

# Impact of extracting strange quark distribution of proton from PVDIS experiment

Nobuo Sato

Winter Hall A (Jan 2023)

$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{\psi}_q (i\gamma_\mu D^\mu - m_q) \psi_q - \frac{1}{2} \text{Tr}[G_{\mu\nu} G^{\mu\nu}]$$

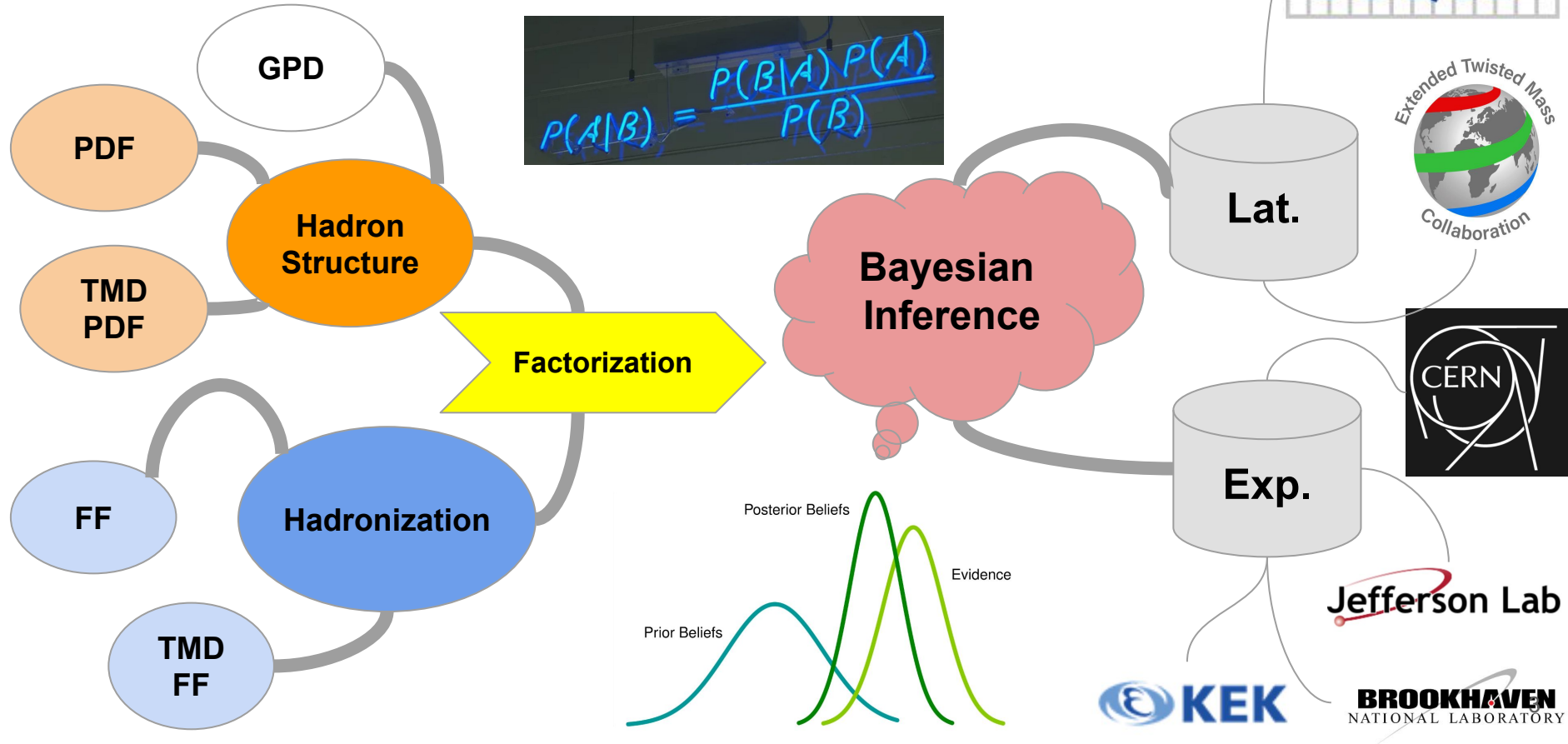
Jefferson Lab

# Outline

1. Recent highlights
2. Status of strange pdf
3. Parity violating DIS
4. PVDIS @ SoLID -22
5. Sin2W
6. QED radiative effects
7. Integrating thy/exp/cs

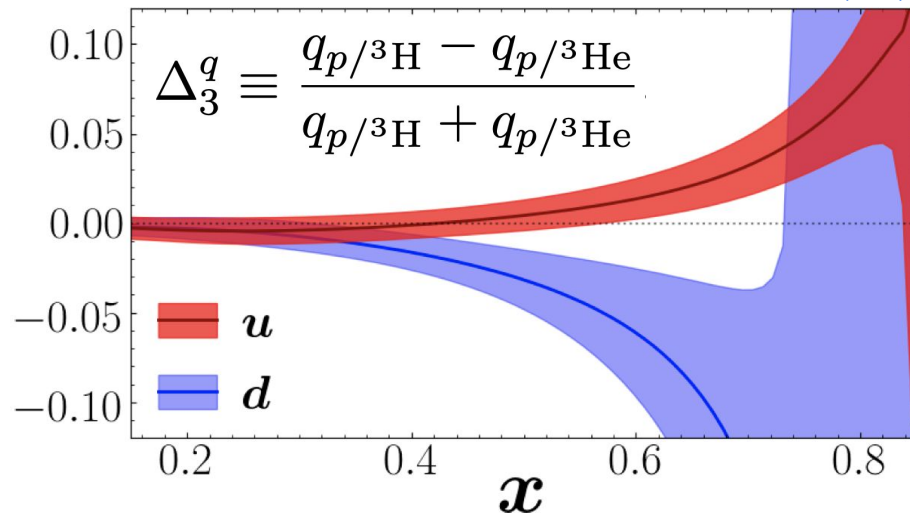
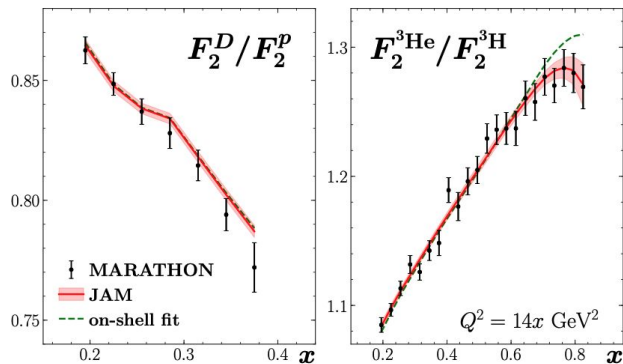
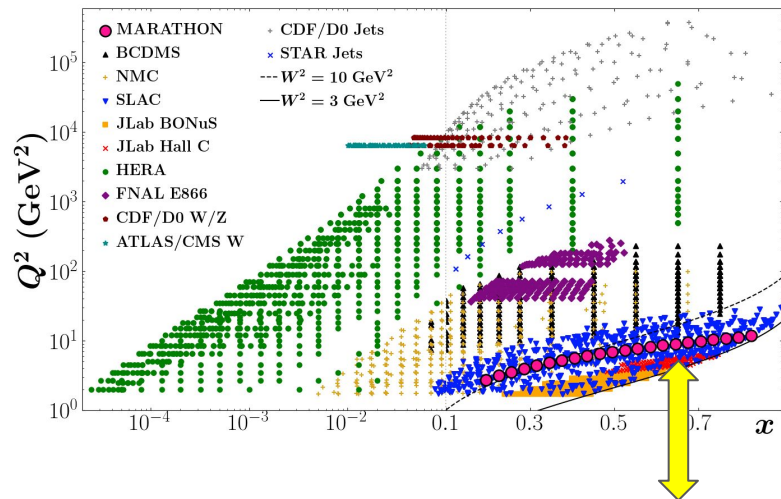


# A holistic approach for global analysis



# Isvector EMC effects from MARATHON data

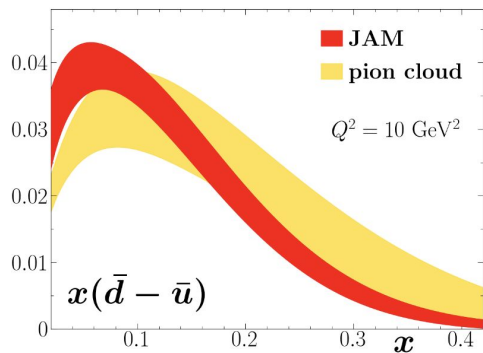
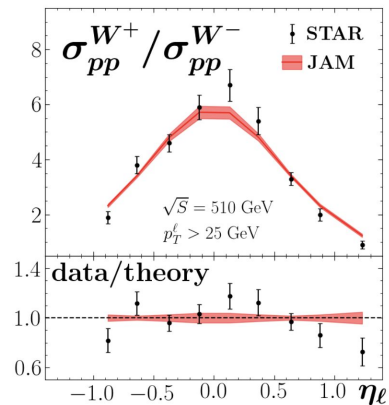
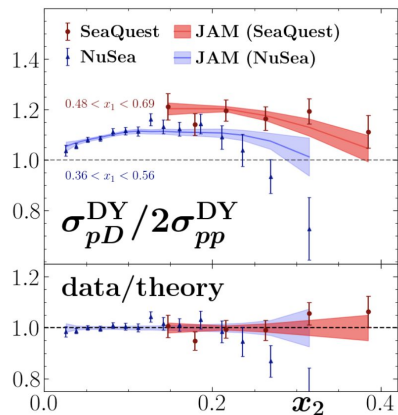
Cocuzza, et al (PRL)



- Global analysis including latest collider  $W/Z$  data and MARATHON  $d/p$ , helium, tritium DIS data
- *Evidence of different medium modifications for  $u$  and  $d$  quarks*
- *Naive modeling of nuclear PDFs, e.g.  $u/pA = d/nA$  (violates isospin for non-isoscalar  $A$ ) is wrong*

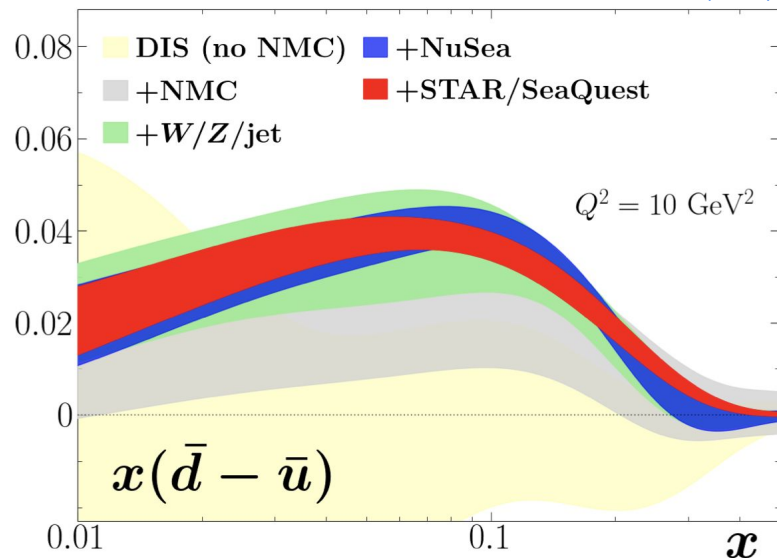


# Antimatter asymmetry



$$(\bar{d} - \bar{u})(x) = [(f_{n\pi^+} + f_{\Delta^0\pi^+} - f_{\Delta^{++}\pi^-}) \otimes \bar{q}_v^\pi](x),$$

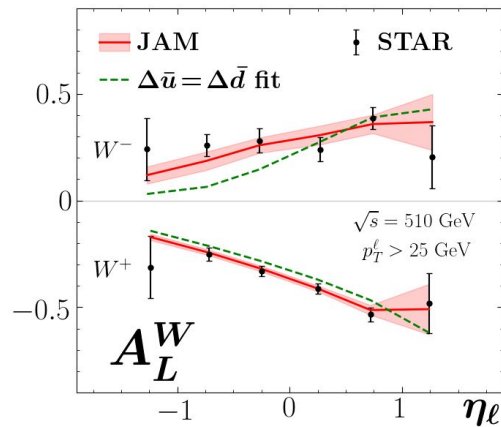
Cocuzza, Melnitchouk, Metz, Sato (PRD)



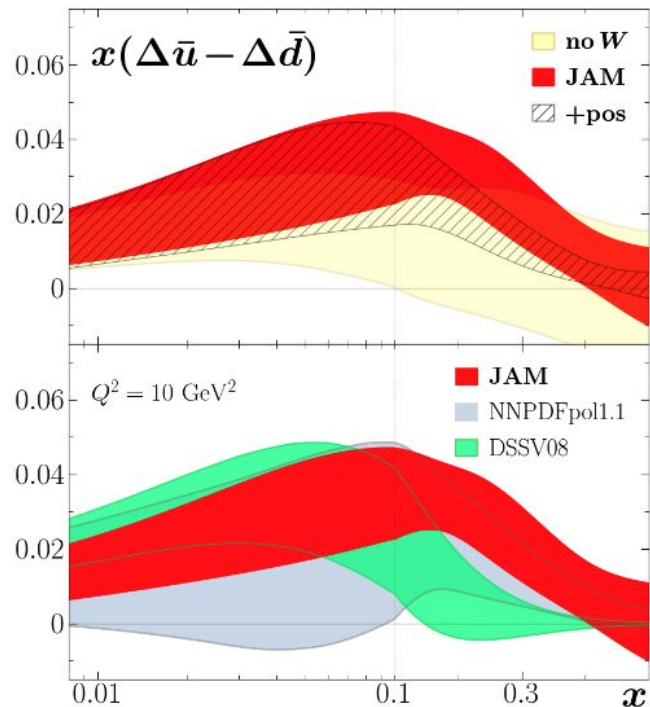
- First global analysis to include latest SeaQuest and STAR data
- Most precise phenomenological extraction of  $d\bar{u}$  asymmetry to date
- Quantitative test of the pion-cloud model

# Polarized antimatter asymmetry

Process	$N_{\text{dat}}$	$\chi^2/N_{\text{dat}}$
<i>Polarized</i>		
Inclusive DIS	365	0.95
SIDIS ( $\pi^+, \pi^-$ )	64	1.05
SIDIS ( $K^+, K^-$ )	57	0.42
SIDIS ( $h^+, h^-$ )	110	0.95
Inclusive jets	83	0.84
STAR $W^\pm$	12	0.65
PHENIX $W^\pm/Z$	6	0.50
<b>Total</b>	<b>697</b>	<b>0.89</b>
<i>Unpolarized</i>		
Inclusive DIS	3908	1.17
SIDIS ( $\pi^+, \pi^-$ )	498	0.94
SIDIS ( $K^+, K^-$ )	494	1.31
SIDIS ( $h^+, h^-$ )	498	0.71
Inclusive jets	198	1.28
Drell-Yan	205	1.21
$W/Z$ production	153	1.01
<b>Total</b>	<b>5954</b>	<b>1.12</b>
SIA ( $\pi^\pm$ )	231	0.91
SIA ( $K^\pm$ )	213	0.70
SIA ( $h^\pm$ )	120	1.07
<b>Total</b>	<b>7215</b>	<b>1.08</b>



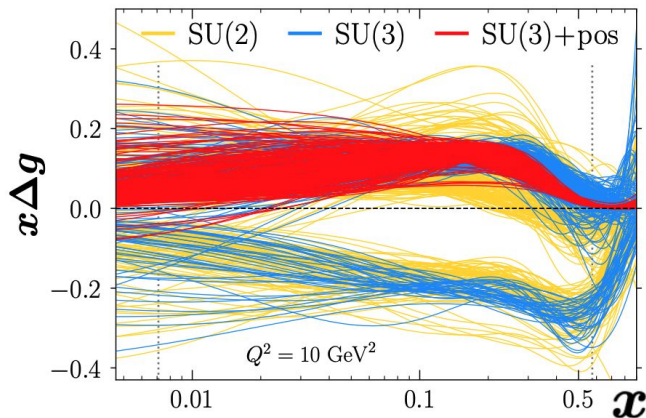
Cocuzza, Melnitchouk, Metz, Sato (PRD)



- First simultaneous extraction of unpolarized and helicity PDFs and FFs in global analysis with inclusion of RHIC spin  $W^{+/-}$  data
- Most precise phenomenological extraction of polarized  $d\bar{u}$ - $u\bar{d}$  asymmetry to date

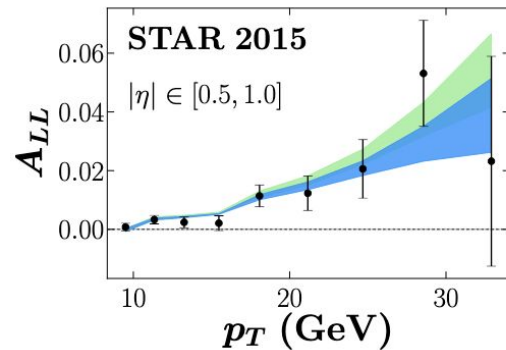
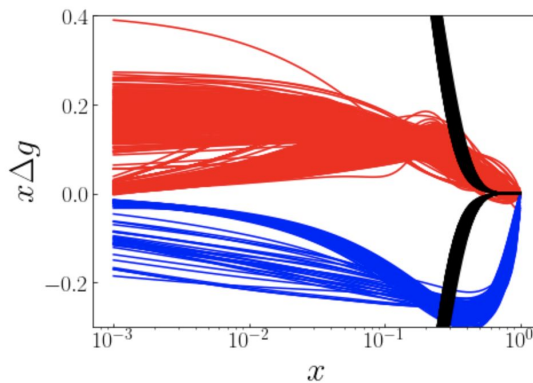
# News on Gluon helicity

Zhou, Melnitchouk, Sato (PRD)

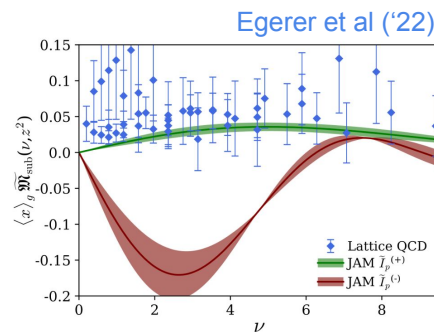
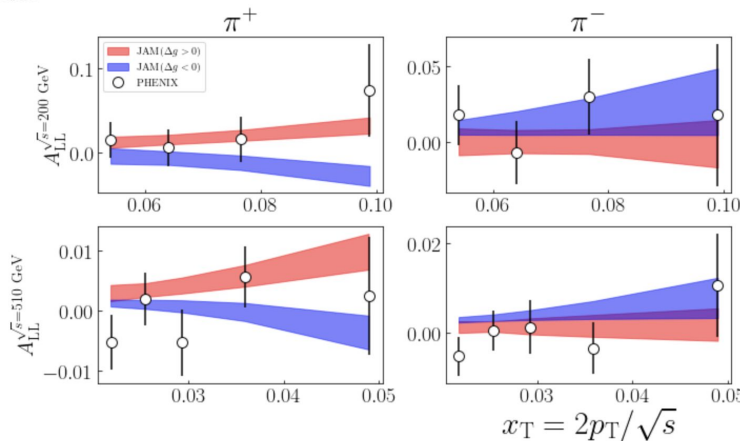


$$\int_{0.05}^1 dx \Delta g(x) = \begin{array}{ll} 0.23 \pm 0.03 & \text{pos} \\ -0.62 \pm 0.03 & \text{neg} \end{array}$$

$\Delta g > 0$   
 $\Delta g < 0$

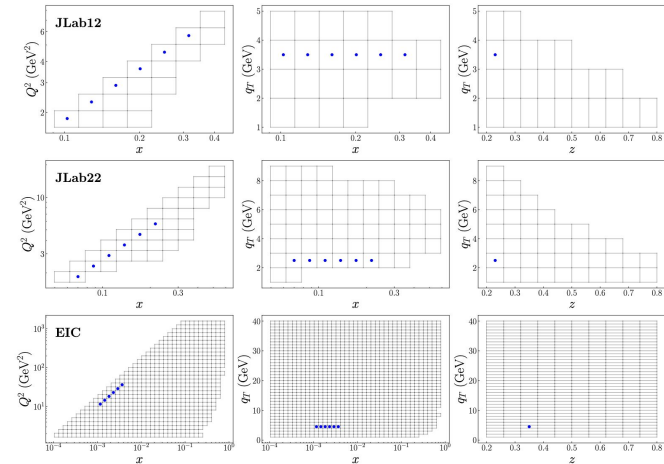
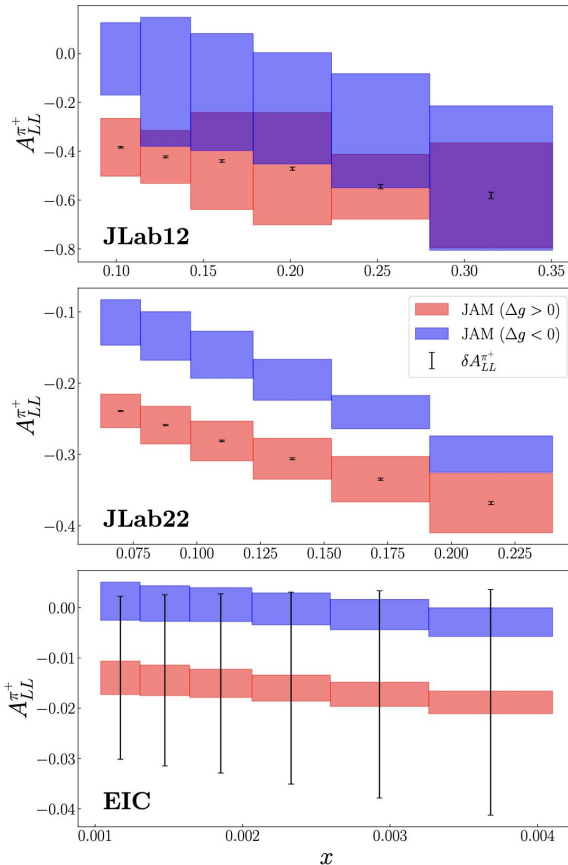
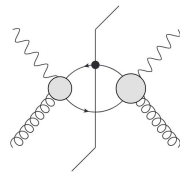
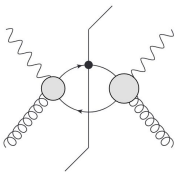
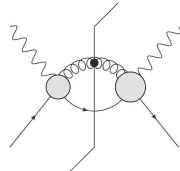
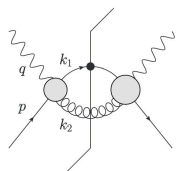
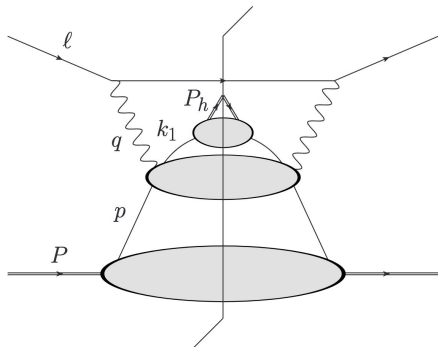


- Inclusion of RHIC polarized jet data allows both positive and negative gluon helicity solutions (in absence of positivity constraints on unpolarized gluon PDF)
- PHENIX has attempted to have empirical confirmation of gluon helicity sign (PRD102.032001, PRD91.032001)



# Accessing gluon polarization with high- $P_T$ hadrons in SIDIS

R.M. Whitehill (Wichita State U.), Yiyu Zhou (South China Normal U. and UCLA), N. Sato (Jefferson Lab), W. Melnitchouk (Jefferson Lab and Adelaide U., Sch. Chem. Phys.)  
Oct 21, 2022



**Bottom line:** sign of gluon polarization might be resolved empirically at JLab 22

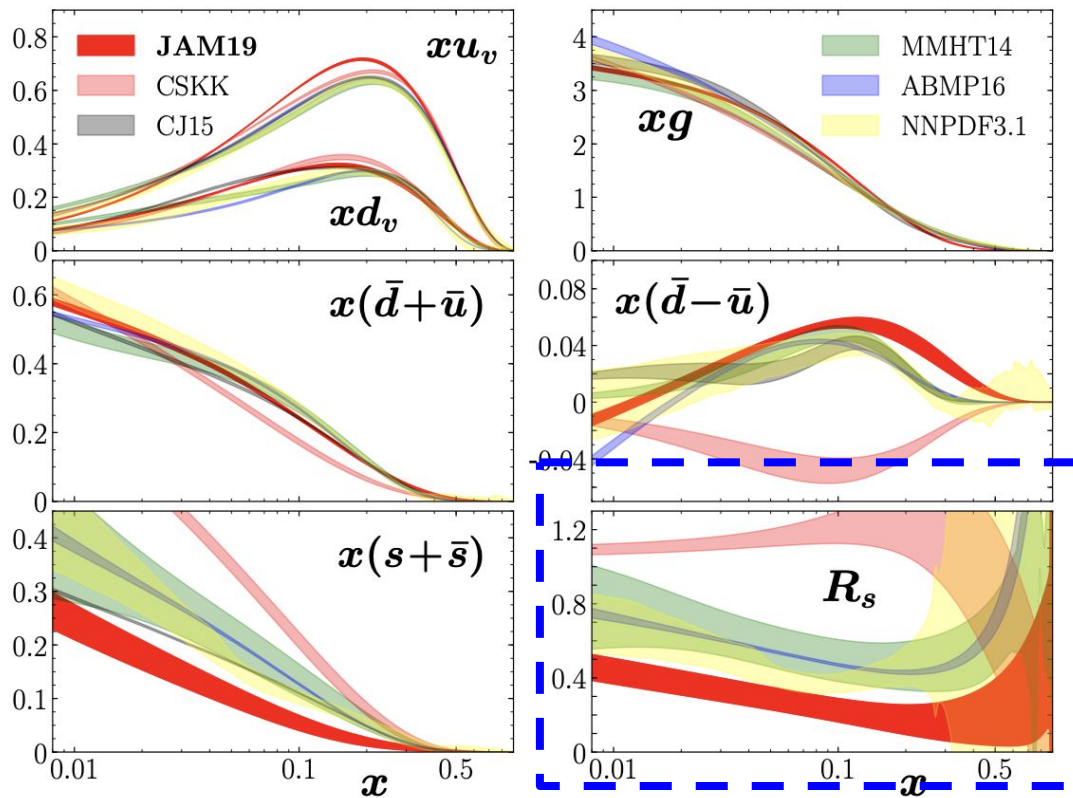


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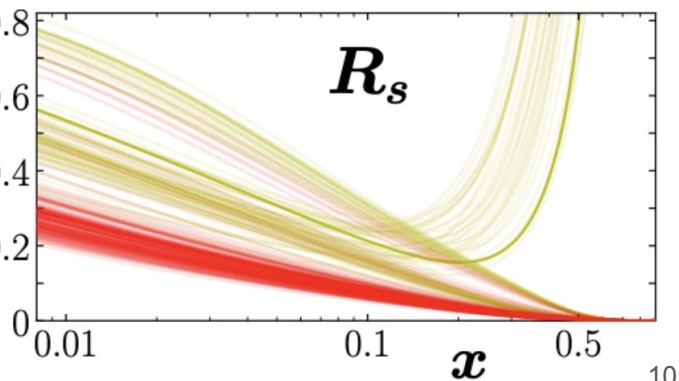
# JAM19



- Strange PDF is one of the least constrained f<sub>1</sub> pdfs
- Different analyses have inferred different sizes of the strangeness

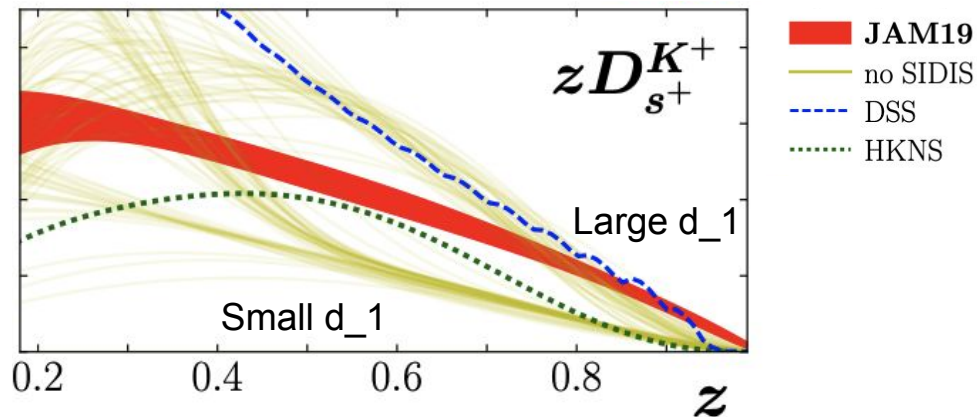
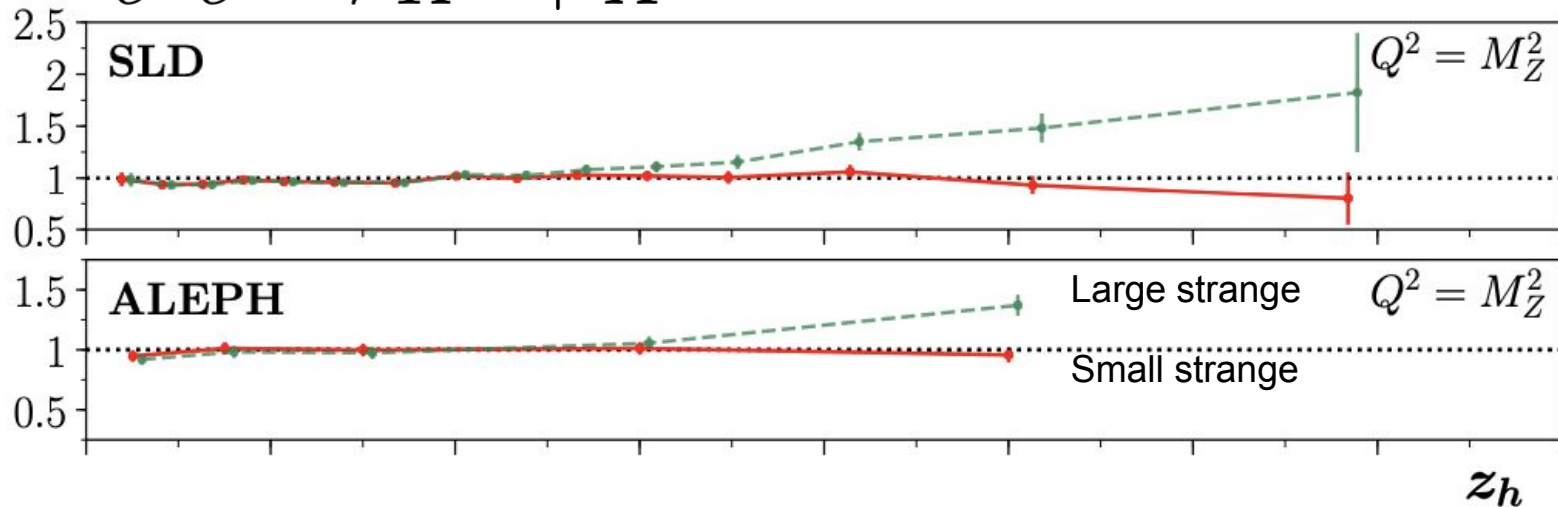
$$R_s = \frac{s + \bar{s}}{\bar{u} + \bar{d}}$$

— JAM19  
— no SIDIS/SIA

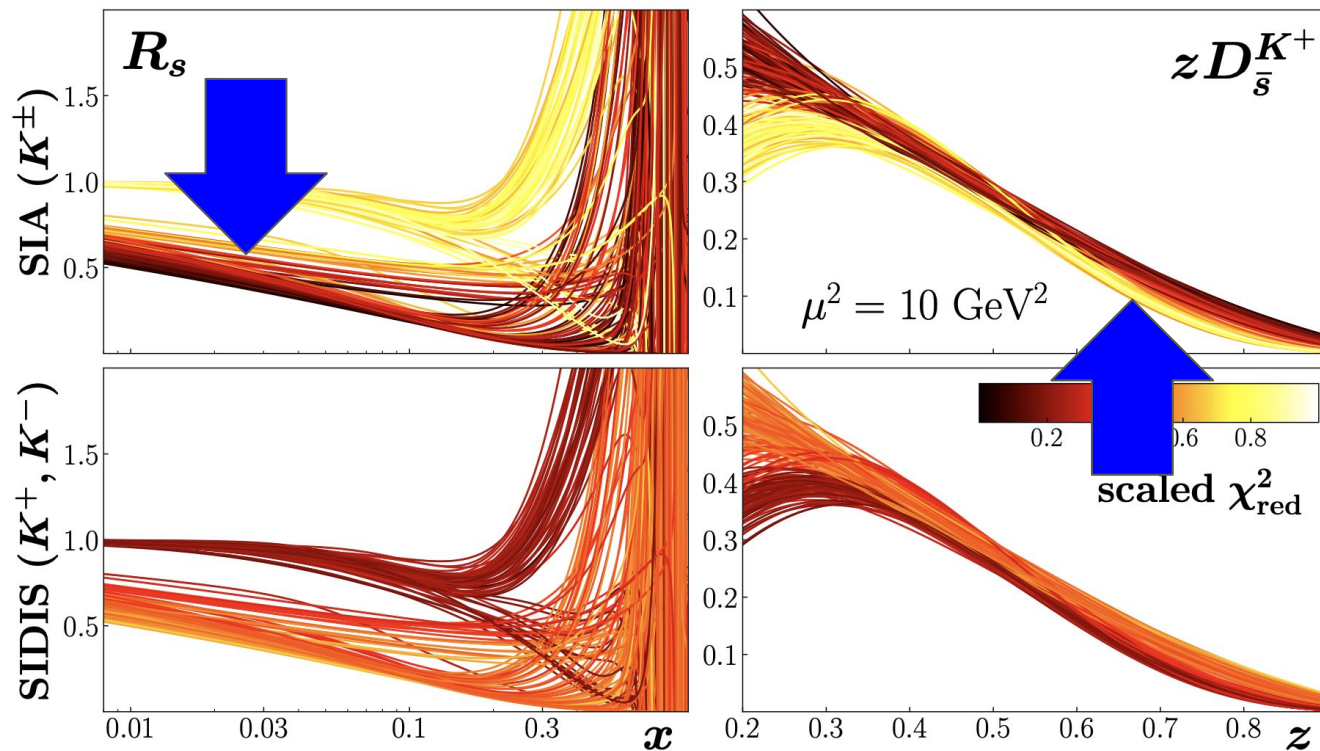


$$e^+e^- \rightarrow K^\pm + X$$

data/theory



- LEP kaon data disfavors small  $s \rightarrow K$  fragmentation
- SIDIS data compensates large strange FF by suppressing strange PDF

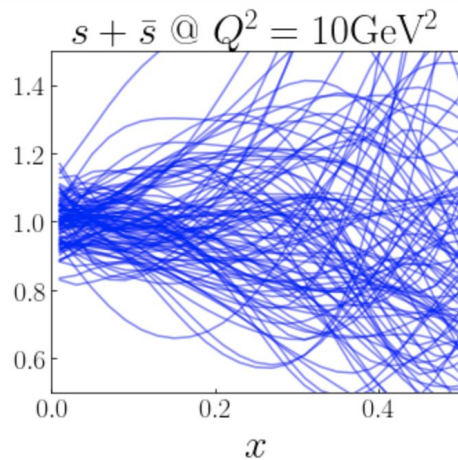
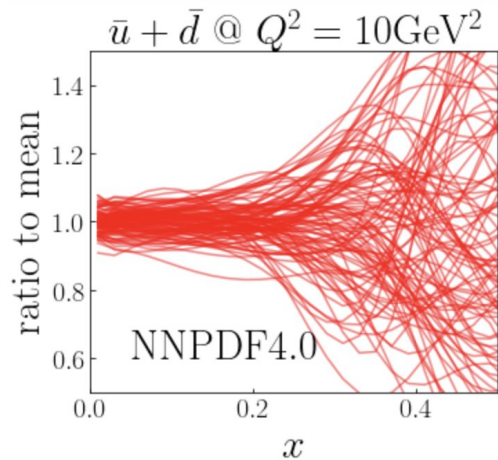
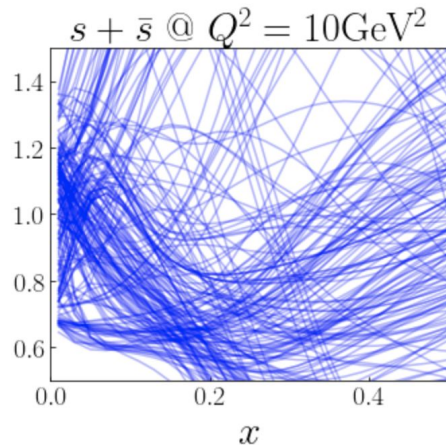
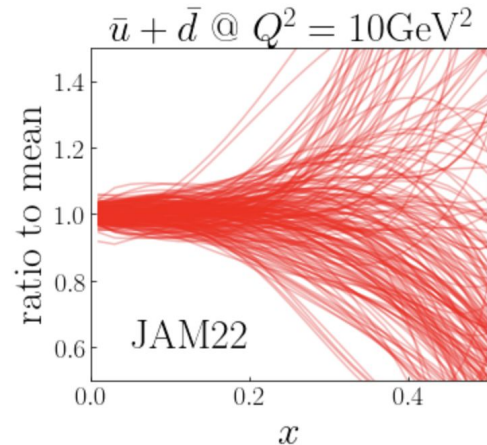


**Bottom line:**

Simultaneous analysis suggests a strong strange suppression,  
and differs from other global analyses using LHC data



# JAM22



Process	$N_{\text{dat}}$	$\chi^2/N_{\text{dat}}$
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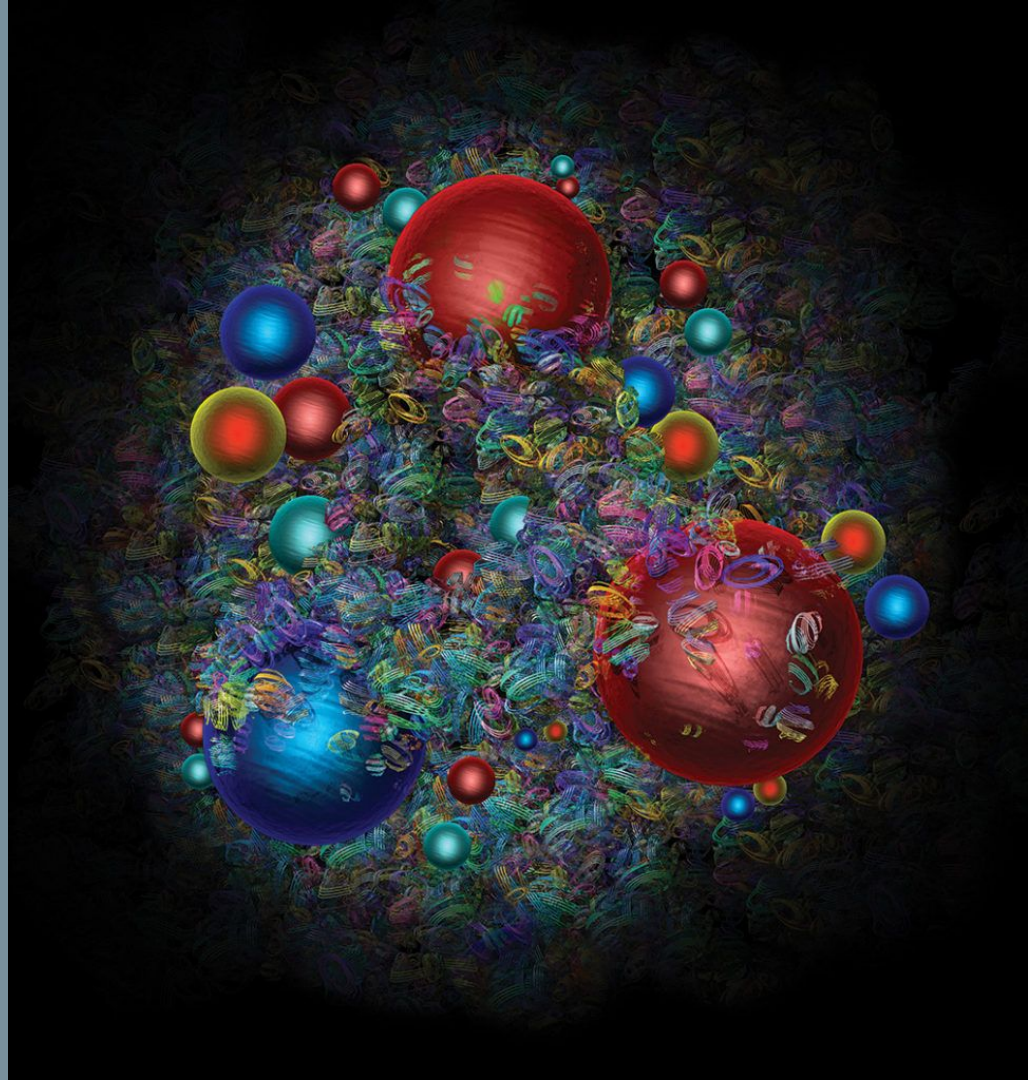
*Strangeness pdf  
is mostly **fiction!***





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# Model-independent remarks on electron-quark parity-violating neutral-current couplings

J. D. Bjorken

Stanford Linear Accelerator Center, Stanford University, Stanford, California 94305

(Received 10 July 1978)

$$\frac{A^{eD}(Q^2, \nu, y)_{AV}}{Q^2} \propto \frac{l_{\mu\nu} \int \langle D | [j^\mu(x) J^\nu(0) + J^\mu(x) j^\nu(0)] | D \rangle e^{iq \cdot x} d^4x}{l_{\mu\nu} \int \langle D | j^\mu(x) j^\nu(0) | D \rangle e^{iq \cdot x} d^4x}$$



$$\left. \frac{A^{eD}}{Q^2} \right|_{y=0} = -\frac{3G}{10\pi\alpha\sqrt{2}} \left[ 2\epsilon_{AV}(e, u) \left(1 + \frac{3}{10} \delta\right) - \epsilon_{AV}(e, d) \left(1 - \frac{6}{5} \delta\right) \right].$$

If  $\delta$  is small then  $A_{PV}$  on deuteron is highly sensitive to  $\sin^2 \theta_W$ .

$$\epsilon_{VA}(e, u) = \frac{1}{2}(1 - 4\sin^2 \theta_W),$$

$$\epsilon_{VA}(e, d) = -\frac{1}{2}(1 - 4\sin^2 \theta_W),$$

$$\epsilon_{AV}(e, u) = \frac{1}{2}\left(1 - \frac{8}{3} \sin^2 \theta_W\right),$$

$$\epsilon_{AV}(e, d) = -\frac{1}{2}\left(1 - \frac{4}{3} \sin^2 \theta_W\right).$$



# 1978 $\rightarrow$ 2014

*From currents to partons*

## Measurement of Parity-Violating Asymmetry in Electron-Deuteron Inelastic Scattering

D. Wang, R. Subedi,<sup>\*</sup> G. D. Cates, M. M. Dalton, X. Deng, D. Jones, N. Liyanage, V. Nelyubin,  
K. D. Paschke, S. Riordan, K. Saenboonruang,<sup>†</sup> R. Silwal, W. A. Tobias, and X. Zheng  
*University of Virginia, Charlottesville, Virginia 22904, USA* ■ ■ ■ ■ ■

## Measurement of Parity-Violating Asymmetry in Electron-Deuteron Inelastic Scattering

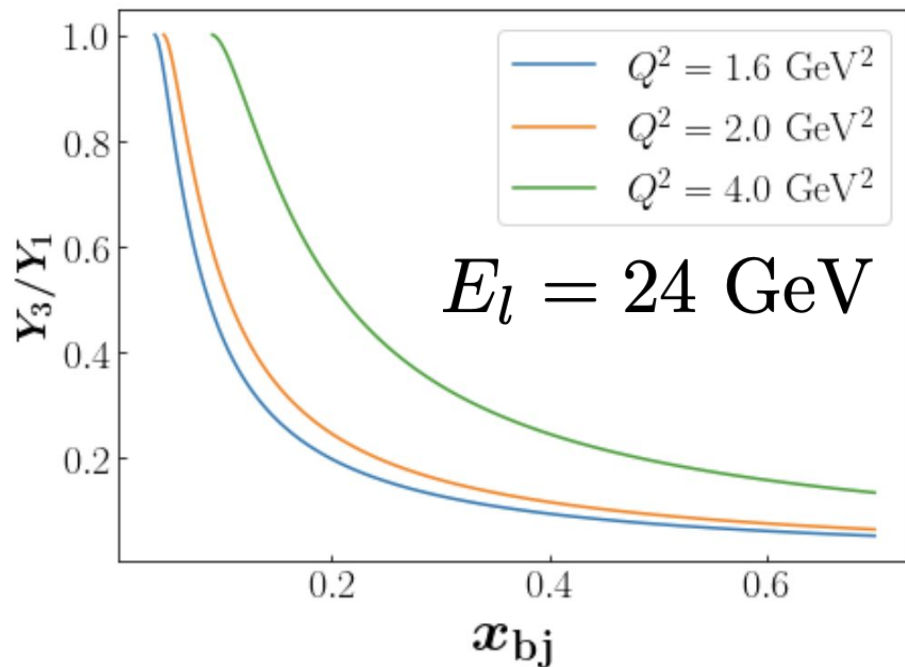
D. Wang, R. Subedi,<sup>✉</sup> G. D. Cates, M. M. Dalton, X. Deng, D. Jones, N. Liyanage, V. Nelyubin,  
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*University of Virginia, Charlottesville, Virginia 22904, USA* ■ ■ ■ ■ ■

$$A_{PV} = -\frac{G_F Q^2}{4\sqrt{2}\pi\alpha(Q^2)} \left[ \underline{a_1(x, Q^2)Y_1(x, y, Q^2)} + \underline{a_3(x, Q^2)Y_3(x, y, Q^2)} \right]$$

$$a_1(x) = 2g_A^e \frac{F_1^{\gamma Z}}{F_1^\gamma} \quad Y_1 = \left[ \frac{1 + R^{\gamma Z}}{1 + R^\gamma} \right] \frac{1 + (1 - y)^2 - \frac{y^2}{2} \left[ 1 + r^2 - \frac{2r^2}{1 + R^{\gamma Z}} \right]}{1 + (1 - y)^2 - \frac{y^2}{2} \left[ 1 + r^2 - \frac{2r^2}{1 + R^\gamma} \right]}$$
$$a_3(x) = g_V^e \frac{F_3^{\gamma Z}}{F_1^\gamma} \quad Y_3 = \left[ \frac{r^2}{1 + R^\gamma} \right] \frac{1 - (1 - y)^2}{1 + (1 - y)^2 - \frac{y^2}{2} \left[ 1 + r^2 - \frac{2r^2}{1 + R^\gamma} \right]}.$$

# y dependence at JLab 20+ GeV

$$A_{PV} = -\frac{G_F Q^2}{4\sqrt{2}\pi\alpha(Q^2)} \left[ \underline{a_1(x, Q^2)Y_1(x, y, Q^2)} + \underline{a_3(x, Q^2)Y_3(x, y, Q^2)} \right]$$



$$Y_1 = \left[ \frac{1 + R^2}{1 + R^2 y} \right] \frac{1 + (1-y)^2 - \frac{y^2}{2} \left[ 1 + r^2 - \frac{2r^2}{1+R^2 y} \right]}{1 + (1-y)^2 - \frac{y^2}{2} \left[ 1 + r^2 - \frac{2r^2}{1+R^2 y} \right]}$$

$$Y_3 = \left[ \frac{r^2}{1 + R^2} \right] \frac{1 - (1-y)^2}{1 + (1-y)^2 - \frac{y^2}{2} \left[ 1 + r^2 - \frac{2r^2}{1+R^2 y} \right]}$$

$$r^2 = 1 + \frac{Q^2}{\nu^2}$$



## Parton level view...

$$a_1(x) = 2g_A^e \frac{F_1^{\gamma Z}}{F_1^\gamma} \qquad a_3(x) = g_V^e \frac{F_3^{\gamma Z}}{F_1^\gamma}$$

At LO in QCD we have

$$\begin{aligned} F_1^\gamma(x, Q^2) &= \frac{1}{2} \sum Q_{q_i}^2 [q_i(x, Q^2) + \bar{q}_i(x, Q^2)] \\ F_1^{\gamma Z}(x, Q^2) &= \sum Q_{q_i} g_V^i [q_i(x, Q^2) + \bar{q}_i(x, Q^2)] \\ F_3^{\gamma Z}(x, Q^2) &= 2 \sum Q_{q_i} g_A^i [q_i(x, Q^2) - \bar{q}_i(x, Q^2)] \end{aligned}$$

$$a_1(x) = 2g_A^e \frac{F_1^{\gamma Z}}{F_1^\gamma}$$



$$a_1 = \frac{2(c^+ + u^+)(8s_w^2 - 3) + (d^+ + s^+)(4s_w^2 - 3)}{4c^+ + d^+ + s^+ + 4u^+}$$



$$a_1 = 4s_w^2 - \frac{9}{5}$$

If we ignore s and c and  
use deuteron target

$$a_3(x) = g_V^e \frac{F_3^{\gamma Z}}{F_1^\gamma}$$

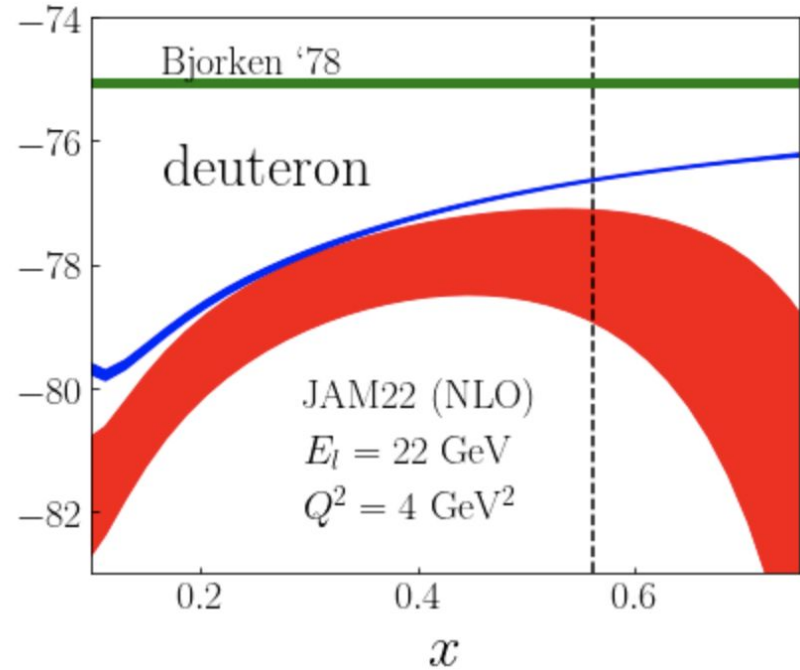
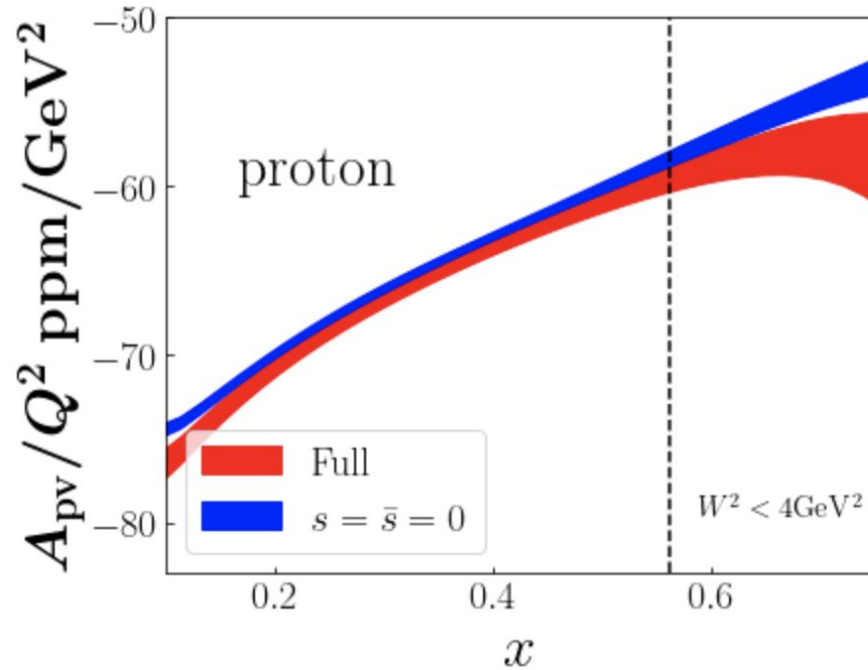


$$a_3 = \frac{3(4s_w^2 - 1)(2c^- + d^- + s^- + 2u^-)}{4c^+ + d^+ + s^+ + 4u^+}$$



$$a_3 = \frac{3(d^- + 2u^-)(4s_w^2 - 1)}{5u^+}$$

# Apv with full QCD theory @ NLO



**Bottom line:** Apvd uncertainties correlates significantly from strangeness

# Outline

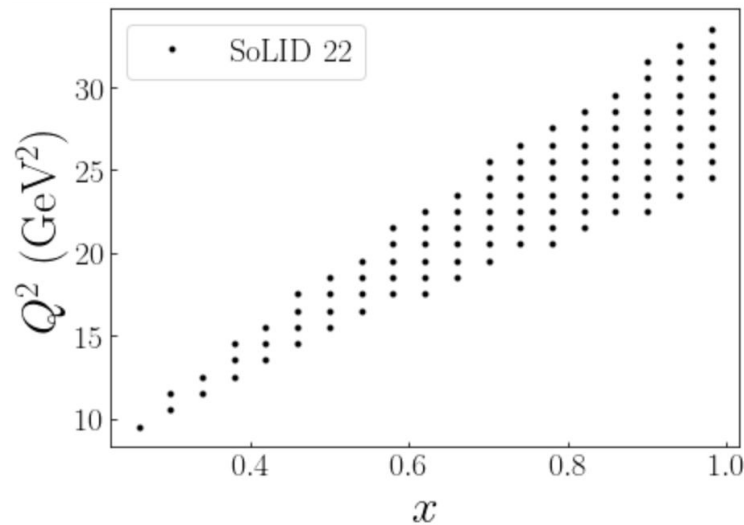
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# Simulation at SoLID 22

	Q2min	Q2max	xmin	xmax	acceptance	xmid	Q2mid	W2	rs
1	9.0	10.0	0.24	0.28	0.150635	0.26	9.5	27.9183055384...	6.42432...
2	10.0	11.0	0.28	0.32	0.208101	0.300000...	10.5	25.3798439999...	6.42432...
3	11.0	12.0	0.28	0.32	0.137562	0.300000...	11.5	27.7131773333...	6.42432...
4	11.0	12.0	0.32	0.36	0.226112	0.339999...	11.5	23.2033734117...	6.42432...
5	12.0	13.0	0.32	0.36	0.267877	0.339999...	12.5	25.1445498823...	6.42432...
6	12.0	13.0	0.36	0.4	0.227464	0.38	12.5	21.2745808421...	6.42432...
7	13.0	14.0	0.36	0.4	0.355759	0.38	13.5	22.9061597894...	6.42432...
8	14.0	15.0	0.36	0.4	0.185517	0.38	14.5	24.5377387368...	6.42432...
9	13.0	14.0	0.4	0.44	0.220877	0.420000...	13.5	19.5227011428...	6.42432...
10	14.0	15.0	0.4	0.44	0.377228	0.420000...	14.5	20.9036535238...	6.42432...
11	15.0	16.0	0.4	0.44	0.305199	0.420000...	15.5	22.2846059047...	6.42432...
12	14.0	15.0	0.44	0.48	0.218991	0.459999...	14.5	17.9015831304...	6.42432...

Thanks to: T. Liu & M. Dalton for providing the table



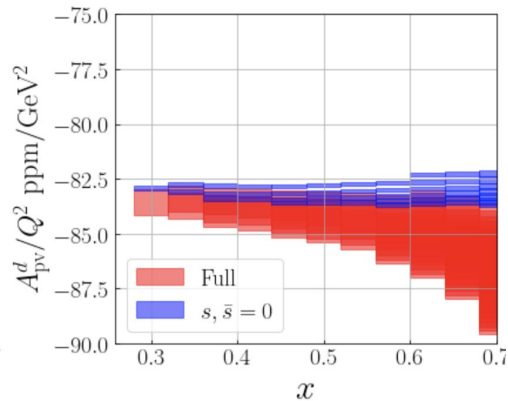
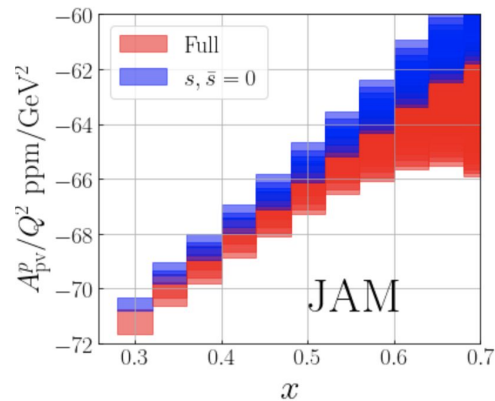
$$N_{\text{bin}} = \mathcal{L} \left. \frac{d\sigma}{dx dQ^2} \right|_{\text{mid}} \times \Delta x \Delta Q^2$$

$$\delta A_{\text{pv}} = \frac{1}{\sqrt{N_{\text{bin}}}}$$

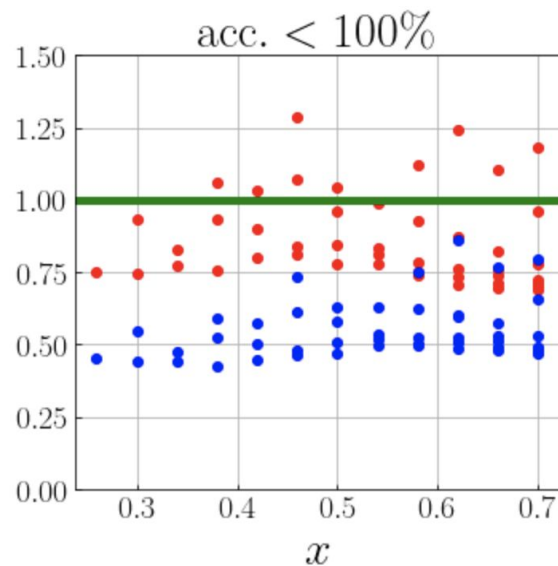
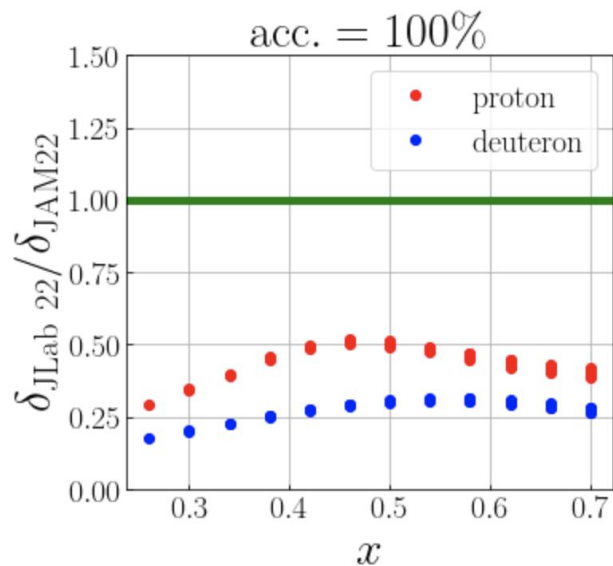
$$\mathcal{L} = 5.0 \times 10^{38} \text{cm}^{-2} \text{s}^{-1}$$

200 days of data taking





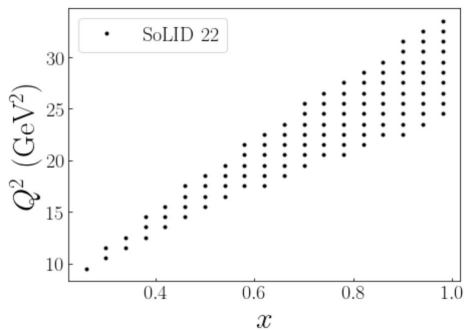
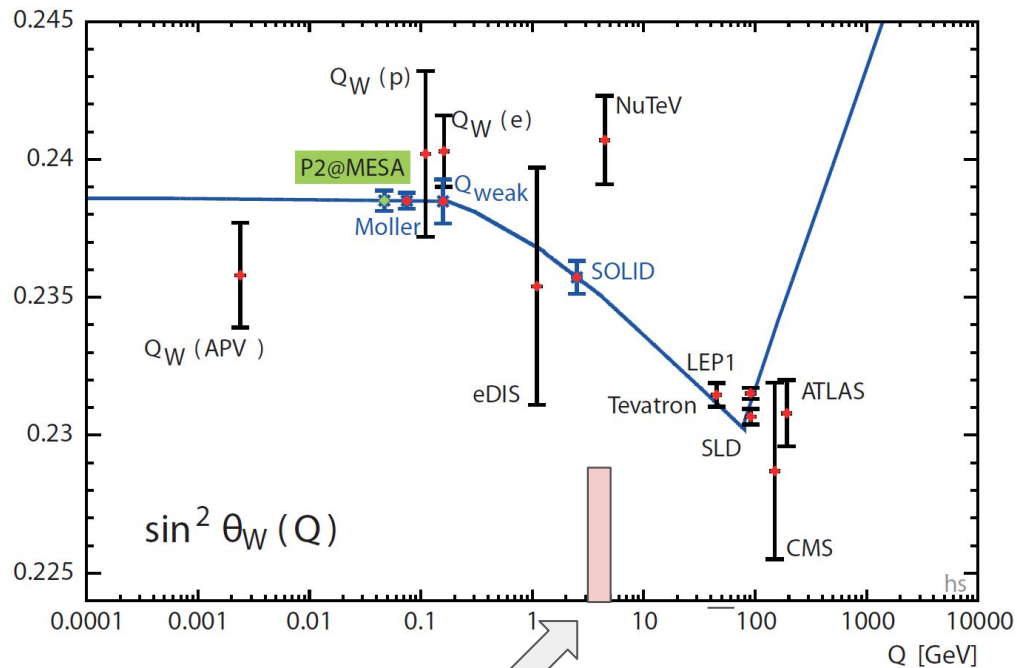
Spread  
represents PDF  
uncertainties



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3. Parity violating DIS
4. PVDIS @ SoLID -22
5.  $\sin^2\theta_W$
6. QED radiative effects
7. Integrating thy/exp/cs





**In principle there are two unknowns:**

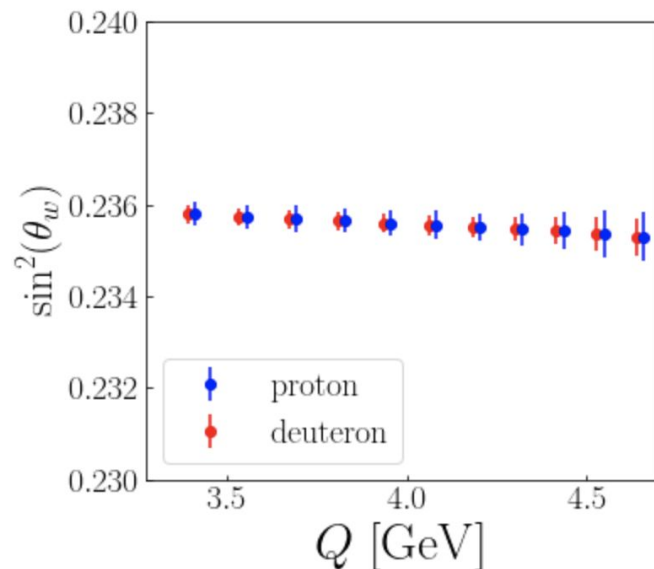
- 1) The running of sin2w as a function of scale
- 2) The value of sin2w at a given scale

# Point-by-point extraction of $\sin^2\theta_W$

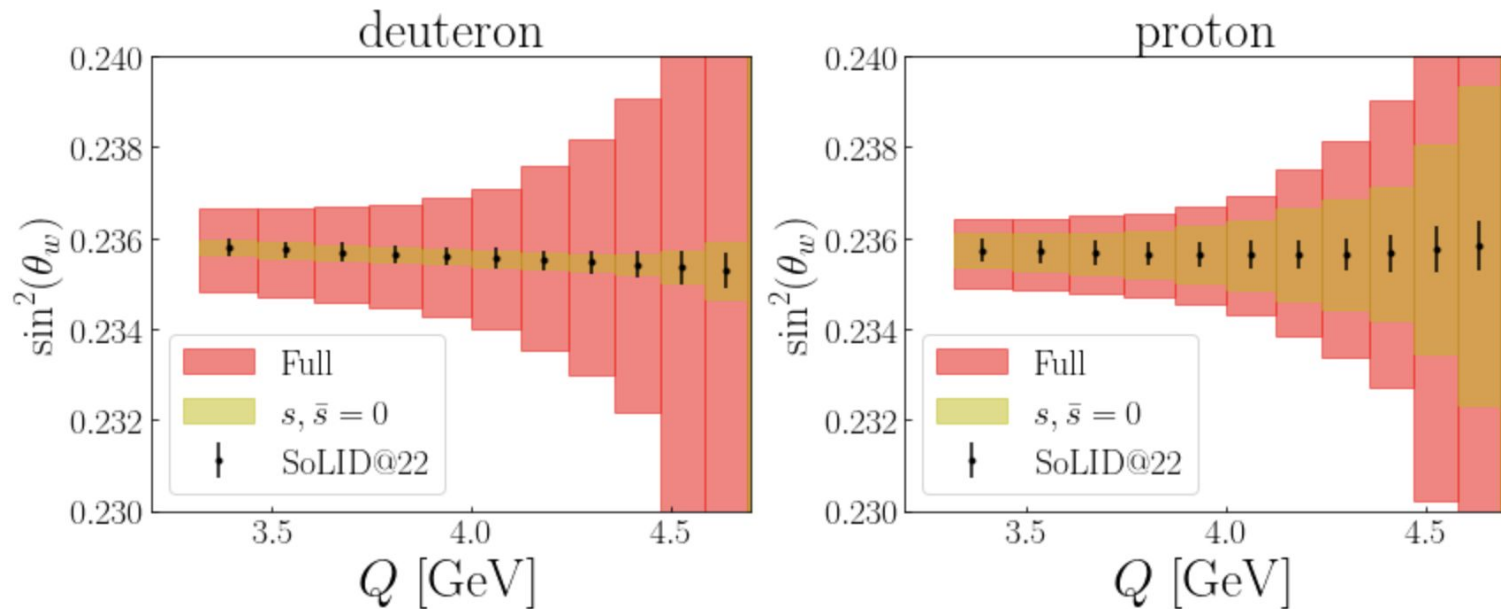
$$\sin^2 \theta_W(Q) = \frac{2\sqrt{2}\pi A_{pv} F_1^{gg} \alpha}{F_3^{gZ} G_F Q^2 Y_3} - \frac{F_1^{gZ} Y_1 g_A}{F_3^{gZ} Y_3} + \frac{1}{4}$$

## Procedure:

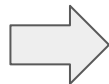
- Measure  $A_{pv}$
- Use existing knowledge of structure functions to solve for  $\sin^2\theta_W$  point-by-point in  $Q$
- Combine different  $x$  values using linear regression.
- Propagate uncertainties from  $A_{pv}$  using gaussian approximation



# Taking into account PDF errors



**Bottom line:** using protons and deuteron, one can solve for the unknowns eg strangeness +  $\sin 2w$

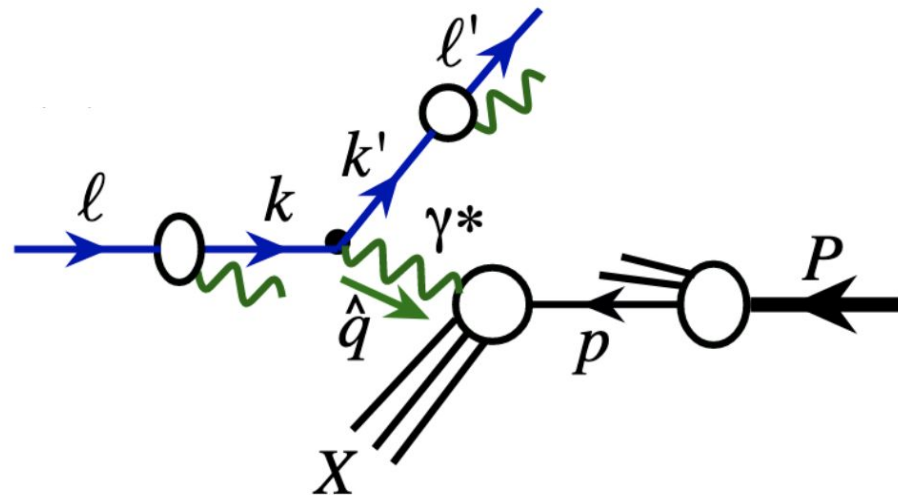


Simultaneous fit of pdfs and eweak parameters!



# Outline

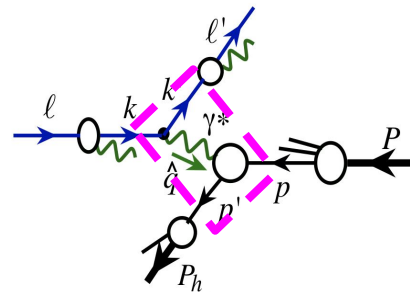
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Factorized approach to radiative corrections for inelastic lepton-hadron collisions

Tianbo Liu, W. Melnitchouk, Jian-Wei Qiu, and N. Sato  
Phys. Rev. D **104**, 094033 – Published 19 November 2021

# Hybrid QED+QCD factorization



$$E' \frac{d\sigma_{\text{DIS}}}{d^3\ell'} = \frac{1}{2s} \sum_{ija} \int_{z_L}^1 \frac{d\zeta}{\zeta^2} \int_{x_L}^1 \frac{d\xi}{\xi} \underline{D_{e/j}(\zeta) f_{i/e}(\xi)}$$

Collinear LDFs  
and LFFs

$$\times \int_{x_h}^1 \frac{dx}{x} f_{a/N}(x) \boxed{\hat{H}_{ia \rightarrow j}}(\xi, \zeta, x) + \mathcal{O}\left(\frac{1}{\ell'^2_T}\right),$$

Collinear PDFs

Short distance  
hard part

# Collinear LDFs and LFFs

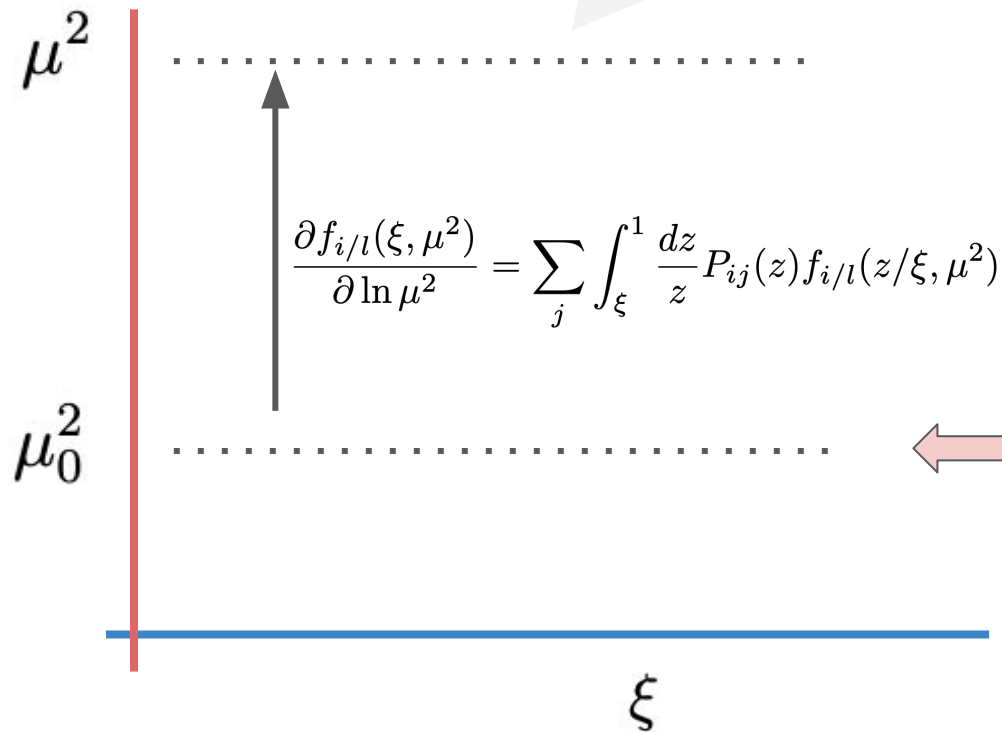
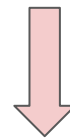
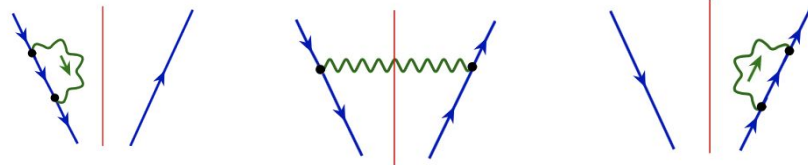
$$f_{i/e}(\xi) = \int \frac{dz^-}{4\pi} e^{i\xi\ell^+ z^-} \langle e | \bar{\psi}_i(0) \gamma^+ \Phi_{[0,z^-]} \psi_i(z^-) | e \rangle$$

$$D_{e/j}(\zeta) = \frac{\zeta}{2} \sum_X \int \frac{dz^-}{4\pi} e^{i\ell'^+ z^- / \zeta} \text{Tr} [\gamma^+ \langle 0 | \bar{\psi}_j(0) \Phi_{[0,\infty]} | e, X \rangle \langle e, X | \psi_j(z^-) \Phi_{[z^-,\infty]} | 0 \rangle].$$

perturbatively calculable if we  
neglect hadronic components

# RGE

Resummation of collinear logs



$$f_{e/e}^{(0)}(\xi) = \delta(\xi - 1)$$

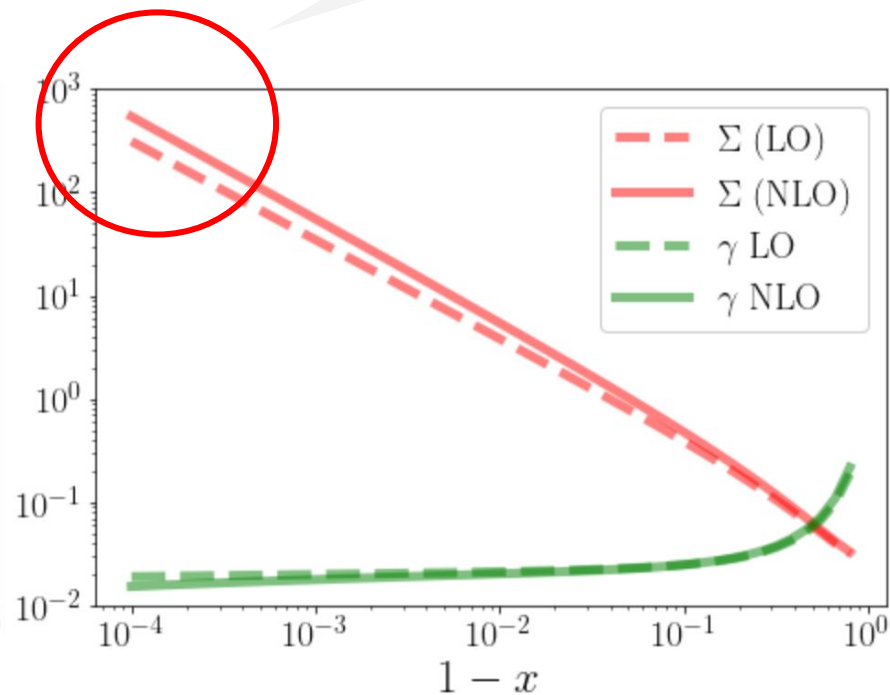
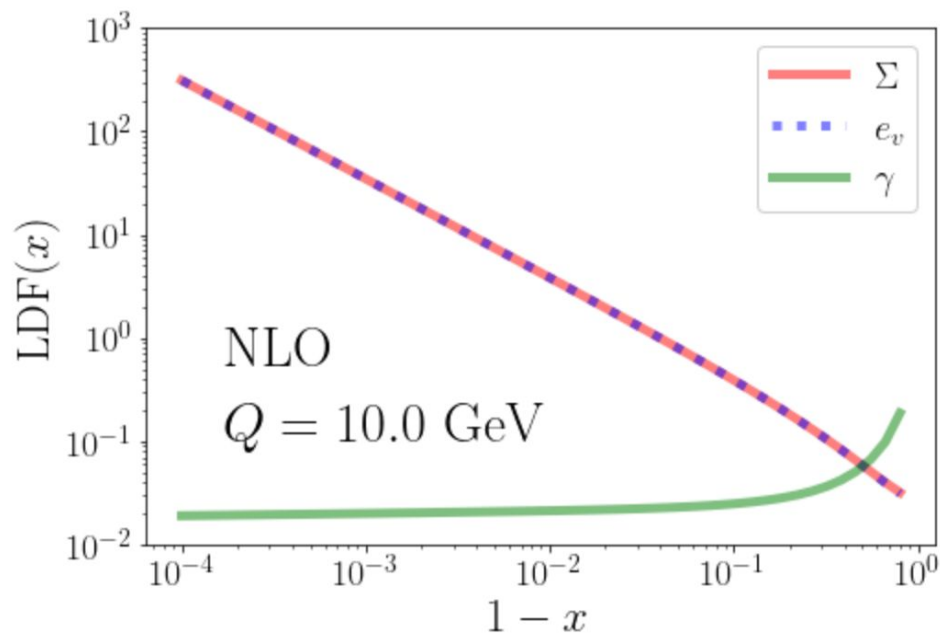
$$f_{e/e}^{(1)}(\xi, \mu_0^2) = \frac{\alpha}{2\pi} \left[ \frac{1 + \xi^2}{1 - \xi} \ln \frac{\mu_0^2}{(1 - \xi)^2 m_e^2} \right]_+$$

$$D_{e/e}^{(0)}(\zeta) = \delta(\zeta - 1)$$

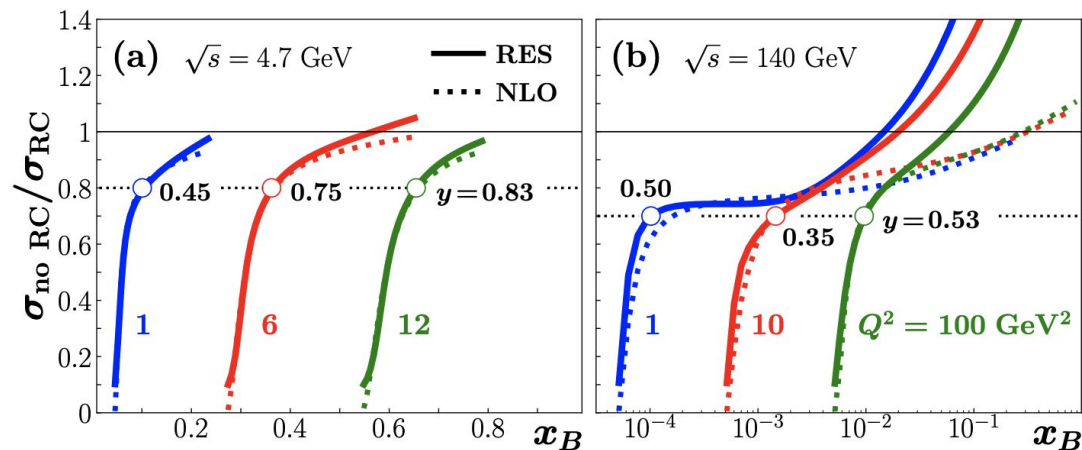
$$D_{e/e}^{(1)}(\zeta, \mu_0^2) = \frac{\alpha}{2\pi} \left[ \frac{1 + \zeta^2}{1 - \zeta} \ln \frac{\zeta^2 \mu_0^2}{(1 - \zeta)^2 m_e^2} \right]_+$$

# Evolution effects

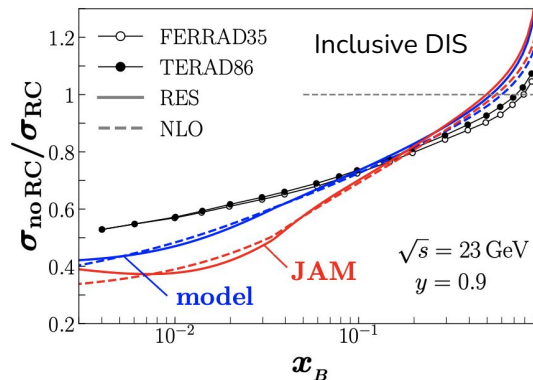
LDFs peaks at the endpoint



# Some examples from inclusive DIS



Liu, Melnitchouk, Qiu, Sato ('20, '21)

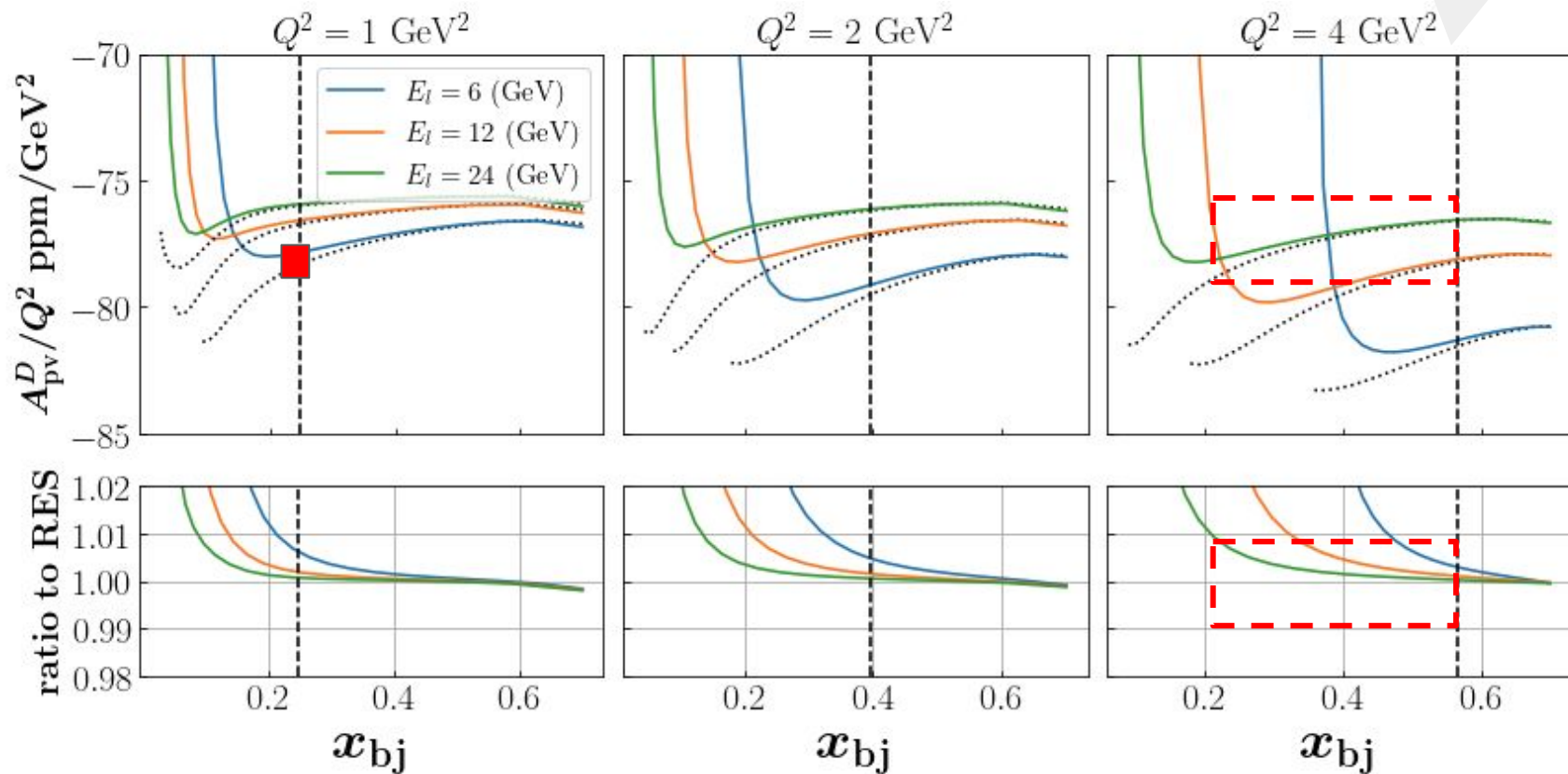


- QED “corrections” depend on the input hadronic tensor
- Not possible to construct model-independent QED RC corrections
- Need to include QED in global analysis



# Apv with QED effects

W2 > 4 GeV2

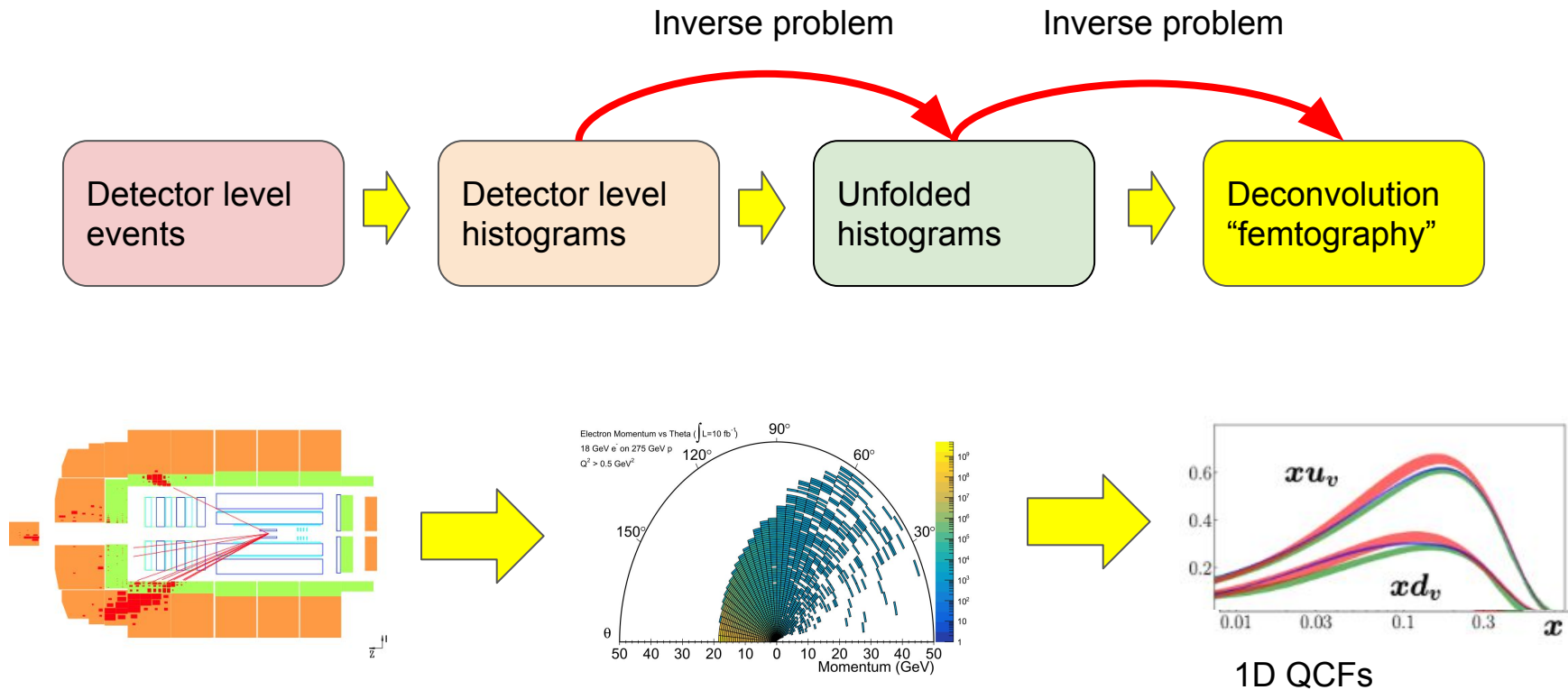


# Outline

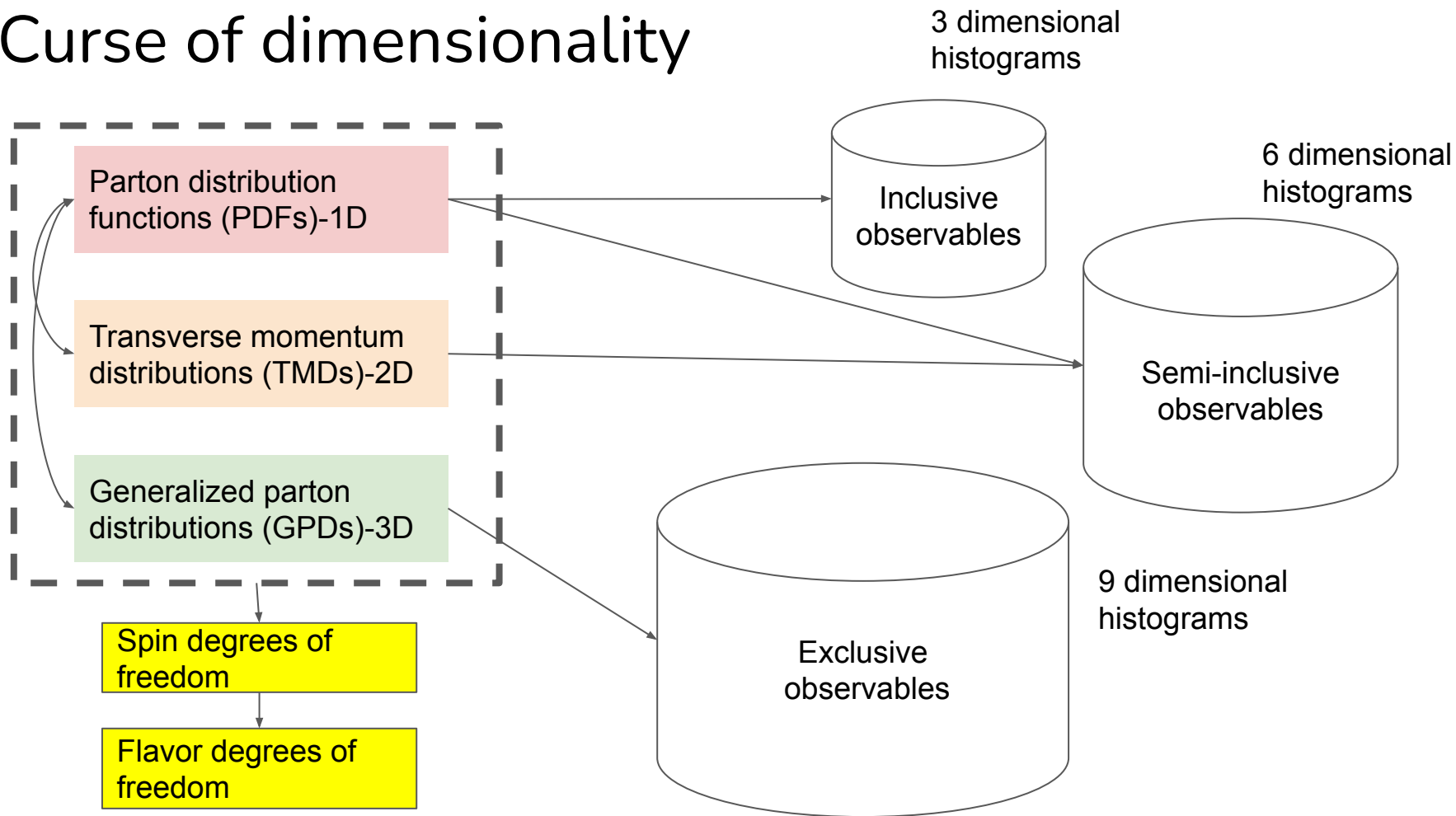
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# Existing paradigm -> histogram approach



# Curse of dimensionality

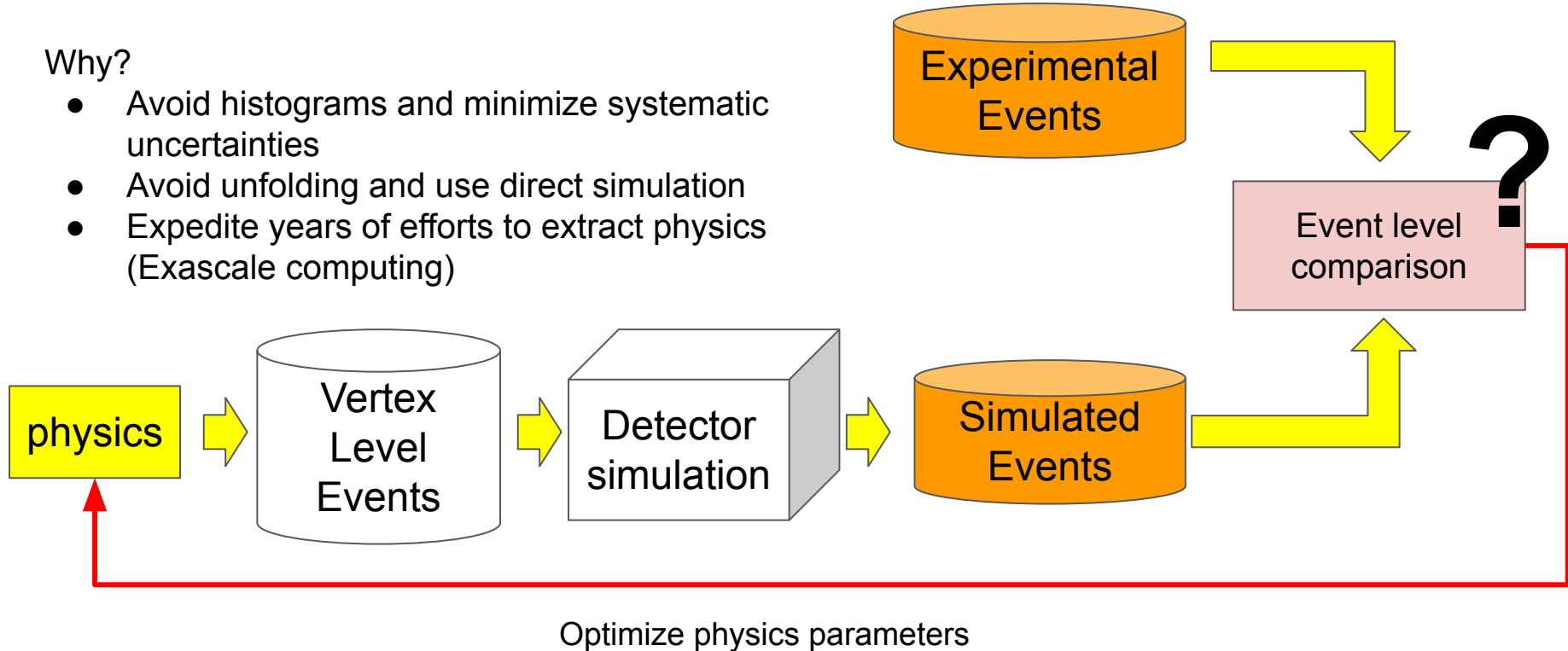


# Event-based analysis for global analysis?

Can we compare real vs synthetic events?

Why?

- Avoid histograms and minimize systematic uncertainties
- Avoid unfolding and use direct simulation
- Expedite years of efforts to extract physics (Exascale computing)





# QuantOm Collaboration



Jefferson Lab



Argonne  
NATIONAL LABORATORY



## Theory

- Jianwei Qiu (PI)
- NS
- Adam Freez (postdoc)

## Experiment

- Markus Diefenthaler
- Daniel Lersch

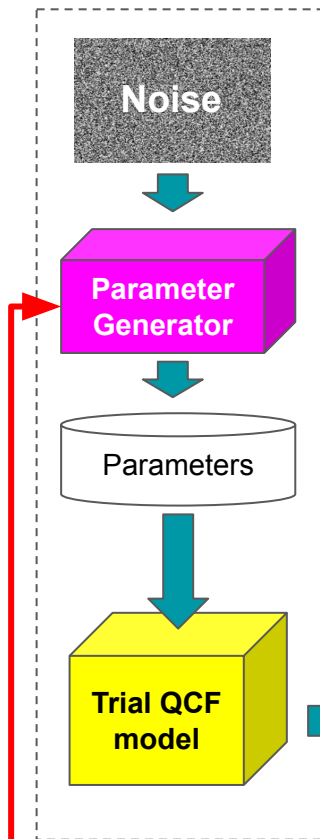
## CS

- Malachi Schram
- Kishan Rajput
- Daniel Lersch
- Yasir Alanazi (postdoc)

**Supported by DOE SciDAC**



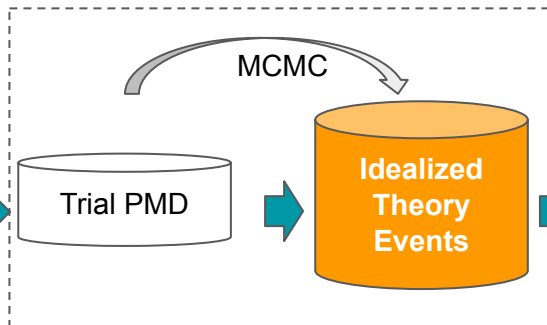
## Module 1



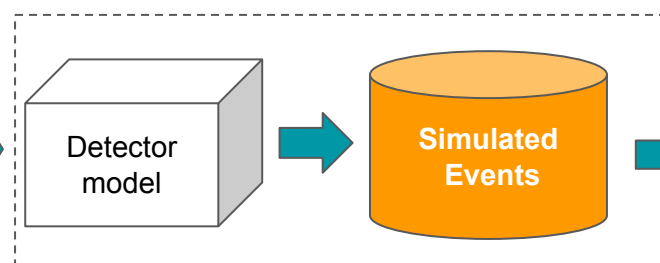
# Event-level QCF inference framework



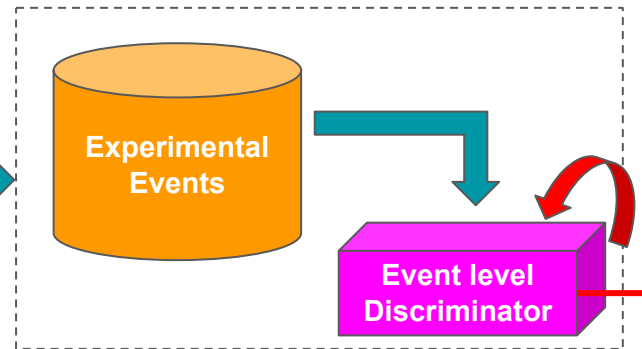
## Module 2



## Module 3



## Module 4



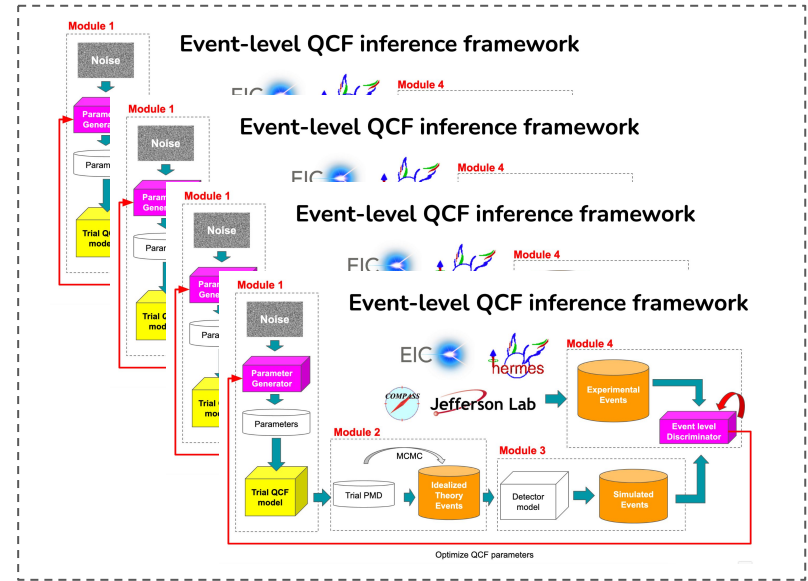
Optimize QCF parameters

## Opportunities

- Unified Theory+Exp analysis framework for hadron structure -> paradigm shift
- Near real time analysis and expedite scientific discovery

## Challenges

- Big event level data processing from JLab/EIC requires large scale computing -> exascale computing
- Dedicated distributed ML workflow needs to be developed



# An example with a proxi theory

$$u(x) = N_u x^{a_u} (1-x)^{b_u}$$

$$d(x) = N_d x^{a_d} (1-x)^{b_d}$$

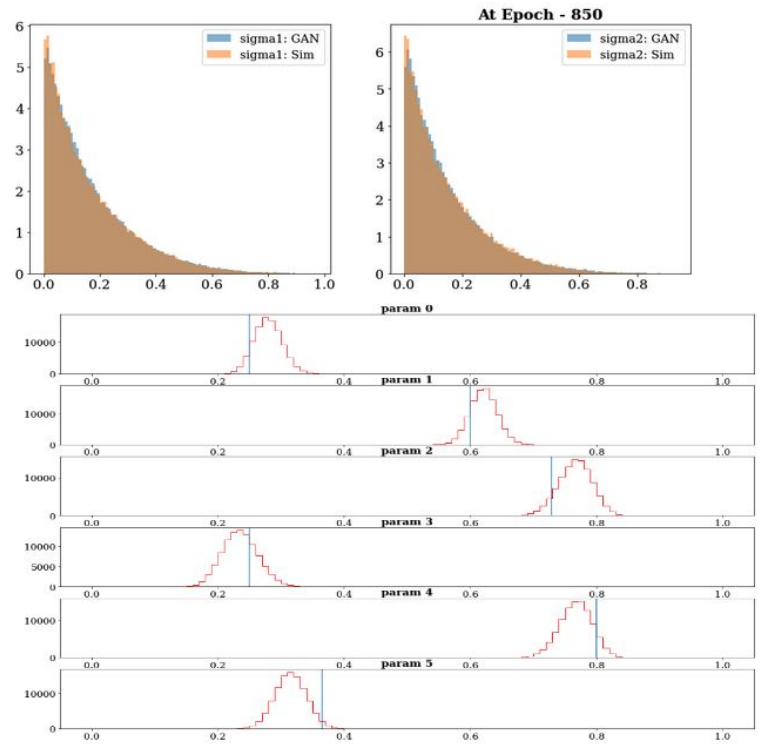
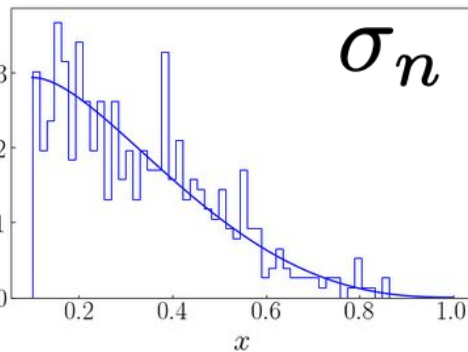
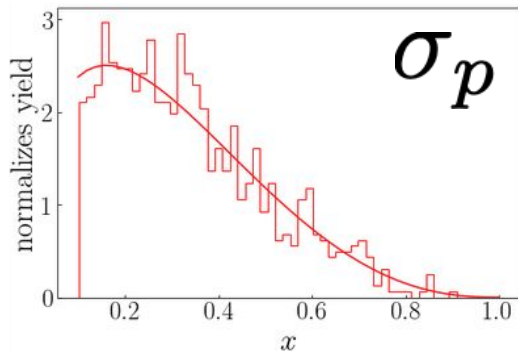
6 parameter  
QCFs model



$$\sigma_p(x) = 4u(x) + d(x),$$

$$\sigma_n(x) = 4d(x) + u(x).$$

Proxi cross  
sections

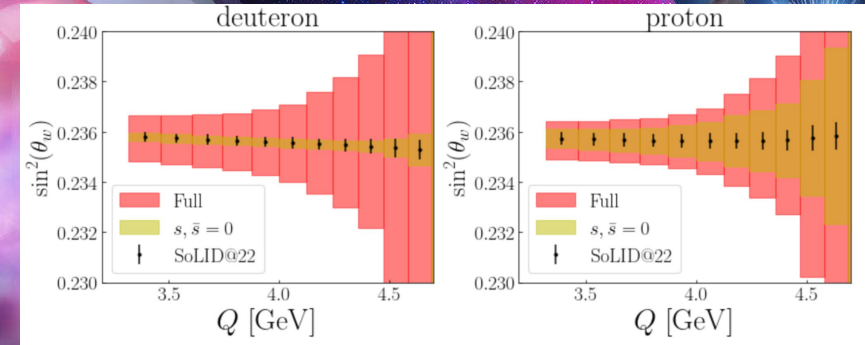
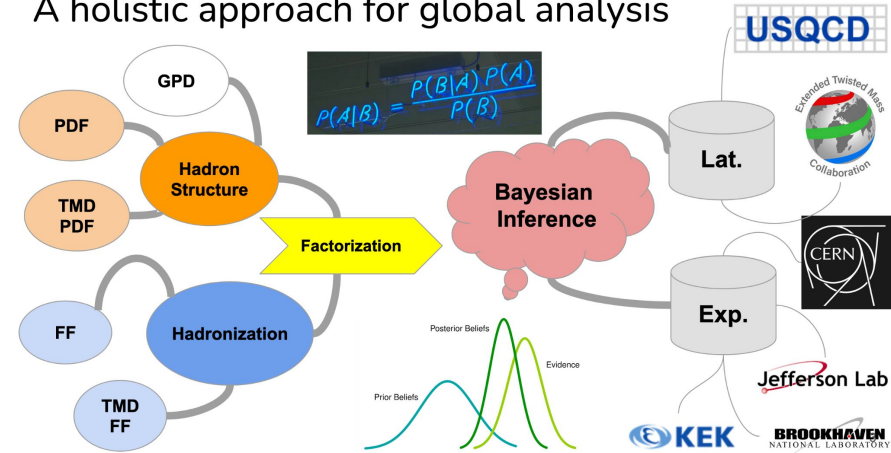


Workflow can produce parameter  
samples close to the truth

# Summary

- Apv program is critical to map finer details of hadron structure -> the strange sector of the nucleon
- Precision electroweak studies requires the use of simultaneous global analysis framework
- We really need JLab 22!

A holistic approach for global analysis



$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{\psi}_q (i\gamma_\mu D^\mu - m_q) \psi_q - \frac{1}{2} \text{Tr}[G_{\mu\nu} G^{\mu\nu}]$$