

DOUBLE-REGGE CONTRIBUTION TO $\eta^{(\prime)}\pi$ PRODUCTION

Sunset in Madrid

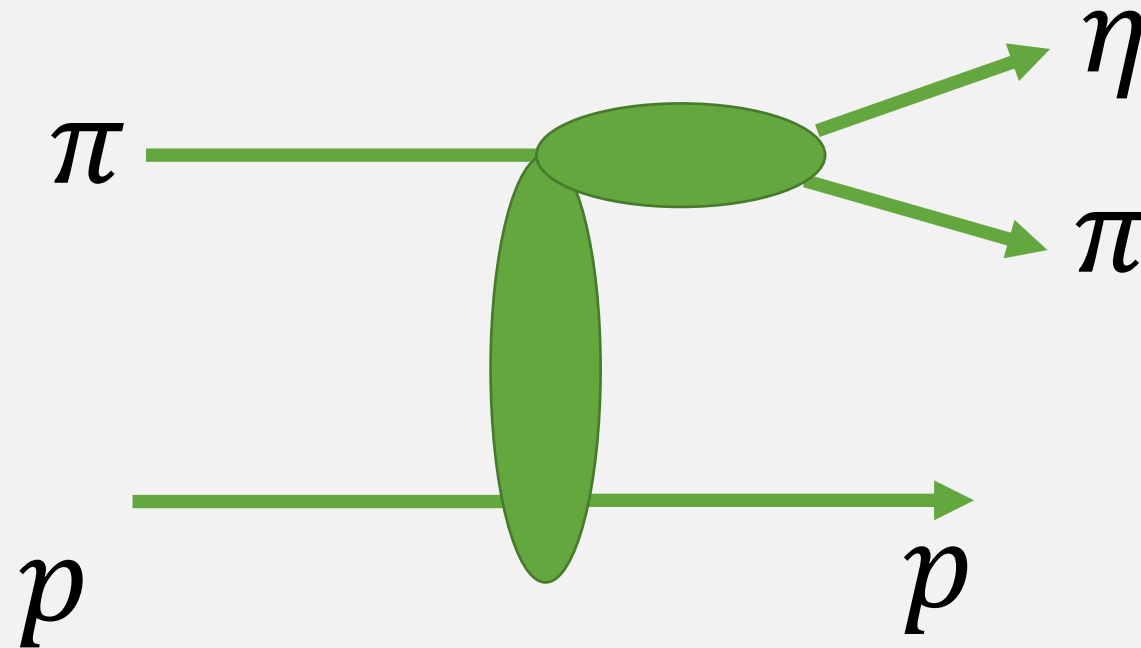
César Fernández Ramírez

#SOMOS2030



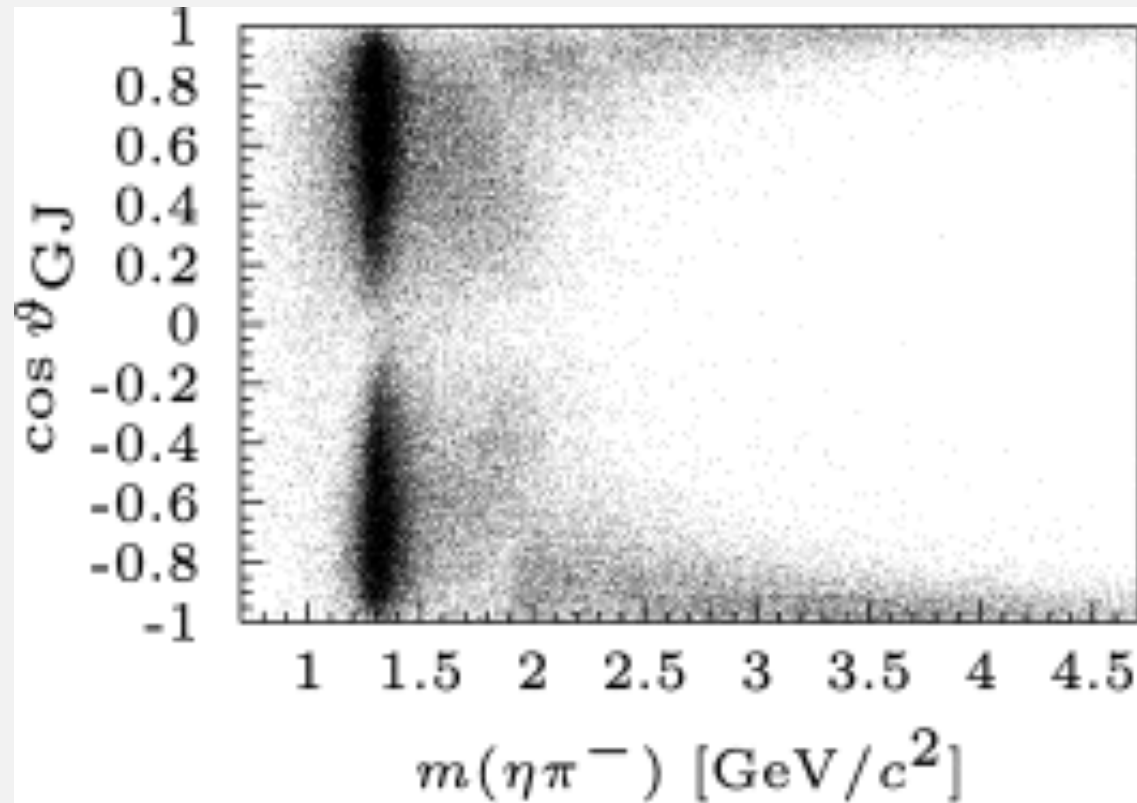
Pion diffractive dissociation

$$\pi^{-}(q) + p(p_1) \rightarrow \eta^{(\prime)}(k_{\eta}) + \pi^{-}(k_{\pi}) + p(p_2)$$

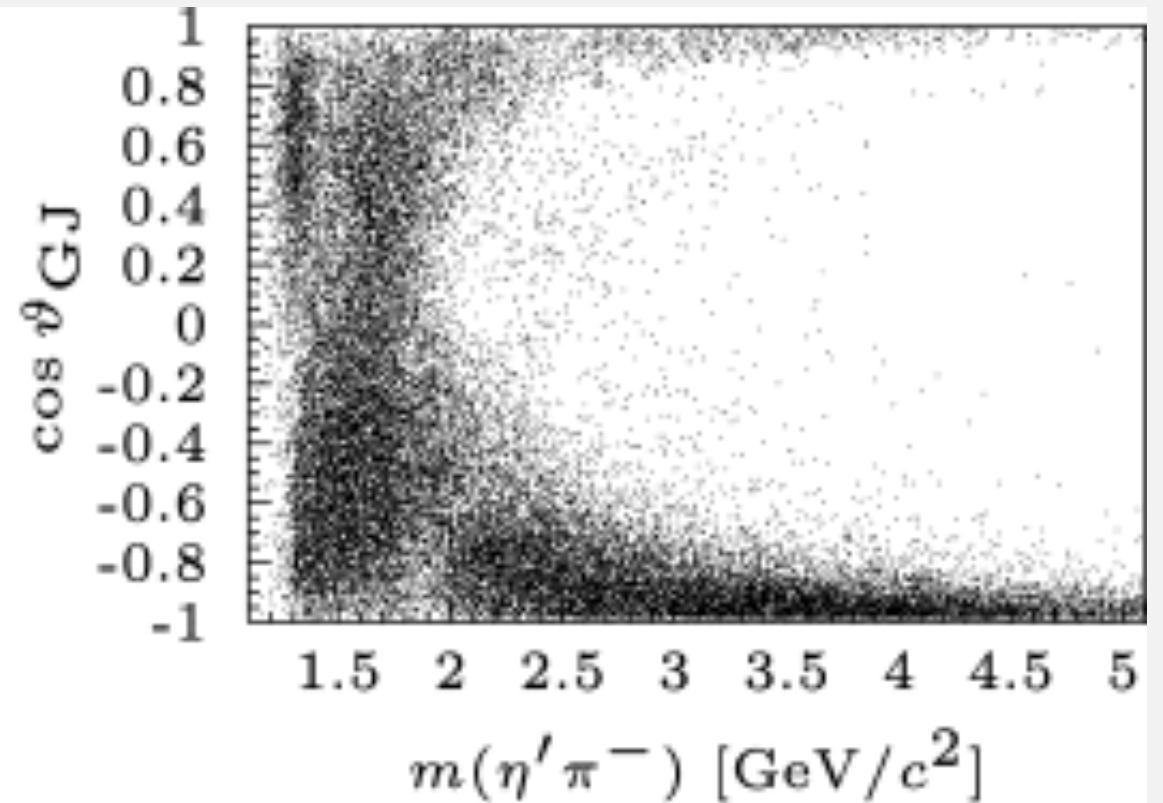


Check out Boris' talk

COMPASS plot

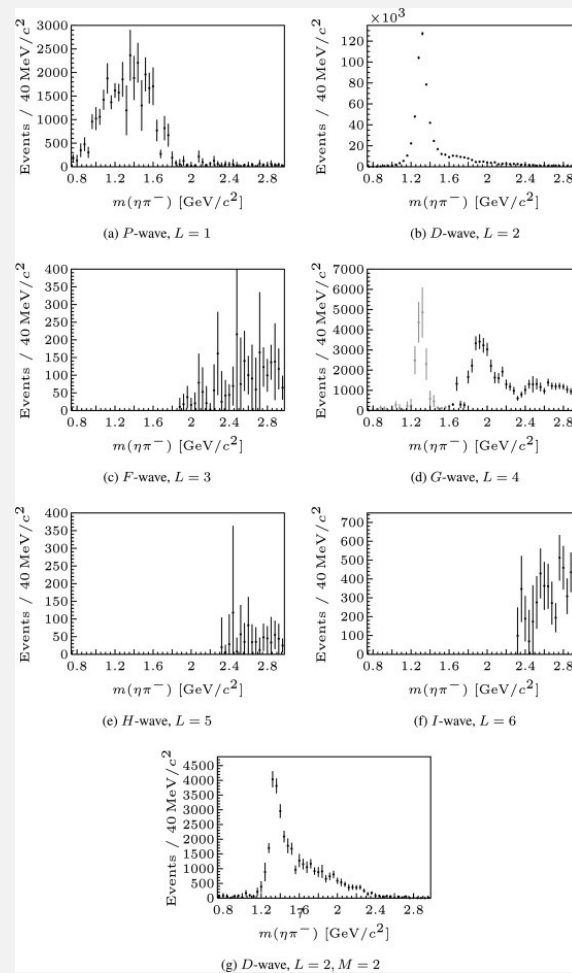
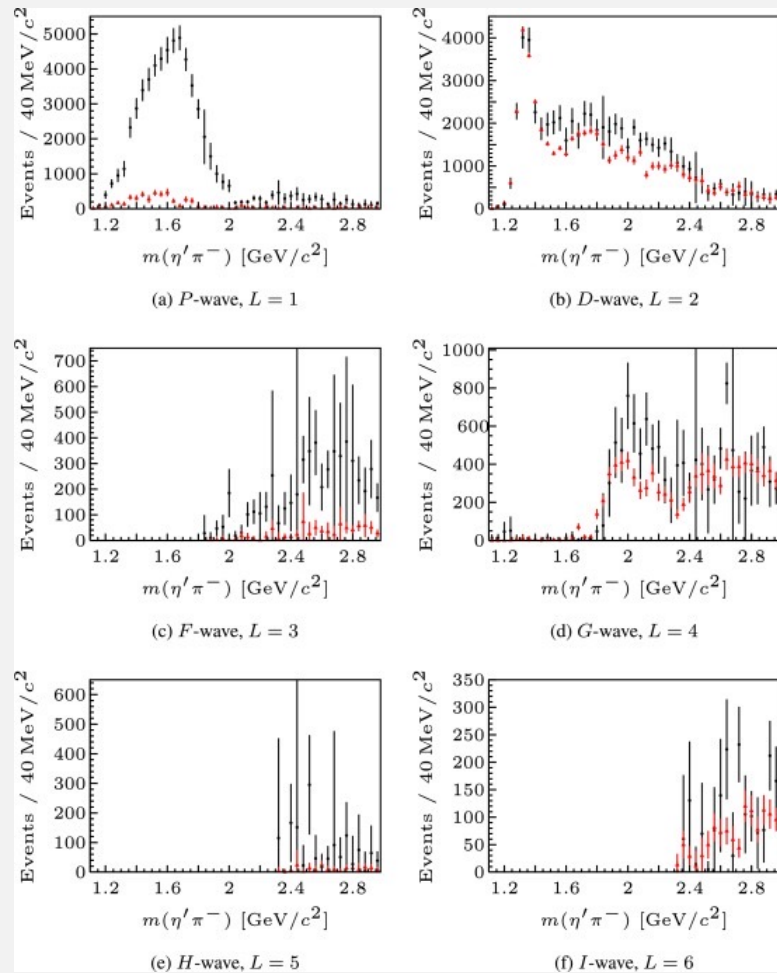


(a) $m(\eta\pi^-)$ vs. $\cos \vartheta_{GJ}$



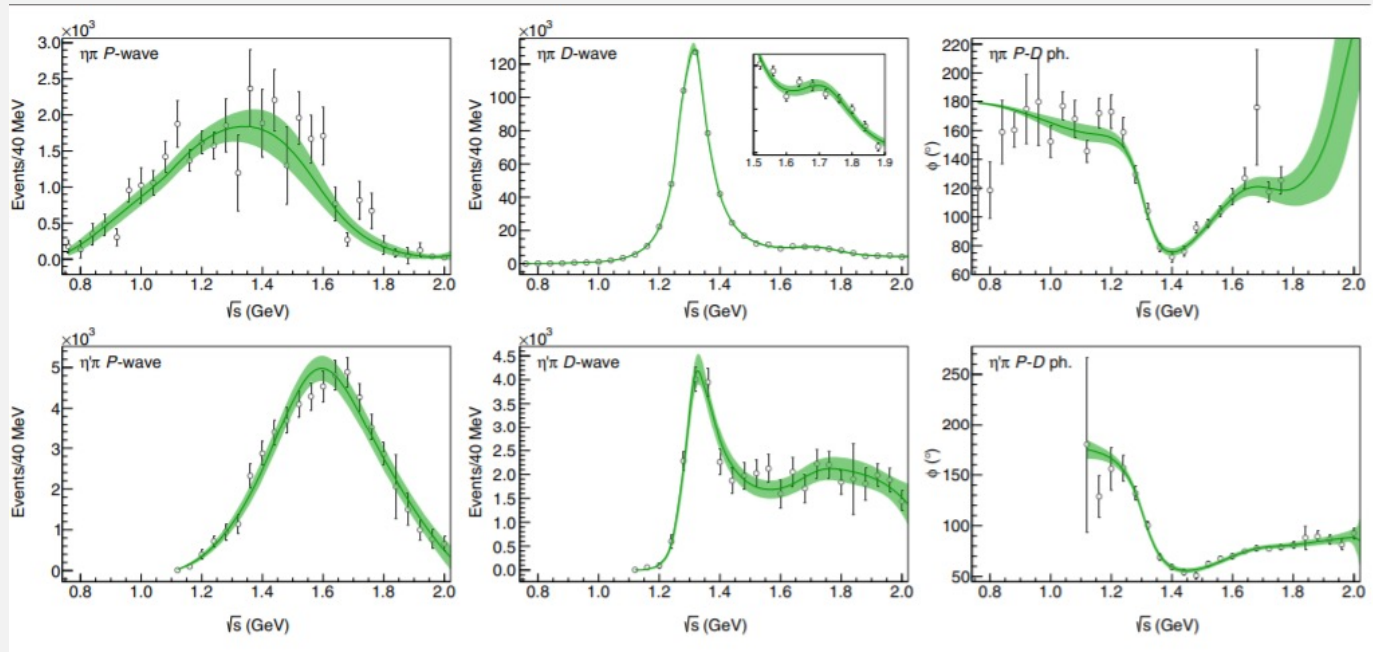
(b) $m(\eta'\pi^-)$ vs. $\cos \vartheta_{GJ}$

“Partial waves”

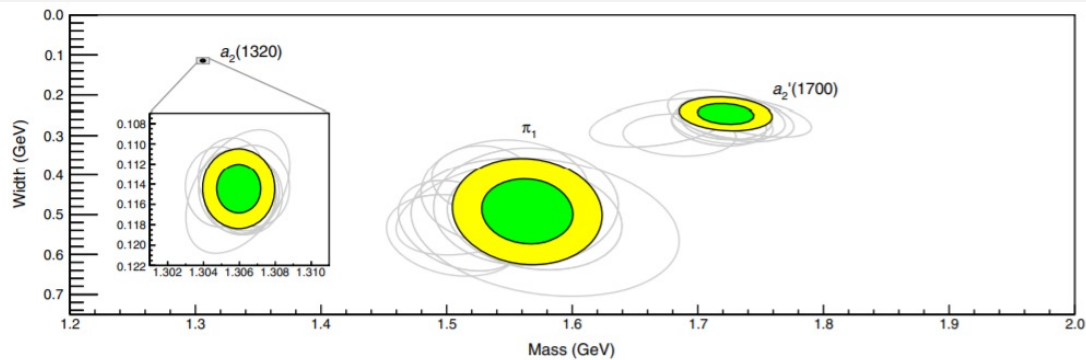


COMPASS, PLB 740 303 (2015)

Exotic in the resonance (low-energy) region

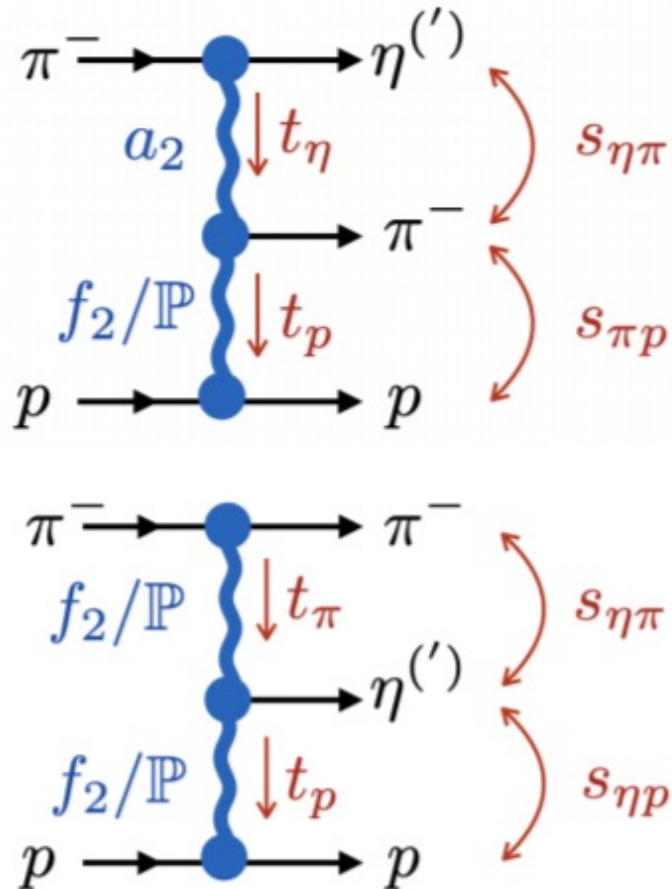


Rodas et al., PRL 122 042002 (2019)



Poles	Mass (MeV)	Width (MeV)
$a_2(1320)$	$1306.0 \pm 0.8 \pm 1.3$	$114.4 \pm 1.6 \pm 0.0$
$a_2'(1700)$	$1722 \pm 15 \pm 67$	$247 \pm 17 \pm 63$
π_1	$1564 \pm 24 \pm 86$	$492 \pm 54 \pm 102$

High-energy region model

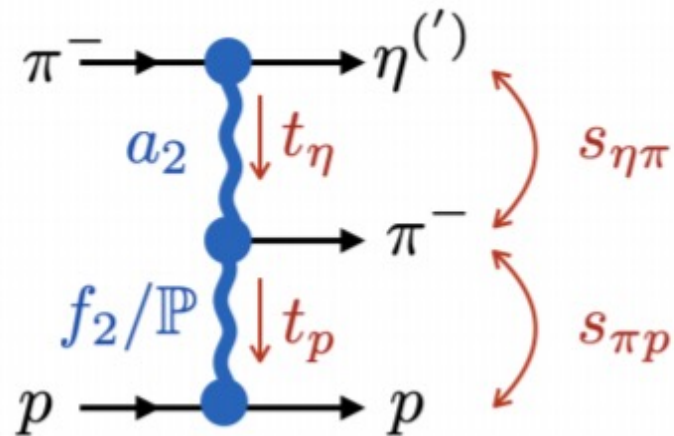
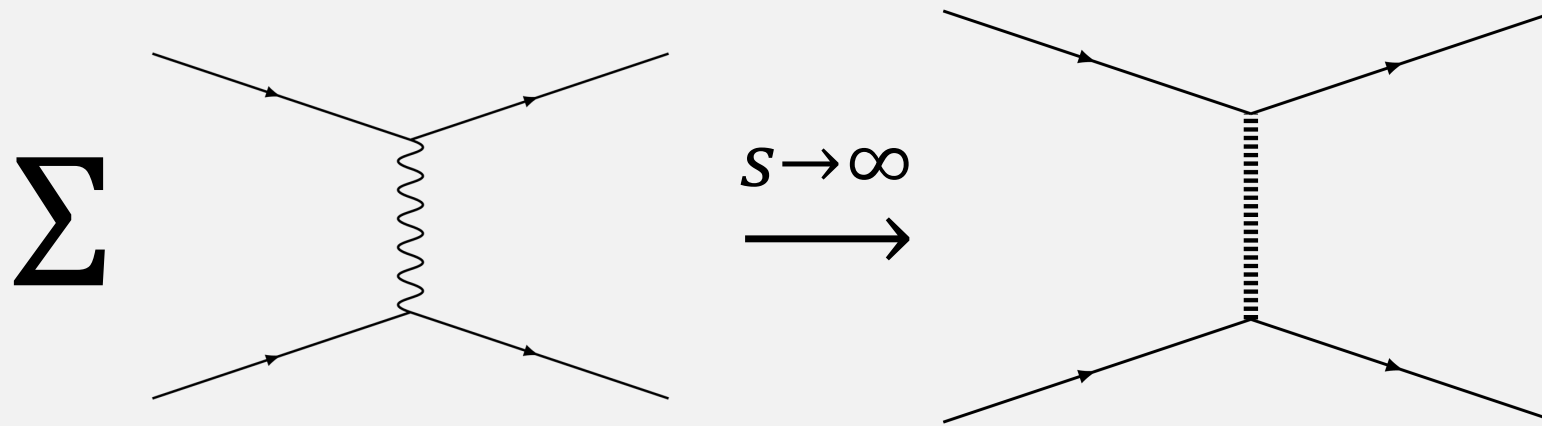


$$I_{\text{Th}}(m_f, \Omega) = k(m_f) |A_{\text{Th}}(m_f, \Omega)|^2$$

$$A_{\text{Th}}(m_f, \Omega) = c_1 A_{a_2 \mathbb{P}}^I + c_2 A_{a_2 f_2}^I + c_3 A_{f_2 \mathbb{P}}^{II} + c_4 A_{f_2 f_2}^{II} + c_5 A_{\mathbb{P} \mathbb{P}}^{II} + c_6 A_{\mathbb{P} f_2}^{II}.$$

Bibrzycki et al., EPJC 81 647 (2021)

Reggeons



$$\begin{aligned}\alpha_{a_2}(t) &= 0.53 + 0.90t, \\ \alpha_{f_2}(t) &= 0.47 + 0.89t, \\ \alpha_{\mathbb{P}}(t) &= 1.08 + 0.25t,\end{aligned}$$

High-energy region models

$$\begin{aligned}
 A_{a_2 P}^I &= T(\alpha_{a_2}(t_\eta), \alpha_P(t_p); s_{\eta\pi}, s_{\pi p}), \\
 A_{a_2 f_2}^I &= T(\alpha_{a_2}(t_\eta), \alpha_{f_2}(t_p); s_{\eta\pi}, s_{\pi p}), \\
 A_{f_2 P}^{II} &= T(\alpha_{f_2}(t_\pi), \alpha_P(t_p); s_{\eta\pi}, s_{\eta p}), \\
 A_{f_2 f_2}^{II} &= T(\alpha_{f_2}(t_\pi), \alpha_{f_2}(t_p); s_{\eta\pi}, s_{\eta p}), \\
 A_{P P}^{II} &= T(\alpha_P(t_\pi), \alpha_P(t_p); s_{\eta\pi}, s_{\eta p}), \\
 A_{P f_2}^{II} &= T(\alpha_P(t_\pi), \alpha_{f_2}(t_p); s_{\eta\pi}, s_{\eta p}).
 \end{aligned}$$

Shimada, Martin, Irving, NPB 142 344 (1978)
 Shi et al., PRD 91 034007 (2015)

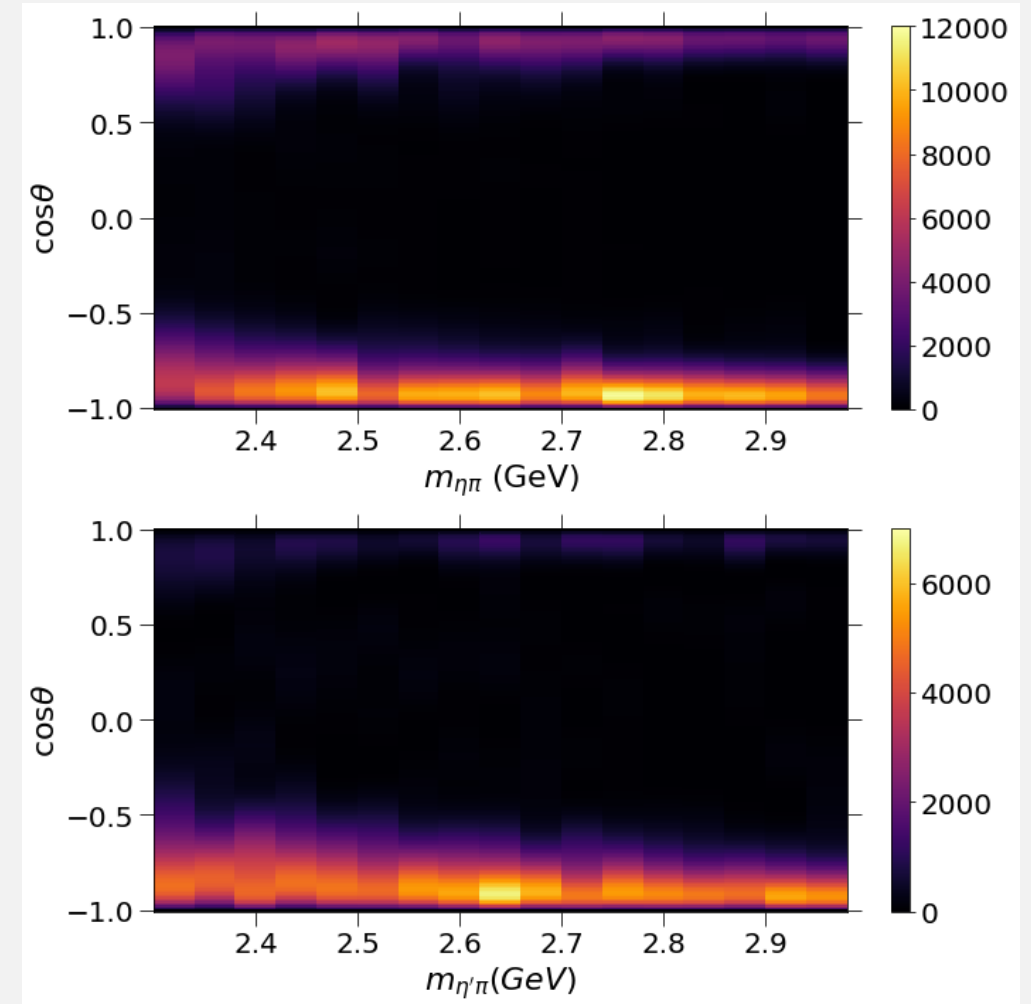
$$T(\alpha_1, \alpha_2; s_1, s_2) = K \Gamma(1 - \alpha_1) \Gamma(1 - \alpha_2) \frac{(\alpha' s_1)^{\alpha_1} (\alpha' s_2)^{\alpha_2}}{\alpha' s} \left[\frac{\xi_1 \xi_{21}}{\kappa^{\alpha_1}} V(\alpha_1, \alpha_2, \kappa) + \frac{\xi_2 \xi_{12}}{\kappa^{\alpha_2}} V(\alpha_2, \alpha_1, \kappa) \right]$$

$$V(\alpha_1, \alpha_2, \kappa) = \frac{\Gamma(\alpha_1 - \alpha_2)}{\Gamma(1 - \alpha_2)} {}_1F_1(1 - \alpha_1, 1 - \alpha_1 + \alpha_2, -\kappa)$$

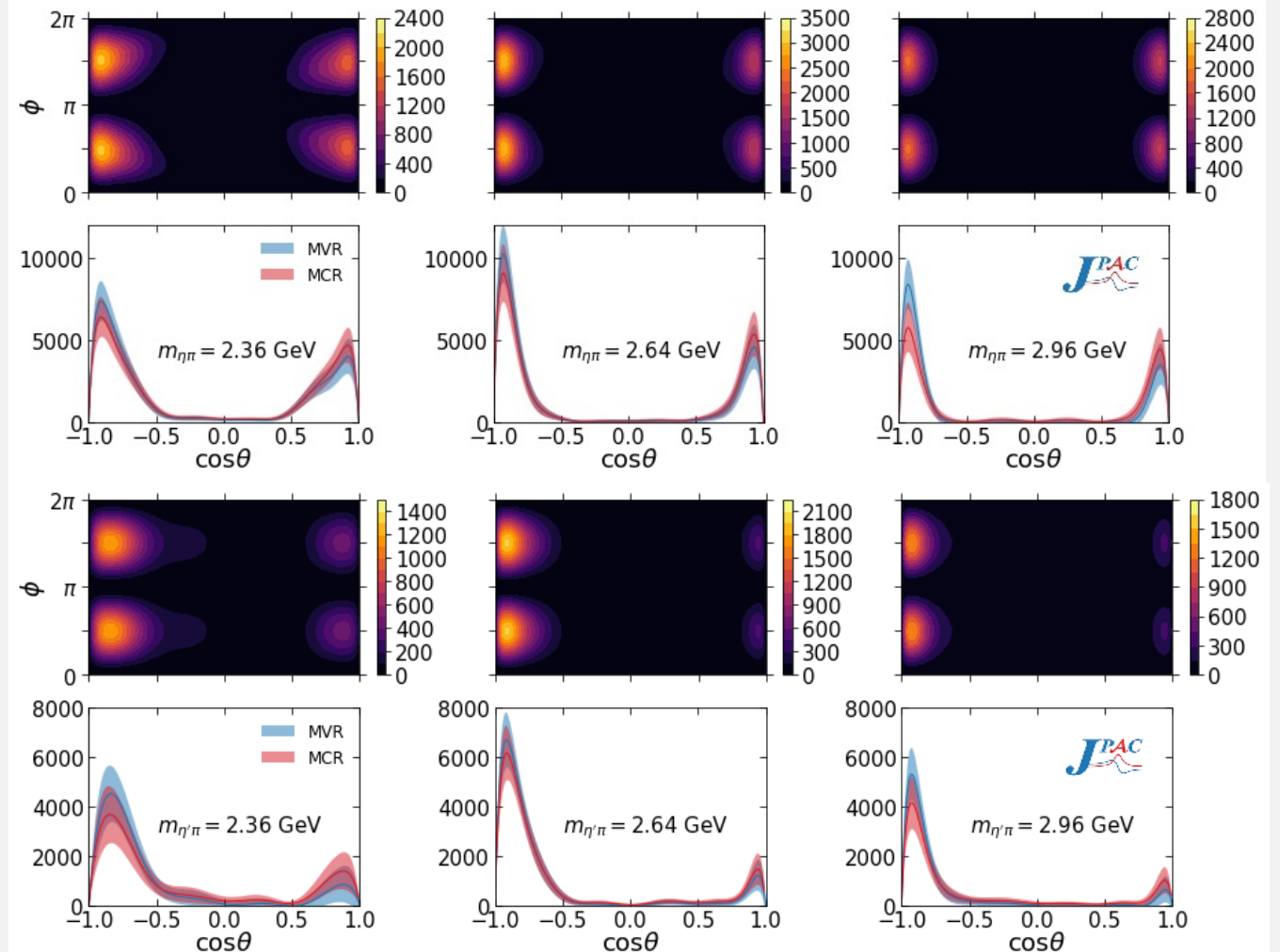
From “partial waves” to intensity distributions

$$I(m_f, \Omega) = \sum_{\epsilon=\pm} \left| \sum_{L,M} f_{LM}^{\epsilon}(m_f) \Psi_{LM}^{\epsilon}(\Omega) \right|^2$$

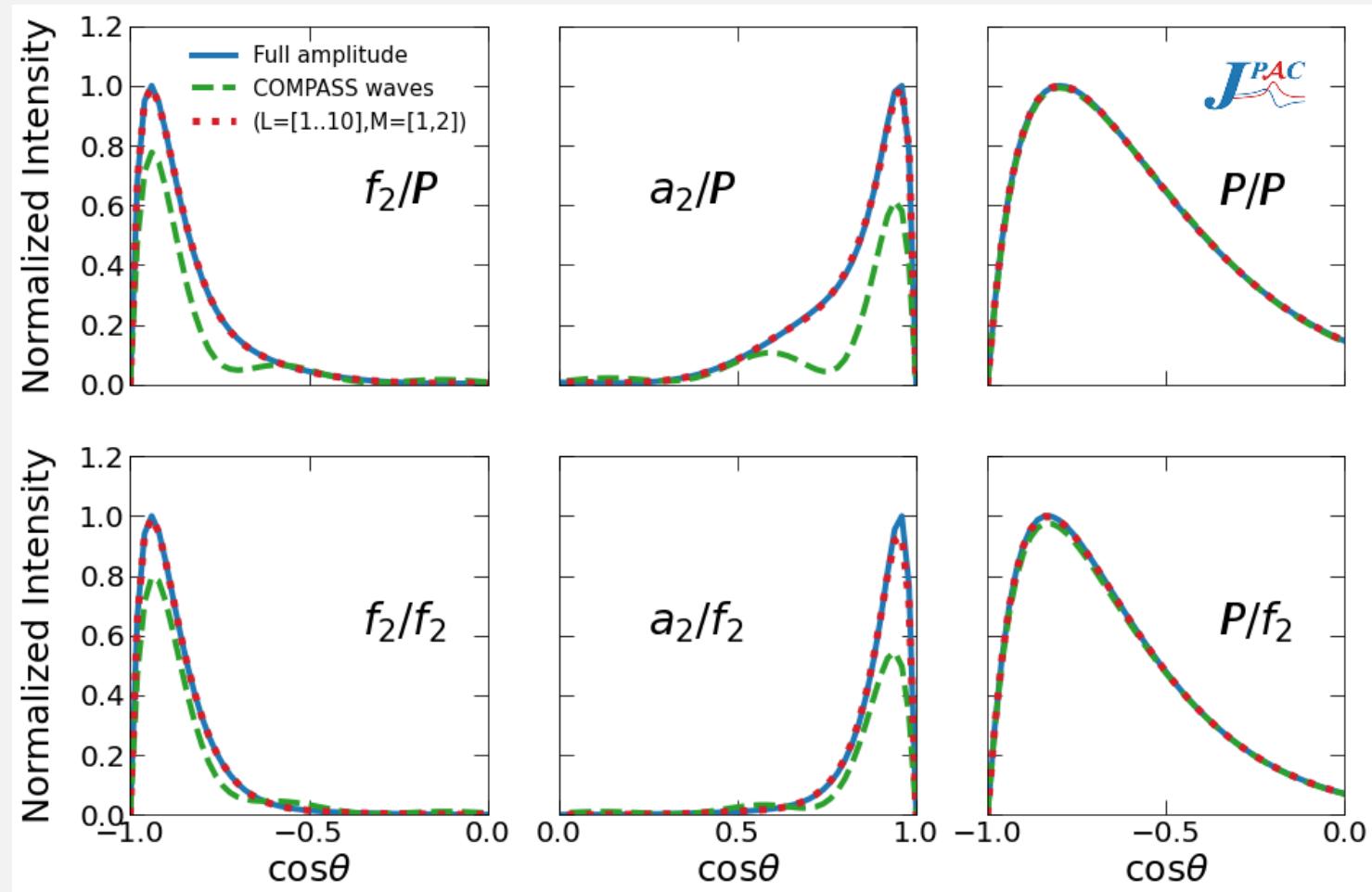
Bibrzycki et al., EPJC 81 647 (2021)



Intensity in the high-energy region



Truncation effect

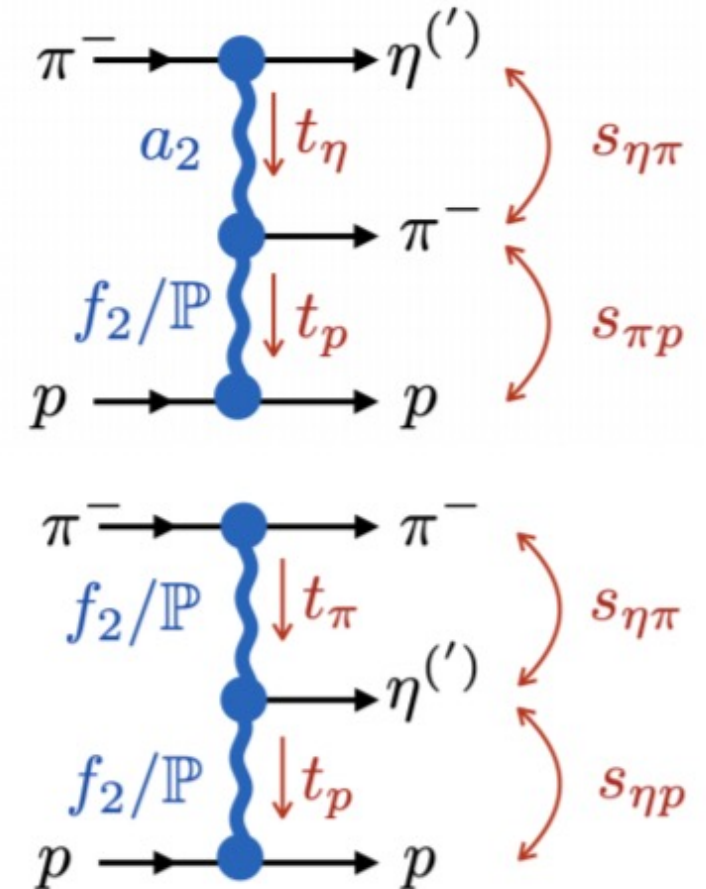


Minimal models

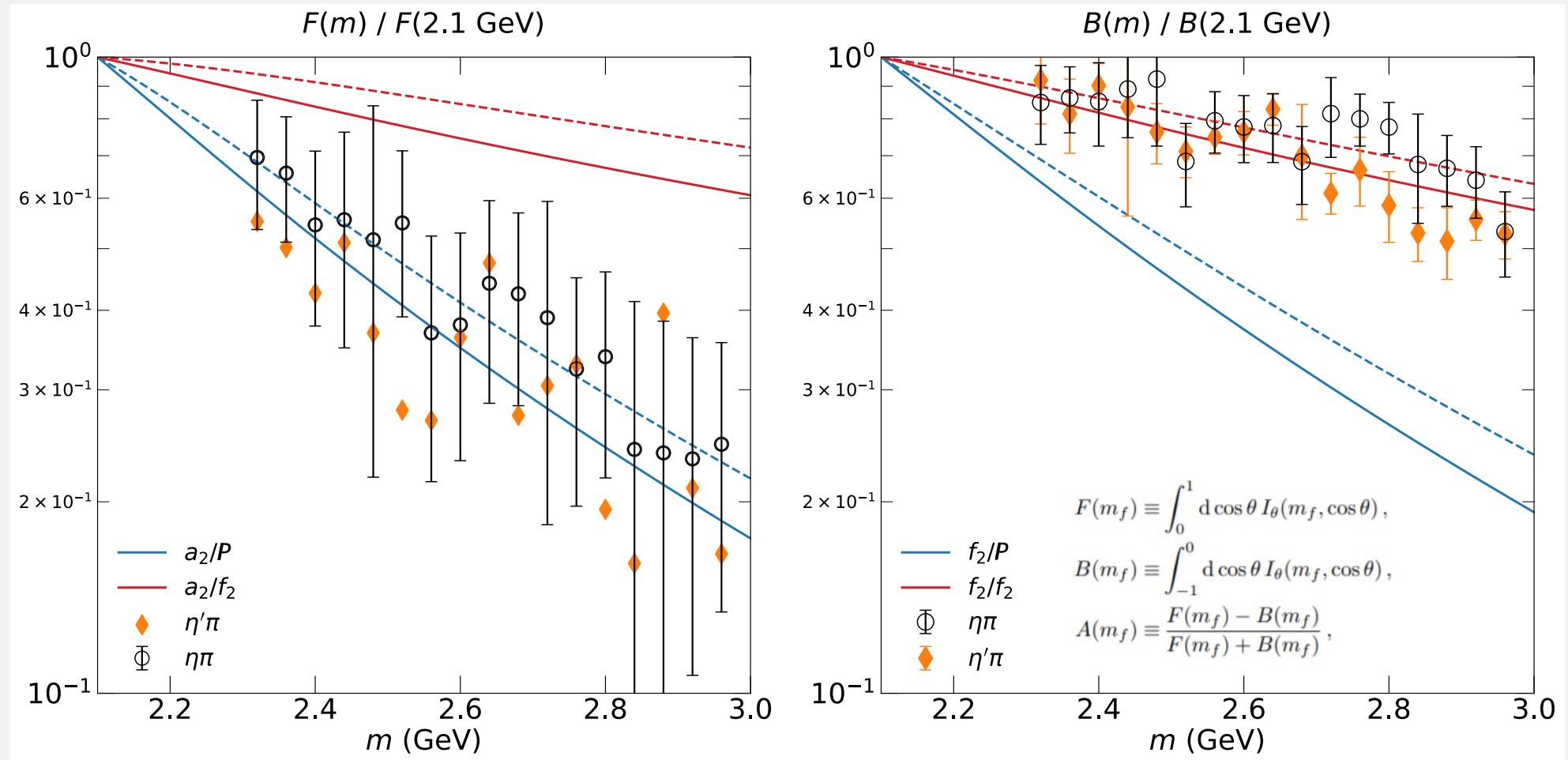
$MIN: a_2/P + a_2/f_2 + f_2/f_2$

$MIN + f/P: a_2/P + a_2/f_2 + f_2/f_2 + f_2/P$

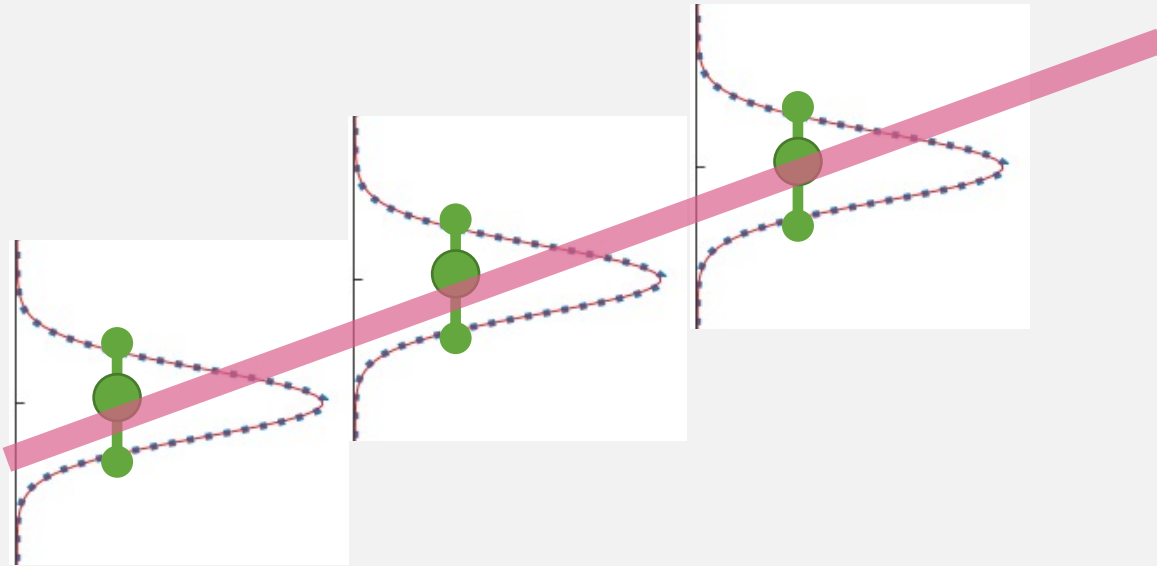
$MIN + P/P: a_2/P + a_2/f_2 + f_2/f_2 + P/P$



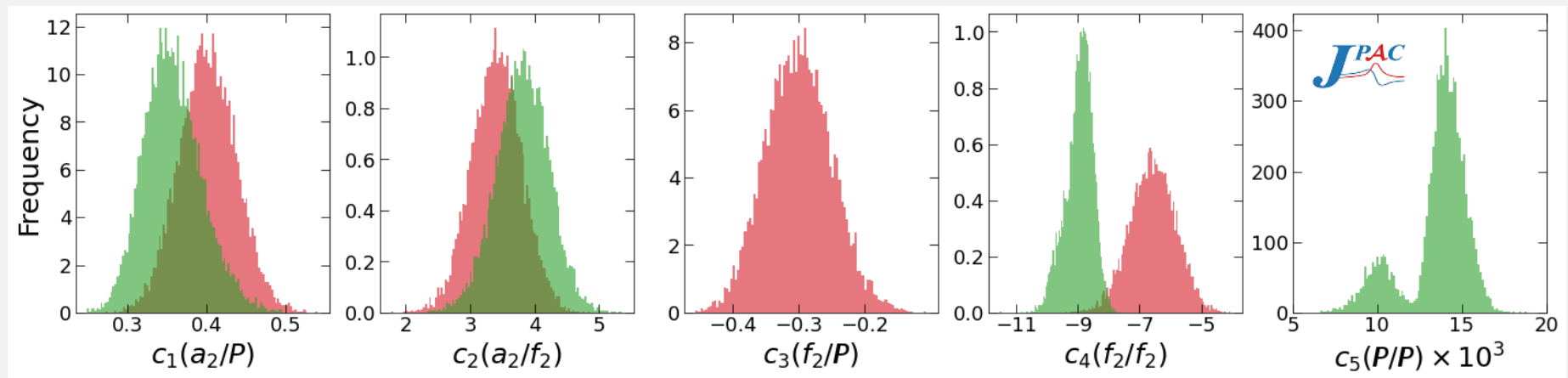
Forward and backward peaks



Bootstrap



You can run a Monte Carlo over the data, generating pseudodata and adjust each pseudo-dataset, obtaining parameter distributions

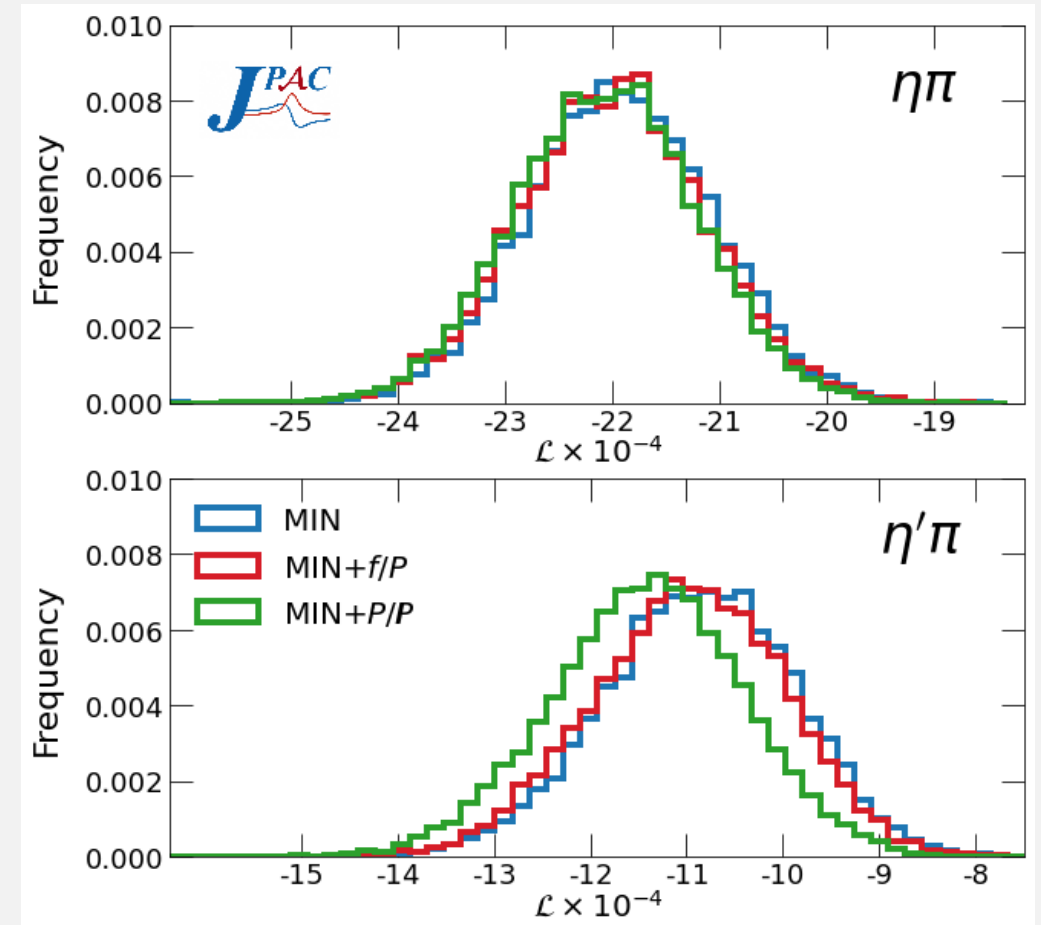


Extended log-likelihood (ELLH)

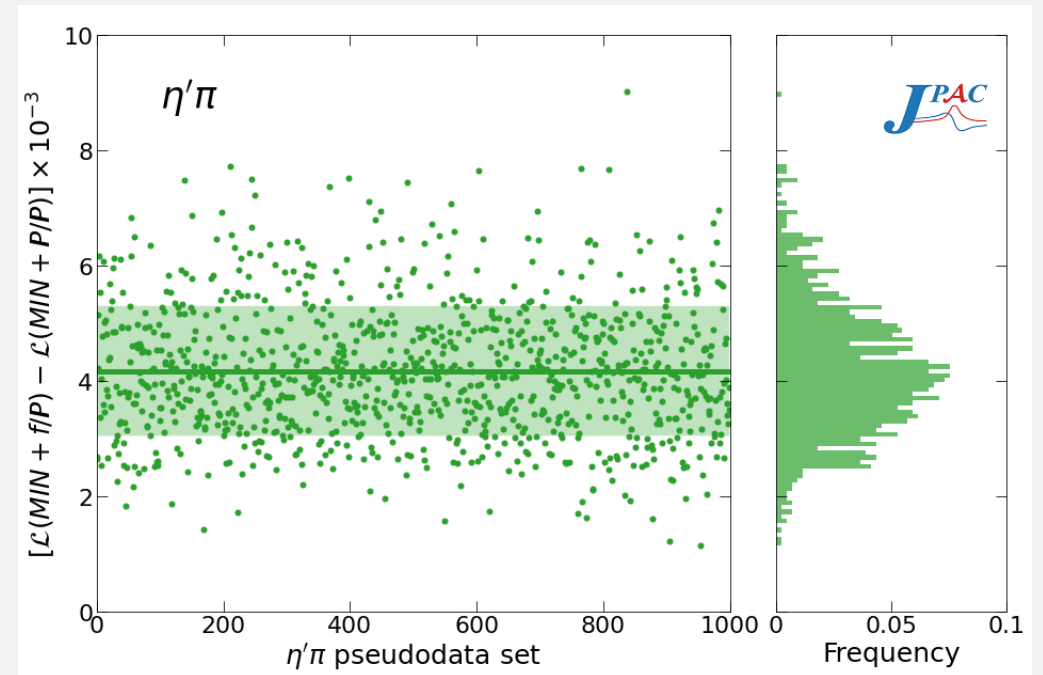
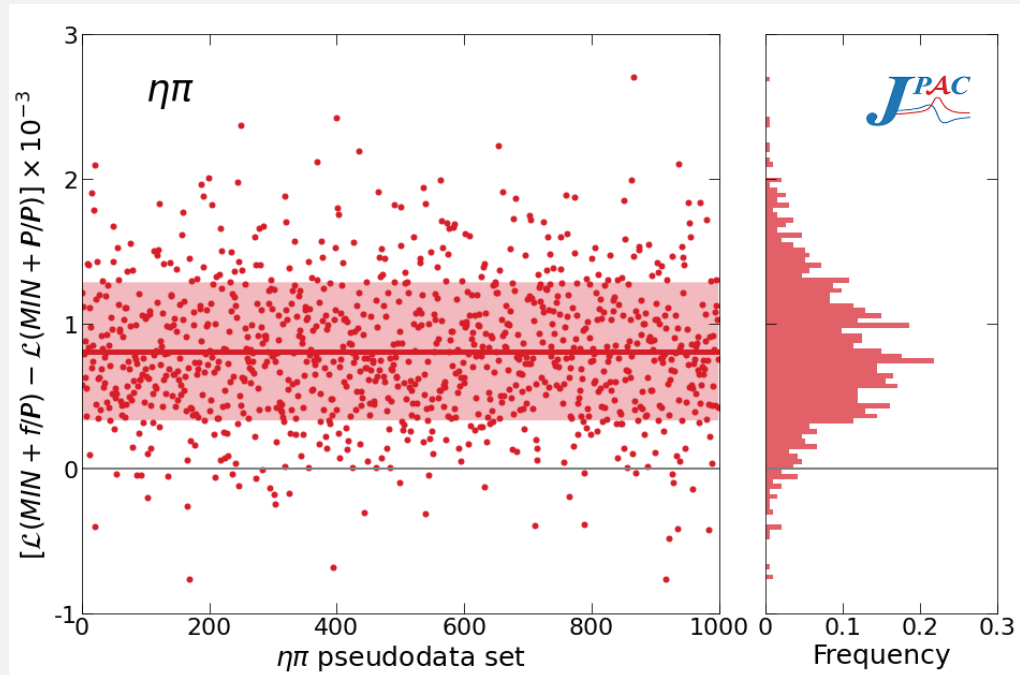
$$\mathcal{L}(\{c\}) = \sum_{m_i} \int d\Omega [I_T(m_i|\{c\}) - I_E(m_i) \log I_T(m_i|\{c\})],$$

Incorporates the restriction
on the total intensity

Barlow, NIMA 297 496 (1990)



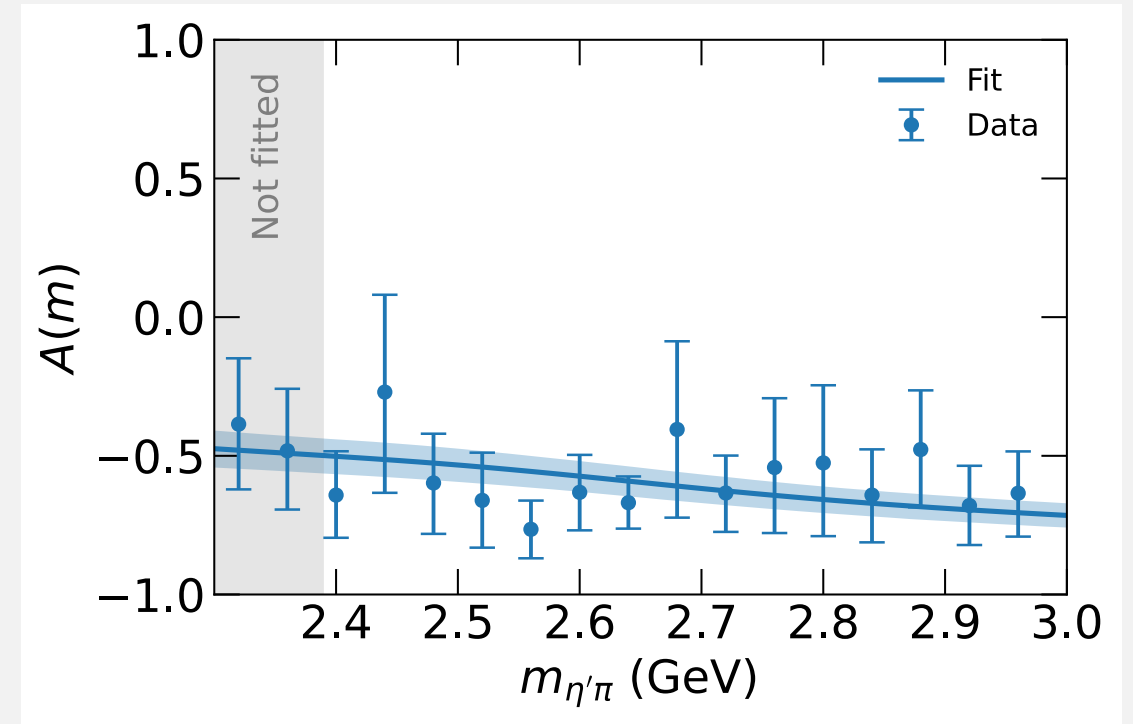
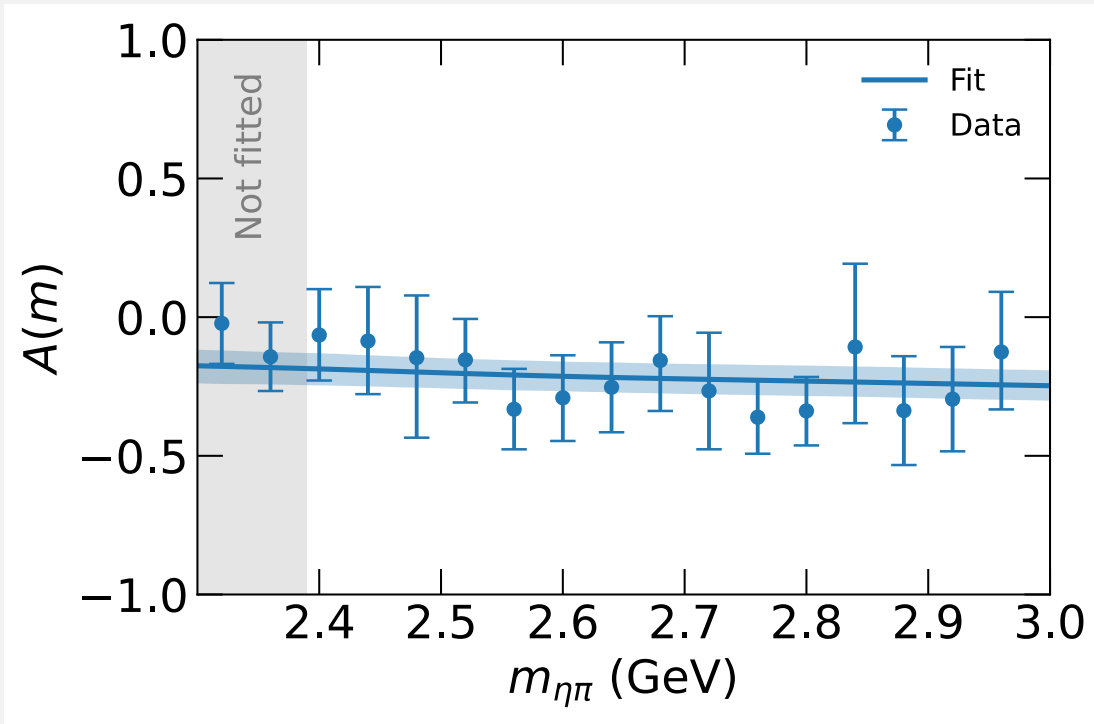
Comparing ELLH



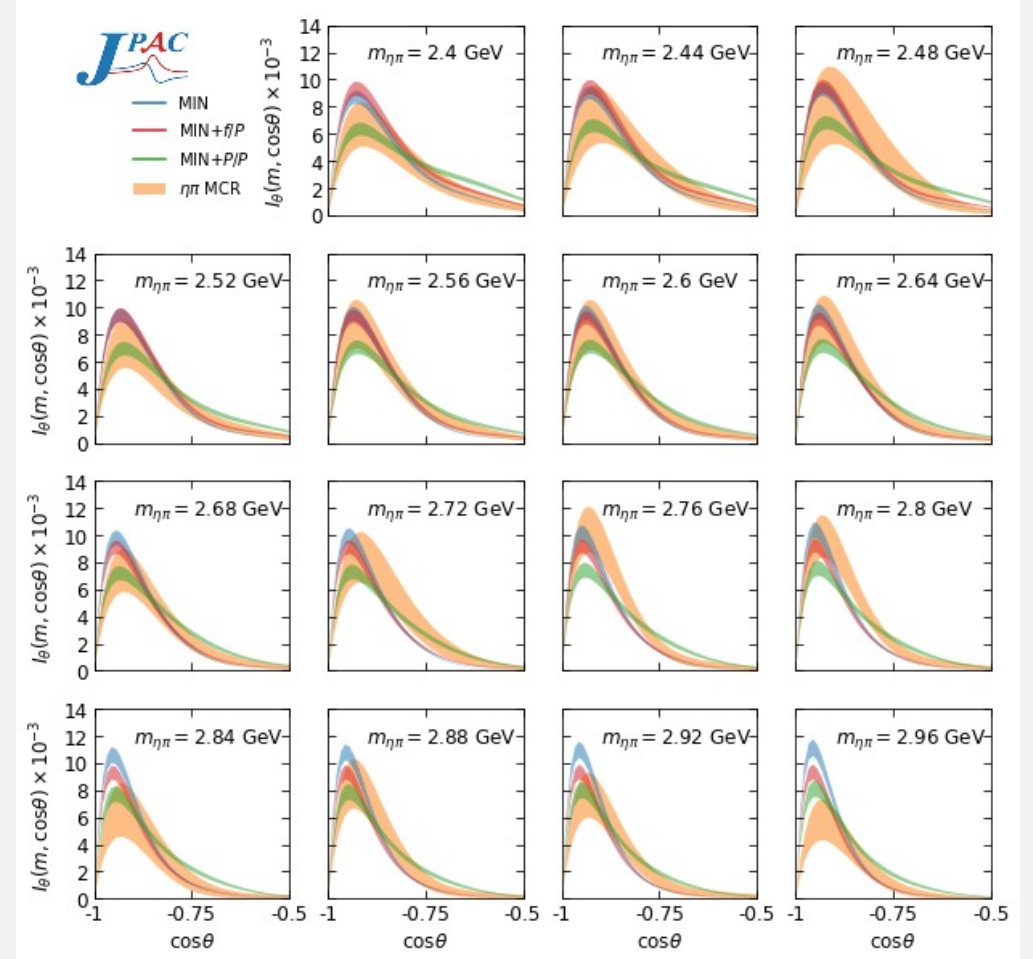
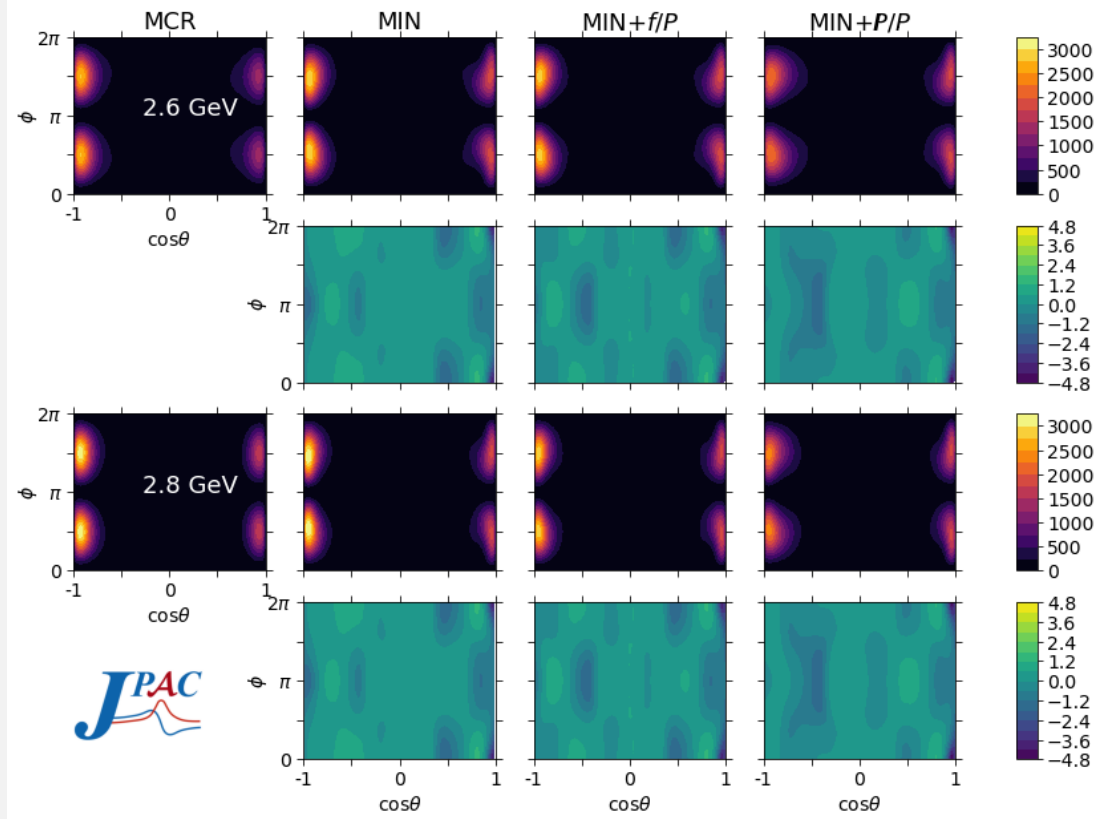
Asymmetry

The asymmetry is proof of exotics

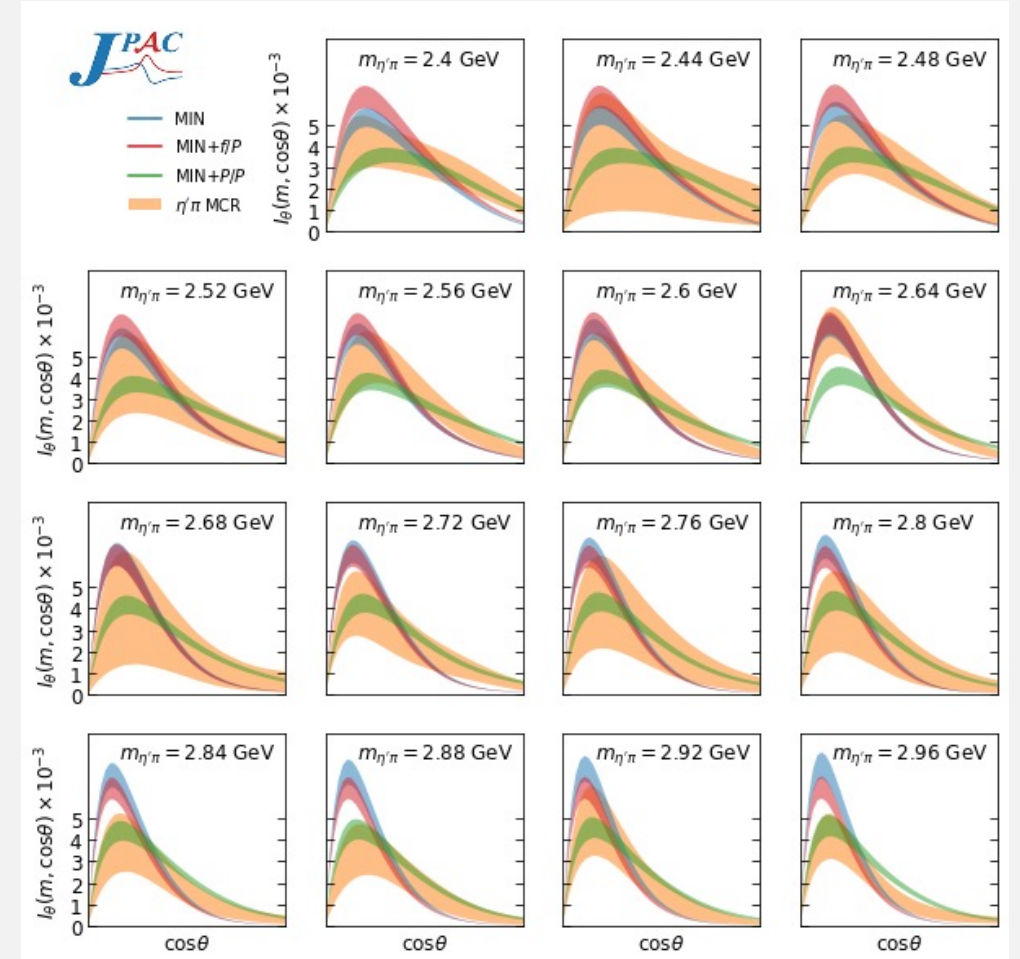
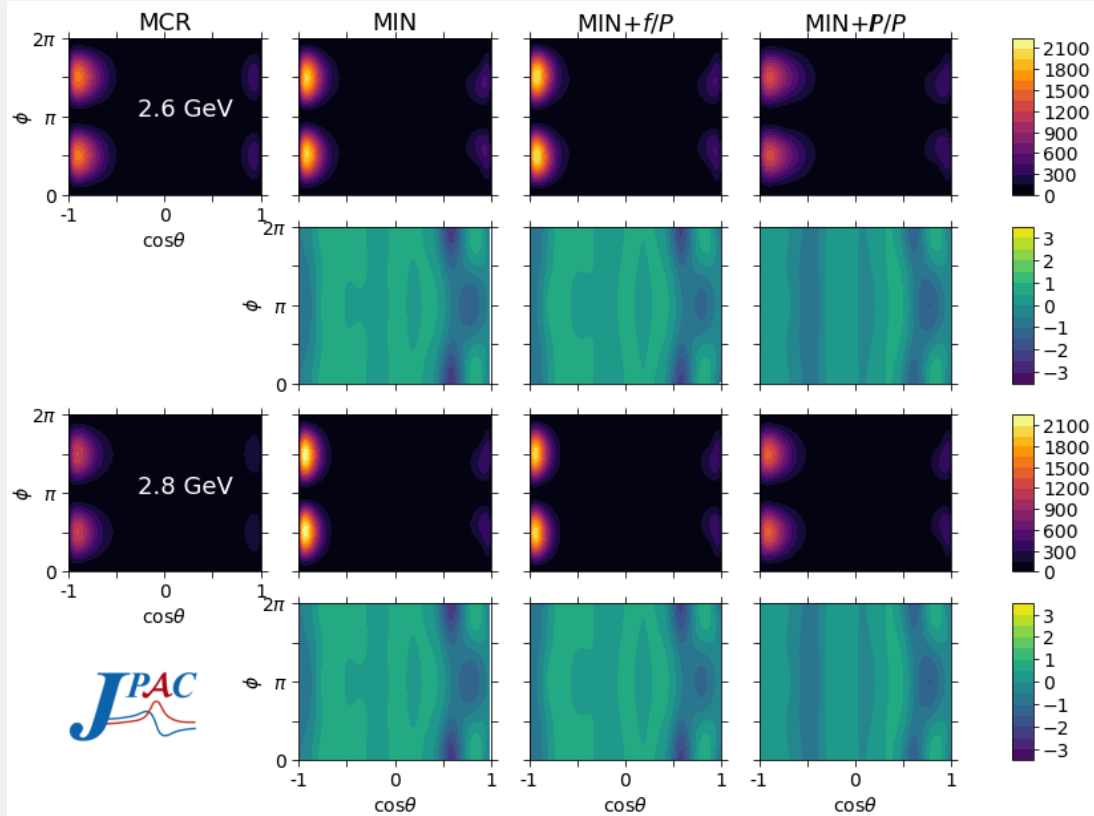
$$F(m_f) \equiv \int_0^1 d \cos \theta I_\theta(m_f, \cos \theta),$$
$$B(m_f) \equiv \int_{-1}^0 d \cos \theta I_\theta(m_f, \cos \theta),$$
$$A(m_f) \equiv \frac{F(m_f) - B(m_f)}{F(m_f) + B(m_f)},$$



Angular distribution



Angular distribution

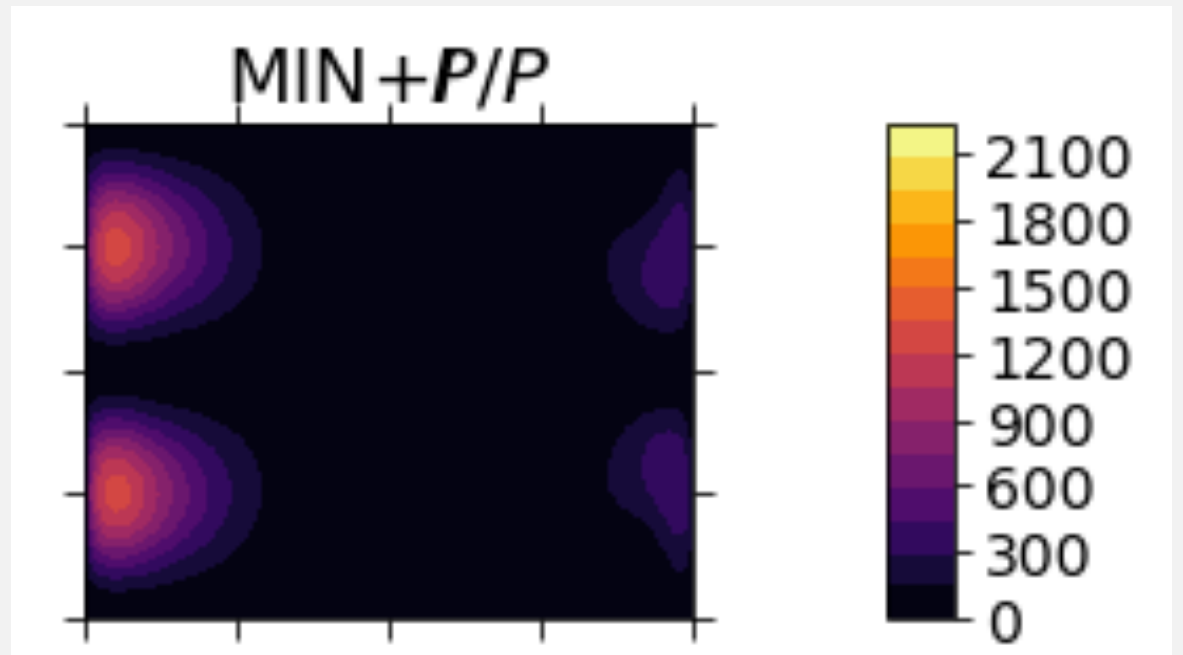


Constrained partial wave analysis (cPW)

$$\hat{\mathcal{L}}_i(\{p\}) = \int d\Omega [I_E(m_i|\{p\}) - I_T(m_i) \log I_E(m_i|\{p\})],$$

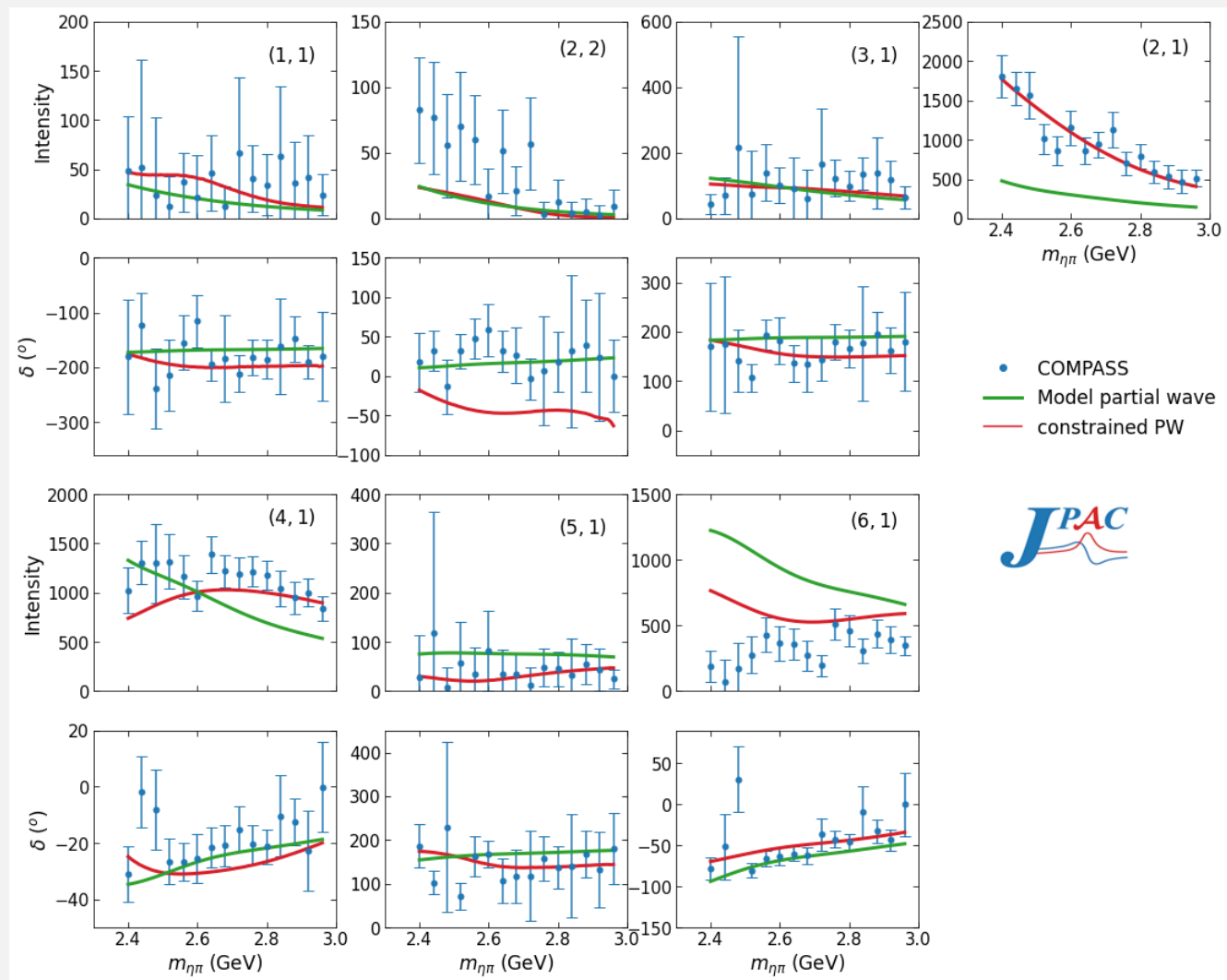
$$I_E(m_i, \Omega|\{p\}) = \left| \sum_{L,M} p_{LM}(m_i) \Psi_{LM}^+(\Omega) \right|^2$$

We treat theory as experimental data and
we extract the partial waves under the
Same approximations made by COMPASS

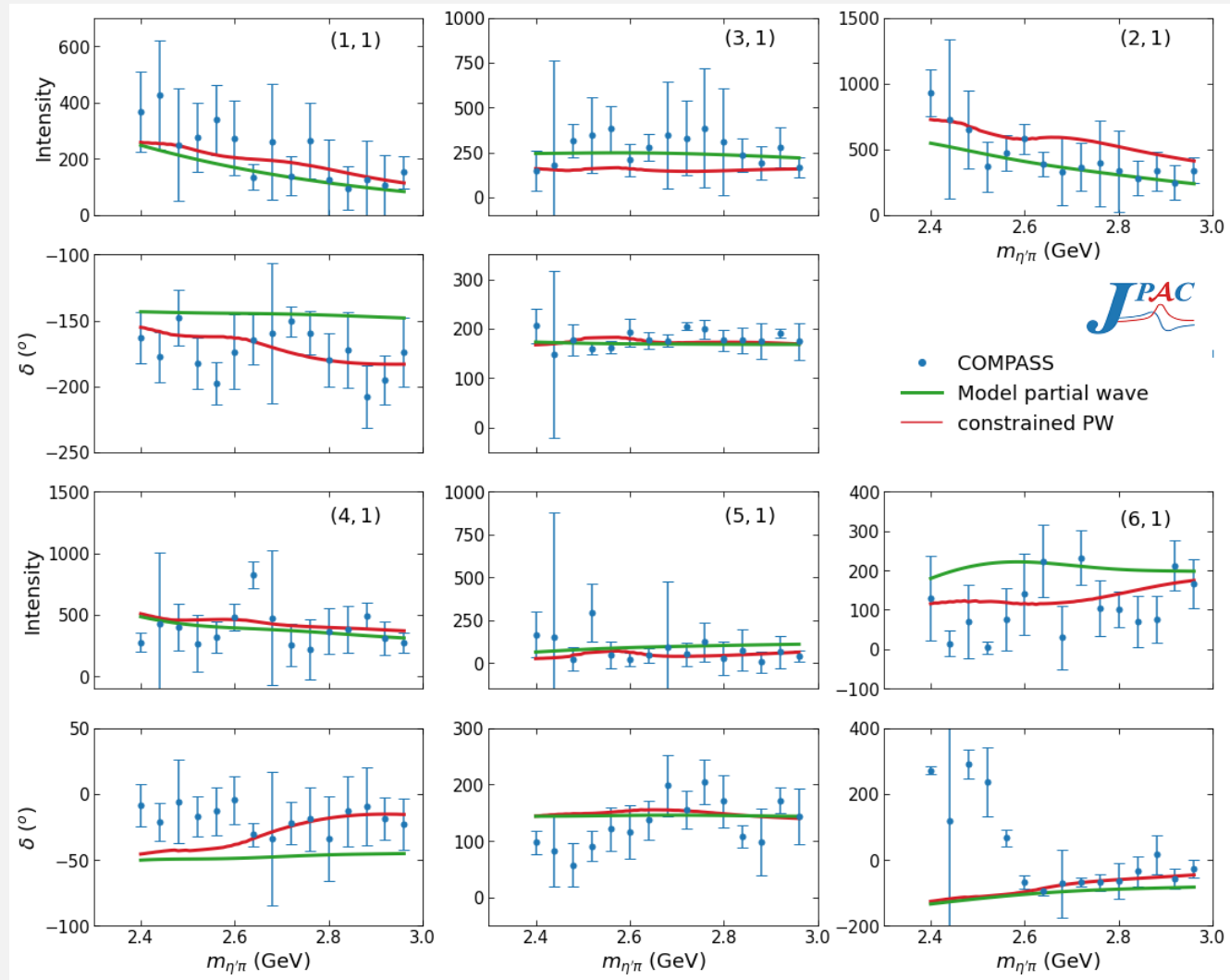


cPW $\eta\pi$

$$\hat{\mathcal{L}}_i(\{p\}) = \int d\Omega [I_E(m_i|\{p\}) - I_T(m_i) \log I_E(m_i|\{p\})],$$



cPW $\eta'\pi$



Results

$\eta\pi$ and $\eta'\pi$ at “high” energy

- *Golden channels to search exotic quantum numbers*

Above 2.5 GeV

- *$\eta\pi$ intensity can be described by four double-Regge exchanges: a_2/P , a_2/f_2 , f_2/f_2 y f_2/P o P/P . Details in Bibrzycki et al., EPJC 81 647 (2021)*

Pomeron is essential for $\cos \theta \sim -1$ data

- *Gluonic degrees of freedom in $\eta'\pi$, potentially related to hybrids*

Results

Asymmetry

- *Proves the existence of exotics*

The dangers of partial wave truncation

- *Highlighted at high-energy because of Reggeons*
- *Shown by the cPW analysis*

The danger of not comparing “apples” to “apples”

- *If we had fitted COMPASS partial waves directly we would have been wrong*