

SHEDDING LIGHT ON SHADOW GENERALIZED PARTON DISTRIBUTIONS (GPDS)

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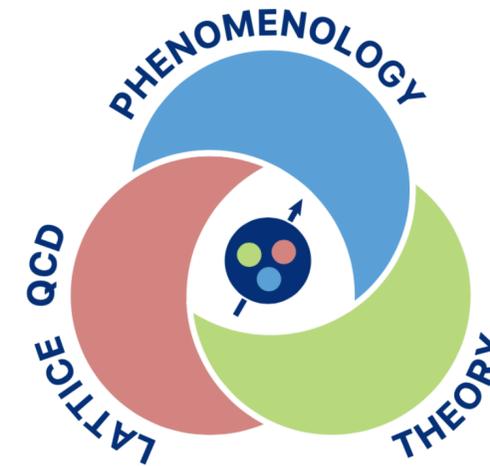
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Introduction

- * Generalized Parton Distributions (GPDs) contain information about many hadron properties:
 - * 3D structure
 - * Spin sum
 - * Pressure and shear force distributions
- * Goal:
 - * Perform a global fit of GPDs using the Jefferson Lab Angular Momentum (JAM) methodologies.
- * Obstacle:
 - * Shadow GPDs (SGPDS) (Bertone, et. al. Phys.Rev.D 103 (2021) 11, 114019):
 - * There is an infinite number of functions that can give the same CFF.



**QUARK-GLUON
TOMOGRAPHY
COLLABORATION**



GPDs

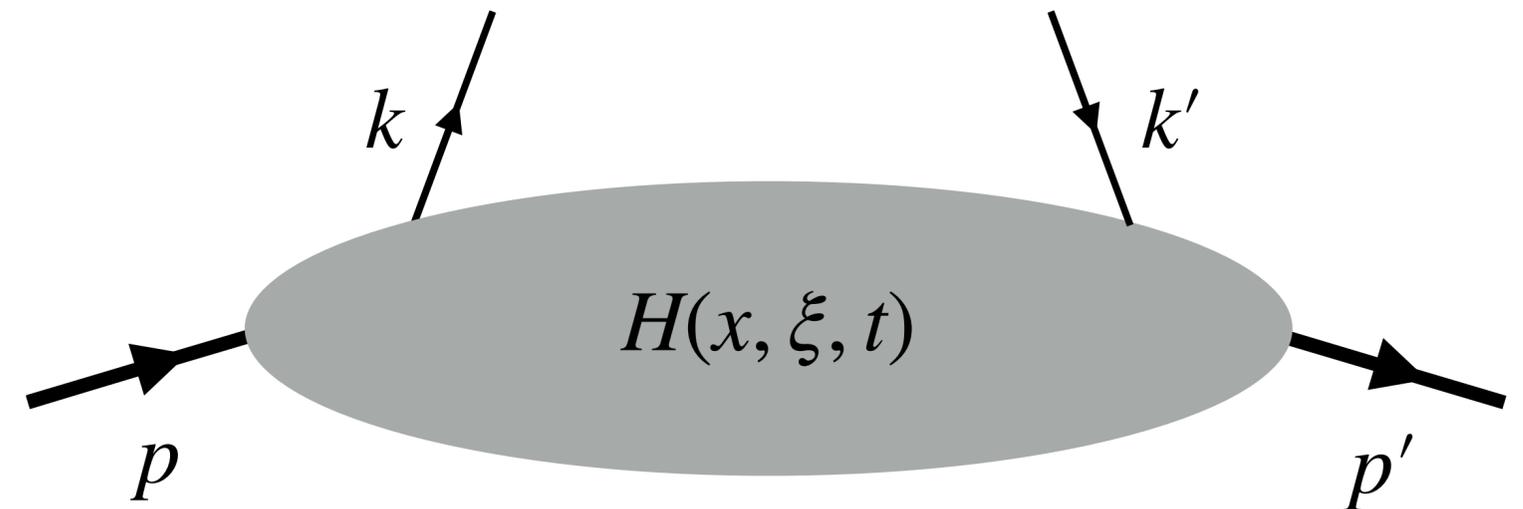
* Definition:

$$P \cdot n \int \frac{d\lambda}{2\pi} e^{ixP \cdot n\lambda} \langle p' | \bar{\psi}^q(-\frac{1}{2}\lambda n) \not{n} \psi^q(\frac{1}{2}\lambda n) | p \rangle = \bar{u}(p') \left[H^q(x, \xi, t; \mu^2) \not{n} + E^q(x, \xi, t; \mu^2) \frac{i\sigma^{n\Delta}}{2M} \right] u(p),$$

$$n_\mu n_\nu \int \frac{d\lambda}{2\pi} e^{ixP \cdot n\lambda} \langle p' | G^{\mu\alpha}(-\frac{1}{2}\lambda n) G_{\alpha\nu}(\frac{1}{2}\lambda n) | p \rangle = \bar{u}(p') \left[x H^g(x, \xi, t; \mu^2) \not{n} + x E^g(x, \xi, t; \mu^2) \frac{i\sigma^{n\Delta}}{2M} \right] u(p),$$

* Functions of x , ξ , and t :

$$x = \frac{k^+ + k'^+}{p^+ + p'^+} \quad \xi = \frac{p'^+ - p^+}{p^+ + p'^+} \quad t = (p' - p)^2$$



GPDs

- * Properties:

- * Polynomiality:
$$\int_{-1}^1 dx x^s H^a(x, \xi, t; \mu^2) = \sum_{i=0}^s (2\xi)^i A_{s+1,i}^a(t, \mu^2) + \text{mod}(s, 2) (2\xi)^{s+1} C_{s+1}^a(t, \mu^2),$$

$$\int_{-1}^1 dx x^s E^a(x, \xi, t; \mu^2) = \sum_{i=0}^s (2\xi)^i B_{s+1,i}^a(t, \mu^2) - \text{mod}(s, 2) (2\xi)^{s+1} C_{s+1}^a(t, \mu^2),$$

- * Forward Limit ($\xi, t \rightarrow 0$):

$$H^q(x, 0, 0) = q(x) \Theta(x) - \bar{q}(-x) \Theta(-x),$$
$$2 H^g(x, 0, 0) = g(x) \Theta(x) - g(-x) \Theta(-x),$$

- * Evolution:

- * GPDs change with the energy scale in accordance with evolution equations of the general form:

$$\frac{dH^a(x, \xi, t)}{d \ln Q^2} = \int dx P^a(x, \xi) H^a(x, \xi, t; Q_0^2)$$

The Inverse Problem

- * Deeply virtual Compton scattering:

- * Compton Form Factors:

$$\mathcal{H}(\xi, t, Q^2) = \int_{-1}^1 dx \sum_a C^a(x, \xi, Q^2, \mu^2) H^a(x, \xi, t; \mu^2)$$

- * x-dependence is lost in the integration:

- * There is an infinite number of functions that can give the same CFF.

Shadow GPDs

- * The difference between one of the multiple solutions to the inverse problem and the true GPD:

$$F_S^a(x, \xi; \mu^2) = F_F^a(x, \xi; \mu^2) - F_T^a(x, \xi; \mu^2)$$

- * Can rule out any F_F^a that do not satisfy the properties of GPDs, therefore SGPDs:

- * Must satisfy polynomiality

- * Zero contribution to CFF:

$$\sum_a C^a(x, \xi, Q^2, \mu^2) \otimes F_S^a(x, \xi; \mu^2) = 0$$

- * Forward Limit: $H_S^a(x, 0, 0) = 0$

Evolution and SGPDs

- * Example SGPDs explored in Bertone, et. al. Phys.Rev.D 103 (2021) 11, 114019 give very small but non-zero CFF after evolution to a different energy scale.
- * SGPDs can be multiplied by any factor and the result would still be a SGPD at the input scale
- * Non-zero CFF after evolution would be multiplied by this factor
- * Data spanning a range of energy scales would give a limit to the possible scaling factors

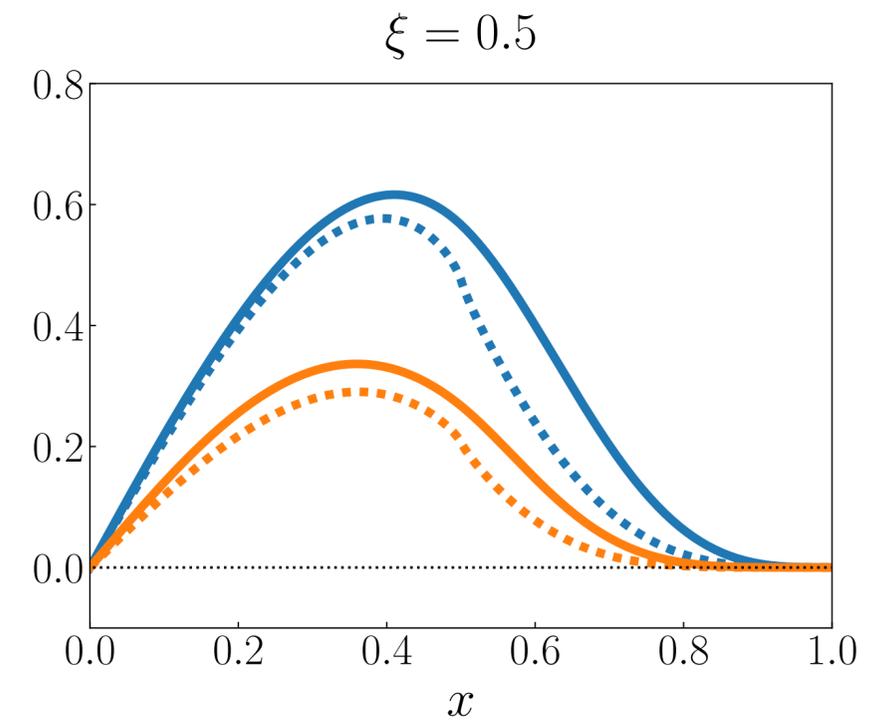
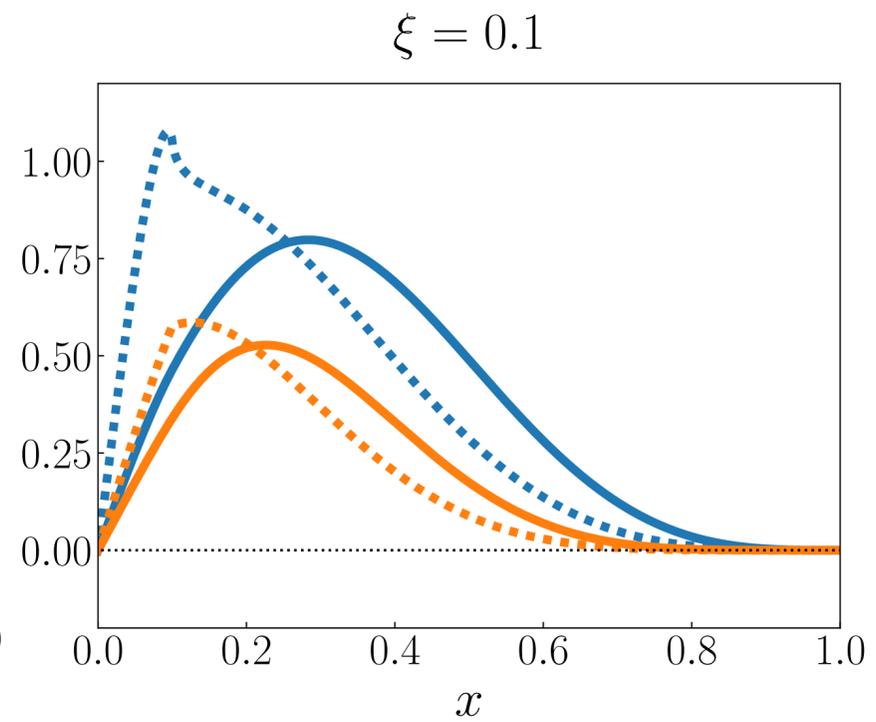
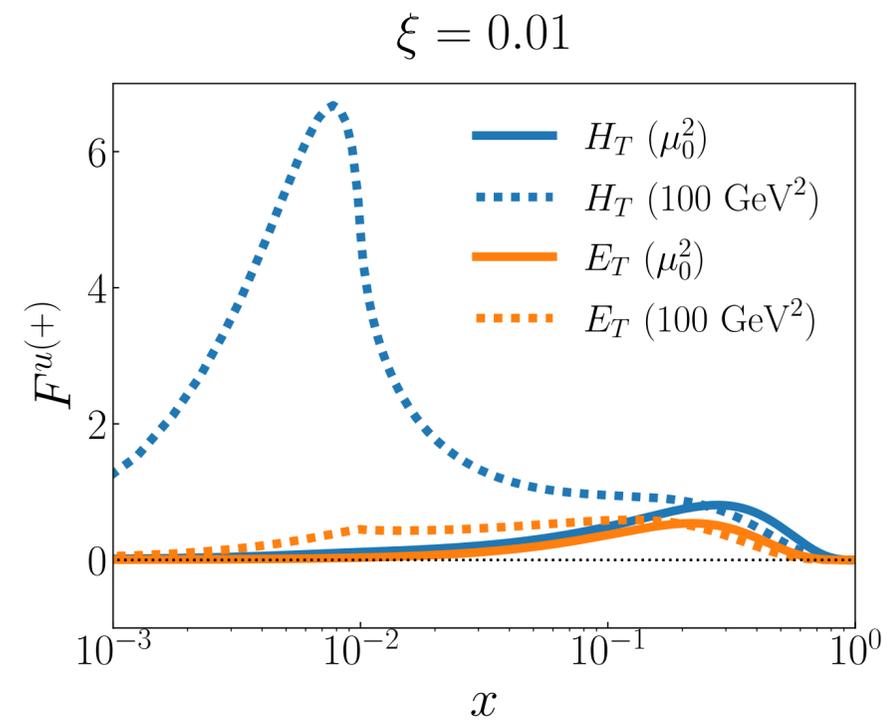
Evolution and SGPDs

- * In this work:
 - * Generate simulated CFF data spanning a range of energy scales and skewness using a model
 - * Calculate how this data constrains a Monte Carlo sampling of SGPDs

“True” GPDs

- * Use VGG model as a proxy for the “true” GPD:
 - * Vanderhaeghen, et. al., Phys. Rev. Lett. 80, 5064 (1998)
 - * Vanderhaeghen, et. al., Phys. Rev. D 60, 094017 (1999)
 - * Goeke, et. al., Prog. Part. Nucl. Phys. 47, 401 (2001)
 - * Guidal, et. al., Phys. Rev. D 72, 054013 (2005)
- * Use PDFs from JAM20-SIDIS (EM, et. al., Phys. Rev. D 104, 016015 (2021))

“True” GPDs



Calculating Shadow GPDs

- * Start from a double distribution (DD):

$$F_{DD}(\alpha, \beta) = \sum_{\substack{m+n \leq N \\ m \text{ even}, n \text{ odd}}} c_{mn} \alpha^m \beta^n$$

- * SGPD is a Radon transform of the DD:

$$H_S(x, \xi) = \int_{-1}^1 d\beta \int_{-1+|\beta|}^{1-|\beta|} d\alpha \delta(x - \beta - \alpha\xi) F_{DD}(\alpha, \beta)$$

- * This guarantees the SGPDs satisfy polynomiality

Calculating Shadow GPDs

- * SGPD conditions give a set of equations that can be solved for the unknowns (c_{mn})
 - * For a given N there are more unknown coefficients than constraining equations:
 - * Assign random values to enough randomly selected coefficients to reduce the number of unknowns so that the equations can be solved
 - * Use $N = 27$
- * SGPDs give zero contribution to the CFF at next-to-leading order

Calculating Shadow GPDs

- * For SGPDs derived this way we can impose the forward limit in two ways:

- * Type A:

- * Consistent with Bertone, et. al. Phys.Rev.D 103 (2021) 11, 114019:

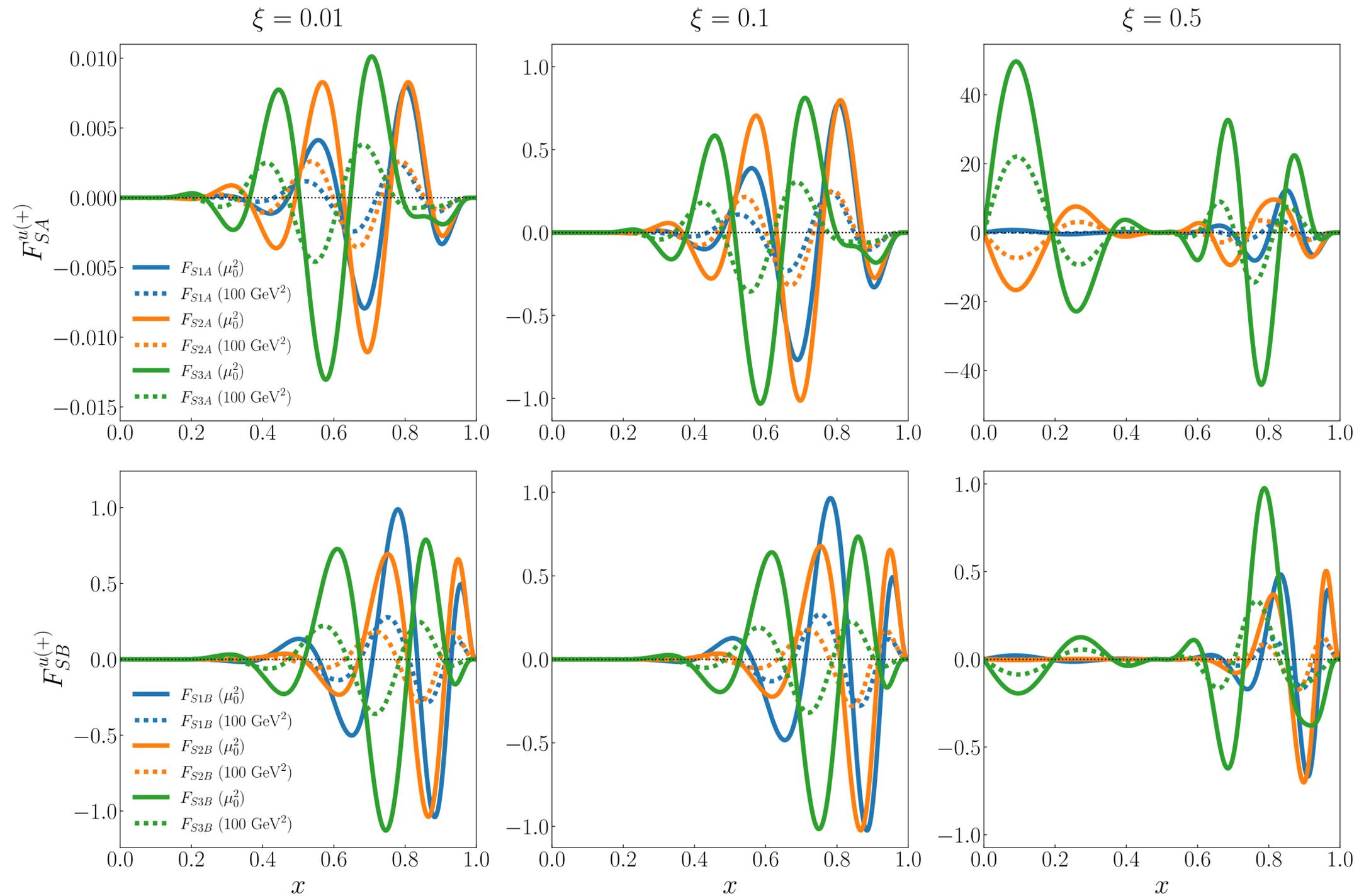
$$H_S^{u(+)}(x,0; \mu_0) = 0$$

- * Type B:

- * Could also multiply F_{DD} by a function of t that is zero when $t = 0$

$$H_S^{u(+)}(x,0; \mu_0) \neq 0$$

Example Shadow GPDs



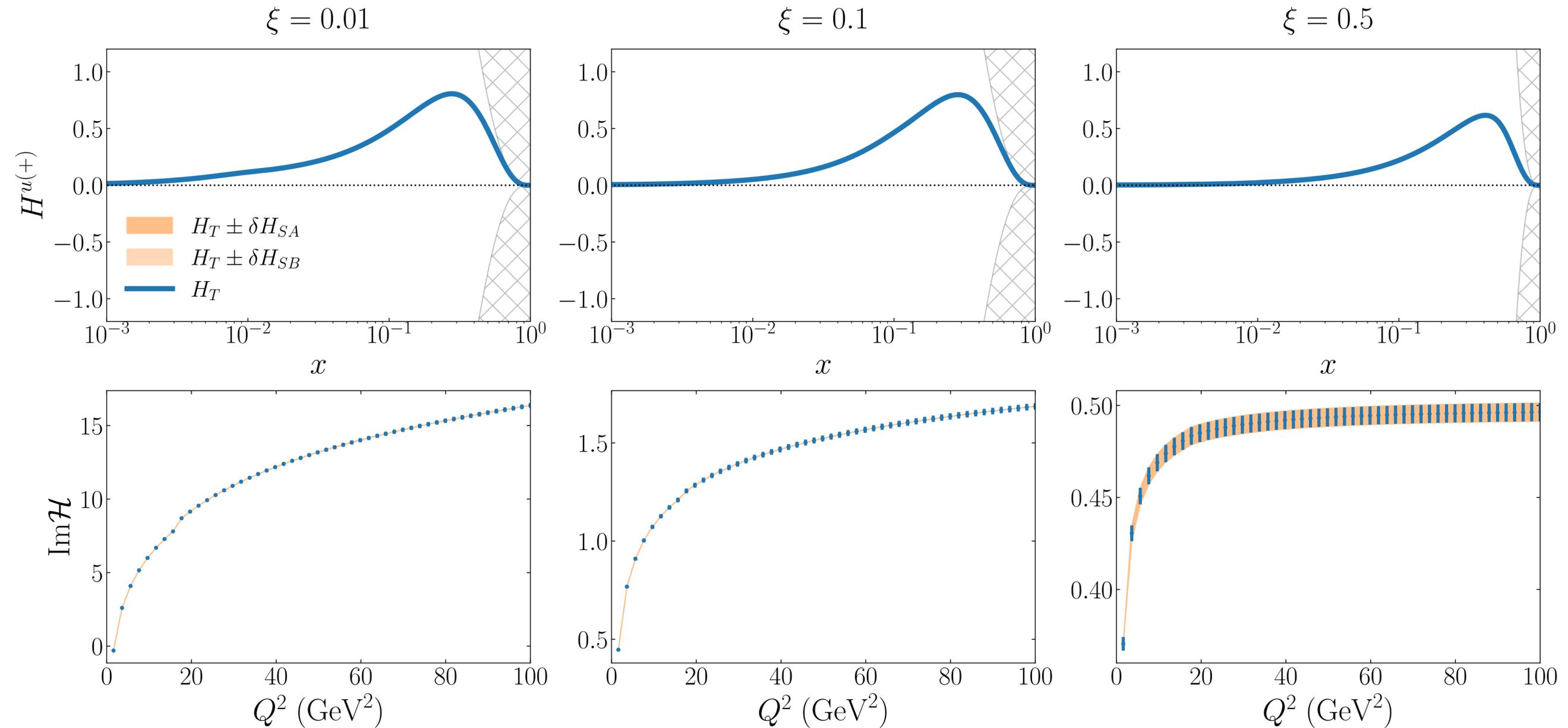
Exploring SGPDs and Evolution

- * Use Monte Carlo sampling to generate replicas that are linear combinations of three SGPDs:

$$H^{u(+)}(x, \xi; \mu^2, \lambda) = H_T^{u(+)}(x, \xi; \mu^2) + \lambda_1 H_{S1}^{u(+)}(x, \xi; \mu^2) + \lambda_2 H_{S2}^{u(+)}(x, \xi; \mu^2) + \lambda_3 H_{S3}^{u(+)}(x, \xi; \mu^2)$$

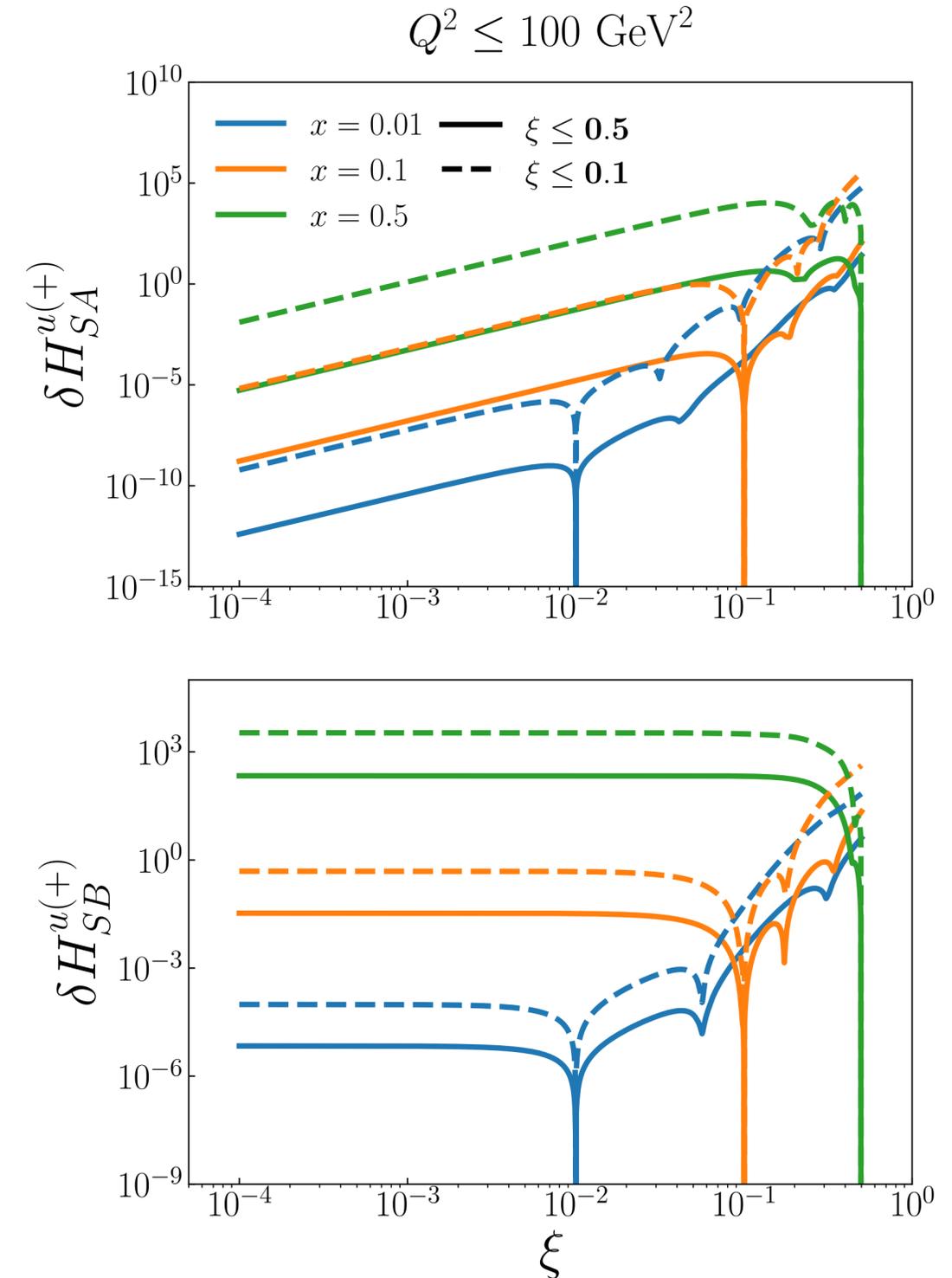
- * Randomly select the scaling factors until we get 10000 replicas that all give CFFs that are within 1% of the simulated data from the model.
- * Plot the region δH_S : Outer boundary of all 10000 replicas

Exploring SGPDs and Evolution



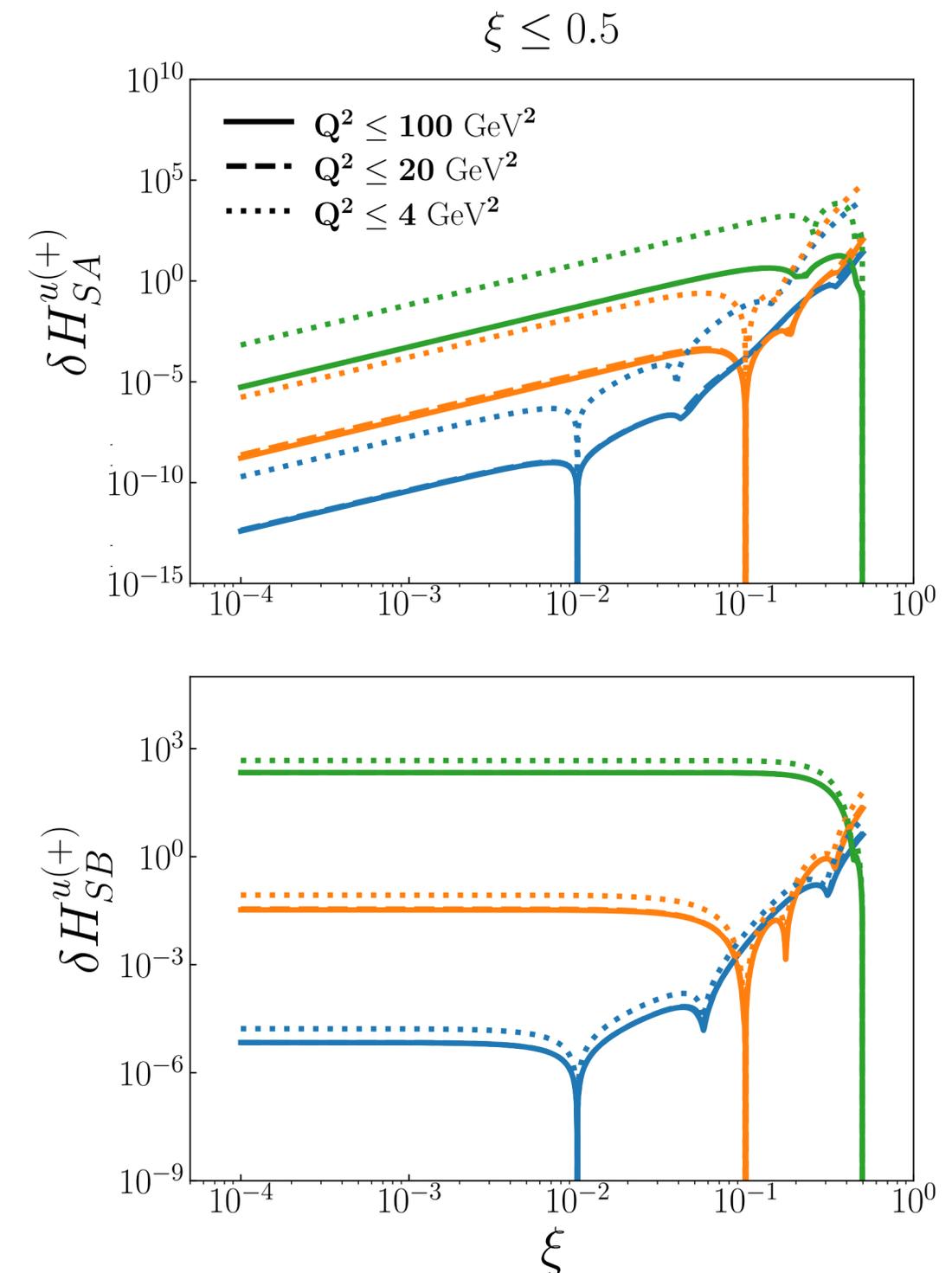
Exploring SGPDs and Evolution

- * Inclusion of higher ξ data leads to better constraint of SGPDs at smaller ξ
- * True over the full range of x when $H_S^{u(+)}(x,0; \mu_0) = 0$
- * Only true for low x when $H_S^{u(+)}(x,0; \mu_0) \neq 0$



Exploring SGPDs and Evolution

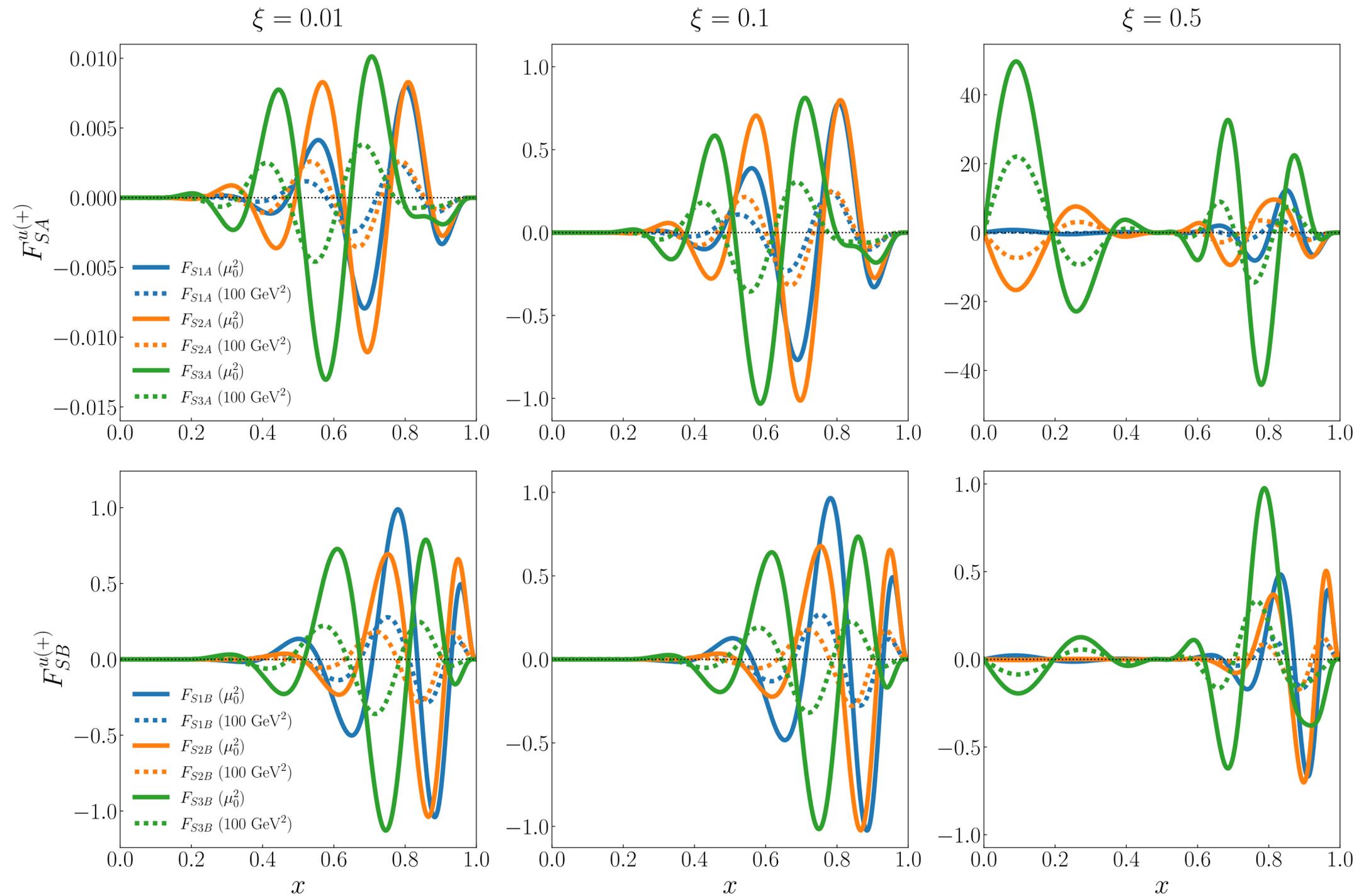
- * Some range of Q^2 is necessary for evolution to constrain the SGPDs but a large range is not as necessary as having large ξ data.



Exploring SGPDs and Evolution

- * The trend of larger ξ data leading to better constrained SGPDs at smaller ξ is a direct result of the ξ dependence of the SGPDs
- * Independent of the model used as a proxy for the “true” GPD
- * Independent of the chosen uncertainty

Example Shadow GPDs



Conclusions

- * Conclusions:
 - * For the SGPDs that have been explored here:
 - * Data spanning a range of Q^2 at larger ξ leads to the SGPDs being better constrained at lower ξ at least in the range of low x
 - * These findings are independent of the model used as the proxy for the “true” GPD.
 - * The SGPDs explored are only a small sampling of all possible SGPDs:
 - * At this point we cannot generalize these results to all SGPDs
 - * Data spanning a range of Q^2 at larger ξ is a necessary but possibly not sufficient condition for extracting GPDs from DVCS data.