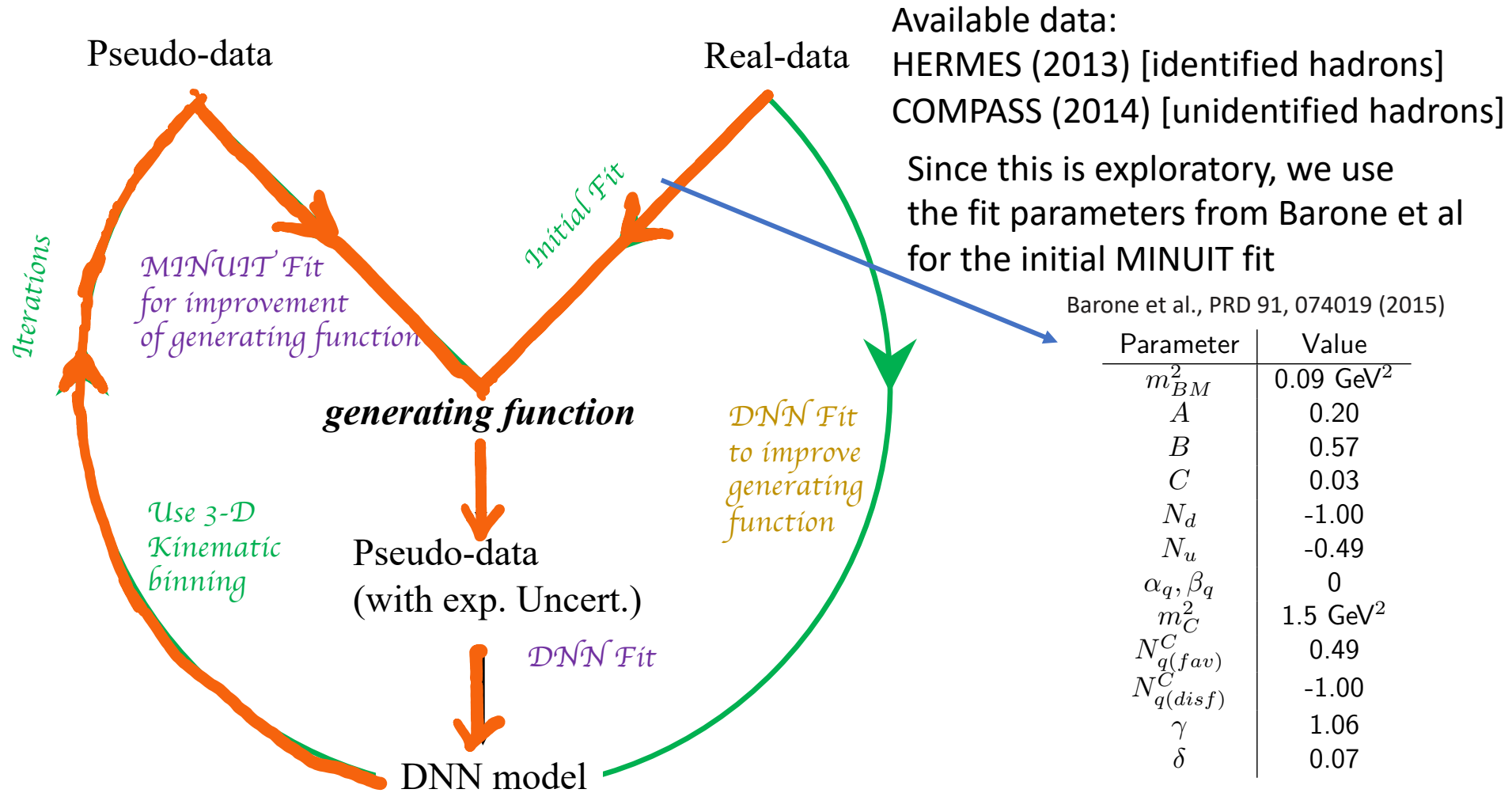
Method Testing

$$\langle p_{\perp}^2 \rangle_C = \frac{m_C^2 \langle p_{\perp}^2 \rangle}{\langle p_{\perp}^2 \rangle + m_C^2}$$

$$\langle p_{\perp}^2 \rangle = A + z^2 B$$

$$\langle k_{\perp}^2 \rangle = C$$

$$\sigma_q(x, k_\perp) = \sqrt{\frac{\sum_i \left(\Delta^N f_{q/p^\uparrow}^{(i)} - \Delta^N f_{q/p^\uparrow}^{(\text{mean})} \right)^2}{N}}.$$



The “DNN Method” for extracting BM

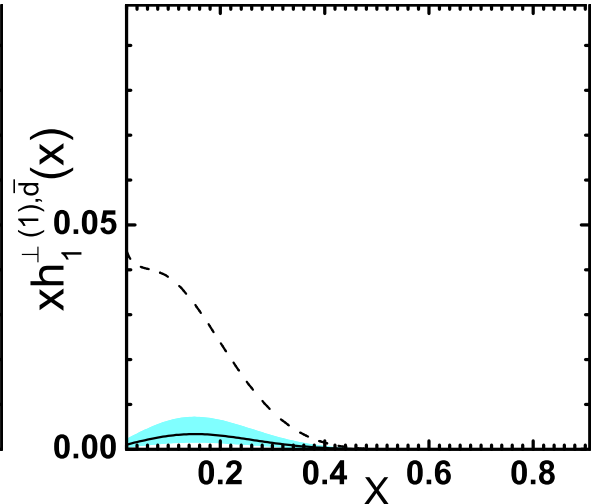
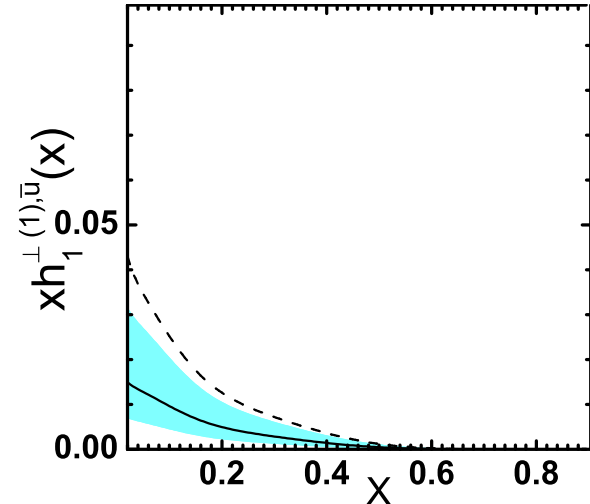
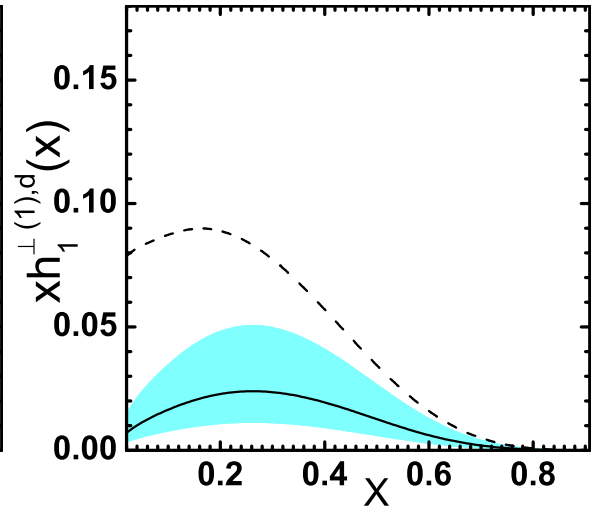
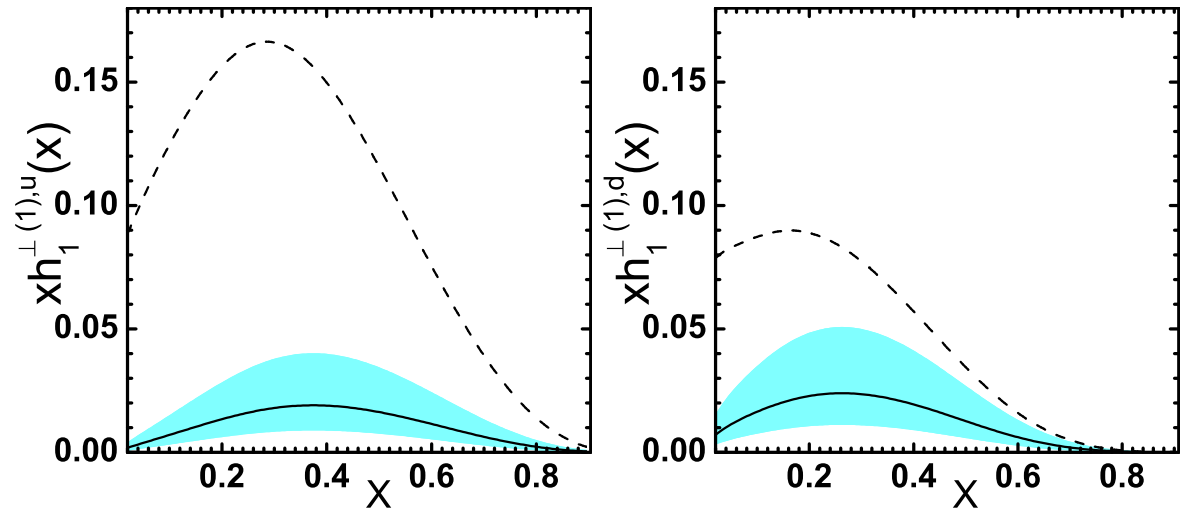
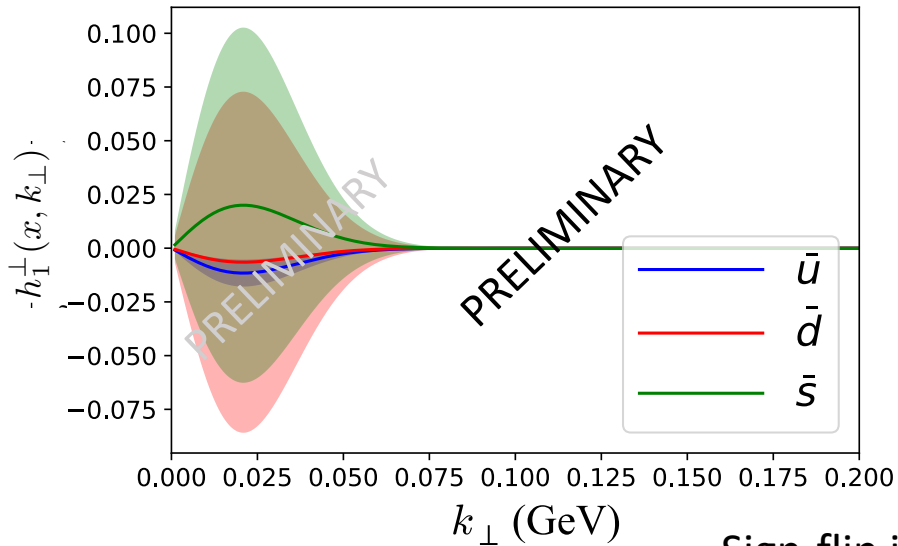
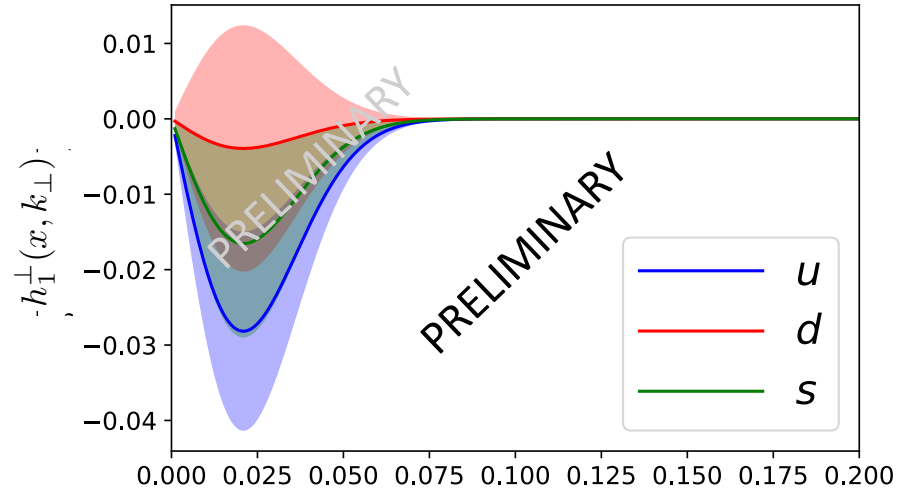
- All quark flavors in SU(3) are considered in a more generic fashion.
- As this is exploratory, the fit results for the coefficients not related to $\mathcal{N}_q(x)$ are fixed by the Barone et al. (2015) fit results, for the purpose of extracting the “generating function”.
- Twist-4 Cahn term was not ignored.
- A linear combination of $\cos(\phi)$ and $\cos(2\phi)$ asymmetries were assumed.
- Training in-progress for two separate DNN models: proton target and deuteron target
- Prospective plan: use inputs x , Q^2 (for evolution), and target type

Boer-Mulders (BM) function

Very very preliminary DNN fits
to HERMES 'proton' SIDIS data

Zhun Lu and Ivan Schmidt PhysRevD.81.034023 (2010)

From the fits to DY data



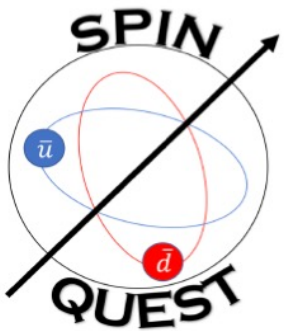
Sign-flip is expected from
QCD gauge invariance as BM is T-odd

Summary & Outlook

- We proposed a **new method** for performing global fits to extract TMDs.
- Our method is based on the integration of AI and the use of generating function to ensure the accuracy and precision.
- We have successfully tested our method with pseudo-data, also a dedicated systematic study.
- We chose Deep Neural Net (DNN) to incorporate all x-dependent features of $\mathcal{N}_q(x)$.
- We performed global fit with experimental data: Separately on polarized SIDIS with Proton target and Deuteron target and obtained reasonably well description and extracted the Siverson functions for all light quark flavors in SU(3).
- We projected SIDIS and DY Siverson asymmetries: for existing (as a validation check) and upcoming experiments (such as SpinQuest).
- Currently working on Boer-Mulders function extraction...

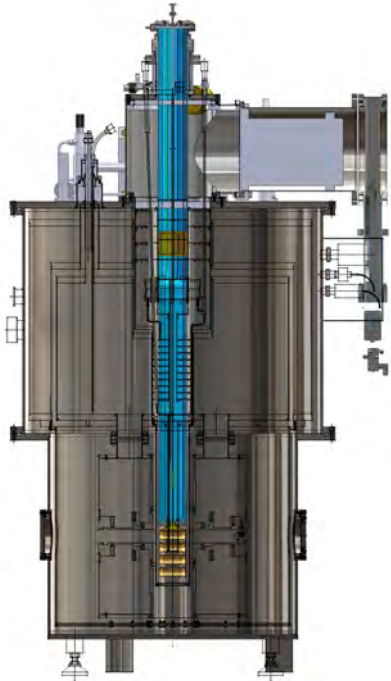
Next:

- Performing simultaneous global fit with SIDIS, DY, W^+/W^- data including a study of TMD evolution with DNN techniques to extract the unpolarized TMD using DNN.
- Applying the “DNN method” to extract other TMDs such as Transversity, Boer-Mulders function, as well as Spin-1 TMDs.



SpinQuest (E1039) Experiment at Fermilab

➤ Measurement of 'sea' quark Sivers function



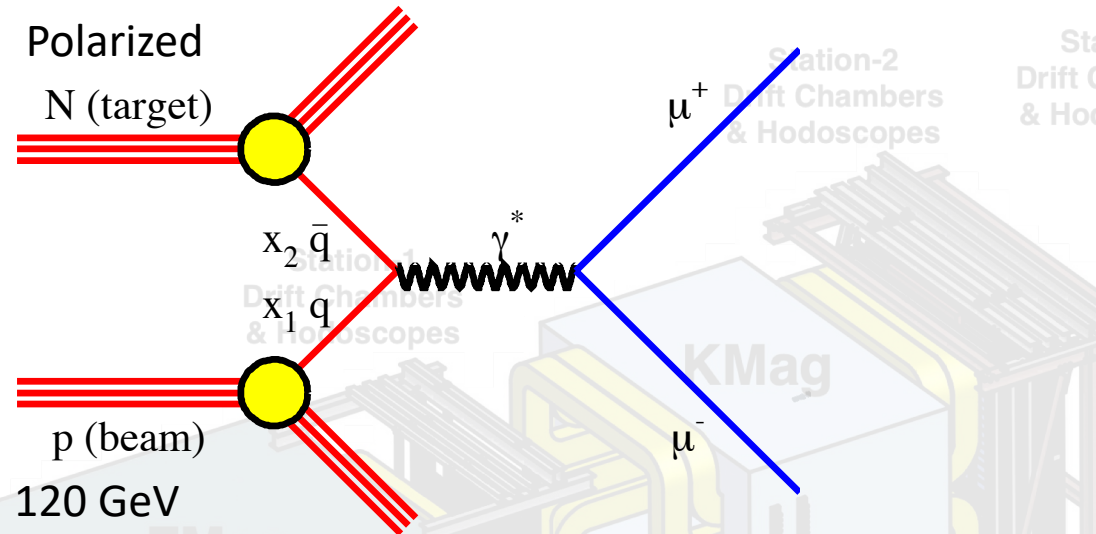
LANL-UVA

Polarized Target

<https://spinqwest.fnal.gov/>

<http://twist.phys.virginia.edu/E1039/>

$$pp \uparrow (d^\uparrow) \rightarrow \mu^+ \mu^- X, 4 < M_{\mu\mu} < 9 \text{ GeV}$$

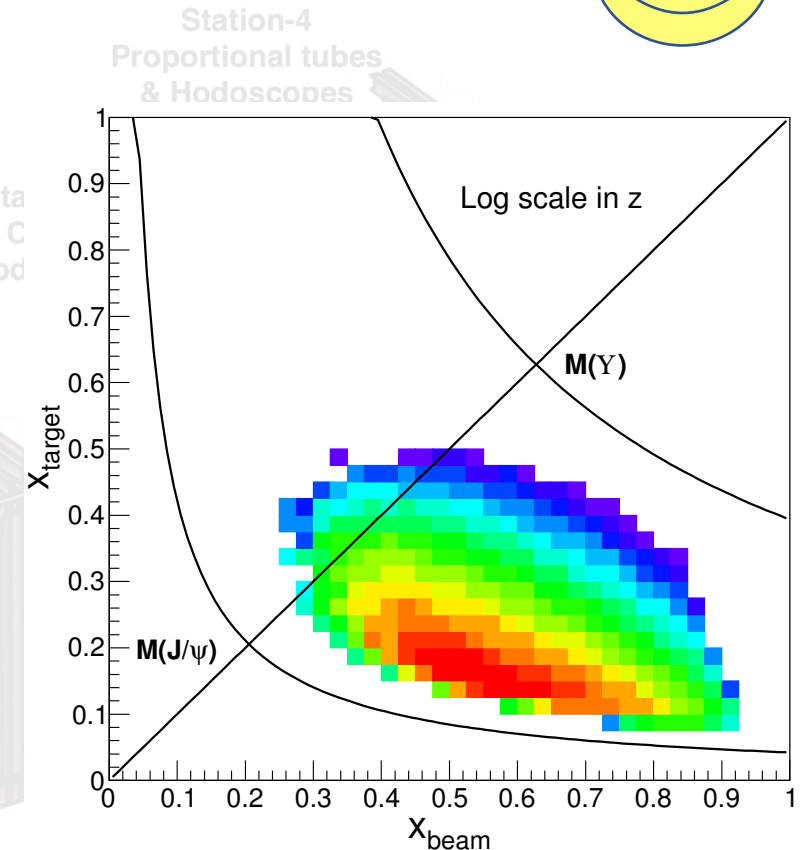


$$\frac{d\sigma}{dx_1 dx_2} = \frac{4\pi\alpha^2}{9sx_1 x_2} \sum_i e_i^2 (q_i^B(x_1, Q^2) \bar{q}_i^T(x_2, Q^2) + \bar{q}_i^B(x_1, Q^2) q_i^T(x_2, Q^2))$$

Please Join The Effort

Dustin Keller (dustin@virginia.edu)[Spokesperson]

Kun Liu (liuk@lanl.gov)[Spokesperson]]



Highest beam intensity on a polarized target ever!

Thank you

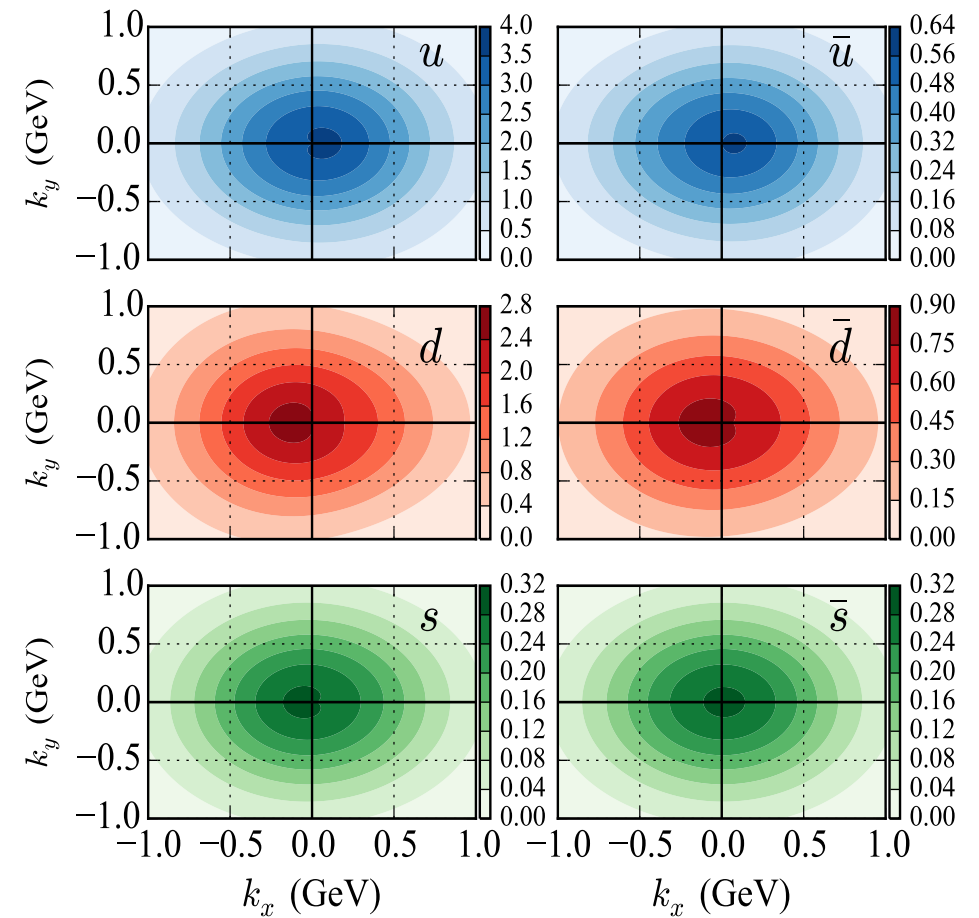


UNIVERSITY
of VIRGINIA



U.S. DEPARTMENT OF
ENERGY

Office of
Science



This work is supported by DOE contract DE-FG02-96ER40950

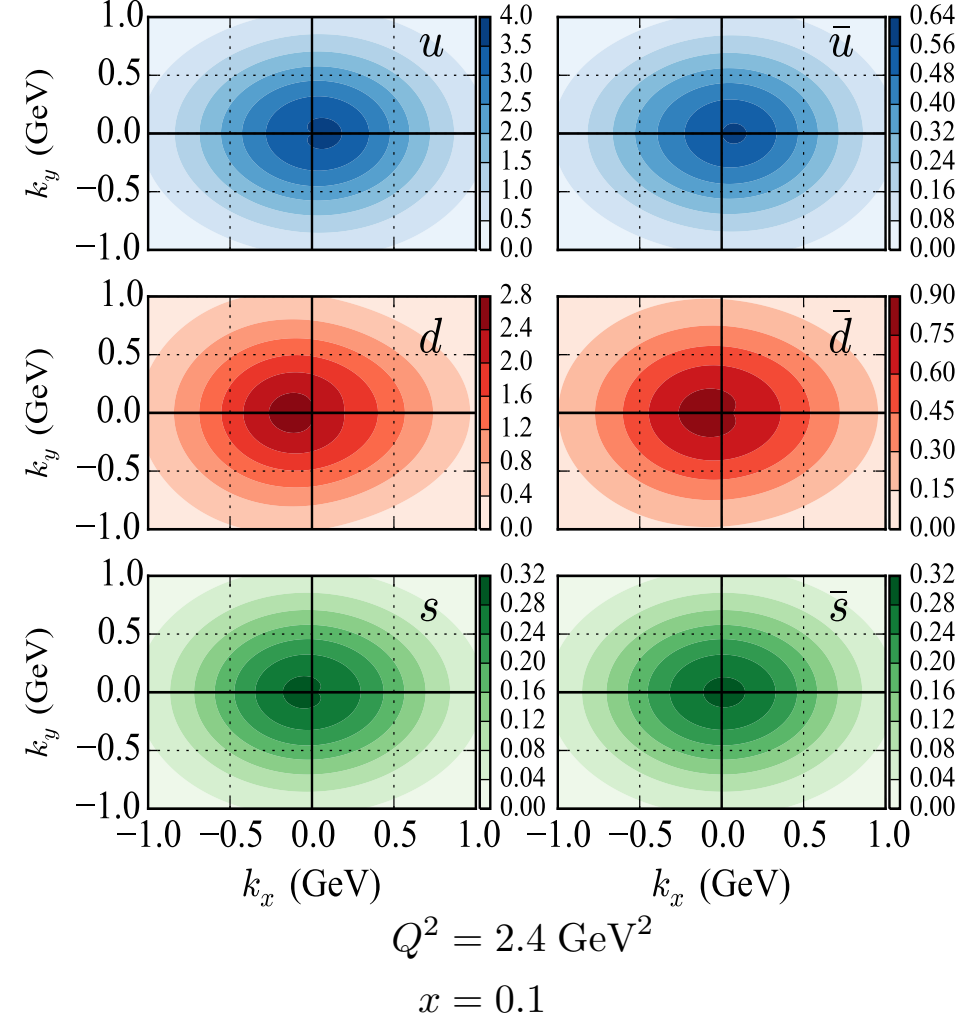
Backup Slides

3D Tomography from the “Proton” DNN Model

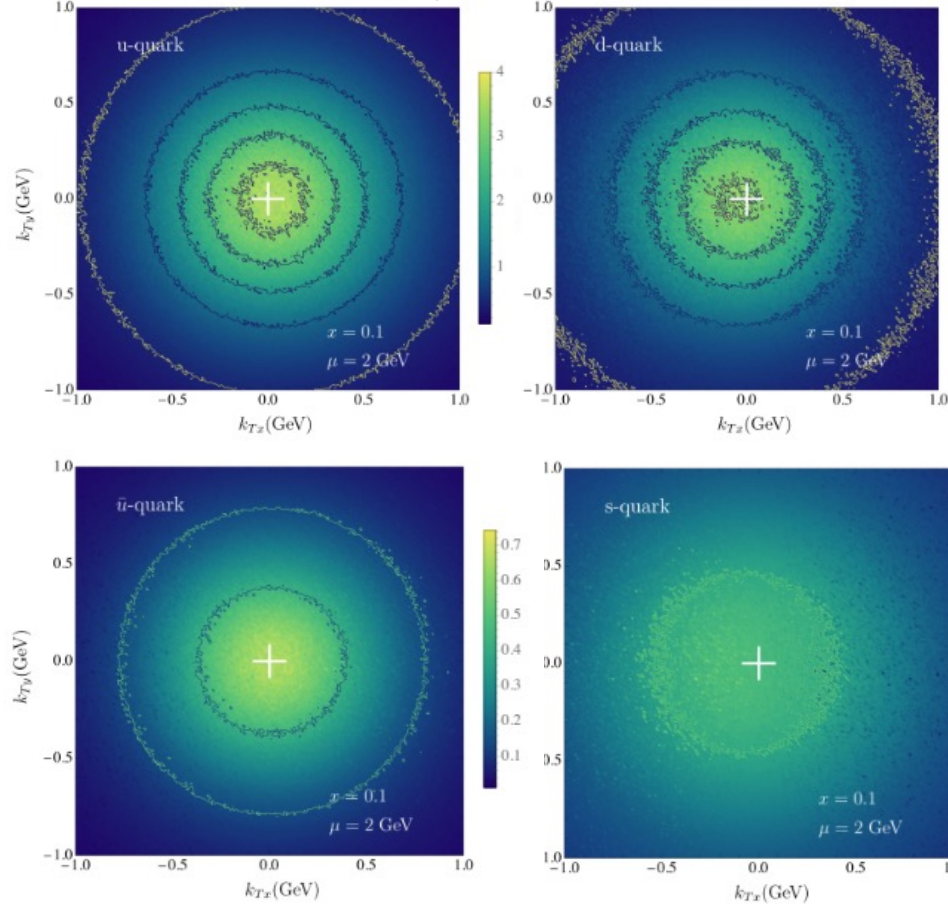
$$\rho_{p\uparrow}^a(x, k_x, k_y; Q^2) = f_1^a(x, k_\perp^2; Q^2) - \frac{k_x}{m_p} f_{1T}^a(x, k_\perp^2; Q^2)$$

I. P. Fernando and D. Keller
Phys. Rev. D.108.054007 (2023)

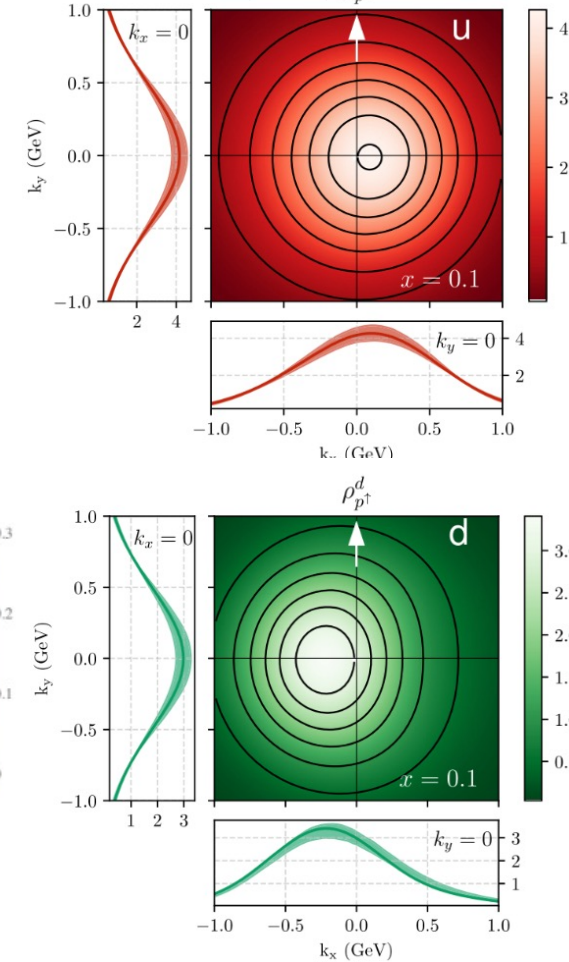
Proton-DNN model



Bury et al (2021)

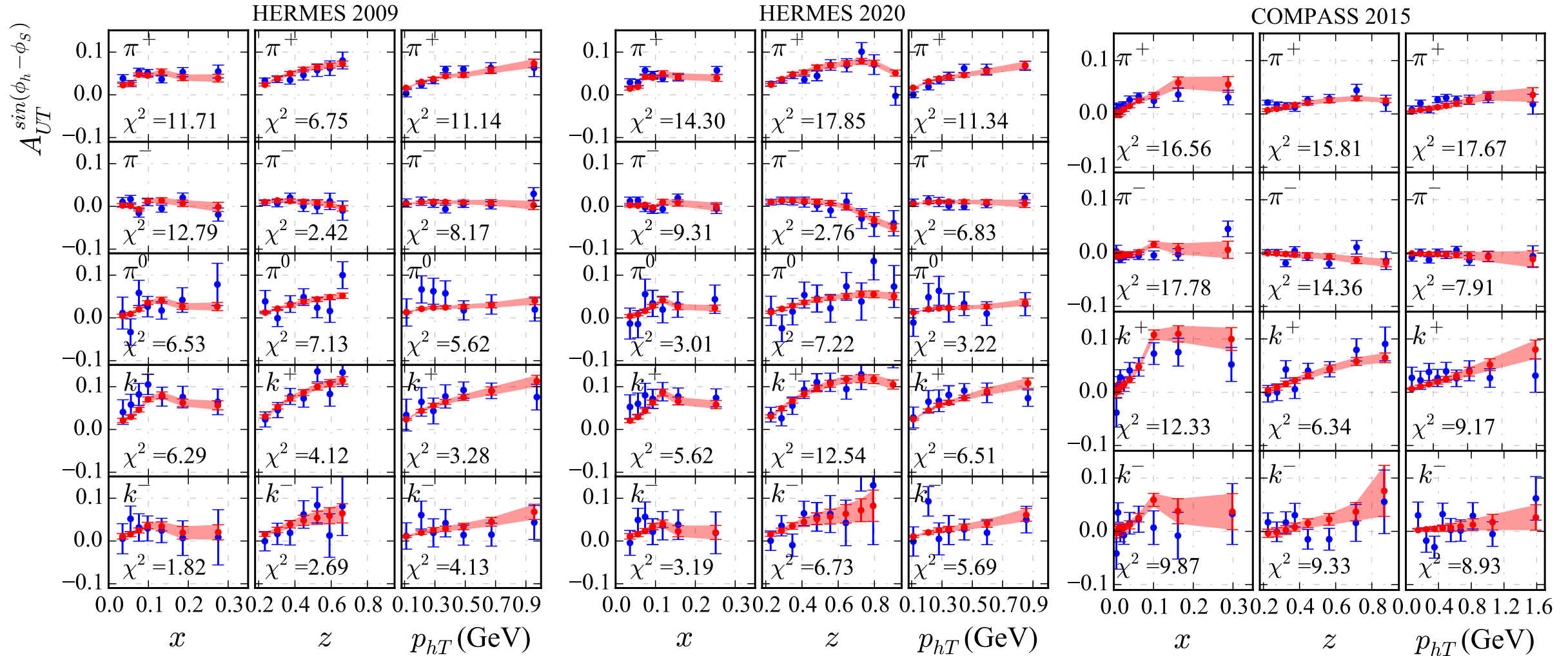


A. Bacchetta et al (2021) $\rho_{p\uparrow}^u$



Proton DNN Fit Results

*I. P. Fernando and D. Keller
Phys. Rev. D.108.054007 (2023)*



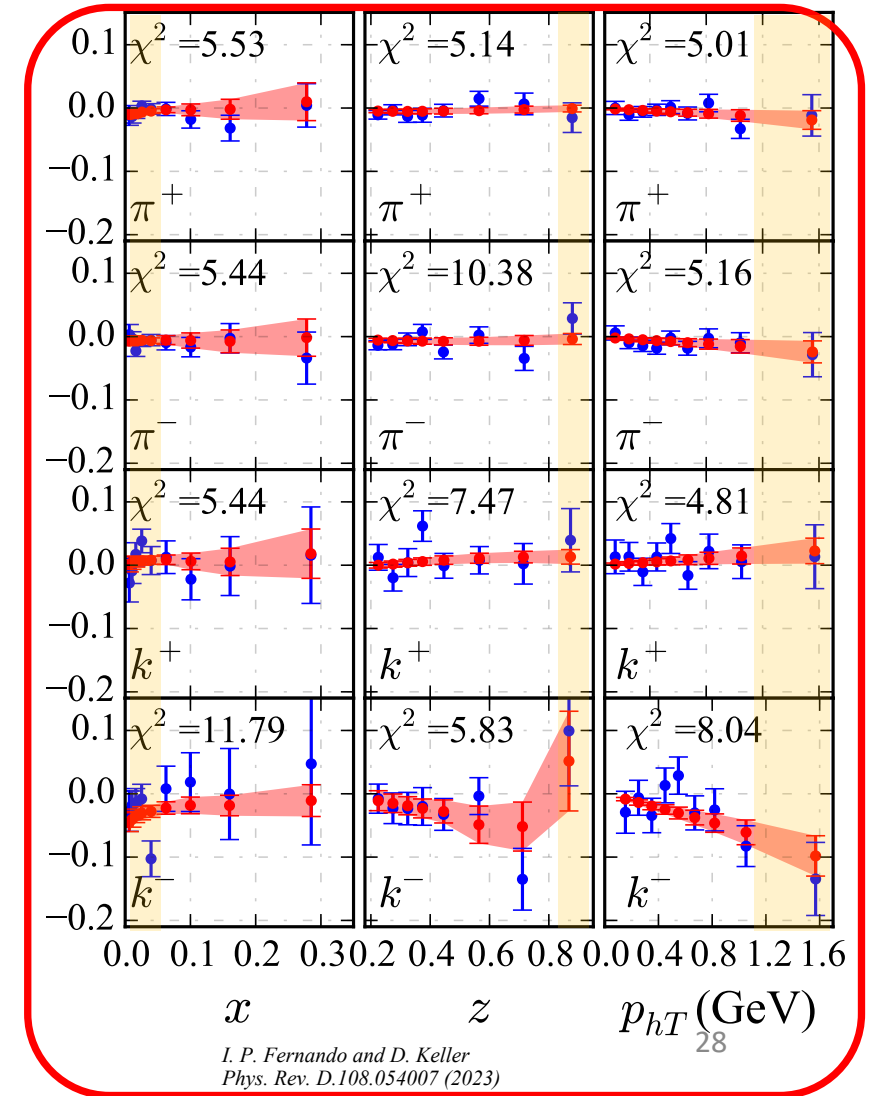
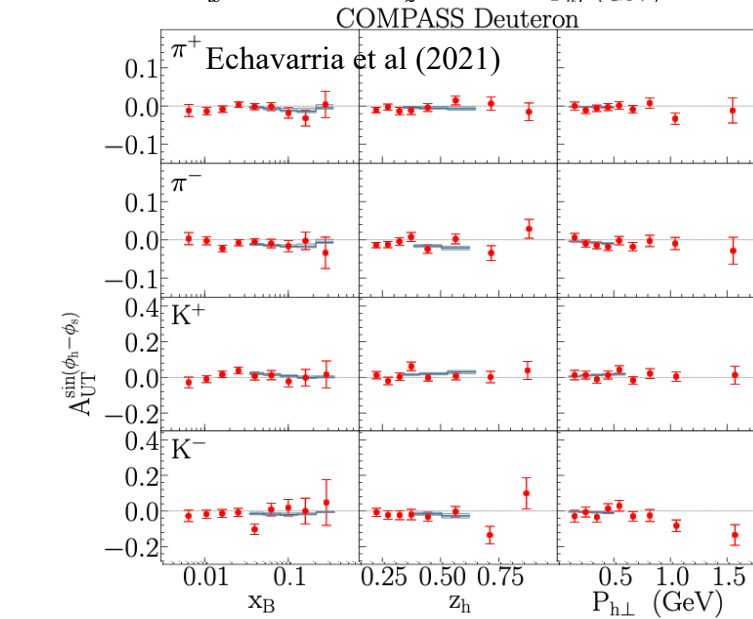
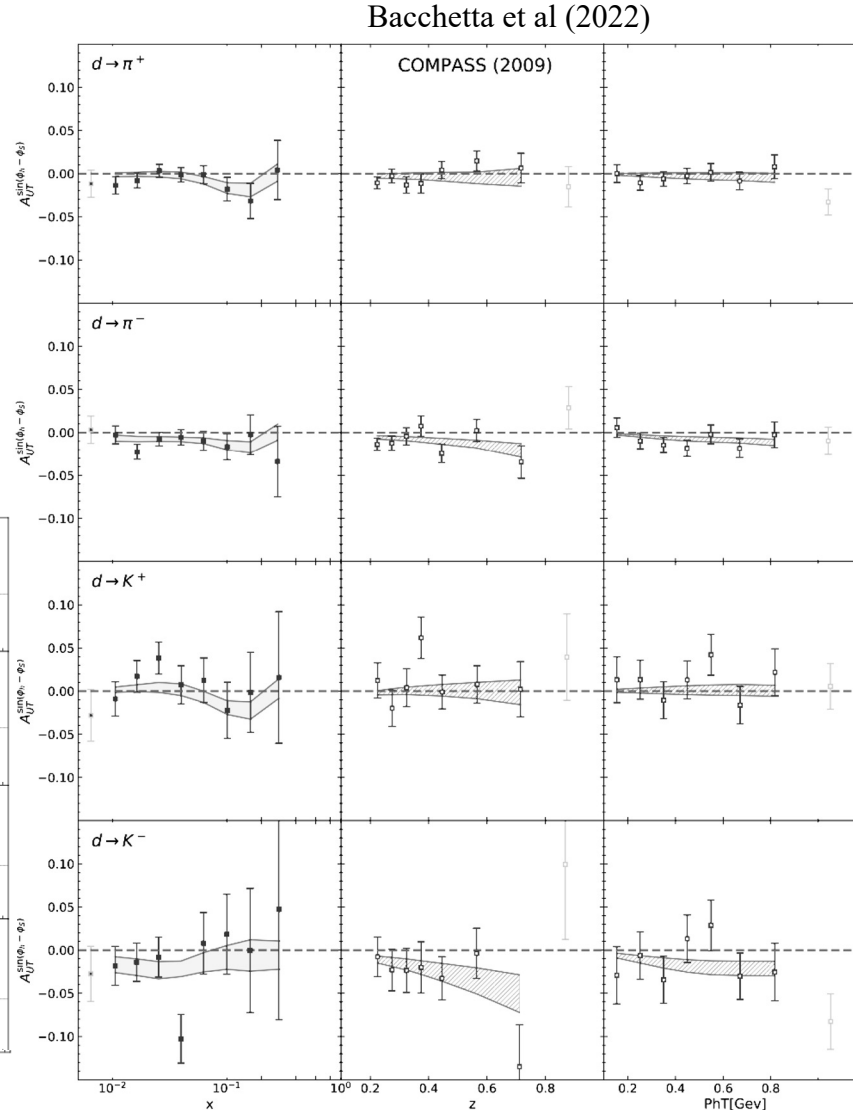
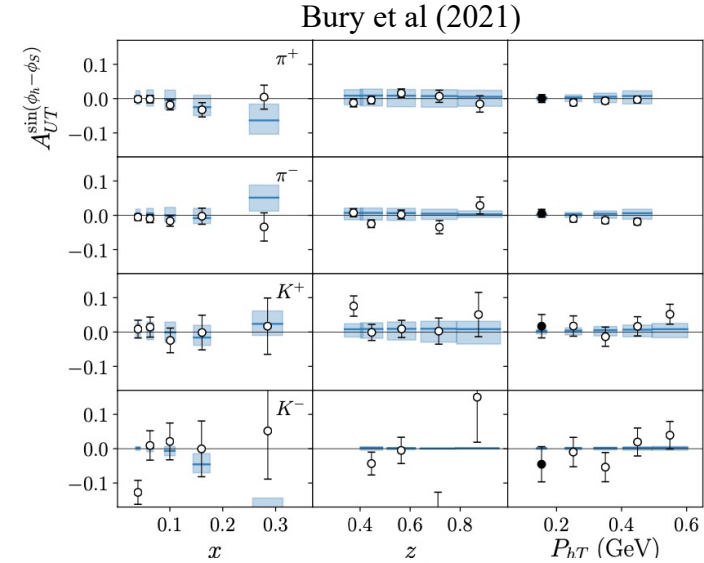
- All data points are well-described by the proton-DNN model.
- No kinematic cuts were implemented.

Calculated $\chi_{\text{total}}^2/N_{\text{pt}} = 1.04$

Deuteron DNN Fit Results

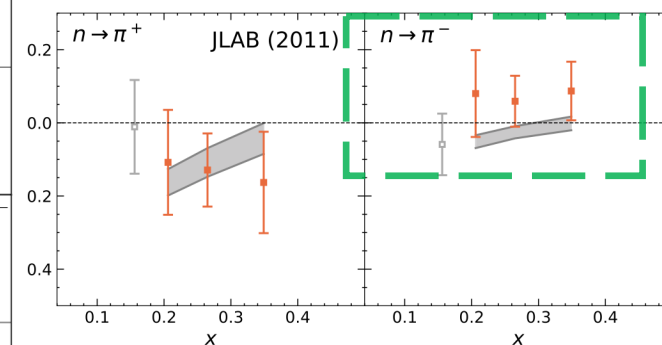
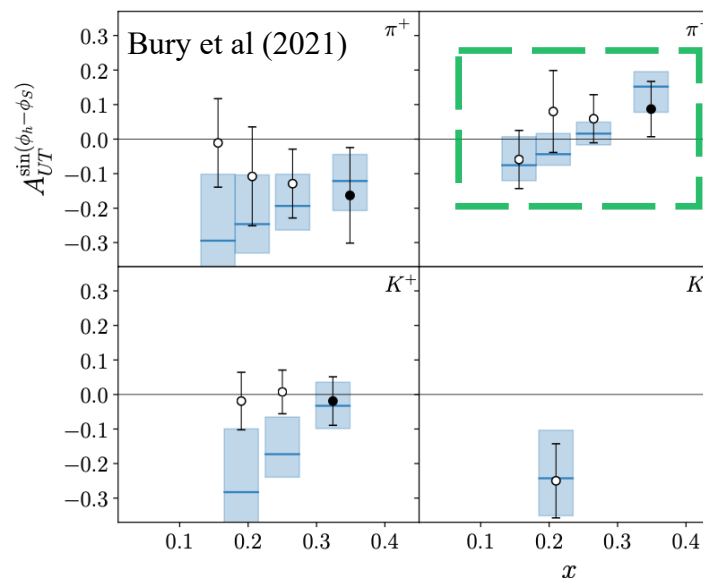
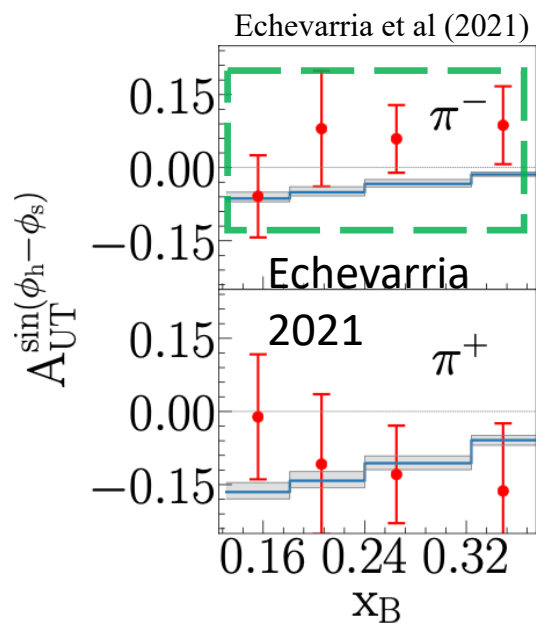
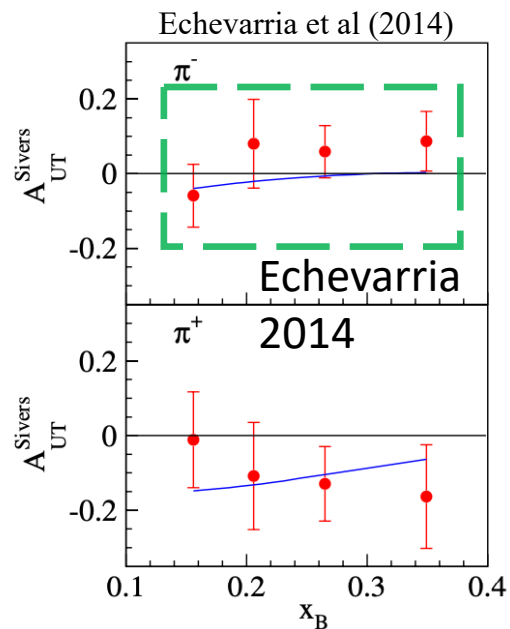
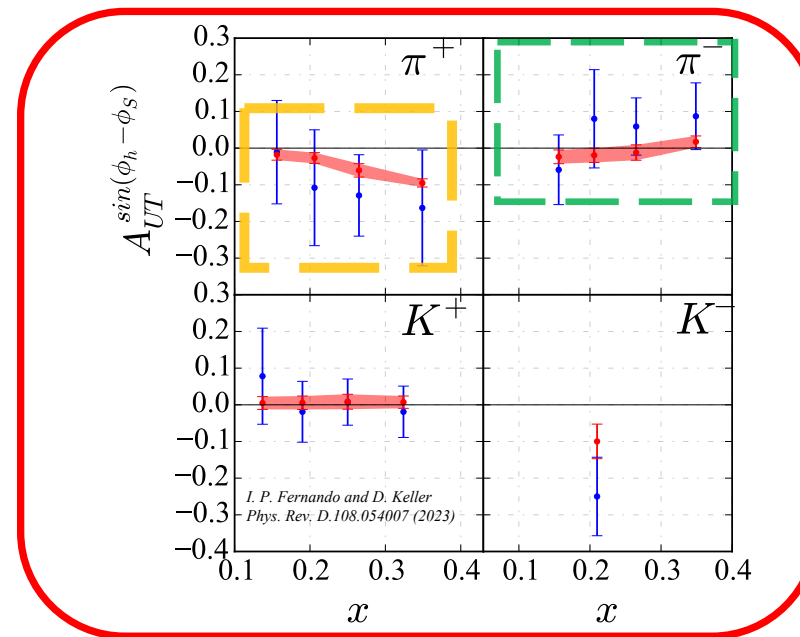
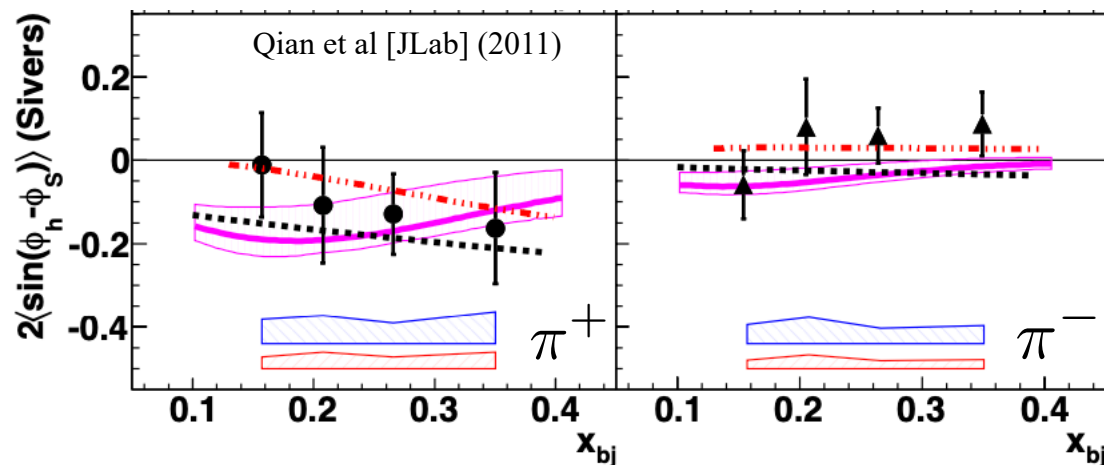
- No kinematic cuts are applied
Deuteron-DNN model can describe data reasonably well
- No iso-spin symmetry conditions are applied

$$f_{1T,u\leftarrow d}^\perp = f_{1T,d\leftarrow d}^\perp = \frac{f_{1T,u\leftarrow p}^\perp + f_{1T,d\leftarrow p}^\perp}{2} \quad \chi^2/N_{\text{pt}} = 0.81$$



Deuteron DNN Projections for JLab Kinematics

Deuteron-DNN



Bacchetta et al (2022)

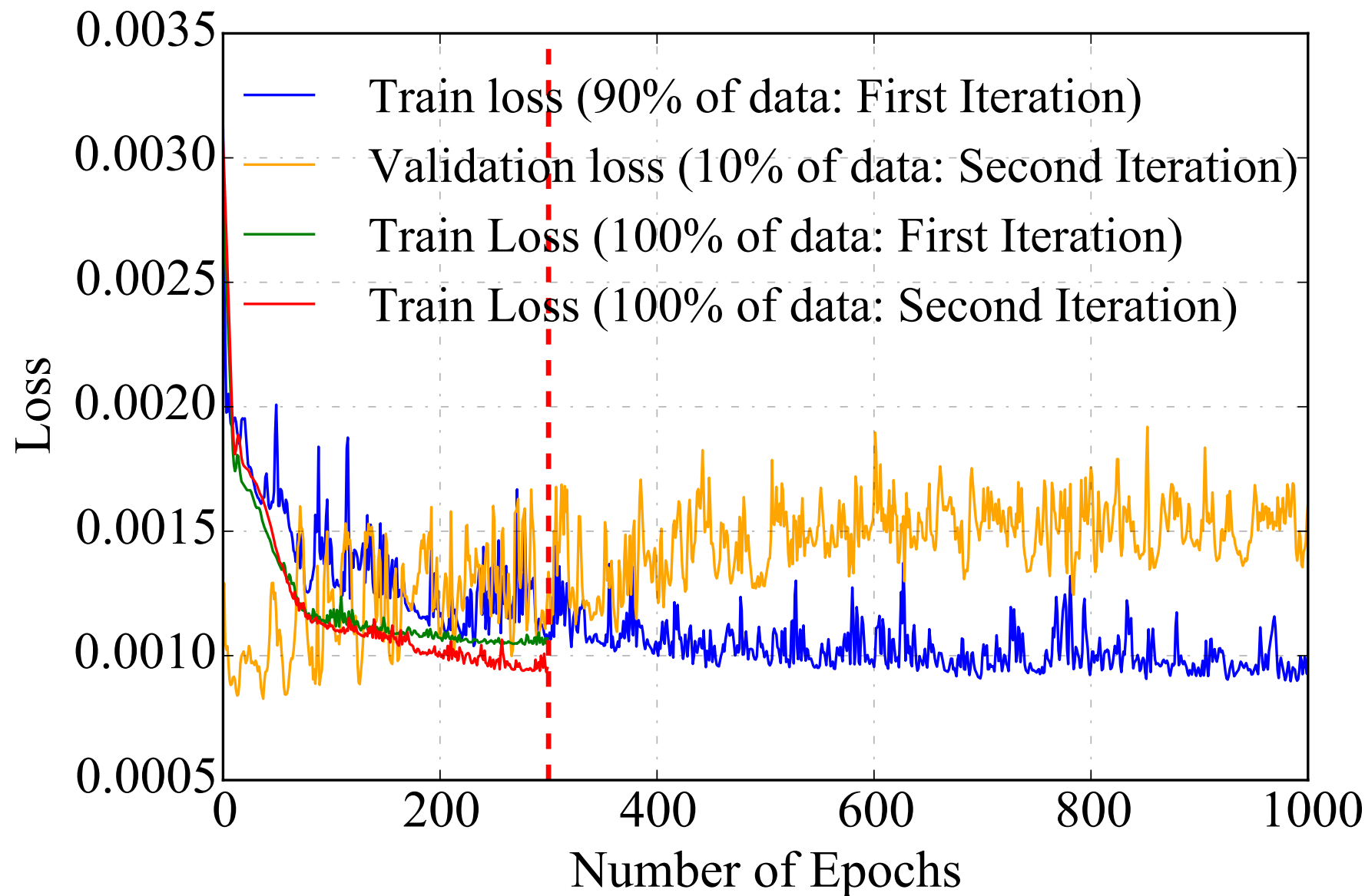


TABLE III. The summary of the optimized sets of hyperparameters: The indications in the table are \mathcal{C}_0^i and \mathcal{C}_0^f for results from the pseudodata from the generating function, \mathcal{C}_p^i , and \mathcal{C}_p^f for results from SIDIS data from experiments associated with the polarized-proton target, and \mathcal{C}_d^i and \mathcal{C}_d^f for results from SIDIS data from experiments associated with the polarized-deuterium target, where i and f indicate the *First Iteration* and *Second Iteration* respectively. The initial learning rate is also listed ($\times 10^{-4}$) as is the final training loss ($\times 10^{-3}$). The accuracy and precision in each case are the maxima over the phase space.

Hyperparameter	\mathcal{C}_0^i	\mathcal{C}_0^f	\mathcal{C}_p^i	\mathcal{C}_p^f	\mathcal{C}_d^i	\mathcal{C}_d^f
Hidden layers	5	7	5	7	5	8
Nodes/layer	256	256	550	550	256	256
Learning rate	1	0.125	5	1	10	1
Batch size	200	256	300	300	100	20
Number of epochs	1000	1000	300	300	200	200
Training loss	0.6	0.05	1.5	1	2	1
ϵ_u^{\max}	95.67	99.27	55.21	94.04	56.80	93.02
$\epsilon_{\bar{u}}^{\max}$	42.62	98.09	52.57	96.70	34.83	91.40
ϵ_d^{\max}	80.46	98.89	55.69	93.13	52.44	89.27
$\epsilon_{\bar{d}}^{\max}$	74.59	97.08	55.37	95.04	46.60	92.58
ϵ_s^{\max}	45.53	79.27	49.54	90.64	36.34	93.41
$\epsilon_{\bar{s}}^{\max}$	59.27	91.13	33.89	82.51	65.57	91.45
σ_u^{\max}	3	0.1	5	2	2	0.4
$\sigma_{\bar{u}}^{\max}$	2	0.2	6	2	8	2
σ_d^{\max}	10	1	20	6	2	1
$\sigma_{\bar{d}}^{\max}$	7	4	20	8	7	1
σ_s^{\max}	2	0.2	4	1	6	2
$\sigma_{\bar{s}}^{\max}$	1	0.1	4	2	6	3

Systematic Studies: data cuts

Backup

$$W^{\mu\nu} = \sum_f |\mathcal{H}_f(Q^2, \mu)|^{\mu\nu} \times \int d^2 k_\perp d^2 p_\perp \delta^{(2)}(z_h k_\perp + p_\perp - p_{hT}) \times F_{f/N^\uparrow}(x, z_h k_\perp, S; \mu, \zeta_F) D_{h/f}(z_h, p_\perp; \mu, \zeta_D) + Y(p_{hT}, Q^2),$$

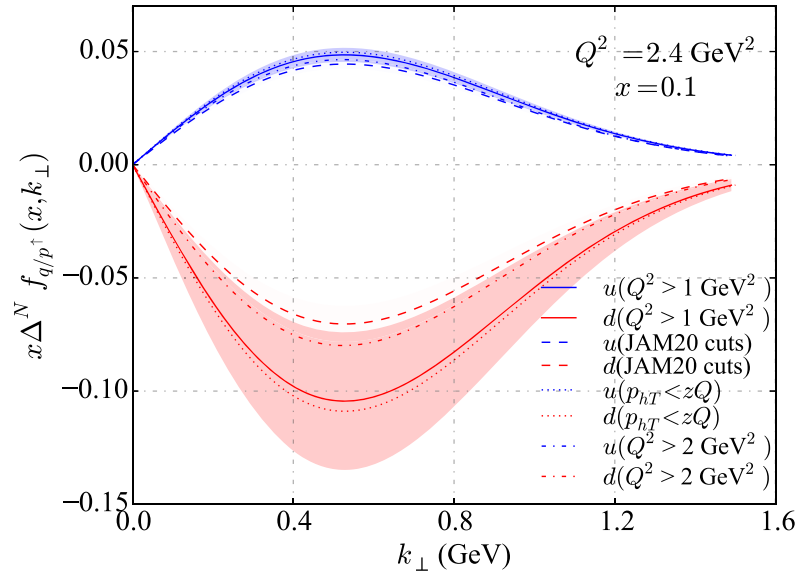


FIG. 17. Solid lines with light band represent the u (in blue), d (in red) Siverts functions using the cut $Q^2 > 1 \text{ GeV}^2$. These resulting DNN models made from the cuts from all tests are also shown.

The applicability of TMD factorization was ensured by applying cuts to SIDIS data based on various criteria in the literature.

In this exploratory effort with DNNs, the power corrections are not directly imposed. In addition to the basic data cut $Q^2 > 1 \text{ GeV}^2$ we performed $Q^2 > 2 \text{ GeV}^2$ and $p_{hT} < zQ$ cuts separately with the proton-DNN model to understand the impact on the extracted Siverts functions.

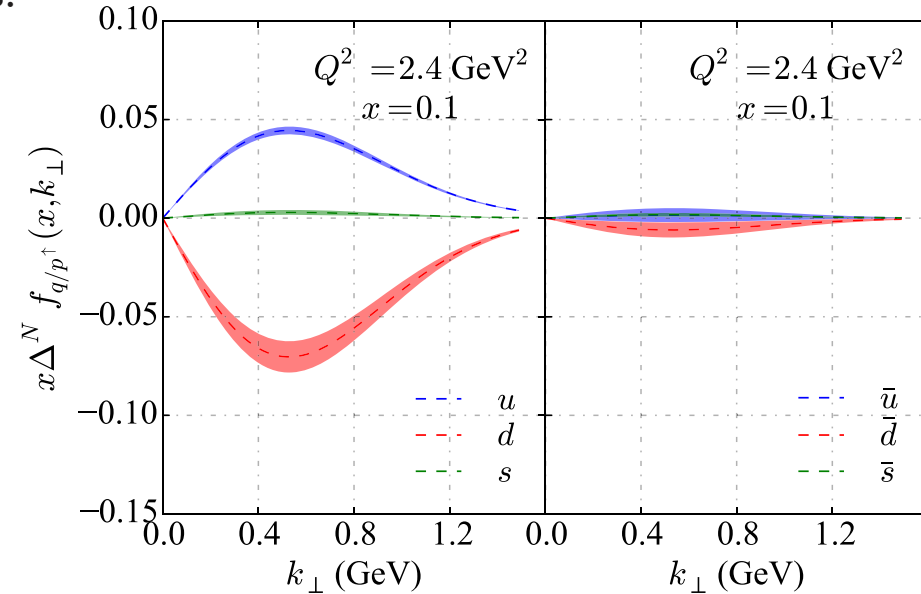


FIG. 18. Siverts functions from a retrained DNN model using the cuts [65] to the data demonstrating that being selective with the data can reduce the error bands of the fit but may also add an unintentional bias.

Systematic Studies: Choice of $h(k)$

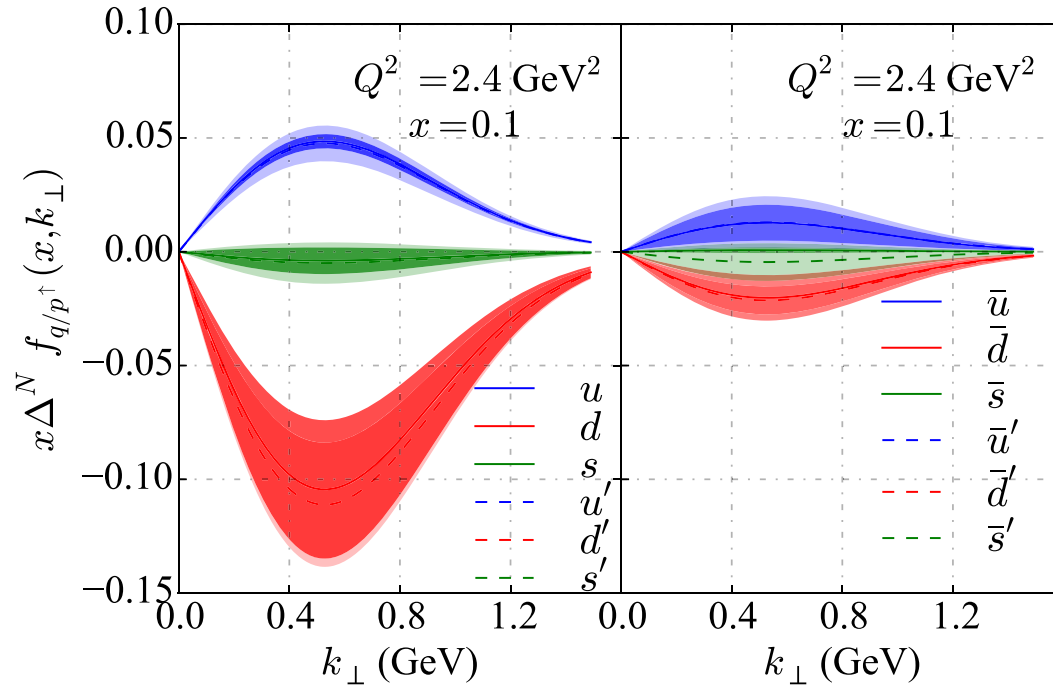


FIG. 19. Using two different $h(k_\perp)$. Solid line with dark band represents the Siverts functions with $h(k_\perp) = \sqrt{2}e \frac{k_\perp}{m_1} e^{-k_\perp^2/m_1^2}$, whereas the dashed line with light band represents the Siverts functions with $h(k_\perp) = \frac{2k_\perp m_1}{m_1^2 + k_\perp^2}$.

$$h(k_\perp) = \sqrt{2}e \frac{k_\perp}{m_1} e^{-k_\perp^2/m_1^2}$$

$$h(k_\perp) = \frac{2k_\perp m_1}{m_1^2 + k_\perp^2}$$

Systematic Studies : TMD Evolution

Backup

The solution of the TMD evolution equations

$$\mu^2 \frac{dF(x, b; \mu, \zeta)}{d\mu^2} = \frac{\gamma_F(\mu, \zeta)}{2} F(x, b; \mu, \zeta)$$

$$\zeta \frac{dF(x, b; \mu, \zeta)}{d\zeta} = -\mathcal{D}(b, \mu) F(x, b; \mu, \zeta),$$

$$F(x, b; \mu, \zeta) = \left(\frac{\zeta}{\zeta_\mu(b)} \right)^{-\mathcal{D}(b, \mu)} F(x, b)$$

$$\mu \sim Q, \quad \zeta_F \zeta_D \sim Q^4, \quad \mu^2 = \zeta^2 = Q^2$$

$$\mathcal{N}_q(x) \longrightarrow \mathcal{N}_q(x, Q^2)$$

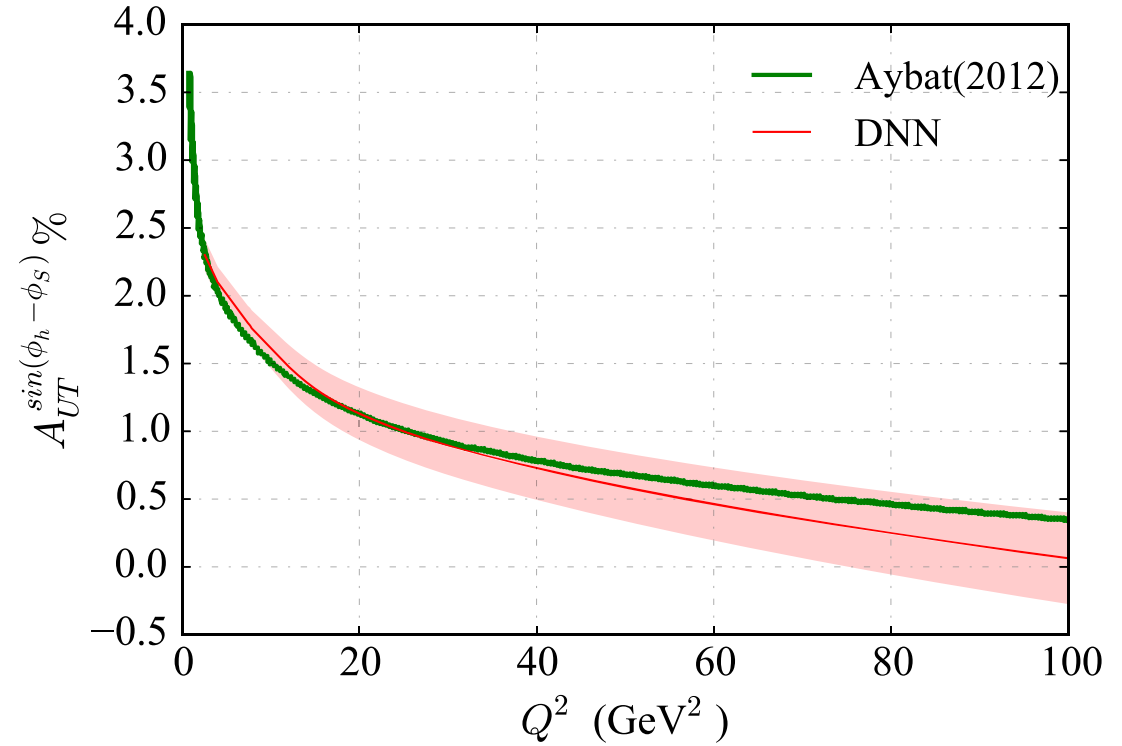


FIG. 21. The Sivers asymmetry evolution in Q^2 compared to the result from [6]. The red-colored solid line and the band represent the mean and standard deviation of the $A_{UT}^{\sin(\phi_h - \phi_S)}$ from 1000 replica models of the proton DNN at $x = 0.12$, $z = 0.32$, $p_{hT} = 0.14$ GeV.